

Suppression of high p_T single hadrons and dihadrons in heavy-ion collisions at $\sqrt{s_{NN}} = 0.2, 2.76, 5.02$ and 5.44 TeV

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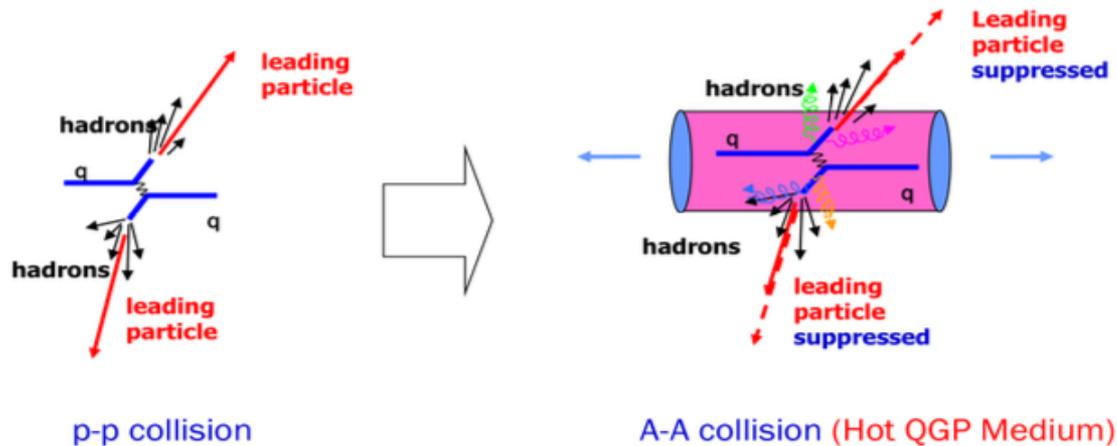
Outline

- Introduction
- NLO pQCD, JQ mFFs, and Hydro models
- Extracted \hat{q}_0 via single and dihadron suppression in HIC
- Prediction for dihadron suppression factor at 5.02 and 5.44TeV
- Summary and outlook

Introduction

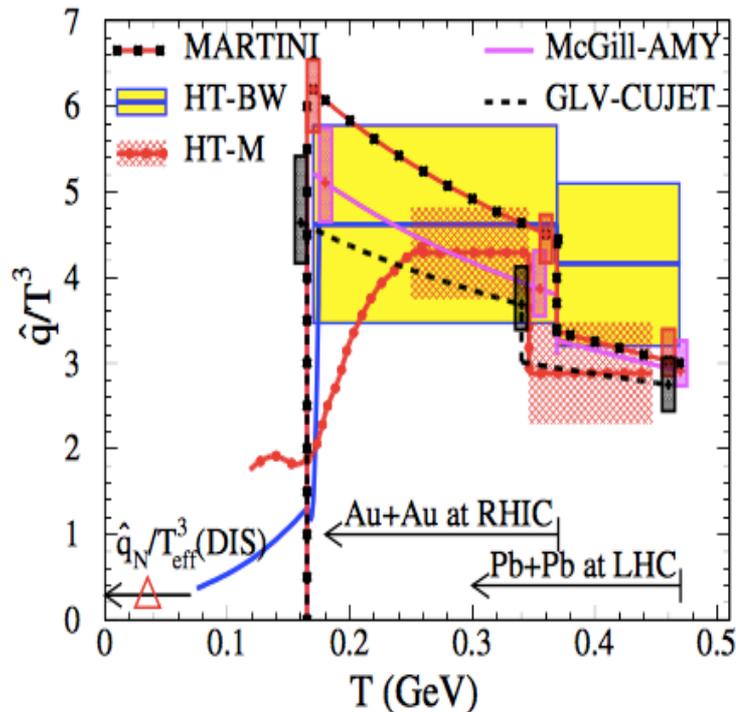
- **Jet quenching:** [X.-N. Wang and M. Gyulassy, Phys. Rev. Lett. 68, 1480(1992)]

The hard jet loses a large amount of its energy via radiating gluon induced by multiple scattering.



- Jet energy loss in the QGP medium: $\Delta E \propto \hat{q}$ --- jet transport coefficient.
- $\hat{q} \equiv \frac{d\langle q_T^2 \rangle}{dL}$: transverse momentum broadening squared per unit length.
[BDMPS, NPB482(1997)291]

Motivation



In the initial time and in the center of the fireball

$$\hat{q} \approx \begin{cases} 1.2 \pm 0.3 \\ 1.9 \pm 0.7 \end{cases} \text{ GeV}^2/\text{fm} \text{ at } \begin{matrix} T=370 \text{ MeV,} \\ T=470 \text{ MeV,} \end{matrix}$$

[The JET collaboration.arXiv:1312.5003,Dec 2013]

- Extracting \hat{q}_0 via single hadron and dihadron suppression in HIC.
- **Whether the dihadron \hat{q}_0 is consistent with the single hadron?**
- Predicting the dihadron suppression factor I_{AA} in central and non-central Pb+Pb collisions at $\sqrt{s_{NN}}=5.02\text{TeV}$ and Xe+Xe collisions at $\sqrt{s_{NN}}=5.44\text{TeV}$.

NLO pQCD parton model

- In A+A collisions, dihadron spectra can be given by:

$$\begin{aligned} \frac{d\sigma}{dy_1 d^2 p_T^{h1} dy_2 d^2 p_T^{h2}} &= \sum_{abcd} \int d^2 b d^2 r t_A(\vec{r}) t_B(|\vec{r} - \vec{b}|) dx_a dx_b dz_c dz_d \\ &\times f_{a/A}(x_a, \mu^2, \vec{r}) f_{b/B}(x_b, \mu^2, (|\vec{r} - \vec{b}|)) \frac{\hat{s}}{2\pi z_c^2 z_d^2} \times \frac{d\sigma}{d\hat{t}}(ab \rightarrow cd) \\ &\times D_{h/c}(z_c, \mu^2, \Delta E_c) D_{h/d}(z_d, \mu^2, \Delta E_d) \delta^4(p_a + p_b - p_c - p_d) + O(\alpha_s^3) \end{aligned}$$

- $t_A(\vec{r})$ → thickness function: Woods-Saxon
- $f_{a/A}$ → PDF: CTEQ6.6M, EPS09 [[Eur. Phys. J. C 12, 375\(2000\)](#)], [[JHEP 0904, 065 \(2009\)](#)]
- $\frac{d\sigma}{d\hat{t}}$ → hard scattering cross section
- $D_{h/c}^0$ → FFs: Kretzer [[PHYSICAL REVIEW D, VOLUME 62, 054001](#)]

Modified fragmentation functions--mFFs

- Modified fragmentation functions in QGP medium:

$$D_{h/c}(z_c, \mu^2, \Delta E_c) = (1 - e^{-\langle N_g \rangle}) \left[\frac{z_c'}{z_c} D_{h/c}^0(z_c', \mu^2) + \langle N_g \rangle \frac{z_g'}{z_c} D_{h/g}^0(z_g', \mu^2) \right] \\ + e^{-\langle N_g \rangle} D_{h/c}^0(z_c, \mu^2)$$

where $z_c' = p_T / (p_{Tc} - \Delta E_c)$, $z_g' = \langle N_g \rangle p_T / \Delta E_c$.

[X.-N. Wang, PRC70 (2004) 031901], [H. Z. Zhang, J.F. Owens, Phys. Rev. Lett. 98.212301 (2007)], and [H. Z. Zhang, J.F. Owens, Phys. Rev. Lett. 103, 032302 (2009)]

- Total energy loss of jet in high-twist method:

$$\frac{\Delta E}{E} = C_A \frac{\alpha_s}{2\pi} \int d\tau \int_0^{Q^2} \frac{dl_T^2}{l_T^4} \int_\epsilon^{1-\epsilon} dz [1 + (1-z)^2] \times \hat{q}_F(y) 4 \sin^2\left(\frac{l_T^2 \tau}{4z(1-z)E}\right)$$

[W.T. Deng and X.-N. Wang, Phys. Rev. C81,024902(2010)], [E. Wang and X.-N. Wang, Phys. Rev. Lett. 87, 142301 (2001); 89, 162301 (2002)]

2+1 D hydrodynamics medium model

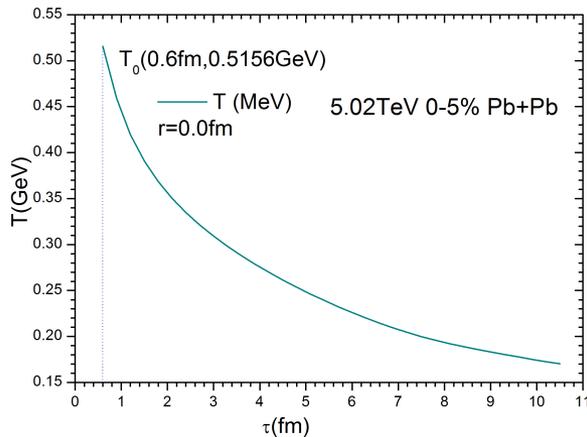
- \hat{q} depends on the local T in the jet trajectory:

$$\hat{q}(\tau, r) = \left[\hat{q}_0 \frac{\rho_g(\tau, r)}{\rho_g(\tau_0, 0)} (1-f) \right] \frac{p^\mu u_\mu}{p_0} \quad \rho_g \propto T^3 \rightarrow \boxed{\frac{\hat{q}}{T^3} = \frac{\hat{q}_0}{T_0^3} = C}$$

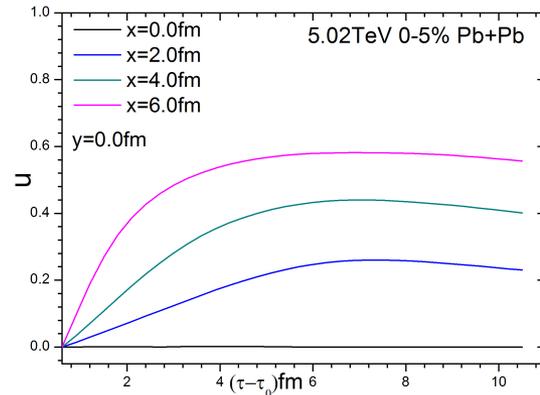
- hadron phase fraction: $f(\tau, r) = \begin{cases} 0 & \text{when } T \geq 170 \text{ MeV} \\ 1 & \text{when } T \leq 170 \text{ MeV} \end{cases}$

[X. Chen, T. Hirano, E. Wang, X.-N. Wang, H. Z. Zhang, Phys. Rev. C84 (2011) 034902]

- T, u are given by 2+1 D hydro to describe the medium evolution. [Z. Qiu, C. Shen and U. Heinz, Phys. Lett. B 707, 151 (2012)]



Medium Temperature



Medium Velocity

Nuclear modification factor

- Single hadron nuclear modification factor:

$$R_{AA} = \frac{dN_{AA}/dyd^2p_T}{T_{AA}(b)d\sigma_{pp}/dyd^2p_T}$$

- Dihadron nuclear modification factor:

$$I_{AA}(z_T) = \frac{D_{AA}(z_T)}{D_{pp}(z_T)} = \frac{D_{AA}(p_T^{assoc})}{D_{pp}(p_T^{assoc})}$$

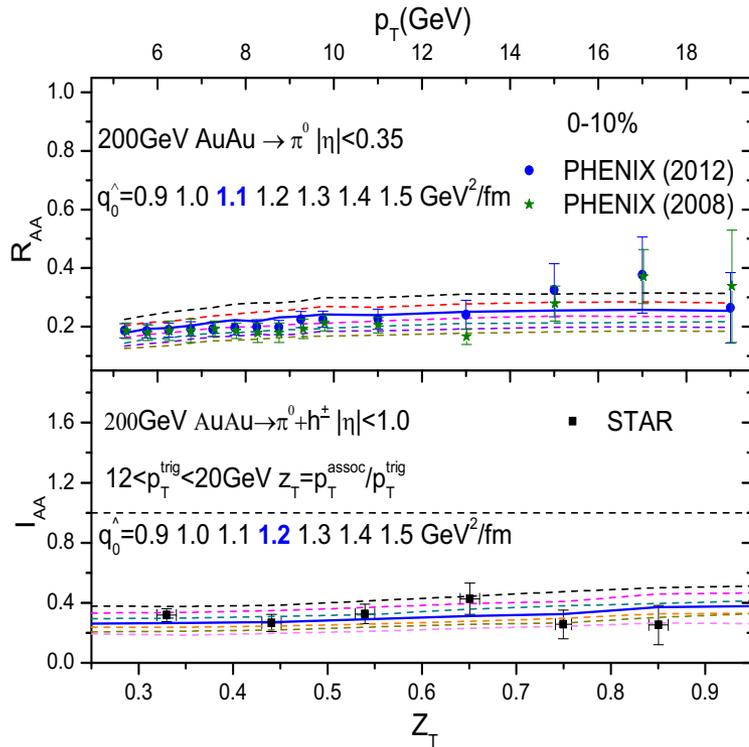
where $z_T = \frac{p_T^{assoc}}{p_T^{trig}}$.

$$D_{AA}(z_T) \equiv \frac{1}{N_{AA}^{trig}} \frac{dN_{AA}^{h_1 h_2}}{dz_T} = \frac{\int dp_T^{trig} dy^{trig} \int dy^{assoc} d\Delta\phi \frac{p_T^{trig} d^5\sigma_{AA}^{h_1 h_2}}{dp_T^{trig} dy^{trig} dp_T^{assoc} dy^{assoc} d\Delta\phi}}{\int dp_T^{trig} dy^{trig} \frac{d^2\sigma_{AA}^{h_1}}{dp_T^{trig} dy^{trig}}}$$

$$\text{or } D_{AA}(p_T^{assoc}) \equiv \frac{1}{N_{AA}^{trig}} \frac{dN_{AA}^{h_1 h_2}}{dp_T^{assoc}} = \frac{\int dp_T^{trig} dy^{trig} \int dy^{assoc} d\Delta\phi \frac{d^5\sigma_{AA}^{h_1 h_2}}{dp_T^{trig} dy^{trig} dp_T^{assoc} dy^{assoc} d\Delta\phi}}{\int dp_T^{trig} dy^{trig} \frac{d^2\sigma_{AA}^{h_1}}{dp_T^{trig} dy^{trig}}}$$

Extracting \hat{q}_0 via dihadron suppression in HIC

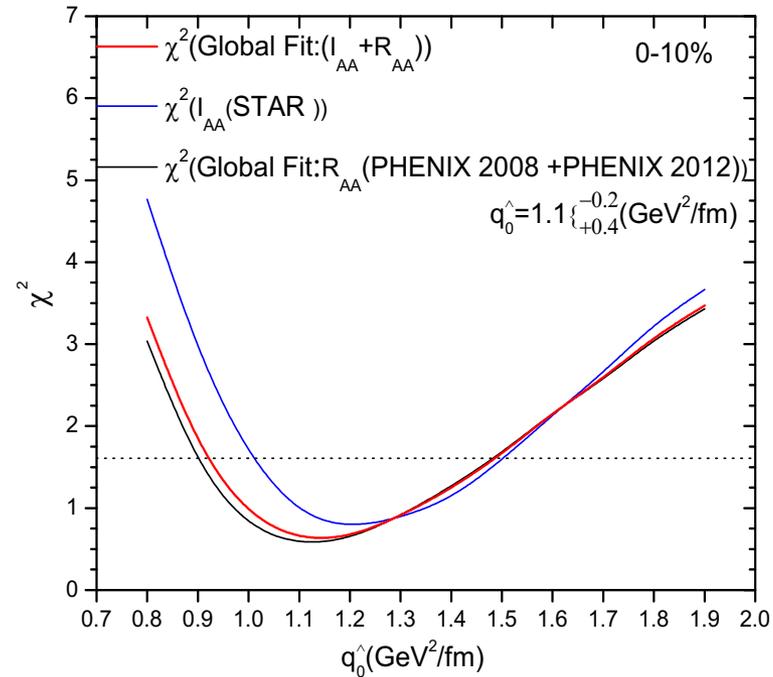
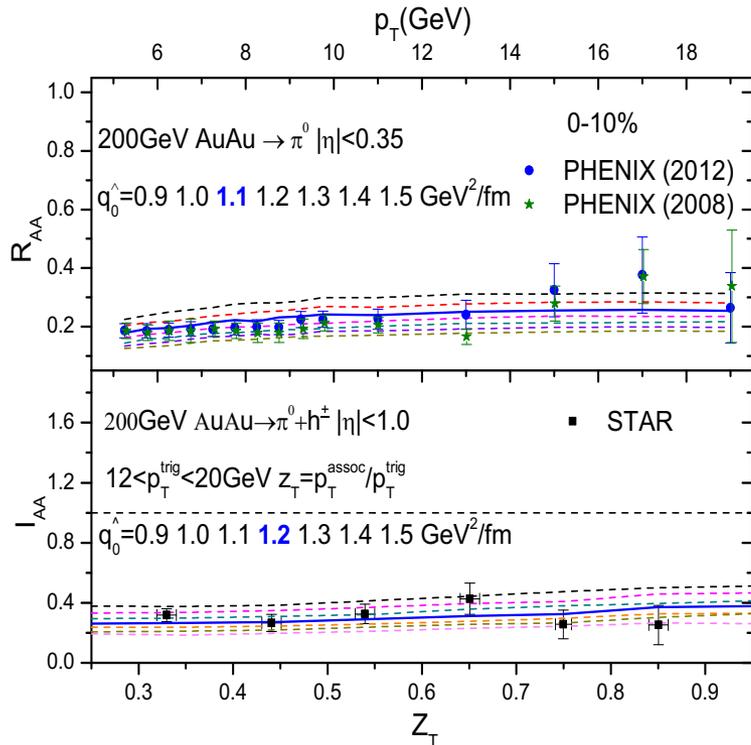
- Extracting parameters \hat{q}_0 in central Au+Au collisions at $\sqrt{s_{NN}}=200$ GeV.



Data from [[PHENIX Collaboration], Phys. Rev. C 87, no. 3,034911(2013),
 Phys. Rev. Lett. 101, 232301(2008)], and [[STAR], arXiv: 1604. 01117v2 [nucl-ex] 7 Apr 2016]

Extracting \hat{q}_0 via dihadron suppression in HIC

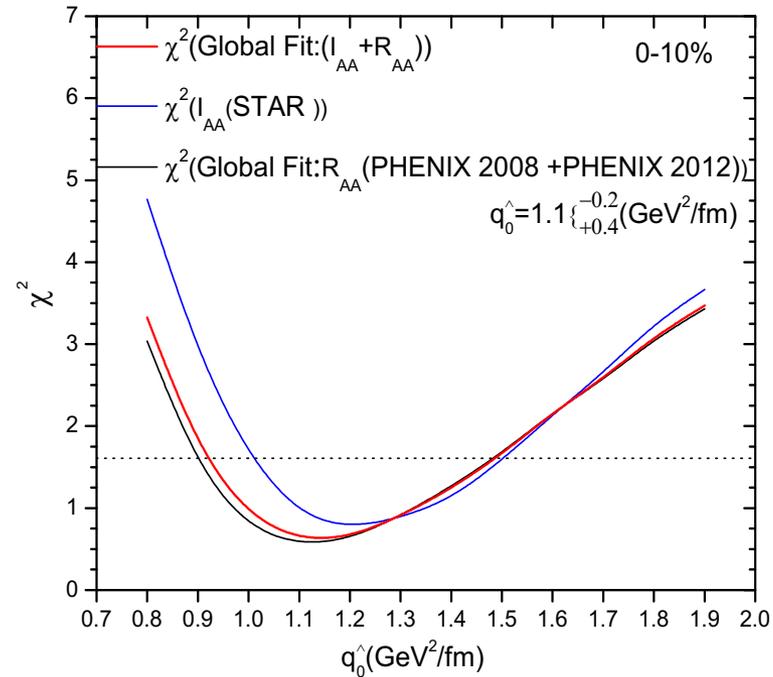
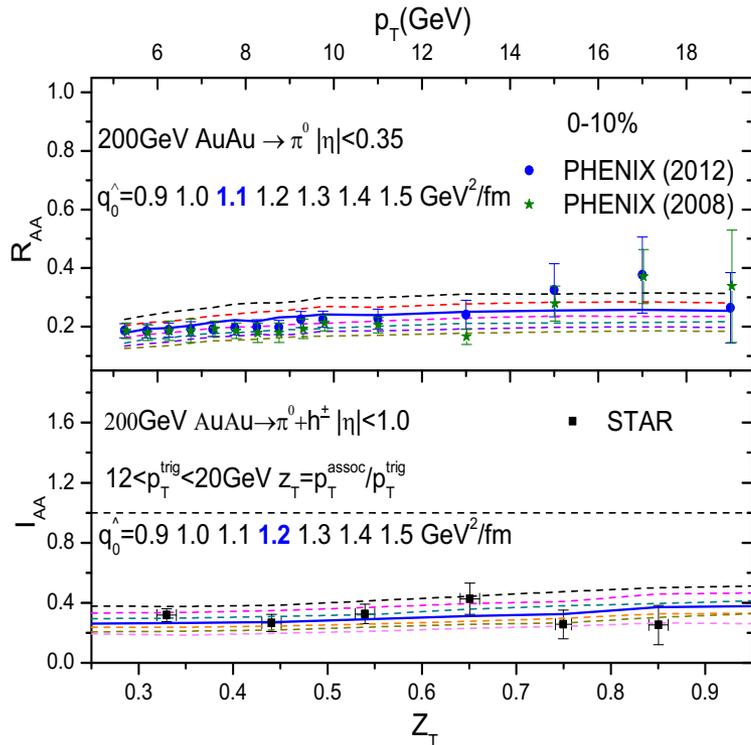
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Extracting \hat{q}_0 via dihadron suppression in HIC

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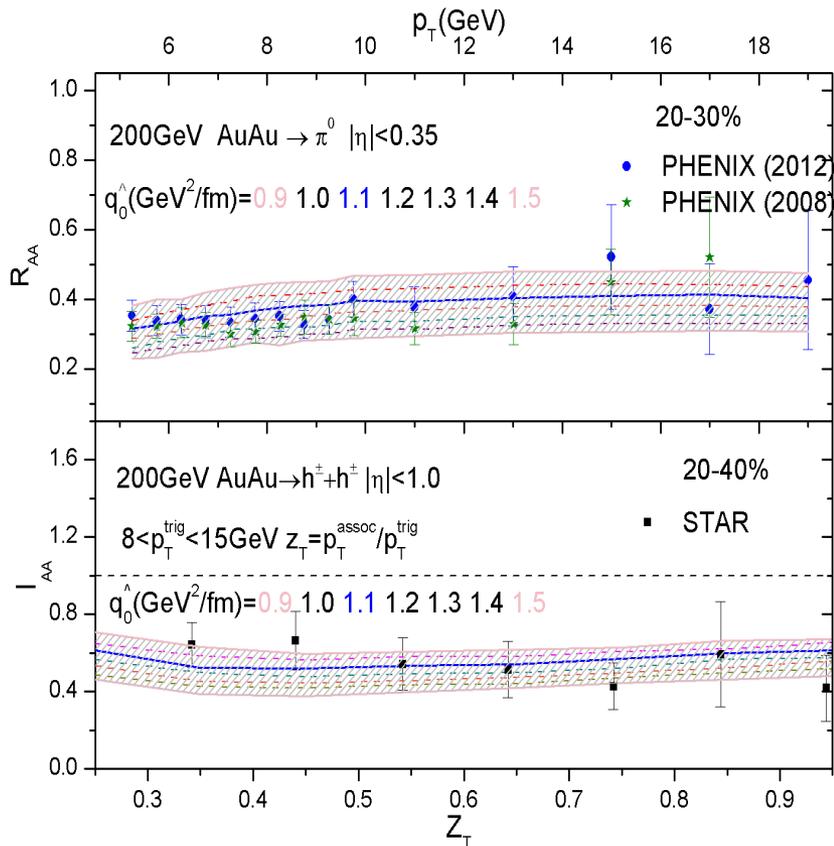


$$\hat{q}_0 = 1.1 \{_{+0.4}^{-0.2} (\text{GeV}^2/\text{fm})$$

Data from [[PHENIX Collaboration], Phys. Rev. C 87, no. 3, 034911(2013), Phys. Rev. Lett. 101, 232301(2008)], and [[STAR], arXiv: 1604.01117v2 [nucl-ex] 7 Apr 2016]

Extracting \hat{q}_0 at RHIC in the 20–40% centrality

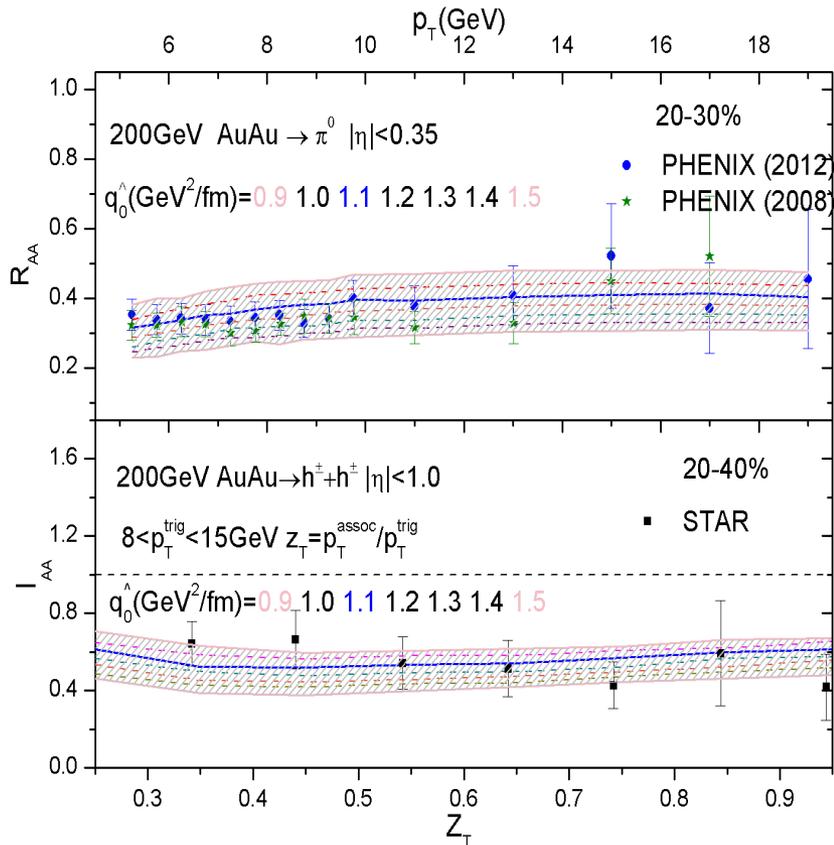
- R_{AA} and I_{AA} in the 20–40% centrality with the \hat{q}_0 values extracted from central Au+Au collisions as compared to experimental data.



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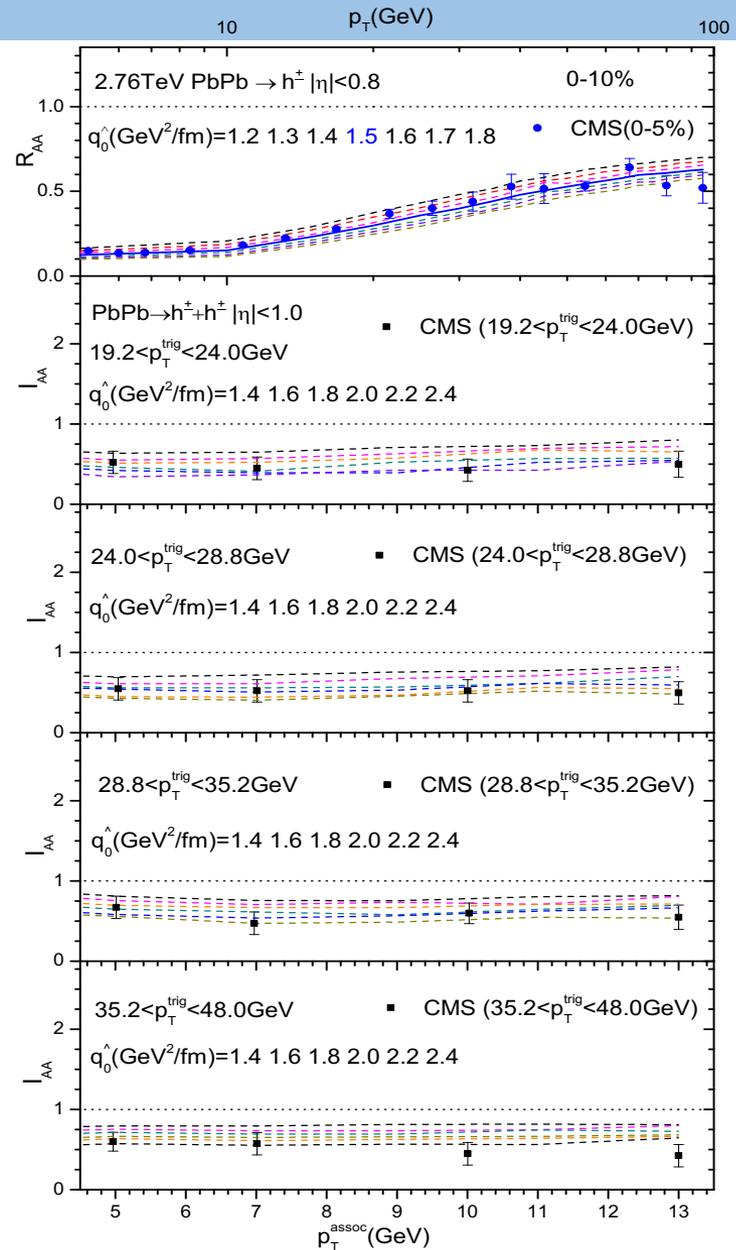
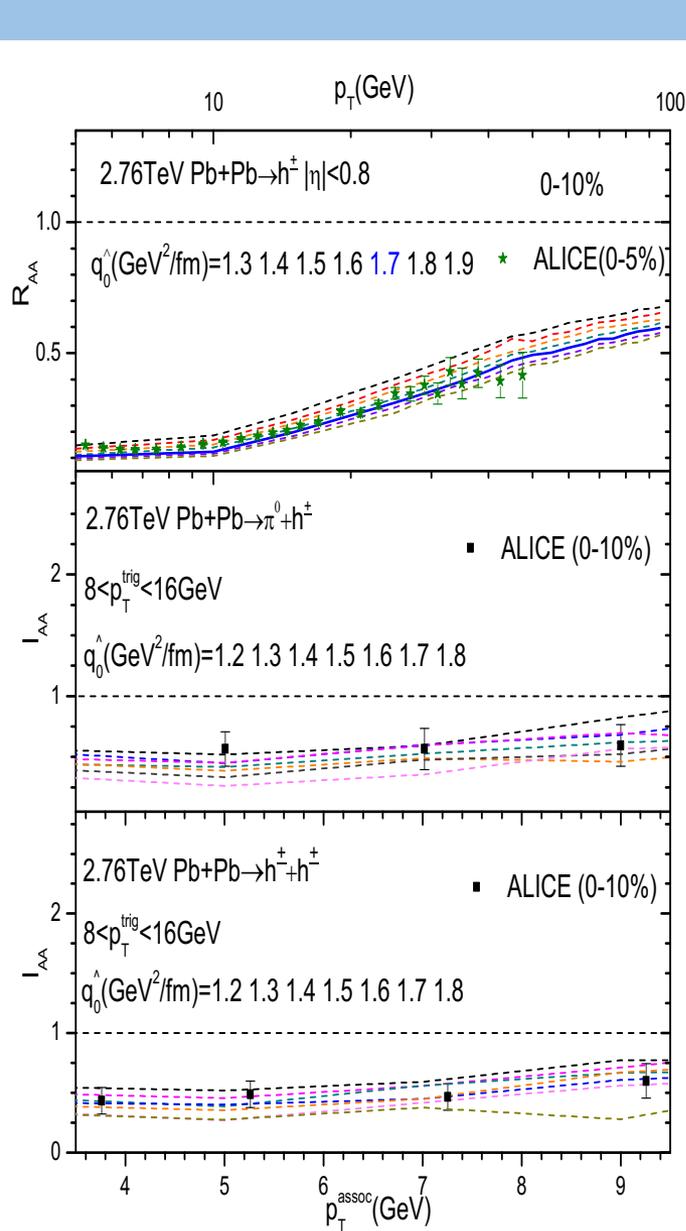
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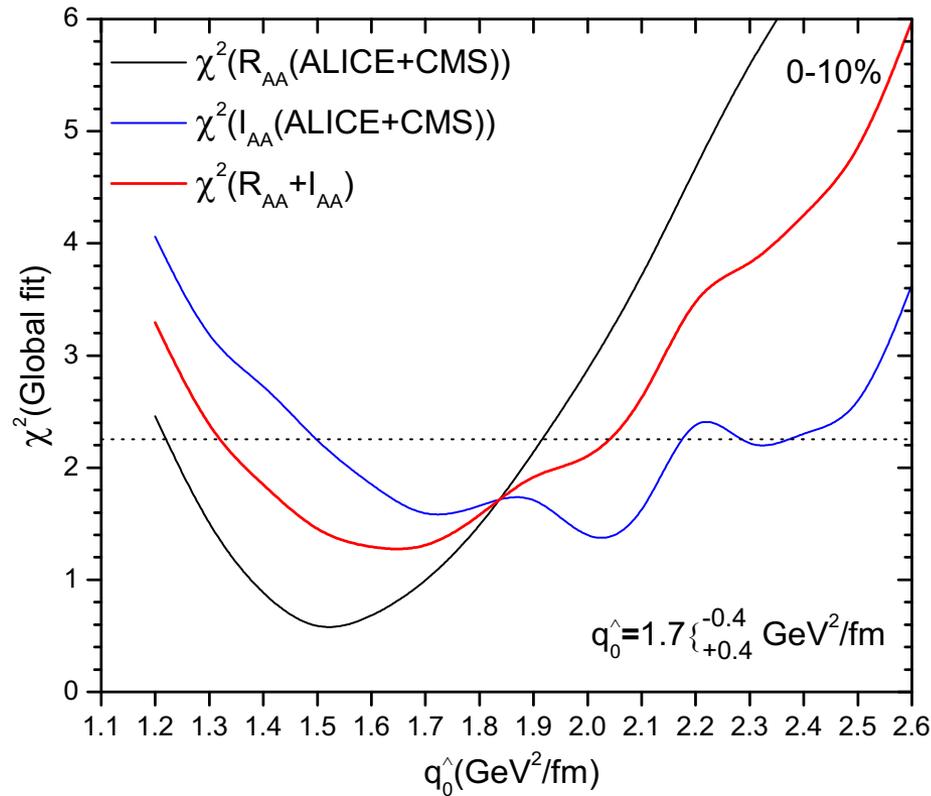
- Dihadron \hat{q}_0 is consistent with single hadron.
- Non central collisions results is consistent with central collisions.

Data from [[PHENIX Collaboration], Phys. Rev. C 87, no. 3,034911(2013), Phys. Rev. Lett. 101, 232301(2008)], and [[STAR], Phys. Rev. Lett. 97, 162301(2006)]

Extracting \hat{q}_0 in central Pb+Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV



Extracting \hat{q}_0 in central Pb+Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV

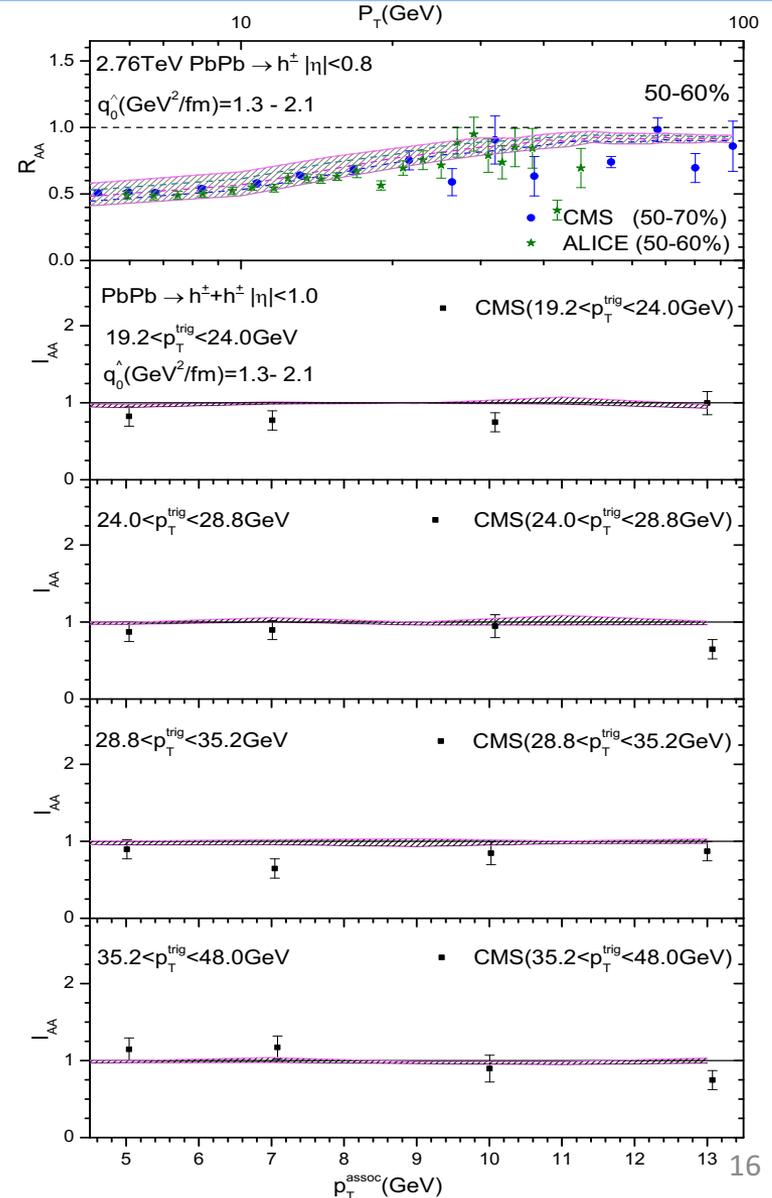


$$\hat{q}_0 = 1.7 \left\{ \begin{matrix} -0.4 \\ +0.4 \end{matrix} \right. (\text{GeV}^2/\text{fm})$$

Data from [[ALICE Collaboration], Phys. Lett. B 720 52-62(2013), Phys. Lett. B 763 238-250(2016)], [CMS Collaboration], Nuclear. Phys. A 904-905 451c-454c(2013), and [Eur. Phys. J. C 72, 1945(2012)]

Extracting \hat{q}_0 in non-central collisions at $\sqrt{s_{NN}}=2.76$ TeV

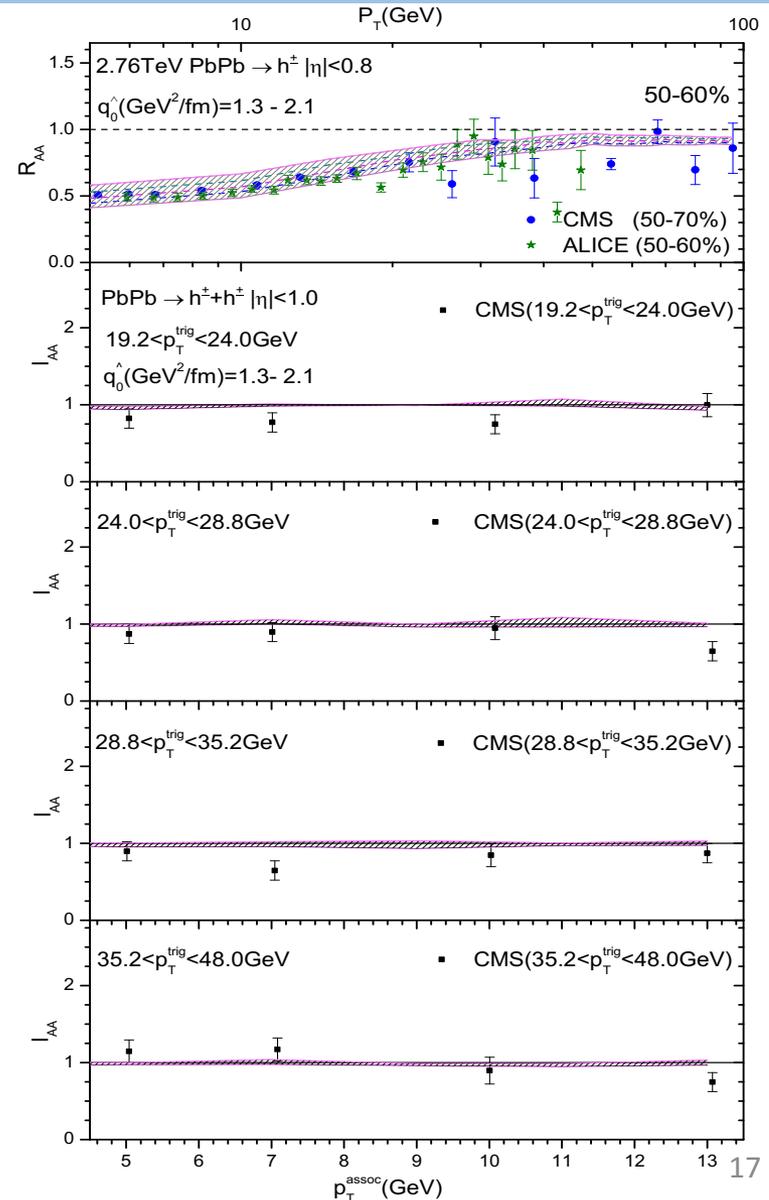
- R_{AA} and I_{AA} in the 50 - 60 % centrality with the \hat{q}_0 values extracted from central Pb+Pb collisions as compared to experimental data.



Extracting \hat{q}_0 in non-central collisions at $\sqrt{s_{NN}}=2.76$ TeV

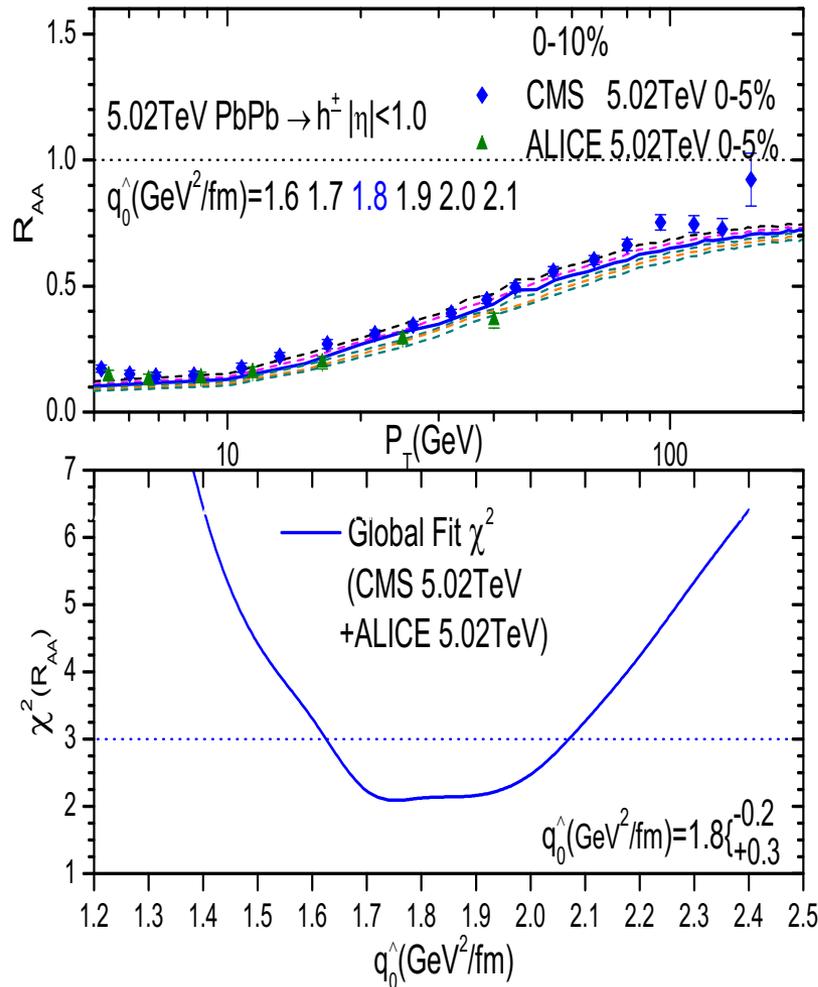
- R_{AA} and I_{AA} in the 50 - 60 % centrality with the \hat{q}_0 values extracted from central Pb+Pb collisions as compared to experimental data.

- \hat{q}_0 extracted from di-hadron suppression is consistent with single hadron.
- The single hadron R_{AA} and dihadron I_{AA} with the \hat{q}_0 extracted from central A+A collisions can give a good description of experimental data for the non-central collisions.



Extracting \hat{q}_0 via single hadron R_{AA} at $\sqrt{s_{NN}}=5.02$ TeV

- Extracting parameters \hat{q}_0 via single hadron suppression in central Pb+Pb collisions at $\sqrt{s_{NN}}=5.02$ TeV.

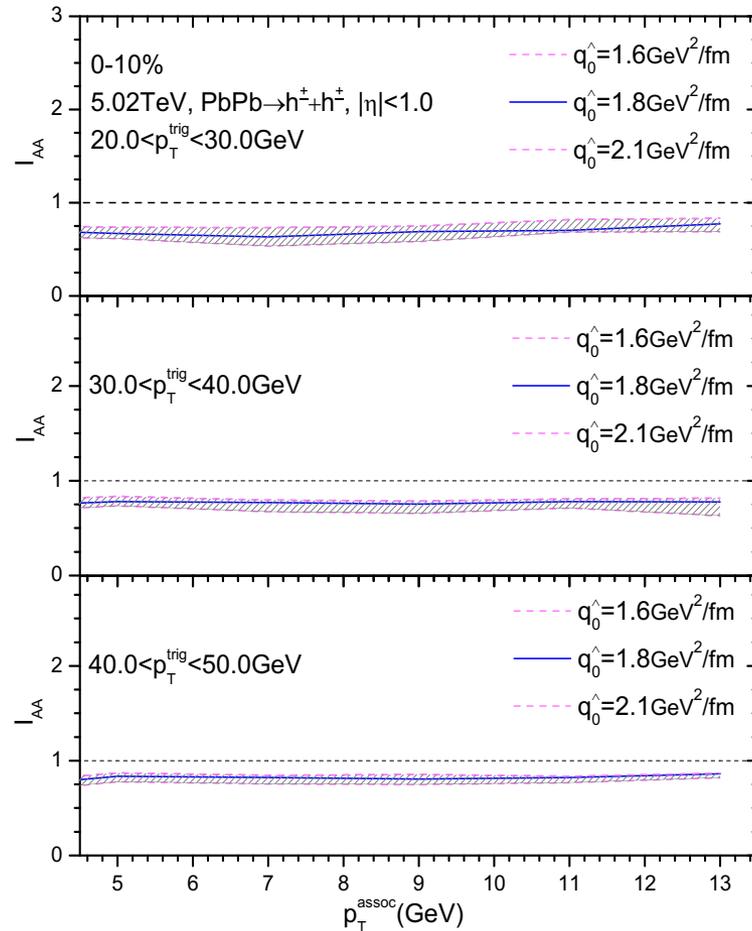
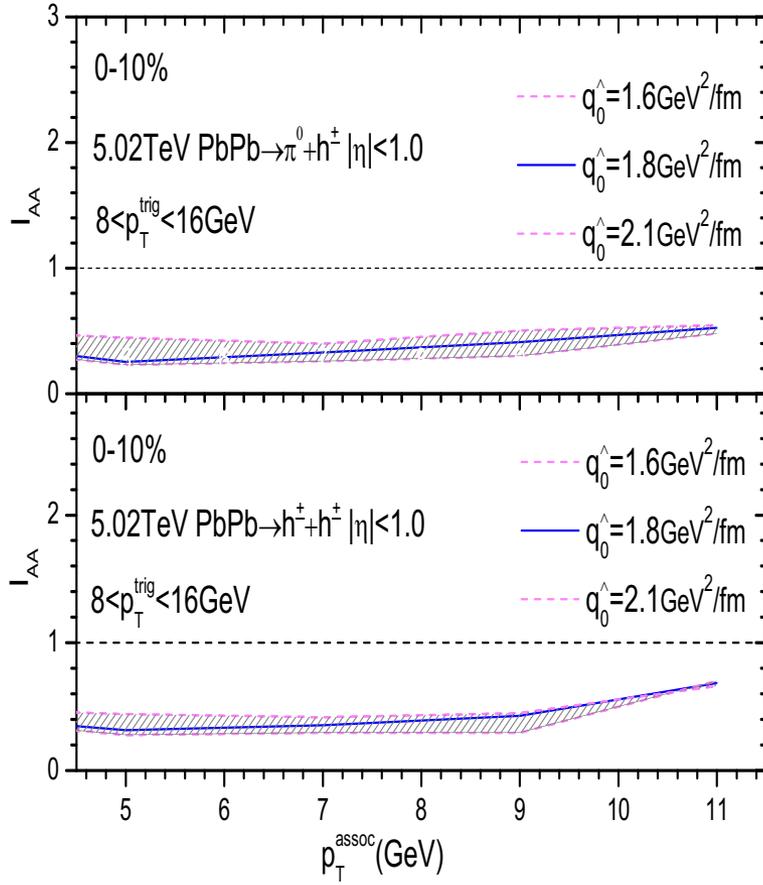


$$\hat{q}_0 = 1.8^{+0.3}_{-0.2} (\text{GeV}^2 / \text{fm})$$

[Data from [CMS collaboration], JHEP 1704 039 (2017), [ALICE collaboration], CERN EP 025 (2018), [arXiv:1802.09145v1[nucl-ex]](2018)]

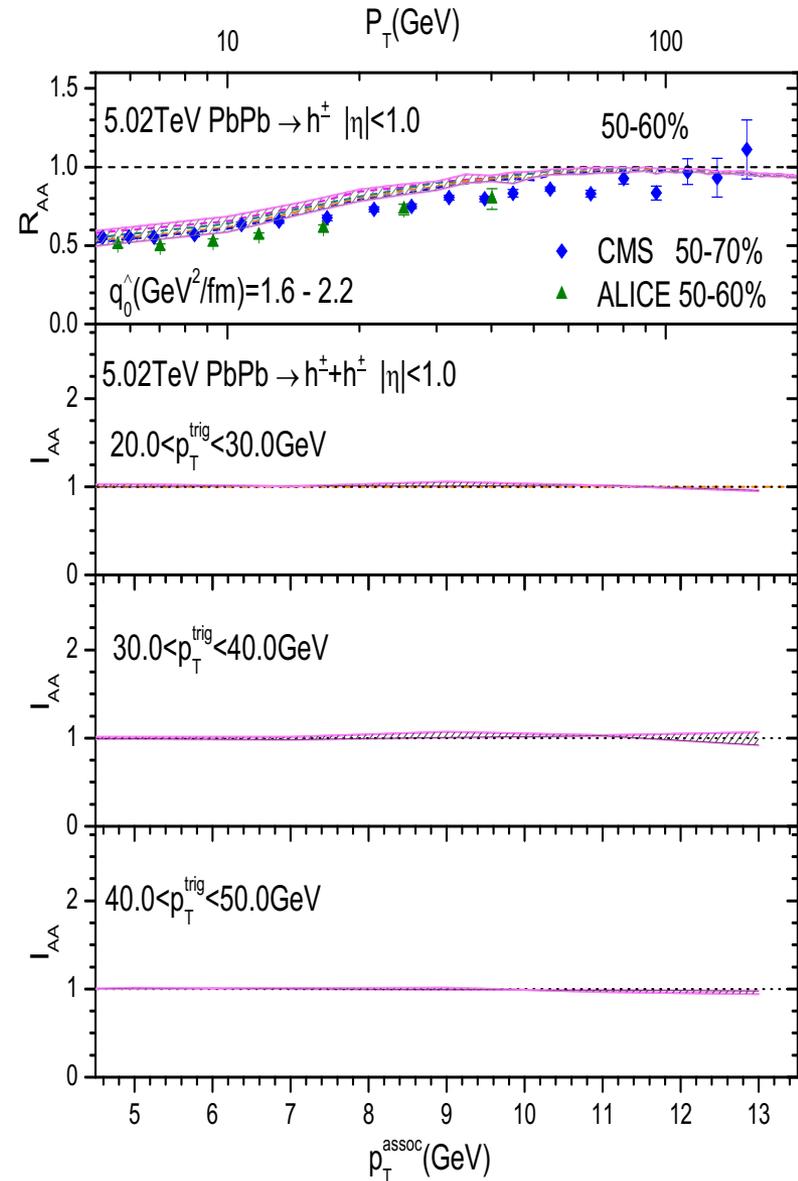
Prediction I_{AA} for central collisions at $\sqrt{s_{NN}}=5.02$ TeV

- The prediction of dihadron I_{AA} for central Pb+Pb collisions at $\sqrt{s_{NN}}=5.02$ TeV with the \hat{q}_0 values extracted from central Pb+Pb collisions.



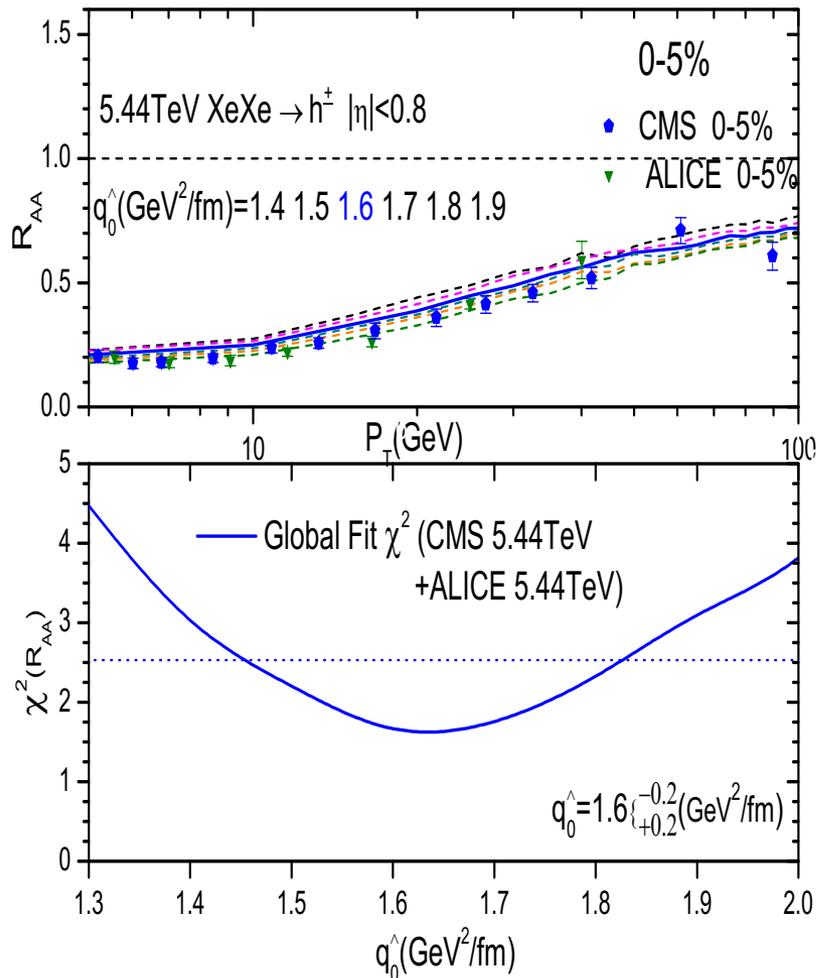
Prediction I_{AA} for non-central collisions at $\sqrt{s_{NN}}=5.02$ TeV

- The prediction of dihadron I_{AA} for Pb+Pb collisions in 50-60% centrality at $\sqrt{s_{NN}}=5.02$ TeV with the \hat{q}_0 values extracted from central Pb+Pb collisions.



Extracting \hat{q}_0 via single hadron R_{AA} at $\sqrt{s_{NN}}=5.44$ TeV

- Extracting parameters \hat{q}_0 via single hadron suppression in central Xe+Xe collisions at $\sqrt{s_{NN}}=5.44$ TeV.

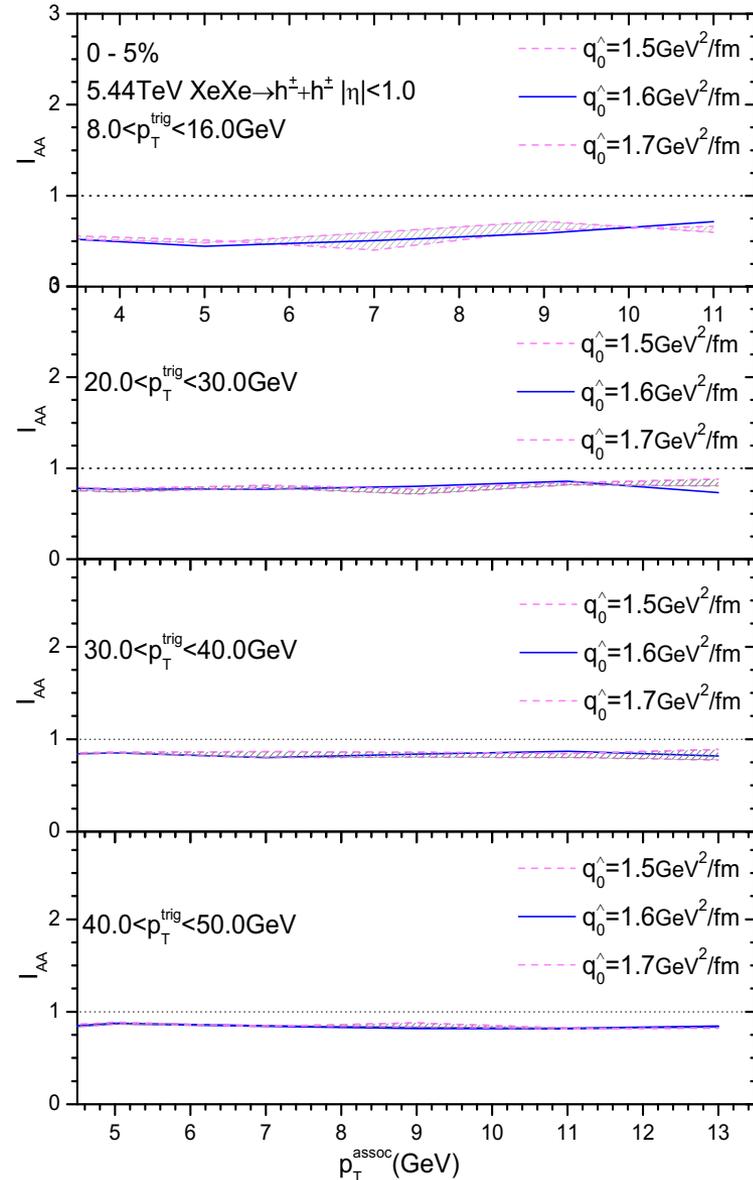


$$\hat{q}_0 = 1.6^{+0.2}_{-0.2} \text{ (GeV}^2/\text{fm)}$$

[Data from [ALICE collaboration], CERN EP 112 (2018), [arXiv:1805.04399v1[nucl-ex]](2018)]

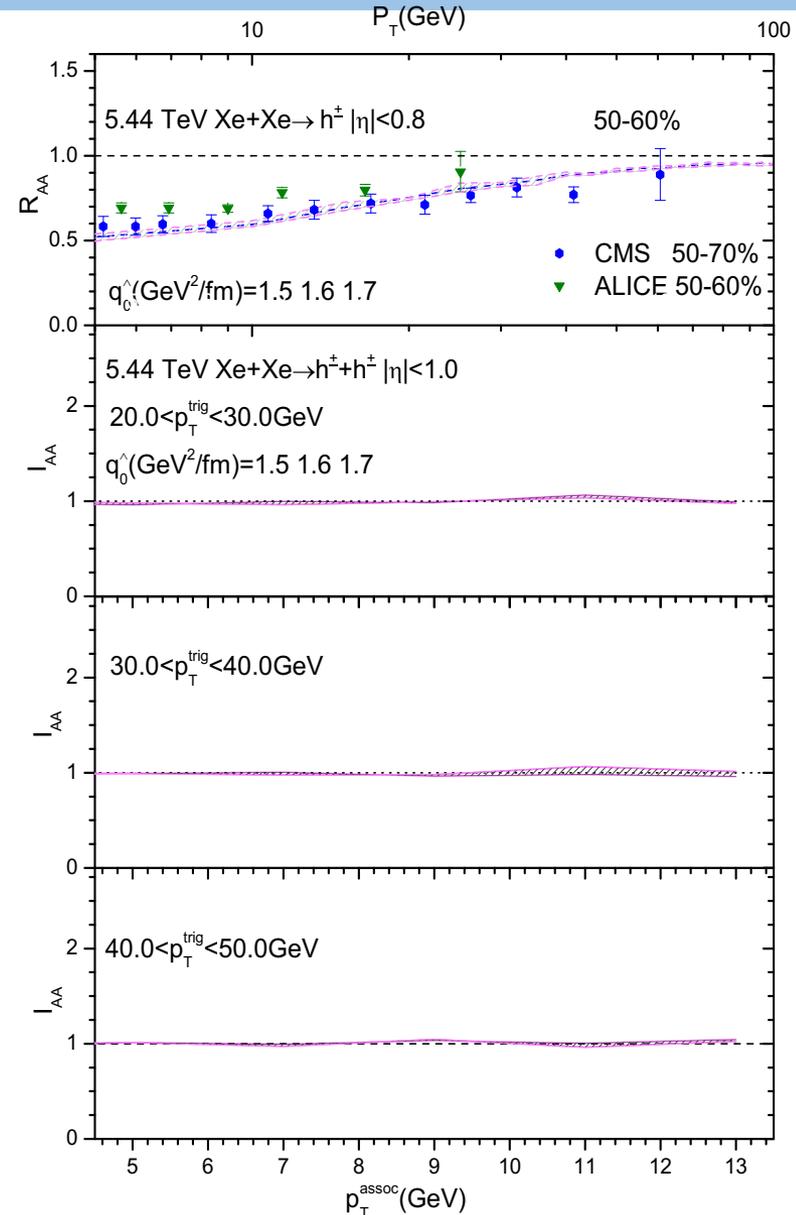
Prediction I_{AA} for central Xe+Xe collisions at $\sqrt{s_{NN}}=5.44$ TeV

- The prediction of dihadron I_{AA} for central Xe+Xe collisions at $\sqrt{s_{NN}}=5.44$ TeV with the \hat{q}_0 values extracted from central Xe+Xe collisions.



Prediction I_{AA} for non-central Xe+Xe collisions

- The prediction of dihadron I_{AA} for Xe+Xe collisions in 50-60% centrality at $\sqrt{s_{NN}}=5.44$ TeV with the \hat{q}_0 values extracted from central Xe+Xe collisions.



Summary and Outlook

- Large p_T hadrons are studied in a NLO pQCD parton model in heavy-ion collisions with mFFs due to jet quenching.
- We obtain $\hat{q}_0 \approx 1.1^{+0.4}_{-0.2} \text{ GeV}^2/\text{fm}$ for central Au+Au collisions at 200GeV, and $\hat{q}_0 \approx 1.7^{+0.4}_{-0.4} \text{ GeV}^2/\text{fm}$ for central Pb+Pb collisions at 2.76 TeV.
- We predicted the dihadron suppression factors for central and non-central Pb+Pb collisions at 5.02TeV with the \hat{q}_0 extracted via single hadron suppression.
- We also used the latest data of Xe+Xe collisions to extract \hat{q}_0 , and predicted dihadron suppression factors for central and non-central Xe+Xe collisions.
- In the future, to study jet quenching we want to do jet-hadron photon-hadron correlations to extract \hat{q}_0 .

THANK YOU!