

External masses effect to $Z\bar{b}s$ coupling

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2018 年 6 月 22 日 @ SJTU, 中国物理学会高能物理分会第十届全国会员代表大会暨学术年会

Overview

Background

$g_{Z\bar{b}s}$
Mass effect

Conclusion

Background

- † Inami-Lim function is widely used to deal with B meson decays
- † large m_t , the limit of vanishing external momentum,
 $B_0, C_0, D_0, E_0, D'_0, E'_0$
- † In B meson decay, the invariant mass of boson is much less than m_t , the limit is good

- † But in b production, ie., BK production in Z factory (CEPC), we should revisit C_0
- † At least the m_Z should be considered, how well the limit works

the first glance at $Z \rightarrow \bar{b}s$

- $g_{Z\bar{b}s} = \sum_i V_{is} V_{ib}^* \left[g_0 + m_i^2/m_W^2 \left(g_1 + g_2 \ln \frac{m_i^2}{m_W^2} \right) + \dots \right]$
- Suppressed by the Glashow-Iliopoulos-Maini (GIM) mechanism.
- do occur in electroweak interaction at one-loop penguin level
- Gauge condition: **The 't' Hooft-Feynman gauge** ($\xi = 1$ in R_ξ)
 - † vector boson is proportional to $g_{\mu\nu}$, the simplest form
 - † the unphysical scalar boson has a pole at $p^2 = M^2$
 - † the negative-metric scalar-boson pole of the vector-boson propagator is cancelled by the unphysical-scalar-boson propagator
 - † the masses of unphysical scalars are taken on their physical partner.
- **Dimension regularization with the completely anticommuting γ_5** and the 't Hooft convention of γ_5

$$Z \rightarrow \bar{b}s$$

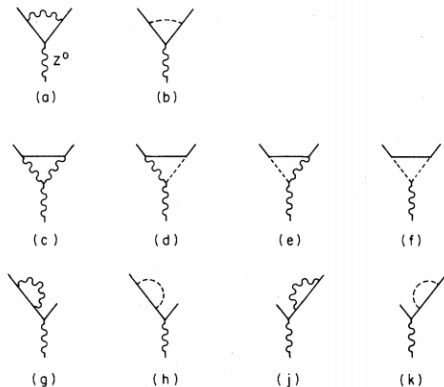


FIG. 1. Feynman diagrams contributing to $Z^0 \rightarrow q\bar{x}$ in the 't Hooft-Feynman gauge. The internal quarks must have weak isospin different from q and x . Wavy internal lines represent W^\pm ; dashed internal lines represent the associated unphysical scalars.

looking at $Z \rightarrow \bar{b}s$

- no divergence appears in Fig1(d,e)
- Unitarity of the charged current mixing matrix, each divergence in Fig1(a,c,j,g) becomes finite
- divergences in Fig1(b,f,h,k) cancels for each internal flavour
- the couplings in Fig1(b,f,h,k) are proportional to $\mathcal{O}(m_i^2/m_W^2)$, zero power ($\propto 1$) in Fig1(a,c,j,g)
- the $\sin \theta_W$ dependence in individual diagram disappear finally

† $\sin \theta_W$ represents the coupling of photon to $\bar{b}\gamma_\mu(1 - \gamma_5)s$

† the related $\phi\bar{b}s$ vertex does not have this dependence

Effects of the external momentum

Taking Fig.1(b) for example, **Preliminary**

$$\begin{aligned}
 Z^\mu \bar{b}(k_1) g_{Zbs}^{(b)} s(k_2) &= Z^\mu \bar{b} \frac{g^3 V_{tb}^* V_{ts}}{4 c_{\theta_W}} \frac{m_t^2}{m_W^2} \left\{ \left([(2-\epsilon) C_{g_{\mu\nu}}(\bar{k}'_2, -\bar{k}'_1) + (m_b^2 + m_s^2) C_{pp} \right. \right. \\
 &\quad \left. \left. - (m_Z^2 - m_b^2) C_{p_1 p_2} - 2m_s^2 C_1 \right) \left(-\frac{4}{3} s_{\theta_W}^2 \right) \right. \\
 &\quad \left. - [m_t^2 C_0 - m_b^2 (C_0 - C_2) + m_s^2 (C_0 - C_1)] \left(1 - \frac{4}{3} s_{\theta_W}^2 \right) \right\} \gamma_\mu L \\
 &\quad + m_b m_s \left[2C_2 \left(-\frac{4}{3} s_{\theta_W}^2 \right) + (C_1 - C_2) \left(1 - \frac{4}{3} s_{\theta_W}^2 \right) \right] \gamma_\mu R \\
 &\quad - 2m_b C_{pp} \left(-\frac{4}{3} s_{\theta_W}^2 \right) k_{1,\mu} L \\
 &\quad - 2m_s \left[C_{p_1 p_2} \left(-\frac{4}{3} s_{\theta_W}^2 \right) - C_2 \left(1 - \frac{4}{3} s_{\theta_W}^2 \right) \right] k_{1,\mu} R \\
 &\quad + 2m_b \left[C_{p_1 p_2} \left(-\frac{4}{3} s_{\theta_W}^2 \right) + C_1 \left(1 - \frac{4}{3} s_{\theta_W}^2 \right) \right] k_{2,\mu} L \\
 &\quad - 2m_s [2C_1 - C_{pp}] \left(-\frac{4}{3} s_{\theta_W}^2 \right) k_{2,\mu} R \\
 &\quad \left. - \frac{\epsilon}{4} B_0 (-p'_s - p'_b) \left(1 - \frac{4}{3} s_{\theta_W}^2 \right) \gamma_\mu L \right\} s
 \end{aligned}$$

Effects of the external momentum

- † C_0 ; $C_\mu(k_1, k_2) = k_{1\mu} C_1 + k_{2\mu} C_2$;
† $C_{\mu\nu}(k_1, k_2) = g_{\mu\nu} C_g + (k_{1\mu} k_{1\nu} + k_{2\mu} k_{2\nu}) C_{kk} + (k_{1\mu} k_{2\nu} + k_{1\nu} k_{2\mu}) C_{k_1 k_2}$
- C_g, B_1 are the leading power of $\mathcal{O}(m_t^2)$
- $m_W = 80.4 \text{ GeV}, m_t = 173 \text{ GeV}$
- **Preliminary**
 - † Varying k_1^2, k_2^2 does not change B_1, C_g
 - † $\gamma_\mu R, k_\mu$ terms are four and five order of magnitude smaller
 - † The mass effect mainly come from the m_Z terms

Masses (GeV)	Figs. (a+c+g+j)	Figs. (b+f+h+k)	Figs. (d+e)	All Figs
$m_Z = 0, m_b = 0, m_s = 0$	1.14	1.71	0.56	3.40
$m_Z = 91.2, m_b = 4.2, m_s = 0.095$	0.78	1.67	0.61	3.07
$m_Z = 0, m_b = 4.2, m_s = 0.095$	1.14	1.71	0.56	3.40

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The End, Thanks.