model-independent determination of Higgs (self-)couplings @ e+e- colliders

Junping Tian (U' of Tokyo) CEPC Theory Group Workshop, July 11-14, 2017 @ IHEP model independence in kappa framework (elementary school)

- recoil mass technique —> inclusive σ_{Zh}
- $\sigma_{Zh} \longrightarrow \kappa_Z \longrightarrow \Gamma(h->ZZ^*)$
- WW-fusion $v_e v_e h \longrightarrow \kappa_W \longrightarrow \Gamma(h \rightarrow WW^*)$
- total width $\Gamma_h = \Gamma(h \rightarrow ZZ^*)/BR(h \rightarrow ZZ^*)$
- or $\Gamma_h = \Gamma(h \longrightarrow WW^*)/BR(h \rightarrow WW^*)$
- then all other couplings

PoS EPS-HEP2013 (2013) 316

Nucl.Part.Phys.Proc. 273-275 (2016) 826-833

the key: inclusive σ_{Zh} (independent of h decay modes)

Yan, et al, Phys.Rev. D94 (2016) 113002; Thomson, Eur.Phys.J. C76 (2016) 72

$H \rightarrow XX$	bb	cc	gg	au au	WW*	ZZ^*	$\gamma\gamma$	γZ
BR (SM)	57.8%	2.7%	8.6%	6.4%	21.6%	2.7%	0.23%	0.16%
Lepton Finder	93.70%	93.69%	93.40%	94.02%	94.04%	94.36%	93.75%	94.08%
Lepton ID+Precut	93.68%	93.66%	93.37%	93.93%	93.94%	93.71%	93.63%	93.22%
$M_{l^+l^-} \in [73, 120] \text{ GeV}$	89.94%	91.74%	91.40%	91.90%	91.82%	91.81%	91.73%	91.47%
$p_{\mathrm{T}}^{\mathrm{l^+l^-}} \in [10, 70] \; \mathrm{GeV}$	89.94%	90.08%	89.68%	90.18%	90.04%	90.16%	89.99%	89.71%
$ \cos heta_{ m miss} < 0.98$	89.94%	90.08%	89.68%	90.16%	90.04%	90.16%	89.91%	89.41%
$\mathrm{BDT}>$ - 0.25	88.90%	89.04%	88.63%	89.12%	88.96%	89.11%	88.91%	88.28%
$M_{\rm rec} \in [110, 155] \text{ GeV}$	88.25%	88.35%	87.98%	88.43%	88.33%	88.52%	88.21%	87.64%

bias < 0.1 in leptonic recoil mode

still need effort to achieve bias in hadronic recoil mode < 1%

question 1: how can we determine λ_{hhh} if there are anomalous hhVV, hVV, hhh couplings?



question 2: can we assume $\sigma(e+e-->Zh) \propto \Gamma(h->ZZ^*)$?







question 3: can we determine hWW precisely at $\sqrt{s} = 250$ GeV?



some quick answers

 measure directly hVV couplings (tensor structure) using σ, dσ/dX, in e+e- —> Zh process

Ogawa, Fujii, Tian, EPS-HEP 2017

• measure hhVV couplings and λ_{hhh} simultaneously using σ , d σ /dX, in e+e- —> Zhh process

determine tensor structure of hVV couplings

 $e^+ + e^- \to Zh \to f\bar{f}h$





determine tensor structure of hVV couplings (full simulation)

$$L_{hZZ} = M_Z^2 (\frac{1}{v} + \frac{a}{\Lambda}) hZ_{\mu} Z^{\mu} + \frac{b}{2\Lambda} hZ_{\mu\nu} Z^{\mu\nu} + \frac{b}{2\Lambda} hZ_{\mu\nu} \tilde{Z}_{\mu\nu}$$

$$\Lambda = 1 \text{ TeV}$$

$$\frac{\sqrt{s}=250 \text{GeV and } \int \text{Ldt}=250 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV}}$$

$$\frac{\sqrt{s}=500 \text{GeV and } \int \text{Ldt}=500 \text{fb}^{-1}}{\sqrt{s}=500 \text{GeV}}$$

hhVV, hVV and λ_{hhh} in e+e- —> Zhh



long answer: SM Effective Field Theory

Model-Independent Determination of the Triple Higgs Coupling at e^+e^- Colliders

TIM BARKLOW^{a1}, KEISUKE FUJII^b, SUNGHOON JUNG^{a1}, MICHAEL E. PESKIN^{a1}, AND JUNPING TIAN^c

Improved Formalism for Precision Higgs Coupling Fits

TIM BARKLOW^a, GAUTHIER DURIEUX^b, KEISUKE FUJII^c, CHRISTOPHE GROJEAN^{b,d}, JIAYIN GU^{b,e}, SUNGHOON JUNG^f, ROBERT KARL^b, JENNY LIST^b, TOMOHISA OGAWA^c, MICHAEL E. PESKIN^a, JUNPING TIAN^g, AND KECHEN WANG^{b,e}

"The leptonic future of the Higgs", Durieux, et al, 1704.02333

SM Effective Field Theory

("Warsaw" basis)

$$\begin{split} \Delta \mathcal{L} &= \frac{c_H}{2v^2} \partial^{\mu} (\Phi^{\dagger} \Phi) \partial_{\mu} (\Phi^{\dagger} \Phi) + \frac{c_T}{2v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\Phi^{\dagger} \overleftrightarrow{D}_{\mu} \Phi) - \frac{c_6 \lambda}{v^2} (\Phi^{\dagger} \Phi)^3 \\ &+ \frac{g^2 c_{WW}}{m_W^2} \Phi^{\dagger} \Phi W^a_{\mu\nu} W^{a\mu\nu} + \frac{4gg' c_{WB}}{m_W^2} \Phi^{\dagger} t^a \Phi W^a_{\mu\nu} B^{\mu\nu} \\ &+ \frac{g'^2 c_{BB}}{m_W^2} \Phi^{\dagger} \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g^3 c_{3W}}{m_W^2} \epsilon_{abc} W^a_{\mu\nu} W^{b\nu}{}_{\rho} W^{c\rho\mu} \\ &+ i \frac{c_{HL}}{v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\overline{L} \gamma_{\mu} L) + 4i \frac{c'_{HL}}{v^2} (\Phi^{\dagger} t^a \overleftrightarrow{D}^{\mu} \Phi) (\overline{L} \gamma_{\mu} t^a L) \\ &+ i \frac{c_{HE}}{v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\overline{e} \gamma_{\mu} e) \;. \end{split}$$

10 operators (h,W,Z, γ): CH, CT, C6, CWW, CWB, CBB, C3W, CHL, C'HL, CHE

- + 4 SM parameters: g, g', v, λ
- + 5 operators modifying h couplings to b, c, τ , μ , g
- + 2 parameters for h->invisible and exotic

EFT input: EWPOs (7)

$$\alpha(m_Z), G_F, m_W, m_Z, m_h, A_{LR}(\ell), \Gamma(Z \to \ell^+ \ell^-)$$

$$\begin{split} 4\pi\alpha(m_Z) &= g_0^2 s_0^2 \left(1 + 2s_0^2 \delta g + 2c_0^2 \delta g' \right. \\ &+ s_0^2 (8c_{WW}) - 2s_0^2 (8c_{WB}) + s_0^2 (8c_{BB}) \right) \\ \frac{G_F}{\sqrt{2}} &= \frac{1}{2v_0^2} \left(1 - 2\delta v + 2c'_{HL} \right) \\ m_W &= \frac{g_0 v_0}{2} \left(1 + \delta g + \delta v + \frac{1}{2} (8c_{WW}) \right) \\ m_Z &= \frac{(g_0^2 + g_0'^2)^{1/2} v_0}{2} \left(1 + c_0^2 \delta g + s_0^2 \delta g' + \delta v - \frac{1}{2} c_T \right. \\ &+ \frac{1}{2} c_0^2 (8c_{WW}) + s_0^2 (8c_{WB}) + \frac{1}{2} (s_0^4/c_0^2) (8c_{BB}) \\ m_h &= \sqrt{2\lambda_0} \ v_0 \left(1 + \delta v + \frac{1}{2} \delta \lambda - \frac{1}{2} c_H + \frac{3}{4} c_6 \right) \end{split}$$

$$\begin{split} A_{\ell} &= \frac{(1-4s_0^2)}{(1-4s_0^2+8s_0^4)} + \frac{32c_0^2s_0^4(1-2s_0^2)}{D^2}\delta g - \frac{32c_0^2s_0^4(1-2s_0^2)}{D^2}\delta g' \\ &+ \frac{16s_0^4(1-2s_0^2)}{D^2}(c_{HL}+c_{HL}') + \frac{8s_0^2(1-2s_0^2)^2}{D^2}c_{HE} \\ &+ \frac{16c_0^2s_0^4(1-2s_0^2)}{D^2}(8c_{WW}) - \frac{16s_0^4(1-2s_0^2)^2}{D^2}(8c_{WB}) - \frac{16s_0^6(1-2s_0^2)}{D^2}(8c_{BB}) \\ \Gamma_{\ell} &= \Gamma_{\ell 0} \Big(1 + \frac{2c_0^2(1-8s_0^2)}{D}\delta g - \frac{2s_0^2(3-16s_0^2+8s_0^4)}{D}\delta g' + \frac{2(1-2s_0^2)}{D}(c_{HL}+c_{HL}') - \frac{4s_0^2}{D}c_{HE} \\ &+ \frac{c_0^2(1-8s_0^2)}{D}(8c_{WW}) - \frac{2s_0^2(1-8s_0^2+8s_0^4)}{D}\delta g' + \frac{2(1-2s_0^2)}{D}(8c_{WB}) - \frac{s_0^4(3-16s_0^2+8s_0^4)}{c_0^2D}(8c_{BB}) \Big) \end{split}$$

EFT input: TGC (3)

$$egin{aligned} g_{1Z} &= 1 + (1 + s_0^2) \delta g - s_0^2 \delta g' + rac{1}{2} (1 + s_0^2) (8 c_{WW}) + rac{s_0^4}{c_0^2} (8 c_{WB}) - rac{1}{2} rac{s_0^4}{c_0^2} (8 c_{BB}) \ \kappa_A &= 1 + (8 c_{WB}) \ \lambda_A &= -6 g_0^2 c_{3W} \end{aligned}$$

2000 fb-1 @ 250 GeV, simultaneous fit

$$\Delta g_{1Z} = 3.8 \times 10^{-4} \qquad \rho(g_{1Z}, \kappa_{\gamma}) = 70.1\%$$

$$\Delta \kappa_{\gamma} = 4.5 \times 10^{-4} \qquad \rho(g_{1Z}, \lambda_{\gamma}) = 41.0\%$$

$$\Delta \lambda_{\gamma} = 3.8 \times 10^{-4} \qquad \rho(\kappa_{\gamma}, \lambda_{\gamma}) = 38.5\%$$

Barklow, Karl, List,

preliminary results, extrapolated from 500 GeV (1TeV) full simulation studies;

EFT input: BR(h-> $\gamma\gamma$)/BR(h->ZZ*), BR(h-> γ Z)/BR(h->ZZ*) (2: HL-LHC)

$$\begin{split} \Gamma(h \to \gamma \gamma) &= \Gamma(h \to \gamma \gamma)_0 \cdot \left(1 + (1 + 2s_w^2) \delta g + 2c_w^2 \delta g' - \delta v - c_H \\ &+ 526.1 \ s_w^2 ((8c_{WW}) - 2(8c_{WB}) + (8c_{BB})] \right) \\ \Gamma(h \to Z\gamma) &= \Gamma(h \to Z\gamma)_0 \cdot \left(1 + [0 \text{ for the moment}] - \delta v - c_H \\ &+ 289.7 \ s_w c_w ((8c_{WW}) - (1 - \frac{s_w^2}{c_w^2})(8c_{WB}) - \frac{s_w^2}{c_w^2}(8c_{BB})] \right) \end{split}$$

$$\Gamma(h \to ZZ^*) = \Gamma(h \to ZZ^*)_0 \cdot (1 - \delta v - c_H - (0.50)[c_w^2(8c_{WW}) + 2s_w^2(8c_{WB}) + \frac{s_w^4}{c_w^2}(8c_{BB})]\Big)$$

EFT coefficients

10: CH, CT, C6, CWW, CWB, CBB, C3W, CHL, C'HL, CHE + 4: g, g', ν, λ

can already be determined, except c₆, с_н

---> Higgs observables @ e+e-

Higgs couplings in EFT

$$\begin{split} \Delta \mathcal{L}_{Zhh} &= -\eta_h \lambda_0 v_0 h^3 + \eta_Z \frac{m_Z^2}{v_0} Z_\mu Z^\mu h + \frac{1}{2} \eta_{2Z} \frac{m_Z^2}{v_0^2} Z_\mu Z^\mu h^2 \\ &\quad + \frac{\theta_h}{v_0} h \partial_\mu h \partial^\mu h + \frac{\zeta_Z}{2v_0} Z_{\mu\nu} Z^{\mu\nu} h + \frac{\zeta_{2Z}}{4v_0^2} Z_{\mu\nu} Z^{\mu\nu} h^2 \\ &\quad + \frac{\zeta_{AZ}}{v_0} A_{\mu\nu} Z^{\mu\nu} h + \frac{\zeta_{2AZ}}{2v_0^2} A_{\mu\nu} Z^{\mu\nu} h^2 \\ &\quad + g_{LZh} (\overline{e}_L \gamma_\mu e_L) Z^\mu (\frac{h}{v_0} + \frac{1}{2} \frac{h^2}{v_0^2}) + g_{RZh} (\overline{e}_R \gamma_\mu e_R) Z^\mu (\frac{h}{v_0} + \frac{1}{2} \frac{h^2}{v_0^2}) \end{split}$$

$$\begin{split} \eta_{Z} &= (1 - c_{T} - \frac{1}{2}c_{H} - c_{HL}') \\ \eta_{2Z} &= (1 - 5c_{T} - c_{H} - 2c_{HL}') \\ \eta_{W} &= (1 - \frac{1}{2}c_{H} - c_{HL}') \\ \eta_{2W} &= (1 - c_{H} - c_{H}') \\ \eta_{2W} &= (1 - c_{H} - c_{H}') \\ \eta_{2W} &$$

EFT input: $\sigma(e+e-->Zh)$, $\sigma(e+e-->Zhh)$

• c_H has to be determined by inclusive σ_{Zh} measurement

• c₆ has to be determined by double Higgs measurement

question 1: how can we determine λ_{hhh} if there are anomalous hhVV, hVV, hhh couplings?



answer to Q1: determine λ_{hhh} in EFT



 $\frac{\sigma_{Zhh}}{\sigma_{SM}} - 1 = 0.565c_6 - 3.58c_H + 16.0(8c_{WW}) + 8.40(8c_{WB}) + 1.26(8c_{BB}) - 6.48c_T - 65.1c'_{HL} + 61.1c_{HL} + 52.6c_{HE},$

question 2: can we assume $\sigma(e+e-->Zh) \propto \Gamma(h->ZZ^*)$?







answer to Q2:

• $\sigma(e+e-->Zh) \propto \kappa^2(hZZ) \propto \Gamma(h->ZZ^*)$ not any more: EFT is more general than kappa-framework

$$\delta \mathcal{L} = (1+\eta_Z) \frac{2m_Z^2}{v} h Z_\mu Z^\mu + \zeta_Z \frac{h}{2v} Z_{\mu\nu} Z^{\mu\nu}$$





answer to Q3: hWW coupling can be as precise as hZZ @ $\sqrt{s} = 250$ GeV

 hWW/hZZ ratio can be determined to <0.1%: feature of a general SU(2) x U(1) gauge theory

$$\Gamma(h \to ZZ^*) = (SM) \cdot (1 + 2\eta_Z - (0.50)\zeta_Z) ,$$

$$\Gamma(h \to WW^*) = (SM) \cdot (1 + 2\eta_W - (0.78)\zeta_W)$$

$$\eta_W = -\frac{1}{2}c_H$$

$$\kappa_W = -\frac{1}{2}c_H - c_T$$

custodial symmetry

$$\eta_Z = -\frac{1}{2}c_H - c_T$$

SM-like hVV

 $C_i \sim O(10^{-4} - 10^{-3})$

$$\zeta_W = (8c_{WW})$$

$$\zeta_Z = c_w^2(8c_{WW}) + 2s_w^2(8c_{WB}) + (s_w^4/c_w^2)(8c_{BB})$$

anomalous hVV

typical precisions by EFT: combined EWPO+TGC+Higgs fit

ILC H20: ∫Ldt = 2 ab⁻¹ @ 250 GeV

coupling $\Delta g/g$	kappa-fit	EFT-fit
hZZ	0.38%	0.63%
hWW	1.9%	0.63%
hbb	2.0%	0.89%
$\Gamma_{\rm h}$	4.2%	2.1%

(for hZZ and hWW couplings: 1/2 of partial width precision)

homework from EFT (limiting factors other than usual Higgs observables)

- TGC: full simulation at 250 GeV
- improve $h\gamma Z$ couplings: using both $h \rightarrow \gamma Z$ and $e + e \rightarrow \gamma h$
- better constrain contact interactions:
 - improve ALR (AFB)
 - improve $\Gamma(Z \rightarrow ee)$
 - improve Γ(W->ev)



Higgs is a window to new physics

coupling deviation patter in BSM benchmark models

	Model	$b\overline{b}$	$c\overline{c}$	gg	WW	au au	ZZ	$\gamma\gamma$	$\mu\mu$
1	MSSM [34]	+4.8	-0.8	- 0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2	Type II 2HD [36]	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3	Type X 2HD [36]	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4	Type Y 2HD [36]	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5	Composite Higgs [38]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6	Little Higgs w. T-parity [39]	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
7	Little Higgs w. T-parity [40]	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8	Higgs-Radion [41]	-1.5	- 1.5	10.	-1.5	-1.5	-1.5	-1.0	-1.5
9	Higgs Singlet [42]	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

Table 4: Deviations from the Standard Model predictions for the Higgs boson couplings, in %, for the set of new physics models described in the text. As in Table 1, the effective couplings g(hWW) and g(hZZ) are defined as proportional to the square roots of the corresponding partial widths. typical parameters of benchmark models

- a PMSSM model with b squarks at 3.4 TeV, gluino at 4 TeV
- a Type II 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
- a Type X 2 Higgs doublet model with $m_A = 450 \text{ GeV}, \tan \beta = 6$
- a Type Y 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
- a composite Higgs model MCHM5 with $f = 1.2 \text{ TeV}, m_T = 1.7 \text{ TeV}$
- a Little Higgs model with T-parity with $f = 785 \text{ GeV}, m_T = 2 \text{ TeV}$
- A Little Higgs model with couplings to 1st and 2nd generation with $f=1.2 \text{ TeV}, m_T=1.7 \text{ TeV}$
- A Higgs-radion mixing model with $m_r = 500 \text{ GeV}$
- a model with a Higgs singlet at 2.8 TeVcreating a Higgs portal to dark matter and large λ for electroweak baryogenesis

model discrimination by EFT

$$(\chi^2)_{AB} = (g_A^T - g_B^T) [VCV^T]^{-1} (g_A - g_B)$$

g_A, g_B: vector of couplings in Model A, B

Vij: linear dependence of coupling gi on EFT coefficient cj

C: covariance matrix of EFT coeffs

 given the coupling deviations in two models, this χ2 gives the most appropriate separation power, taking into account all correlations

discrimination between BSM models (ILC250 stage)



once find deviation against SM —> can tell which BSM

backup

comments on beam polarizations

- not changed: important for systematics control, nature of new particle (once found), e.g. Higgsino, WIMPs
- new roles in EFT
 - -> separate hZZ and h γ Z couplings



important to constrain contact interaction

Ζ



evolution of discovery potential (against SM)



discrimination between BSM models



pin down the story after 250 + 500 full ILC

evolution of coupling precisions: H-20 (-CD/E/F)

