Examining the model dependence of extracting the kinetic freeze-out temperature and transverse flow velocity in small collision system

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Content







The interacting system at the kinetic freeze-out (the last stage of collisions) stays at a thermodynamic equilibrium state or local equilibrium state, when the particle emission process is influenced not only by the thermal motion but also the flow effect.
By analyzing the transverse momentum spectra of final-state particles, one can obtain the kinetic freeze-out temperature of interacting systems (emission sources) and the transverse flow velocity of produced particles.



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The chemical freeze-out temperature describes the excitation degree of the interacting system

The kinetic freeze-out temperature describes the excitation degree of the interacting system momentum spectra

The effective temperature can be extracted from the transverse

3

flow effect

4

describes the kinetic expansion characteristics

at the stage of chemical equilibrium

at the stage of kinetic and thermal equilibrium

by using some distribution laws

of the interacting system



(i) The BGBW model results in the *p*_T distribution to be

$$f_1(p_T) = C_1 p_T m_T \int_0^R r dr \times I_0 \left[\frac{p_T \sinh(\rho)}{T_0} \right] K_1 \left[\frac{m_T \cosh(\rho)}{T_0} \right]$$

E. Schnedermann, J. Sollfrank, and U. Heinz, Phys. Rev. C 48, 2462 (1993)

E. Schnedermann and U. Heinz, Phys. Rev. C 47, 1738 (1993).

transverse mass $\longrightarrow m_T = \sqrt{p_T^2 + m_0^2}$

boost angle $\longrightarrow \rho = \tanh^{-1}[\beta(r)]$

self-similar flow profile

 $\beta(r) = \beta_S (r/R)^{n_0}$

$$\beta_T = (2/R^2) \int_0^R r\beta(r) dr = 2\beta_S/(n_0 + 2)$$

the flow velocity on the surface of the thermal source

$$n_0 = 2$$

(ii) The TBW model results in the p_T distribution to be

 $f_2(p_T) = C_2 p_T m_T \int_{-\pi}^{\pi} d\phi \int_{-\pi}^{R} r dr \{1+$

Z. B. Tang, Y. C. Xu, L. J. Ruan, G. van Buren, F. Q. Wang, and Z. B. Xu, Phys. Rev. C 79, 051901(R) (2009).

$$\frac{q-1}{T_0} \left[m_T \cosh(\rho) - p_T \sinh(\rho) \cos(\phi) \right] \Big\}^{-q/(q-1)}$$

 $n_0 = 1 \longrightarrow \beta_T = 2\beta_S / (n_0 + 2) = (2/3)\beta_S$

 $-1/(q-1) \longrightarrow -q/(q-1)$

H. Zheng and L. L. Zhu, Adv. High Energy Phys.2016, 9632126 (2016).

(iii) The Alternative method based on Boltzmann statistics

 $T = T_0 + am_0$ $\langle p_T \rangle = b_1 + \beta_T \overline{m}$ $\langle p \rangle = b_2 + \beta \overline{m}$

H.-L. Lao, H.-R. Wei, F.-H. Liu, and R. A. Lacey, Eur. Phys. J. A 52, 203 (2016).

the form of Boltzmann distribution

J. Cans and D. Worku, Eur. Phys. J. A 48, 160 (2012).

$$f_3(p_T) = \frac{1}{N} \frac{dN}{dp_T} = C_3 p_T m_T \exp\left(-\frac{m_T}{T}\right)$$

(iv) The Alternative method based on Tsallis statistics

 $T = T_0 + am_0$ $\langle p_T \rangle = b_1 + \beta_T \overline{m}$ $\langle p \rangle = b_2 + \beta \overline{m}$

H.-L. Lao, H.-R. Wei, F.-H. Liu, and R. A. Lacey, Eur. Phys. J. A 52, 203 (2016).

> H. Zheng and L. L. Zhu, Adv. High Energy Phys. 2016, 9632126 (2016).

the form of Tsallis distribution

$$f_4(p_T) = \frac{1}{N} \frac{dN}{dp_T} = C_4 p_T m_T \left(1 + \frac{q-1}{T} m_T \right)^{-q/(q-1)}$$

For the spectra in a wide PT range, we have to consider the contribution of hard scattering process

The contribution of hard process is parameterized to an inverse power-law

$$f_H(p_T) = Ap_T \left(1 + \frac{p_T}{p_0}\right)^{-n}$$

G. Arnison et al. (UA1 Collaboration), Phys. Lett. B 118, 167 (1982).

We can use a superposition of both contributions of soft and hard processes

$$f_0(p_T) = k f_S(p_T) + (1 - k) f_H(p_T)$$





The symbols are taken from literature, the curves are our fitting results



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 $T = T_0 + am_0$

H.-L. Lao, F.-H. Liu, B.-C. Li, M.-Y. Duan, arXiv:1708.07749 [nucl-th] (2017).





 $\langle p_T \rangle = b_1 + \beta_T \overline{m}$

H.-L. Lao, F.-H. Liu, B.-C. Li, M.-Y. Duan, arXiv:1708.07749 [nucl-th] (2017).

 $p\rangle = b_2 + \beta \overline{m}$





pp is similar to **Peripheral collisions**

4 Conclusions

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The p_T spectra of π^{\pm} , K^{\pm} , p and p produced in pp and d-Au collisions at the RHIC, as well as pp and p-Pb collisions at the LHC, have been analyzed by four methods.

The four methods present similar results, and in some cases these results are in agreement with each other within errors.

 T_0 in central d-Au and p-Pb collisions is relatively larger than that in peripheral d-Au and p-Pb collisions, and β_T in central d-Au and p-Pb collisions is slightly larger than or equal to that in peripheral d-Au and p-Pb collisions.



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Central d-Au and p-Pb collisions are similar to central Au-Au and Pb-Pb collisions, and peripheral d-Au and p-Pb collisions are similar to peripheral Au-Au and Pb-Pb collisions.

Comparing with central nucleus-nucleus collisions, pp collisions are closer to peripheral nucleus-nucleus collisions due to similar numbers of participant nucleons.

In central collisions, the excitation degree at the kinetic freezeout is mainly determined by the maximum nucleus and collision energy, but not the numbers of participant nucleons and binary Collisions.

Thank you for your attention!

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