Global and local Λ polarization in heavy-ion collisions

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Contents

• Λ global polarization

- why global polarization decreases with increasing $\sqrt{s_{\rm NN}}$.

• Λ local polarization

- rich information for local vorticity field.

Summary

Li, Pang, Wang, XLX, Phys. Rev. C 96, no. 5, 054908 (2017) XLX, Li, Tang, Wang, arXiv:1803.00867

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Λ global polarization

Brief history of Λ global polarization

► Liang, Wang 2004, Voloshin 2004 Orbital angular momentum in non-central collisions, spin-orbital coupling ⇒ global polarization.



global polarization: average effect of ω field on hyper-surface (**net vorticity**)

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Brief history of Λ global polarization





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Λ polarization probes $\boldsymbol{\omega}$ and \mathbf{B}

Statistical-hydro model
 Becattini etal 2008, 2013
 Fang, Pang, Wang, Wang 2016

For small fields

$$\mathbf{P}_{\Lambda} \simeq rac{oldsymbol{\omega}}{2 \, T} + rac{\mu_{\Lambda} \mathbf{B}}{T} \ \mathbf{P}_{ar{\Lambda}} \simeq rac{oldsymbol{\omega}}{2 \, T} - rac{\mu_{\Lambda} \mathbf{B}}{T}$$

Becattini, Karpenko, Lisa, Upsal, Voloshin 2017

• Important inputs for studies on effects of ω and B:

- ► CME/CVE/CMW ...
- effects on QCD phase diagram.

I focus on $\omega\text{-induced }\Lambda$ polarization in this talk.

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▶

STAR result 2017



- global polarization: effect of net vorticity
- the most vortical fluid:

 $\langle \omega_{\rm y} \rangle \sim 0.01 \, {\rm fm}^{-1} \sim 10^{21} \, {\rm s}^{-1}$

$$\mathbf{P}_{\Lambda} \simeq rac{oldsymbol{\omega}}{2T} + rac{\mu_{\Lambda}\mathbf{B}}{T}$$
 $\mathbf{P}_{ar{\Lambda}} \simeq rac{oldsymbol{\omega}}{2T} - rac{\mu_{\Lambda}\mathbf{B}}{T}$

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Λ global polarization vs $\sqrt{s_{\rm NN}}$



- Λ global polarization decreases with increasing $\sqrt{s_{NN}}$.
- In contrast to total angular momentum vs $\sqrt{s_{\rm NN}}$.

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Λ global polarization in varied models





Baznat etal 2017

Sun, Ko 2017

(s_{NN})^{1/2} (GeV)

Au+Au, 20-50%

– ∧ and ⊼

100

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Consistent with $\langle \omega_y \rangle (\sqrt{s_{\rm NN}})$





- (ω_y): averaged vorticity, weighted by matter or energy density (net vorticity)
- $\langle \omega_y \rangle$ decreases with $\sqrt{s_{\rm NN}}$.

Vorticity distribution

 ω_v on zx plane at t = 5 fm/c in 20-30% Au+Au collisions:



• Net vorticity $\langle \omega_{\gamma} \rangle$ is larger at lower $\sqrt{s_{NN}}$.

Quadrupole pattern: local vorticity (see next section).

How to understand?

Net vorticity and global polarization are generated from the fireball's *titled shape* on the reaction plane.



Central vs non-central collisions

In central collisions, net vorticity vanishes due to the reflection symmetry x → −x and z → −z.



 \blacktriangleright In non-central collisions, the reflection symmetry breaks $\langle \omega_{\rm y} \rangle \neq 0$

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Vorticity and matter distribution

 ω_y on zx plane at t = 5 fm/c in 20-30% Au+Au collisions:



Initial parton distribution:



Λ global polarization vs centrality



STAR 1805.04400

- Λ global polarization increases with centrality,
- consistent to the tilted shape scenario.

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Local vorticity and Λ local polarization

Circular transverse vorticity $\boldsymbol{\omega}_{\perp} = (\omega_x, \omega_y)$





On xz plane: quadrupole ω_y In 3D view: circular ω_{\perp}

$$\mathbf{v}_{\perp} = \mathbf{v}_{\perp}(\mathbf{z})\mathbf{e}_{\mathbf{r}}$$

$$oldsymbol{\omega}_{\perp} \equiv rac{1}{2} (
abla imes \mathbf{v})_{\perp} = rac{1}{2} \partial_z \mathbf{v}_{\perp}(z) \mathbf{e}_{\phi}$$

 $\partial \mathbf{v}_{\perp} / \partial |z| < 0.$

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Circular transverse polarization $\mathbf{P}_{\perp} = (P_x, P_y)$





- $$\begin{split} \langle \mathbf{P}_{\perp} \cdot \operatorname{sign}(\mathbf{Y}) \rangle &= \frac{\mathbf{P}_{\perp}(\phi_p, \mathbf{Y} > 0) \mathbf{P}_{\perp}(\phi_p, \mathbf{Y} < 0)}{2}, \ |\mathbf{Y}| < 1 \\ \phi_p: \text{ azimuthal angle of } \Lambda \text{'s momentum.} \\ \mathbf{Y}: \ \Lambda \text{'s rapidity.} \end{split}$$
 - ▶ insensitive with √s_{NN}, can be tested at 200 GeV and LHC energy.

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▶ quadrupole P_z from ∂_φv_⊥
 ▶ circular P_⊥ from ∂_zv_⊥
 ▶ Becattini, Karpenko 2017
 ▶ Voloshin 2017
 ▶ above slide from Zhoudunming Tu's talk at QM2018

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bove slide from Zhoudunming Tu's talk at QM2 CHEP2018

Quadrupole pattern of P_z



XLX, Li, Tang, Wang 2018 AMPT Becattini, Karpenko 2017 hydro+UrQMD

consistent with Becattini and Karpenko's result.

Local polarization, harmonic behaviors



 $(P_x, P_y) \cdot \operatorname{sign}(Y) \sim -\mathbf{e}_{\phi} \sim (\sin \phi_p, -\cos \phi_p);$ $P_z \sim -\sin(2\phi_p)$

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Local polarization, Fourier coefficients



 Local polarization is sizable, c.f. global polarization < 0.003 at 200 GeV and smaller at LHC energies.

• P_z and $|P_x - P_y|$ increases with centrality: elliptic flow.

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Summary

- The net vorticity and global polarization are from the fireball's *titled shape* on reaction plane.
 - Decreases with $\sqrt{s_{NN}}$ in mid-rapidity;
 - Increases with centrality.
- Local vorticity and local polarization:
 - circular structure of ω_{\perp} and \mathbf{P}_{\perp} , from $\partial_z v_{\perp}$.
 - quadrupole pattern of ω_z and P_z , from $\partial_{\phi} v_{\perp}$.
- Local polarization is sizable at 200 GeV and LHC energies, can be tested.

Thanks for listening

Backup

How to measure Λ polarization?

 $\blacktriangleright~\Lambda$ is "self-analyzing"



Weak decay $\Lambda \to {\it p} + \pi^-$

$$\frac{dN}{d\cos\theta^*} = \frac{1}{2}(1 + \alpha_H |\mathbf{P}_H| \cos\theta^*)$$

$$\label{eq:alpha} \begin{split} \alpha_\Lambda &= -\alpha_{\bar{\Lambda}} = 0.642 \pm 0.013 \\ \text{Xiao-Liang Xia} \qquad \qquad \text{CHEP2018} \end{split}$$

Λ polarization induced by $\boldsymbol{\omega}$ and \mathbf{B}

- ► Polarization from $\boldsymbol{\omega}$ and \mathbf{B} $\rho \sim \exp\left(\frac{\boldsymbol{\omega} \cdot \widehat{\mathbf{S}} + \mathbf{B} \cdot \mu_{\Lambda} \widehat{\mathbf{S}} / S}{T}\right)$ $\mathbf{P} \equiv \langle \mathbf{S} \rangle / S = \operatorname{tr}(\rho \widehat{\mathbf{S}}) / S$ For small fields $\mathbf{P}_{\Lambda} \simeq \frac{\boldsymbol{\omega}}{2T} + \frac{\mu_{\Lambda} \mathbf{B}}{T}$ $\mathbf{P}_{\bar{\Lambda}} \simeq \frac{\boldsymbol{\omega}}{2T} - \frac{\mu_{\Lambda} \mathbf{B}}{T}$ $\mu_{\Lambda} = -\mu_{\bar{\Lambda}} = -0.6138 \mu_N$
 - $\boldsymbol{\omega}$ -induced $\Lambda/\bar{\Lambda}$ polarizations are along $\boldsymbol{\omega}$;
 - **B**-induced $\Lambda/\bar{\Lambda}$ polarizations are along \mp **B**.

$$\boldsymbol{\omega} = (\mathbf{P}_{\Lambda} + \mathbf{P}_{\bar{\Lambda}}) T$$
$$\mathbf{B} = (\mathbf{P}_{\Lambda} - \mathbf{P}_{\bar{\Lambda}}) T / (2\mu_{\Lambda})$$

 $\frac{\Lambda/\bar{\Lambda} \text{ polarization can be used to probe } \omega \text{ and } B.}{\text{Becattini, Karpenko, Lisa, Upsal, Voloshin 2017}}$ Xiao-Liang Xia CHEP2018

B-field from Λ global polarization





above picture from Lisa's talk at Florence workshop, 2018

Uncertainties on both experimental and theoretical sides.