

3-BODY PHYSICS IN (IN)FINITE VOLUME

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I. Introduction

2-BODY EXAMPLE: $\sigma(500)$

- parameters (existence) debated for decades
- **dispersive techniques:** most precise resonance parameters

→ [Review Pelaez \(2015\)](#)

[Colangelo et al. \(2001\)](#) [Garcia-Martin et al. \(2011\)](#) [Moussallam \(2011\)](#) ..



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$$T_{IAM} = \frac{(T_2)^2}{T_2 - T_4 + A_m(s)}$$

Truong (1988)... Hanhart et al. (2008)

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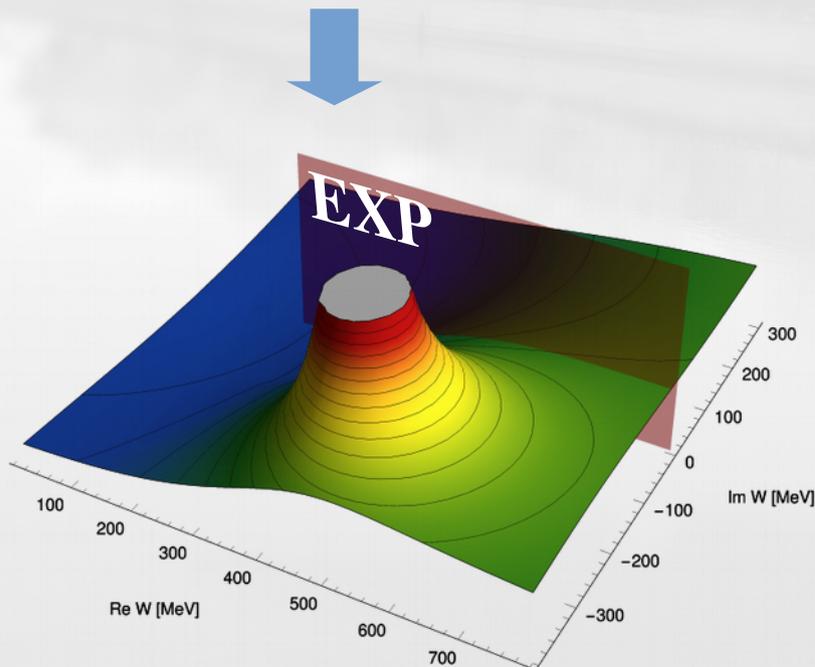
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→ *Scalar-isoscalar $\pi\pi$ interaction generates σ -resonance*

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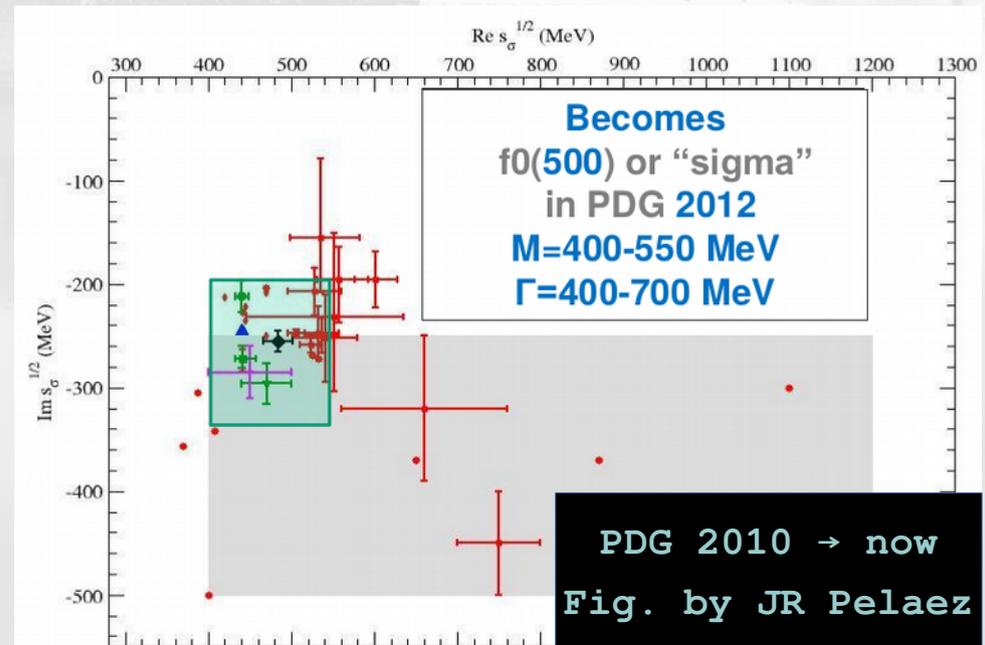
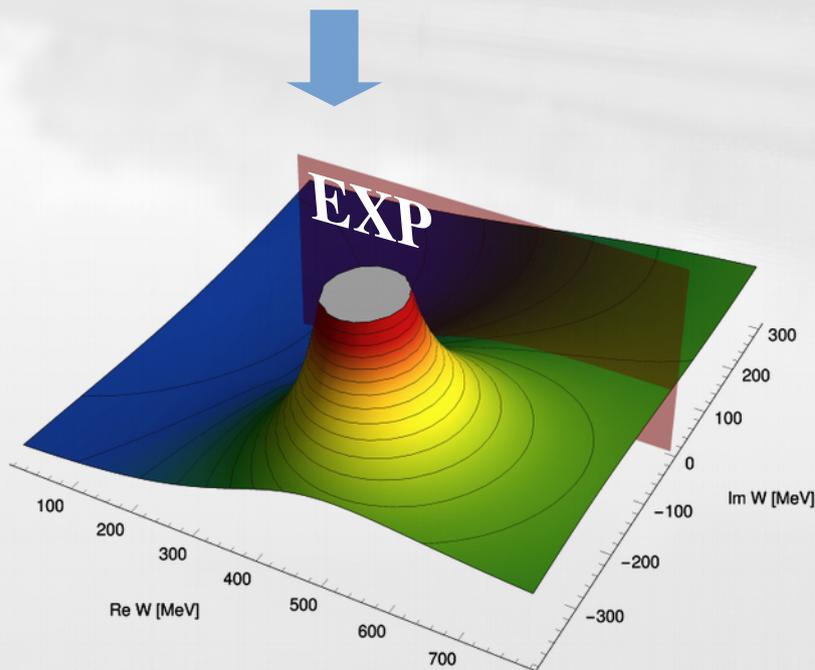
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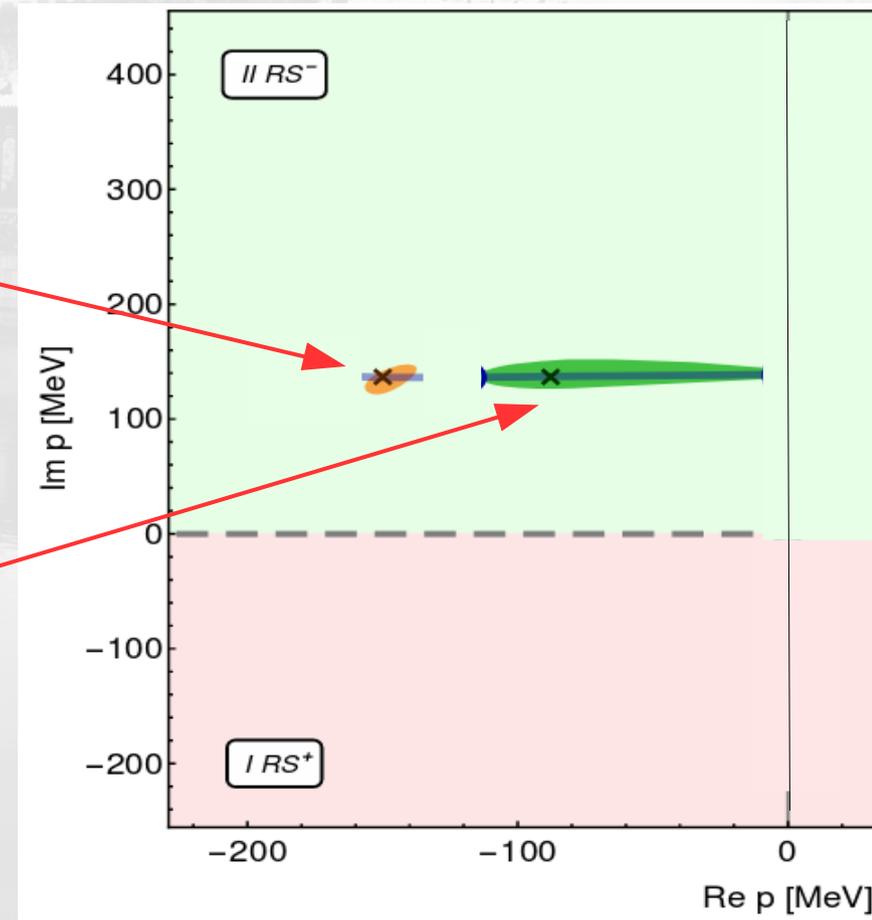
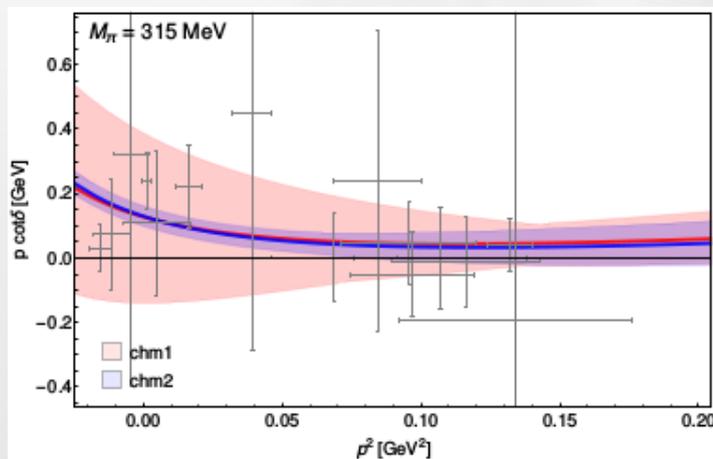
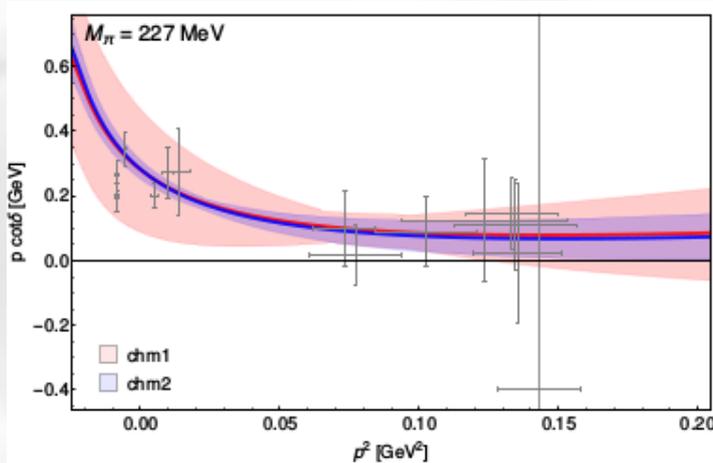
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Briceno et al. (2016) / ETMC (2017) / Fu, ... (2018)

Example: $N_f=2$, $m_\pi=227/315$ MeV

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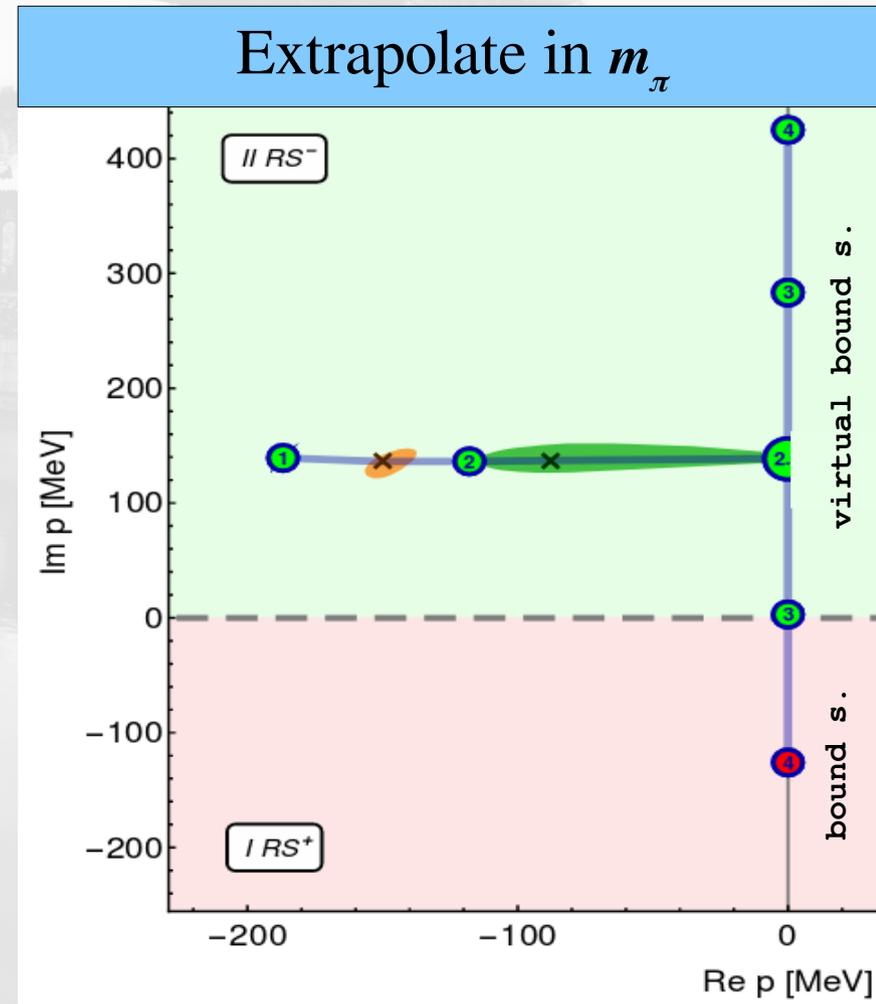
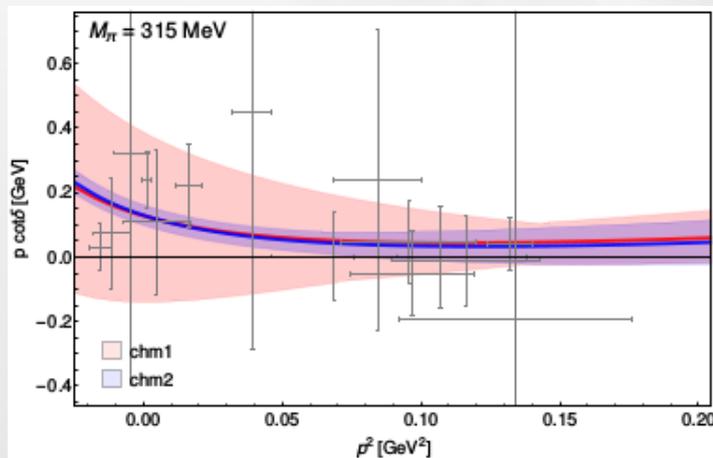
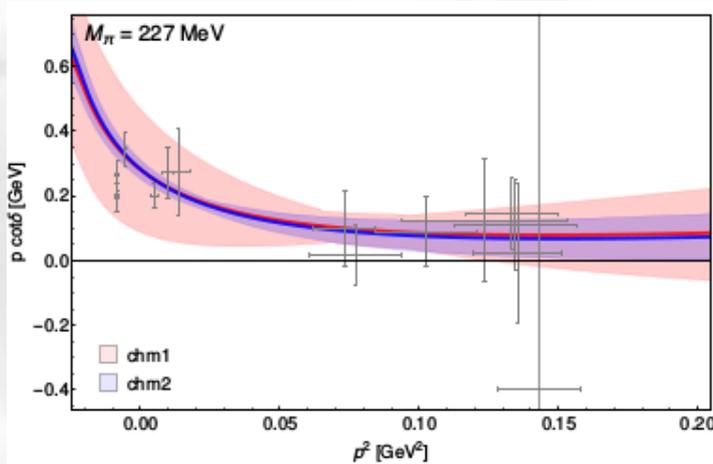
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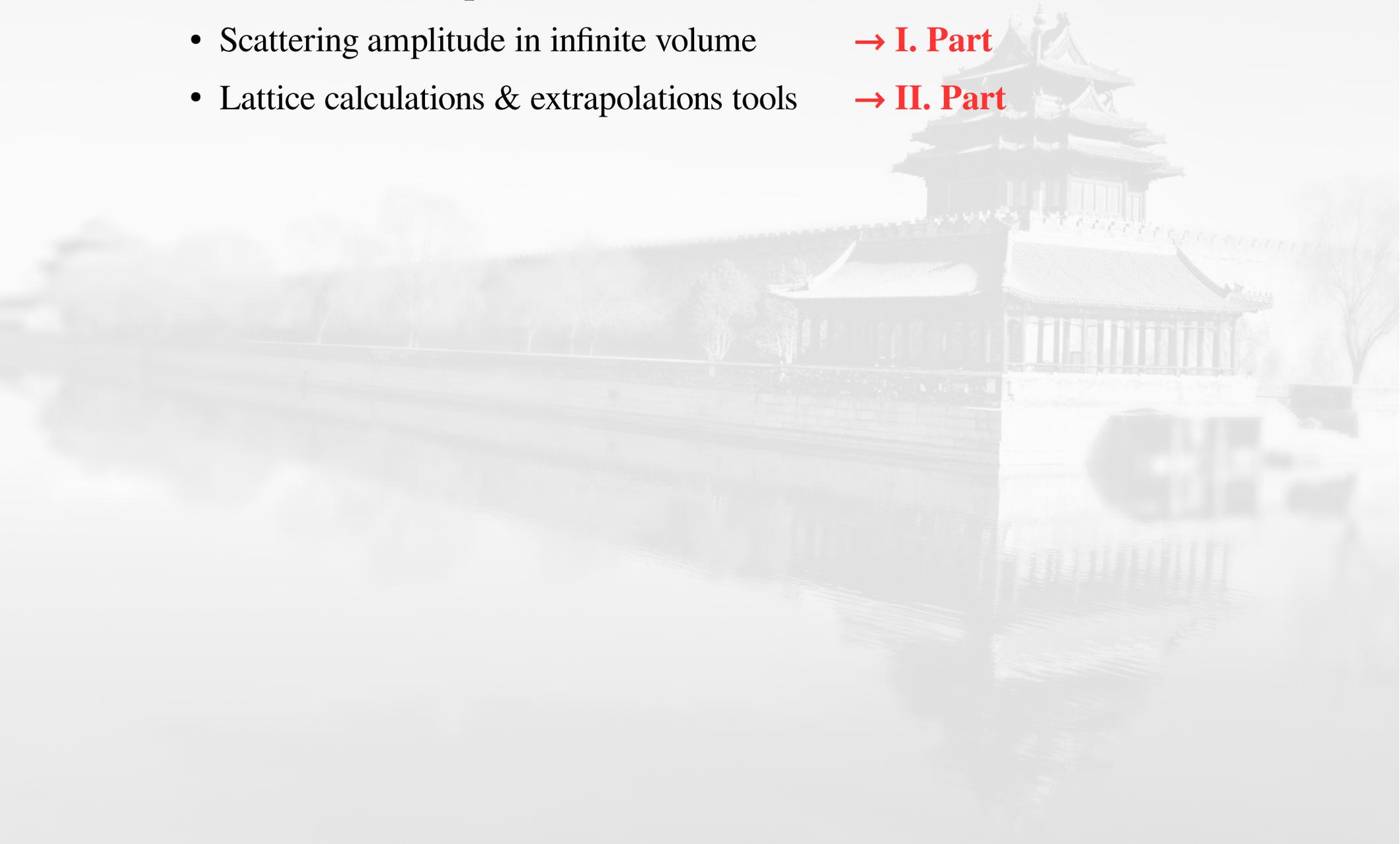
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- Scattering amplitude in infinite volume
- Lattice calculations & extrapolations tools

→ I. Part

→ II. Part



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 - gluonic degrees of freedom
 - Cannot decay in 2 but in 3 pions, **BUT...**

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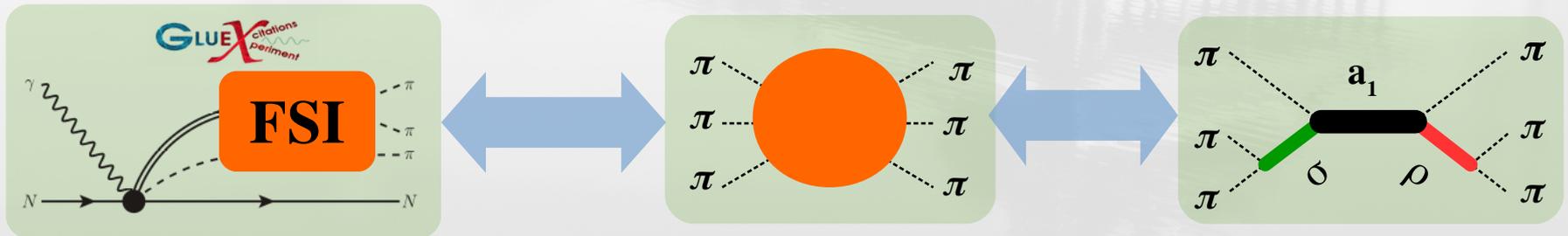
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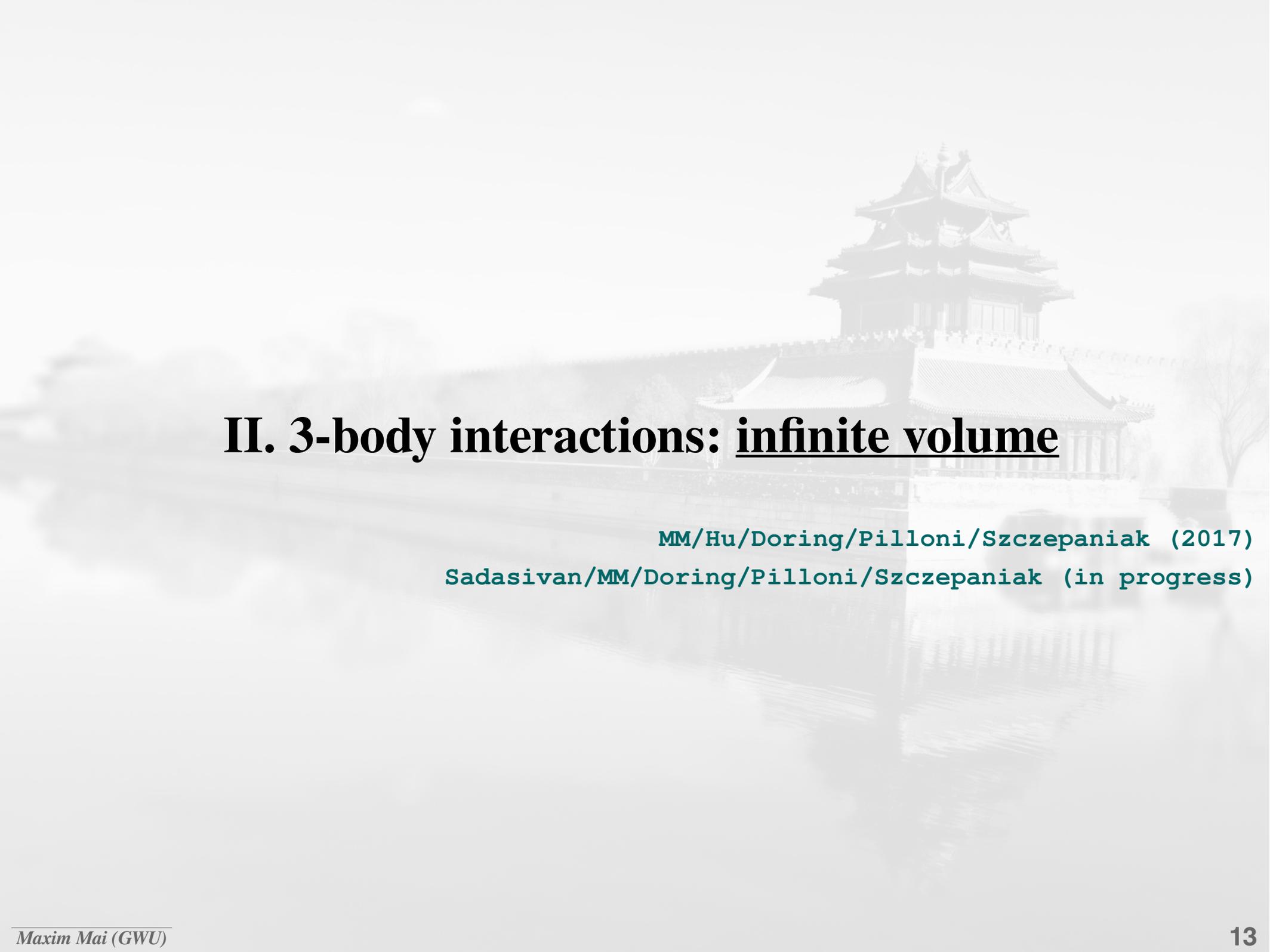
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→ Need to pin down 3-hadron interaction



II. 3-body interactions: infinite volume

MM/Hu/Doring/Pilloni/Szczepaniak (2017)

Sadasivan/MM/Doring/Pilloni/Szczepaniak (in progress)

SCATTERING MATRIX (T)

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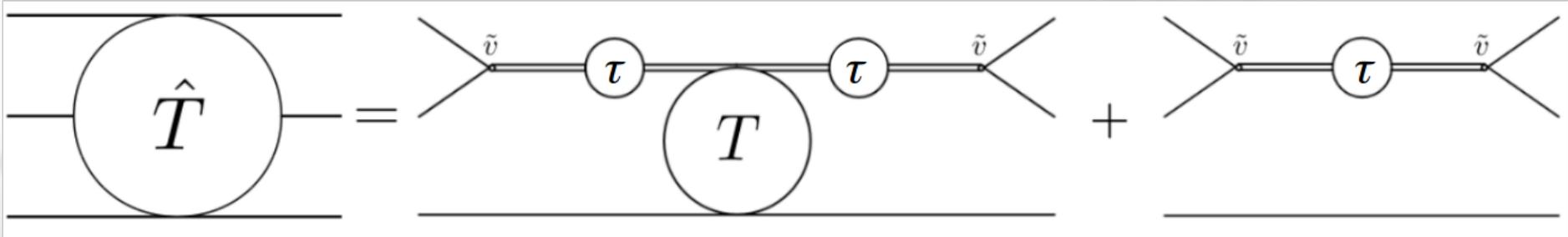
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* “isobars” = $\tau (M_{inv})$ for definite QN & correct right-hand-singularities

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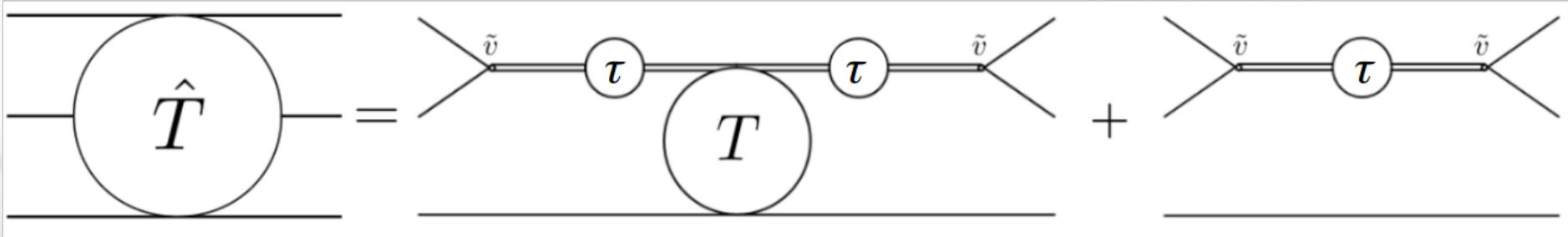
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3 unknown functions & 8 kinematic variables

SCATTERING MATRIX (T)

- **3-body unitarity**

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

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$$\text{Disc } \tau(\sigma(k)) = \frac{-i}{64\pi^2 K_{\text{cm}}} \int d^3 \bar{\mathbf{K}} \frac{\delta(|\bar{\mathbf{K}}| - K_{\text{cm}})}{\sqrt{(\bar{\mathbf{K}})^2 + m^2}} v^2$$

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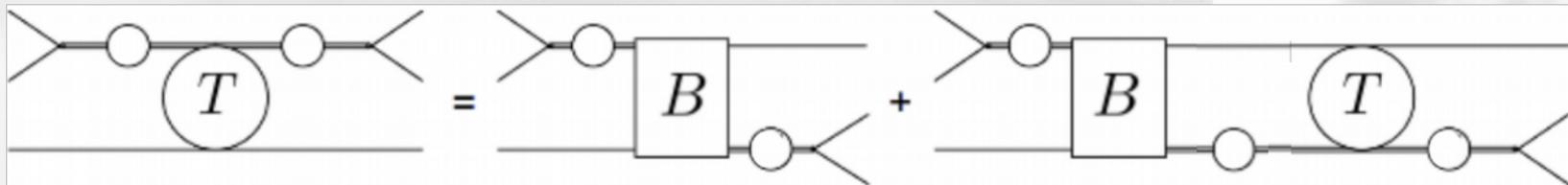
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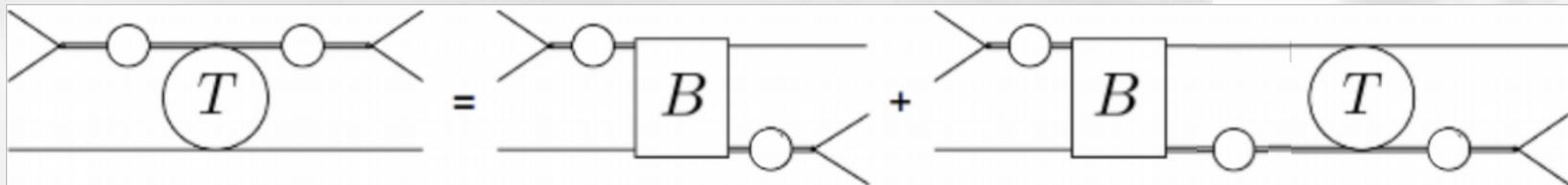
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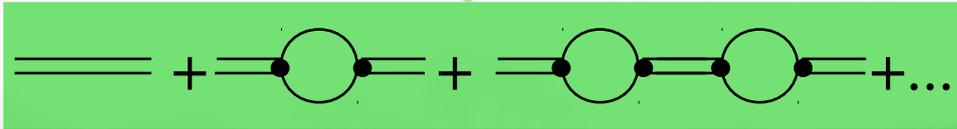
→ gives access to $\text{Disc}[B]$

$$\text{Disc } B(u) = 2\pi i \frac{\delta(E_Q - \sqrt{m^2 + \mathbf{Q}^2})}{2\sqrt{m^2 + \mathbf{Q}^2}} v^2$$

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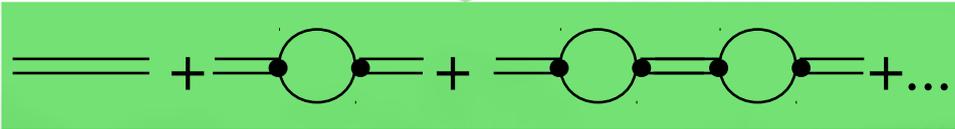
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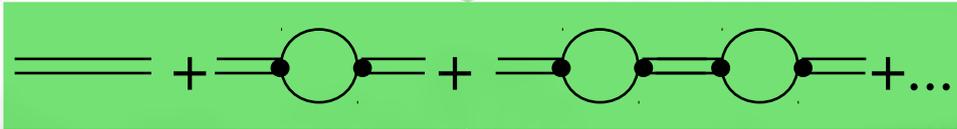


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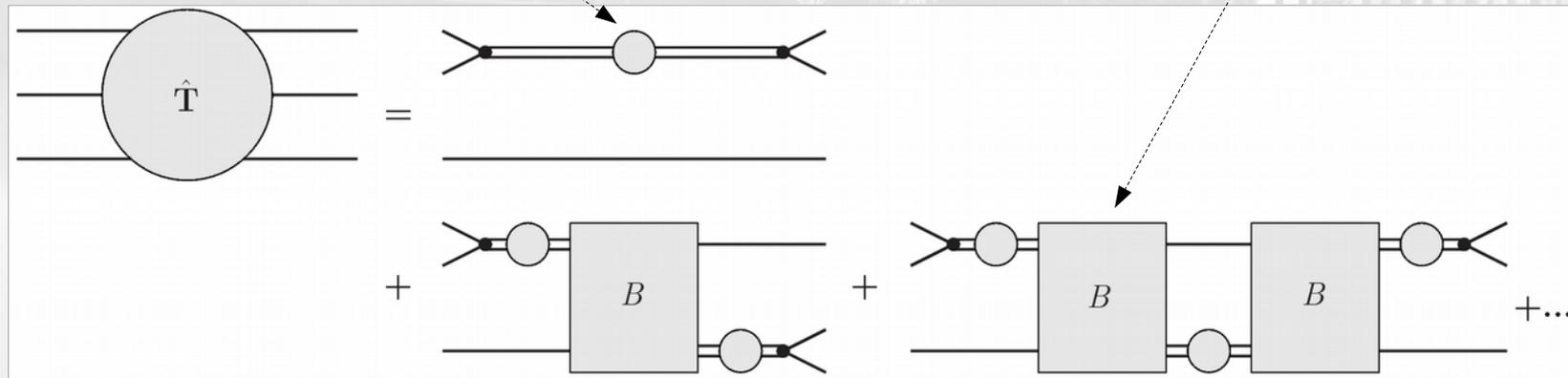


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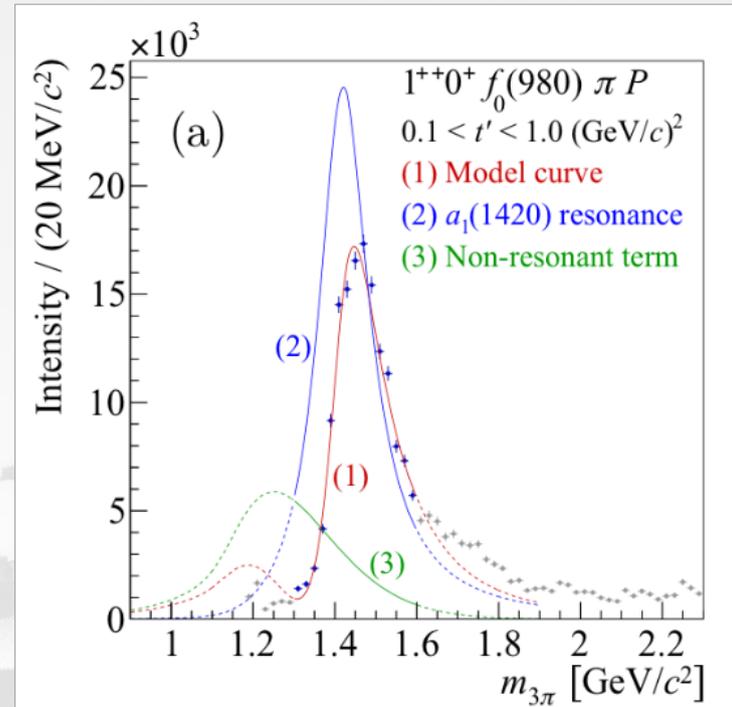
→ one-”pion” exchange is required



- **Manifestly unitary & covariant 3d integral equation**
- An infinite series of isobars (two-body subamplitudes) interacting with spectator:
- Unknown parameter: C, v , *parameters of the isobar (subtraction constants)* ← **II. Part**

Interesting application: $a_1(1420)$

- Lineshape from COMPASS @ CERN
- in $f_0(980)\pi$ final state



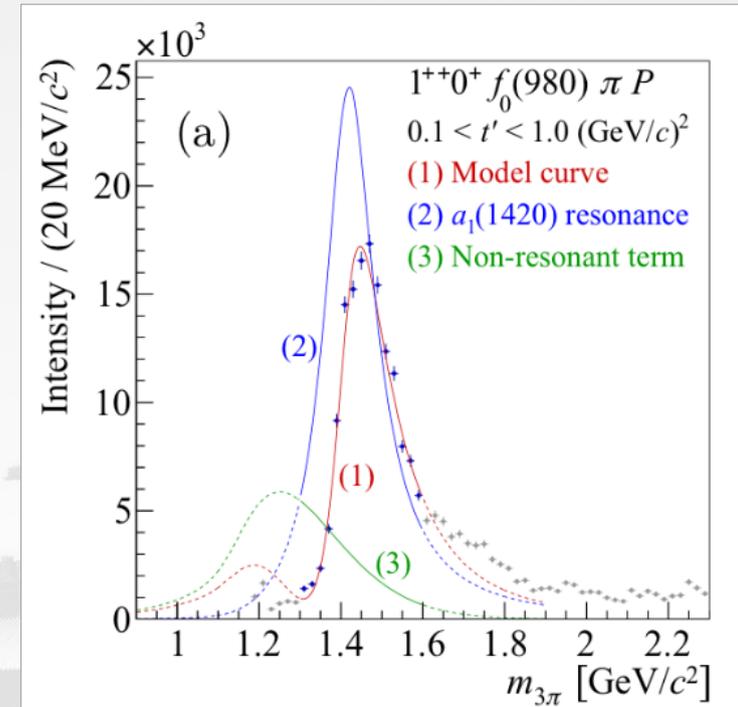
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Mikhasenko/Ketzer/Sarantsev (2015) Aceti/Dai/Oset (2016)



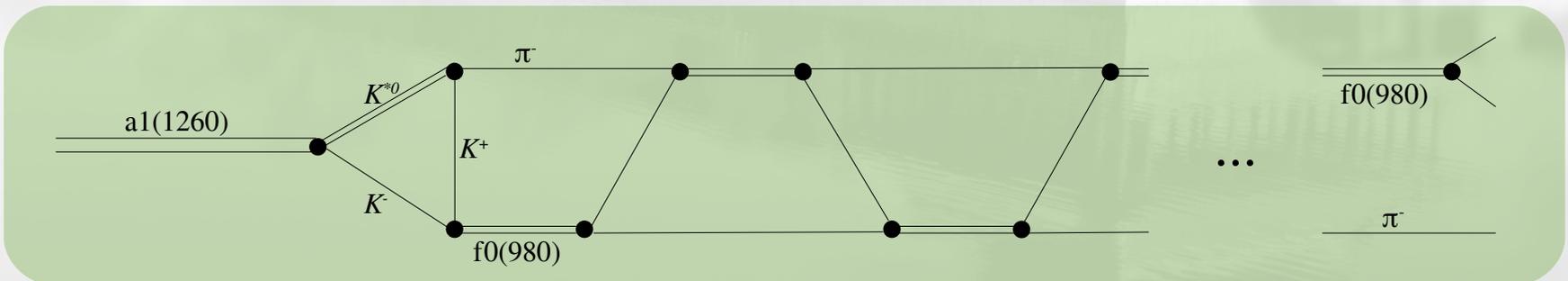
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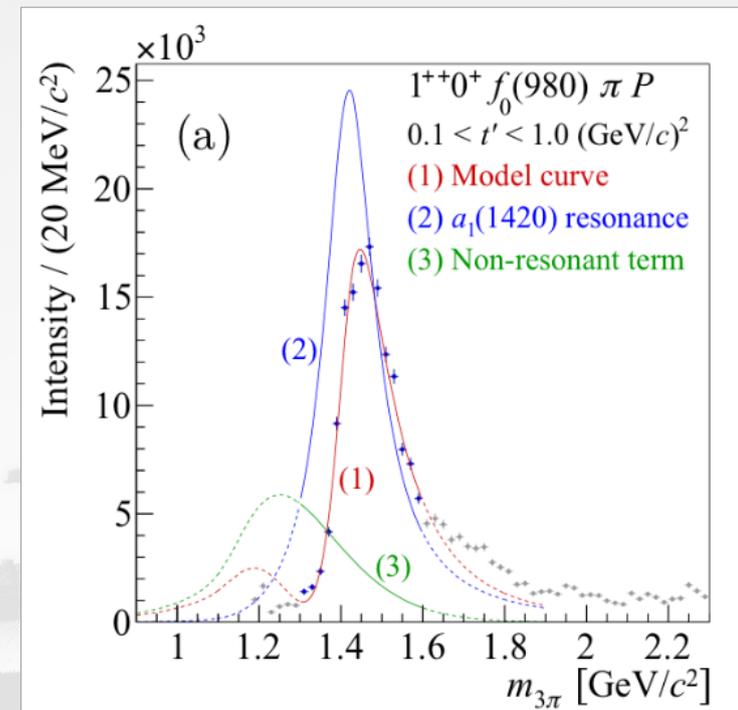
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- **Q:** Does such a feature exist in full 3b-unitary FSI?
→ Check with unitary isobar approach



Sadasivan/MM/... (in progress)



III. 3-body interactions: finite volume

MM/Doring (2017)

MM/Doring (2018)

→ last week

LATTICE QCD

Ab-initio numerical calculations of QCD Greens functions

- Countless **new insights** into properties of hadrons

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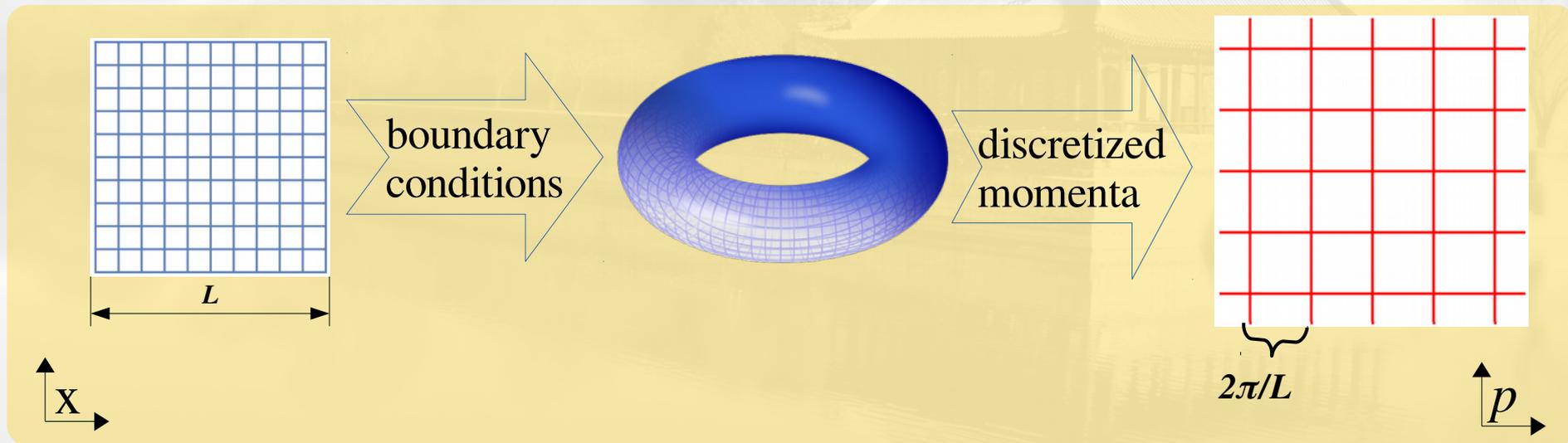
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- Technical necessity: **finite volume**



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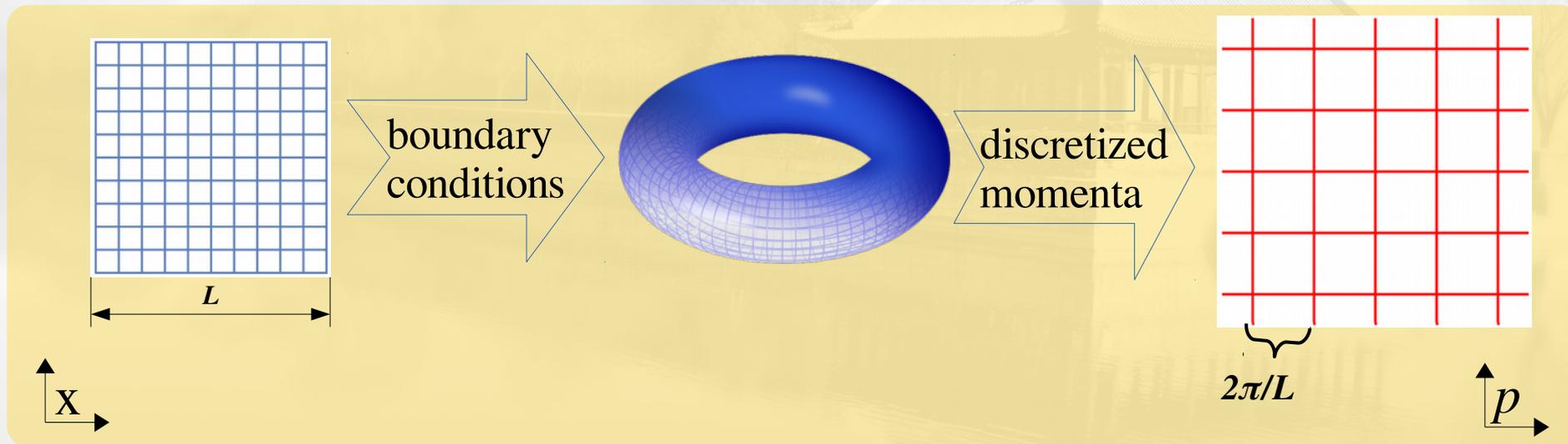
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- Technical necessity: **finite volume**



→ momenta & spectra are discretized

→ extrapolation to infinite volume \leftrightarrow **Quantization Condition**

QUANTIZATION CONDITION

2-body case

- well understood
- multi-channels, spin, ...

Lüscher (1986)

Gottlieb, Rummukainen, Feng, Li, Liu, Doring, Briceno, Rusetsky, Bernard..

3-body case

- important theoretical developments

Sharpe, Hansen, Briceno, Rusetsky, Polejaeva, Davoudi, Guo, MM, Doring..

- pilot numerical investigation

Pang/Hammer/Rusetsky/Wu (2017) MM/Doring (2017) Hansen/Briceno/Sharpe (2018)

- Finite volume spectrum of $(\pi^+\pi^+)$ and $(\pi^+\pi^+\pi^+)$

→ comparison with Lattice QCD results

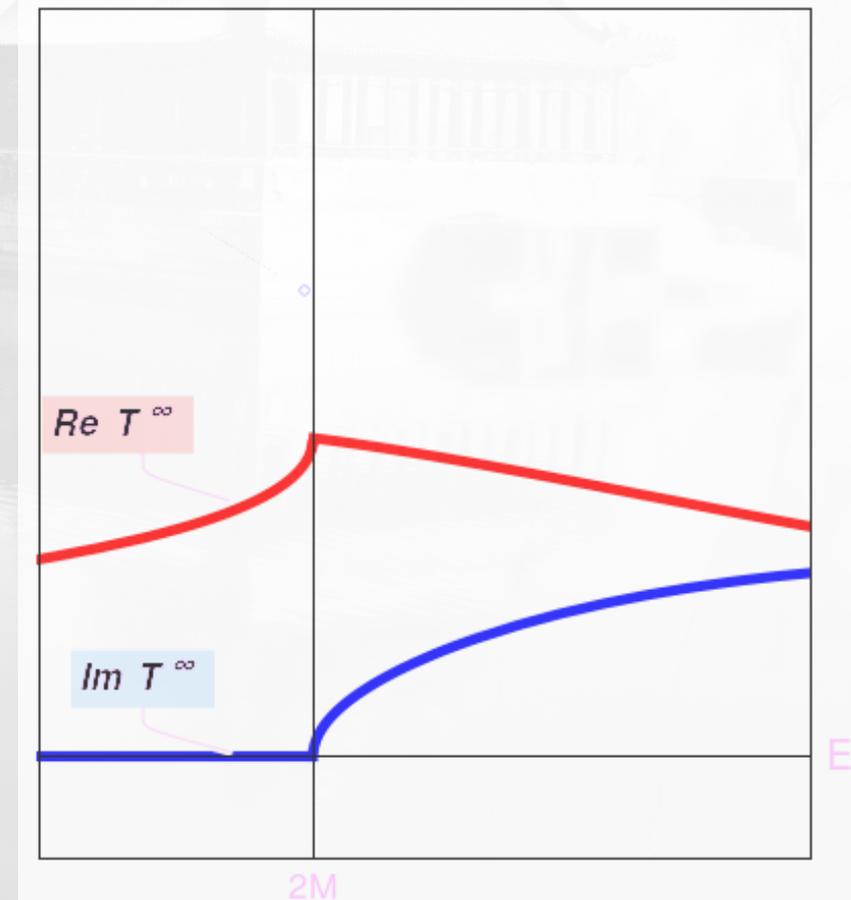
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EXAPMLE: 2-BODY QUANTIZATION CONDITION

One way of thinking:

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$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$



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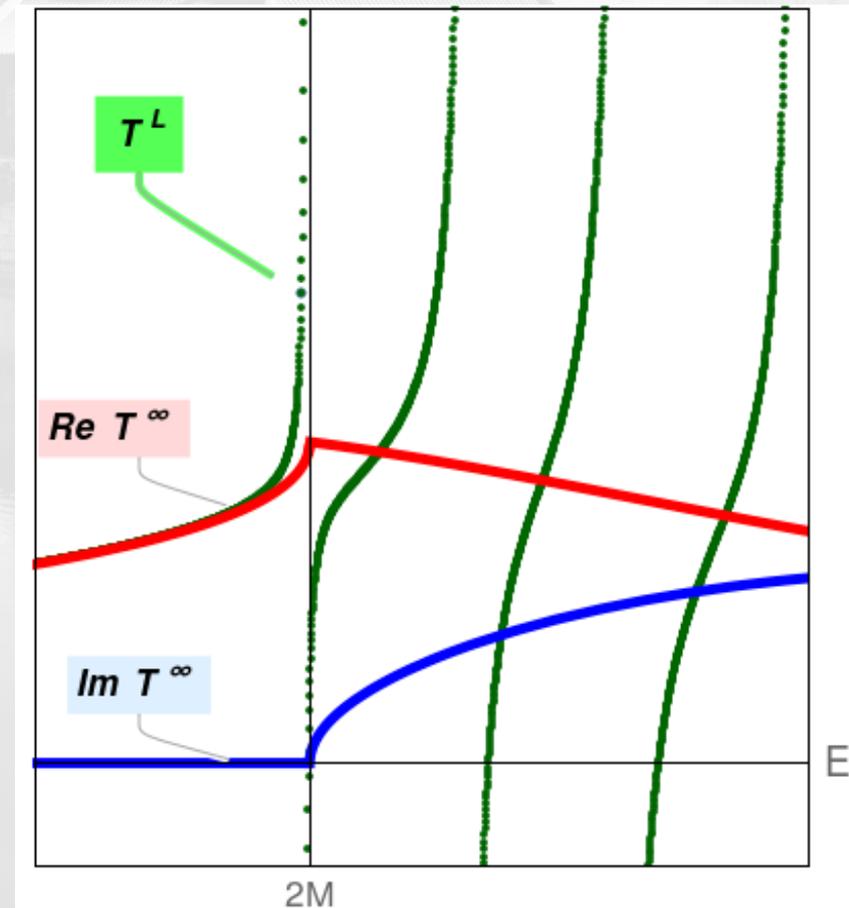
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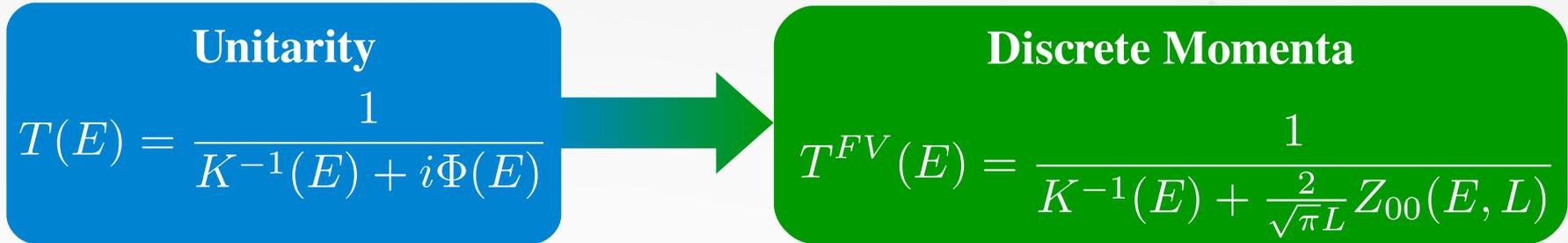
Discrete Momenta

$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi}L} Z_{00}(E, L)}$$



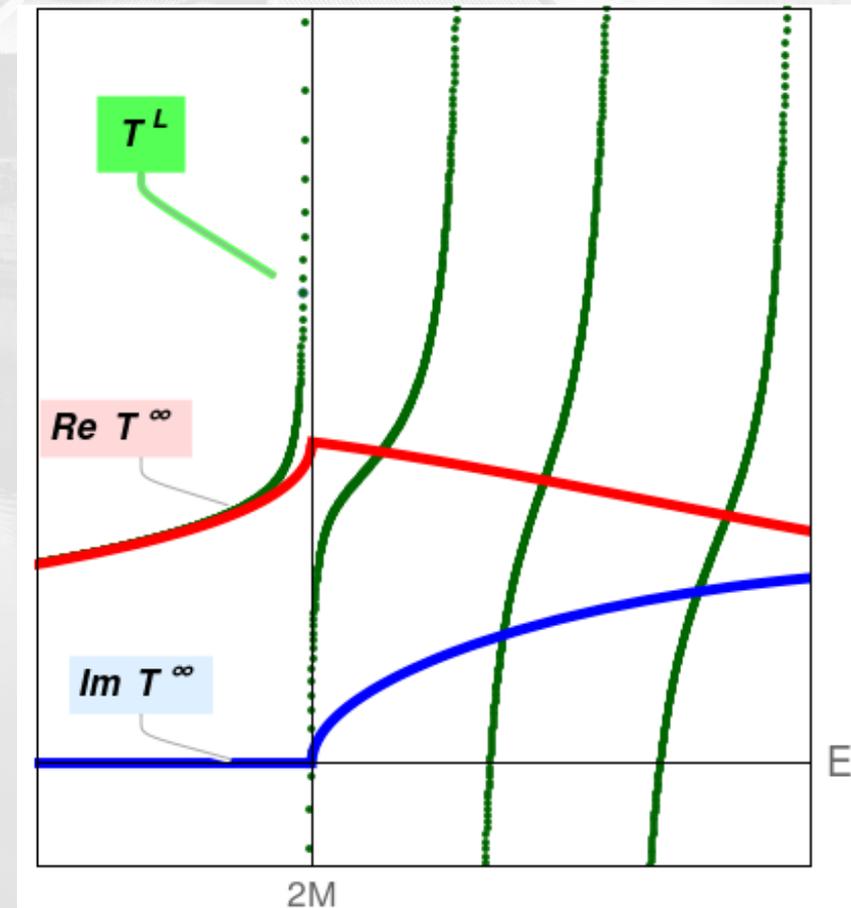
EXAPMLE: 2-BODY QUANTIZATION CONDITION

One way of thinking:



- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ is singular
- LSZ formalism relates Greens fct. to S-matrix
(pole positions) ↔ (energy eigenvalues)

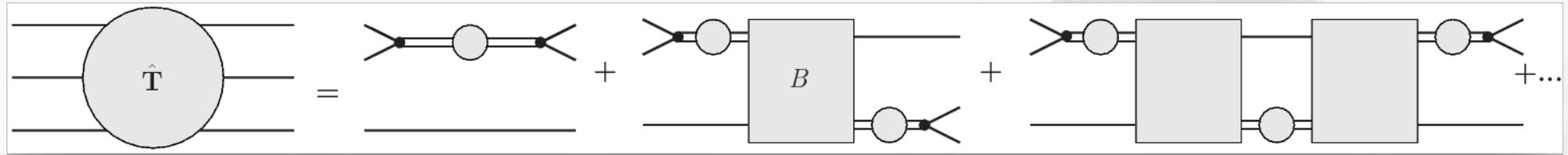
$$K^{-1}(E^*) = -\frac{2}{\sqrt{\pi}L} Z_{00}(E^*, L)$$



3-BODY QUANTIZATION CONDITION

Q: Can we repeat this for 3-body case?

1) Unitary 3-body amplitude

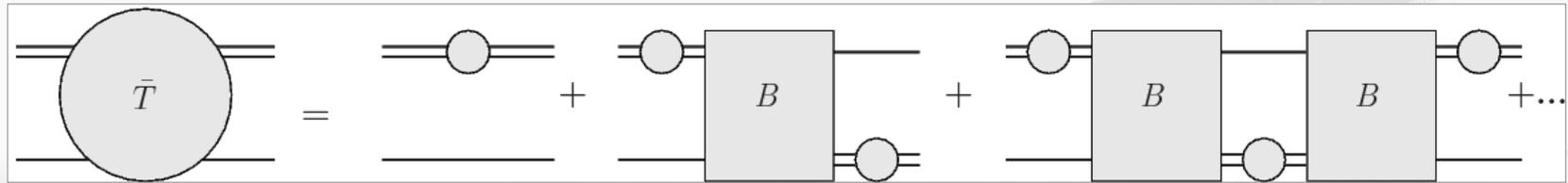


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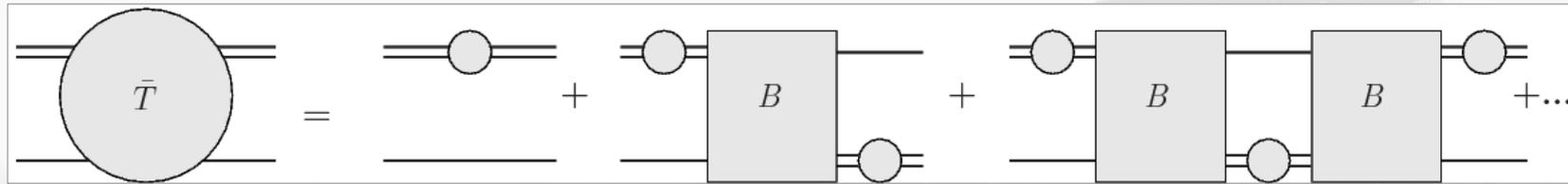
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integral equation \rightarrow matrix equation:

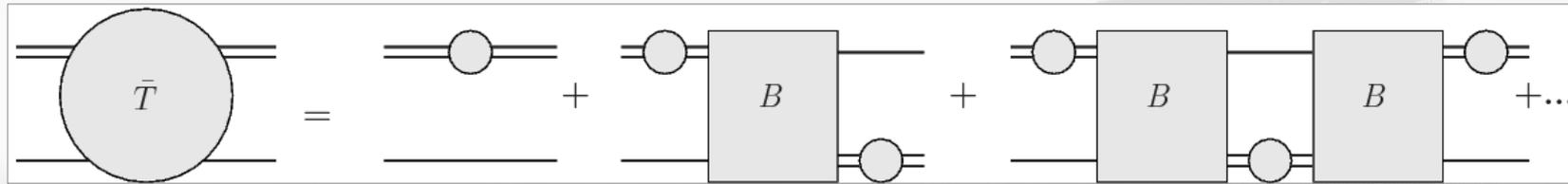
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4) Project to irreducible representations of the cubic group: Γ

\rightarrow Condition for poles of \bar{T} :

W – total energy

s/s' - shell index

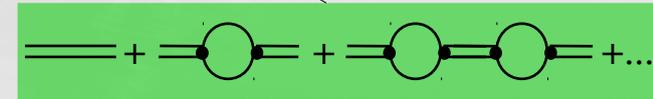
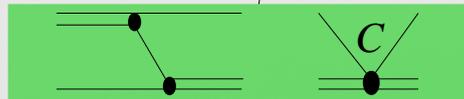
u/u' - basis index

ϑ – multiplicity

L – lattice volume

E_s – 1p. energy

$$\text{Det} \left(\mathbf{B}_{uu'}^{\Gamma ss'}(W^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s (W^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$



REALISTIC CASE STUDY

MM/Doring
arXiv:1807.04746



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Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

- Ground state levels available from NPLQCD

$L=2.5\text{fm}$ and $m_\pi=291/352/491/591\text{ MeV}$

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- *non-resonant* $\pi^+\pi^+$ subsystem in S -wave

→ one isobar required

Q: Can one describe this with an “isobar”?

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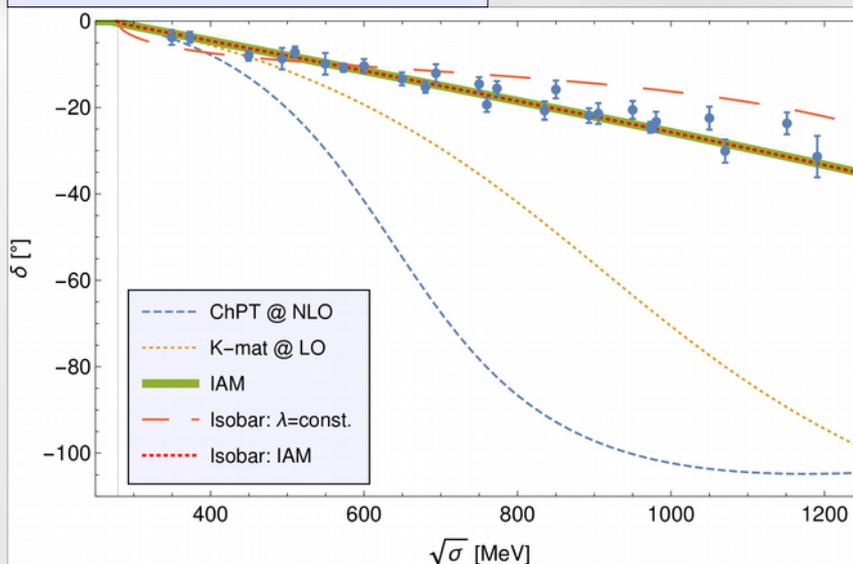
- *non-resonant* $\pi^+\pi^+$ subsystem in S -wave

→ one isobar required

Q: Can one describe this with an “isobar”?

A: YES! Choose $\mathbf{v} = \mathbf{v}(\sigma)$

2-body
phase shift ($I=2, L=0$)



REALISTIC CASE STUDY

Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

- Ground state levels available from NPLQCD

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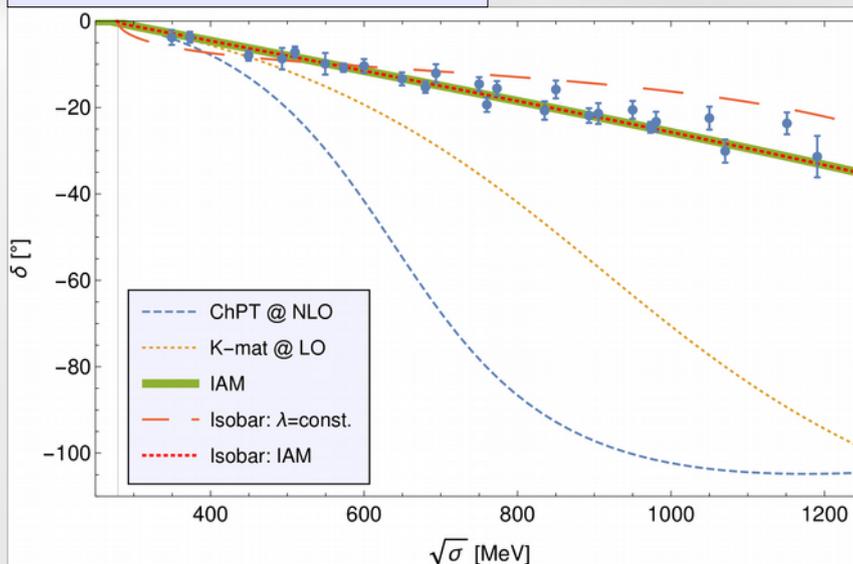
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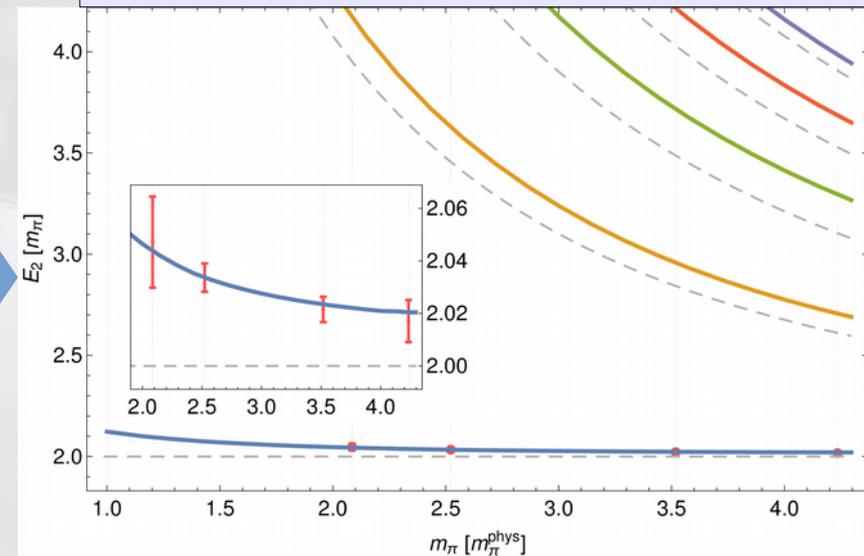
A: YES! Choose $\mathbf{v} = \mathbf{v}(\sigma)$

2-body
phase shift ($I=2, L=0$)



discretize

2-body finite vol. spectrum
PARAMETER FREE PREDICTION



REALISTIC CASE STUDY

MM/Doring
arXiv:1807.04746

Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

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$L=2.5\text{fm}$ and $m_\pi=291/352/491/591\text{ MeV}$

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$$\text{Det} \left(\mathbf{B}_{\mathbf{uu}'}^{\Gamma_{\mathbf{ss}'}}(\mathbf{W}^2) + \frac{2\mathbf{E}_s L^3}{\vartheta(\mathbf{s})} \tau_{\mathbf{s}}(\mathbf{W}^2)^{-1} \delta_{\mathbf{ss}'} \delta_{\mathbf{uu}'} \right) = 0$$

REALISTIC CASE STUDY

Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

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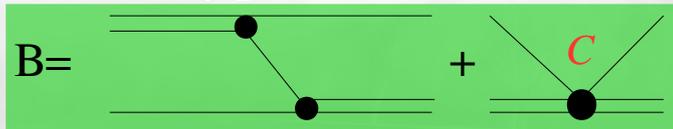
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- Remaining parameter:



~ 3-body force

Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

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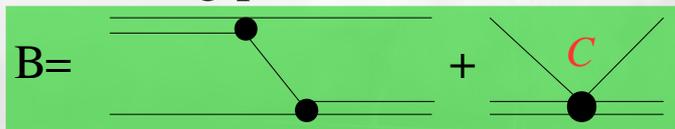
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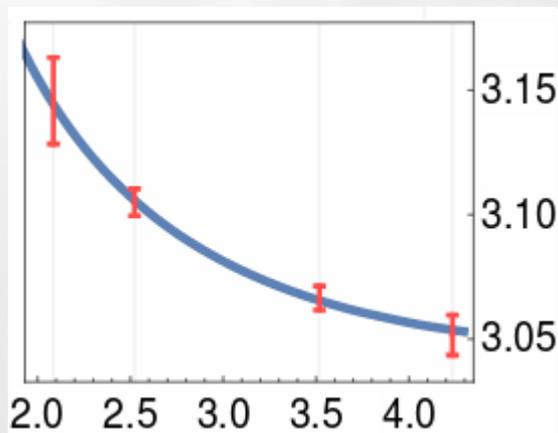
- Remaining parameter:



~ 3-body force

→ fit to the ground state levels by NPLQCD

→ $C=0.2 \pm 1.5 \cdot 10^{-10}$



REALISTIC CASE STUDY

Test bed for the Quantization condition: $\pi^+\pi^+\pi^+$

- Ground state levels available from NPLQCD

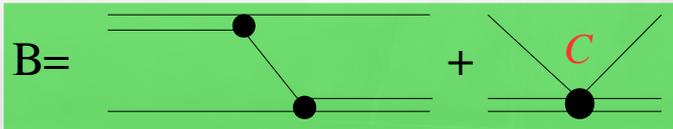
$L=2.5\text{fm}$ and $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- 3-body spectrum from the QC:

$$\text{Det} \left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\Gamma_{\mathbf{s}\mathbf{s}'}}(\mathbf{W}^2) + \frac{2\mathbf{E}_s \mathbf{L}^3}{\vartheta(\mathbf{s})} \tau_{\mathbf{s}}(\mathbf{W}^2)^{-1} \delta_{\mathbf{s}\mathbf{s}'} \delta_{\mathbf{u}\mathbf{u}'} \right) = 0$$

- Remaining parameter:

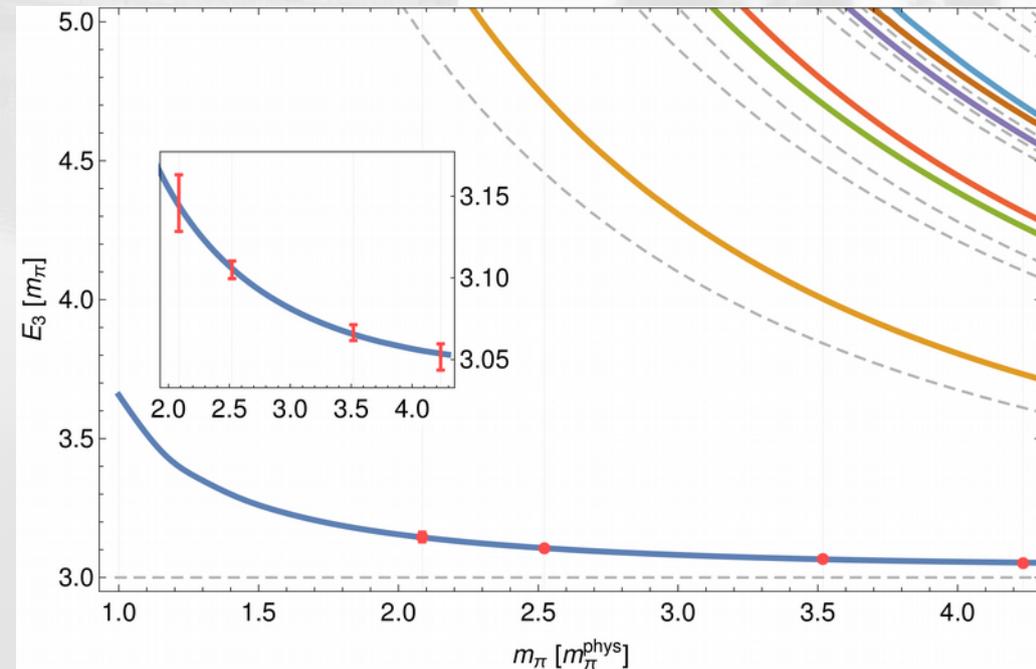


~ 3-body force

→ fit to the ground state levels by NPLQCD

→ $C=0.2 \pm 1.5 \cdot 10^{-10}$

- Exited states = prediction



3-BODY: INFINITE VOLUME

Phenomenology (exotics, etc...)

✓ Unitary isobar amplitude derived

→ 2b sub-amplitudes = tower of isobars

→ 3-dim. relativistic integral equation

Applications to $a_1(1260)$ and $a_1(1420)$

... in progress

3-BODY: FINITE VOLUME

✓ 3-body quantization condition derived

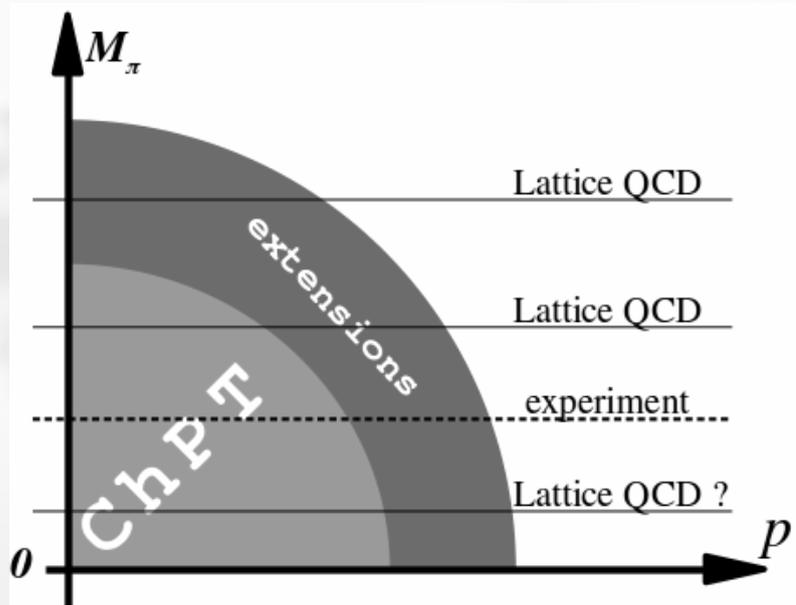
✓ Finite volume spectrum is investigated:

- 2-body sepectrum predicted
- 3b force fitted to NPLQCD ground state results ($c \sim 0$)
- excited 3-body spectrum is predicted

→ Future applications: $N^*(1440)$, $a_1(1260)$, ...

2-BODY EXAMPLE: $\sigma(500)$

New insights from a Lattice QCD calculations



- *phase-shifts* ($N_f=3$; $m_\pi=236/391$ MeV)
Briceno et al. (2016)

- *sc. length* ($N_f=2$; $m_\pi=139/240/330$ MeV)
ETMC (2017)

- *sc. length* ($N_f=3$; $m_\pi=247/249/314$ MeV)
Fu/Chen (2018)

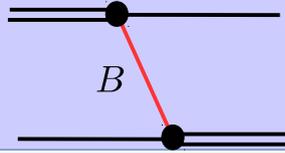
- *phase-shifts* ($N_f=2$; $m_\pi=227/315$ MeV)
Guo/.../MM/... (2018)

- *Mass & width around chiral limit*
Li/Pagels (1971) Bruns/MM (2017)

Briceno et al. (2016) / ETMC (2017) Fu/Chen (2018)

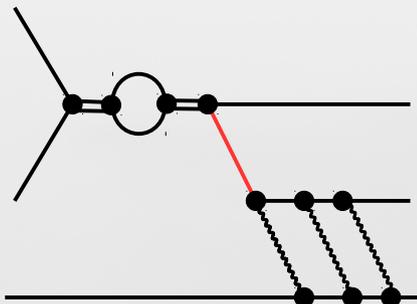
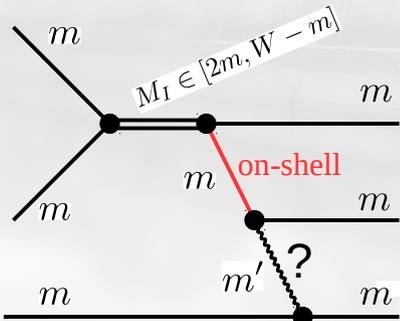
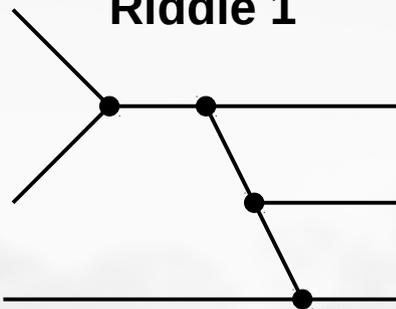
The Power of Unitarity

Question: Does

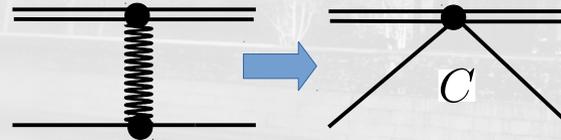
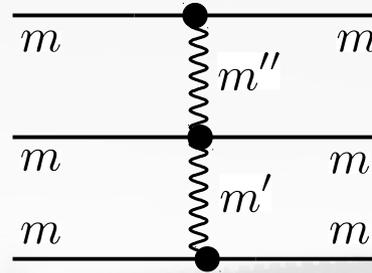


provide full imaginary part of all possible $3 \rightarrow 3$ transitions?

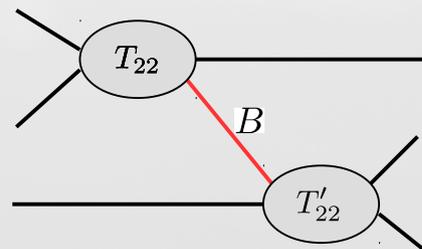
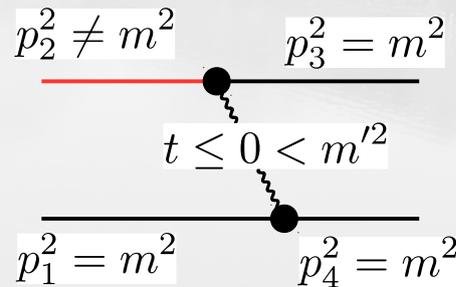
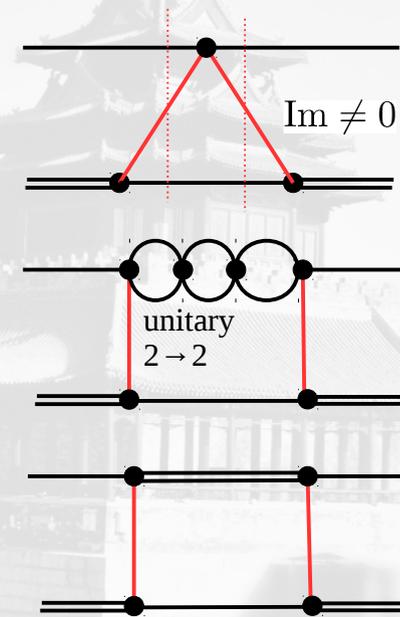
Riddle 1



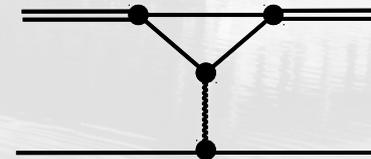
Riddle 2



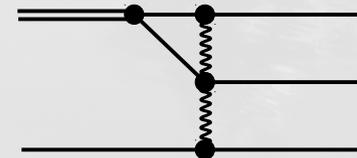
Riddle 3



Riddle 4

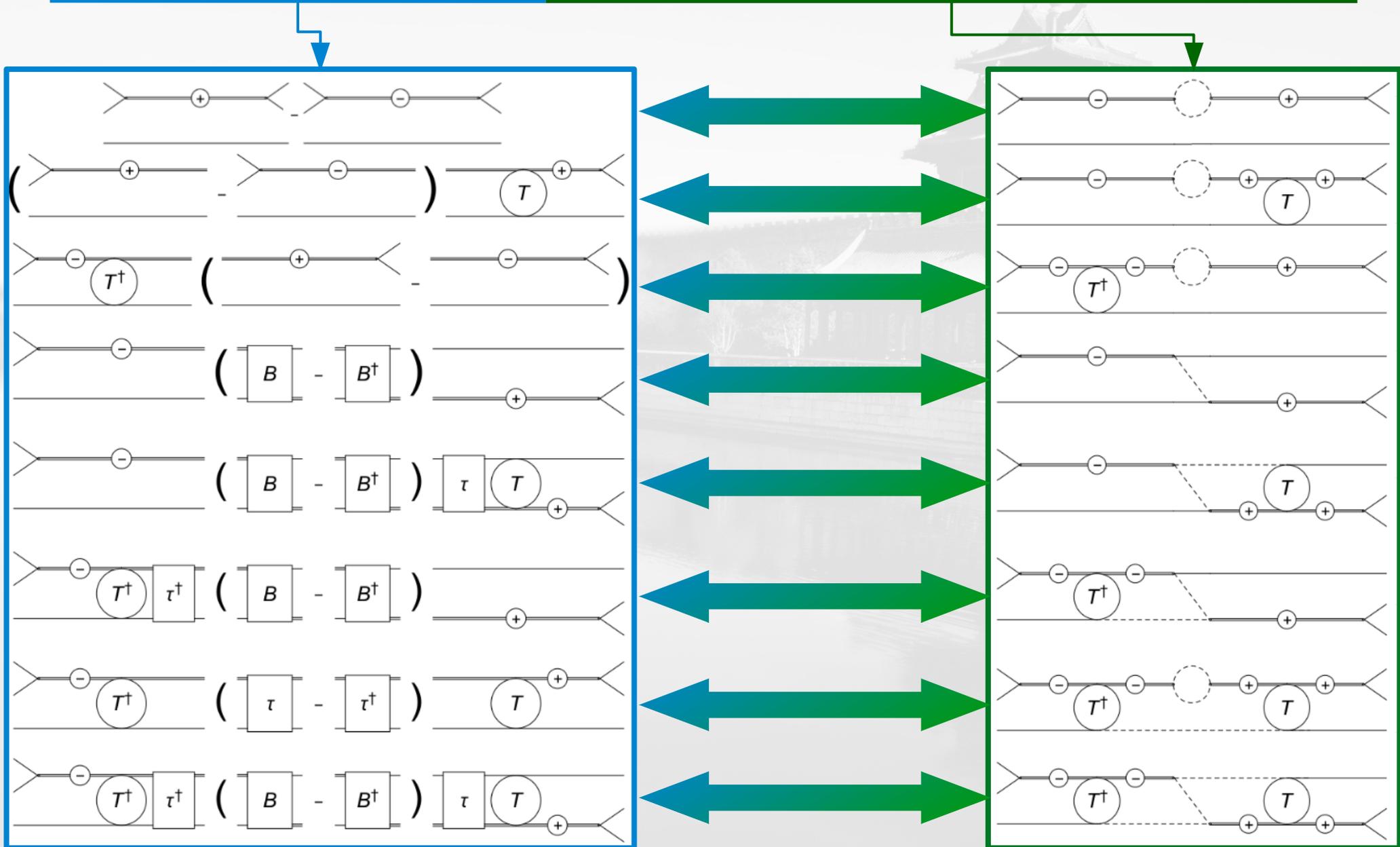


Riddle 5



3-body Unitarity (phase space integral)

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$



- Projection of T

$$T^{ss'}(\hat{\mathbf{p}}_j, \hat{\mathbf{p}}_{j'}) = 4\pi \sum_{\Gamma\alpha} \sum_{uu'} \chi_u^{\Gamma\alpha s}(\hat{\mathbf{p}}_j) T_{uu'}^{\Gamma ss'} \chi_{u'}^{\Gamma\alpha s'}(\hat{\mathbf{p}}_{j'}),$$

$$T_{uu'}^{\Gamma ss'} = \frac{4\pi}{\vartheta(s)\vartheta(s')} \sum_{j=1}^{\vartheta(s)} \sum_{j'=1}^{\vartheta(s')} \chi_u^{\Gamma\alpha s}(\hat{\mathbf{p}}_j) T^{ss'}(\hat{\mathbf{p}}_j, \hat{\mathbf{p}}_{j'}) \chi_{u'}^{\Gamma\alpha s'}(\hat{\mathbf{p}}_{j'})$$

Cancellations:

→ fin. vol. normalization of δ -distribution!

$$\bar{T}_{nm}^{A_1^+}(s) = \tau_n(s) T_{nm}^{A_1^+}(s) \tau_m(s) - 2E_n \tau_n(s) \frac{L^3}{\vartheta(n)} \delta_{nm}$$

$$T_{nm}^{A_1^+}(s) = B_{nm}^{A_1^+}(s) - \frac{1}{L^3} \sum_{x \in \text{sets}_8} \vartheta(x) B_{nx}^{A_1^+}(s) \frac{\tau_x(s)}{2E_x} T_{xm}^{A_1^+}(s)$$

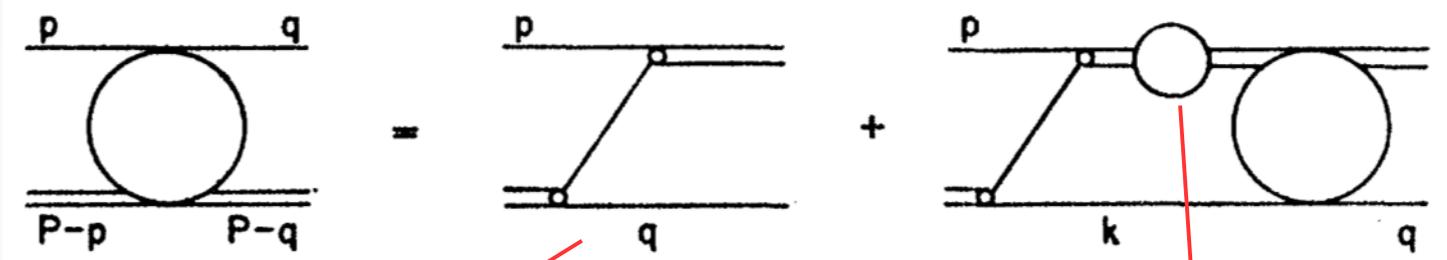
$B^{A_1^+}$ singular at $W^+ = E_m + E_n + E(\mathbf{q}_{nj} + \mathbf{p}_{mi})$

τ_m^{-1} singular at $W^{\pm\pm} = E_m \pm E((2\pi/L)\mathbf{y}) \pm E((2\pi/L)\mathbf{y} + \mathbf{p}_{mi})$ for $\mathbf{y} \in \mathbb{Z}^3$

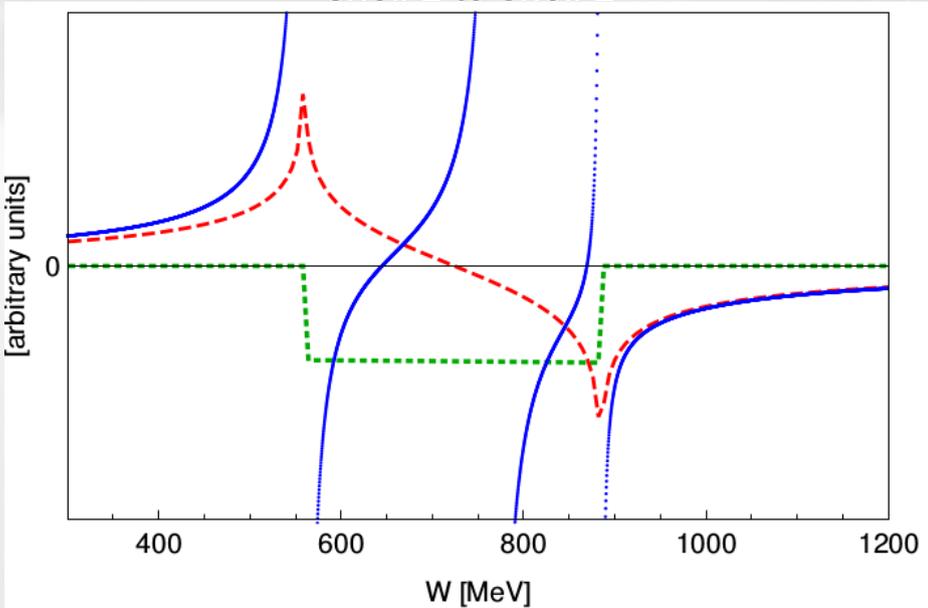
– when isobar-momenta are discretized in the 3-body cms momenta

$$\tau = \sigma(k) - M_0^2 - \frac{1}{(2\pi)^3} \int d^3\ell \frac{\lambda^2}{2E_\ell(\sigma(k) - 4E_\ell^2 + i\epsilon)}$$

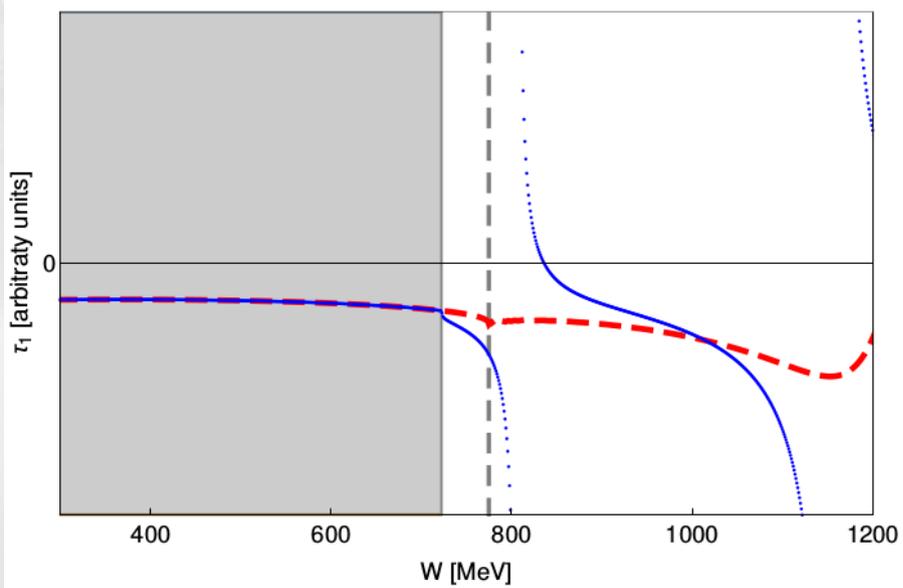
Power-law finite-volume effects dictated by three-body unitarity



S-wave infinite volume vs. A_1^+ finite volume shell 1 to shell 1

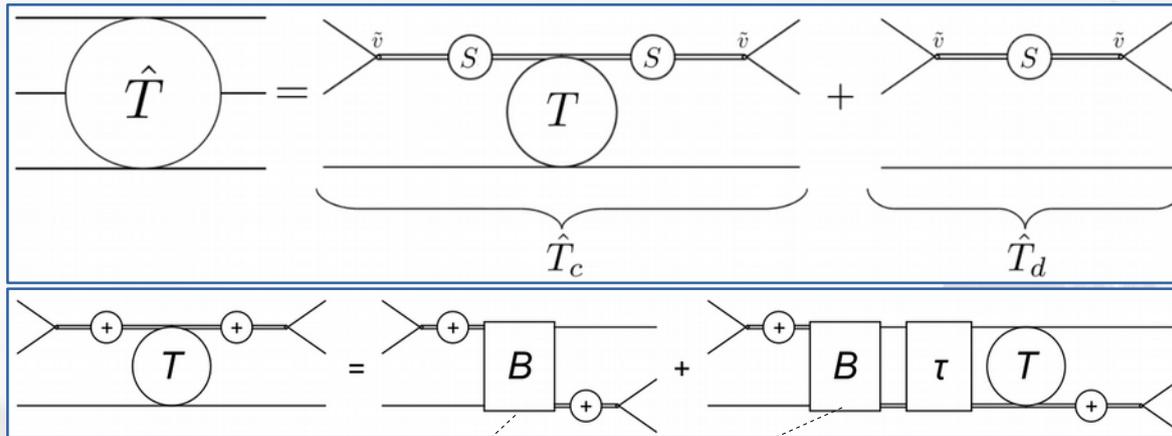


Tower of boosted 2 → 2 amplitudes to implement 3-body quantization condition



SCATTERING AMPLITUDE

3 → 3 scattering amplitude is a 3-dimensional integral equation



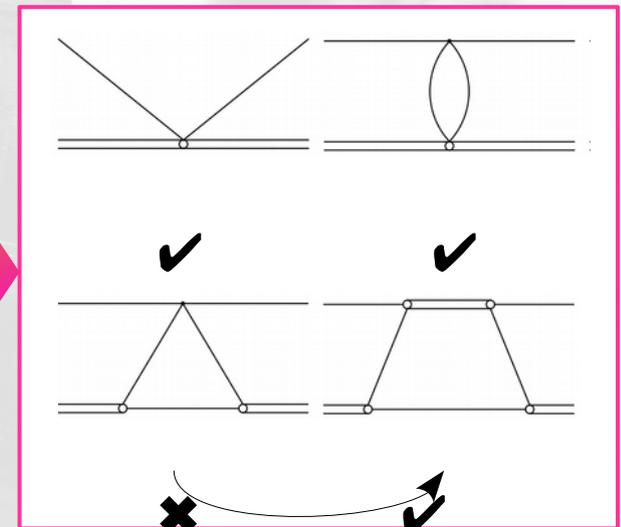
- Imaginary parts of B , S are fixed by **unitarity/matching**
- For simplicity $v=\lambda$ (full relations available)

$$\text{Disc } B(u) = 2\pi i \lambda^2 \frac{\delta(E_Q - \sqrt{m^2 + Q^2})}{2\sqrt{m^2 + Q^2}}$$

- un-subtracted dispersion relation

$$\langle q|B(s)|p\rangle = -\frac{\lambda^2}{2\sqrt{m^2 + Q^2} (E_Q - \sqrt{m^2 + Q^2} + i\epsilon)}$$

- one- π exchange in TOPT → **RESULT!**



Unitarity & Matching

- 3-body Unitarity (normalization condition \leftrightarrow phase space integral)

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

