



"Decay constants" of doubly heavy baryons in QCD sum rules

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Outline



- Introduction
- QCD sum rules study
- Numerical results
- Summary and outlook



Introduction



Observation of the Doubly Charmed Baryon Ξ_{cc}^{++}

R. Aaij *et al.**

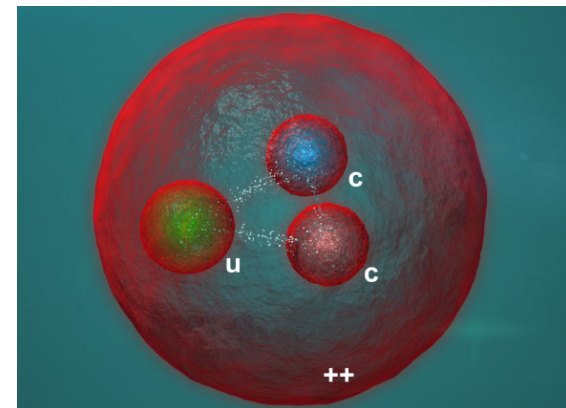
(LHCb Collaboration)

(Received 6 July 2017; revised manuscript received 2 August 2017; published 11 September 2017)

A highly significant structure is observed in the $\Lambda_c^+ K^- \pi^+ \pi^+$ mass spectrum, where the Λ_c^+ baryon is reconstructed in the decay mode $p K^- \pi^+$. The structure is consistent with originating from a weakly decaying particle, identified as the doubly charmed baryon Ξ_{cc}^{++} . The difference between the masses of the Ξ_{cc}^{++} and Λ_c^+ states is measured to be $1334.94 \pm 0.72(\text{stat.}) \pm 0.27(\text{syst.}) \text{ MeV}/c^2$, and the Ξ_{cc}^{++} mass is then determined to be $3621.40 \pm 0.72(\text{stat.}) \pm 0.27(\text{syst.}) \pm 0.14(\Lambda_c^+) \text{ MeV}/c^2$, where the last uncertainty is due to the limited knowledge of the Λ_c^+ mass. The state is observed in a sample of proton-proton collision data collected by the LHCb experiment at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 1.7 fb^{-1} , and confirmed in an additional sample of data collected at 8 TeV.

DOI: [10.1103/PhysRevLett.119.112001](https://doi.org/10.1103/PhysRevLett.119.112001)

- Properties**
- Mass
 - Lifetime
 - Transition form factors
 - Decay modes
 - ...



Quantum numbers and quark content

Baryon	Quark Content	s_h^π	J^P	Baryon	Quark Content	s_h^π	J^P
Ξ_{cc}	$\{cc\}q$	1^+	$1/2^+$	Ξ_{bb}	$\{bb\}q$	1^+	$1/2^+$
Ξ_{cc}^*	$\{cc\}q$	1^+	$3/2^+$	Ξ_{bb}^*	$\{bb\}q$	1^+	$3/2^+$
Ω_{cc}	$\{cc\}s$	1^+	$1/2^+$	Ω_{bb}	$\{bb\}s$	1^+	$1/2^+$
Ω_{cc}^*	$\{cc\}s$	1^+	$3/2^+$	Ω_{bb}^*	$\{bb\}s$	1^+	$3/2^+$
Ξ'_{bc}	$\{bc\}q$	0^+	$1/2^+$	Ω'_{bc}	$\{bc\}s$	0^+	$1/2^+$
Ξ_{bc}	$\{bc\}q$	1^+	$1/2^+$	Ω_{bc}	$\{bc\}s$	1^+	$1/2^+$
Ξ_{bc}^*	$\{bc\}q$	1^+	$3/2^+$	Ω_{bc}^*	$\{bc\}s$	1^+	$3/2^+$

Doubly charmed baryon spectrum ☺

- Many models have been applied to calculate masses: quark models or QCD sum rules

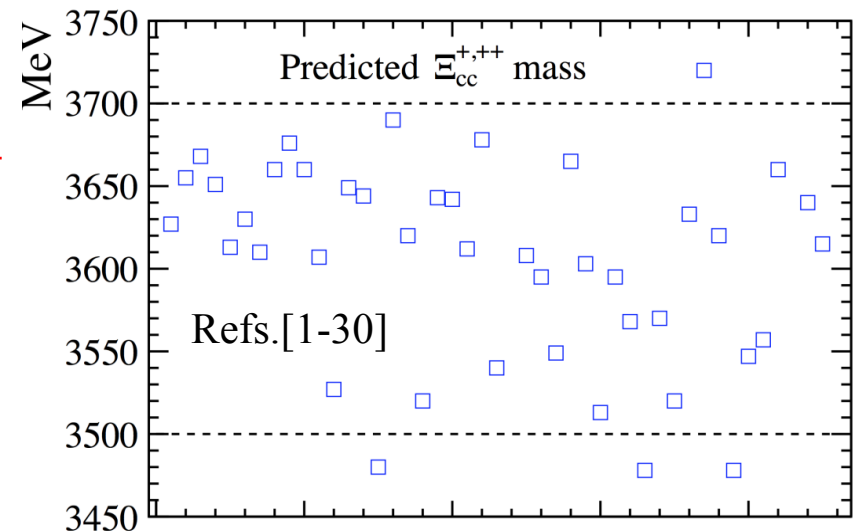
✓ Predicted Mass:

$$m(\Xi_{cc}^{++}) \sim m(\Xi_{cc}^+) \sim (3.5 - 3.7) \text{ GeV}$$

$$m(\Omega_{cc}) \sim m(\Xi_{cc}^+) + 0.1 \text{ GeV}$$

✓ Mass splitting between

Ξ_{cc}^{++} and Ξ_{cc}^+ is a few MeV



- Lattice QCD Calculation:

$$m(\Xi_{cc}) \sim 3.6 \text{ GeV}, m(\Omega_{cc}) \sim 3.7 \text{ GeV}$$

References can be found in 1703.09086

Lifetime ☹️

- Large uncertainties: differ by a factor of 4~8
- $\tau(\Xi_{cc}^{++}) \gg \tau(\Xi_{cc}^+) \sim \tau(\Omega_{cc}^+)$

literature	Ξ_{cc}^{++}	Ξ_{cc}^+	Ω_{cc}^+
Karliner, Rosner, 2014	185	53	
Chang, Li, Li, Wang, 2007	670	250	210
Kiselev, Likhoded, 2002	460 ± 50	160 ± 50	270 ± 60
Kiselev, Likhoded, 1998	430 ± 100	110 ± 10	
Guberina, Melic, Stefancic, 1998	1550	220	250

$$\langle 0 | J_H | H(q, s) \rangle = \lambda_H u(q, s)$$



dimension-3

$$\langle 0 | J_H | H(q, s) \rangle = f_H m_H^2 u(q, s)$$



dimension-1

- Lifetime
- Transition form factors
- Ξ_{cc}, Ω_{cc}
- Ξ_{bb}, Ω_{bb}
- Ξ_{bc}, Ω_{bc}

Applications of QCDSR

- Mass spectrum
- Decay constants
- Transition form factors
- Mixing matrix elements of K-meson and B-meson systems
- ...



QCD sum rules study

The interpolating current

$$J_{\Xi_{QQ}} = \epsilon_{abc} (Q_a^T C \gamma^\mu Q_b) \gamma_\mu \gamma_5 q_c,$$

$$J_{\Omega_{QQ}} = \epsilon_{abc} (Q_a^T C \gamma^\mu Q_b) \gamma_\mu \gamma_5 s_c,$$

$$J_{\Xi_{bc}} = \frac{1}{\sqrt{2}} \epsilon_{abc} (b_a^T C \gamma^\mu c_b + c_a^T C \gamma^\mu b_b) \gamma_\mu \gamma_5 q_c,$$

$$J_{\Omega_{bc}} = \frac{1}{\sqrt{2}} \epsilon_{abc} (b_a^T C \gamma^\mu c_b + c_a^T C \gamma^\mu b_b) \gamma_\mu \gamma_5 s_c.$$

The correlation function

$$\Pi(q) = i \int d^4x e^{iq \cdot x} \langle 0 | T[J(x), \bar{J}(0)] | 0 \rangle,$$

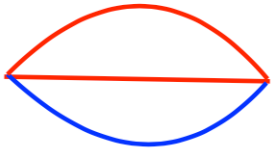
Hadronic level

$$\Pi(q) = \lambda_H^2 \frac{\not{q} + m_H}{m_H^2 - q^2} + \frac{1}{\pi} \int_{s_0}^{\infty} ds \frac{\text{Im}\Pi(s)}{s - q^2},$$

only $1/2^+$ are considered

OPE level

$$\Pi(q) = \not{q}\Pi_1(q^2) + \Pi_2(q^2)$$



$$\Pi_i(q^2) = \int_{(m_Q+m_{Q'})^2}^{\infty} ds \frac{\rho_i(s)}{s - q^2}, \quad i = 1, 2,$$

$$\rho_i(s) = \frac{1}{\pi} \text{Im}\Pi_i^{\text{OPE}}(s)$$

Sum rules

$$\lambda_H^2 e^{-m_H^2/M^2} = \int_{(m_Q+m_{Q'})^2}^{s_0} ds \rho_1(s) e^{-s/M^2},$$

$$\lambda_H^2 m_H e^{-m_H^2/M^2} = \int_{(m_Q+m_{Q'})^2}^{s_0} ds \rho_2(s) e^{-s/M^2}.$$

QCD sum rules study



$$\rho_1^{\text{pert}}(s) = \frac{6}{(2\pi)^4} \int_{\alpha_{\min}}^{\alpha_{\max}} \frac{d\alpha}{\alpha} \int_{\beta_{\min}}^{1-\alpha} \frac{d\beta}{\beta} \left([\alpha\beta s - \alpha m_Q^2 - \beta m_{Q'}^2]^2 + (1 - \alpha - \beta) m_Q m_{Q'} [\alpha\beta s - \alpha m_Q^2 - \beta m_{Q'}^2] \right), \quad (21)$$

$$\rho_1(s) = \rho_1^{\text{pert}}(s) + \frac{\langle g_s^2 G^2 \rangle}{72} \left(m_Q^2 \frac{\partial^3}{(\partial m_Q^2)^3} + m_{Q'}^2 \frac{\partial^3}{(\partial m_{Q'}^2)^3} \right) \rho_1^{\text{pert}}(s) \quad (22)$$

$$+ \frac{4m_Q m_{Q'} \langle g_s^2 G^2 \rangle}{(4\pi)^4} \left(\frac{\partial^2}{(\partial m_Q^2)^2} + \frac{\partial^2}{(\partial m_{Q'}^2)^2} \right) \int_{\alpha_{\min}}^{\alpha_{\max}} \frac{d\alpha}{\alpha} \int_{\beta_{\min}}^{1-\alpha} \frac{d\beta}{\beta} (1 - \alpha - \beta) (\alpha\beta s - \alpha m_Q^2 - \beta m_{Q'}^2) \\ + \frac{2\langle g_s^2 G^2 \rangle}{(4\pi)^4} \left(\frac{\partial}{\partial m_Q^2} + \frac{\partial}{\partial m_{Q'}^2} \right) \int_{\alpha_{\min}}^{\alpha_{\max}} d\alpha \int_{\beta_{\min}}^{1-\alpha} d\beta (3\alpha m_Q^2 + 3\beta m_{Q'}^2 - m_Q m_{Q'} - 4\alpha\beta s) \quad (23)$$

$$\rho_2(s) = -\frac{\langle \bar{q}q \rangle}{2\pi^2} \int_{\alpha_{\min}}^{\alpha_{\max}} d\alpha (3\alpha(1 - \alpha)s - 2\alpha m_Q^2 - 2(1 - \alpha)m_{Q'}^2 + 2m_Q m_{Q'}) \\ - \frac{\langle \bar{q}q g_s \sigma G q \rangle}{8\pi^2} \left(1 + \frac{s}{M^2} \right) A(s) - \frac{2\langle \bar{q}q g_s \sigma G q \rangle}{(4\pi)^2} \left((\alpha_{\max} - \alpha_{\min}) \right. \\ \left. + \frac{1}{2s(\alpha_{\max} - \alpha_{\min})} [\alpha_{\max}(1 - \alpha_{\max})s + \alpha_{\min}(1 - \alpha_{\min})s + 4m_Q m_{Q'}] \right), \quad (24)$$

QCD Sum rules with the negative parity baryon

D. Jido, N. Kodama and M. Oka, Phys. Rev. D 54, 4532 (1996)

Hadronic level

$$\Pi(q) = \lambda_+^2 \frac{\not{q} + m_+}{m_+^2 - q^2} + \lambda_-^2 \frac{\not{q} - m_-}{m_-^2 - q^2} + \frac{1}{\pi} \int_{s_0}^{\infty} ds \frac{\text{Im}\Pi(s)}{s - q^2},$$

\downarrow
 $1/2^+$

\downarrow
 $1/2^-$

$$\frac{1}{\pi} \text{Im}\Pi(q) = \lambda_+^2 (\not{q} + m_+) \delta(q^2 - m_+^2) + \lambda_-^2 (\not{q} - m_-) \delta(q^2 - m_-^2) + \dots$$

Taking $\vec{q} = 0$

$$\frac{1}{\pi} \text{Im}\Pi(q_0) = \gamma_0 A(q_0) + B(q_0) + \dots,$$

$$A(q_0) = \frac{1}{2} [\lambda_+^2 \delta(q_0 - m_+) + \lambda_-^2 \delta(q_0 - m_-)],$$

$$B(q_0) = \frac{1}{2} [\lambda_+^2 \delta(q_0 - m_+) - \lambda_-^2 \delta(q_0 - m_-)].$$

$$A + B \Rightarrow \lambda_+$$

OPE level

Taking $\vec{q} = 0$

$$\frac{1}{\pi} \text{Im}\Pi(q_0) = \gamma_0 \rho^A(q_0) + \rho^B(q_0) + \dots$$

$$\int_{\Delta}^{\sqrt{s_0}} dq_0 \times \exp \left[-q_0^2/M^2 \right] \longrightarrow \text{weight function}$$

$A + B$

$\rho^A + \rho^B$

$$\lambda_+^2 \exp \left[-\frac{m_+^2}{M^2} \right] = \int_{\Delta}^{\sqrt{s_0}} dq_0 (\rho^A(q_0) + \rho^B(q_0)) \exp \left[-\frac{q_0^2}{M^2} \right]$$

New sum rules



Numerical results

Numerical results

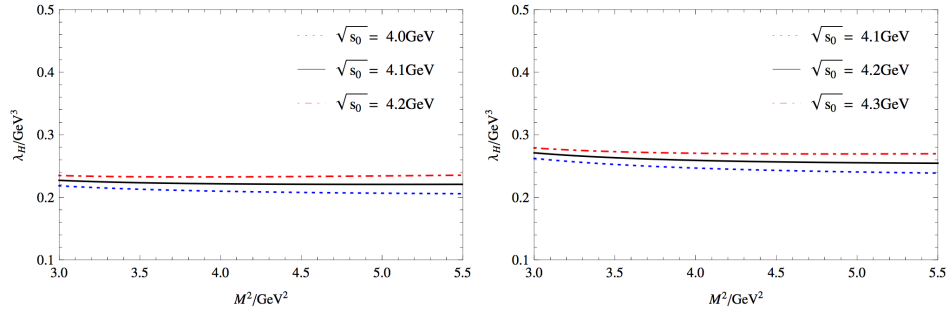


FIG. 1: The dependence on M^2 of the decay constants of Ξ_{cc} and Ω_{cc} at the scale $\mu = 2.1$ GeV. The continuum threshold are taken as $\sqrt{s_0} = 4.0 \sim 4.2$ GeV and $\sqrt{s_0} = 4.1 \sim 4.3$ GeV, respectively.

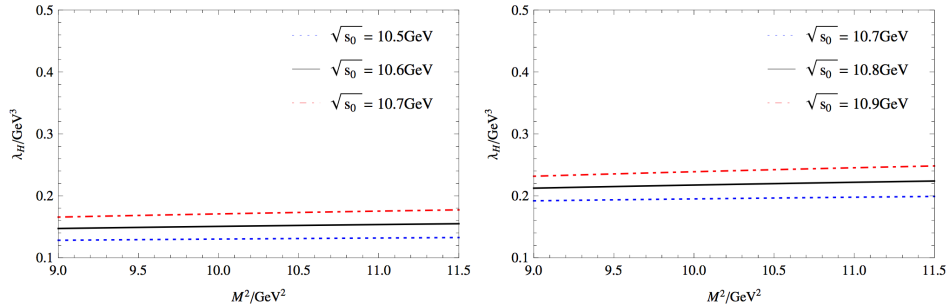


FIG. 2: The dependence on M^2 of the decay constant of Ξ_{bb} and Ω_{bb} at the scale $\mu = 2.1$ GeV. The continuum threshold is taken as $\sqrt{s_0} = 10.5 \sim 10.7$ GeV and $\sqrt{s_0} = 10.7 \sim 10.9$ GeV, respectively.

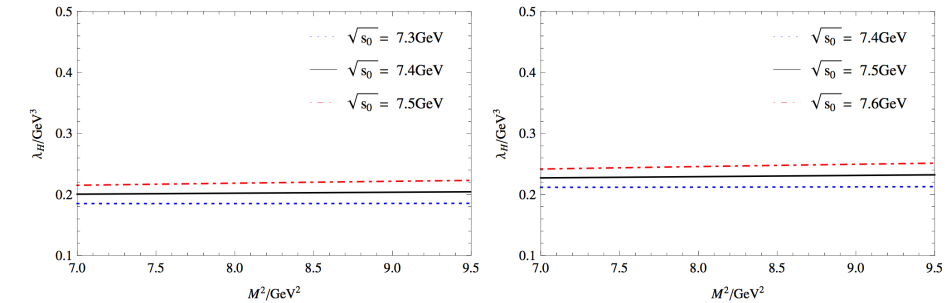


FIG. 3: The dependence on M^2 of the $\lambda_{\Xi_{bc}}$ and $\lambda_{\Omega_{bc}}$ at the scale $\mu = 2.1$ GeV. In the left and right panel, the continuum threshold are taken as $\sqrt{s_0} = 7.3 \sim 7.5$ GeV and $\sqrt{s_0} = 7.4 \sim 7.6$ GeV, respectively.

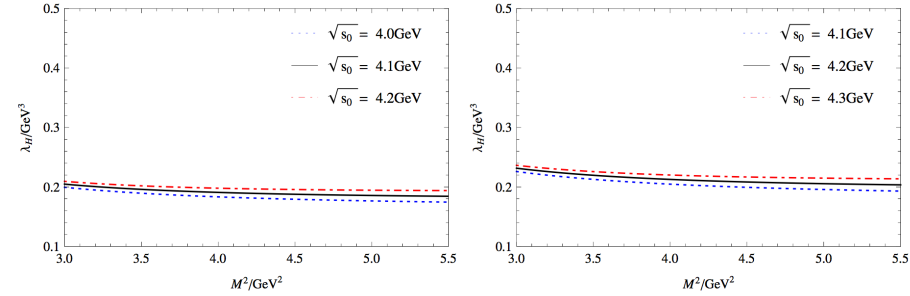


FIG. 5: The dependence on M^2 of the decay constant of Ξ_{cc} and Ω_{cc} at the scale $\mu = 2.1$ GeV. In the left and right panel, the continuum thresholds are taken as $\sqrt{s_0} = 4.0 \sim 4.2$ GeV and $\sqrt{s_0} = 4.1 \sim 4.3$ GeV, respectively.

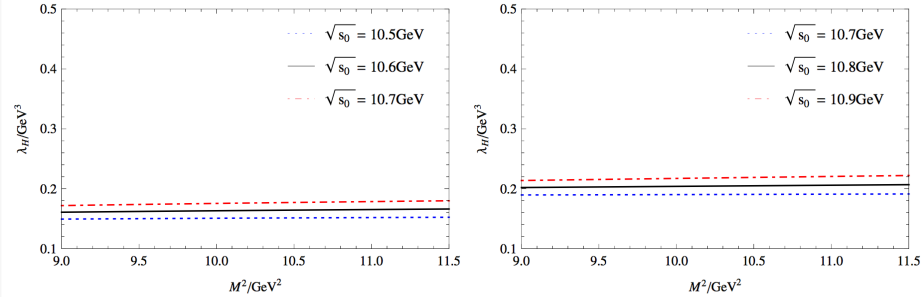


FIG. 6: The dependence on M^2 of the decay constant of Ξ_{bb} and Ω_{bb} at the scale $\mu = 2.1$ GeV. The continuum thresholds are taken as $\sqrt{s_0} = 10.5 \sim 10.7$ GeV and $\sqrt{s_0} = 10.7 \sim 10.9$ GeV for the left and right panel, respectively.

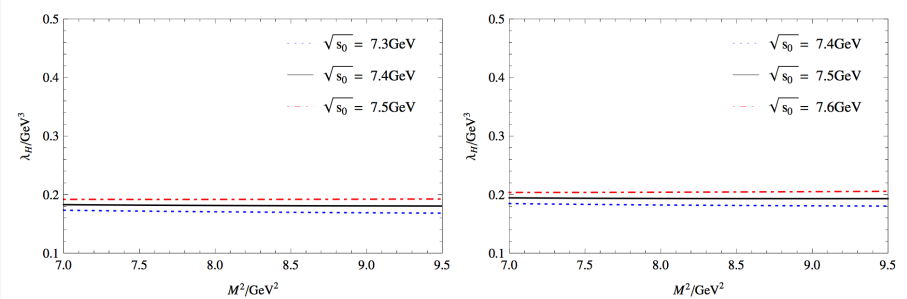


FIG. 7: The dependence on M^2 of the decay constant of Ξ_{bc} and Ω_{bc} at the scale $\mu = 2.1$ GeV. The continuum thresholds are taken as $\sqrt{s_0} = 7.3 \sim 7.5$ GeV and $\sqrt{s_0} = 7.4 \sim 7.6$ GeV for the left and right panel, respectively.



Numerical results

TABLE III: Decay constants λ_H (in units of GeV^3) and f_H (in units of 10^{-3} GeV) for the doubly heavy baryons at the scale $\mu = 2.1 \text{ GeV}$ and $\mu = 1 \text{ GeV}$: $\lambda_H = f_H m_H^2$. The first and second errors come from the uncertainties of $\sqrt{s_0}$ and M^2 respectively. The sum rule in Eq. (19) has been used.

Baryon	$\lambda_H(\text{GeV}^3)[\mu = 2.1 \text{ GeV}]$	$f_H(10^{-3} \text{ GeV})[\mu = 2.1 \text{ GeV}]$	$\lambda_H(\text{GeV}^3)[\mu = 1 \text{ GeV}]$	$f_H(10^{-3} \text{ GeV})[\mu = 1 \text{ GeV}]$
Ξ_{cc}	$0.221 \pm 0.012 \pm 0.002$	$16.9 \pm 1.0 \pm 0.2$	$0.196 \pm 0.011 \pm 0.002$	$15.0 \pm 0.8 \pm 0.1$
Ω_{cc}	$0.258 \pm 0.013 \pm 0.006$	$18.4 \pm 0.9 \pm 0.4$	$0.229 \pm 0.012 \pm 0.005$	$16.4 \pm 0.8 \pm 0.4$
Ξ_{bb}	$0.151 \pm 0.021 \pm 0.002$	$1.47 \pm 0.20 \pm 0.02$	$0.134 \pm 0.018 \pm 0.002$	$1.31 \pm 0.18 \pm 0.02$
Ω_{bb}	$0.219 \pm 0.023 \pm 0.004$	$2.07 \pm 0.22 \pm 0.03$	$0.194 \pm 0.020 \pm 0.003$	$1.84 \pm 0.19 \pm 0.03$
Ξ_{bc}	$0.202 \pm 0.017 \pm 0.001$	$4.20 \pm 0.36 \pm 0.02$	$0.180 \pm 0.015 \pm 0.001$	$3.73 \pm 0.32 \pm 0.02$
Ω_{bc}	$0.230 \pm 0.018 \pm 0.002$	$4.69 \pm 0.36 \pm 0.03$	$0.204 \pm 0.016 \pm 0.001$	$4.16 \pm 0.32 \pm 0.03$

TABLE IV: “Decay constants” λ_H (in units of GeV^3) and f_H (in units of 10^{-3} GeV) for the doubly heavy baryons at the scale $\mu = 2.1 \text{ GeV}$ and $\mu = 1 \text{ GeV}$: $\lambda_H = f_H m_H^2$. The first and second errors come from the uncertainties of $\sqrt{s_0}$ and M^2 respectively. The sum rule in Eq. (34) has been used.

Baryon	$\lambda_H(\text{GeV}^3)[\mu = 2.1 \text{ GeV}]$	$f_H(10^{-3} \text{ GeV})[\mu = 2.1 \text{ GeV}]$	$\lambda_H(\text{GeV}^3)[\mu = 1 \text{ GeV}]$	$f_H(10^{-3} \text{ GeV})[\mu = 1 \text{ GeV}]$
Ξ_{cc}	$0.189 \pm 0.008 \pm 0.007$	$14.4 \pm 0.6 \pm 0.5$	$0.168 \pm 0.007 \pm 0.006$	$12.8 \pm 0.5 \pm 0.5$
Ω_{cc}	$0.210 \pm 0.008 \pm 0.009$	$15.0 \pm 0.6 \pm 0.7$	$0.186 \pm 0.008 \pm 0.008$	$13.3 \pm 0.5 \pm 0.6$
Ξ_{bb}	$0.164 \pm 0.013 \pm 0.002$	$1.59 \pm 0.12 \pm 0.02$	$0.145 \pm 0.011 \pm 0.002$	$1.41 \pm 0.11 \pm 0.01$
Ω_{bb}	$0.204 \pm 0.014 \pm 0.001$	$1.94 \pm 0.13 \pm 0.01$	$0.181 \pm 0.012 \pm 0.001$	$1.72 \pm 0.12 \pm 0.01$
Ξ_{bc}	$0.181 \pm 0.011 \pm 0.001$	$3.76 \pm 0.23 \pm 0.02$	$0.161 \pm 0.010 \pm 0.001$	$3.33 \pm 0.20 \pm 0.02$
Ω_{bc}	$0.193 \pm 0.011 \pm 0.001$	$3.94 \pm 0.23 \pm 0.01$	$0.171 \pm 0.010 \pm 0.000$	$3.50 \pm 0.21 \pm 0.01$

Numerical results

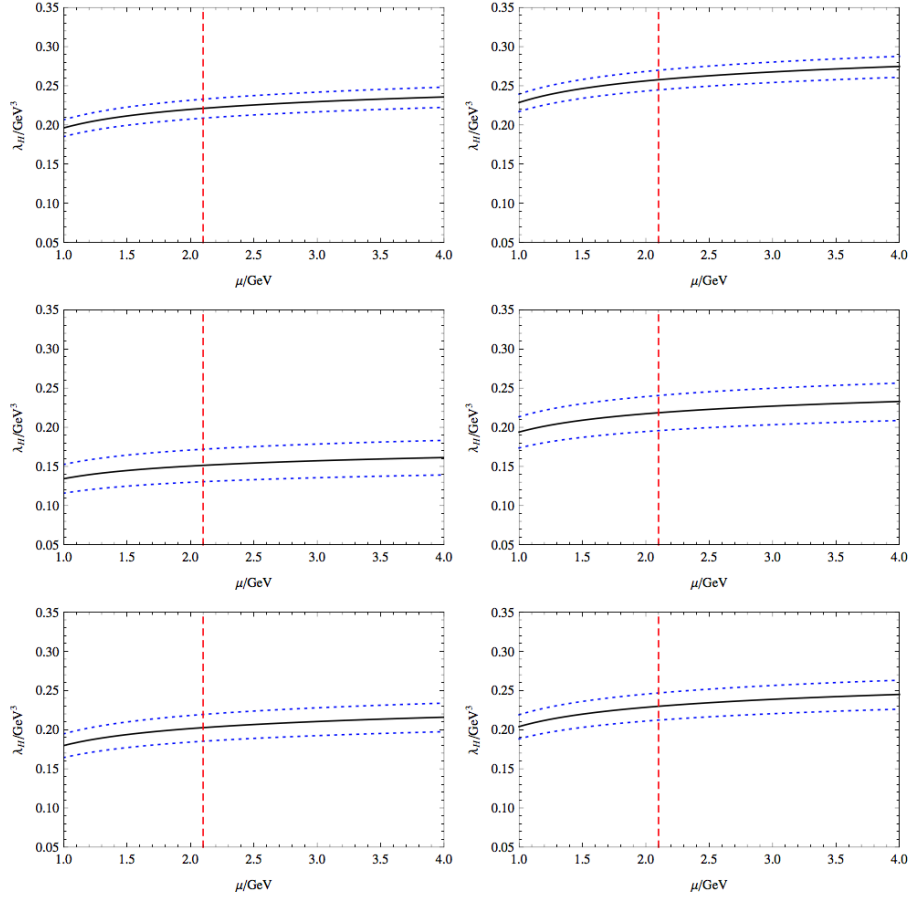


FIG. 4: The scale dependence (see Eq. (38)) of the “decay constants” for Ξ_{cc} , Ω_{cc} (the top two figures), Ξ_{bb} , Ω_{bb} (the middle two figures), Ξ_{bc} and Ω_{bc} (the bottom two figures) with the scale μ ranging from 1 GeV to 4 GeV. The solid line corresponds to $\sqrt{s_0} = 4.1$ GeV for Ξ_{cc} , $\sqrt{s_0} = 4.2$ GeV for Ω_{cc} , $\sqrt{s_0} = 10.6$ GeV for Ξ_{bb} , $\sqrt{s_0} = 10.8$ GeV for Ω_{bb} , $\sqrt{s_0} = 7.4$ GeV for Ξ_{bc} , $\sqrt{s_0} = 7.5$ GeV for Ω_{bc} , respectively. The dotted curves are obtained by varying the $\sqrt{s_0}$ by 0.1 GeV. The vertical line corresponds to the scale $\mu = 2.1$ GeV. The sum rule in Eq. (19) is considered.

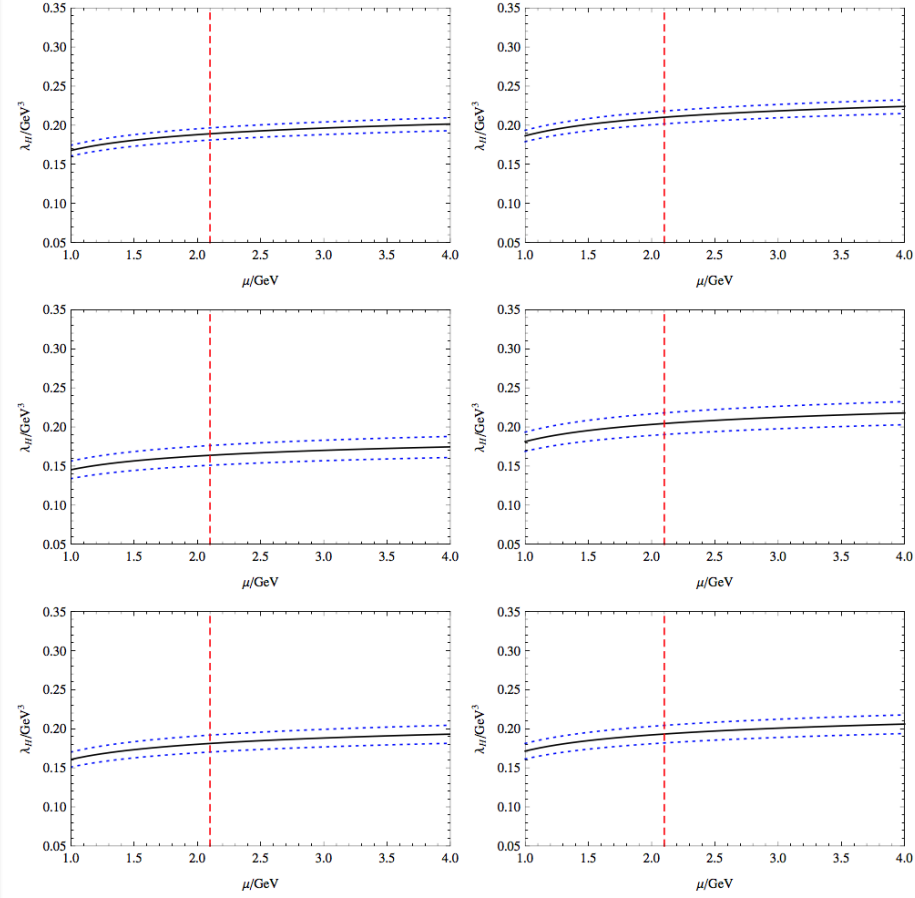


FIG. 8: Same as Figure 4 but for the sum rule in Eq. (34) is considered.

$$\lambda_H(\mu) = \lambda_H(\mu_0) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{\gamma_{\lambda_0}}{2\beta_0}}$$

$$\frac{\gamma_{\lambda_0}}{2\beta_0} = \frac{-6}{25}$$



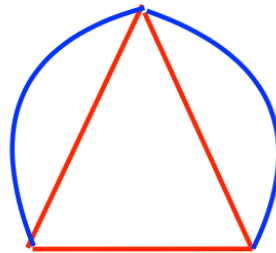
- It is necessary to point out that when including the contributions from the $1/2^-$ baryons the threshold parameter might be somewhat higher. In this analysis, we have approximately use the same values.
- Comparing the results in Table III and Table IV, one can see that the negative parity baryons do not provide significant modifications.
- From Table III and Table IV, we can see that, the decay constant of $\Omega_{QQ'}$ are slightly larger than that of $\Xi_{QQ'}$.
- The decay constants increases with the energy scale.



Summary and outlook

- QCD sum rules
- $\Xi_{cc}, \Omega_{cc}, \Xi_{bb}, \Omega_{bb}, \Xi_{bc}, \Omega_{bc}$
- Decay constants
- Positive parity and negative parity

- Lifetime
- ...





Thank you for your attention!

backup



$$\begin{aligned}\text{Subscript}[\gamma, m_0] &= -4, \\ \text{Subscript}[\beta, 0] \text{ (f = 4)} &= 25/3, \\ \text{Subscript}[\gamma, m_0]/(2 \text{ Subscript}[\beta, 0]) &= -6/25\end{aligned}$$