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Higgs Flavor-violating Decay Mediated by Flavored Dark Matter

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Outline

SM: Higgs & flavor

Beyond SM: Higgs flavor violating decay

HLFV & lepton flavored DM

Conclusion & outlook

SM: Higgs & flavor

Higgs may be a key to understand new physics, it is

Related to neutrino mass origin

Related to dark matter

Related to baryon asymmetry

Related to neutrino mass origin

In the extended theory to establish the above connections, Higgs properties are likely to deviate from the SM predictions

Thus precision of Higgs is crucial to hunt for new physics hints

Flavor changing decays are of particular interest, considering their great suppression in SM

SM: Higgs & flavor

Higgs (one doublet) & quarks (3 generations, chiral, 12 flavors)

Simultaneously diagonalizing the mass and Higgs Yukawa interactions



No flavor changing neutral current (FCNC) mediated by Higgs boson

At loop level? GIM-suppression in the charge flavor-violating loop!



S. Bejar, F. Dilme, J. Guasch, and J. Sola, JHEP 08, 018 (2004)

SM: Higgs & flavor

Higgs & leptons (3 generations, chiral, missing RHN)

Again, no Higgs mediated lepton flavor decay

Neutrinos masses are strong evidence of new physics!!

Neutrino in the SM are massless, due to the absence of right-handed neutrinos

However, neutrinos are found to have tiny mass < 10⁻¹¹ GeV!

▶ Higgs can still account for their masses, via the d=5 Weinberg operators:

Leading to Majorana neutrinos

Beyond SM: Higgs flavor violating decay

Higgs violating decay: effective Lagarangian

► Higgs mediated FCNC in the quark sector VS Δ F=2 bounds

Meson-antimes	son mixing & CPV	Br(h \rightarrow q _i q _j)<10 ⁻³ ! Too small		G. Blankenburg, etc., PLB, 712 (2012) 386–390
Operator	Eff. couplings	95% C.L. Bound		Observables
		c _{eff}	$ \text{Im}(c_{\text{eff}}) $	
$(\overline{s}_R d_L)(\overline{s}_L d_R)$ $(\overline{s}_R d_L)^2, \ (\overline{s}_L d_R)^2$	$c_{sd}c_{ds}^*$ c_{ds}^2, c_{sd}^2	$1.1 imes 10^{-10}$ $2.2 imes 10^{-10}$	$\begin{array}{c} 4.1 \times 10^{-13} \\ 0.8 \times 10^{-12} \end{array}$	Δm_K ; ϵ_K
$(\bar{c}_R u_L)(\bar{c}_L u_R)$ $(\bar{c}_R u_L)^2, \ (\bar{c}_L u_R)^2$	$\begin{array}{c} c_{cu}c_{uc}^{*}\\ c_{uc}^{2}, \ c_{cu}^{2} \end{array}$	$0.9 imes 10^{-9} \ 1.4 imes 10^{-9}$	$1.7 imes 10^{-10}$ $2.5 imes 10^{-10}$	Δm_D ; $ q/p , \phi_D$
$(\bar{b}_R d_L)(\bar{b}_L d_R)$ $(\bar{b}_R d_L)^2, \ (\bar{b}_L d_R)^2$	$egin{array}{llllllllllllllllllllllllllllllllllll$	$0.9 imes 10^{-8}$ $1.0 imes 10^{-8}$	$\begin{array}{c} 2.7 \times 10^{-9} \\ 3.0 \times 10^{-9} \end{array}$	$\Delta m_{B_d}; S_{B_d \to \psi K}$
$(\bar{b}_R s_L)(\bar{b}_L s_R)$ $(\bar{b}_R s_L)^2, \ (\bar{b}_L s_R)^2$	$c_{bs}c_{sb}^*$ c_{sb}^2, c_{bs}^2	$2.0 imes 10^{-7}$ $2.2 imes 10^{-7}$	$2.0 imes 10^{-7}$ $2.2 imes 10^{-7}$	Δm_{B_s}

Beyond SM: Higgs flavor violating decay

► The left corner: Higgs lepton flavor violating (HLFV) decay

anomalous magnetic moments and electric dipole moments of charged leptons for i=j

charged lepton flavor violating decay (CLFV) for i≠j, leading to individual upper bounds: Related to the (g-2)µ puzzle?



In 2012 with 4.7/fb of 7TeV ATLAS data gave $Br(h \rightarrow \tau \mu/e) \leq 10\%$ In 2014 with 19.7/fb of 8 TeV CMS data gave upper bound & best fit (2.4 σ excess) at 95% C.L: $Br(h \rightarrow \tau \mu) \leq 1.57\%$ & $Br(h \rightarrow \tau \mu) \approx 0.84\%$ Prospects 300/fb of 13 TeV LHC: $Br(h \rightarrow \tau \mu) \leq 0.01\%$!

J. Kopp and M. Nardecchia, JHEP 1410, 156 (2014) ; Y. n. Mao and S. h. Zhu, arXiv:1505.07668

Beyond SM: Higgs flavor violating decay

▶ LFV: already observed in active neutrino oscillation such as $v_{\mu} \rightleftharpoons v_{e}$

Can LFV leave hints in other processes?

for example, LFV insertion in charged lepton-flavor-violating (CLFV) decay

JMarciano, W., and A. Sanda, PLB 67 (3), 303, 1977

 Very clear probe to new physics, in particular neutrino mass origin

> J. Bjorken and S. Weinberg, Phys.Rev.Lett. 38, 622 (1977).

Currently, the most stringent bounds: Br($\mu \rightarrow e\gamma$)<10⁻¹³, Br($\tau \rightarrow \mu\gamma$)<4.3×10⁻⁸

Plights of weakly interacting massive particle (WIMP) dark matter

PandaX, LUX, XEON etc., are approaching the limit of DM direct detection, via DMnucleon spin-independent (SI) scattering signals, from interactions: (DM)²(quark)²

Indirect detections from cosmic ray such as γ also begins to rule out typical WIMP with typical annihilation cross section(~1pb) today



Seungwon Bae<mark>k and ZF</mark> Kang, JHEP, 1603 (2016) 106

A way out the plight: leptonpllic dark matter

DM mainly interacts with SM through the lepton sector, thus no (treelevel) DM-nucleon scattering, e.g., **via a leptonic Z' or Higgs portal or,**

> A natural realization: lepton flavored DM (Majorana χ with Z_2)

$$\begin{aligned} -\mathcal{L} = & \mathcal{L}_{\mathrm{SM}} + \left(-y_{La} \bar{l}_a P_R \chi \widetilde{\phi}_{\ell} + y_{Ra} \bar{e}_a P_L \chi \phi_e + h.c. \right) \\ & + \left(-\mu H^{\dagger} \widetilde{\phi}_{\ell} \phi_e^* + c.c. \right) + \lambda_{-1} |\phi_e|^2 |\phi_{\ell}|^2 + \lambda_0 |H|^2 |\phi_e|^2 + V_{2\mathrm{HDM}}, \end{aligned}$$

G. Blankenburg, etc., PLB, 712 (2012) 386–390

mediator quantum numbers specify interactions of the Majorana DM dominantly with leptons Ma's radiative neutrino model (RHN+Higgs doublet) is a "derivation" of leptonflavored DM

Seungwon Ba<mark>ek and ZF</mark> Kang, JHEP, 1603 (2016) 10<mark>6</mark>

> The singlet mediator ϕ_e : **new and crucial**

Admits a large HLFV without giving a too large CLFV, without significant tuning

HLFV is enhanced by μ and survives in the decoupling limit $\theta \rightarrow 0$





CLFV is not proportional to μ and goes to zero if the singlet-doublet mediator mixing angle vanishes

$$\Gamma(\tau \to \mu \gamma) = \frac{(m_{\tau}^2 - m_{\mu}^2)^3}{4\pi m_{\tau}^3} \left[|C_L|^2 + |C_R|^2 \right]$$

in the decoupling limit: F_L is not suppressed while $C_L \rightarrow 0^{-1}$

Seungwon Baek and ZF Kang, JHEP, 1603 (2016) 106

> The singlet mediator ϕ_e : **new and crucial**

Estimation of ratio $R\tau = Br(h \rightarrow \tau \mu)/Br(\tau \rightarrow \mu \gamma)$ in the decoupling limit



Seungwon Baek and ZF Kang, JHEP, 1603 (2016) 106

between DM annihilation & CLFV

The singlet mediator ϕ_e : **new and crucial**

> Admits a singlet fermionic & leptophlic dark matter candidate difficult in the Ma's model because of the strong tension

Annihilation without p-wave suppression: $\sigma v = a + bv^2$

Chiral flipping

2. Generically $a \neq 0$ (in particular for a large HLFV), leading to distinguishable signals

 $a = \frac{1}{16\pi M^2} \frac{1}{(1+x_i^{-1})^2} \left(|\lambda_{ia}^L \lambda_{ib}^R|^2 + |\lambda_{ia}^R \lambda_{ib}^L|^2 \right), \qquad \text{HLFV}, \text{ leading to dis}$ $b = \frac{1}{96\pi M^2} \frac{1}{(1+x_i^{-1})^4} \left[2|\lambda_{ia}^L|^2 |\lambda_{ib}^L|^2 (1+x_i^{-2}) - |\lambda_{ia}^L|^2 |\lambda_{ib}^R|^2 \left(1+4x_i^{-1}-3x_i^{-2} \right) + (L\leftrightarrow R) \right].$

Seungwon Baek and ZF Kang, JHEP, 1603 (2016) 106

DM direct detection at radiation

Higgs mediated & photon mediated: NOT completely undetectable



 $\sigma_{\rm SI}^p \approx 4.0 \times 10^{-8} \left(\lambda_{hN}(0) / 0.1 \right)^2 {\rm pb}$

$$\lambda_{hN}(0) \approx 0.01 \times \left(\frac{\sin\theta}{0.2}\right) \left(\frac{|y_{La}|^2 + |y_{Ra}|^2}{1}\right) \left(\frac{\mu}{5\text{TeV}}\right) \left(\frac{0.3\text{TeV}}{M}\right)$$



$$\mathcal{A} \approx -\frac{e\left(|\lambda_{ia}^L|^2 + |\lambda_{ia}^R|^2\right)}{192\pi^2 M^2} \left(-3\log(x_i\epsilon_a) - \frac{x_i + 3}{1 - x_i}\log\frac{x_i^{-1} - 1}{\sqrt{\epsilon_a}}\right)$$

 $\mathcal{A}/\left(|\lambda_{ia}^L|^2+|\lambda_{ia}^R|^2\right) ~\sim~ \mathcal{O}(10^{-7}) \mathrm{GeV^{-2}}.$

Conclusion & outlook

Higgs flavor violating decay is a good probe to new physics

HLVF is the left playground

We establish a connection between HLFV and lepton-flavored dark matter

The other playground: top flavor violating decay to Higgs & top flavored DM.....