

Test of NRQCD with Quarkonium Production in Jets at the LHC

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with

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Prashant Shrivastava (CMU)

[arXiv:1603.06981, 1702.05525, 1707.08629]

(QWG2017, Beijing, Nov 9, 17)

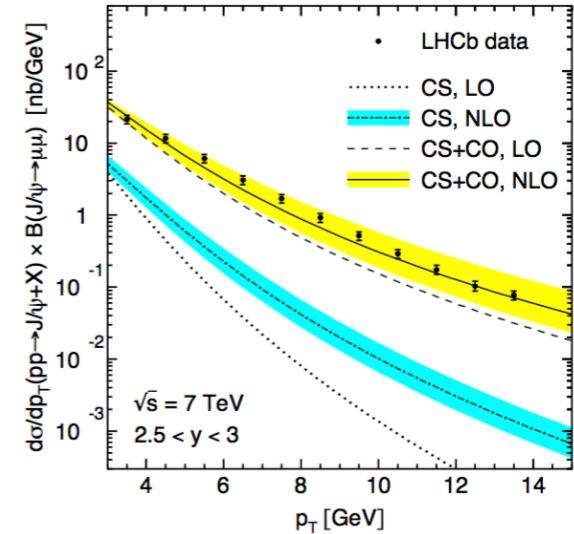
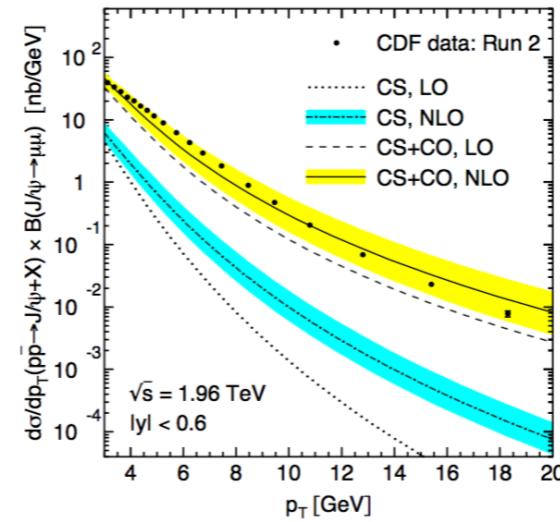


Outline

- ◆ Motivations
- ◆ NRQCD factorization formalism (extraction LDMEs)
- ◆ Fragmenting Jet Functions (analytic tool)
- ◆ e^+e^- collisions: B , J/ψ production in jets (test framework)
- ◆ pp collisions: J/ψ production in jets, compare with LHCb data
- ◆ Summary

Some Motivations

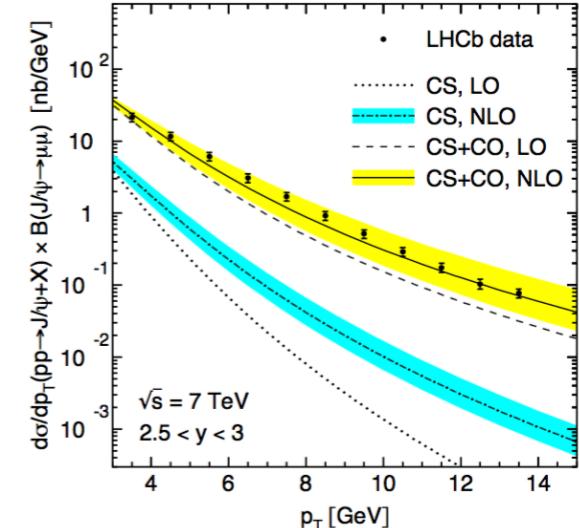
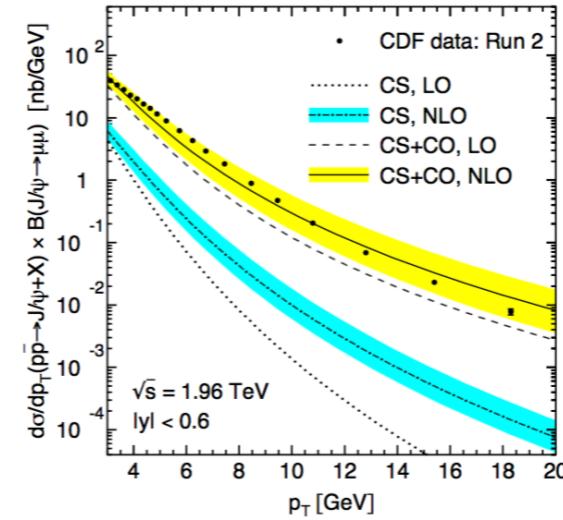
♦ Much of quarkonium production pheno. based on NRQCD factorization formalism



[Butenschon, Kniehl, 1105.0820]

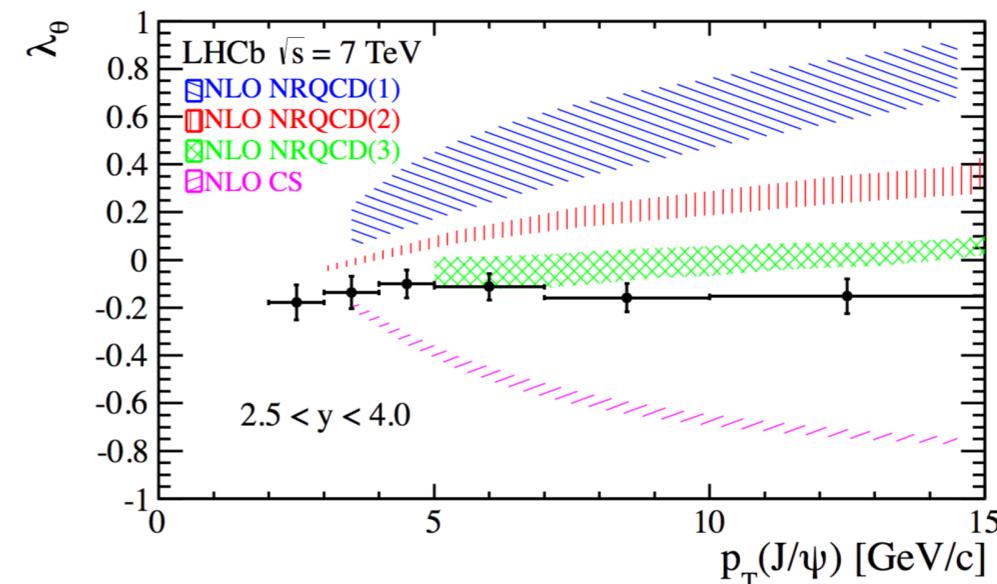
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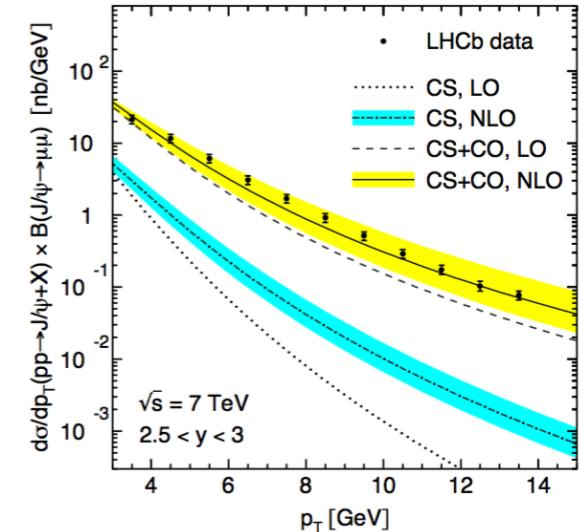
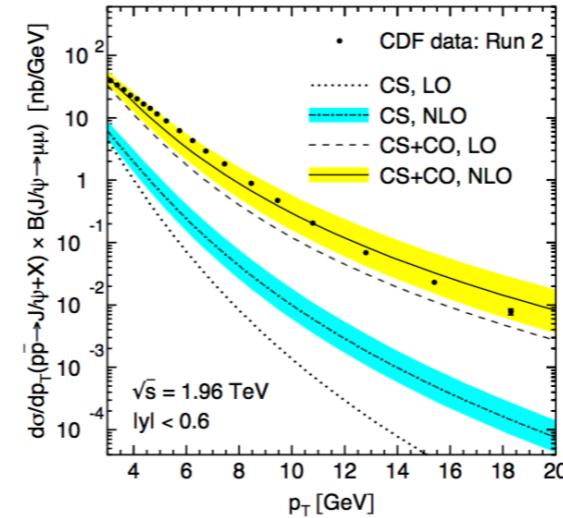
♦ In NRQCD, LDMEs supposed to be universal (however...)



[LHCb, 1307.6379]

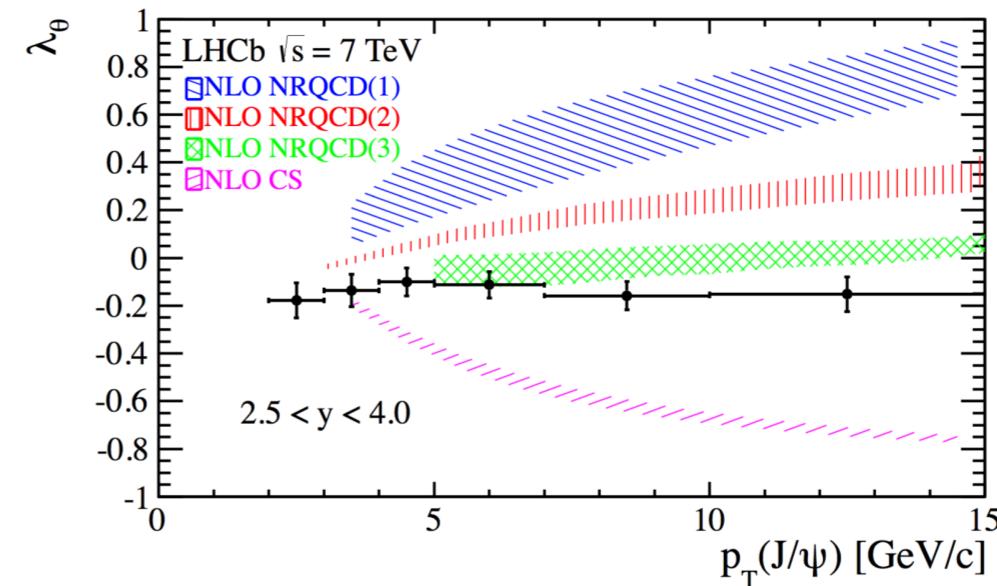
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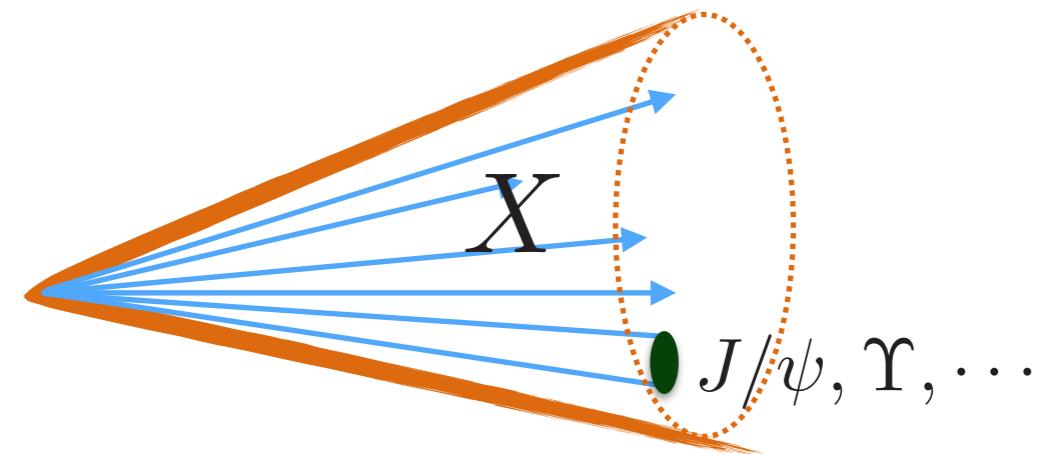
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♦ In NRQCD, LDMEs supposed to be universal (however...)



[LHCb, 1307.6379]

♦ Quarkonium production in jets provides us new way of studying these issues



NRQCD Factorization

Cross section:

SDCs (perturbative, scale $\geq m_H$)

$$\sigma(H) = \sum_n d_n \langle 0 | \mathcal{O}_n^H | 0 \rangle$$

LDMEs (nonperturbative, scale $< m_H$)

$$n : {}^1S_0^{[8]}, {}^3S_1^{[1,8]}, {}^3P_J^{[8]}, \dots$$

e.g.

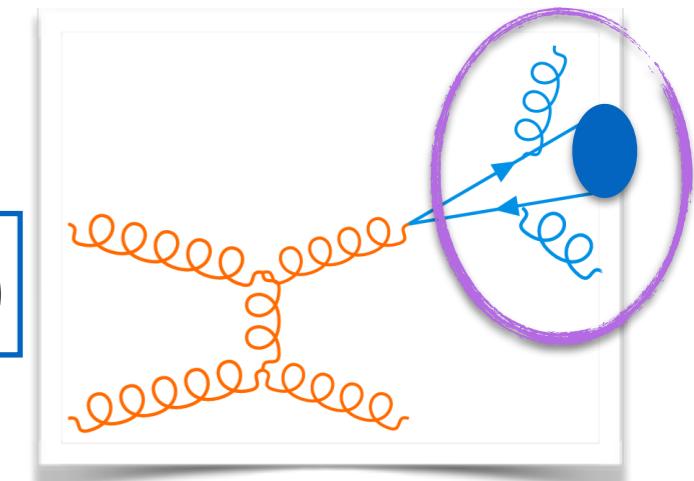
$$\langle 0 | \mathcal{O}_{^3S_1^{[8]}}^H | 0 \rangle = \sum_X \langle 0 | \chi^\dagger \sigma^i T^a \psi | H X \rangle \langle H X | \psi^\dagger \sigma^i T^a \chi | 0 \rangle$$

NRQCD Factorization

Factorized form of FF:

SDCs (perturbative, $\sim m_H$)

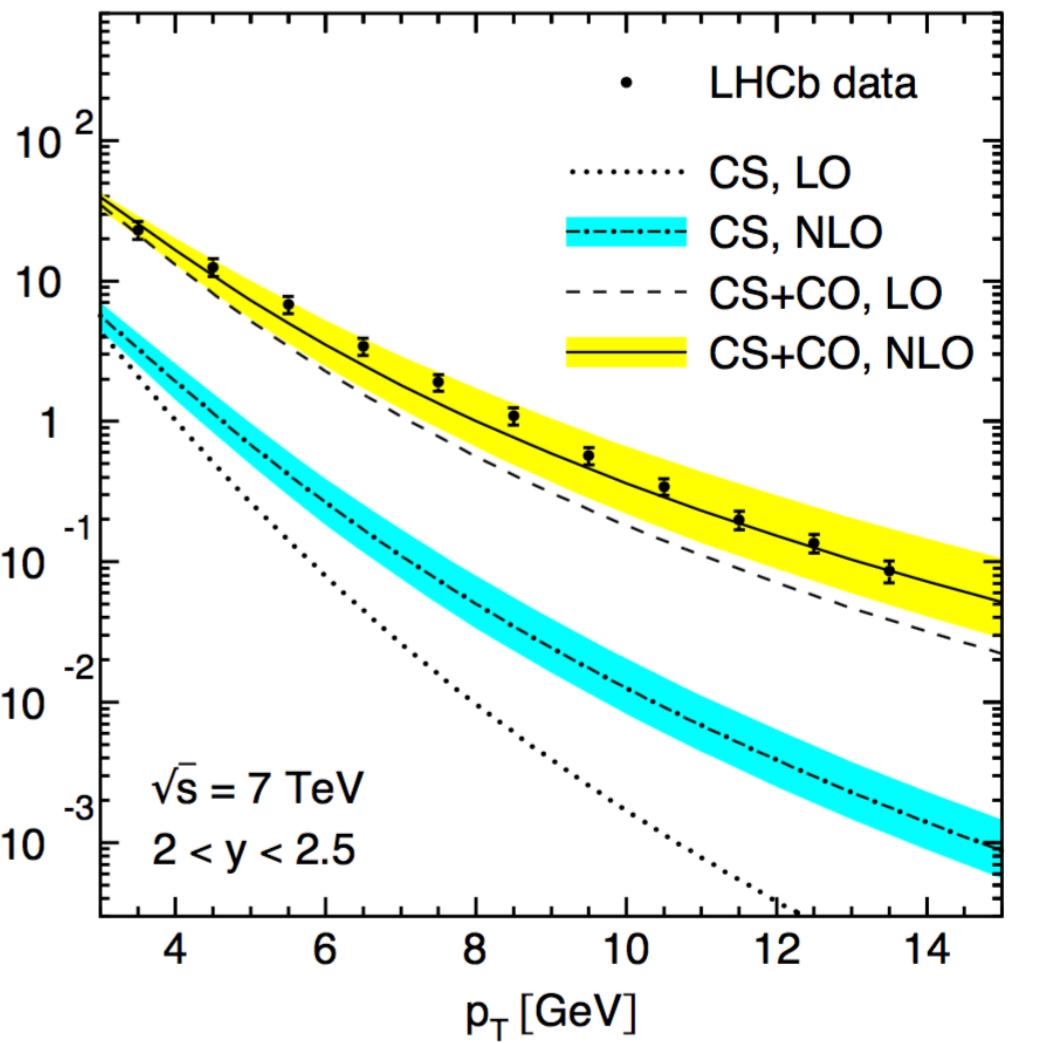
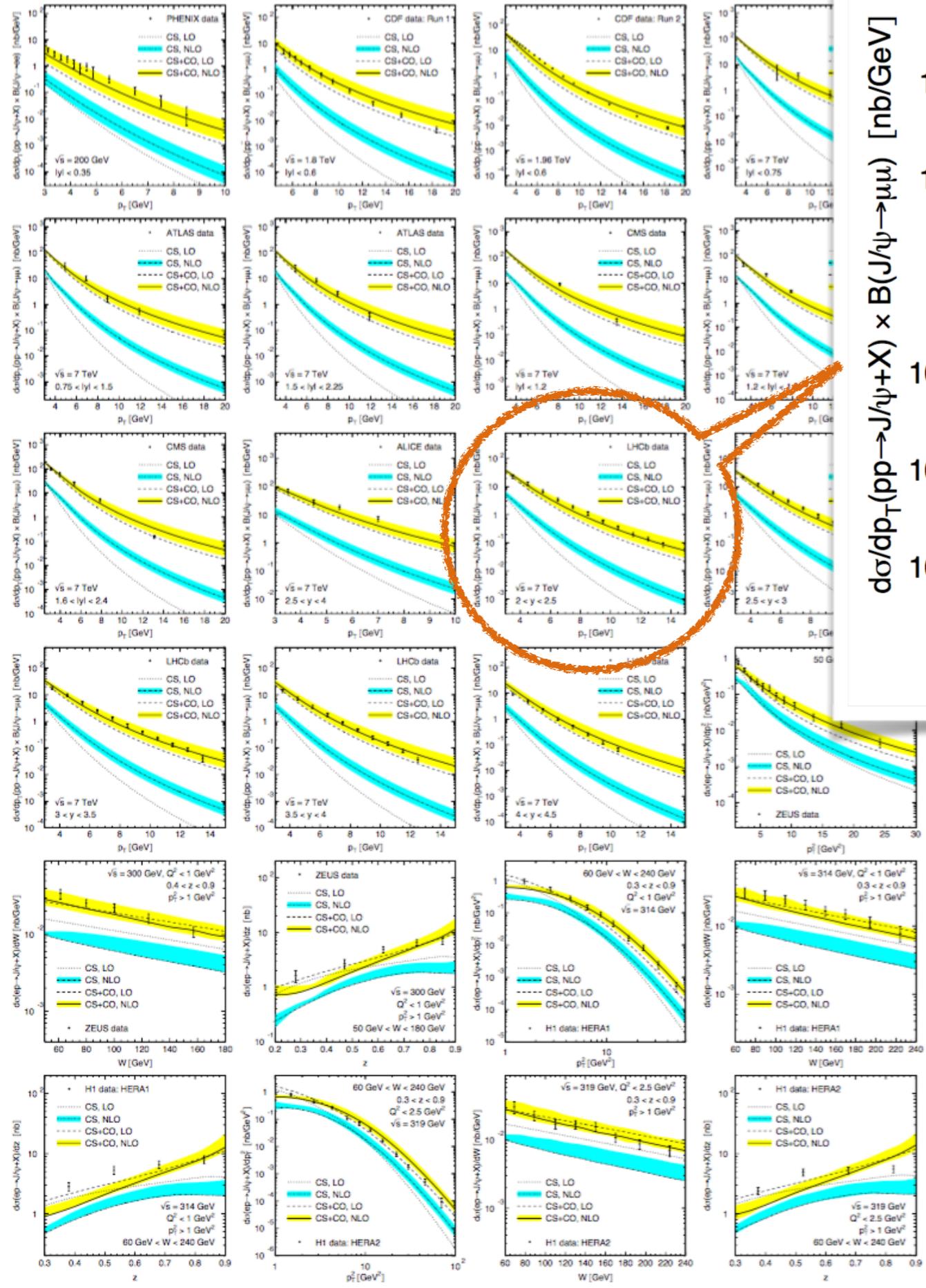
$$D_{q/g}^H = \sum_n d_{q/g,n} \langle 0 | \mathcal{O}_n^H | 0 \rangle$$



LDMEs (nonperturbative, $< m_H$)
The same as those in X-section

e.g. $D_{g \rightarrow \psi}^{1S_0^{(8)}}(z, 2m_c) = \frac{5\alpha_s^2(2m_c)}{96m_c^3} \langle \mathcal{O}^\psi(1S_0^{(8)}) \rangle (3z - 2z^2 + 2(1-z)\log(1-z))$

The channels we need: $1S_0^{[8]}$, $3S_1^{[1,8]}$, $3P_J^{[8]}$
(power counting up to v^4)



NLO Global Fit:

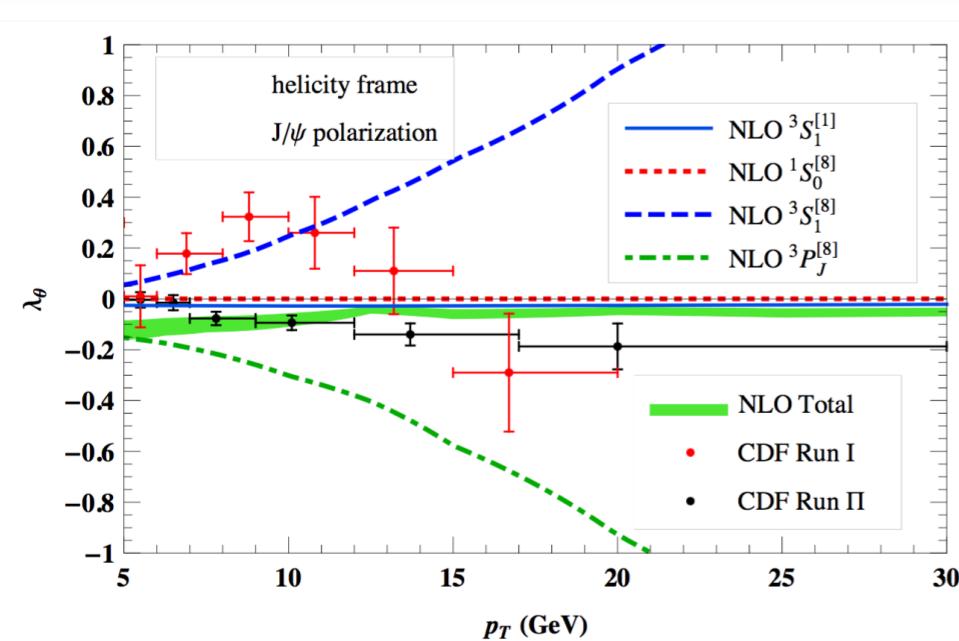
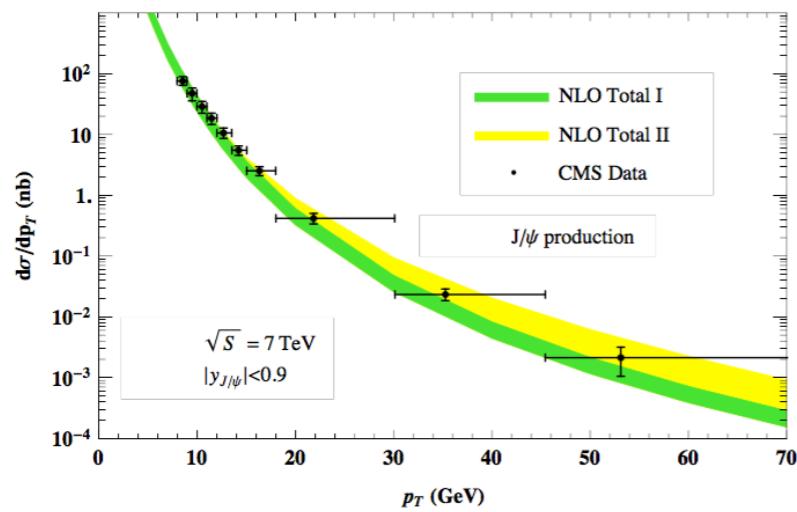
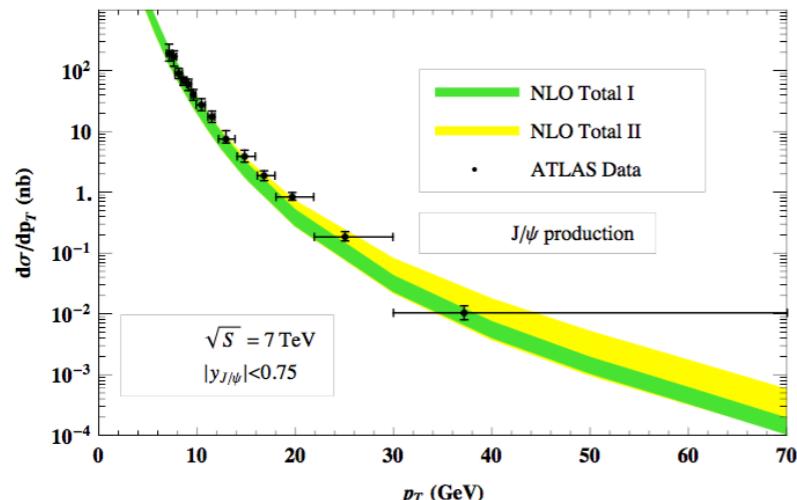
26 sets, 196 data points
 $\chi^2_{\text{d.o.f.}} = 4.42$

$p_T, p_T^2, w, z \dots$

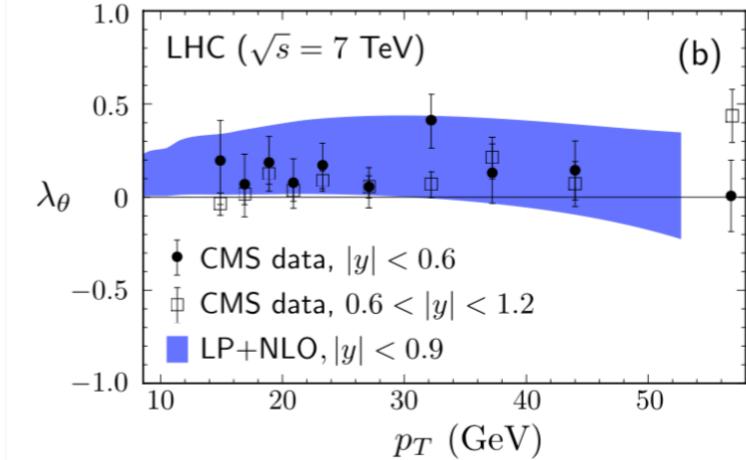
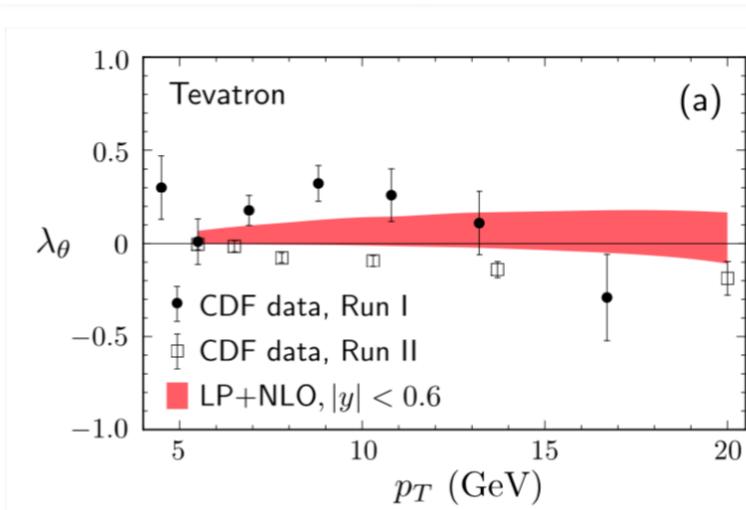
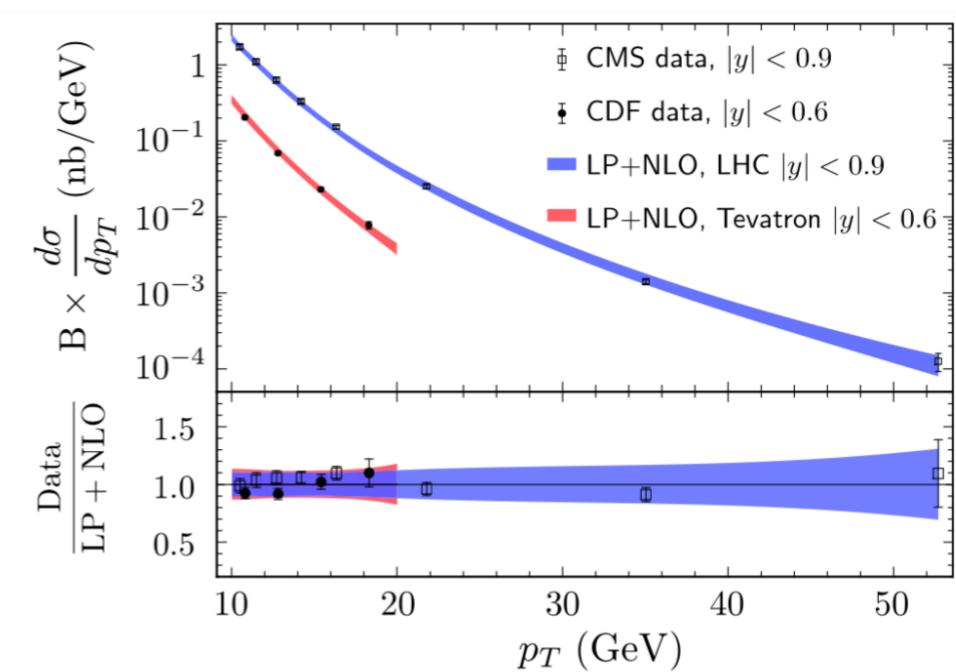
Includes low p_T region

[Butenschon, Kniehl, 1105.0820]

[Chao et al., 1201.2675]



[Bodwin et al., 1403.3612]



Fragmenting Jet Function (FJF)

$$\mathcal{G}_i^H(s, z, \mu) = \sum_j \int_z^1 \frac{dz'}{z'} \mathcal{J}_{ij}(s, z', \mu) D_j^H(z/z', m_H, \mu)$$

[M. Procura, I. Stewart, PRD, 81, 074009]

[M. Baumgart, A. Leibovich, T. Mehen, I. Rothstein, JHEP, 11, 003 (2014)]

Jet scale

Hadron scale

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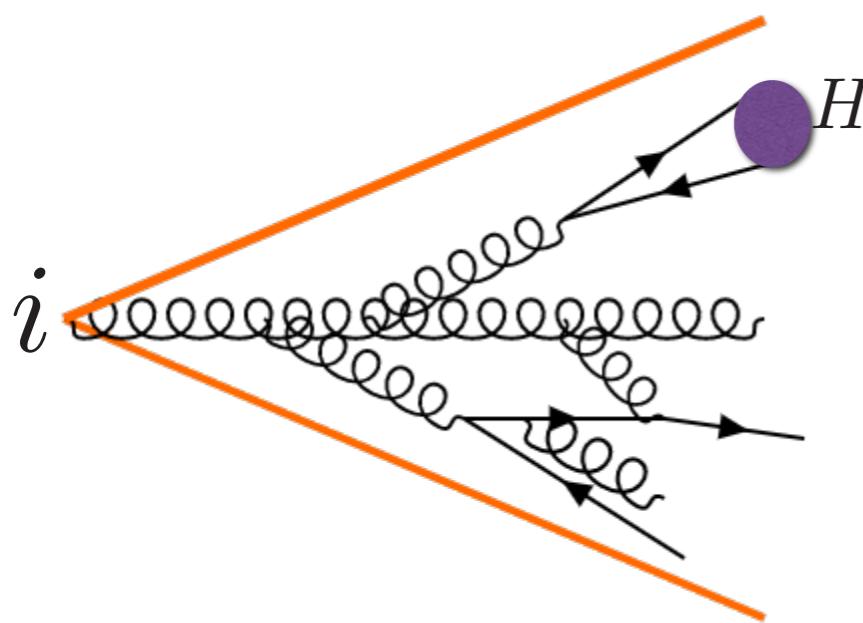
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Physical meaning:



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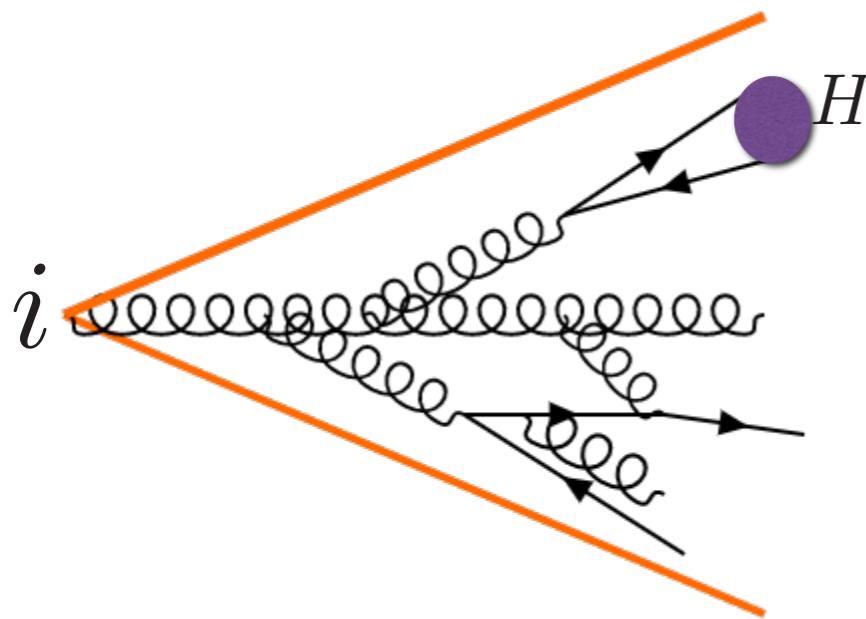
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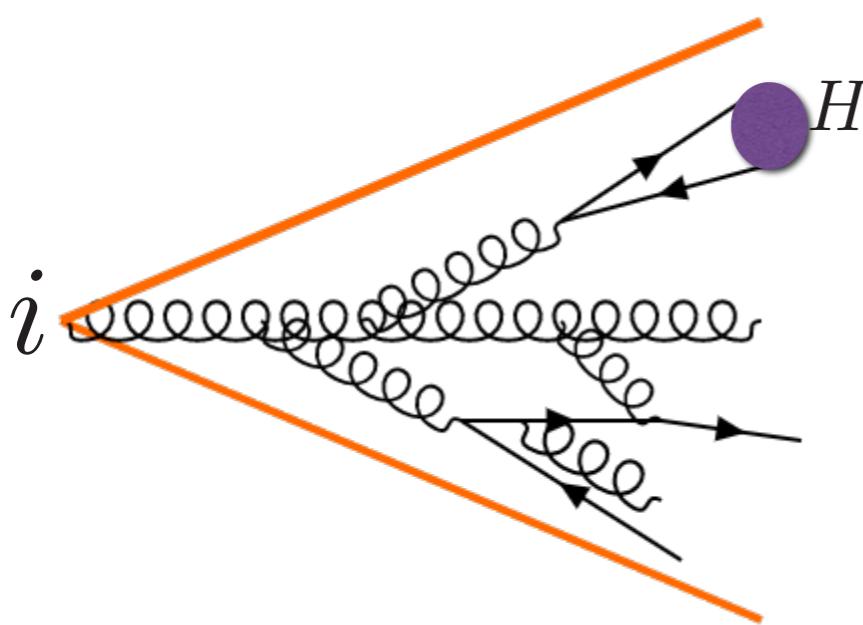
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Physical meaning:

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$$D_q^H(z) \equiv \text{Tr} \langle 0 | \frac{\vec{n}}{2} \delta(\frac{p_-}{z} - \mathcal{P}_-) \psi | HX \rangle \langle HX | \psi | 0 \rangle$$

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$$\mathcal{G}_i^H(s, z, \mu) = \sum_j \int_z^1 \frac{dz'}{z'} \mathcal{J}_{ij}(s, z', \mu) D_j^H(z/z', m_H, \mu)$$

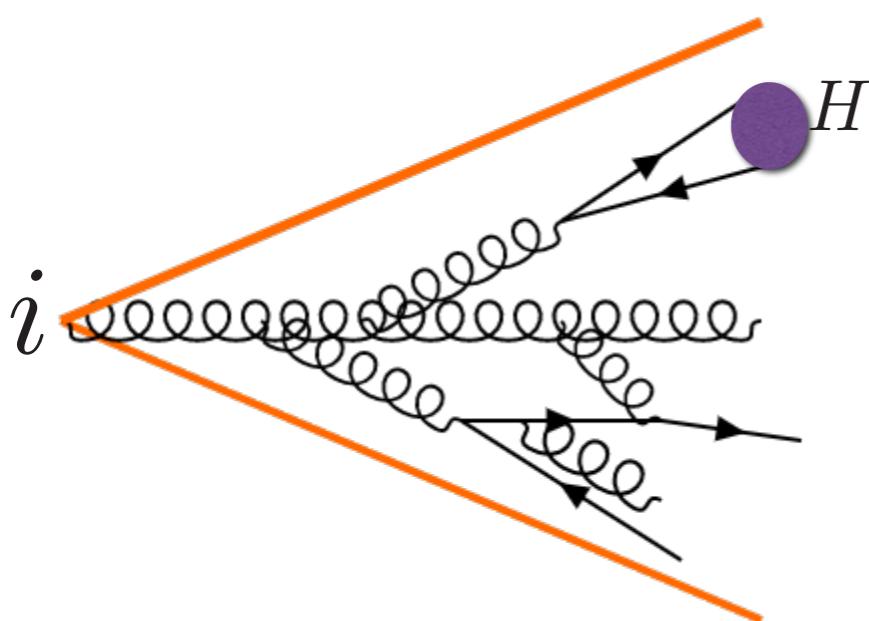
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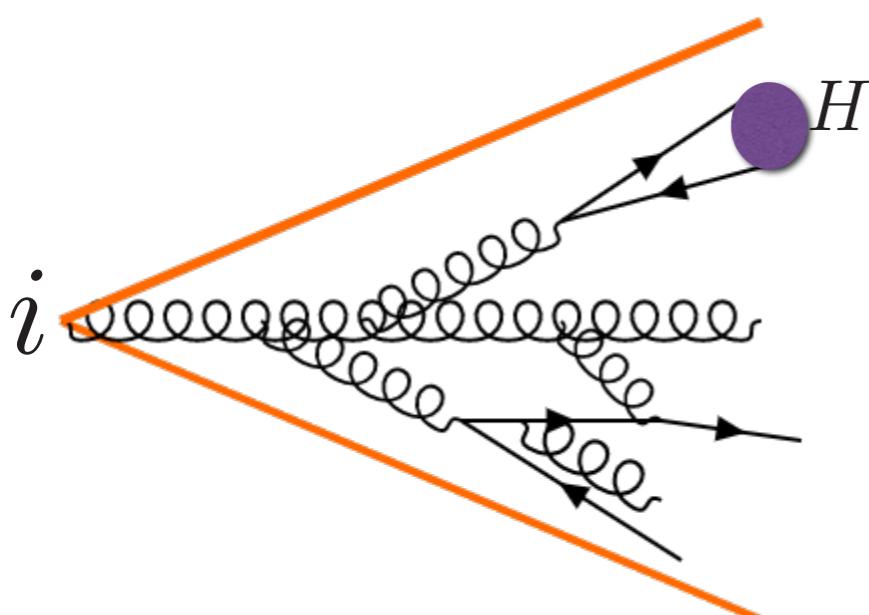
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Insert meas. Op. $\delta(s - \mathcal{S})$ for FFJ

Practical: resum $\log(\frac{m_H}{s})$

Factorize Everything

$$\sigma(H) = \sum_i \mathcal{H} \otimes \mathcal{S} \otimes \mathcal{G}_i^H \otimes \mathcal{J}_1 \otimes \mathcal{J}_2 \otimes \dots$$

$$\mathcal{G}_i^H(s, z, \mu) = \sum_j \int_z^1 \frac{dz'}{z'} \mathcal{J}_{ij}(s, z', \mu) D_j^H(z/z', m_H, \mu)$$

B meson, or J/ ψ FF.

$$e^+ e^- \rightarrow b\bar{b}$$


B Fragmentation Function

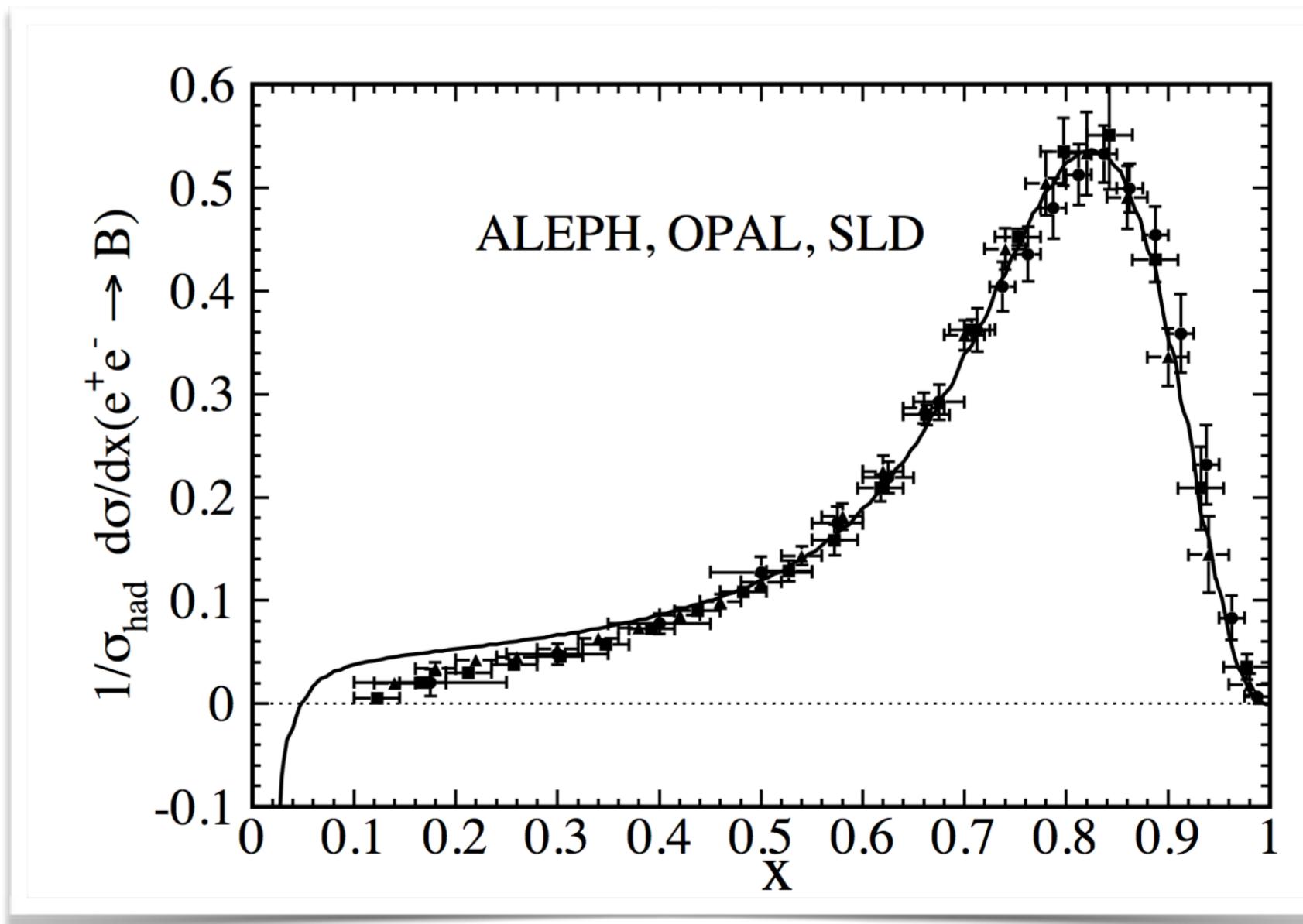
Power Model FF $D_b^B(z) = N z^\alpha (1 - z)^\beta$

- $\alpha = 16.87, \beta = 2.02$

2 parameters,
N fixed by:

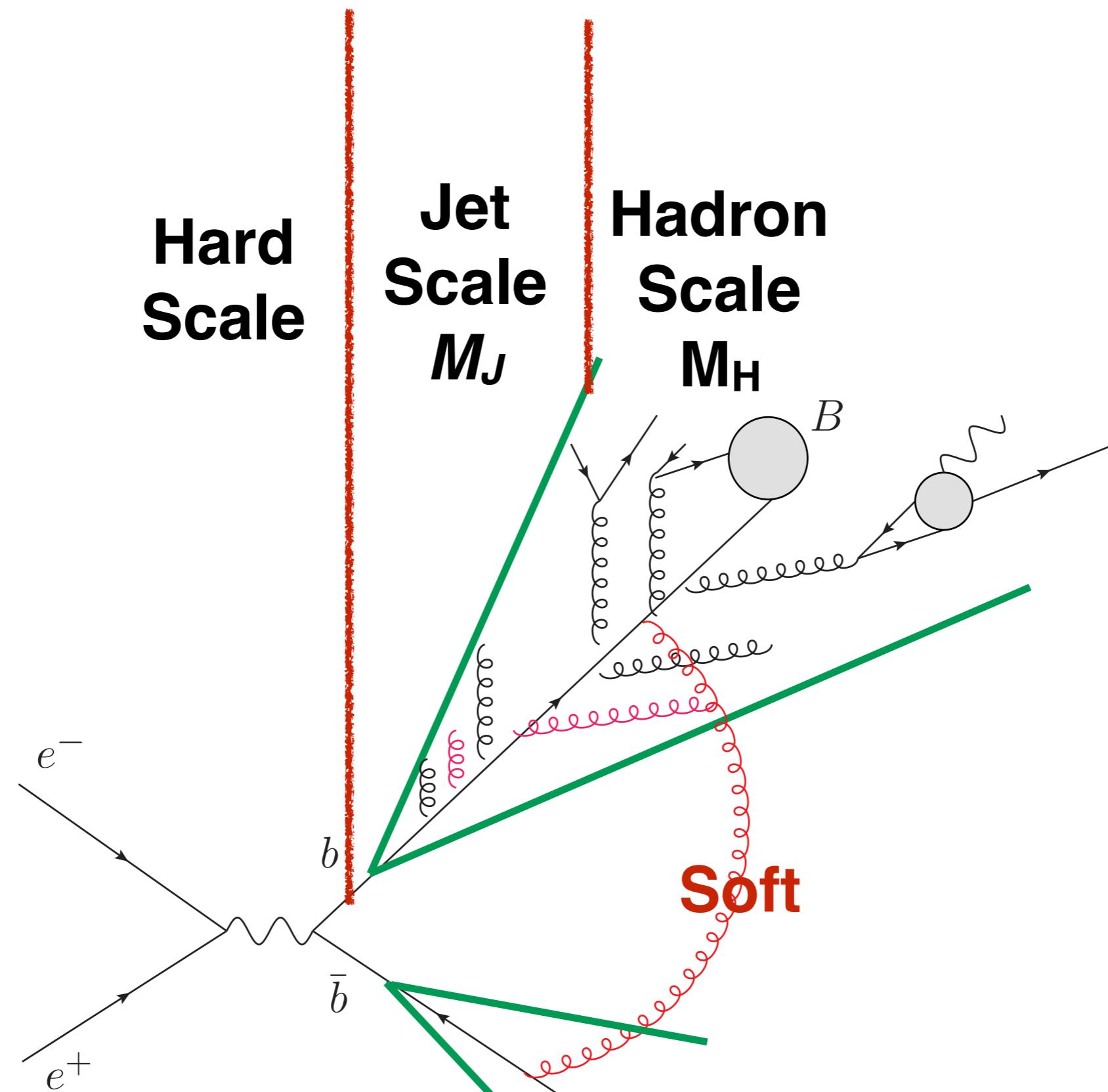
$$\int_0^1 dz z D(z) = 1$$

- Fit at m_B scale



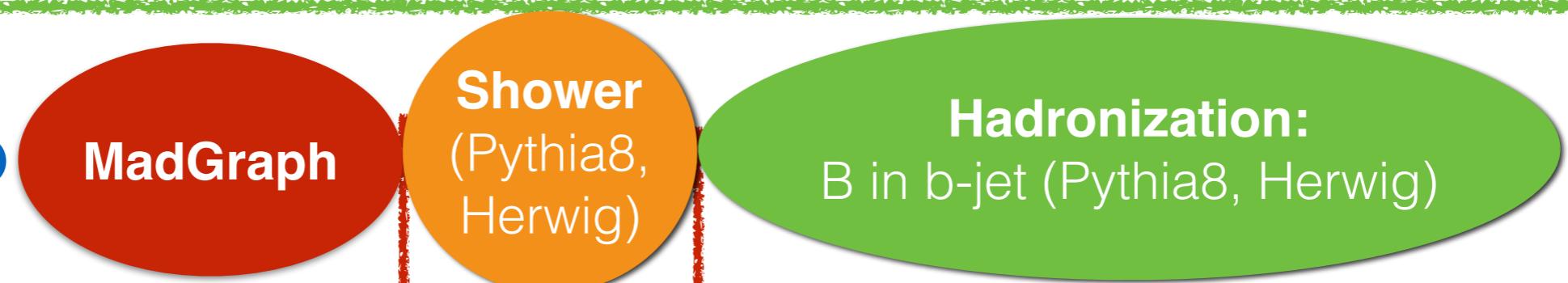
[Kniehl and Kramer, 0705.4392]

Monte Carlo & Analytic (B in jet)



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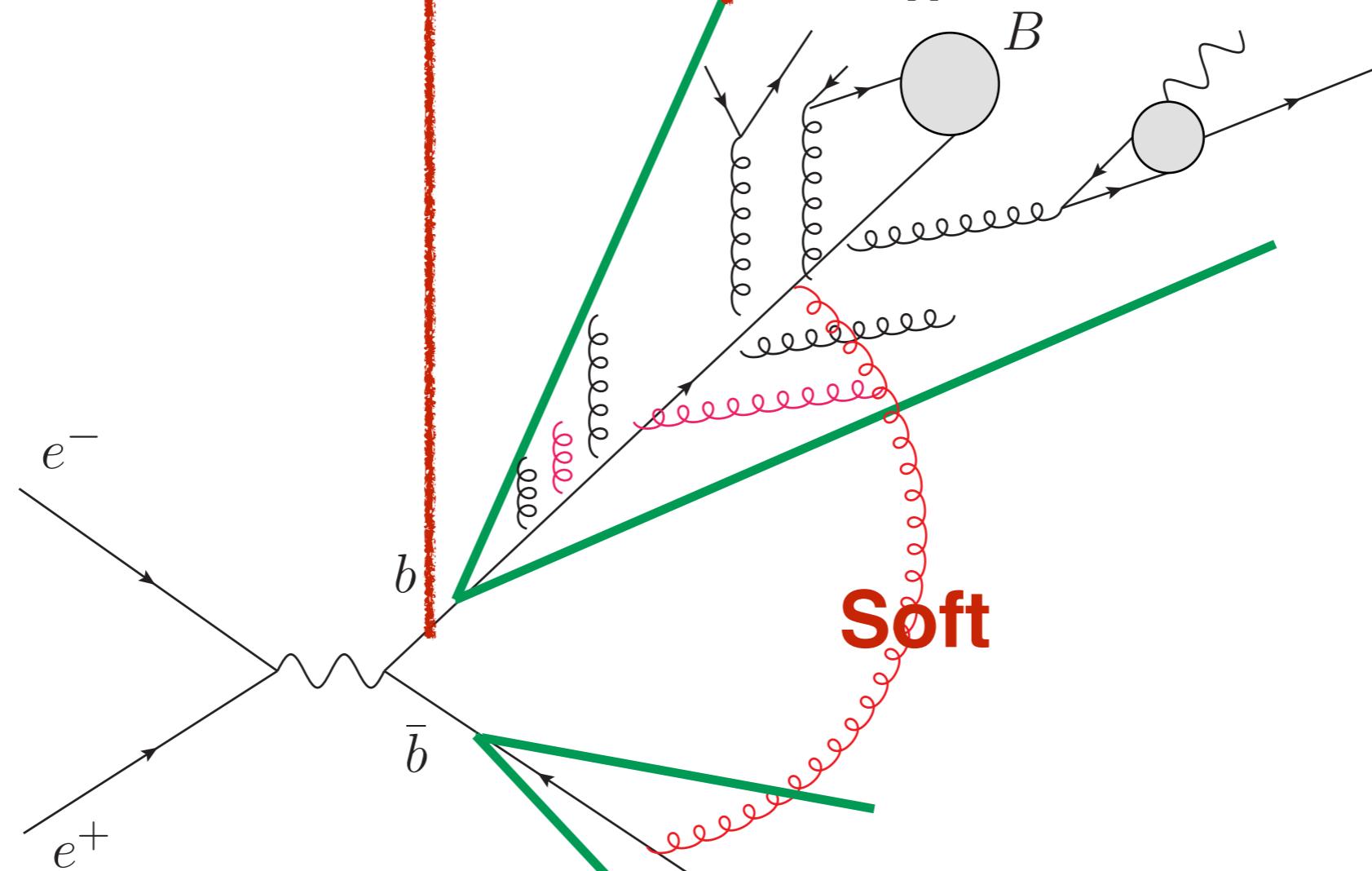
Monte Carlo



Hard
Scale

Jet
Scale
 M_J

Hadron
Scale
 M_H



Monte Carlo & Analytic (B in jet)

Monte Carlo

MadGraph

Shower
(Pythia8,
Herwig)

Hadronization:

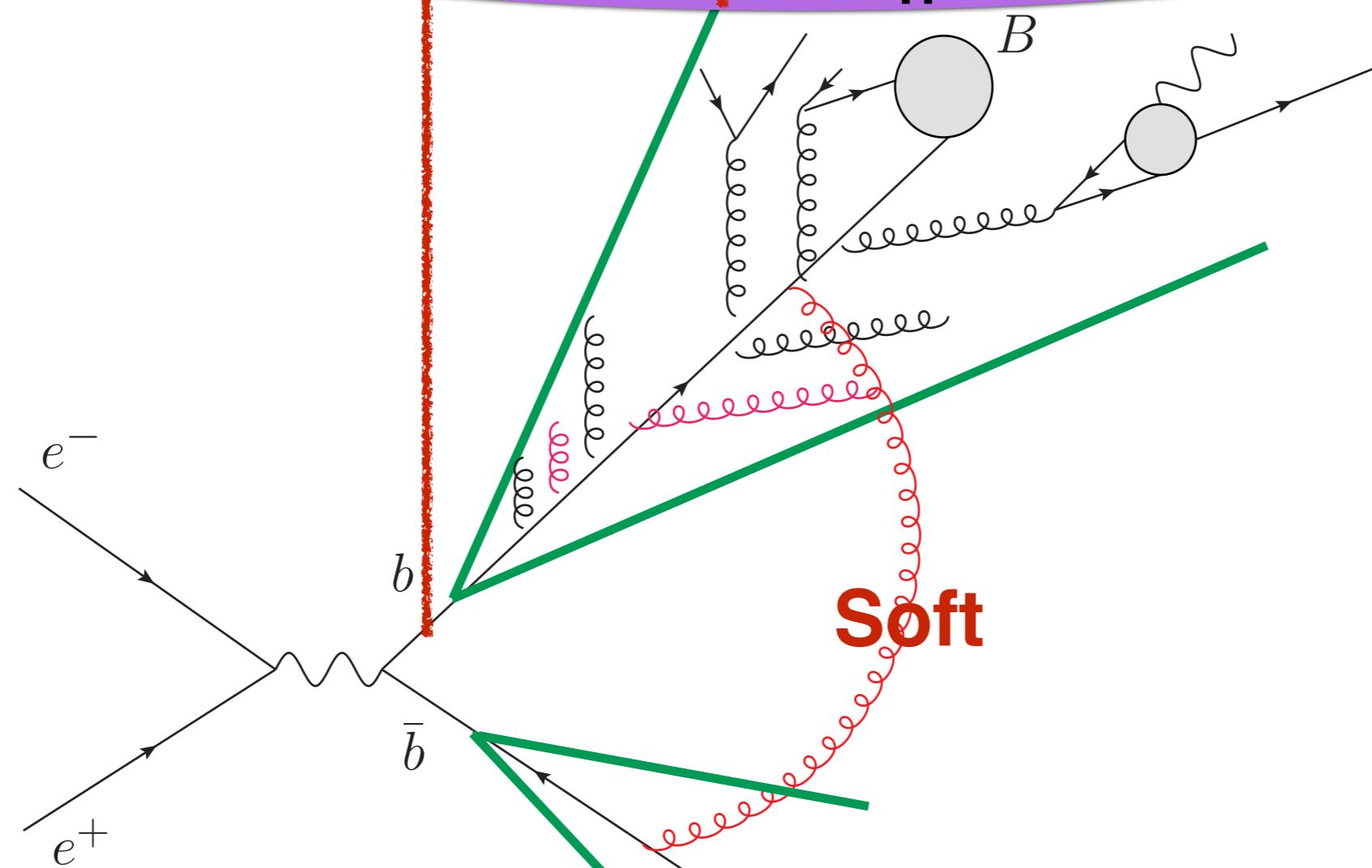
B in b-jet (Pythia8, Herwig)

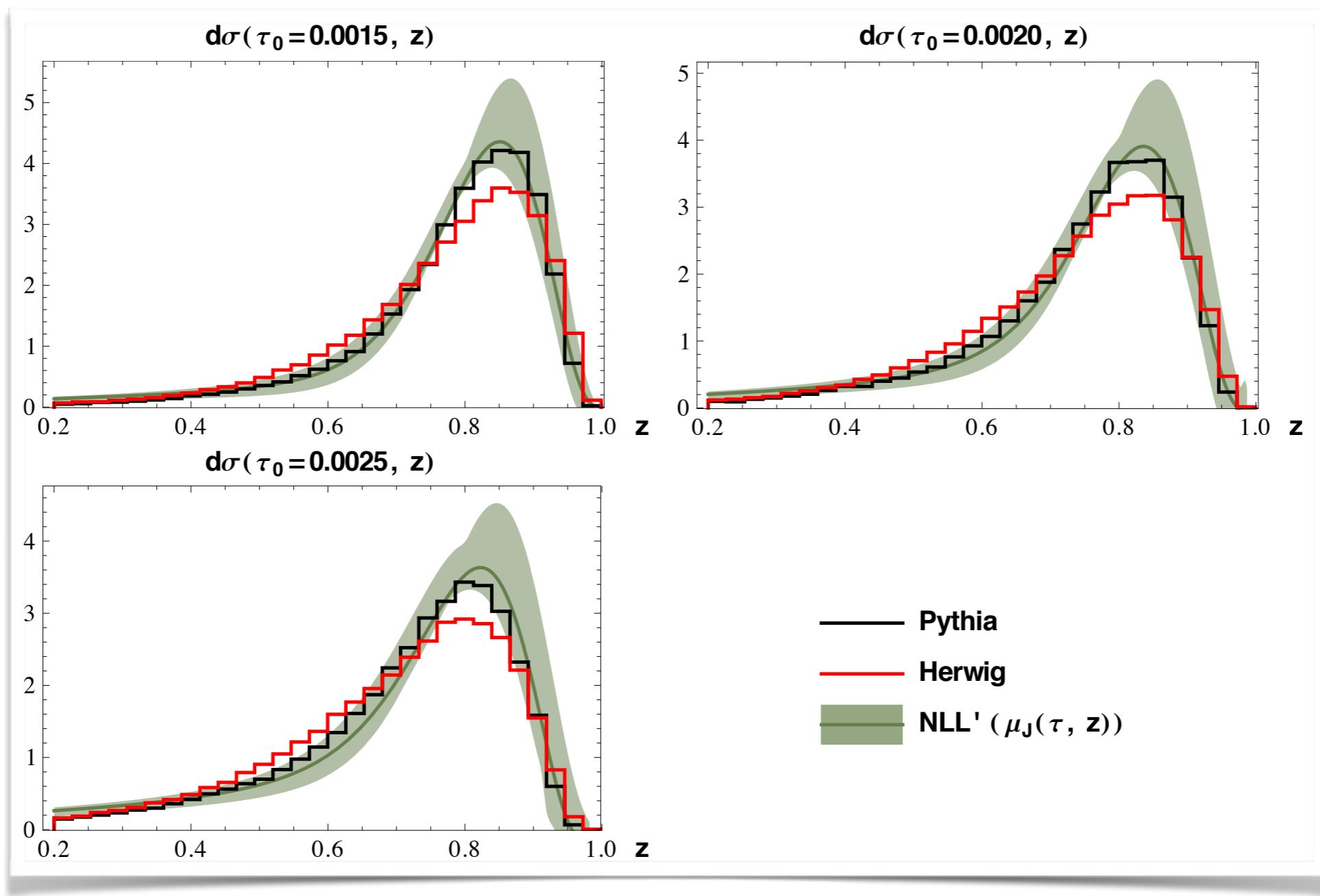
Analytic

Hard Scale

Jet Scale
 M_J

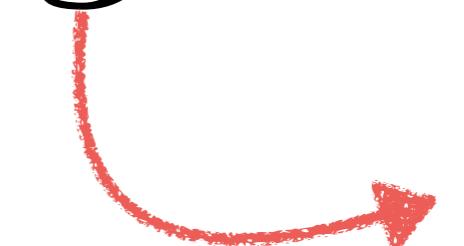
Hadron Scale
 M_H





$$z \equiv \frac{E_B}{E_J} \quad \tau_0 \propto m_J$$

FF: Pythia8 & Analytic: Power Model
 Herwig: Cluster hadronization

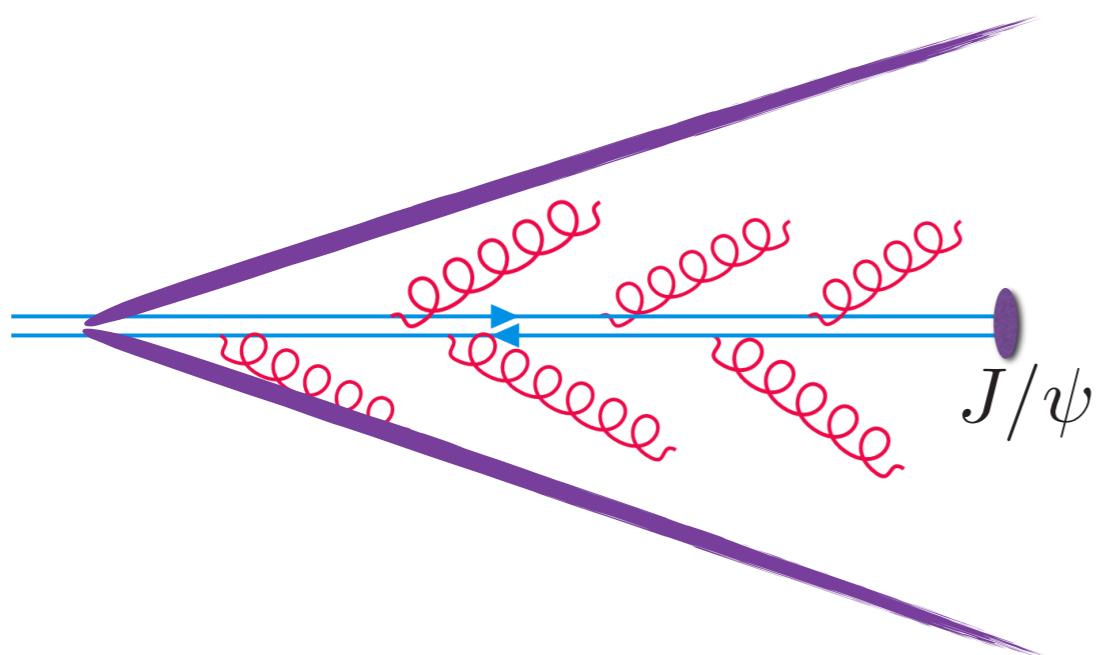
$$e^+ e^- \rightarrow b\bar{b}g$$

$$J/\psi$$

J/ ψ Production with Pythia

Pythia default:

$Q\bar{Q}(^3S_1^{[8]}, \dots)$ Produced in hard process

Pythia Shower with $2P_{qq}$

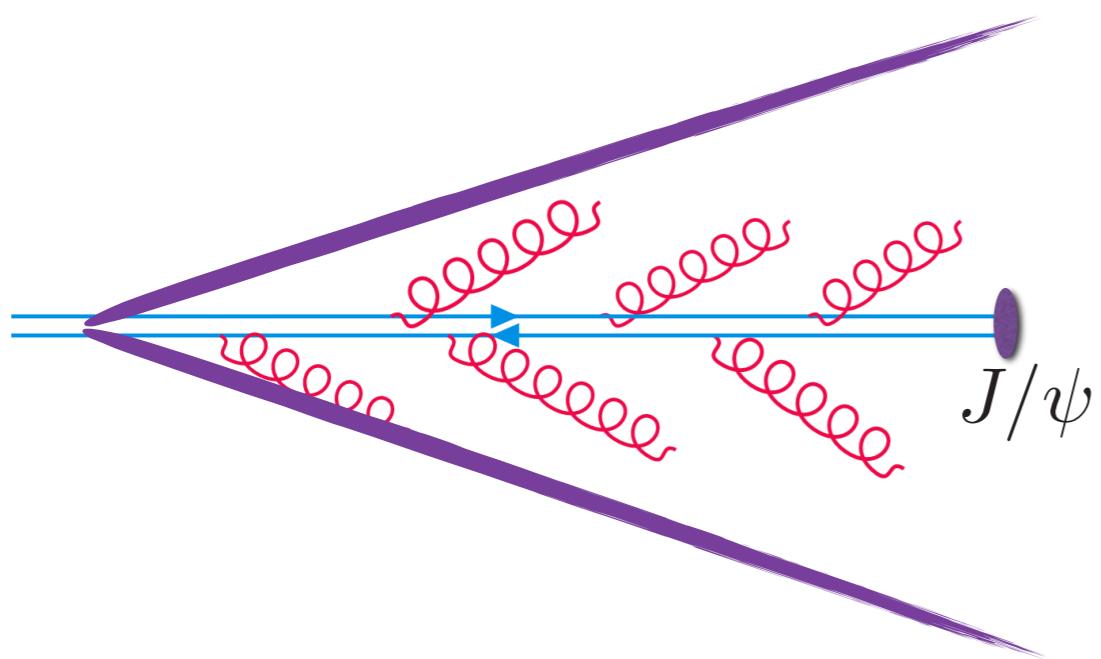


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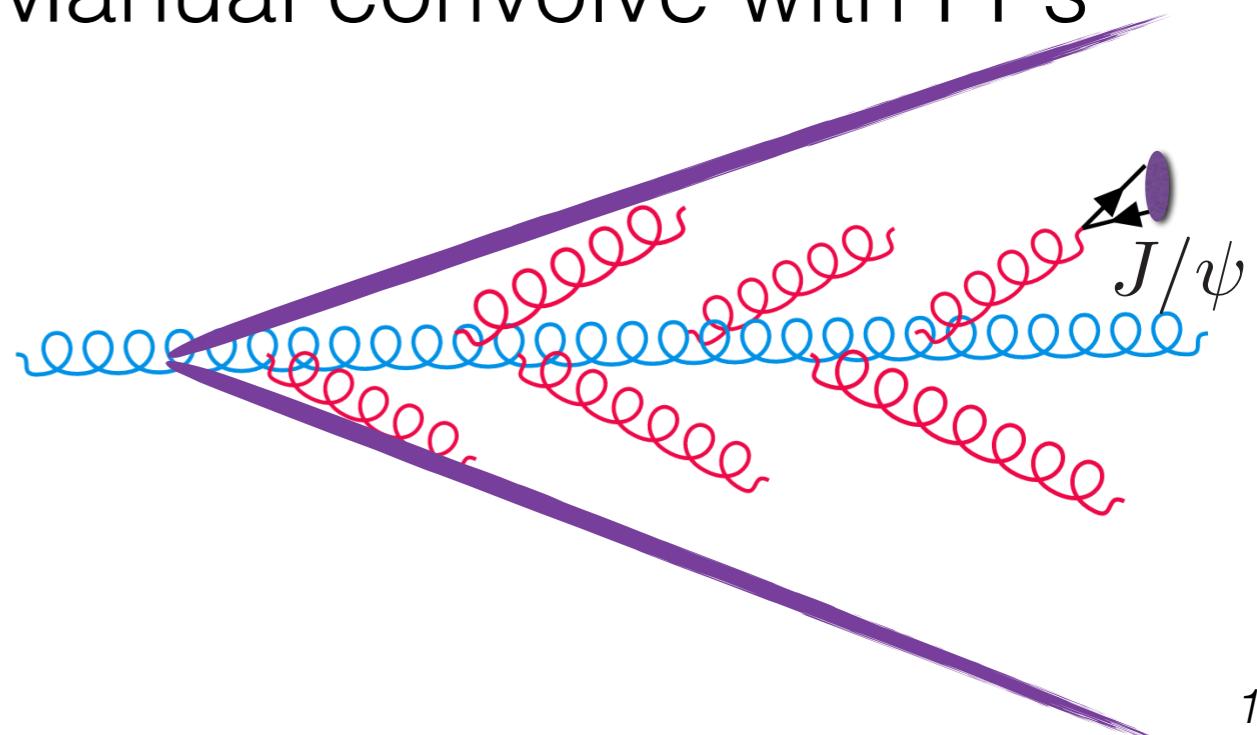


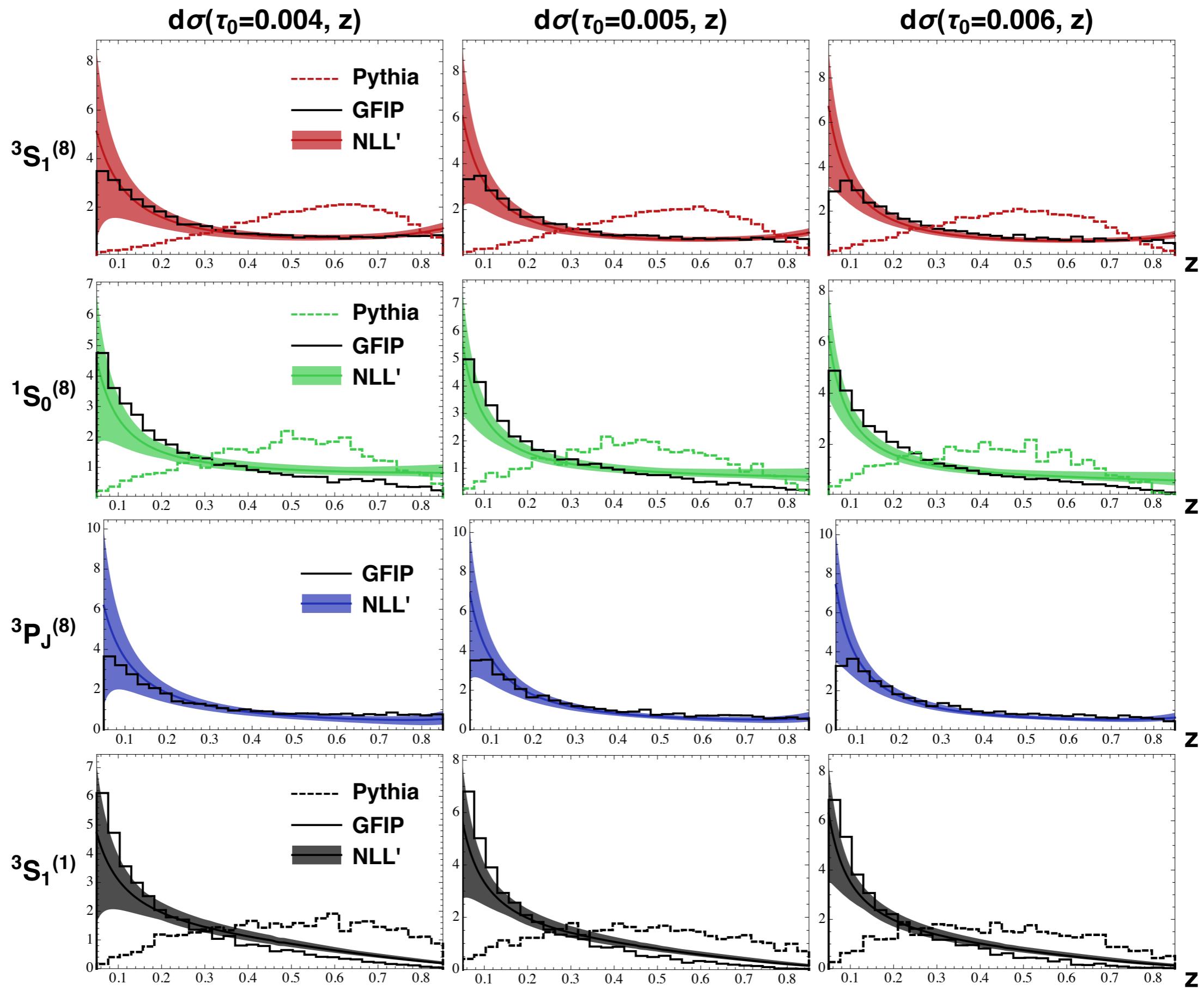
Gluon Fragmentation Improved Pythia (GFIP):

Initiating gluon produced in hard process

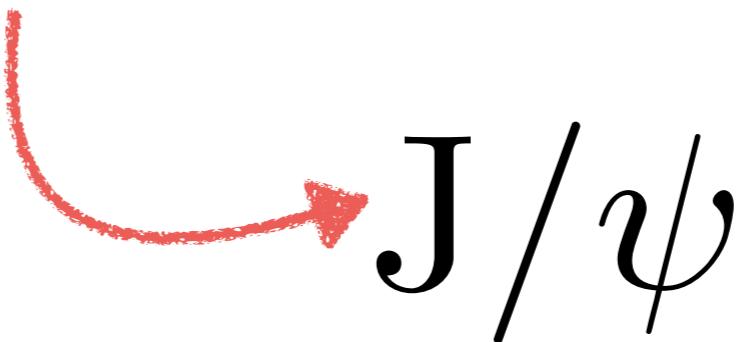
Pythia Shower
(to $2m_c$, Hadronization off)

Manual convolve with FFs





[R. Bain, L. Dai, A. Hornig, A. Leibovich, Y. Makris, T. Mehen, JHEP, 06, 121 (2016)]

$$pp \rightarrow \text{jet} + X$$


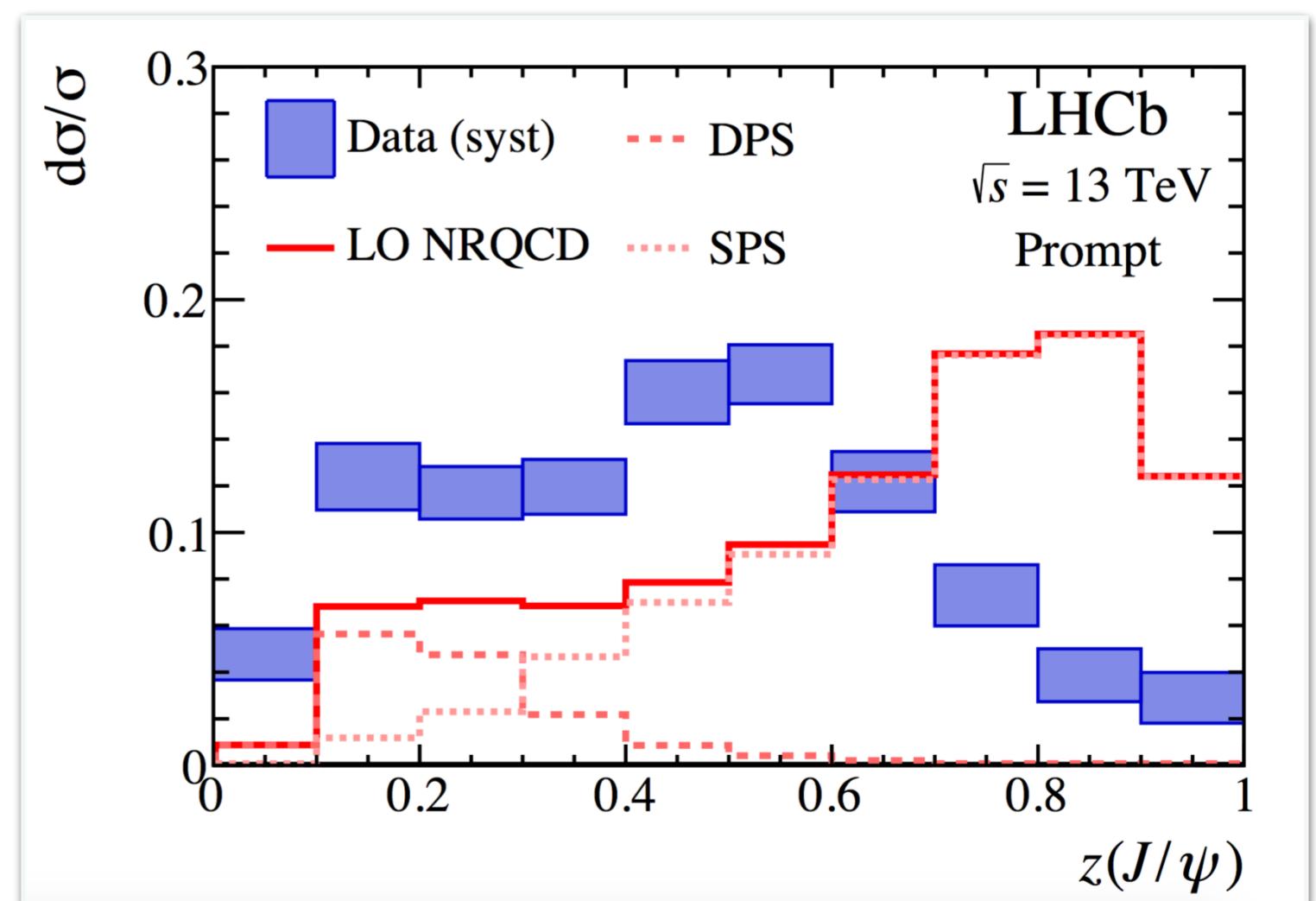
J/psi in jets @ LHCb

Cuts:

$p_T(\text{jet}) > 20 \text{ GeV}$
 $2 > \eta(\text{jet}) > 2.5$
 $p_T(\mu) > 0.5 \text{ GeV}$
 $p(\mu) > 5 \text{ GeV}$
 $4.5 > \eta(\mu) > 2$

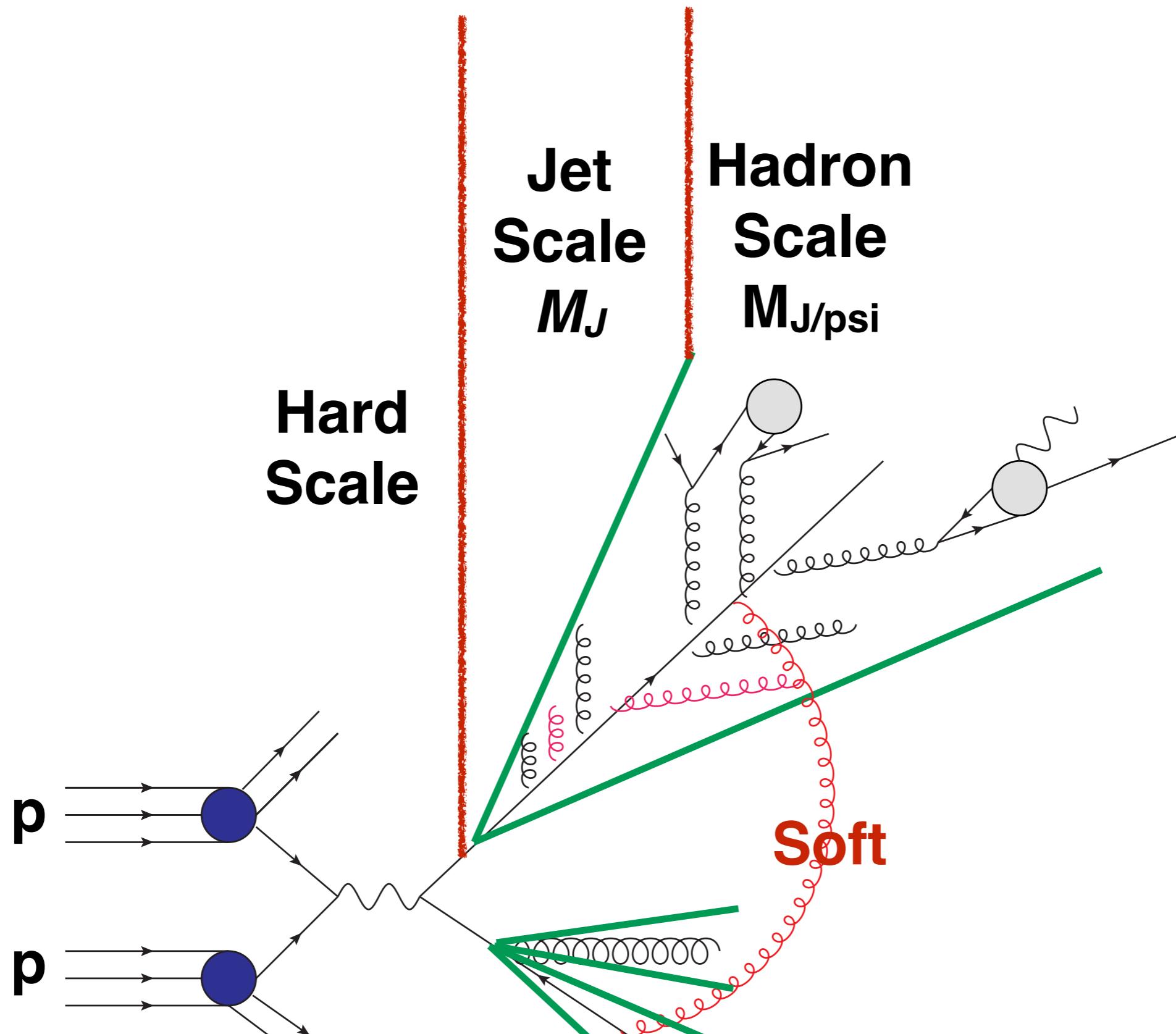
Similar to ee>J/psi:

J/psi too hard in jets

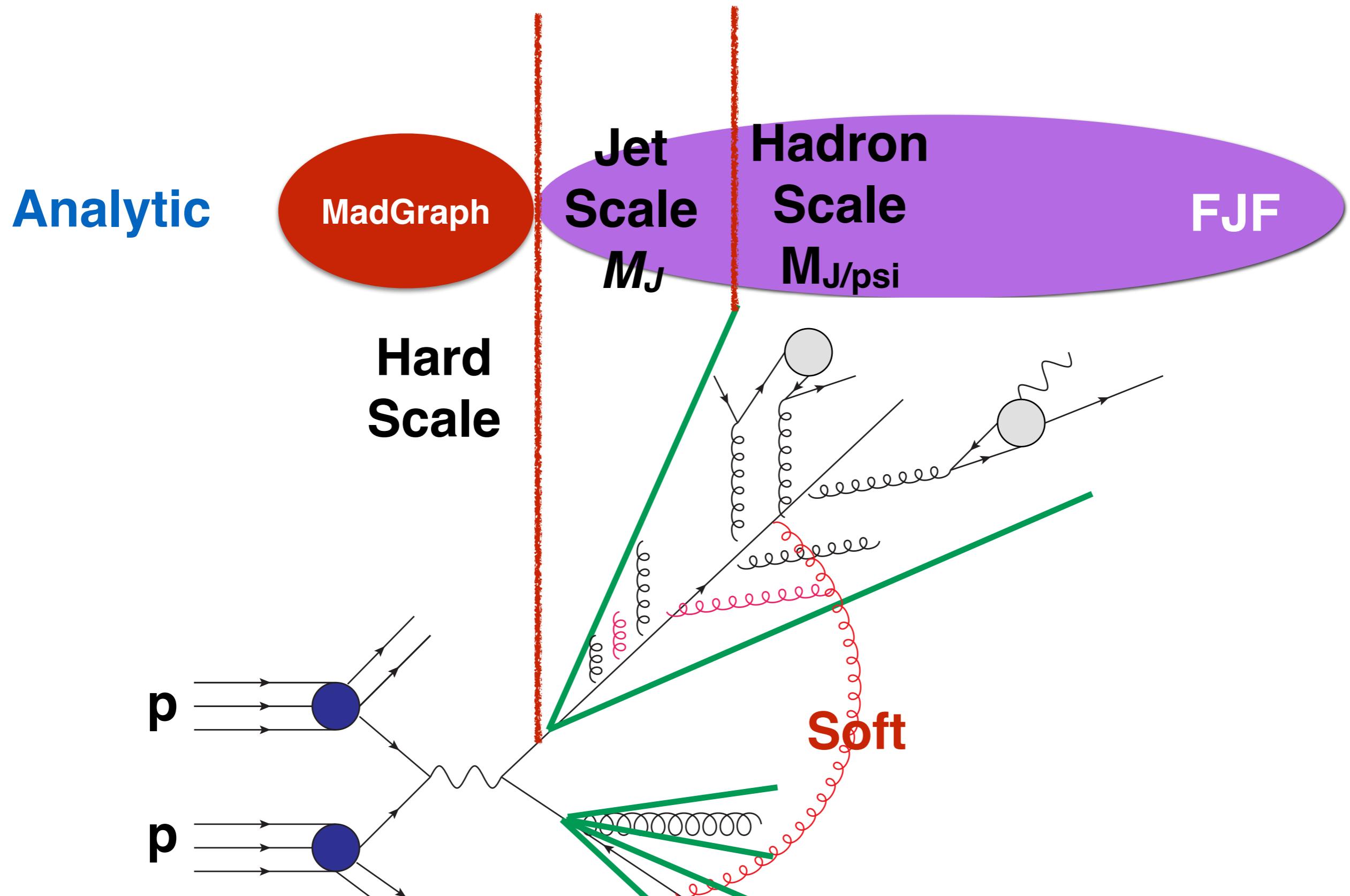


[LHCb, Phys. Rev. Lett. 118, 192001 (2017)]

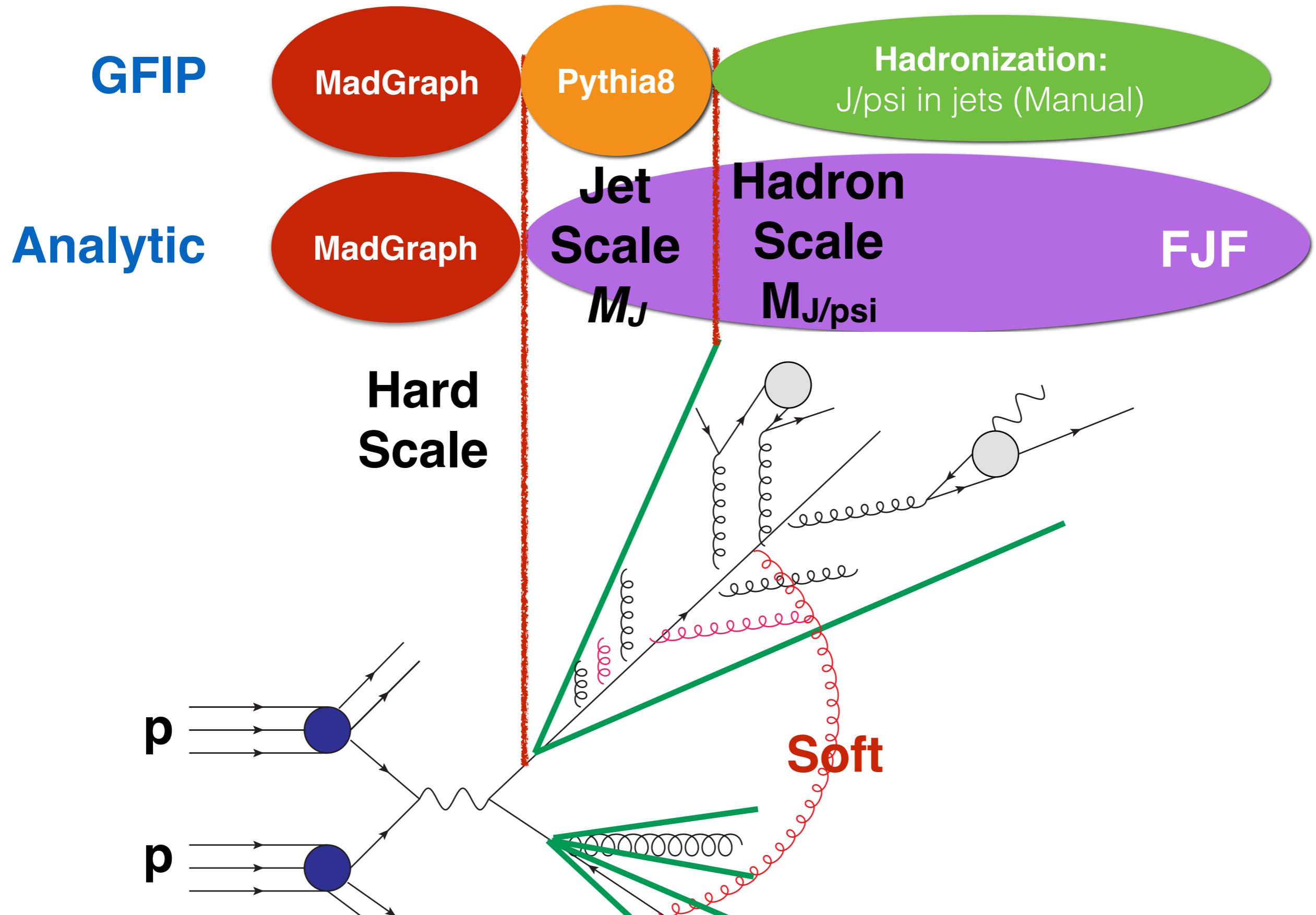
Analytic & GFIP (J/ψ in jets @ LHCb)



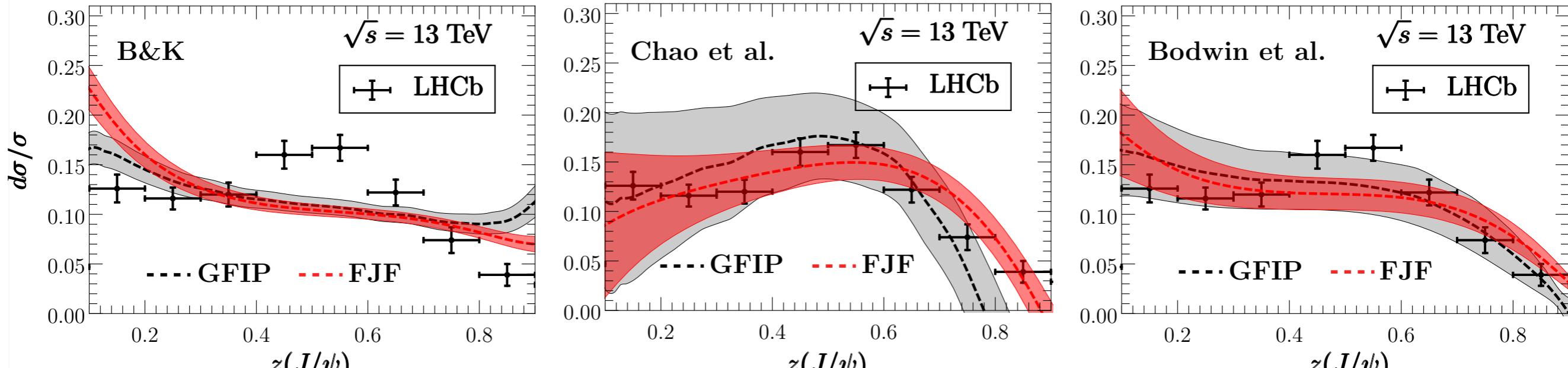
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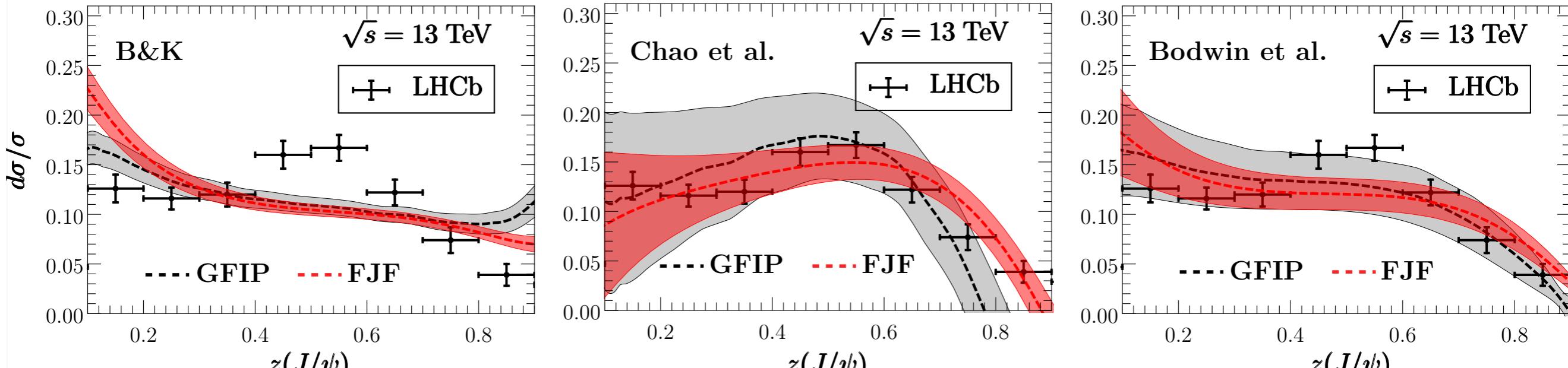
Comparison with LHCb Data



[R. Bain, L. Dai, A. Leibovich, Y. Makris, T. Mehen, Phys. Rev. Lett. 119, 032002 (2017)]

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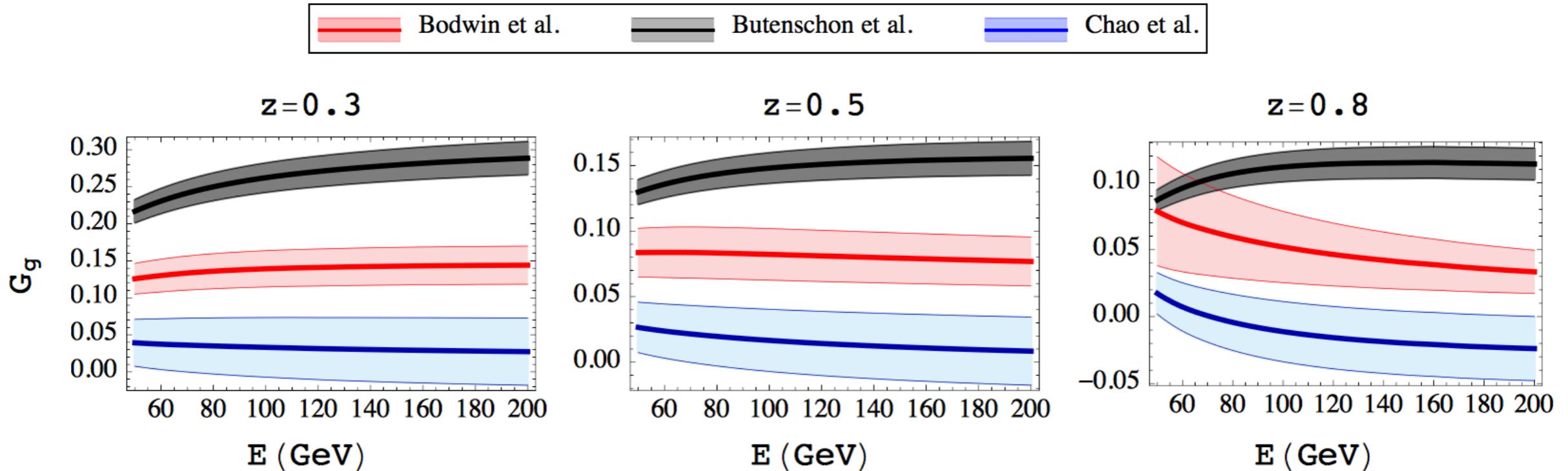


[R. Bain, L. Dai, A. Leibovich, Y. Makris, T. Mehen, Phys. Rev. Lett. 119, 032002 (2017)]

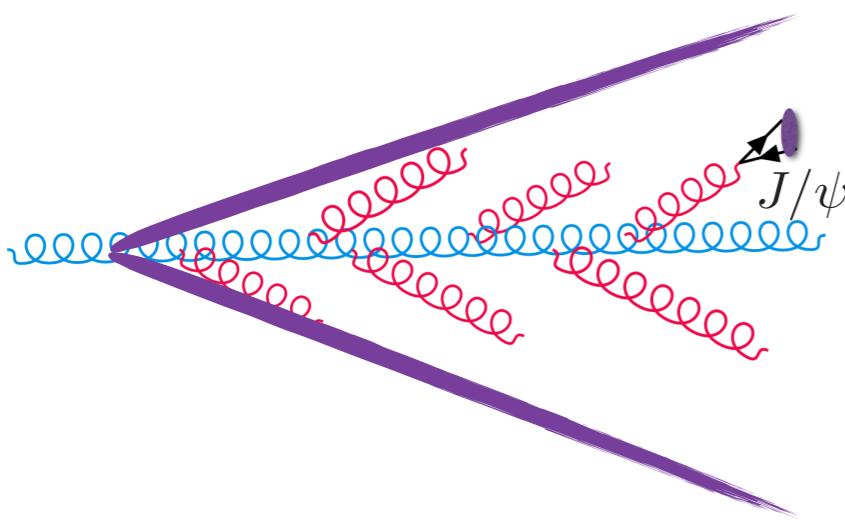
- Both FJF & GFIP work better than default Pythia;
- High pT LDMEs give better description of data than the global fit ones
- Error band only from LDMEs (B&K narrower, used more data)

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 $2 > \eta(\text{jet}) > 2.5$
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E(jet) distribution



[M. Baumgart, A. Leibovich, T. Mehen, I. Rothstein, JHEP, 11, 003 (2014)]



- ◆ E is jet energy
- ◆ LDME from Global and high-pt extractions different
- ◆ G_g is not observable...
- ◆ Evaluate at the cross section level (similar feature)

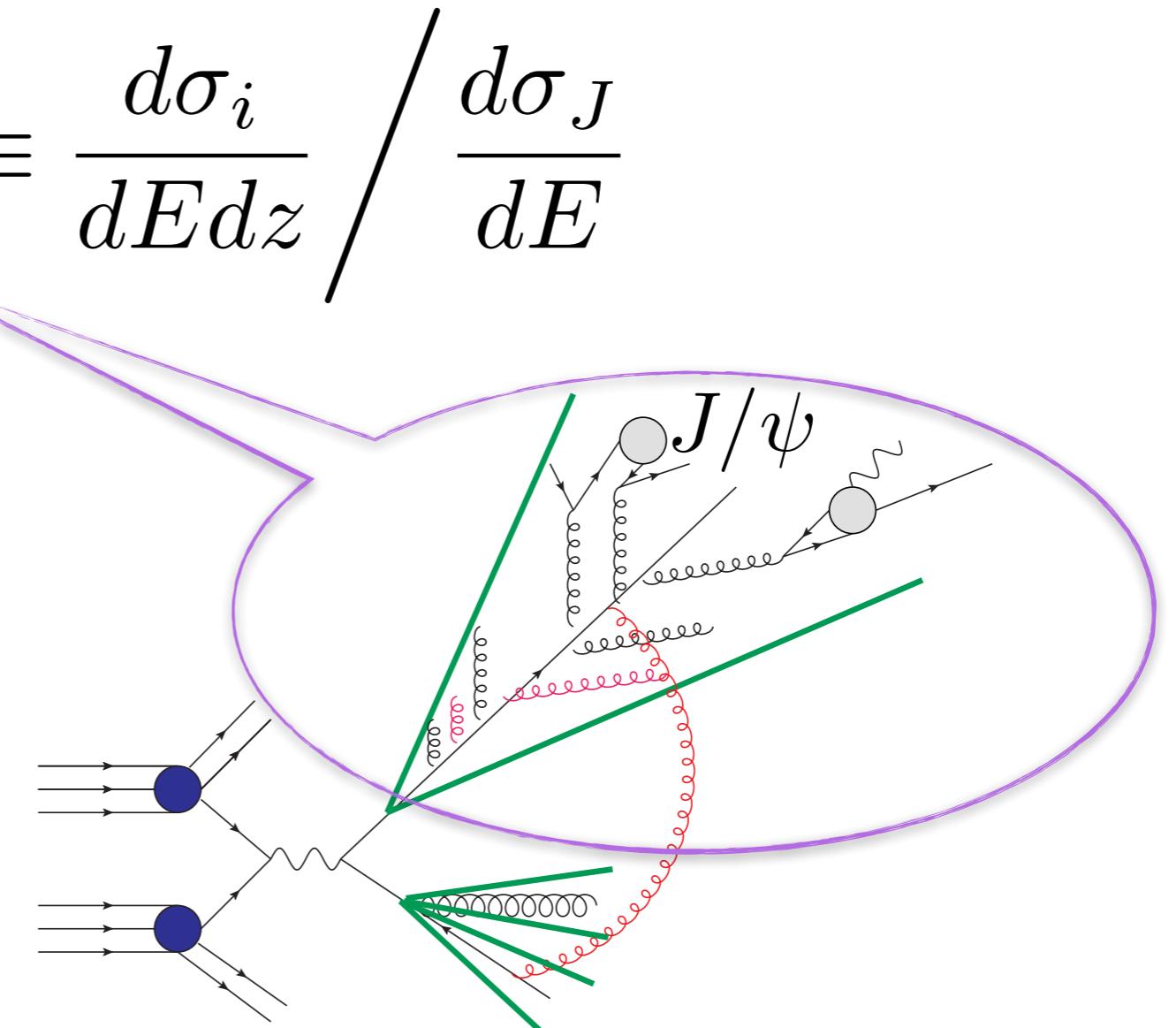
E(jet) distribution

- ♦ Normalized cross sections for the LHC (i denote production channels):

$$\frac{d\sigma_i}{dEdz} = \sum_{a,b,l,m} H_{ab \rightarrow lm} \otimes f_{a/p} \otimes f_{b/p} \otimes J_m \otimes S \otimes \mathcal{G}_l^{\psi(i)}(E, R, z, \mu),$$

$$\frac{d\hat{\sigma}_i}{dEdz} \equiv \frac{d\sigma_i}{dEdz} \Bigg/ \frac{d\sigma_J}{dE}$$

- ♦ Normalized by inclusive 1-jet cross section (σ_J)



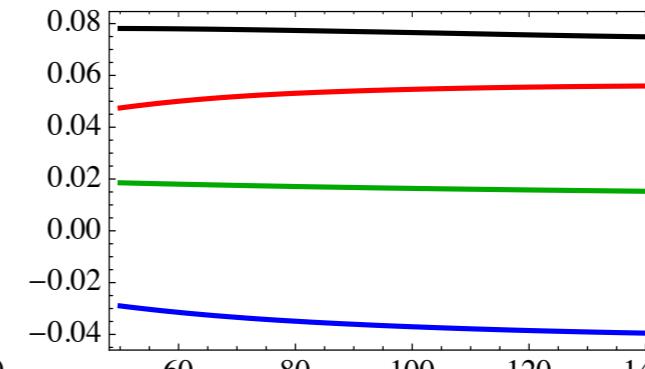
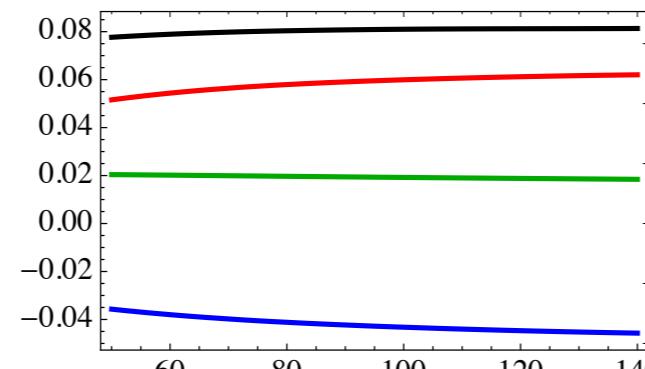
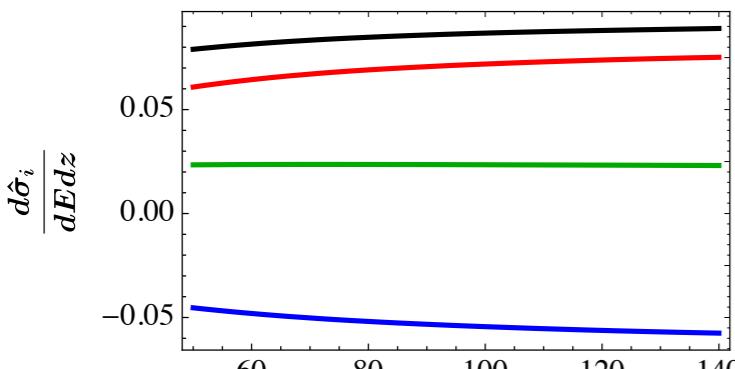


$z = 0.4$

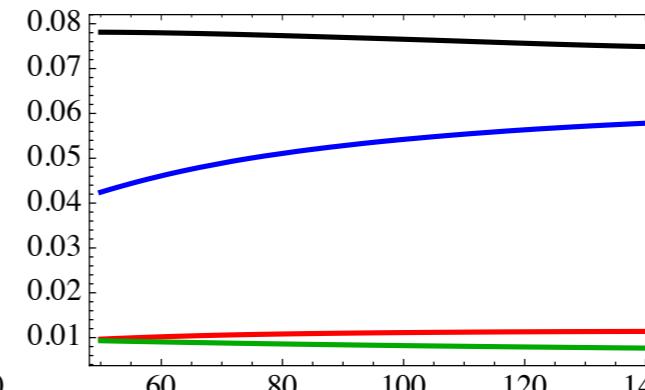
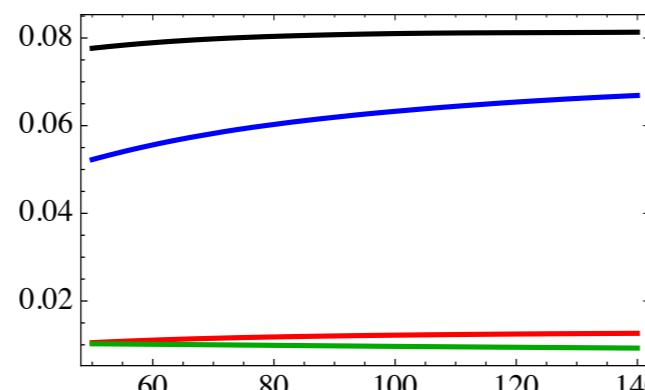
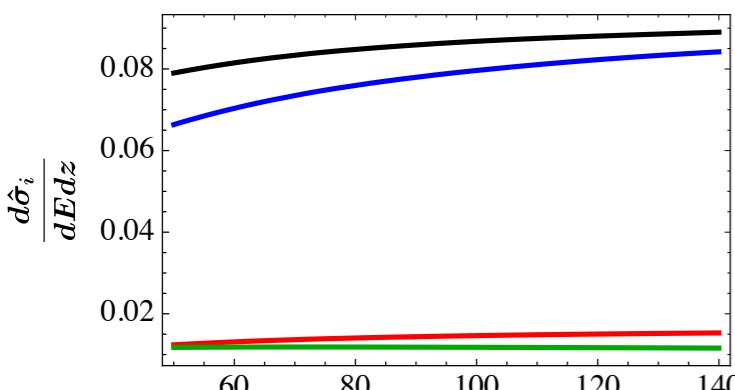
$z = 0.5$

$z = 0.6$

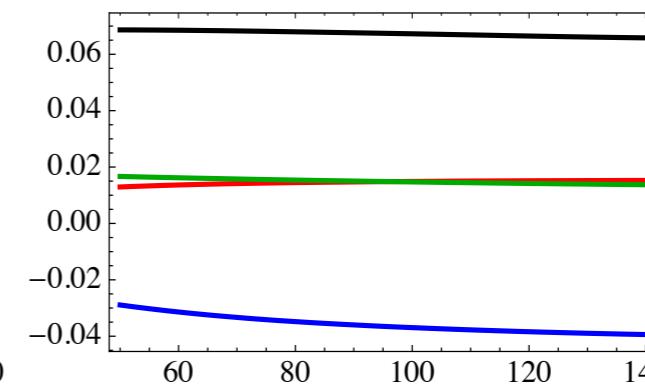
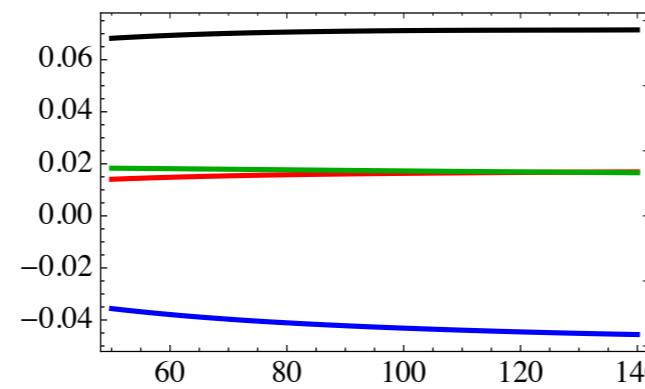
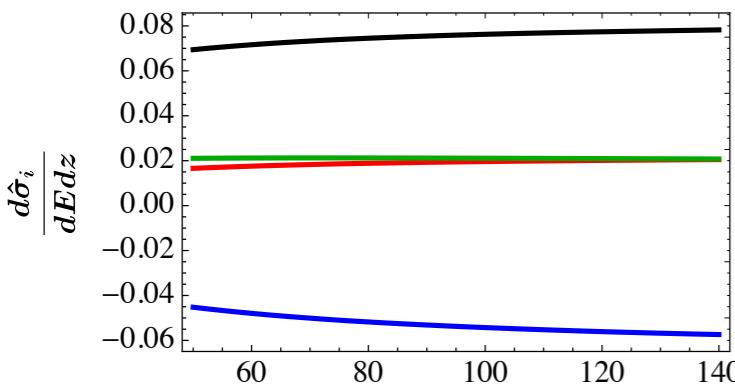
Bodwin et al.



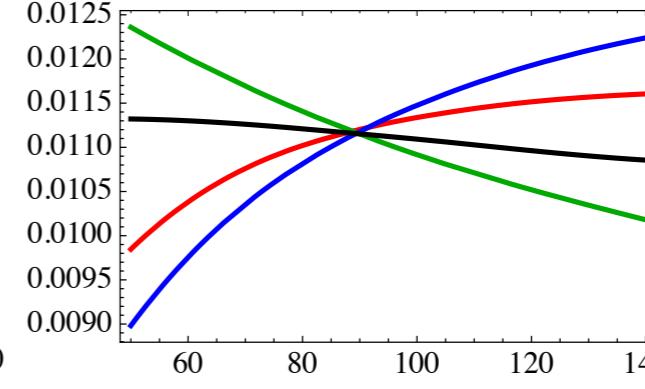
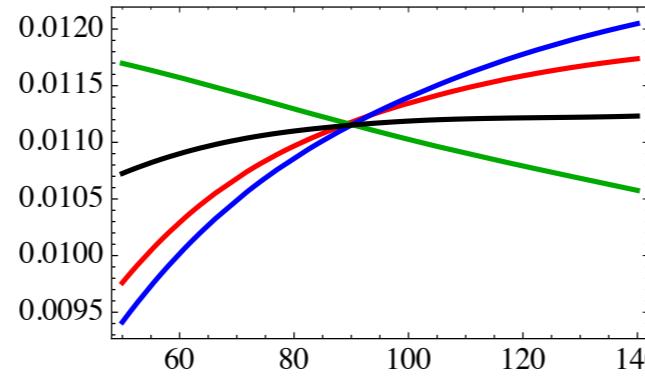
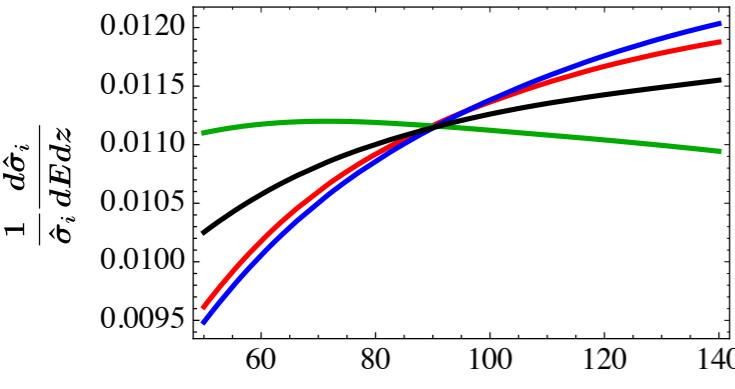
Butenschoen et al.



Chao et al.



Normalized to 1

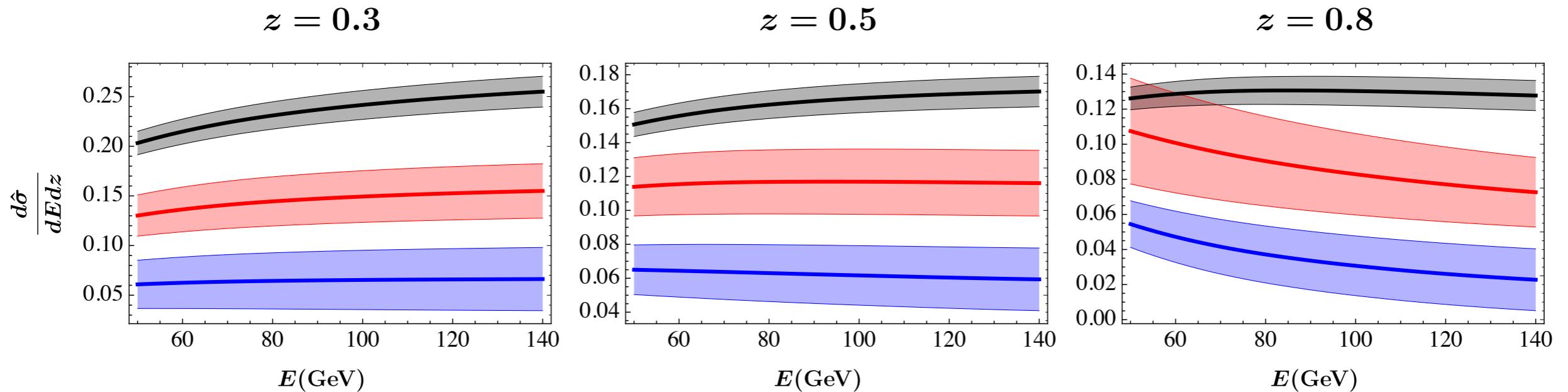


**LDME
Independent
!!!**

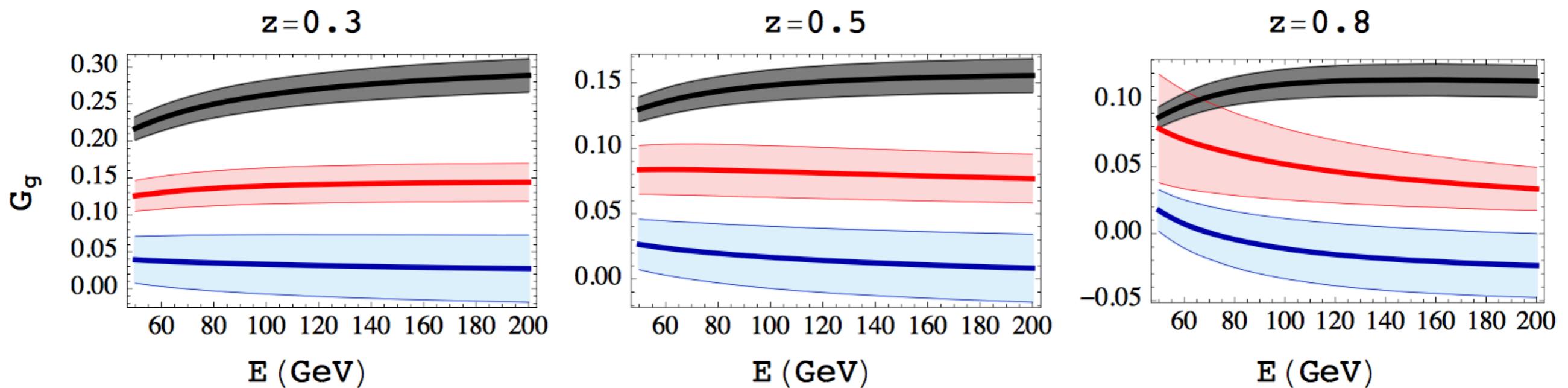
E(jet) distribution

◆ Total Cross Section:

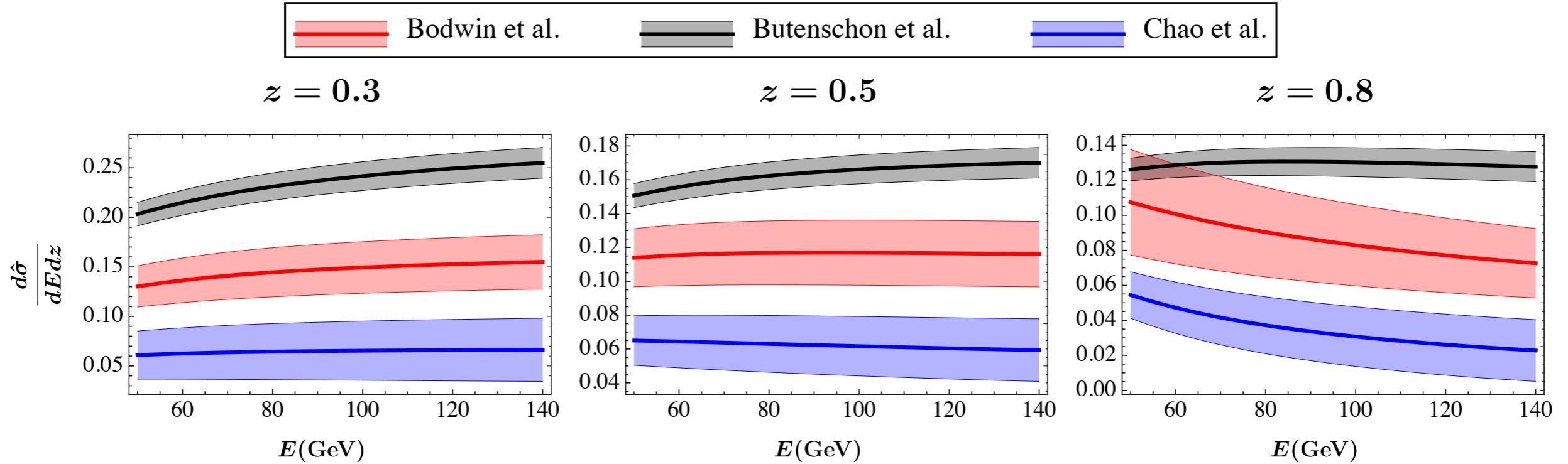
$$\frac{d\hat{\sigma}}{dEdz} \equiv \sum_i \frac{d\hat{\sigma}_i}{dEdz}$$



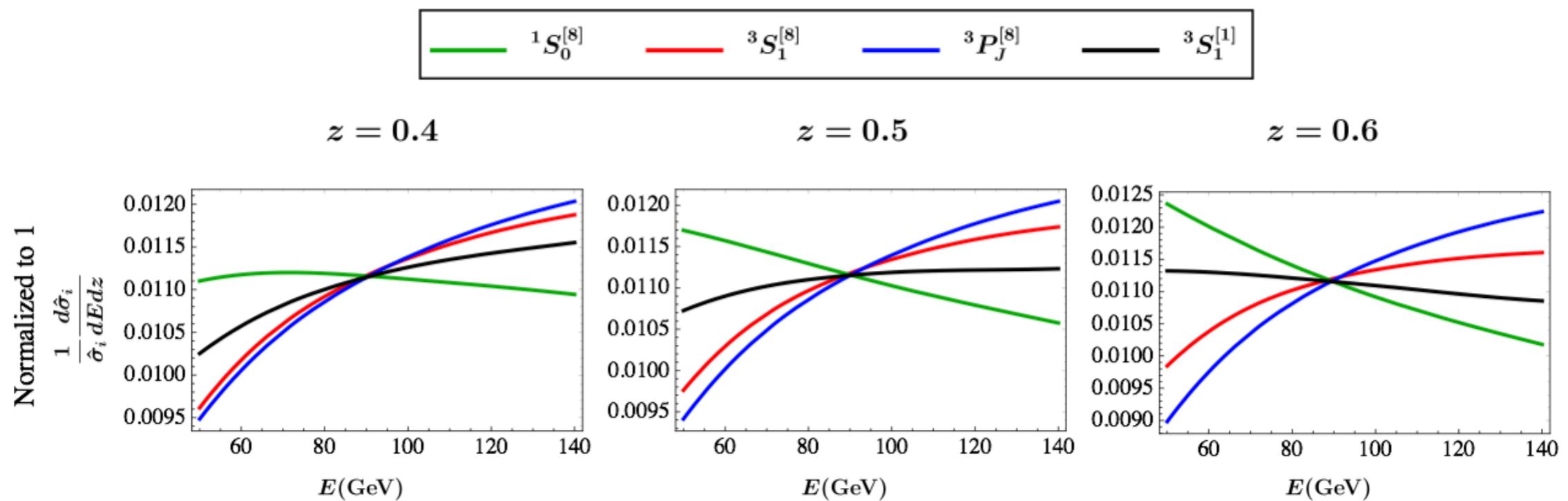
[L. Dai, P. Shrivastava, PRD 96, 036020 (2017)]



[M. Baumgart, A. Leibovich, T. Mehen, I. Rothstein, JHEP, 11, 003 (2014)] 25



Decreasing at $z > 0.5$ (high pT), consistent with the dominance of $^1S_0^{[8]}$ at high pT



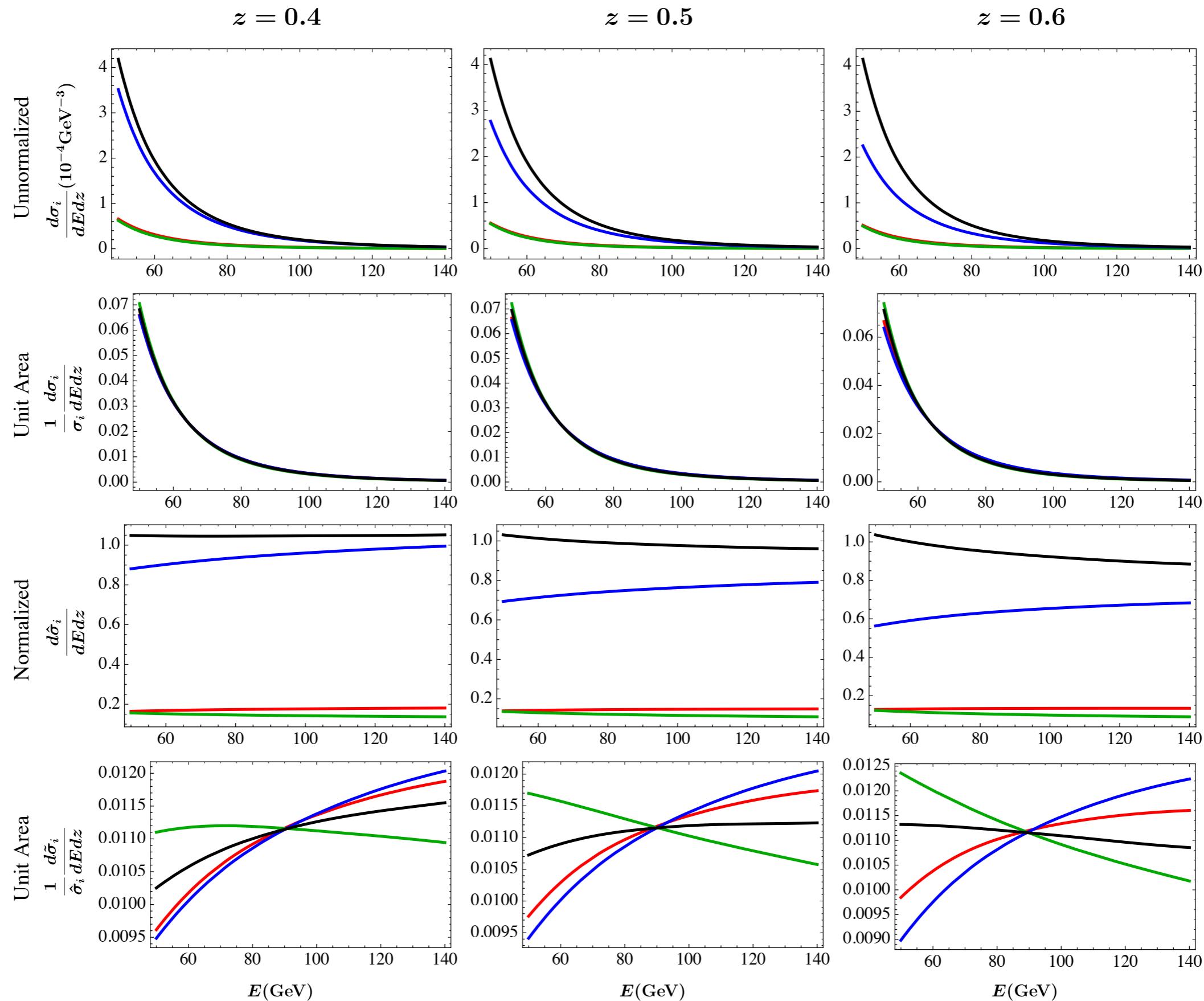
Summary

- ◆ We study quarkonium production in jets using FJFs
- ◆ Use e^+e^- collisions to test our calculations:
 - ◆ B meson production agrees well with Pythia
 - ◆ J/ ψ production, Pythia is not sufficient for our purpose
- ◆ Generalize to pp collision
 - ◆ Compare with LHCb data, high pT extractions give better description of data,
 - ◆ Study E(jet) distributions (the potential to distinguish production mechanisms and LDME extractions)

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Back Up



	$\langle \mathcal{O}^{J/\psi}(^3S_1^{[1]}) \rangle$ $\times \text{GeV}^3$	$\langle \mathcal{O}^{J/\psi}(^3S_1^{[8]}) \rangle$ $\times 10^{-2} \text{ GeV}^3$	$\langle \mathcal{O}^{J/\psi}(^1S_0^{[8]}) \rangle$ $\times 10^{-2} \text{ GeV}^3$	$\langle \mathcal{O}^{J/\psi}(^3P_0^{[8]}) \rangle / m_c^2$ $\times 10^{-2} \text{ GeV}^3$
B & K	1.32 ± 0.20	0.224 ± 0.59	4.97 ± 0.44	-0.72 ± 0.88
Chao et al.	1.16 ± 0.20	0.30 ± 0.12	8.9 ± 0.98	0.56 ± 0.21
Bodwin et al.	1.32 ± 0.20	1.1 ± 1.0	9.9 ± 2.2	0.49 ± 0.44