# SAID for Baryon Spectroscopy 

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- PWA for Baryon Spectroscopy
- Pion Photoproduction.
- Pion Photoproduction on Neutron.
- Pion-Proton Elastic.
- Nucleon-Nucleon Elastic.
- Strange Hadron Spectroscopy.
- Summary.


## PWA for Baryons

- Originally PWA arose as technology to determine amplitude of reaction via fitting scattering data.
That is non-trivial mathematical problem - looking for solution of ill-posed problem following to Hadamard \& Tikhonov.
- Resonances appeared as by-product
[bound states objects with definite quantum numbers, mass, lifetime, \& so on].
- Standard PWA
$\Rightarrow$ Reveals only wide Resonances, but not too wide ( $\Gamma<500 \mathrm{MeV}$ ) \& possessing not too small $B R$ ( $B R>4 \%$ ).
$\Rightarrow$ Tends (by construction) to miss narrow Res with $\Gamma<20 \mathrm{MeV}$.
Most of our current knowledge about bound states of three light quarks
has come mainly from $\pi \mathrm{N} \rightarrow \pi \mathrm{N}$ PWAs:
Karlsruhe-Helsinki,
Carnegie-Mellon-Berkeley,
\& GW.
Main source of EM couplings is GW, BnGa, \& JuBo analyses.


## Road Map to Baryon Spectroscopy



## PNPDG



# SAID Database below 4 GeV 

SAID: http://gwdac.phys.gwu.edu/

- We update SAID databases, develop \& study PWAs, \& keep current versions of phenomenological \& theoretical models, both those of CNS/DAC \& other research groups, on continual basis for relevant two- \& three-body reactions of interest.
- In the full database, one will occasionally find experiments which give conflicting results.
- Some data with very large $\chi^{2}$ contributions have been excluded from our fits.
- Redundant data are also excluded [these include $\sigma_{\text {tot }}$ based on $d \sigma / d \Omega$ already contained in database]
- Measurements of pol observables ( $P$, for instance) with uncertainties more than 0.2 are not included as they have little influence in our fits.
- However, all available data have been retained in database (excluded data labeled as "flagged") so that comparisons can be made through our on-line facility
- Data Analysis Center Institute for Nuclear Studies THE GEORGE WASHINGTON UNIVERSITY WASHINGTON, OC

INS DAC Home<br>- INS DAC [SAID]<br>INS Home<br>Pi-N Newsletters

Partial-Wave Analyses at GW


| [See Instructions] |
| :--- |
| Pion-Nucleon |
| Pi-Pi-N |
| Kaon(+)-Nucleon |
| Nucleon-Nucleon |
| Pion Photoproduction |
| Pion Electroproduction |
| Kaon Photoproduction |
| Eta Photoproduction |
| Eta-Prime Photoproduction |
| Pion-Deuteron (elastic) |
| Pion-Deuteron to Proton+Proton |

[ $\mathrm{W}=1320$ to 1930 MeV ]


For $\pi \rightarrow 2 \pi$, we use log-likelihood while for rest -least-squares technologies.

## Pron Plillo Pronulion

World Progress in Pion PhotoProduction


# Direct Amplitude Reconstruction 

## in Pion PhotoProduction

## $\gamma \mathrm{N} \rightarrow \mathrm{N} \pi$

spin: $1 \quad \frac{1}{2} \rightarrow \frac{1}{2} \quad 0$ helicities: $2 \times 2 \times 2 / 2=$ (4)
parity conservation 个


- In particle physics, helicity is projection of the spin $\vec{S}$ onto direction of momentum, $\hat{p}$ :

$$
\begin{array}{r}
h=\vec{J} \cdot \hat{p}=\vec{L} \cdot \hat{p}+\vec{S} \cdot \hat{p}=\vec{S} \cdot \hat{p} \\
\hat{p}=\frac{\vec{p}}{|\vec{p}|}
\end{array}
$$

Therefore, there are 4 independent invariant amplitudes

- In order to determine pion photoproduction amplitude [4 modules and 3 relative phases], one has to carry ou $\mathcal{P}$ independent measurements at fixed $(\mathbf{W}, \mathrm{t})$ or $(\mathrm{E}, \theta)$.

8) This extra observable is necessary to eliminate sign ambiguity.

Ambiguities in the partial-wave analysis of pseudoscalar-meson photoproduction

## Complete Experiment for Pion PhotoProduction

- There are 16 non-redundant observables.
- They are not completely independent from each other.


1 un-pol measurement: d $\sigma / d \Omega$
3 single pol measurements: $\Sigma$, T, $\mathbf{P}$
12 double pol measurements: $\mathbf{E}, \mathrm{F}, \mathrm{G}, \mathrm{H}$,
$\mathrm{C}_{\mathrm{x}}, \mathrm{C}_{\mathbf{z}}, \mathrm{O}_{\mathrm{x}}, \mathrm{O}_{\mathbf{z}}, \mathrm{L}_{\mathrm{x}}, \mathrm{L}_{\mathbf{z}}, \mathrm{T}_{\mathrm{x}}, \mathrm{T}_{\mathbf{z}}$
18 triple polarization asymmetries
[ 9 for linear pol beam]
[ 9 for circular pol beam]
13 of them are non-vanishing
A. Sandorfi et al. AIP Conf. Proc. 1432, 219 (2012)
K. Nakayama, private communication, 2014

## Importance of $\mathcal{N}$ eutron Data

- EM interaction do not conserve isospin, so multipole amplitudes contain isoscalar \& isovector contributions of EM current.

Proton

$$
A_{\pi^{0} p}=A^{0}+\frac{1}{3} A^{1 / 2}+\frac{2}{3} A^{3 / 2}
$$

$$
A_{\pi^{+} n}=\sqrt{2}\left(A^{0}+\frac{1}{3} A^{1 / 2}-\frac{1}{3} A^{3 / 2}\right)
$$

$$
\begin{aligned}
& A_{\pi^{0} n}=-A^{0}+\frac{1}{3} A^{1 / 2}+\frac{2}{3} A^{3 / 2} \\
& A_{\pi^{-} p}=\sqrt{2}\left(A^{0}-\frac{1}{3} A^{1 / 2}+\frac{1}{3} A^{3 / 2}\right)
\end{aligned}
$$

- Proton data alone/does not allow separation of isoscalar \& isovector components.
- Need data on both proton \& neutron!
- Data Analysis Center Institute for Nuclear Studies THE GEORGE WASHINGTON UNIVERSITY WASHINGTON. DC

SAID for Pion Photo Production
P. Mattione et al, Phys. Rev. C 96, 035204 (2017)

```
- Data driven (model independent) analysis
- Energy dependent MA27
- E = 145-2700 MeV [W = 1080-2460 MeV]
- PWs = 60 [EM multipoles] [J < 6]
- Prms = 210
Constraint: Born [no free parameters to fit] \piN-PWA [no theoretical input]
```




## Photo-Decay Amplitudes in BW © Pole Forms

## - Pole is main signature of resonance.

 $1 / 2$ and $3 / 2$ photo-decay amplitudes in units of $10^{-3}(\mathrm{GeV})^{-1 / 2}$. Errors on the phases are generally $2-5 \mathrm{~d}$ grees. For isospin $1 / 2$ resonances the values of the proton target are given.

| Resonance | Breit-Wigner values |  |  |  | Pole values |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | (Mass, width) | $\Gamma_{x} / 2$ | $A_{1 / 2}$ | $A_{3 / 2}$ | $\left(\operatorname{Re} W_{P},-2 \operatorname{Im} W_{P}\right)$ | $R_{\pi}$ | $A_{1 / 2}$ | $A_{3 / 2}$ |
| $\Delta(1232) 3 / 2^{+}$ | $(1233,119)$ | 60 | $-141 \pm 3$ | $-258 \pm 5$ | (1211, 99) | $52\left[-47^{\circ}\right]$ | $-136 \pm 5\left[-18^{\circ}\right]$ | $-255 \pm 5\left[-6^{\circ}\right]$ |
| $N(1440) 1 / 2^{+}$ | (1485, 284) | 112 | $-60 \pm 2$ |  | $(1359,162)$ | 38 [ $-98^{\circ}$ ] | $-66 \pm 5\left[-38^{\circ}\right]$ |  |
| $N(1520) 3 / 2^{-}$ | (1515, 104) | 33 | $-19 \pm 2$ | $+153 \pm 3$ | (1515, 113) | $38\left[-5^{\circ}\right]$ | $-24 \pm 3\left[-7^{\circ}\right]$ | $+157 \pm 6\left[+10^{\circ}\right]$ |
| $N(1535) 1 / 2^{-}$ | (1547, 188) | 34 | $+92 \pm 5$ |  | $(1502,95)$ | $16\left[-16^{\circ}\right]$ | $+77 \pm 5\left[+4^{\circ}\right]$ |  |
| $N(1650) 1 / 2^{-}$ | $(1635,115)$ | 58 | $+35 \pm 5$ |  | $(1648,80)$ | 14 [ $-69^{\circ}$ ] | $+35 \pm 3\left[-16^{\circ}\right]$ |  |

R.L. Workman et al, Phys Rev C 87, 068201 (2013)
A. Svarc et al, Phys Rev C 89, 065208 (2014)


## Measurement of $\pi^{0}$ photoproduction on the proton at MAMI C

${ }^{6} \mathrm{~A}$. Braghieri $^{7}$ W. J. Briscoe, ${ }^{8}$ S. Cherepnya, ${ }^{9}$ F. Cividini, ${ }^{1}$ C. Collicott, ${ }^{10,11}$ S. Costanza, ${ }^{7,12}$ A. Denig, ${ }^{1}$ E. J. Downie, ${ }^{1,8}$ M. Dieterle, ${ }^{13}$ M. I. Ferretti Bondy, ${ }^{1}$ L. V. Fil'kov, ${ }^{9}$ A. Fix, ${ }^{14}$ S. Gardner, ${ }^{4}$ S. Garni, ${ }^{13}$ D. I. Glaziar ${ }^{4}$ ' ${ }^{\text {o }}$ owa ${ }^{15}$ W. Gradl, ${ }^{1}$ G. Gurevich, ${ }^{16}$ D. J. Hamilton, ${ }^{4}$ D. Hornidge, ${ }^{17}$ G. M. Huber, ${ }^{18}$ A. Käser, ${ }^{13}$ V. L. KashM. Korolija, ${ }^{19}$ B. Krusche, ${ }^{13}$ V. V. Kulikov, ${ }^{20}$ A. Larom
D. M. Manley, ${ }^{3}$ P. P. Martel, ${ }^{1,21}$ M. Martan-
A. Mushkarenkov, ${ }^{7,21}$ A. Negan-
P. Pedroni, ${ }^{7}$ A. Polonski, ${ }^{16}$ $\qquad$
$\qquad$
$\square$ $1) \begin{gathered}\text { Keshelashvili, }{ }^{13} \text { R. Kondratiev, } \\ \text { ston, }{ }^{4} \text { I. J. D. MacGregor }{ }^{4}\end{gathered}$ ston, ${ }^{4}$ I. J. D. MacGregor, ${ }^{4}$
G. Middleton,, 17 R. Miskimen ${ }^{21}$ A. Sarty, ${ }^{11}$ D. M. Schott, ${ }^{8}$ S. Sch $\quad \mid \quad, \ldots,{ }^{4}$ A. Rajabi, ${ }^{21}$ G. Reicherz, ${ }^{22}$ G. Ron, ${ }^{24}$ T. Rostomyan, ${ }^{13}$ I. Supek, ${ }^{19}$ M. F. Taragin, ${ }^{8}$ t.. a mel, ${ }^{2}$ M. Thiel, ${ }^{1}$ L. Tiator ${ }^{1}$ A ${ }^{1}$ K. Spieker, ${ }^{2}$ O. Steffen, ${ }^{1}$ I. I. Strakovsky, ${ }^{8, \dagger}$ Th. Strub D. P. Watts, ${ }^{15}$ D. Werthmüller, ${ }^{4,13}$ J. Wettig, ${ }^{1}$ L. Witthawe ${ }^{13}{ }^{13}$ M. W, M. Unverzagt, ${ }^{1}$ Yu. A. Usov, ${ }^{6}$ S. Wagner, (A2 Collaboration at MAMI)


$$
\begin{aligned}
& W=1135-1957 \mathrm{MeV} \\
& \theta=15-165^{0} \\
& \pi^{0} \mathrm{p}: 7978 \mathrm{~d} \sigma / \mathrm{d} \Omega
\end{aligned}
$$

## W = 1216-1448 MeV <br> $\theta=31-158^{0}$ <br> $\pi^{0} \mathrm{p}: 1403 \Sigma$

THE EUROPEAN PHYSICAL JOURNAL A

Photon asymmetry measurements of $\vec{\gamma} p \rightarrow \pi^{0} \mathbf{p}$ for $\mathrm{E}_{\gamma}=320-650 \mathrm{MeV}$


## The A2 Collaboration at MAMI

S. Gardner ${ }^{1, a}$, D. Howdle ${ }^{1}$, M.H. Sikora ${ }^{6}$, Y. Wunderlich ${ }^{4}$, S. Abt $^{10}$, P. Achenbach ${ }^{2}$, F. Afzal ${ }^{4}$, P. Aguar-Bartolome ${ }^{2}$, Z. Ahmed ${ }^{14}$, J.R.M. Annand ${ }^{1}$, H.J. Arends ${ }^{2}$, K. Bantawa ${ }^{3}$, M. Bashkanov ${ }^{6}$, R. Beck ${ }^{4}$, M. Biroth ${ }^{2}$, N.S. Borisov ${ }^{15}$, A. Braghieri'5, W.J. Briscoe ${ }^{7}$, S. Cherepnya ${ }^{9}$, F. Cividini ${ }^{2}$, S. Costanza ${ }^{5}$, C. Collicott ${ }^{7}$, B.T. Demissie ${ }^{7}$, A. Denig ${ }^{2}$, M. Dieterle ${ }^{10}$, E.J. Downie ${ }^{7}$, P. Drexler ${ }^{2}$, M.I. Ferretti-Bondy ${ }^{2}$, L.V ${ }^{-}{ }^{-}{ }^{\text {ov }}{ }^{9}$, D.I. Glazier ${ }^{1}$, S. Garni ${ }^{10}$, W. Gradl ${ }^{2}$, M. Günther ${ }^{10}$, G.M. Gurevich ${ }^{\text {i2 }}$, P. Hall Barrientos ${ }^{6}$, D. O. Jahn ${ }^{2}$, T.C. Jude ${ }^{6}$, A. Käser ${ }^{10}$, S. Kay ${ }^{6}$, V.L. ${ }^{\boldsymbol{V}}$, B. Krusche ${ }^{10}$, J.M. Linturi ${ }^{2}$, V. Lisin
D.M. Manley ${ }^{3}$, P.P. Martel ${ }^{2}$, J. D.M. Manley ${ }^{3}$, P.P. Martel ${ }^{2}$, J..
$\qquad$

 ${ }^{2}$, D. Paudyal ${ }^{14}$, P. P ${ }^{1}$. olonski ${ }^{12}$, S. Prakhov ${ }^{8}$, A. Rajabi ${ }^{17}$, J. Robinson ${ }^{1}$, G. Rosner ${ }^{1}$, ${ }^{1}$, ${ }^{\mathrm{an}^{10}}$, A. Sarty ${ }^{16}$, S. Sc _umann ${ }^{2}$, V. Sokhoyan ${ }^{2}$, K. Spieker ${ }^{4}$, O. Steffen ${ }^{2}$, C. Sfienti ${ }^{2}$, I.I. Strakovsky ${ }^{7}$ erg $^{1}$, Th. Strub ${ }^{10}$, I. Supek ${ }^{13}$, C.M. Tarbert ${ }^{6}$, A. Thiel ${ }^{4}$, M. Thiel ${ }^{2}$, A. Thomas ${ }^{2}$, M. Unverzagt ${ }^{2}$ ${ }^{5}$, D.P. Watts ${ }^{6}$, D. Werthmüller ${ }^{1,10}$, J. Wettig ${ }^{2}$, M. Wolfes ${ }^{2}$, L. Witthauer ${ }^{10}$, and L. Zana ${ }^{6}$

## PHYSICAL REVIEW C 93, 055209 (2016)

## $T$ and $F$ asymmetries in $\pi^{0}$ photoproduction on the proton

J. R. M. Annand, ${ }^{1}$ H. J. Arends, ${ }^{2}$ R. Beck, ${ }^{3}$ N. Borisov, ${ }^{4}$ A. Braghie
J. R. M. Annand, ${ }^{1}$ H. J. Arends, ${ }^{2}$ R. Beck, ${ }^{3}$ N. Borisov, ${ }^{4}$ A. Braghieri, ${ }^{5} \mathrm{~W}$ r
S. Costanza, ${ }^{5}$ E. J. Downie, ${ }^{2,6}$ M. Dieterle, ${ }^{9}$ A. Fix,${ }^{10}$ L. V. Fil'kov ${ }^{7}-$
P. Hall Barrientos, ${ }^{10}$ D. Hamilton, ${ }^{1}$ D. Hornidge, ${ }^{13}$ D $v$ R. Kondratiev, ${ }^{12}$ M. Korolija, ${ }^{15}$ B. Krusche ${ }^{9}$. D. M. Manley, ${ }^{16}$ P. P. Martel, ${ }^{2,17} \mathrm{E} .{ }^{-}$. $\quad$ POT Neganov, ${ }^{4}$ A. Nikolaev ${ }^{3}$ M ${ }^{\circ}$ berle, V. V. Polyanski, ${ }^{7}$ S. Prakhovv, ${ }^{19}$ G. 1 Th. Strub, ${ }^{9}$ I. Supek, ${ }^{15}$ L. Tiator; 4. 1 r. B. Otte, ${ }^{2}$ B. Oussena, ${ }^{2,6}$ P. Pedroni, ${ }^{5}$ A. Polonski, ${ }^{12}$
_.tyan, ${ }^{4}$ A. Sarty, ${ }^{8}$ S. Schumann, ${ }^{2}$ O. Steffen, ${ }^{2}$ I. I. Strakovsky, ${ }^{6}$
,.tas, ${ }^{2}$ M. Unverzagt, ${ }^{2}$ Yu. A. Usov, ${ }^{4}$ D. P. Watts, ${ }^{11}$ D. Werthmüller, ${ }^{1,9}{ }^{1,9}$
L. Witthaynverzagt, ${ }^{2}$ Yu. A. Usov, ${ }^{4}$ D. P. Watts, ${ }^{11}$ D. Werthmüller, ${ }^{1,9}$
L. Witthauer, ${ }^{9}$ and M. Wolfes ${ }^{2}$
(A2 Collaboration at MAMI)





$W=1322-1841 \mathrm{MeV}$
$\theta=75-140^{0}$
$\pi^{0} \mathrm{p}: 45 \mathrm{Cx}^{\prime}$




Measurement of the ${ }^{1} H(\vec{\gamma}, \vec{p}) \boldsymbol{\pi}^{0}$ Reaction Using a Novel Nucleon Spin Polarimeter
M. H. Sikora, ${ }_{4}{ }^{, *}$ D.P. Watts, ${ }^{1}$ D. I. Glazier, ${ }^{1}$ P. Aguar-Bartolomé, ${ }^{2}$ L. K. Akasoy, ${ }^{2}$ J. R. M. Annand, ${ }^{3}$ H. J. Arends, ${ }^{2}$ K. Bantawa, ${ }^{4}$ R. Beck, ${ }^{5}$ V.S. Bekrenev, ${ }^{6}$ H. Berghäuser, ${ }^{7}$ A. Braghieri, ${ }^{8}$ D. Branford ${ }^{1}$, ${ }^{1}$ J. Briscoe, ${ }^{9}$ J. Brudvik, ${ }^{1}$ S. Cherepnya, ${ }^{11}$ R.E. B. Codling, ${ }^{3}$ B.T. Demissie, E. J. Downie, ${ }^{2,39}$ P n-
D. Hamilton, ${ }^{3}$ E. Heid, ${ }^{29}$ D. Homidge, ${ }^{12}$ D. A. How. ${ }^{2}$, ${ }^{11}$ B. Freehart, ${ }^{9}$ R. Gregor,

 B. McKinnon, ${ }^{3}$ E. F. McNicoll, ${ }^{3}$ D. W. Micanovic ${ }^{15}$ D. G. Middleton, ${ }^{12}$ A. Mushkarenkov, B. M. K. Nefkens, ${ }^{10}$ A. Nikolaev,
 Uuy, M. Ostrick, ${ }^{2}$ P. B. Otte, ${ }^{2}$ B. Oussena ${ }^{2,9}$ P. Pedroni, ${ }^{8}$ F. Pheron, ${ }^{13}$
I. Strakovsky, ${ }^{9}$ I. M. Suarez ${ }^{\text {i0 }}$ I Supek ${ }^{15}$ M. Thiel, ${ }^{1}$ A. Thomas, ${ }^{2}$ M. Unverzagt ${ }^{2}$ D. Werthmüller ${ }^{1 / 3}$ R. L. Workman, ${ }^{2}$, I. Zamboni, ${ }^{15}$ and F Zehr ${ }^{13}$ (A2 Collaboration at MAMI)


7/16/2018
PWA10/ATHOS5 2018, Beijing, China, July 2018
William Briscoe 13


Courtesy of Steffen Strauch, 2018

PHYSICAL REVIEW C 88,065203 (2013)
Beam asymmetry $\Sigma$ for $\pi^{+}$and $\pi^{0}$ photoproduction on the proton for ph
from 1.102 to 1.862 GeV from 1.102 to 1.862 GeV
clos

$$
\left(\frac{d \sigma}{d \Omega}\right)=\left(\frac{d \sigma}{d \Omega}\right)_{0}\left(1-P_{z} P_{\odot} E\right)
$$

M. Dugger, ${ }^{2}$ B. G. Ritchie, ${ }^{2}$ P. Collins, ${ }^{2, *}$ E. Pasyuk, ${ }^{2, \dagger}$ W. J. Briscoe, ${ }^{14}$ I. I. Strakovsky, ${ }^{14}$ R. L. Workman, ${ }^{14}$ Y. Azimov, ${ }^{29}$ K. P. Adhikari, ${ }^{28}$ D. Adikaram, ${ }^{28}$ M. Aghasyan, ${ }^{17}$ M. J. Amaryan, ${ }^{28}$ M. D. Anderson, ${ }^{37}$ S. Anefalos Pereira, ${ }^{17}$ H. Avakian, ${ }^{15}$ J. Ball, ${ }^{6}$ N. A. Baltzell, ${ }^{1,34}$ M. Battaglieri, ${ }^{18}$ V. Batourine, ${ }^{23,35}$ I. Bedlinskiy, ${ }^{21}$ A. S. Biselli, ${ }^{4,10}$ S. Boiarinov, ${ }^{35}$ V. D. Burkert, ${ }^{35}$ D. S. Carman, ${ }^{35}$ A. Celentano, ${ }^{18}$ S. Chandavar, ${ }^{27}$ P. L. Cole, ${ }^{15}$ M. Contalbrigo, ${ }^{16}$ O. Cortes, ${ }^{15}$ V. Crede, ${ }^{12}$ A. D'Angelo, ${ }^{19,32}$





 R. A. Montgomery,${ }^{37} \mathrm{H} . \mathrm{M} 2$, . Munevar,,${ }^{35} \mathrm{C}$. Munoz Camacho, ${ }^{20}$ P. Nadel-Turonski, ${ }^{14,35} \mathrm{C}$. S. Nepali, ${ }^{28}$ S. Niccolai, ${ }^{20}$ G. Niculescu, ${ }^{22}$ I. Niculescu, ${ }^{22}$ M. Osipenko, ${ }^{18}$ A. I. Ostrovidov, ${ }^{12}$ L. L. Pappalardo, ${ }^{16}$ R. Paremuzyan, ${ }^{40,8}$ K. Park, ${ }^{23,35}$ S. Park, ${ }^{12}$ E. Phelps, ${ }^{34}$ J. J. Phillips, ${ }^{37}$ S. Pisano, ${ }^{17}$ O. Pogorelko, ${ }^{21}$ S. Pozdniakov, ${ }^{21}$ J. W. Price, ${ }^{3}$ S. Procureur, ${ }^{6}$ Y. Prok, ${ }^{28,35,38}$ D. Protopopescu, ${ }^{37}$ B. A. Raue, ${ }^{11,35}$ D. Rimal, ${ }^{11}$ M. Ripani, ${ }^{18}$ G. Rosner, ${ }^{37}$ P. Rossi, ${ }^{17,35}$ F. Sabatié, ${ }^{6}$ M. S. Saini, ${ }^{12}$
C. Salgado, ${ }^{26}$ D. Schott, ${ }^{14}$ R. A. Schumacher, ${ }^{4}$ E. Seder, ${ }^{8}$ H. Seraydaryan,${ }^{28}$ Y. G. Sharabian, ${ }^{35}$ G. D. Smith, ${ }^{37}$ D. I. Sober, ${ }^{5}$ D. Sokhan, ${ }^{37}$ S. S. Stepanyan, ${ }^{23}$ P. Stoler ${ }^{30}$ S. Strauch, ${ }^{14,34}$ M. Taiuti, ${ }^{13,5}$ W. Tang, ${ }^{27}$ Ye Tian, ${ }^{34}$ S. Tkachenko, ${ }^{28,38}$ B. Torayev ${ }^{28}$ H. Voskanyan, ${ }^{40}$ E. Voutier,,$^{24}$ N. K. Walford, ${ }^{5}$ D. P. Watts, ${ }^{9}$ D. P. Weygand,,$^{35}$
N. Zachariou, ${ }^{34}$ L. Zana, ${ }^{25}$ J. Zhang, ${ }^{28,35}$ Z. W. Zhao, ${ }^{38}$ and I. Zonta ${ }^{19, \|}$



7/16/2018
PWA10/ATHOS5 2018, Beijing, China, July 2018


First measurement of the polarization observable $E$ in the $\vec{p}\left(\vec{\gamma}, \pi^{+}\right) n$ reaction up to 2.25 GeV chas collaboration








 KL Givanemi Ex Girod ${ }^{2}$













 x. Wei

1. Zonat ${ }^{\text {m. }}$. Wood

G for $\vec{\gamma} \vec{p} \rightarrow \pi^{+} n$


Courtesy of Lorenzo Zana 2018
PWA10/ATHOS5 2018, Beijing, China, July 2018
$\mathrm{T} \& \mathrm{~F}$ for $\vec{\gamma} \vec{p} \rightarrow \pi^{+} n$



Courtesy of Mike Dugger, 2018

P \& H for $\vec{\gamma} \vec{p} \rightarrow \pi^{+} n$







Courtesy of Mike Dugger, 2018

## Pion PRok Prourtion

## © <br>  <br> 

Single Pion PhotoProduction on

## "Neutron" Target

- Accurate evaluation of EM couplings $\mathbf{N}^{*} \rightarrow \gamma \mathbf{N} \& \Delta^{*} \rightarrow \gamma \mathbf{N}$ from meson photoproduction data remains paramount task in hadron physics.
- Only with good data on both proton \& neutron targets, one can hope to disentangle isoscalar \& isovector EM couplings of various $N^{*} \& \Delta^{*}$ resonances, K.M. Watson, Phys Rev 95, 228 (1954): R.L. Walker, Phys Rev 182, 1729 (1969) as well as isospin properties of non-resonant background amplitudes.
- The lack of $\mathbf{Y n} \rightarrow \boldsymbol{\pi}^{-} \boldsymbol{p}$ \& $\mathbf{Y n} \rightarrow \boldsymbol{\pi}^{0} \boldsymbol{n}$ data does not allow us to be as confident about determination of neutron couplings relative to those of proton.
- Radiative decay width of neutral baryons may be extracted from $\pi^{-} \& \pi^{0}$ photoproduction off neutron, which involves bound neutron target \& needs use of model-dependent nuclear (FSI) corrections.
A.B. Migdal, JETP 1, 2 (1955); K.M. Watson, Phys Rev 95, 228 (1954)


# FSI for $\gamma d \rightarrow \pi p \mathcal{N} \Longrightarrow \gamma n \rightarrow \pi \mathcal{N}$ 

V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, IS, Phys Rev C 84, 035203 (2011) V. Tarasov, A. Kudryavtsev, W. Briscoe, B. Krusche, IS, M. Ostrick, Phys At Nucl 79, 216 (2016)

- FSI plays critical role in state-of-the-art analysis of $\gamma \mathrm{n} \rightarrow \pi \mathrm{N}$ data.
$\bullet$ For $\gamma \mathrm{n} \rightarrow \pi \mathrm{N}$, effect is $5 \%-60 \%$. It depends on ( $\mathrm{E}, \theta$ ).


Input: SAID: $\gamma \mathrm{N} \rightarrow \pi \mathrm{N}, \pi \mathrm{N} \rightarrow \pi \mathrm{N}, \mathrm{NN} \rightarrow \mathrm{NN}$ amplitudes for 3 leading terms.
DWF: full Bonn NN Potential (there is no sensitivity to DWF).

$$
R=\left(d \sigma / d \Omega_{\pi p}\right) /\left(d \sigma^{I A} / d \Omega_{\pi p}\right) \quad \square \quad \frac{d \sigma}{d \Omega}(\gamma n)=R^{-1} \frac{d \sigma}{d \Omega}(\gamma d)
$$

# THTF 

V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, IS, Phys Rev C 84, 035203 (2011)



# CL.AS g13 Impact for $\mathcal{N}$ veutron $S=0 \mathcal{Q} I=1 / 2$ Couplings 

P. Mattione et al, Phys. Rev. C 96, 035204 (2017)

- Selected photon decay amplitudes $\mathbf{N}^{*} \rightarrow \mathbf{\gamma n}$ at resonance poles are determined for the first time.

$\overline{\text { GW \&ing }}$


MAID


10PDG
CL.AS g14 Impact for $\mathcal{N}$ eutron $S=0 \mathcal{Z} I=1 / 2$ Couplings
D. Ho et al, Phys Rev Lett 118, 242002 (2017)

| BW | $\mathrm{A}^{1 / 2}$ | $\left(10^{-3} \mathrm{GeV}^{-1 / 2}\right)$ | $\mathrm{A}^{3 / 2}$ | $\left(10^{-3} \mathrm{GeV}^{-1 / 2}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
|  | g14 PRL | previous | g14 PRL | previous |
| SAID |  |  |  |  |
| $N(1720) 3 / 2^{+}$ | $-9 \pm 2$ | $-21 \pm 4$ | +19 $\pm 2$ | $-38 \pm 7$ |
| $N(2190) 7 / 2^{-}$ | $-6 \pm 9$ | --- | -28 $\pm 10$ | --- |
| BnGa |  |  |  |  |
| $N(1720) 3 / 2^{+}$ | tbd | $-80 \pm 50$ | tbd | $-140 \pm 65$ |
| $\mathrm{N}(2190) 7 / 2^{-}$ | +30 $\pm 7$ | $-15 \pm 12$ | $-23 \pm 8$ | $-33 \pm 20$ |

$\bullet$ I $=3 / 2$ waves $\sim$ unchanged $\Longleftrightarrow$ determined by proton data.

- Inclusion of these g14 data in new PWA calculations has resulted in revised $\gamma \mathbf{N}^{*}$ couplings $\&$, in case of $\mathbf{N}(\mathbf{2 1 9 0}) 7 / \mathbf{2}^{-}$, convergence among different PWA groups.
- Such couplings are sensitive to dynamical process of $\mathrm{N}^{*}$ excitation \& provide important guides to nucleon structure models.


## MAX-lab

- It is difficult task to measure $\pi^{-} p$ final state close to threshold.
- We measured $\pi^{0}$ decay in to $2 \gamma$ from $\gamma n \rightarrow \pi^{-} p \rightarrow \pi^{0} n$.


Courtesy of Bruno Strandberg, 2018


Courtesy of Daria Sokhan, 2018

# FSI for $\gamma d \rightarrow \pi^{0} n p \Longrightarrow \gamma n \rightarrow \pi^{0} n \mathcal{A} \gamma p \rightarrow \pi^{0} p$ 

V. Tarasov, A. Kudryavtsev, W. Briscoe, B. Krusche, IS, M. Ostrick, Phys At Nucl 79, 216 (2016)

- $\gamma \mathbf{n} \rightarrow \pi^{0} \mathbf{n}$ case is much more complicated vs. $\gamma \mathbf{n} \rightarrow \pi^{-} \mathbf{p}$ because $\pi^{0}$ can come from both $\gamma \mathbf{n}$ \& $\gamma \mathbf{p}$ initial interactions.

$$
\begin{aligned}
& A\left(\gamma p \rightarrow \pi^{0} p\right)=A_{v}+A_{s} \\
& A\left(\gamma n \rightarrow \pi^{0} n\right)=A_{v}-A_{s}
\end{aligned}
$$

- The corrections for both target nucleons are practically identical for $\pi^{0}$ production in energy range of $\Delta(1232) 3 / 2^{+}$due to isospin structure of $\gamma \mathbf{N} \rightarrow \pi \mathbf{N}$ amplitude. $A_{s}=0$ or $A_{v}=0$

$$
R_{n}=R_{p}
$$



- In general case, $R_{n} \neq R_{p}$


# Meson Production off(Deuteron at CB@MAMMI 

V. Kulikov et al, in progress

- Differential cross sections for $\gamma \mathbf{n} \rightarrow \pi^{0} n$.


Data up to $\mathrm{E}=1500 \mathrm{MeV}$ are coming.



Courtesy of Beatrice Ramstein, 2018

## PHYSICAL REVIEW C 93, 062201(R) (2016)

Search for narrow resonances in $\pi p$ elastic scattering from the EPECUR experiment
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## (a) Curlean (9) Curleon



Evidence for a New Resonance from Polarized Neutron-Proton Scattering
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$$ $11 \mathrm{~A} \quad \stackrel{\text { Marciniewski, }}{ }{ }^{3} \mathrm{~B}$. Marianski,

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(WASA-at-COSY Collaboration)



## Sensitivity of the COSY dibaryon candidate to $n p$ elastic scattering measurements

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Institute for Nuclear Studies, Department of Physics, The George Washington University, Washington, DC 20052, USA (Received 24 January 2016; published 1 April 2016)
$7 / 16 / 2018$
$\left.\operatorname{Re}\left({ }^{3} D_{3}\right)\right)^{0.0}$
PWA10/ATHOS5 2018, Beijing, China, July 2018

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## Aims of JГab KL.F Project



- KLF project has to establish secondary $K_{L}$ beam line at Jefferson Lab with flux of three order of magnitude higher than SiAC
- for scattering experiments on both proton \& neutron (first time !) targets in order to determine differential cross sections \& self-polarization of strange hyperons with Guble detector to enable precise PWA in order to determine all resonances up to 3 GeV in spectra of $\Lambda^{*}, \Sigma^{*}, \Xi^{*}, \& \Omega^{*}$.
- In addition, we intend to do strange meson spectroscopy by studies of $\pi$-K interaction to locate pole positions in $\mathrm{I}=1 / 2 \& 3 / 2$ channels.
- KLF has link to ion-ion high energy facilities as
 formation of our world in several microseconds after Big Bang.


## Gule

## Hall D Beam Line Set up for K-Congs



$$
\begin{array}{ll}
\mathrm{I}_{\mathrm{e}} & =5 \mu \mathrm{~A} \\
\mathrm{~W} \text {-radiator } & =0.1 \mathrm{R} . \mathrm{L} . \\
\text { Be-target } & =1.7 \mathrm{R} . \mathrm{L} .
\end{array}
$$

- Electrons are hitting W-radiator at CPS.
- Photons are hitting Be-target at cave.
$-\mathrm{K}_{\mathrm{L}} \mathrm{s}$ are hitting the $\mathrm{LH}_{2} / \mathrm{LD}_{2}$ target within GLueX setting.


## Expected Cross Sections vs Buб6โe Chamber Data

$\bullet$ GlueX measurements will span $\cos \theta$ from $\mathbf{- 0 . 9 5}$ to $\mathbf{0 . 9 5}$ in CM above $\mathbf{W}=\mathbf{1 4 9 0} \mathrm{MeV}$.

- $\mathrm{K}_{\mathrm{L}}$ rate is $10^{4} \mathrm{~K}_{\mathrm{l}} / \mathrm{s}=2500 \mathrm{x}$ SLAC
- Uncertainties (statistics only) correspond to 100 days of running time for:
$K_{L} p \rightarrow K_{s} p$
5 Geant4



Expected GlueX Data

$$
\mathrm{K}_{\mathrm{L}} \mathrm{p} \rightarrow \pi^{+} \Lambda
$$



## Why We Have to Measure

## Double-Strange Cascades in JLab

- Heavy quark symmetry (Isgur-Wise symmetry) suggests that multiplet splittings in strange, charm, \& bottom hyperons should scale as approximately inverses of corresponding quark masses:

$$
1 / m_{s}: 1 / m_{c}: 1 / m_{b}
$$

N. Isgur \& M.B. Wise, Phys Rev Lett 661130 (1991)

- If they don't, that scaling failure implies that structures of corresponding states are anomalous, \& very different from one another.
- So far only hyperon resonance multiplet, where this scaling can be "tested" \& seen is lowest negative parity multiplet:


## $\Lambda(1405) 1 / 2^{-}-\Lambda(1520) 3 / 2^{-}, \Lambda_{c}(2595) 1 / 2^{-}-\Lambda_{c}(2625) 3 / 2^{-}, \Lambda_{b}(5912) 1 / 2^{-}-\Lambda_{b}(5920) 3 / 2^{-}$

- It works approximately (30\%) well for those $\Lambda$-splittings. It would work even better for $\Xi, \Xi_{c}, \Xi_{b}$ splittings, \& should be very good for $\Omega, \Omega_{c}, \Omega_{\mathrm{b}}$ splittings.
- Jefferson Lab
As $L H C b$
HCGP
 is doing double charm cascade spectrum. $\Xi_{\mathrm{c}}(2790) 1 / 2^{-}-\Xi_{\mathrm{c}}(2815) 3 / 2^{-}$

R. Aaij et al, Phys Rev Lett 119, 112001 (2017)


## on $K \pi$ Scattering with Regards to $\kappa$ Meson in JLab

- KLF proposal will have very significant impact in our knowledge of $K \pi$ scattering amplitudes in scalar I = ½ channel.
- It will reduce by more than factor of two uncertainty in mass determination \& by a factor of five uncertainty on its width (and therefore on its coupling) of controversial or $k(800)$.
- Neutral kaon beam scattering on both proton \& neutron targets at low t-Mandelstam will allow to produce \& identify all four isospin partners of $\kappa(800)$.

- 203 researchers from 61 institutes
are co-authors.

Strange Hadron Spectroscopy with Secondary $K_{L}$ Beam at GlueX

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    ## Full Proposal was submitted for JLab PAC46

