Pion photoproduction on the deuteron and extraction of neutron-target observables

arXiv:1804.04757

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Introduction

Why neutron target data ?

First motivation

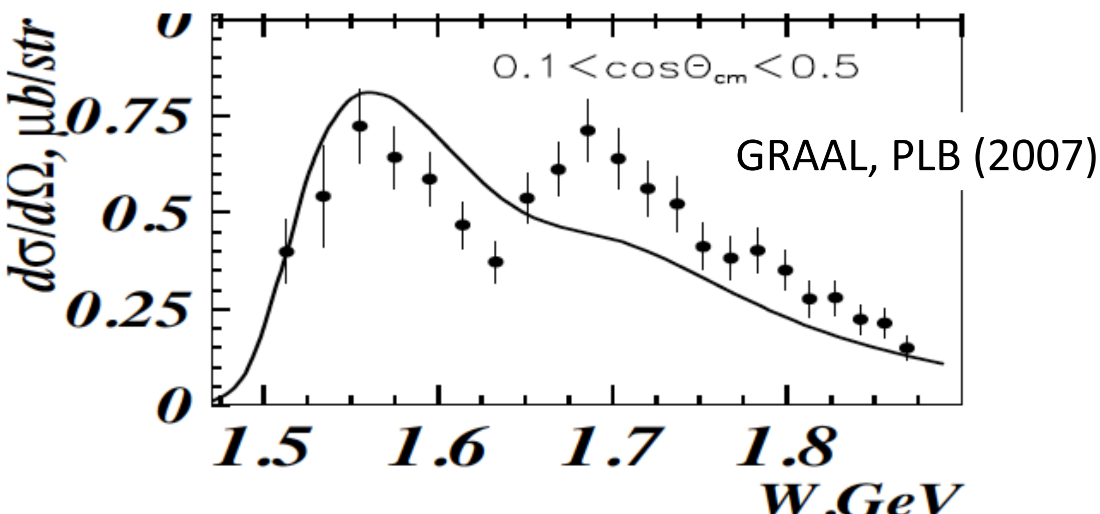
$\gamma^{(*)}\, p \rightarrow N^*$ AND $\gamma^{(*)}\, n \rightarrow N^*$ form factors $\Rightarrow$ Isospin structure of $\gamma^{(*)} N \rightarrow N^*$ form factors

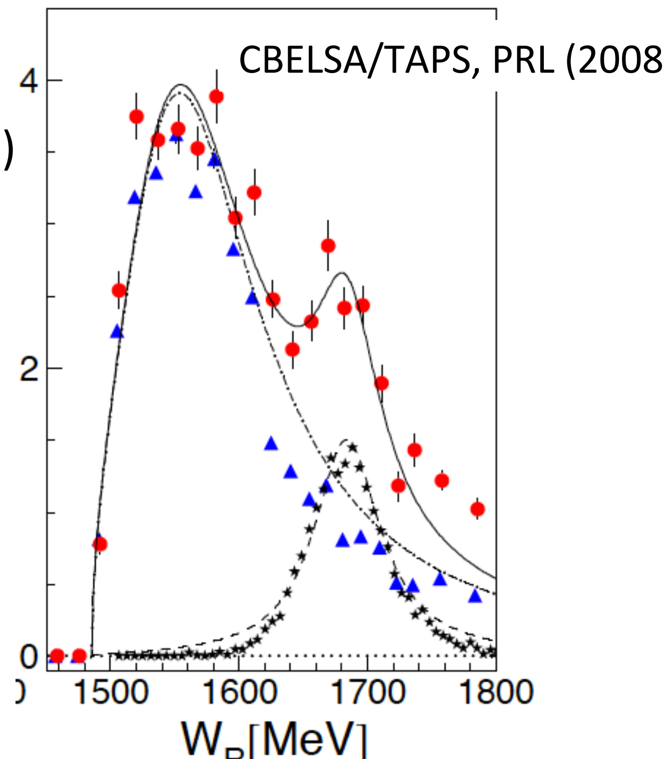
- Interesting quantities for understanding hadron structures
- Necessary for application to neutrino-induced reactions

  → needed for analyzing data from neutrino-oscillation experiments SXN et al. PRD (2015)

Unexpected surprises

Narrow bump at $W \sim 1.68$ GeV in $\gamma\, n \rightarrow \eta n$ (not in πN data)



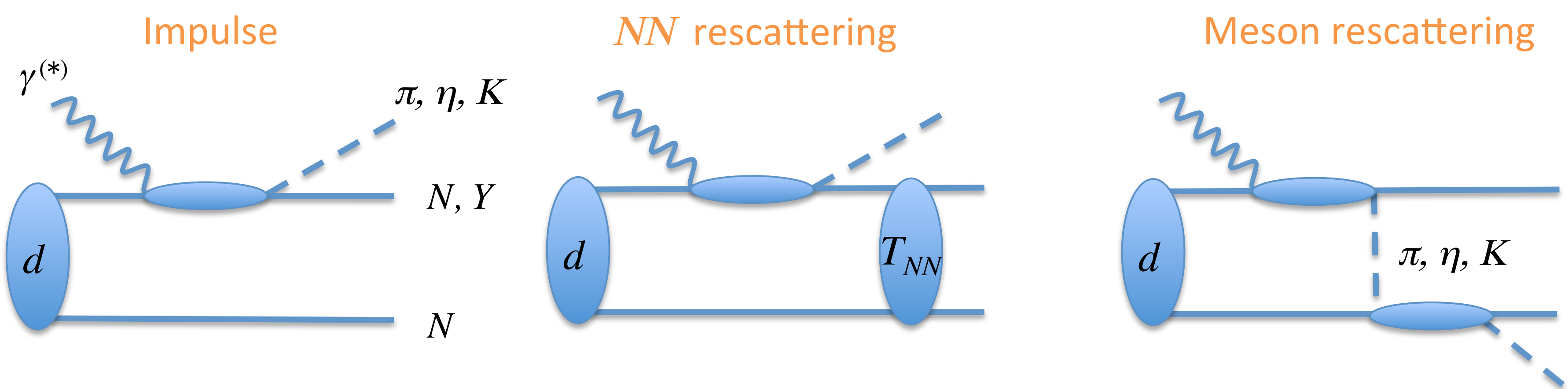


Photon and electron-induced meson productions off deuteron

To obtain $\gamma^{(*)}$ $'n' \rightarrow \pi N, \eta N, K Y$ data, Deuteron is the primary target; we need to understand:

How to extract $\gamma^{(*)}$ $'n' \rightarrow MB$ cross sections from $\gamma^{(*)} d \rightarrow MBN$ data

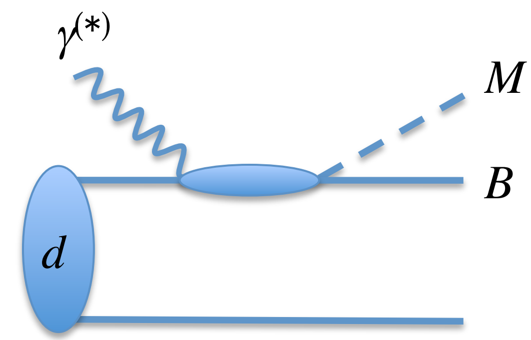
To address this question theoretically, we first need understand meson productions on deuteron



Features

- Photo- and electro-excitations of bound nucleons
- initial nucleons are in Fermi motion
- Final state interactions (FSI)

Q : How to extract $\sigma(\gamma^{(*)}\ 'n' \rightarrow MB)$ from $\sigma(\gamma^{(*)}\ d \rightarrow MBN)$?



Common (and practical) procedure

Kinematical cuts ➔ quasi-free $\gamma^{(*)}\ 'n' \rightarrow MB$ events selected

Concerns : FSI and/or cuts could distort $\sigma(\gamma^{(*)}\ 'n' \rightarrow MB)$ from true one ?

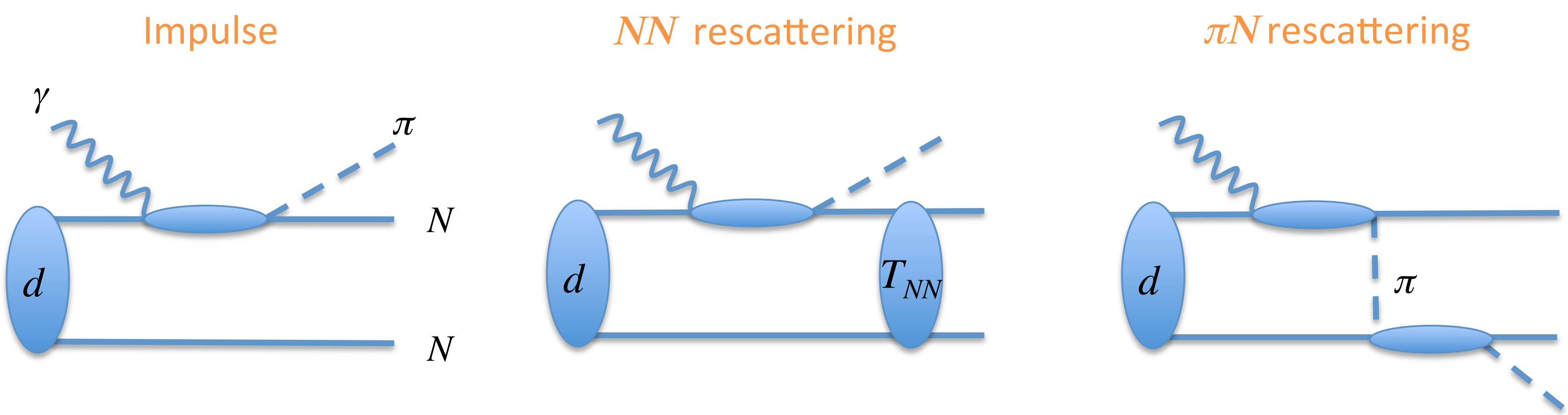
This work

- A dynamical model for $\gamma\, d$ ➔ πNN with FSI is developed
- $d\sigma/d\Omega_\pi$ and Σ, E, G for $\gamma\, d$ ➔ πNN are calculated; FSI effects examined
- $d\sigma/d\Omega_\pi$ and Σ, E, G for $\gamma\ 'n'$ ➔ $\pi^0 n, \pi^- p$ are extracted from $\gamma\, d$ ➔ πNN
- We address how FSI and cuts could distort $\gamma\ 'n'$ observables

$\gamma d \rightarrow \pi NN$ reaction model based on dynamical coupled-channels model

Model for $\gamma\, d \rightarrow \pi N N$

Multiple scattering theory truncated at the first-order rescattering



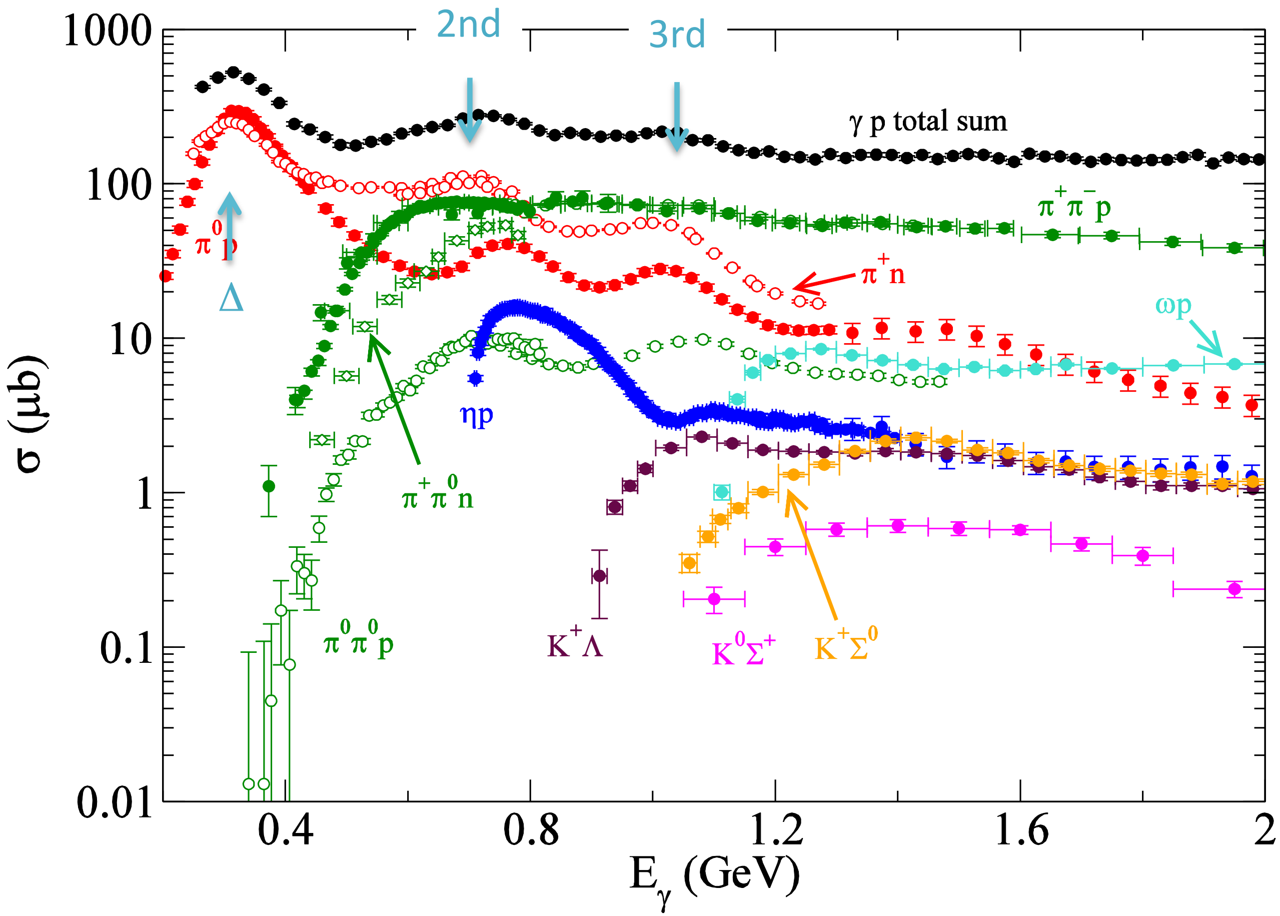
Elementary amplitudes

$\gamma N \rightarrow \pi N,\quad \pi N \rightarrow \pi N$ amplitude ← DCC model (Kamano et al., PRC94 (2016))

T_{NN}, deuteron w.f. ← CD-Bonn potential (Machleidt et al., PRC 63 (2001))

3-dim. loop integral with off-shell amplitudes are numerically evaluated

ANL-Osaka Dynamical coupled-channels model for meson productions in resonance region



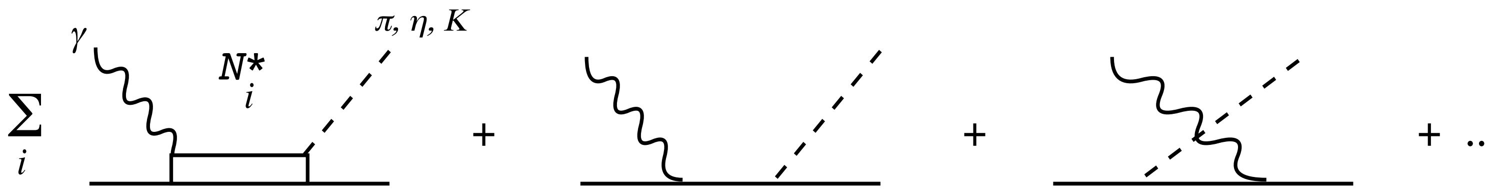
Data for $\gamma p \rightarrow X$

Need develop a model to describe these reactions in a unified manner

- Several nucleon resonances form characteristic peaks
- 2π production is comparable to 1π
- η, K productions (multi-channel couplings are important physics)

Dynamical coupled-channels model for resonance region

Resonance excitation + non-resonant meson-exchange mechanisms



Theoretically sound model should also account for:

- Channel-couplings required by unitarity (πN, ηN, $K\Lambda$, $K\Sigma$ stable channels)
- 2 π production mechanisms (ρN, σN , $\pi\Delta$ $\leftarrow\rightarrow$ $\pi\pi N$ channels)

Dynamical Coupled-Channels (DCC) model accounts for these features

developed through analyzing data for γN, πN $\rightarrow$ πN, ηN, $K\Lambda$, $K\Sigma$

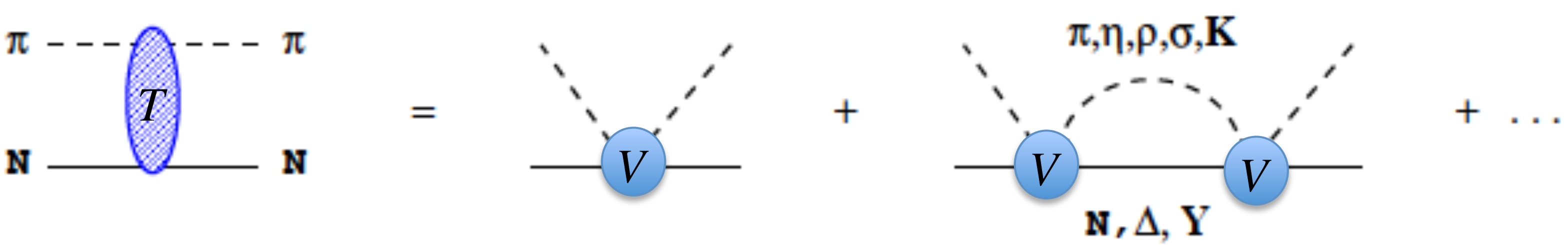
$\sim$ 26,000 data points

DCC (Dynamical Coupled-Channels) model

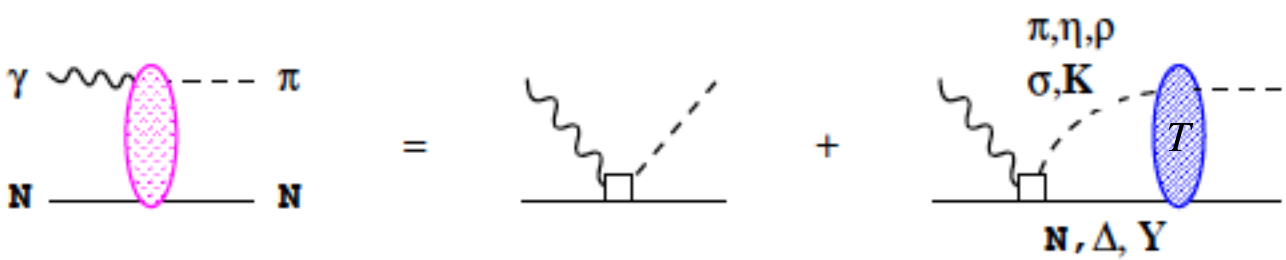
Matsuyama et al., Phys. Rept. 439, 193 (2007)
Kamano et al., PRC 88, 035209 (2013)

Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_{c} V_{ac} G_c T_{cb} \qquad \{a,b,c\} = \{\pi N, \eta N, \pi\pi N, \pi\Delta, \sigma N, \rho N, K\Lambda, K\Sigma\}$$



In addition, γN channel is included perturbatively

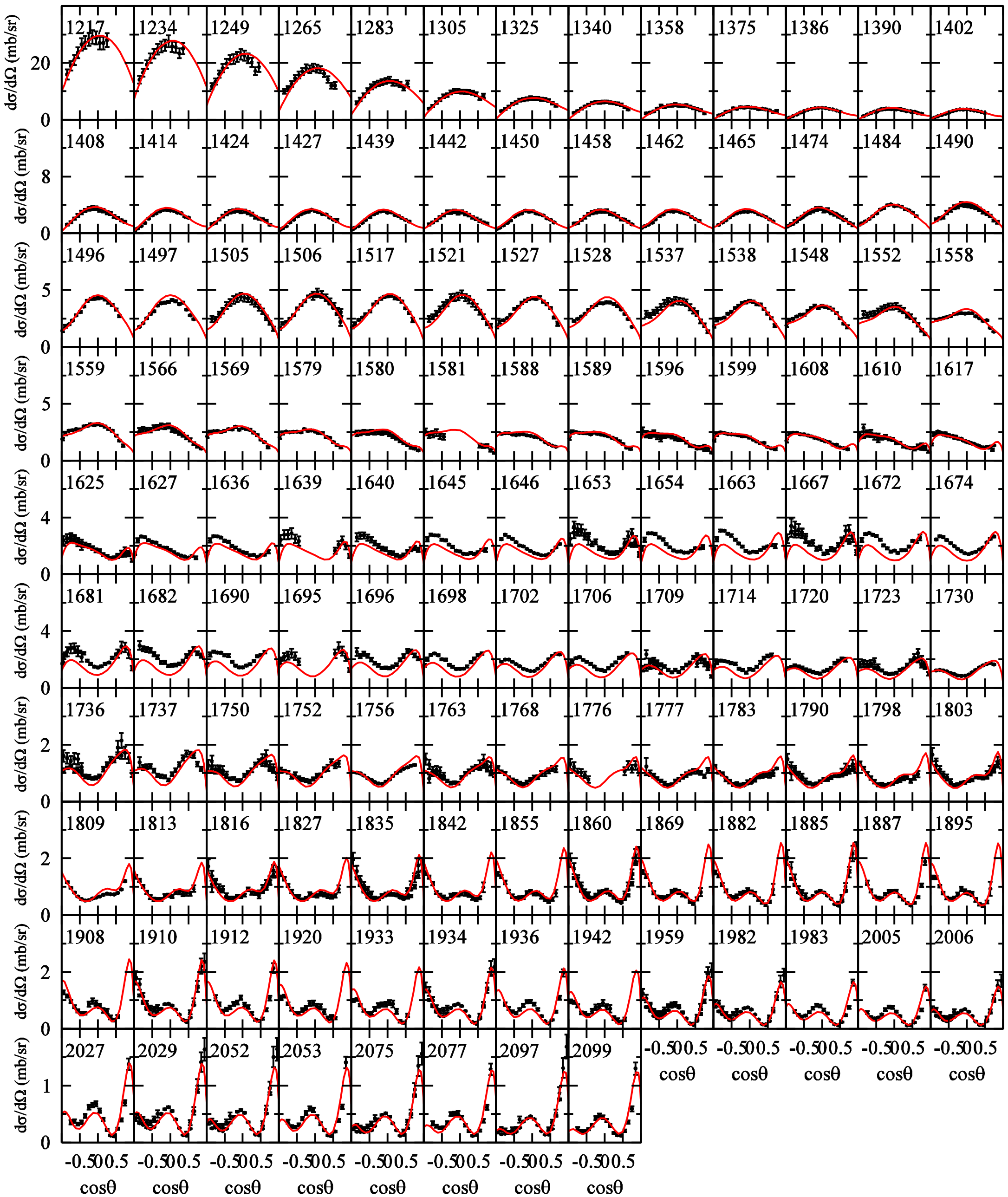


γp → π^0p

dσ/dΩ for W < 2.1 GeV

Comparison of DCC model with data

Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

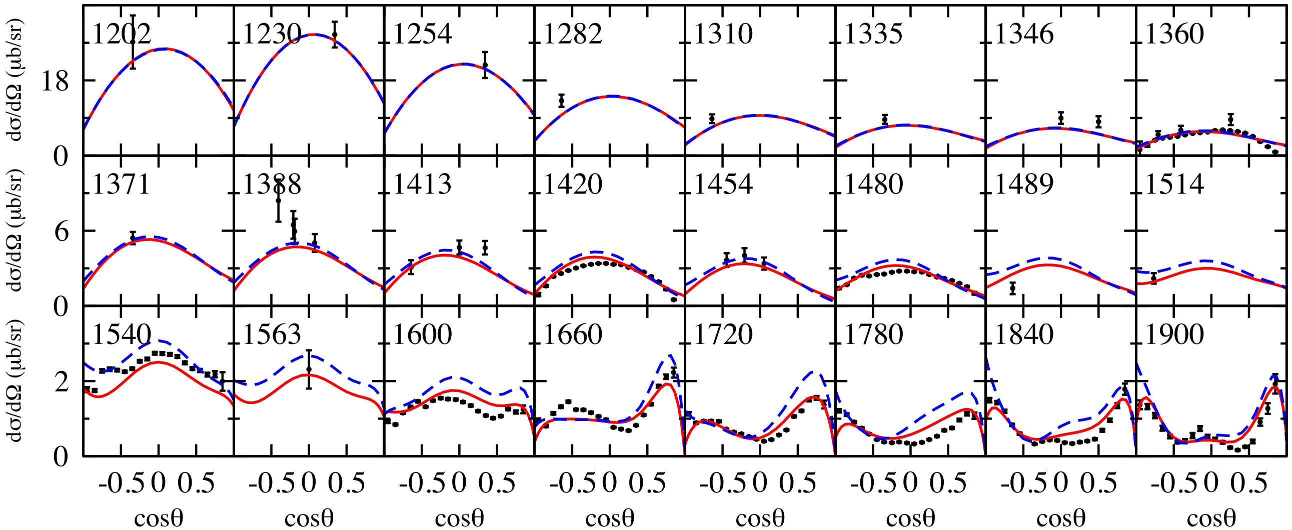


Reasonable fit to data

in the whole resonance region

γn → π^0n **dσ/dΩ for W < 2 GeV**

SXN, Kamano, Lee, Sato, arXiv:1804.04757

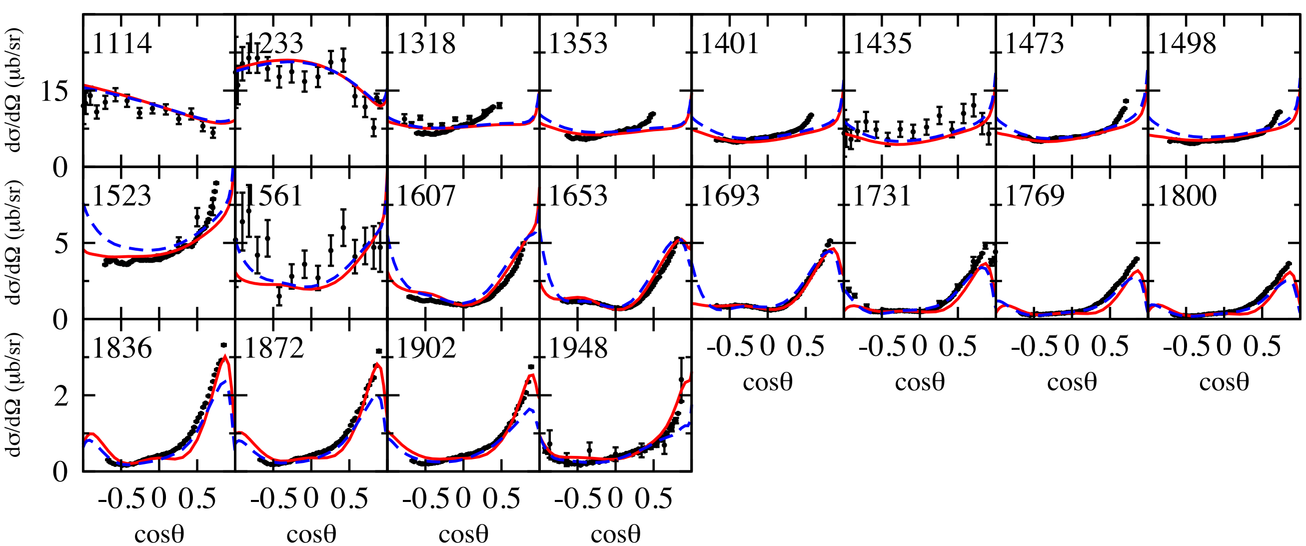


: new fit

: previous fit
PRC 88 (2013)

Recent MAMI data included
PRL 112, 142001 (2014)

γn → π^-p

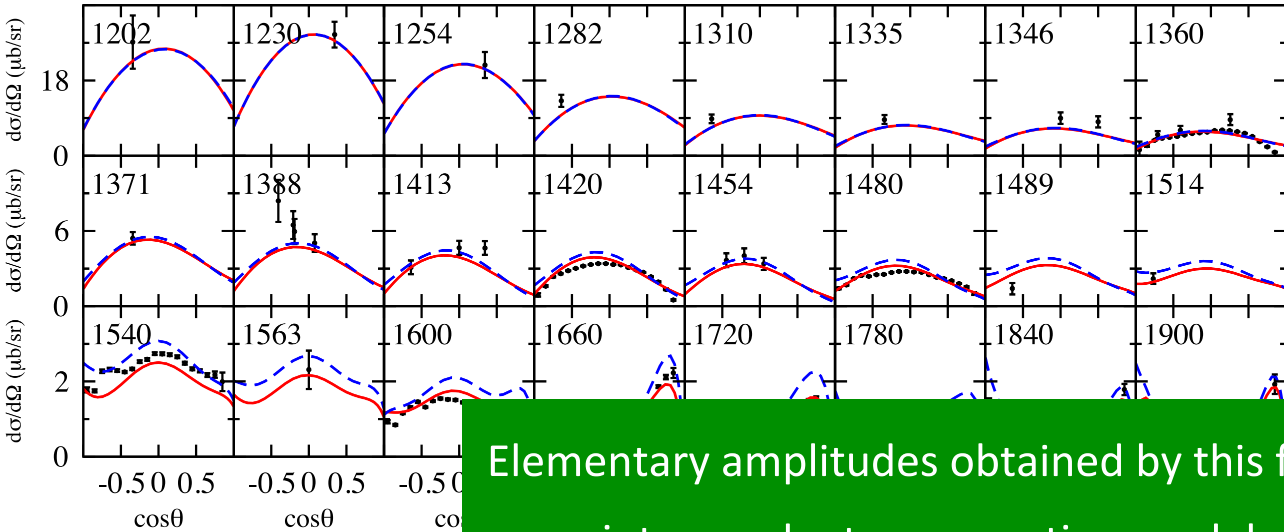


Recent JLab data included
PRC 96, 035204 (2017)

γn → π^0n

dσ/dΩ for W < 2 GeV

SXN, Kamano, Lee, Sato, arXiv:1804.04757



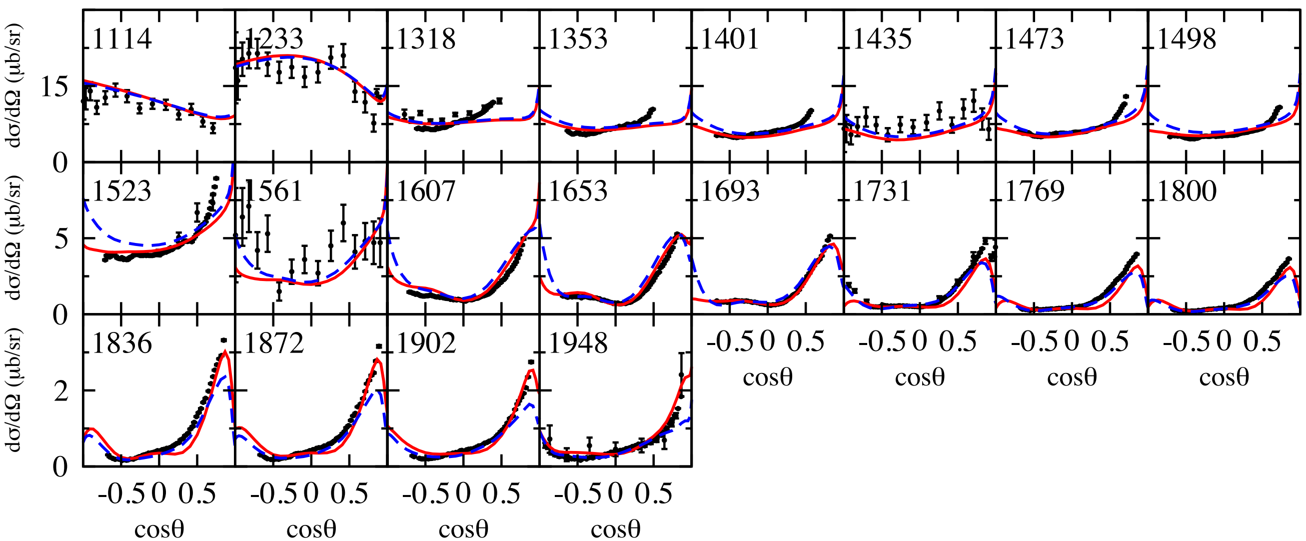
: new fit

: previous fit PRC 88 (2013)

Recent MAMI data included PRL 112, 142001 (2014)

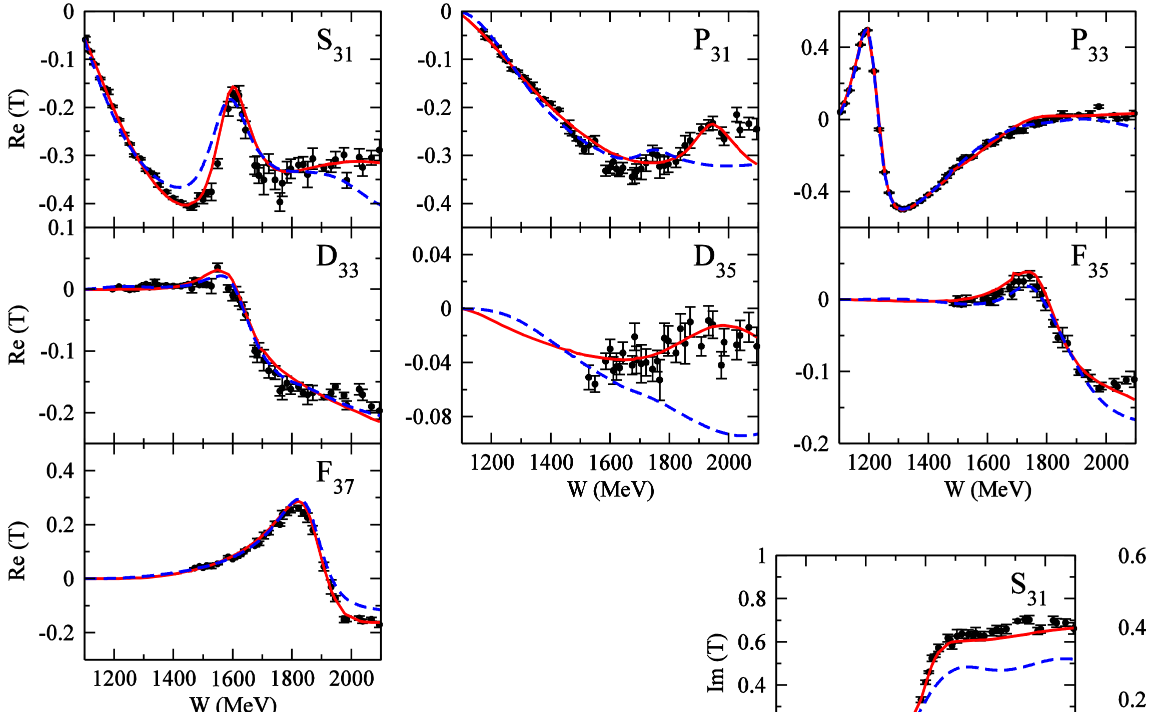
Elementary amplitudes obtained by this fit go into our deuteron reaction model

γn → π^-p



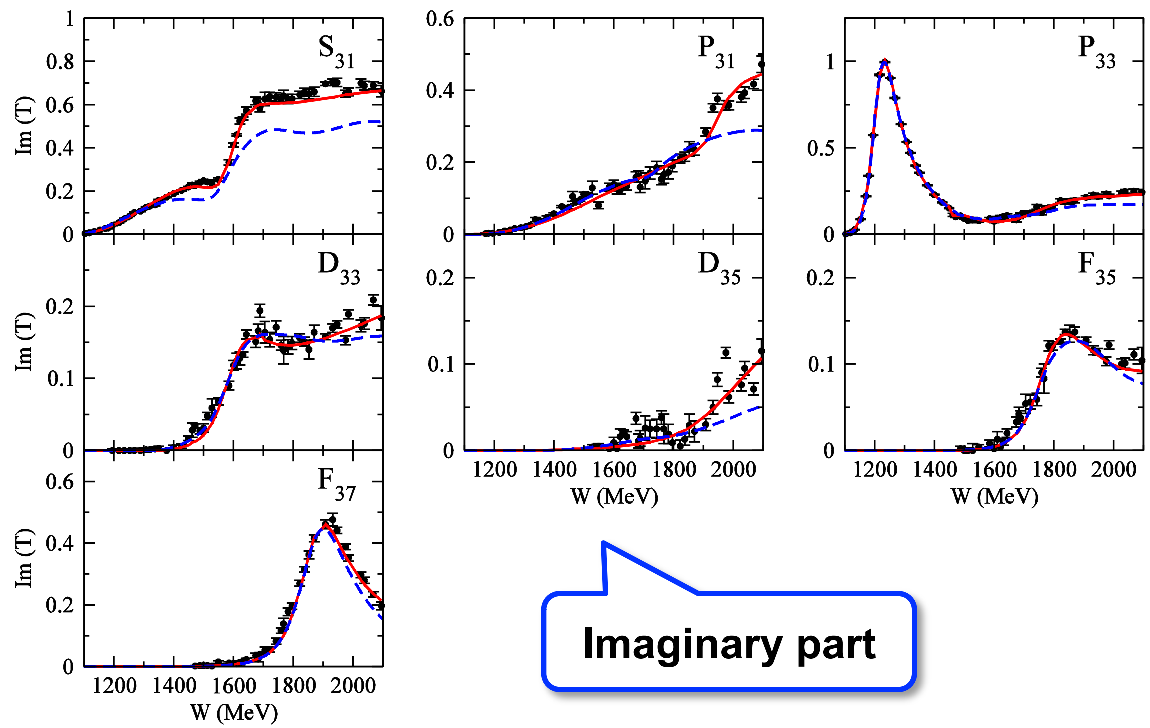
Recent JLab data included PRC 96, 035204 (2017)

Partial wave amplitudes of π N scattering



Real part

$$I = \frac{3}{2}$$



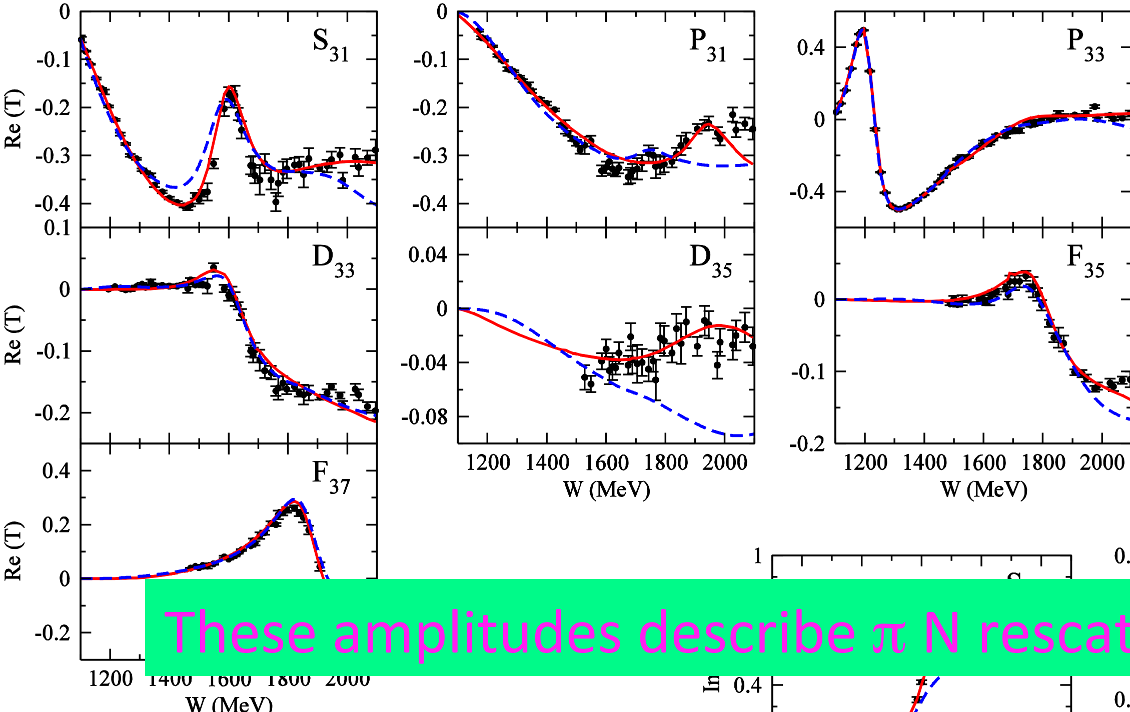
Imaginary part

Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Previous model
(fitted to πN → πN data only)
[PRC76 065201 (2007)]

Data: SAID πN amplitude

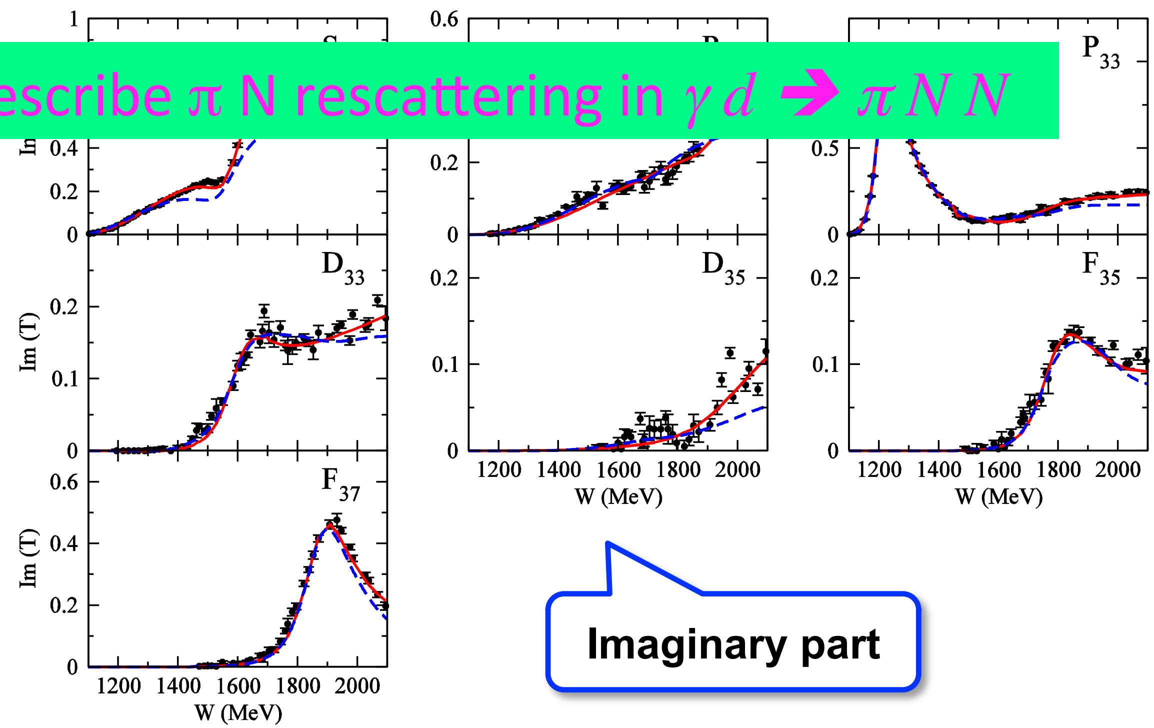
Partial wave amplitudes of π N scattering



Real part

$$I = \frac{3}{2}$$

These amplitudes describe π N rescattering in $\gamma\, d \rightarrow \pi\, N\, N$



Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Previous model
(fitted to πN → πN data only)
[PRC76 065201 (2007)]

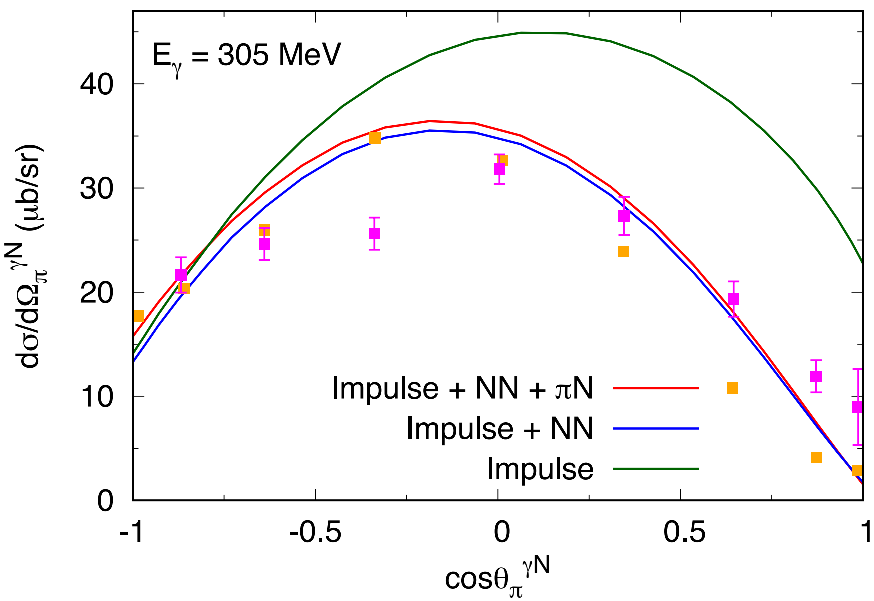
Imaginary part

Data: SAID πN amplitude

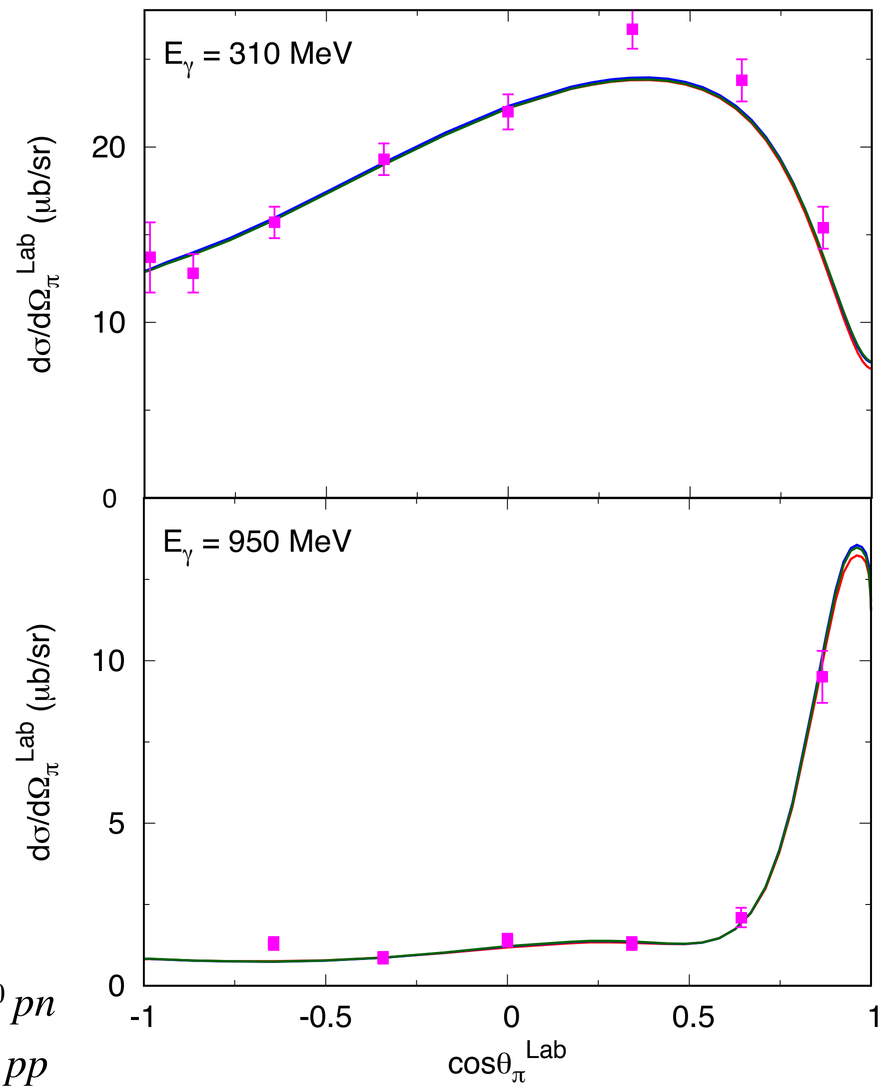
Results for $\gamma d \rightarrow \pi N N$

$\gamma\, d$ ➔ $\pi\, N\, N$: model predictions and data

$\gamma\, d$ ➔ $\pi^0\, pn$



$\gamma\, d$ ➔ $\pi^-\, pp$



- Large NN FSI effect for π^0 productions
  ← NN and deuteron wave fn. are orthogonal
- FSI effects are small for π^- productions
- Reasonable agreement with data

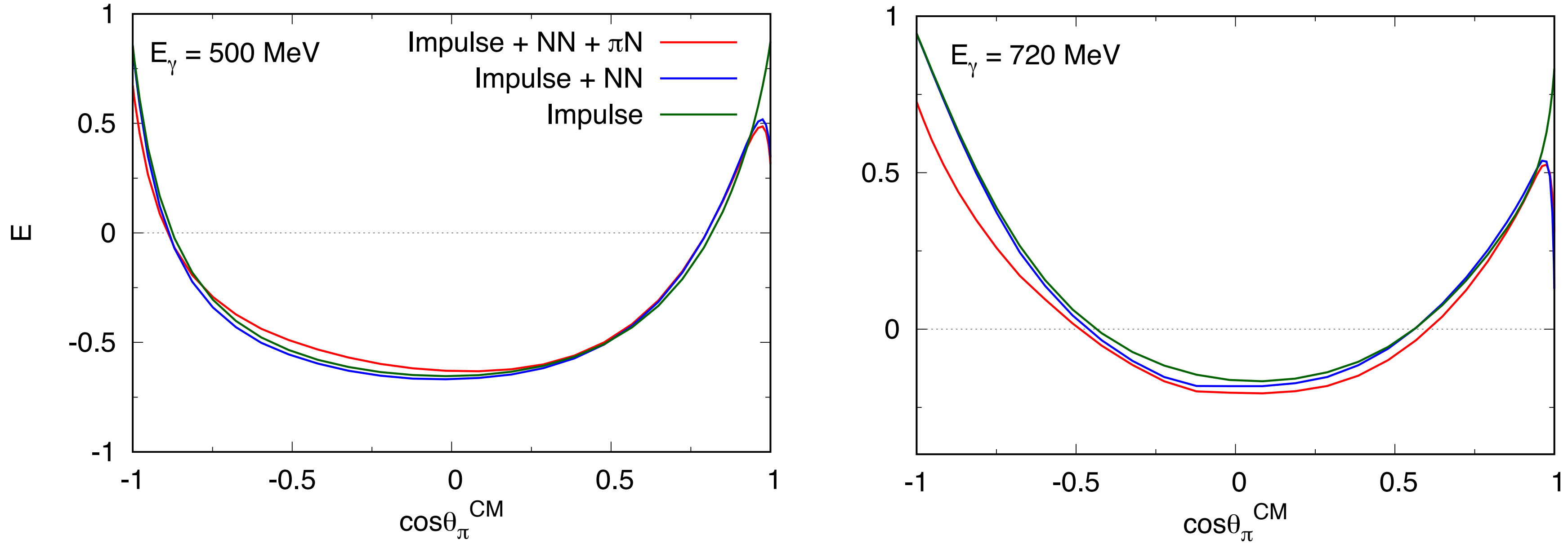
Data: EPJA 6, 309 (1999); 10, 365 (2001) for $\gamma\, d$ ➔ $\pi^0\, pn$

NPB 65, 158 (1973) for $\gamma\, d$ ➔ $\pi^-\, pp$

Model predictions for E of $\gamma\, d \rightarrow \pi N N$

Polarization asymmetry E : $$E \equiv \frac{\sigma_{+-} - \sigma_{++}}{\sigma_{+-} + \sigma_{++}} \qquad \sigma_{+\pm} \equiv \sigma(s_\gamma^z = +1, s_d^z = \pm 1)$$

$\gamma\, d \rightarrow \pi^0 pn$



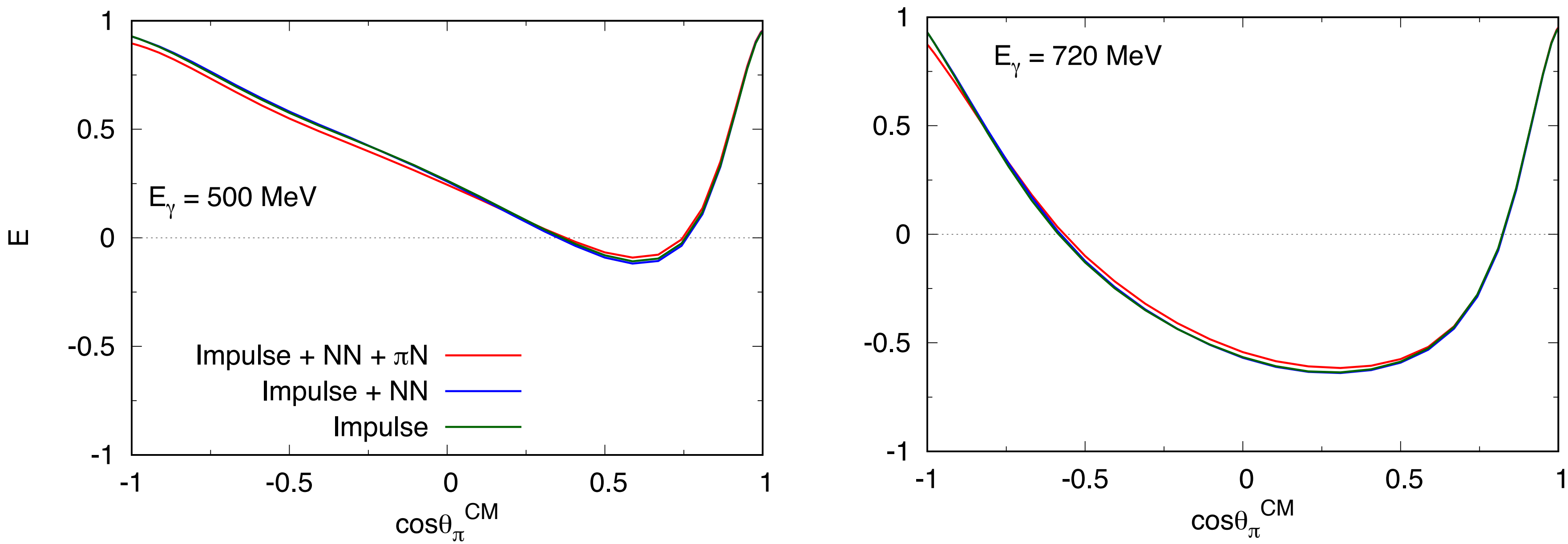
- FSI effects are smaller than $d\sigma/d\Omega_\pi$, but still visible

  → not completely cancelled out in the ratio

Model predictions for E of $\gamma\, d \rightarrow \pi\, N\, N$

Polarization observable E : $$E \equiv \frac{\sigma_{+-} - \sigma_{++}}{\sigma_{+-} + \sigma_{++}} \qquad \sigma_{+\pm} \equiv \sigma(s_\gamma^z = +1, s_d^z = \pm 1)$$

$\gamma\, d \rightarrow \pi^- pp$



- FSI effects are small as the same for $d\sigma/d\Omega_\pi$

Quasi-free γ 'n' cross sections extracted from γd cross sections

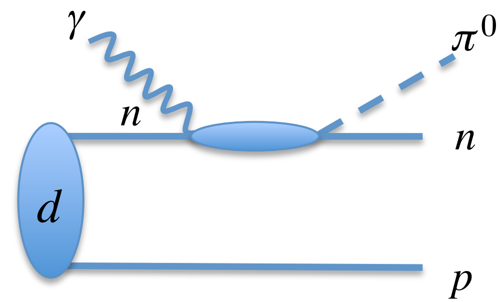
How to extract $\sigma(\gamma\,\text{'}n\text{'} \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$

First we need to establish a formula to relate $\sigma(\gamma\,\text{'}n\text{'} \rightarrow \pi^0 n)$ and $\sigma(\gamma d \rightarrow \pi^0 np)$

Ideal situation : only this mechanism contributes →

No FSI, No exchange terms, No deuteron D-wave

"quasi-free mechanism"



↓

$$\frac{d^2\sigma(\gamma d \rightarrow \pi^0 np; E_\gamma)}{dW\, d\cos\theta^*_\pi} = \phi(W)\, \frac{\tilde{E}_\gamma}{E_\gamma}\, \frac{d\sigma(\gamma n \rightarrow \pi^0 n; W)}{d\cos\theta^*_\pi}$$

E_γ ($\tilde{E}_\gamma$) : γ–d (γ–'n') Lab frame

W : $\pi^0 n$ invariant mass $\cos\theta^*_\pi \equiv \hat{q}_\gamma \cdot \hat{k}_\pi$ in π^0-n CM frame

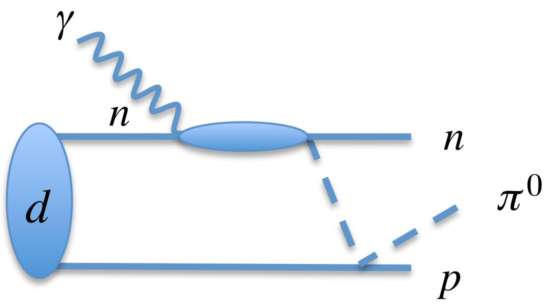
$\phi(W)$: Probability of a photon of E_γ interacting with the nucleon with W

$$\phi(W) \equiv \int dp^3_p\, \delta(W - W(\vec{p}_p, E_\gamma))\, \left|\psi_d(\vec{p}_p)\right|^2 , \qquad W(\vec{p}_p, E_\gamma) \equiv \sqrt{(E_\gamma + m_d - E_p(\vec{p}_p))^2 - (\vec{q}_\gamma - \vec{p}_p)^2}$$

How to extract $\sigma(\gamma\,'n' \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$

In reality, photon hits the other nucleon and FSI contribute

→ kinematical cuts are applied to remove them



Counterpart from our γd model : $\left.\dfrac{d^2\sigma(\gamma d \to \pi^0 np)}{dW\, d\cos\theta^*_\pi}\right|_{\text{cuts}}$

Assuming that this is solely from quasi-free contribution integrated over the same phase-space

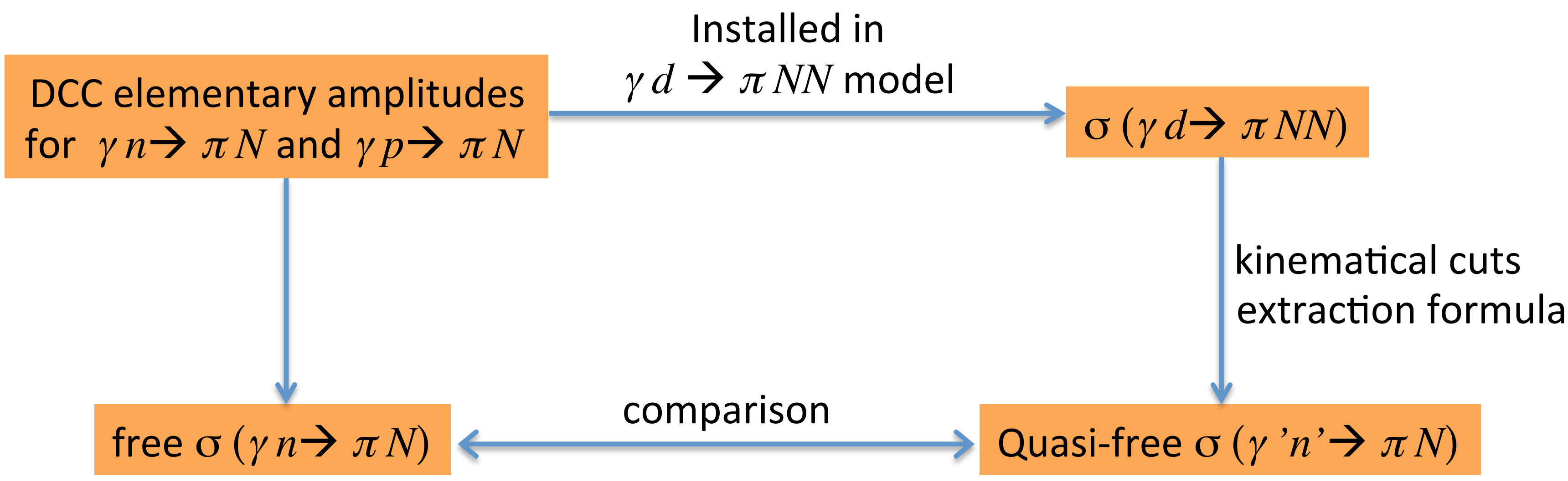
$$\left.\frac{d^2\sigma(\gamma d \to \pi^0 np)}{dW\, d\cos\theta^*_\pi}\right|_{\text{cuts}} = \phi_{\text{mod}}(W)\,\frac{\tilde{E}_\gamma}{E_\gamma}\,\frac{d\sigma(\gamma n \to \pi^0 n)}{d\cos\theta^*_\pi}\,, \qquad \phi_{\text{mod}}(W) \equiv \int^{\text{cuts}} dp_p^3\; \delta(W - W(\vec{p}_p, E_\gamma))\,\left|\psi_d(\vec{p}_p)\right|^2$$

↑

Formula to extract $\sigma(\gamma\,'n' \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$ from either experiment or model

Examine within our model how much $\sigma(\gamma\,'n' \rightarrow \pi^0 n)$ deviates from 'free' $\sigma(\gamma n \rightarrow \pi^0 n)$ due to FSI

Q: How FSI and kinematical cuts could distort γ 'n' observables ($d\sigma/d\Omega$, Σ, E, G) ?



Kinematical cuts

For extracting γ 'n' → $\pi^- p$ from γd → $\pi^- pp$, we use :

[1] CLAS@JLab, PRC 86 (2012)
[2] CLAS@JLab, PRC 96 (2017)
[3] CLAS@JLab, PRL 118 (2017)

Ref.	[1] ($d\sigma/d\Omega$)	[2] ($d\sigma/d\Omega$)	[3] (E)
E_γ (MeV)	301 – 455	445 – 2510	700 - 2400
π^- momentum (MeV)	> 80	> 100	> 400
Faster proton momentum (MeV)	> 270	> 360	> 400
Slower proton momentum (MeV)	< 270	< 200	< 100
$\Delta\phi_\pi = \lvert \phi_\pi - \phi_{\text{faster proton}} \rvert$	–	–	$160^\circ < \Delta\phi < 200^\circ$

For extracting γ 'n' → $\pi^0 n$ from γd → $\pi^0 pn$, the same cuts are used after modification

"π^-" → "π^0", "Faster proton" → "neutron", "Slower proton" → "proton"

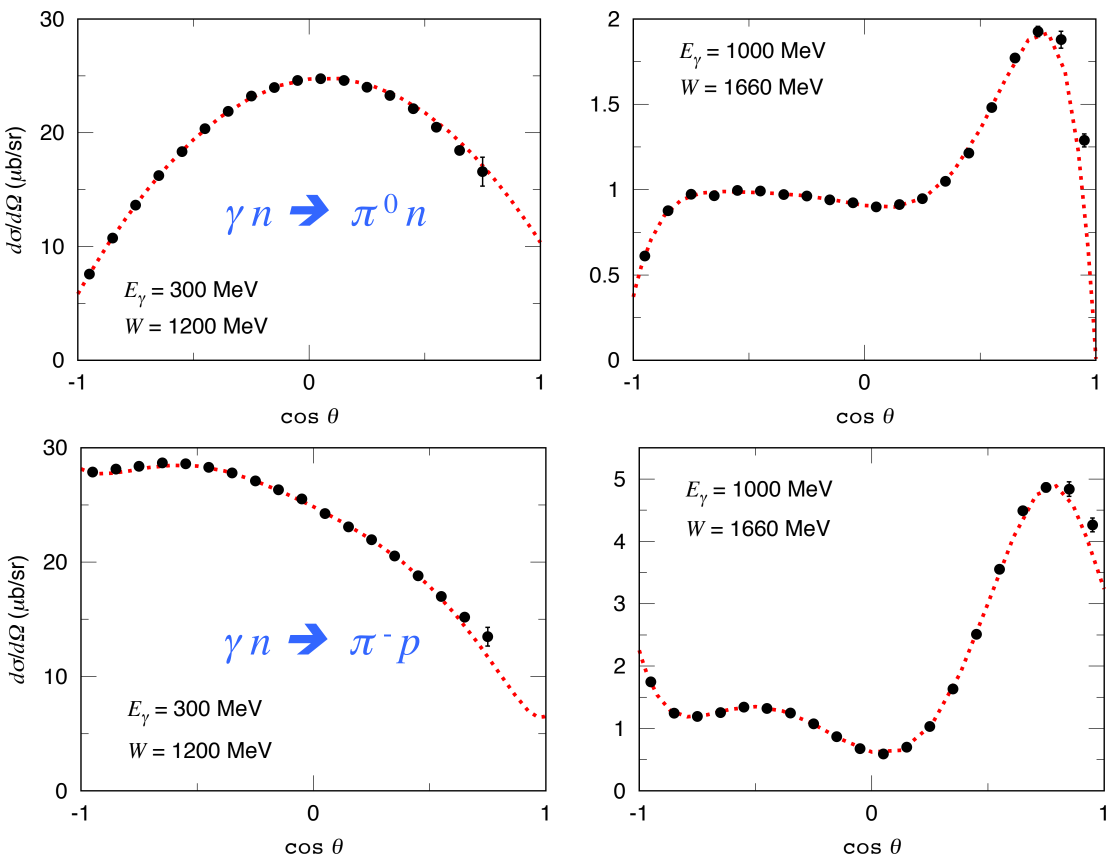
NOTE: Different kinematical cuts were used in A2@MAMI analysis on $\gamma n \to \pi^0 n$

Check the extraction method in ideal case

Q: With quasi-free mechanism only, examine if the formula $\left.\frac{d^2\sigma(\gamma d \rightarrow \pi^0 np)}{dW\, d\cos\theta_\pi^*}\right|_{\text{cuts}} = \phi_{\text{mod}}(W)\, \frac{\tilde{E}_\gamma}{E_\gamma}\, \frac{d\sigma(\gamma n \rightarrow \pi^0 n)}{d\cos\theta_\pi^*}$ returns $\sigma(\gamma\, 'n' \rightarrow \pi N)$ that agree with 'free' $\sigma(\gamma\, n \rightarrow \pi N)$

Check the extraction method in ideal case

Q: With quasi-free mechanism only, examine if the formula $\left.\frac{d^2\sigma(\gamma d \to \pi^0 np)}{dW\, d\cos\theta^*_\pi}\right|_{\text{cuts}} = \phi_{\text{mod}}(W)\,\frac{\tilde{E}_\gamma}{E_\gamma}\,\frac{d\sigma(\gamma n \to \pi^0 n)}{d\cos\theta^*_\pi}$ returns $\sigma(\gamma\, 'n' \to \pi N)$ that agree with 'free' $\sigma(\gamma n \to \pi N)$



● : extracted $\sigma(\gamma\ 'n' \to \pi N)$

······· : free $\sigma(\gamma n \to \pi N)$

Note:

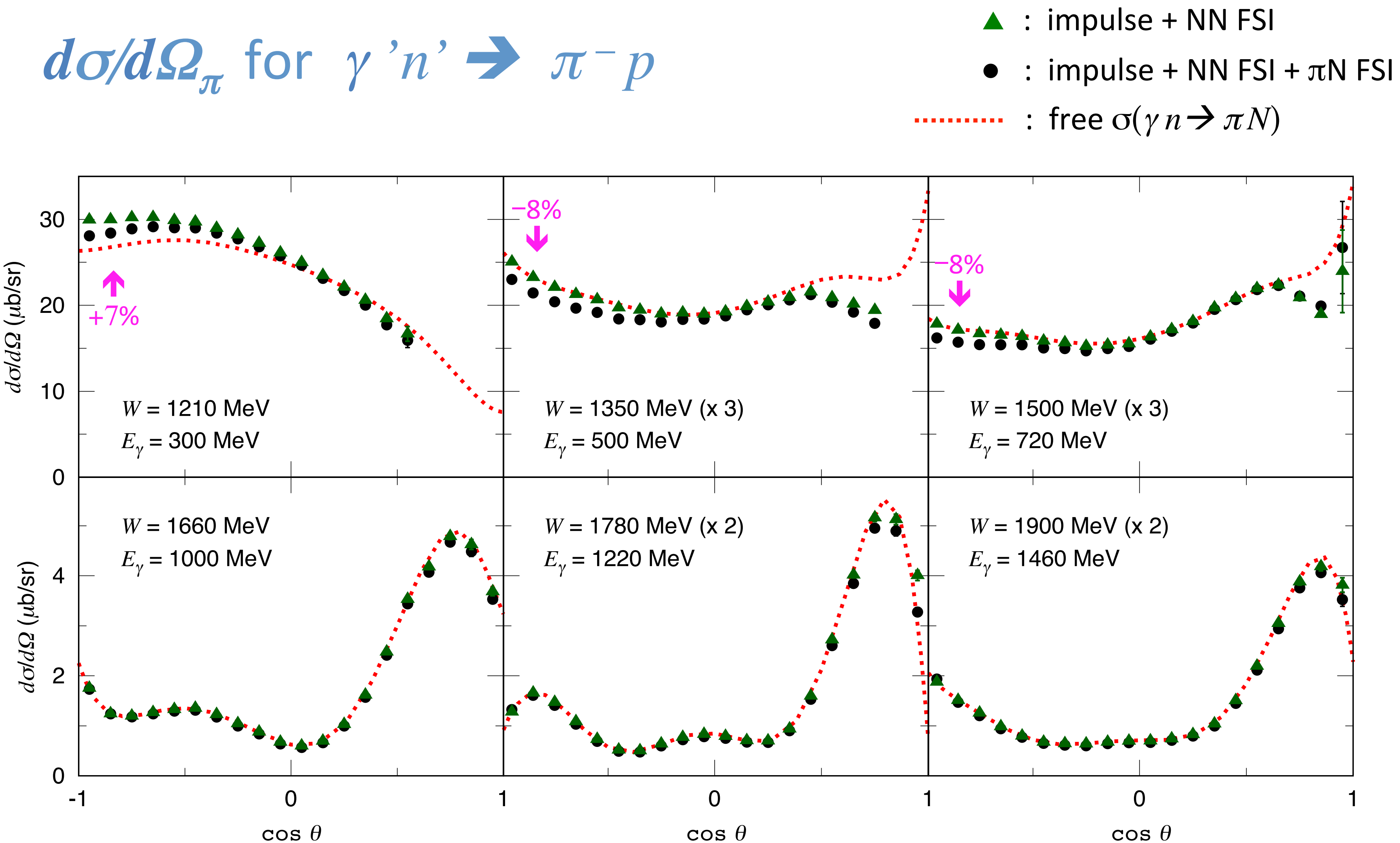
* Phase-space integrals for $\sigma(\gamma d)$ are done with Monte-Carlo method to easily implement any cuts; statistical error bars are very small

* Forward pion productions are invisible (left figures) due to the kinematical cuts

Free $\sigma(\gamma n \to \pi N)$ are well reproduced as expected for the ideal case

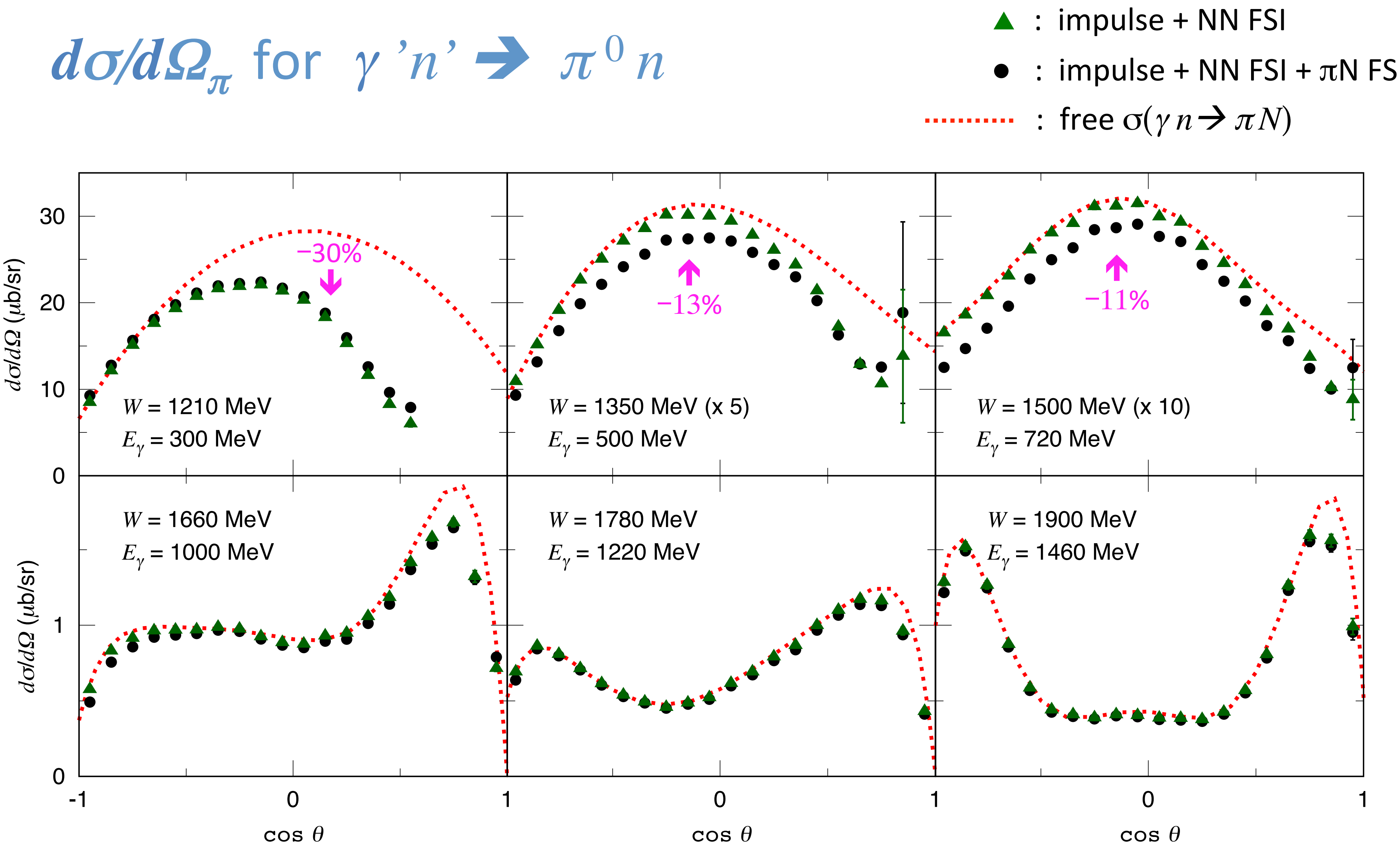
Now go on to realistic case !

$d\sigma/d\Omega_\pi$ for γ 'n' ➔ $\pi^- p$



- NN (π N) FSI reduce $d\sigma/d\Omega_\pi$ of forward (backward) pion
- kinematical cuts cannot remove the FSI effect
- Larger FSI effects for smaller E_γ

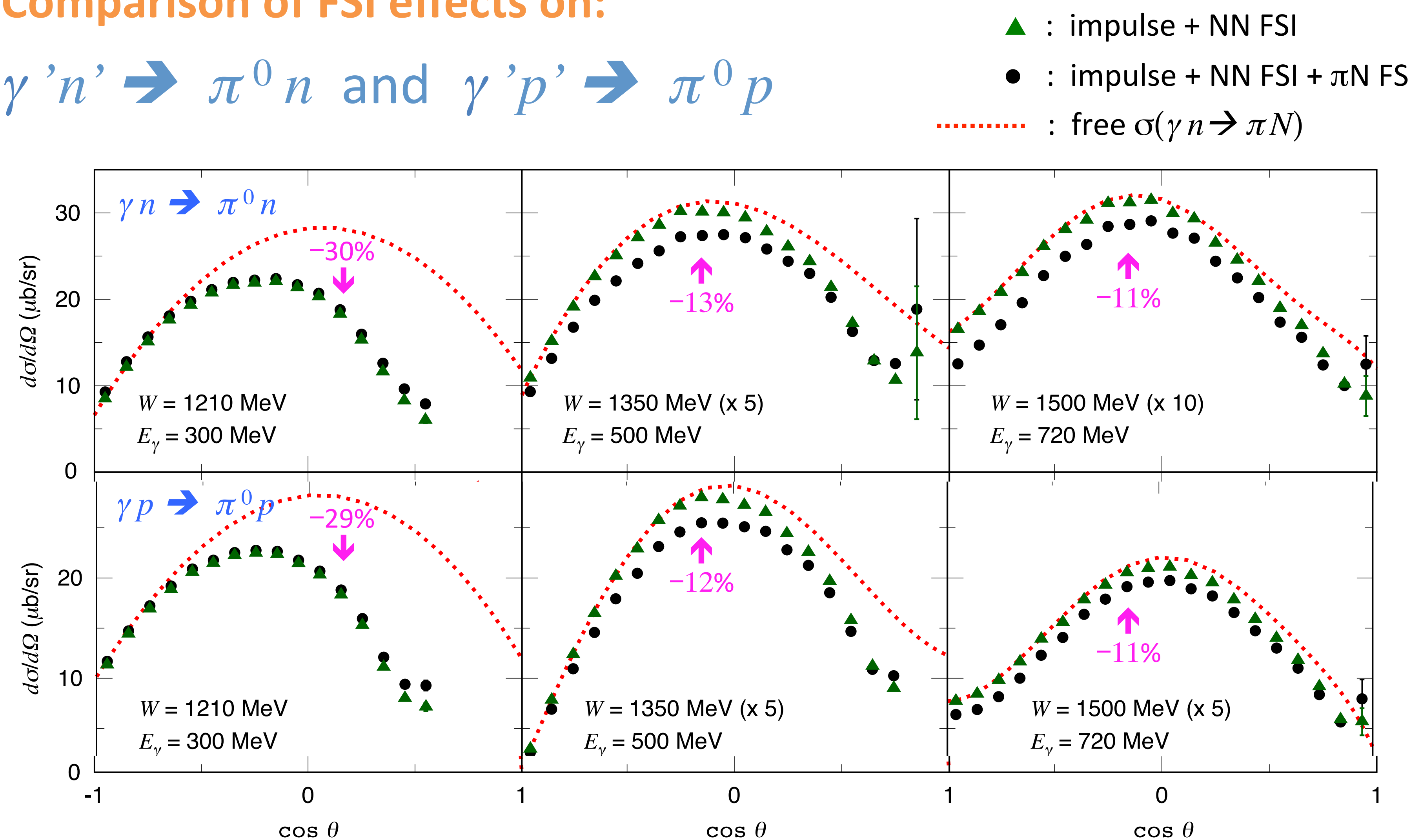
$d\sigma/d\Omega_\pi$ for γ 'n' ➔ $\pi^0 n$



- Significant FSI effects reduce $d\sigma/d\Omega_\pi$; πN and NN FSI are comparably important
- kinematical cuts cannot remove FSI effect
- Larger FSI effects for smaller E_γ

Comparison of FSI effects on:

$\gamma\,\text{'}n\text{'} \rightarrow \pi^0 n$ and $\gamma\,\text{'}p\text{'} \rightarrow \pi^0 p$

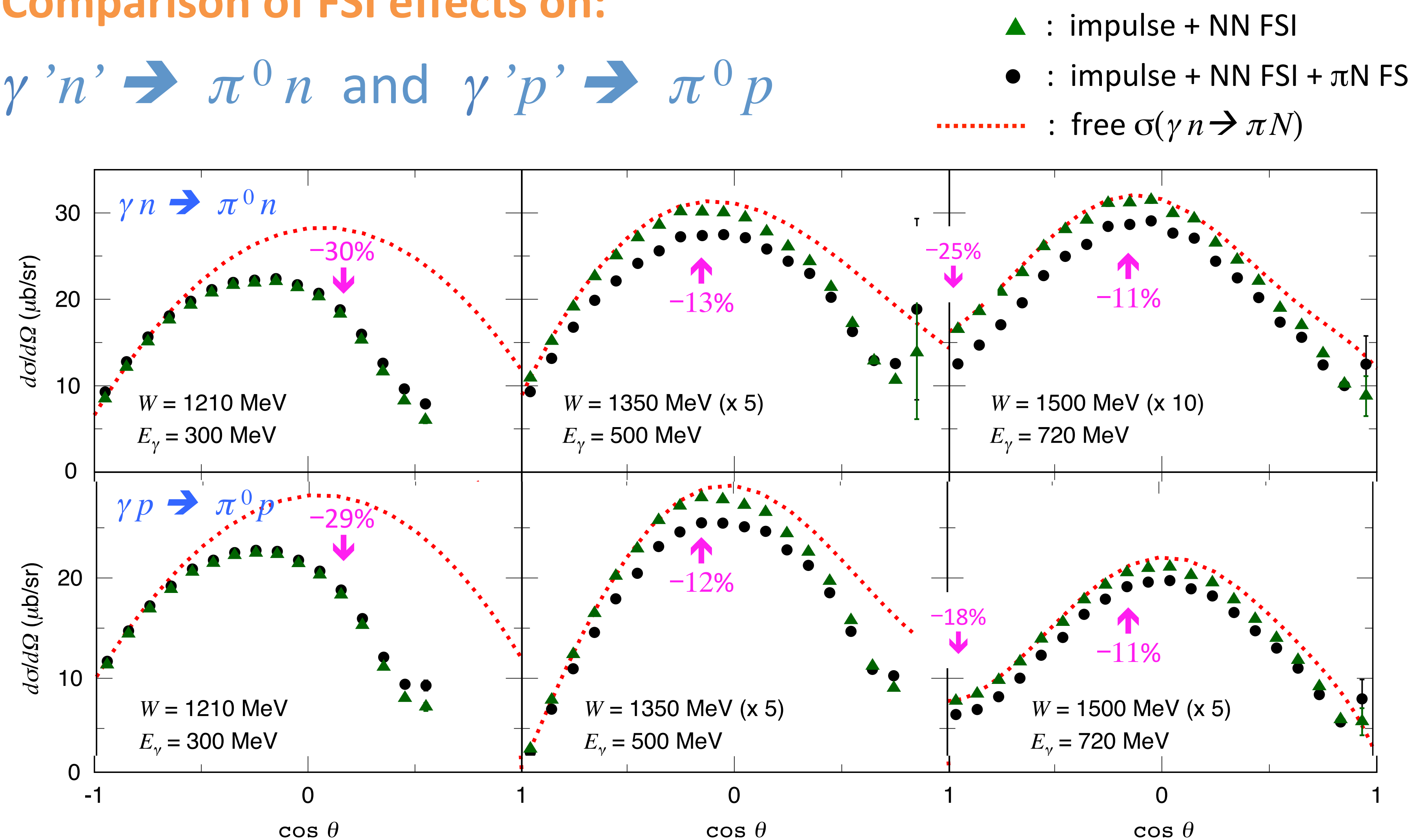


- FSI effects on $\gamma\,\text{'}n\text{'} \rightarrow \pi^0 n$ and $\gamma\,\text{'}p\text{'} \rightarrow \pi^0 p$ are generally similar (a few % difference)

  ⇔ Same FSI effects are assumed in A2@MAMI analysis [PRL 112 (2014)]
- But sometimes more different

Comparison of FSI effects on:

γ 'n' ➔ $\pi^0 n$ and γ 'p' ➔ $\pi^0 p$



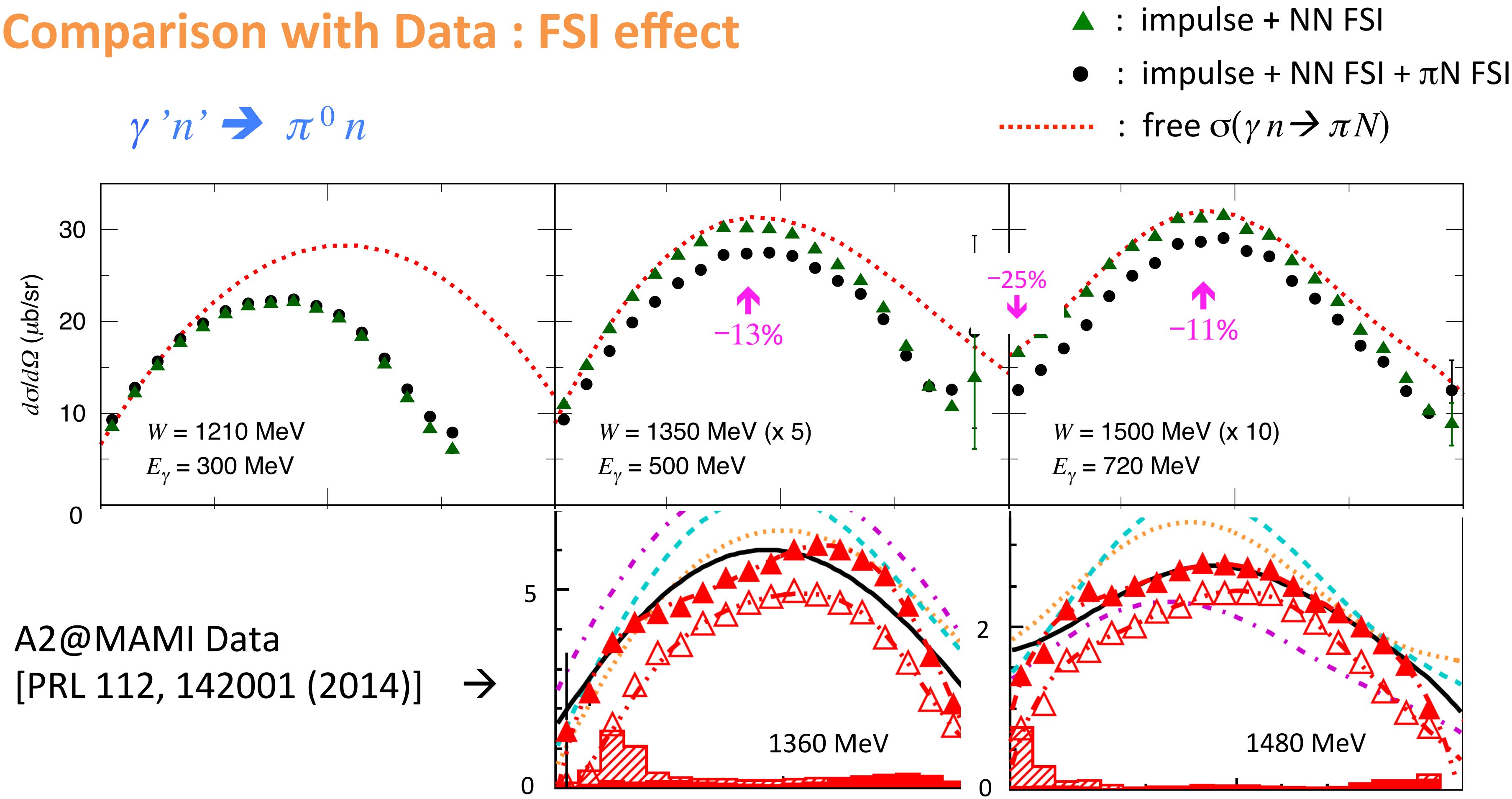
- FSI effects on γ 'n' ➔ $\pi^0 n$ and γ 'p' ➔ $\pi^0 p$ are generally similar (a few % difference)

  ⇔ Same FSI effects are assumed in A2@MAMI analysis [PRL 112 (2014)]

- But sometimes more different

Comparison with Data : FSI effect

γ 'n' → $\pi^0 n$



A2@MAMI Data
[PRL 112, 142001 (2014)] →

- Good agreement on the FSI effects estimated by A2@MAMI (first theoretical explanation)
- A2@MAMI analysis assumed $\frac{\sigma(\gamma' n' \to \pi^0 n)}{\text{free } \sigma(\gamma n \to \pi^0 n)} = \frac{\sigma(\gamma' p' \to \pi^0 p)}{\text{free } \sigma(\gamma p \to \pi^0 p)}$ no theoretical estimate
- FSI effects can depend on the cuts; MAMI analysis uses different cuts

(More) common extraction method

Extraction formula used so far in this presentation

$$\left.\frac{d^2\sigma(\gamma d \to \pi^0 np;\, E_\gamma)}{dW\, d\cos\theta_\pi^*}\right|_{\text{cuts}} = \phi_{\text{mod}}(W)\, \frac{\tilde{E}_\gamma}{E_\gamma}\, \frac{d\sigma(\gamma n \to \pi^0 n;\, W)}{d\cos\theta_\pi^*}, \quad \phi_{\text{mod}}(W) \equiv \int^{\text{cuts}} dp_p^3\ \delta(W - W(\vec{p}_p, E_\gamma))\, \left|\psi_d(\vec{p}_p)\right|^2$$

(More) common extraction method

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Integrate both sides with respect to W, assuming dominant contribution from nucleon-at-rest in deuteron

(and/or fairly weak W-dependence)

(More) commonly used extraction formula

$$\left.\frac{d^2\sigma(\gamma d \to \pi^0 np;\, E_\gamma)}{d\cos\theta^*_\pi}\right|_{\text{cuts}} = \frac{d\sigma(\gamma n \to \pi^0 n;\, \bar{W})}{d\cos\theta^*_\pi} \int^{\text{cuts}} dp_p^3\, \frac{\tilde{E}_\gamma}{E_\gamma}\, \left|\psi_d(\vec{p}_p)\right|^2, \qquad \bar{W}^2 = (E_\gamma + m_N)^2 - E_\gamma^2$$

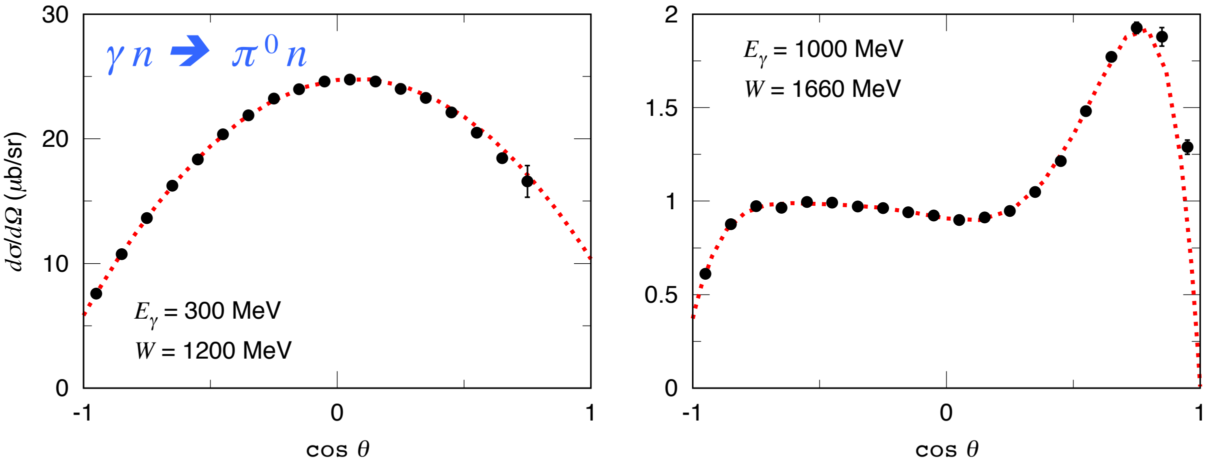
Used in recent CLAS@JLab analyses: PRC 96 (2017) for $d\sigma/d\Omega_\pi$

PRL 118 (2017) for E

also in theoretical analysis Tarasov et al., PRC 84, 035203 (2011)

Examine the validity of this formula without W-cut

(More) common extraction method



Quasi-free mechanism only;

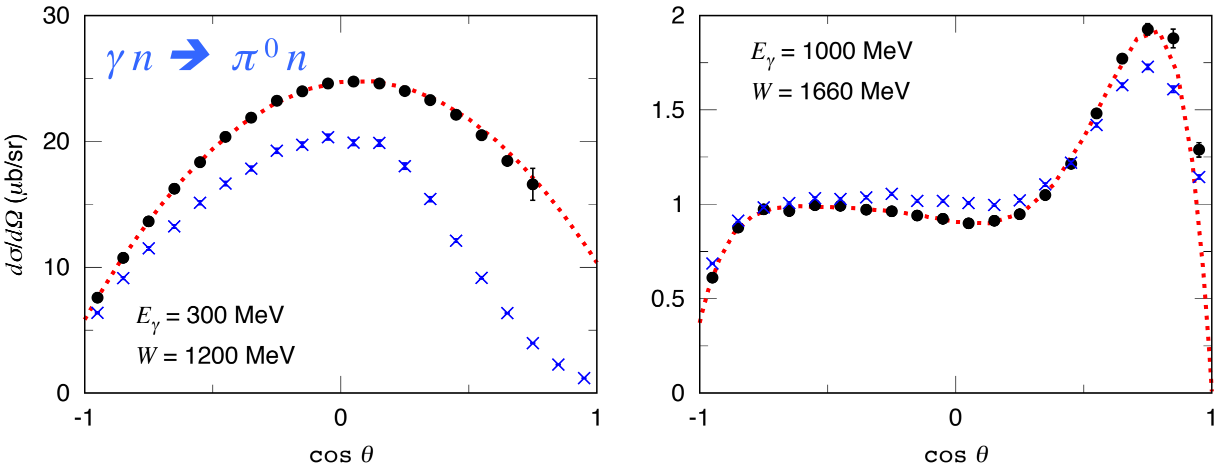
no FSI, no exchange terms

● : extracted with W-cut

x : without W-cut

⋯⋯ : free $\sigma(\gamma\, n \rightarrow \pi N)$

(More) common extraction method

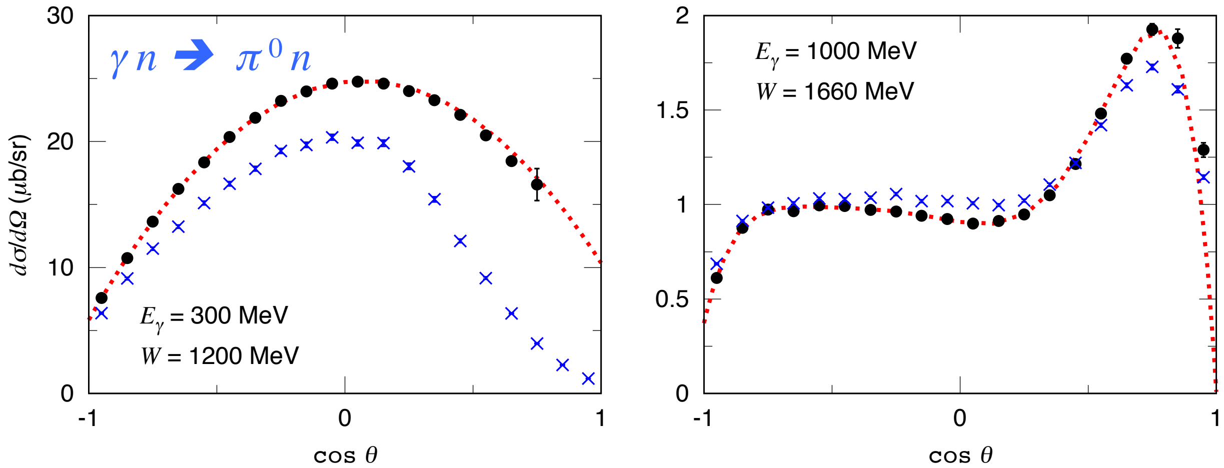


Quasi-free mechanism only;

no FSI, no exchange terms

- ● : extracted with W-cut
- x : without W-cut
- ....: free $\sigma(\gamma\, n \rightarrow \pi N)$

(More) common extraction method



Quasi-free mechanism only;

no FSI, no exchange terms

● : extracted with W-cut

x : without W-cut

⋯⋯: free $\sigma(\gamma n \rightarrow \pi N)$

- E_γ=300 MeV : Significant difference in $\Delta(1232)$-region

  1. Sharp peak of $\sigma_{\gamma n}$ at $W\sim1.2$ GeV
  2. $\sigma_{\gamma d}$ is an average of $\sigma_{\gamma n}$ over $W\sim1.18-1.21$ GeV ($W<1.2$ GeV in $\cos\theta \gtrsim 0.7$) because of Fermi motion

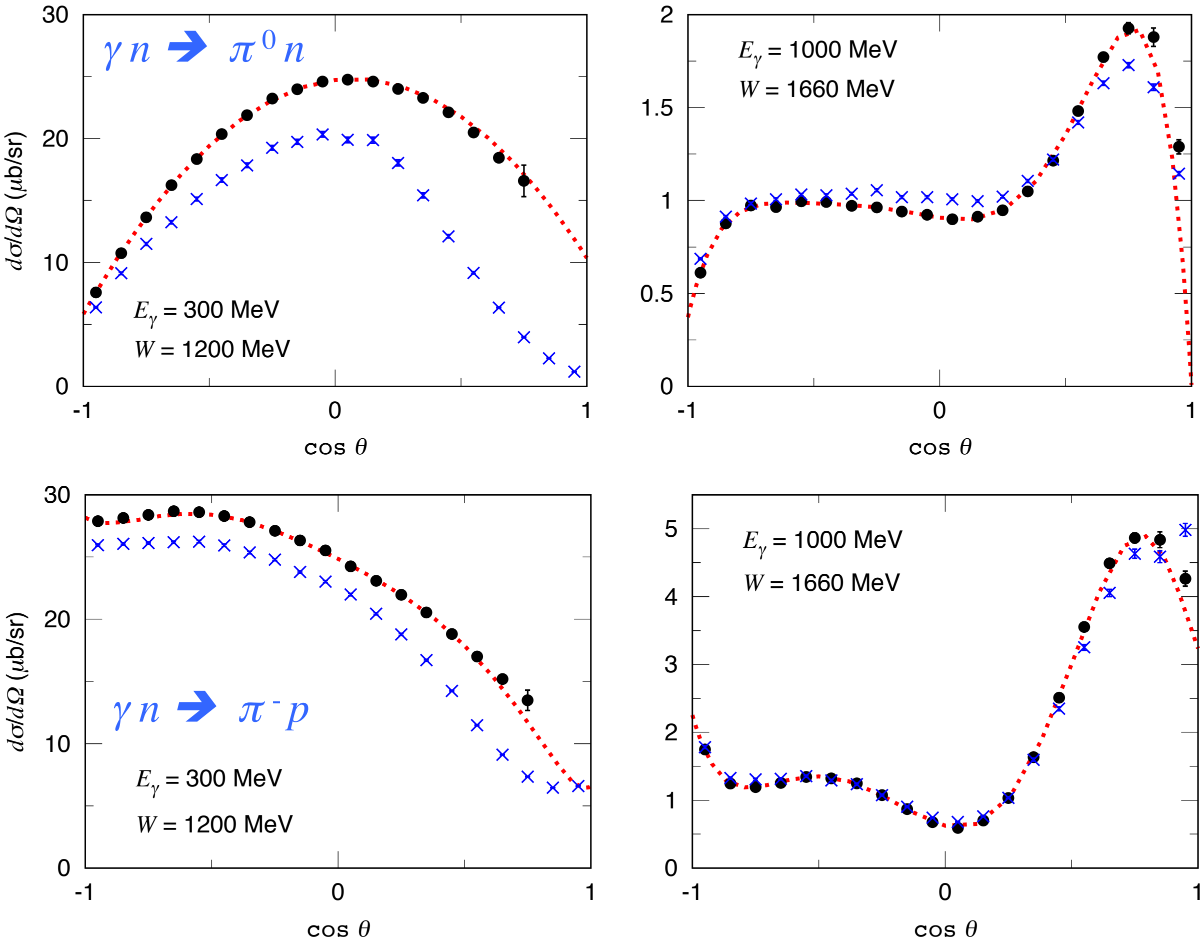
  → extracted $\sigma_{\gamma n}$ without W-cut is necessarily smaller than $\sigma_{\gamma n}$ at W=1.2 GeV

- E_γ=1GeV : better agreement between with and without W-cut

  Milder and monotonic W-dependence of $\sigma_{\gamma n}$ for $W\sim1.6-1.7$ GeV

  → W-dependence is canceled in $\sigma_{\gamma d}$ by W-average, but not completely

(More) common extraction method



Quasi-free mechanism only;

no FSI, no exchange terms

● : extracted with W-cut

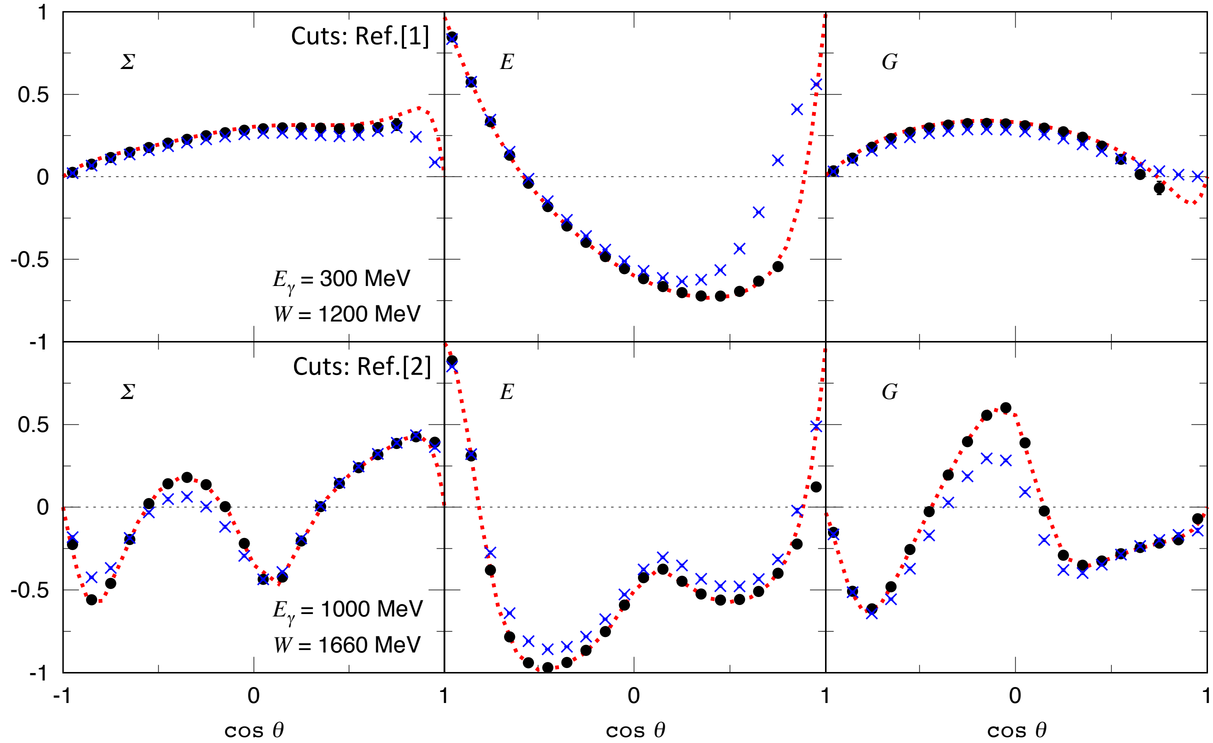
x : without W-cut

: free $\sigma(\gamma n \rightarrow \pi N)$

Similar result for $\gamma n \rightarrow \pi^- p$

(More) common extraction method

Polarization observables for $\gamma n \rightarrow \pi^- p$



Quasi-free mechanism only;

no FSI, no exchange terms

● : extracted with W-cut

x : without W-cut

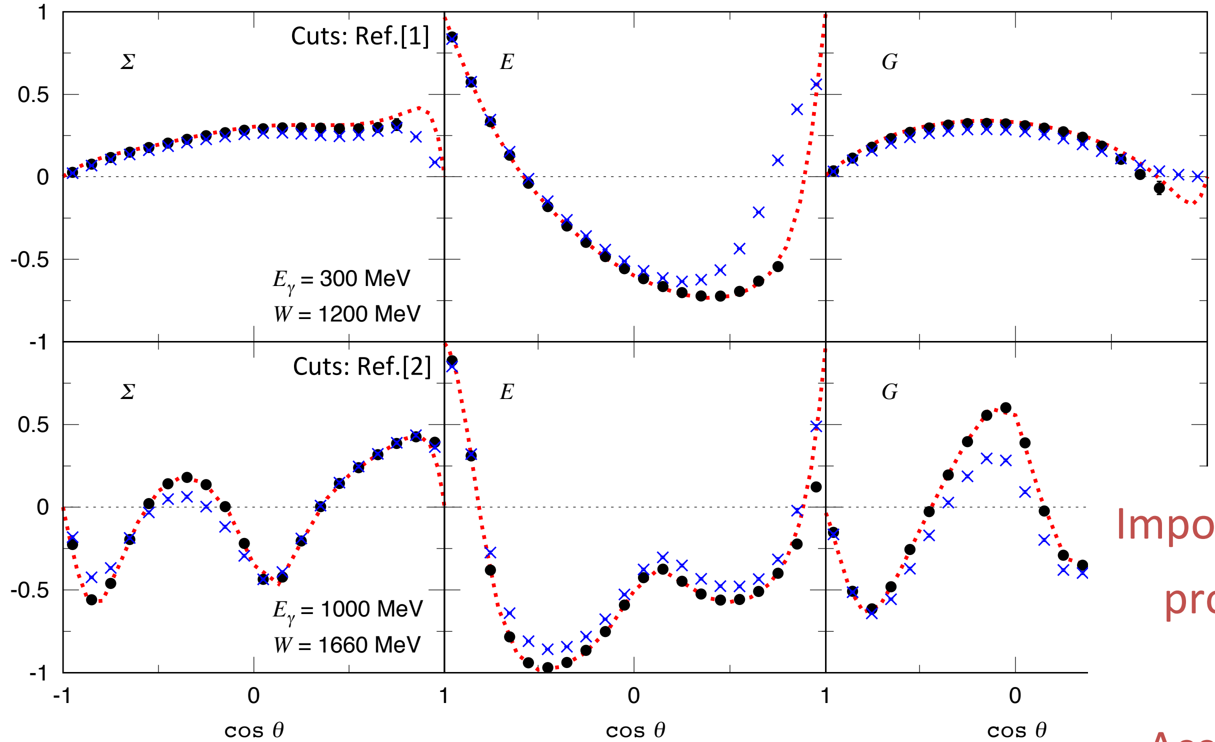
: free $\sigma(\gamma n \rightarrow \pi N)$

Reasons for deviations from free ones

- E at E_γ=300MeV, $\cos\theta \gtrsim 0.5$: nucleon-at-rest kinematics is largely restricted by cuts; contributions are from different W
- G at E_γ=1GeV : Average non-monotonic W-dependence → W-dependence not cancelled

(More) common extraction method

Polarization observables for $\gamma n \rightarrow \pi^- p$



Quasi-free mechanism only;

no FSI, no exchange terms

● : extracted with W-cut

x : without W-cut

······ : free $\sigma(\gamma n \rightarrow \pi N)$

Important to apply W-cut to suppress problematic Fermi motion effect

↓

Accuracy of extracted observables

Reasons for deviations from free ones

- E at E_γ=300MeV, $\cos\theta \gtrsim 0.5$: nucleon-at-rest kinematics is largely restricted by cuts; contributions are from different W
- G at E_γ=1GeV : Average non-monotonic W-dependence → W-dependence not cancelled

Conclusion

Conclusions

- $d\sigma/d\Omega_\pi$ and Σ, E, G for γ 'n' ➔ $\pi^0 n, \pi^- p$ are extracted from γd ➔ πNN using kinematical cuts of CLAS@JLab analyses
- $d\sigma/d\Omega_\pi$ for γ 'n' ➔ πN are significantly distorted by FSI even after the cuts;

  FSI effects on γ 'n' ➔ $\pi^0 n$ data of A2@MAMI explained with πN and NN FSI
- Σ, E, G for γ 'n' ➔ πN are mostly unaffected by FSI effects if W-cut is applied
- Without W-cut, extracted observables can be largely distorted by Fermi motion

Future plans

- Similar study for γ 'n' → $\eta n, K\Lambda, K\Sigma$ and electroproductions
- More polarization observables including recoil polarization

Thank you very much for your attention

Acknowledgments

- Financial support for this work

  FAPESP 2016/15618-8

  KAKENHI JP25105010

- Computing resource

  The Yukawa Institute Computer Facility

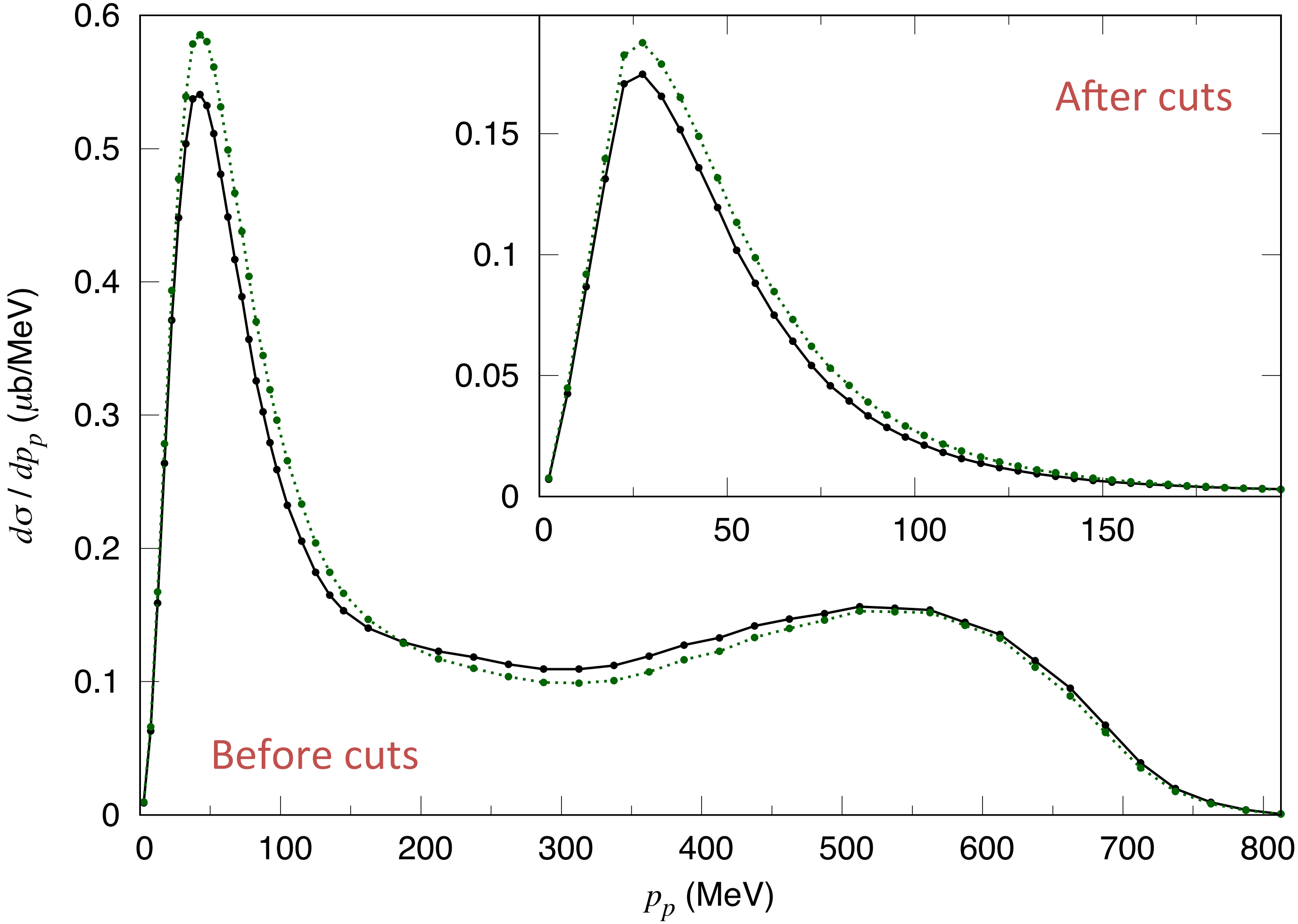
  The Research Center of Nuclear Physics

  LCRC at Argonne National Lab

  NERSC

BACKUP

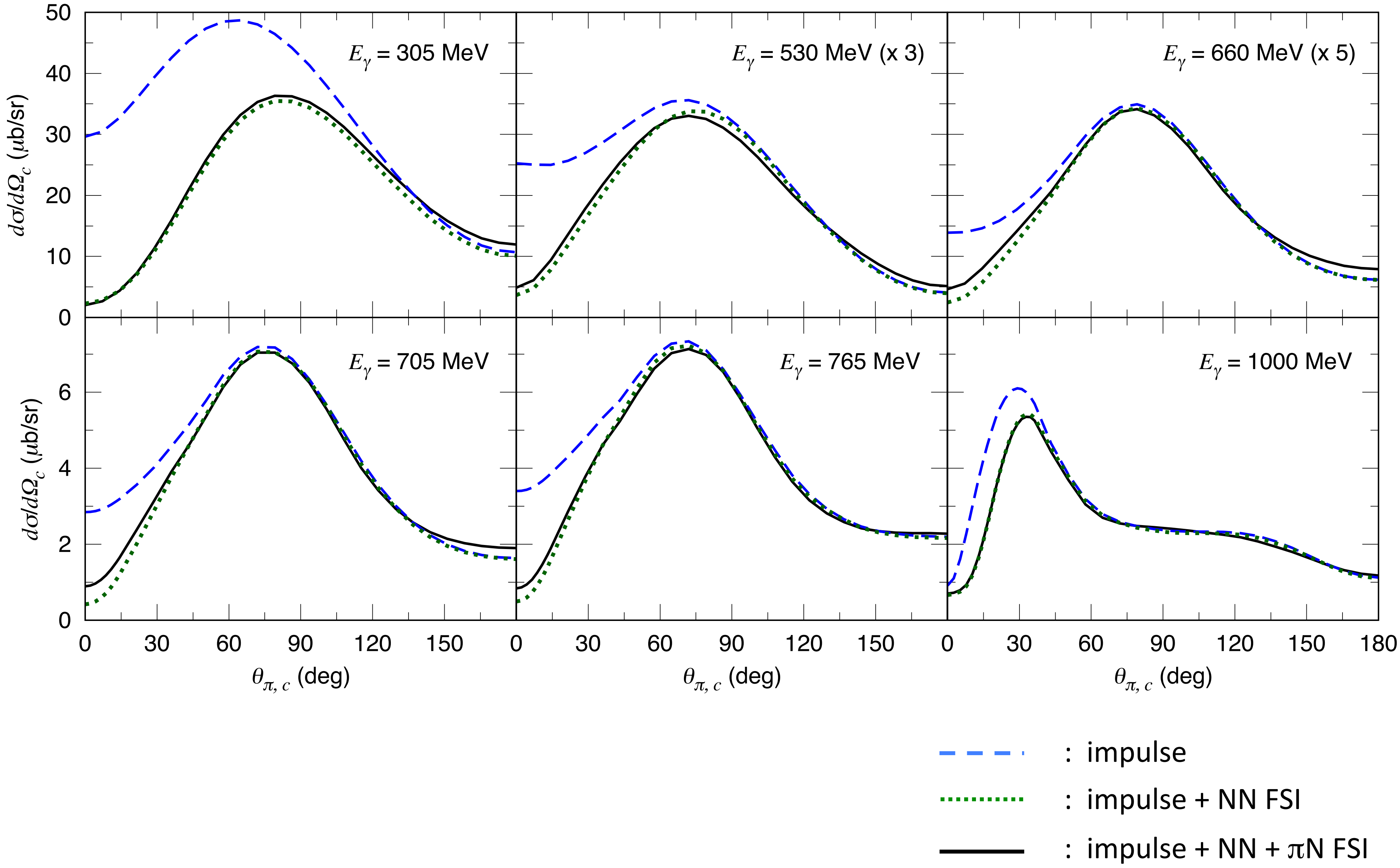
Proton momentum distribution in $\gamma\, d \rightarrow \pi^0\, np$



● : impulse + NN FSI

● : impulse + NN + πN FSI

Differential cross sections for $\gamma d \rightarrow \pi^0 np$



Meson productions induced by photon and electron beams

Interests

- Information on N^* properties (mass, width, missing resonances, etc.)

  complementary to πN scattering data ← "complete" measurements
- Electromagnetic $N \rightarrow N^*$ transition form factors (Q^2 −dependence)

  ➔ important information for hadron structures

Measurements of neutron target data (incl. polarization obs.) are recently active

$\gamma\, n \rightarrow \pi^0\, n$	$d\sigma/d\Omega$	A2@MAMI	PRL 112, 142001 (2014)
	E	A2@MAMI	PLB 770, 523 (2017)
$\gamma\, n \rightarrow \pi^-\, p$	$d\sigma/d\Omega$	CLAS@JLab	PRC 96, 035204 (2017)
	E	CLAS@JLab	PRL 118, 242002 (2017)
$\gamma\, n \rightarrow \eta\, n$	$d\sigma/d\Omega$	A2@MAMI	PRC 90, 015205 (2014)
	E	A2@MAMI	PRC 95, 055201 (2017)

More data are available, and more are coming soon

Extraction of polarization observables for γ 'n' → $\pi^0 n$ from γd → $\pi^0 np$

A possible complication : relativistic nucleon spin rotation

← Lorentz boost from γn CM frame to γd Lab frame

We confirmed this effect is very small for Σ, E, G to be discussed here → ignored

Polarized cross sections can be extracted using formula similar to the unpolarized one

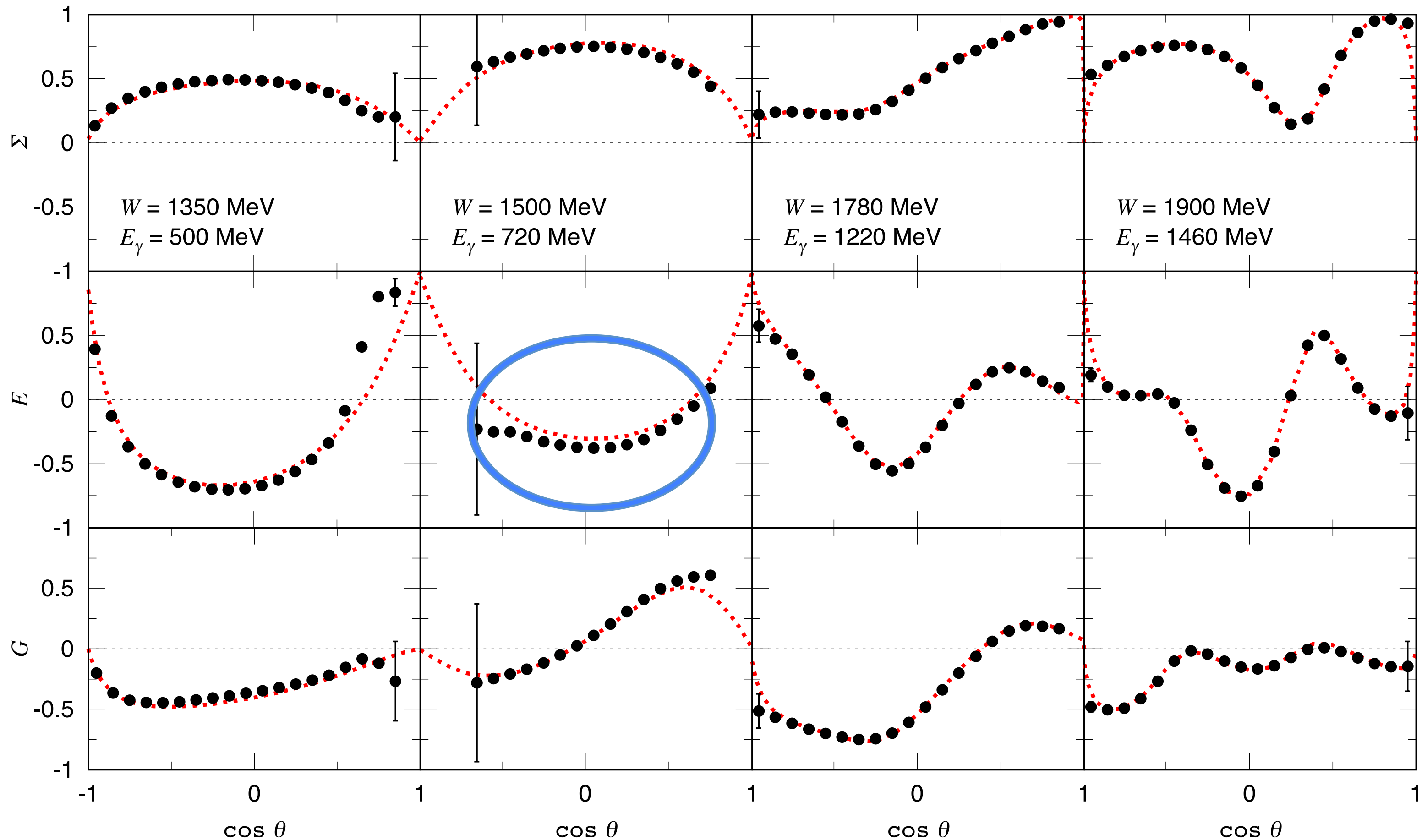
$$\left.\frac{d^2\sigma(\gamma d \to \pi^0 np;\ \lambda, s_d = +1)}{dW\, d\cos\theta^*_\pi}\right|_{\mathrm{cuts}} = \phi_{\mathrm{mod}}(W)\, \frac{\tilde{E}_\gamma}{E_\gamma}\, \frac{d\sigma(\gamma n \to \pi^0 n; \lambda, s_N = +1/2)}{d\cos\theta^*_\pi}$$

$$E = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{d\sigma(\lambda = -1, s_N = +1/2) - d\sigma(\lambda = +1, s_N = +1/2)}{d\sigma(\lambda = -1, s_N = +1/2) + d\sigma(\lambda = +1, s_N = +1/2)}$$

Σ, E, G for γ 'n' ➔ $\pi^0 n$

● : impulse + NN FSI + πN FSI

⋯⋯ : free $\sigma(\gamma n \rightarrow \pi N)$

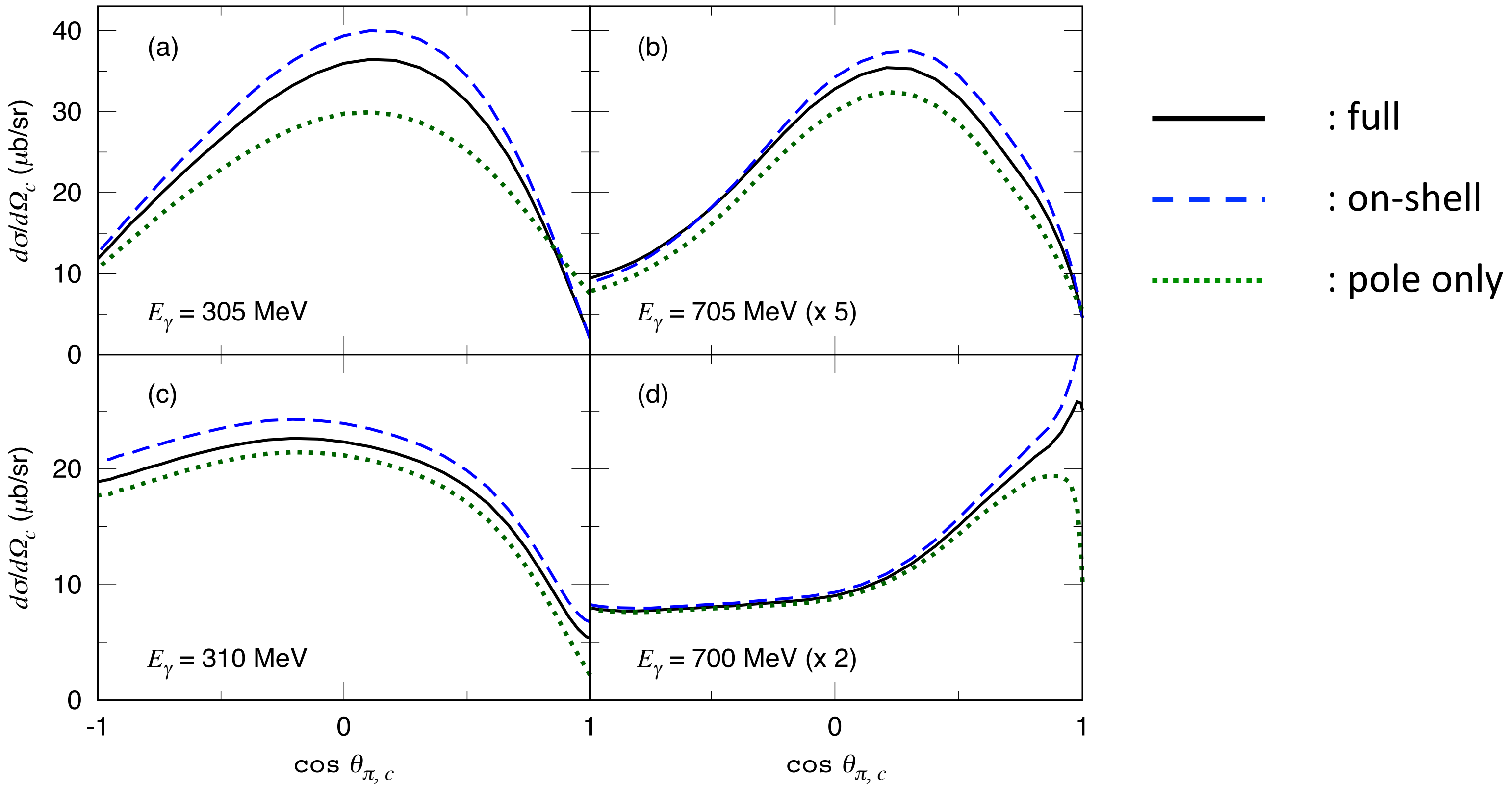


- Free $\gamma n \rightarrow \pi N$ polarization asymmetries are well reproduced
- FSI effects are very small ; canceled by taking the ratios; One exception
- No FSI effects are assumed in A2@MAMI analysis for E [PLB 770 (2017)]

Off-shell effects

$\gamma n \rightarrow \pi^0 n$

$\gamma n \rightarrow \pi^- p$

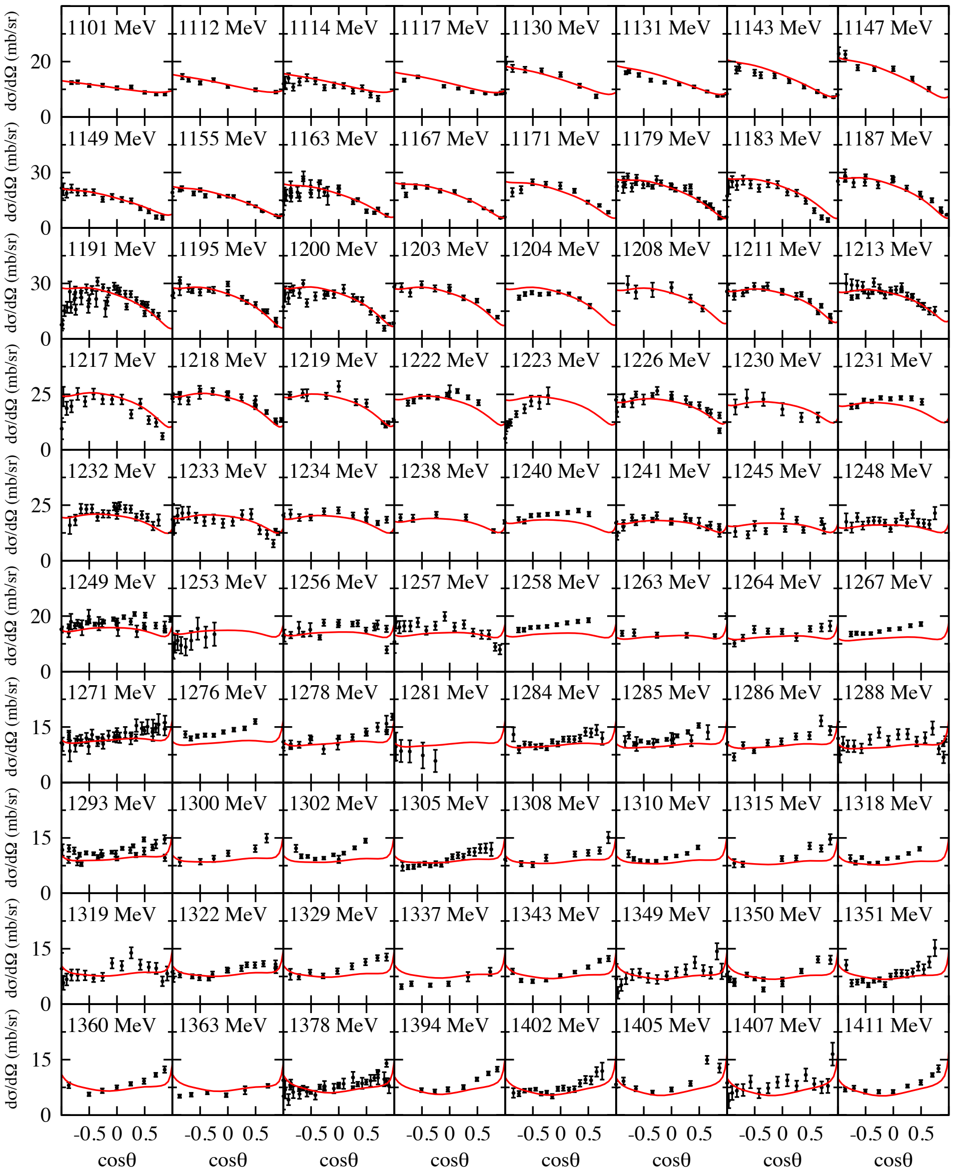


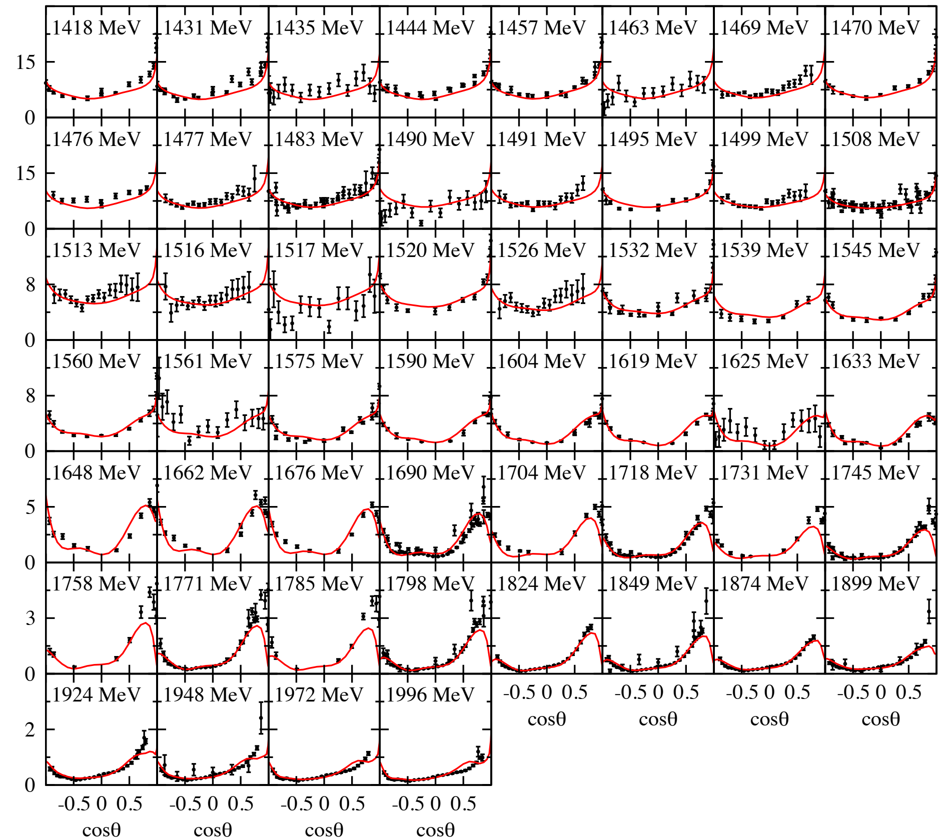
γn → π⁻p

dσ/dΩ for W < 2.1 GeV

Comparison of DCC model with data

Kamano, Nakamura, Lee, Sato, PRC 94 (2016)





Elementary amplitudes obtained by this fit go into our deuteron reaction model

Predicted πN → ππN total cross sections with our DCC model