New results of R-parity-violating MSSM contributions to neutral mesons' mixing

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- Softly-broken supersymmetric (SUSY) extensions of the Standard Model (SM) have long been regarded as a leading class of candidates for the resolution of the Hierarchy Problem, as well as a possible framework in view of understanding the nature of dark matter or the unification of gauge-couplings.
- R-parity-violating Minimal Supersymmetric Standard Model (RPV-MSSM) leads to a distinctive phenomenology related to LHC searches.
- Low-energy flavor observables can place stringent bounds on parameters that arise in New Physics (NP) theories.

• Existing studies in the literature are far from being comprehensive.

MSSM and R-parity

- Despite the fact that no predicted superpartner has been found in the LHC, the MSSM remains one of the leading candidates of NP theories
- It introduces for each Standard-Model (SM) field a superpartner field, thus solving the Hierarchy Problem in the SM
- R-parity for particles is usually implied to ensure proton stability and obtain automatically a lightest supersymmetric particle (LSP) as dark matter (DM) candidate if it is neutral

R-parity and RPV-MSSM

For each field, a discrete symmetry called R-parity is defined as follows:

R-parity

 $R_P = (-1)^{3\mathbf{B}+L+2\mathbf{S}}$

- B: baryon number, L: lepton number, S: spin.
 - $R_P = +1$ for the SM fields and $R_P = -1$ for their superpartners.

RPV-MSSM superpotential

$$W_{\mathcal{R}_p} = \mu_i H_u L_i + \frac{1}{2} \lambda_{ijk} L_i L_j \bar{E}_k + \lambda'_{ijk} L_i Q_j \bar{D}_k + \frac{1}{2} \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k$$

Flavor violation! $\Rightarrow \Delta M_d, \Delta M_s, \Delta M_K$, experimentally well known.

Deficiencies in the literature studies

 $\Delta M_d, \Delta M_s, \Delta M_K$ have been studied under RPV-MSSM in the past. However, deficiencies exist.

- Diagrams beyond the tree-level and box contributions have been ignored
- Sfermion mixings have been routinely ignored.
- RPV-induced mixings have also been routinely ignored.

RPV-induced mixings $\mu_i H_u L_i$:

- neutral Higgs-sneutrino
- charged Higgs-slepton

- neutralino-neutrino
- chargino-charged lepton

Computation procedure

- Framework: Effective Field Theory (EFT) where short-distance effects intervene via the Wilson coefficients of dimension-6 flavor-changing ($\Delta F=2$) operators
- Amplitude: calculate Feynman diagrams using RPV-MSSM.
- Matching: match the full-theory amplitudes to the effective Lagrangian and obtain the corresponding Wilson coefficients.
- ΔM's: Use software to evaluate the effects of the Wilson coefficients on the ΔM's. Explicit formulas are also well known.

$$\Delta M = \frac{| < f | \mathcal{L}_{eff} | i > |}{M}$$

Low-energy EFT ΔF = 2 operators

Effective Lagrangian:

$$\mathcal{L}_{\text{eff}} = \sum_{i} C_i O_i + h.c.$$

 C_i : Wilson coefficients. O_i : dim-6 operators.

$$\begin{array}{ll} O_{1} = (\bar{d}_{j}\gamma^{\mu}P_{L}d_{i})(\bar{d}_{j}\gamma_{\mu}P_{L}d_{i}), & \tilde{O}_{1} = (\bar{d}_{j}\gamma^{\mu}P_{R}d_{i})(\bar{d}_{j}\gamma_{\mu}P_{R}d_{i}), \\ O_{2} = (\bar{d}_{j}P_{L}d_{i})(\bar{d}_{j}P_{L}d_{i}), & \tilde{O}_{2} = (\bar{d}_{j}P_{R}d_{i})(\bar{d}_{j}P_{R}d_{i}), \\ O_{3} = (\bar{d}_{j}^{a}P_{L}d_{i}^{b})(\bar{d}_{j}^{b}P_{L}d_{i}^{a}), & \tilde{O}_{3} = (\bar{d}_{j}^{a}P_{R}d_{i}^{b})(\bar{d}_{j}^{b}P_{R}d_{i}^{a}), \\ O_{4} = (\bar{d}_{j}P_{L}d_{i})(\bar{d}_{j}P_{R}d_{i}), & O_{5} = (\bar{d}_{j}^{a}P_{L}d_{i}^{b})(\bar{d}_{j}^{b}P_{R}d_{i}^{a}). \end{array}$$

i, j = d, s, b down-type quarks, a, b = 1, 2, 3 three colors.

Compute the full-theory amplitudes

- Feynman t'Hooft gauge and dimensional regularization
- \overline{DR} -renormalization consistent with numerical tools
- Consider only short-distance effects: discarding QED and QCD loops. Photons and gluons are active fields in the EFT.
- Different topologies of Feynman diagrams: tree-level and its one-loop corrections, and one-loop box diagrams
- Crosscheck between four-component and two-component spinor formalisms

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Topologies of Feynman diagrams I



The tree-level contribution is purely due to the λ' couplings of LQD operator.

Topologies of Feynman diagrams II



(a) S/F/S/F "straight" box



(b) S/F/S/F "scalar-cross" box



(c) S/F/S/F × "fermion-cross" box



(d) V/F/S/F "straight" box



(e) V/F/S/F "cross" boxes

(f) V/F/S/F "fermion-cross" box

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(g) V/F/V/F "straight" box

S: Scalar, F: Fermion, V: Vector

Different one-loop contributions

- SM-like: box diagrams with internal u, c, t quarks, W and Goldstone bosons
- 2-Higgs-doublet-model-like: box diagrams with internal u, c, t quarks, charged-Higgs bosons and possibly W or Goldstone bosons
- R-parity conserving: box diagrams with chargino/sup, neutralino/sdown or gluino/sdown particles in the loop
- RPV: self-energies and vertex corrections, box diagrams with sneutrino/quark, slepton/quark, lepton/squarks, neutrino/squark or quark/squark internal lines)
- RPV-driven mixing further mixes these contributions

Comparing analytic results with the literature

- Self-energy and vertex corrections were not considered before, but at least the scalar self-energy is consistent with Higgs self-energy calculation in the literature.
- R-parity conserving MSSM result is recovered.
- In no-mixing limit, comparing with the literature shows some difference:

 C_5 from $[\nu/\tilde{D}/\nu/\tilde{D}]$: difference in prefactor and in sfermion chiralities.

• Cross-check with amplitudes generated from public code FlavorKit

Software tools introduction

Tools: Mathematica packages PreSARAH, SARAH, FlavorKit, SPheno and Python3 software Flavio

- PreSARAH: incorporates quark self-energy and vertex corrections; scalar-self energy included by hand
- SARAH: uses *MSSM_TriRpV* model file modified by PreSARAH
- FlavorKit: generates all the amplitudes for general meson mixing $\bar{d}_j d_i \Leftrightarrow \bar{d}_i d_j$
- SPheno: spectrum generator, giving in particular Wilson coefficients and predictions for ΔM 's.
- Flavio: Better handle with hadronic contributions for ΔM_d and ΔM_s .

Reproduction of plots in the literature I

Figure 7 of Altmannshofer, Buras, Guadagnoli, 2007. Using *MSSM* model file in SARAH.



Reproduction of plots in the literature II

Figure 4 of de Carlos, White, 1997. Using *MSSMTriBpV* model file in SARAH.



 $A_0 = 0, tan\beta = 10, \mu < 0.$

Different boxes contributions to ΔM_{κ} . black dashed: χ^0 red dot-dashed: $\chi^0 - \tilde{g}$ magenta dashed: \tilde{g} red solid: experimental upper limit blue solid: full cvan dot: direct RPV boxes without W green solid: direct RPV boxes with W

Numerics To-Do: Quark Flavor Violation (QFV)

Explore the size of and the interplay between different sources of $\mathsf{QFV}.$

- λ' tree-level: $\lambda'_{i12} \cdot \lambda'_{i21}$, $\lambda'_{i13} \cdot \lambda'_{i31}$, $\lambda'_{i23} \cdot \lambda'_{i32}$, (i = 1, 2, 3)
- λ' boxes, unaligned with tree-level λ' 's..
- λ'' boxes
- CKM (charged Higgs) v.s RPV
- RPC-MSSM v.s RPV

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Plots of tree-level LQD couplings I



Plots of tree-level LQD couplings II



$$\Delta M_K$$
 w.r.t. $\lambda'_{112}\lambda'_{121}$ at
SUSY scale
 $M_1 = M_2 = M_3 = 500$
GeV
 $\mu = 1000$ GeV,
 $tan\beta = 10$.

Errorbar: 15% SM + 30% NP

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- Studies on neutral mesons mixing from RPV-MSSM have huge room for improvement.
- Analytically, we obtained a comprehensive list of tree-level and one-loop contributions to these processes excluding QED and QCD loops, and there is some difference w.r.t. some results in the literature
- Numerically, we are trying to compare the effects on Quark Flavor Violation from different sources: RPV-MSSM, RPC-MSSM, CKM, etc.

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Thank You! 謝謝!

Matching amplitudes with Wilson coefficients

Procedure to determine the Wilson coefficients:

- calculate the scattering amplitude in the full theory;
- find the appropriate corresponding dim-6 operators;
- determine the Wilson coefficients of the operators by equating the amplitude calculated from the effective operators with that from the full theory.

Illustrate this by C_4^{tree} : $A_4^{\text{tree}} = \frac{i}{m_5^2} \left(\left[g_L^{Sd_jd_i} P_L \right] \otimes \left[g_R^{Sd_jd_i} P_R \right] + \left[g_R^{Sd_jd_i} P_R \right] \otimes \left[g_L^{Sd_jd_i} P_L \right] \right)$ corresponding to O_4 : $C_4^{\text{tree}} = \frac{i}{2m_5^2} (2g_L^{Sd_jd_i} g_R^{Sd_jd_i})$. Factor 2 in the denominator arises because O_4 is symmetrical: $O_4 = (\bar{d}_j P_L d_i)(\bar{d}_j P_R d_i) = (\bar{d}_j P_R d_i)(\bar{d}_j P_L d_i)$