

DISPERSIVE APPROACH TO THREE BODY SCATTERING

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Ψ

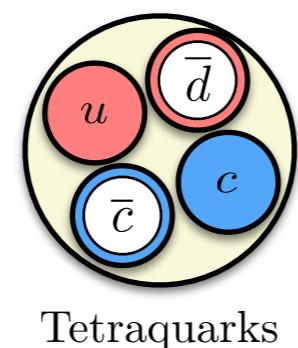
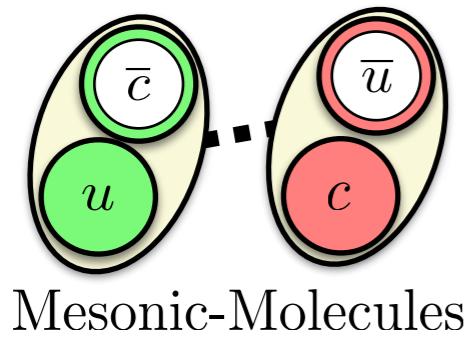
JPAC

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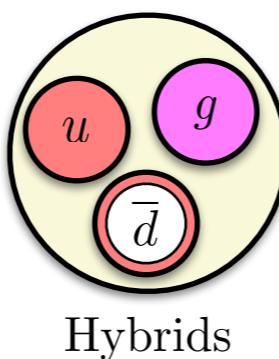
A. Jackura, Indiana University and the Joint Physics Analysis Center

Hadrons and QCD

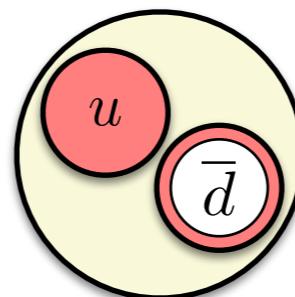
- Signatures of hadrons in data
 - Peaks and dips on **Real** energy axis
 - Patterns in mass/ spin
- Theoretical description
 - **Complex** energy plane singularities
- Interpretation
 - Constituent quarks, exotica



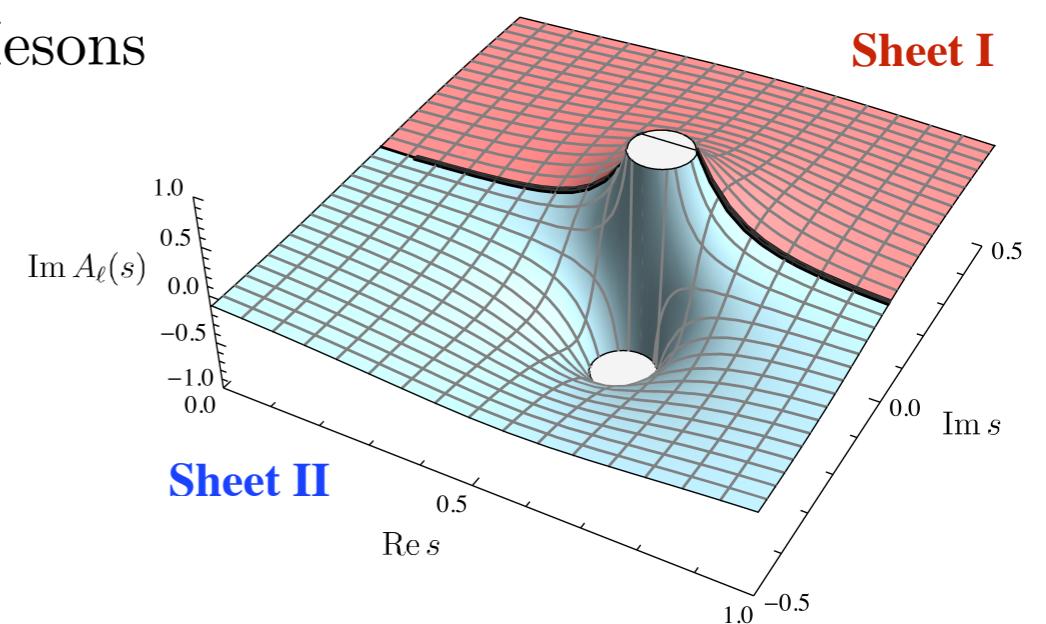
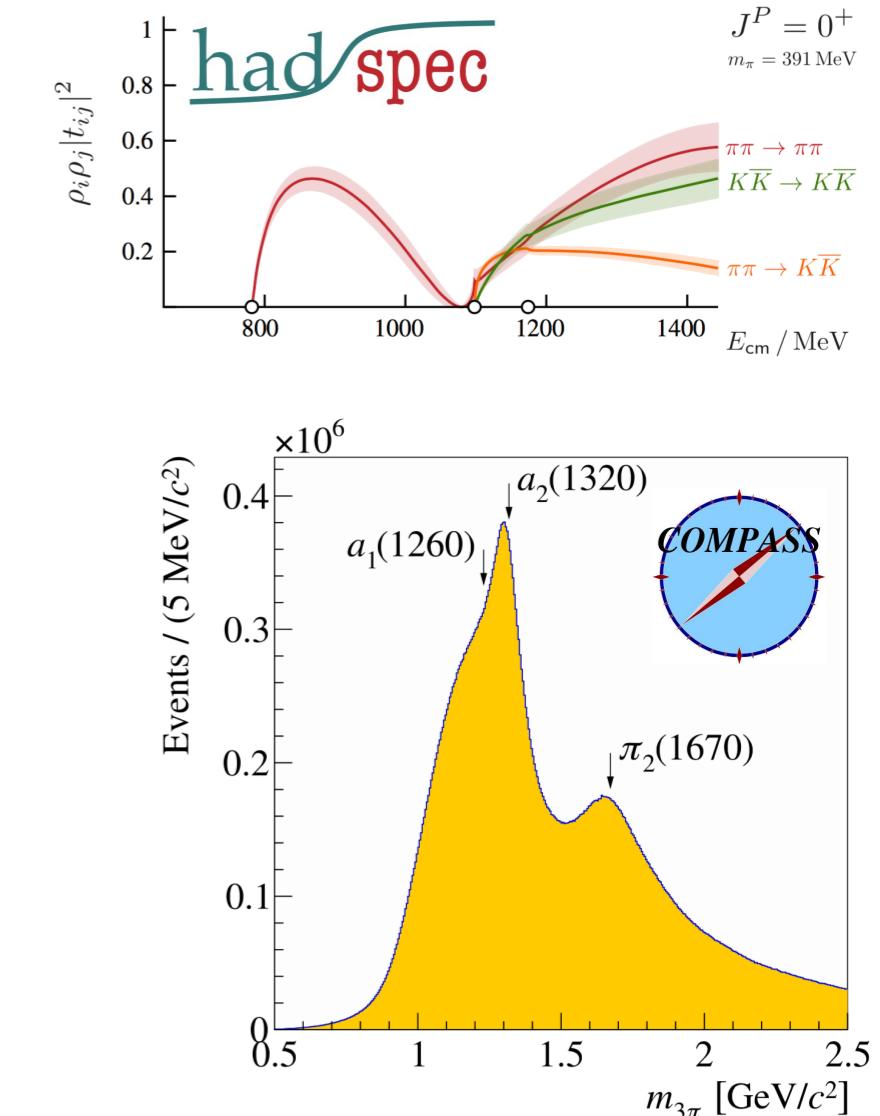
Tetraquarks



Hybrids

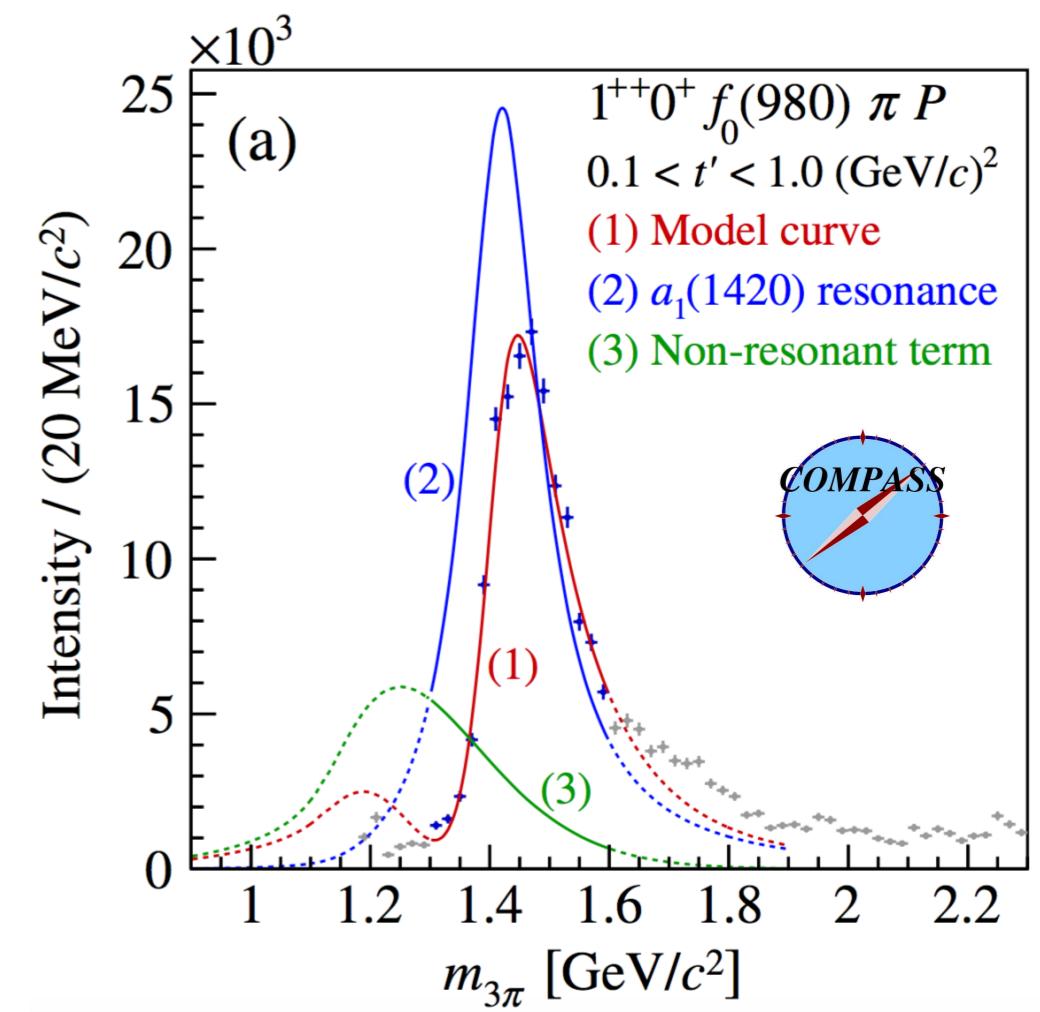
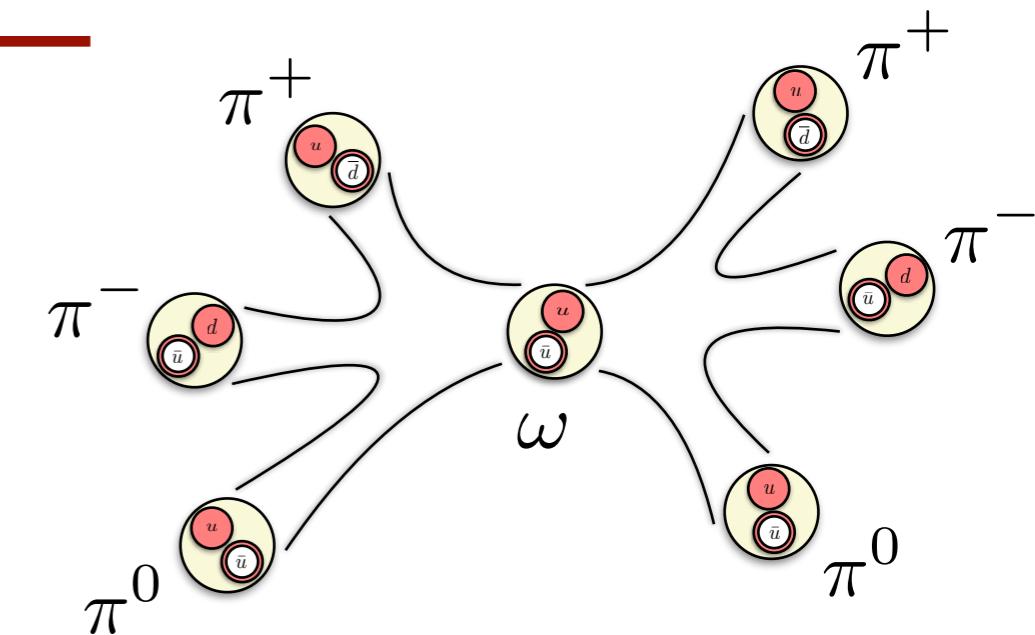
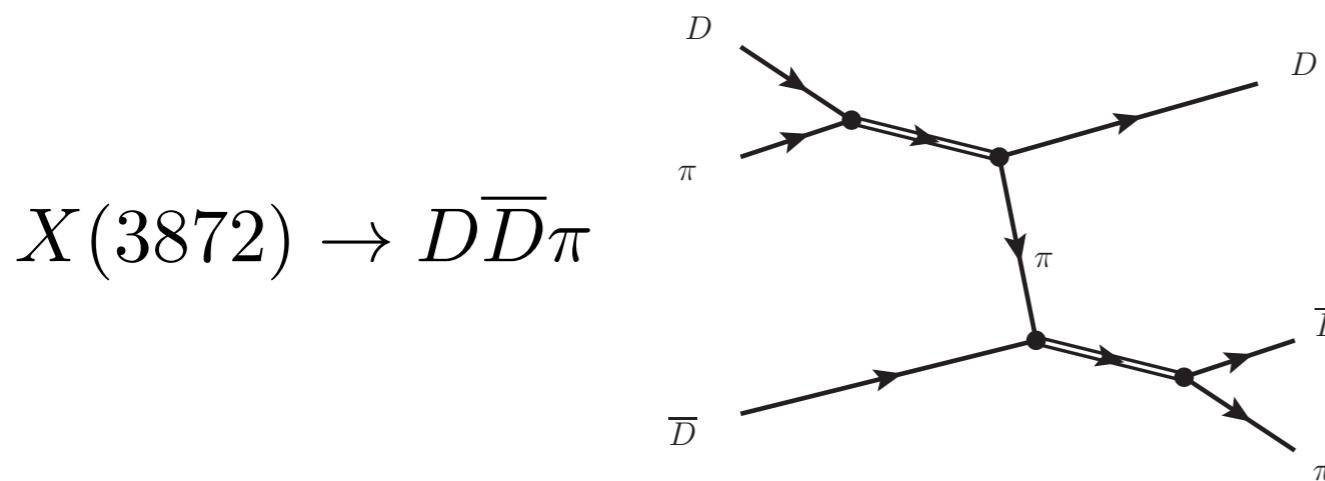


Mesons



3-Body Phenomenology

- Many Resonances decay to 3-particles
 - New high-statistics experimental results
 - Advancements in Lattice Formalism
- Signatures of new resonances
 - High-Precision PWA of 3π yields $a_1(1420)$, Incompatible with quark model/ Regge expectations
- XYZ States
 - Lie near thresholds



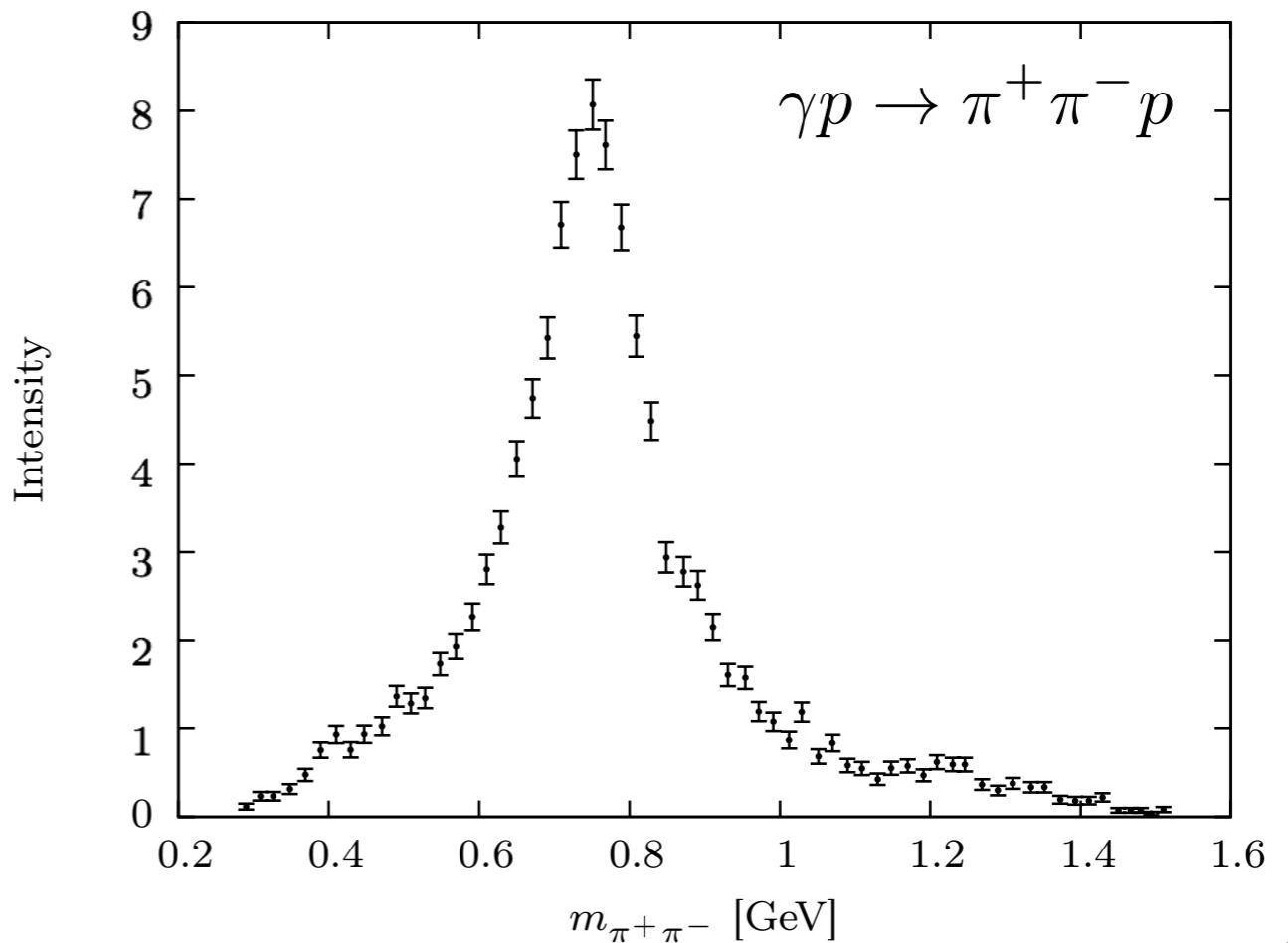
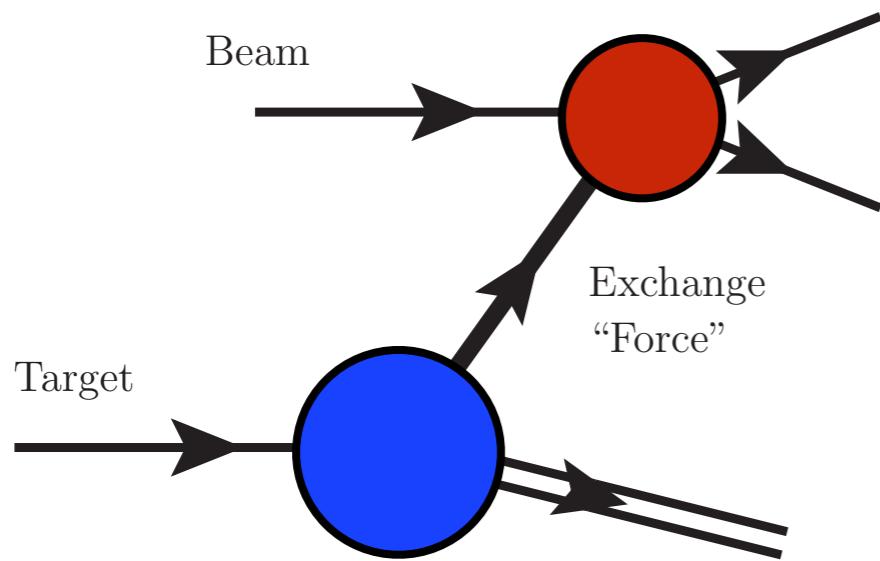
Scattering Phenomenology

- Model independent constraints - **S-matrix theory**
 - **Unitarity, Analyticity, Crossing, and Poincaré Symmetry**

Scattering Phenomenology

- Model independent constraints - **S-matrix theory**
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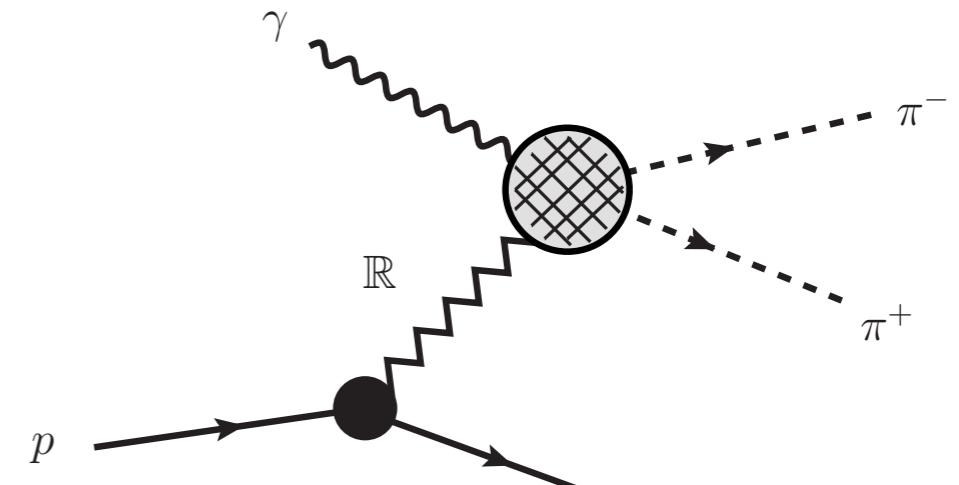
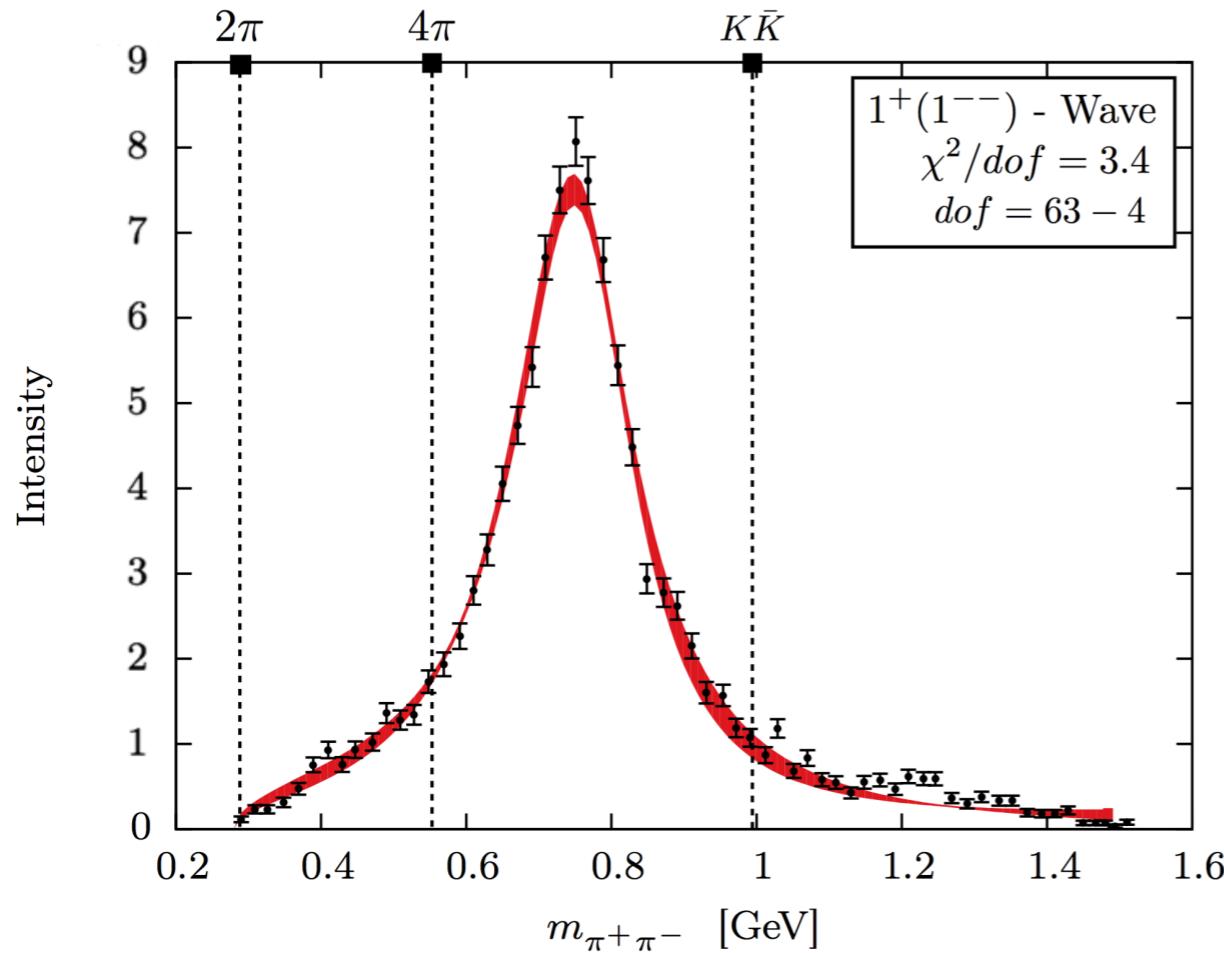
Data (Experimental/Lattice)



Scattering Phenomenology

- Model independent constraints - **S-matrix theory**
 - **Unitarity**, **Analyticity**, **Crossing**, and **Poincaré Symmetry**

Amplitude Analysis



Amplitude Model

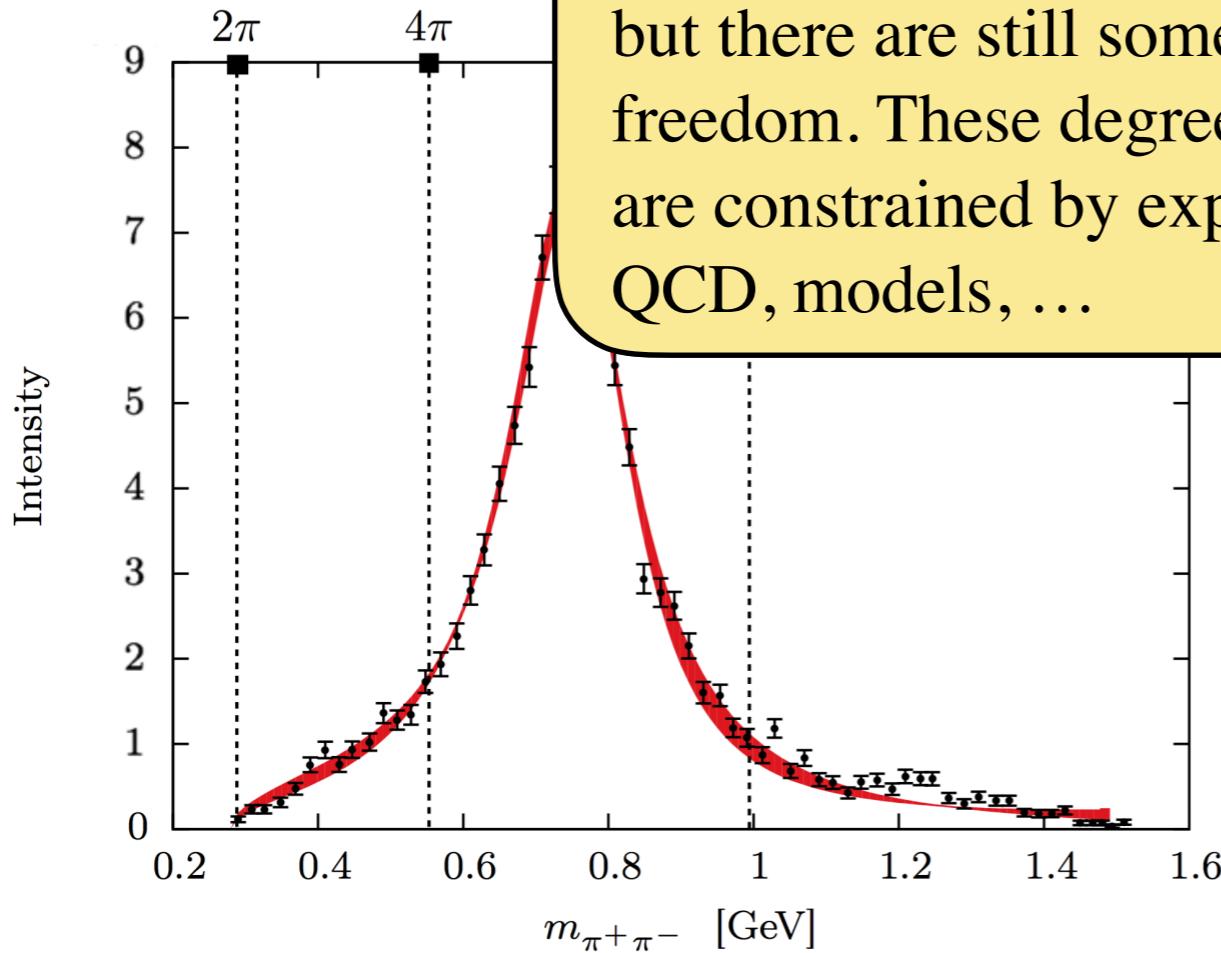
$$\mathcal{A}_\ell(s) = K_\ell^{-1}(s) - \frac{s}{\pi} \int_{s_{th}}^{\infty} ds' \frac{\rho(s')}{s' - s}$$

Functional form fixed
from S-matrix constrains

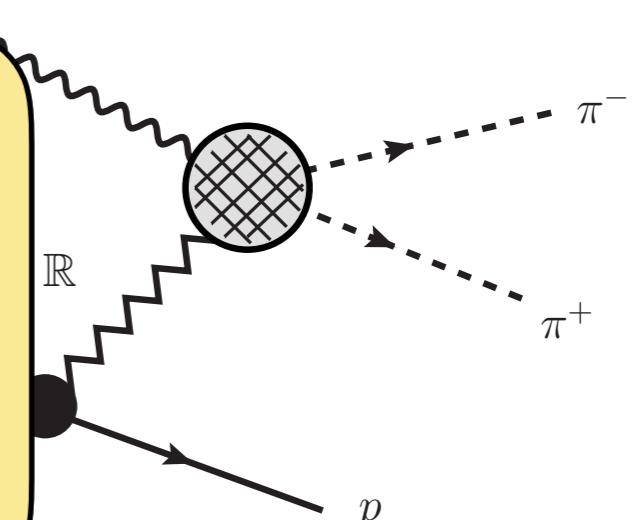
Scattering Phenomenology

- Model independent constraints - **S-matrix theory**
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Amplitude



Imposing these constraints gives us general forms of reaction amplitudes, but there are still some degrees of freedom. These degrees of freedom are constrained by experiment, lattice QCD, models, ...



Amplitude Model

$$\mathcal{A}_\ell(s) = K_\ell^{-1}(s) - \frac{s}{\pi} \int_{s_{th}}^{\infty} ds' \frac{\rho(s')}{s' - s}$$

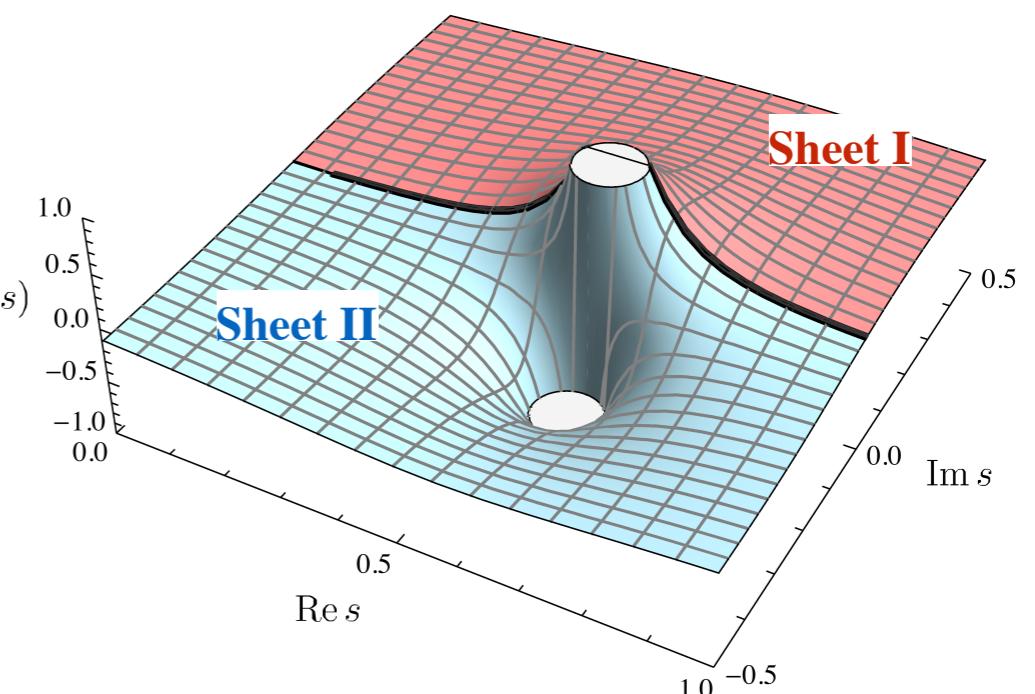
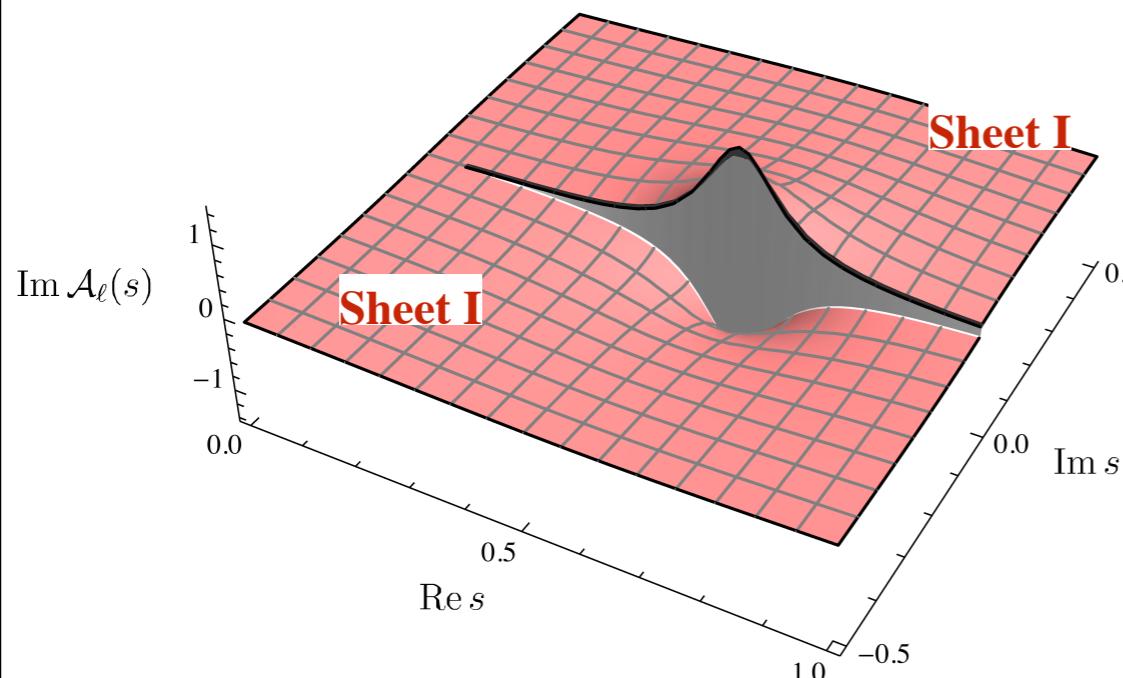
Functional form fixed
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Scattering Phenomenology

- Model independent constraints - **S-matrix theory**
 - **Unitarity**, **Analyticity**, **Crossing**, and **Poincaré Symmetry**

Analytic Continuation

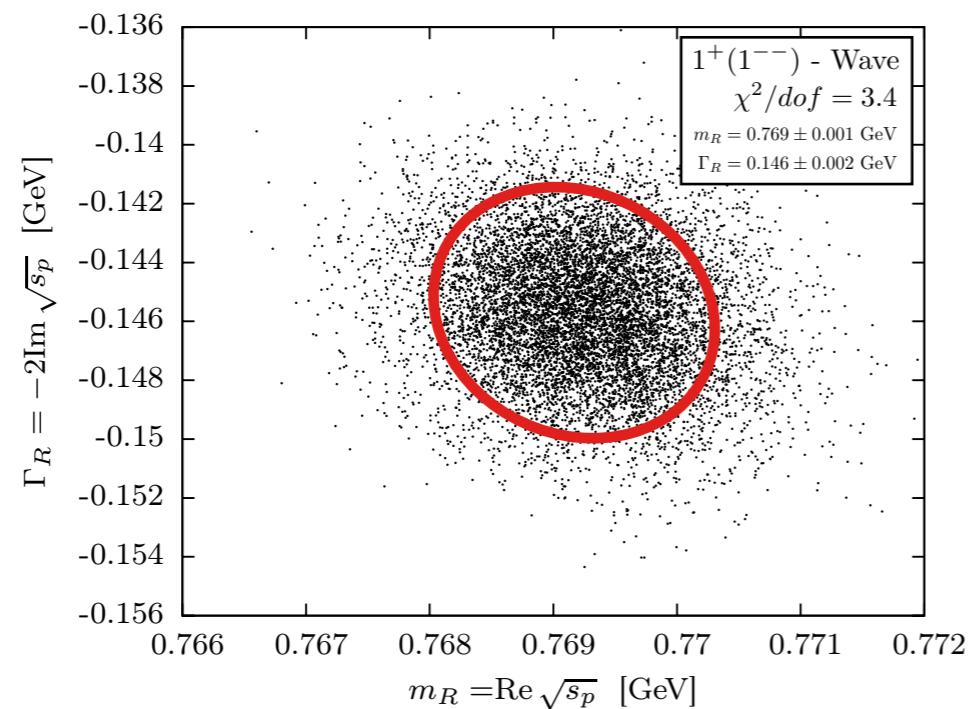
$$\mathcal{A}_\ell(s) \rightarrow \mathcal{A}_\ell^{\text{II}}(s) = \frac{\mathcal{A}_\ell(s)}{1 + 2i\rho(s)\mathcal{A}_\ell(s)}$$



Scattering Phenomenology

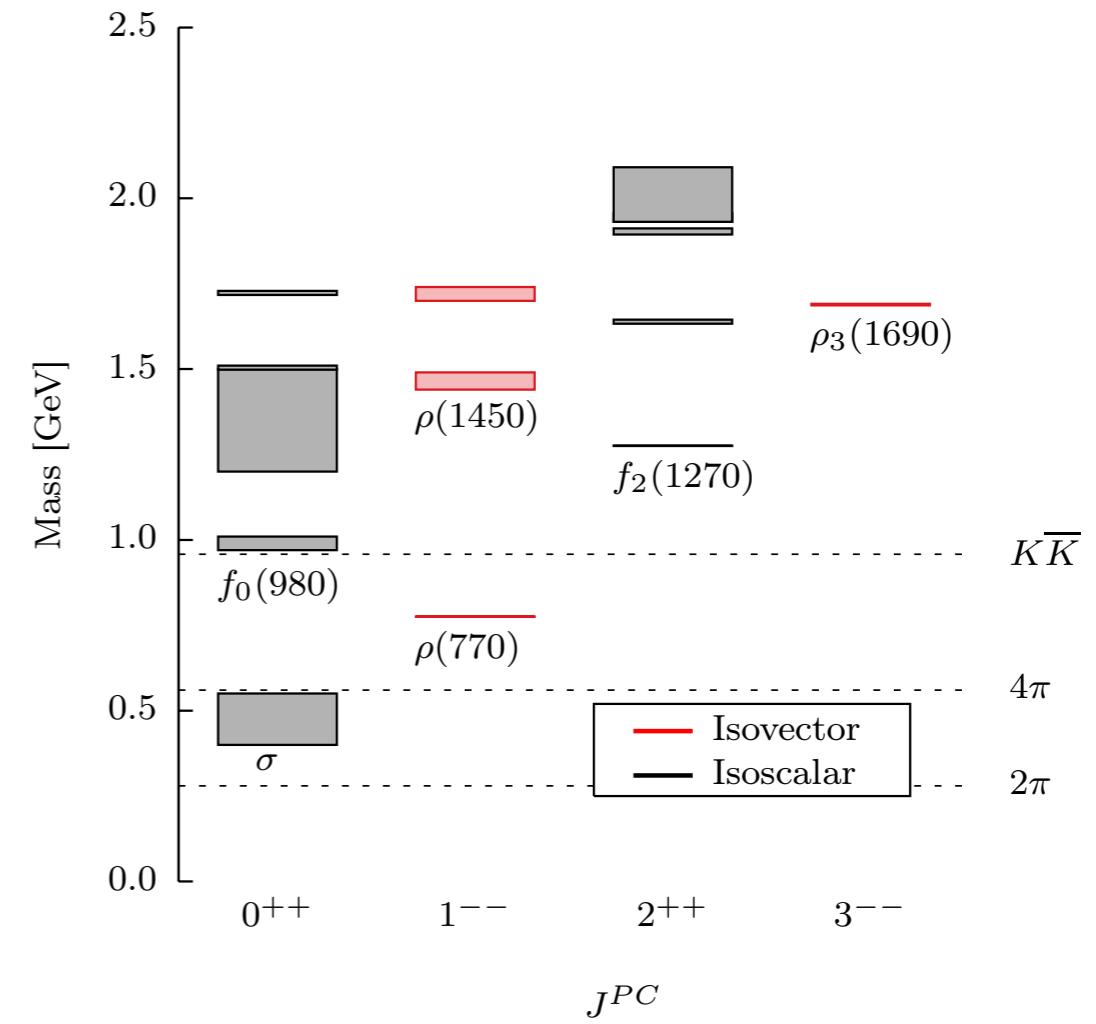
- Model independent constraints - **S-matrix theory**
 - **Unitarity, Analyticity, Crossing, and Poincaré Symmetry**

Resonance Parameters



$$m_R = 769 \pm 1 \text{ MeV}$$

$$\Gamma_R = 149 \pm 2 \text{ MeV}$$



$3 \rightarrow 3$ Scattering

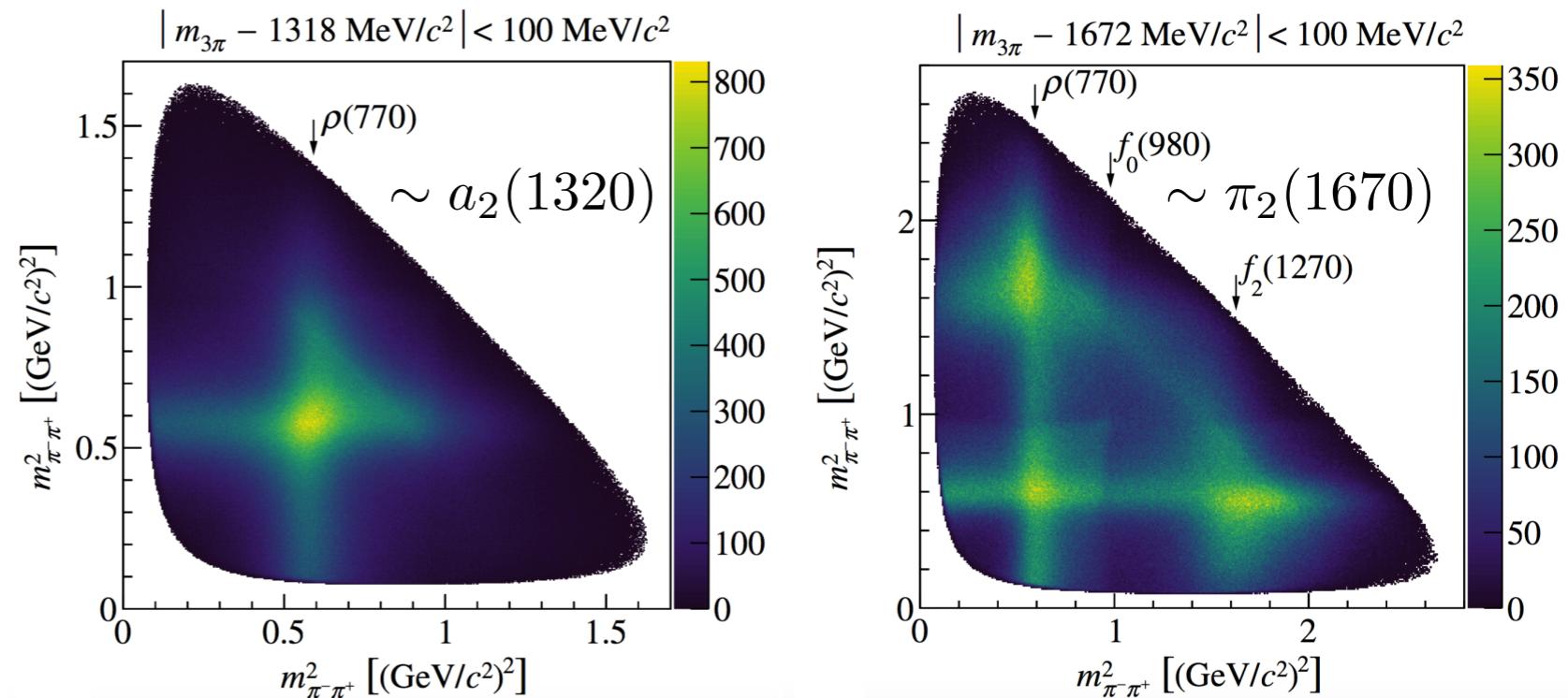
Unitarity

Analyticity

Dispersion Relations

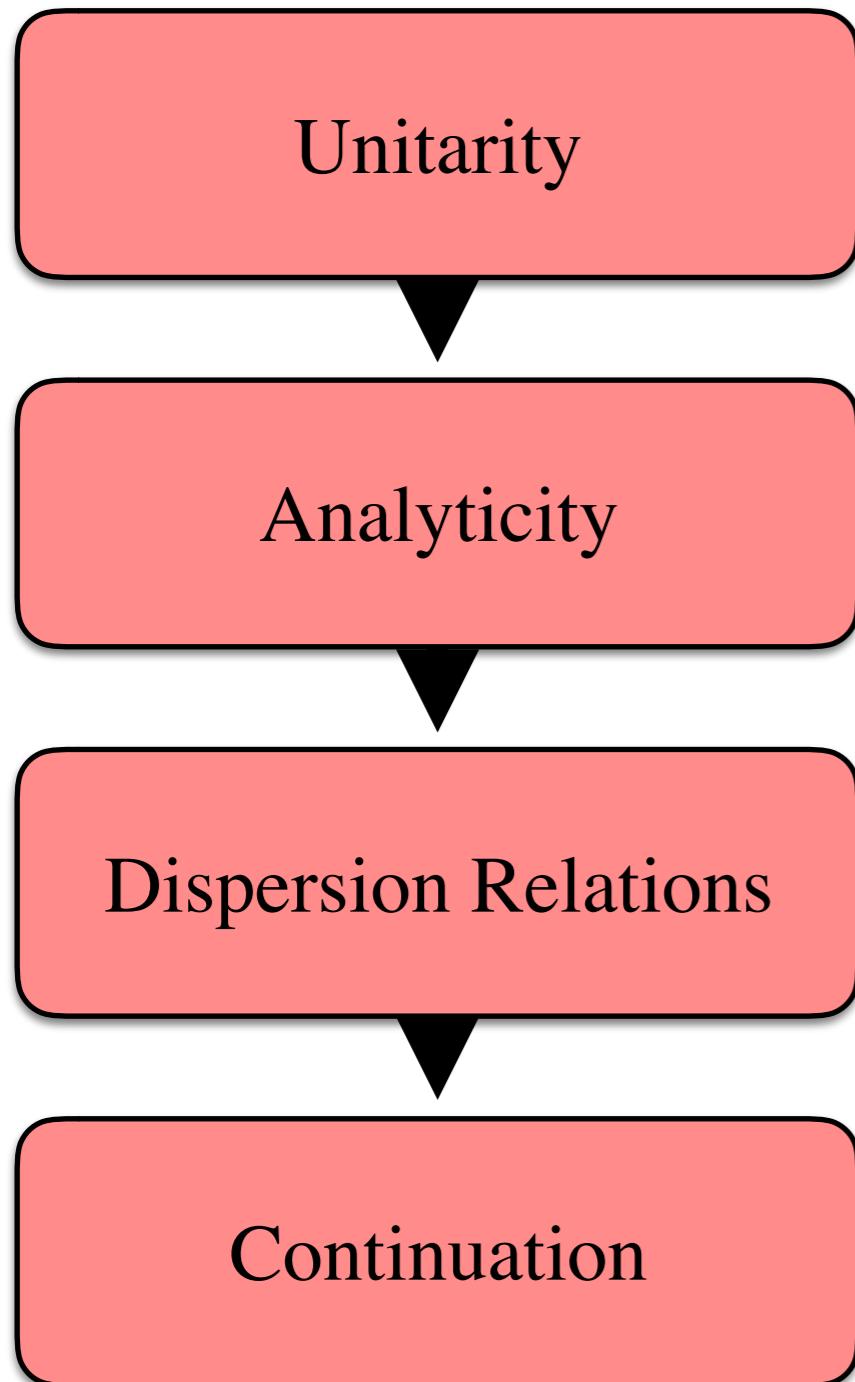
Continuation

- Consider the elastic scattering $123 \rightarrow 123$
 - Distinguishable, spinless particles
- Isobar Model
 - Interacting two-particle sub-channel (**Isobar**) interacts with third particle (**Spectator**)

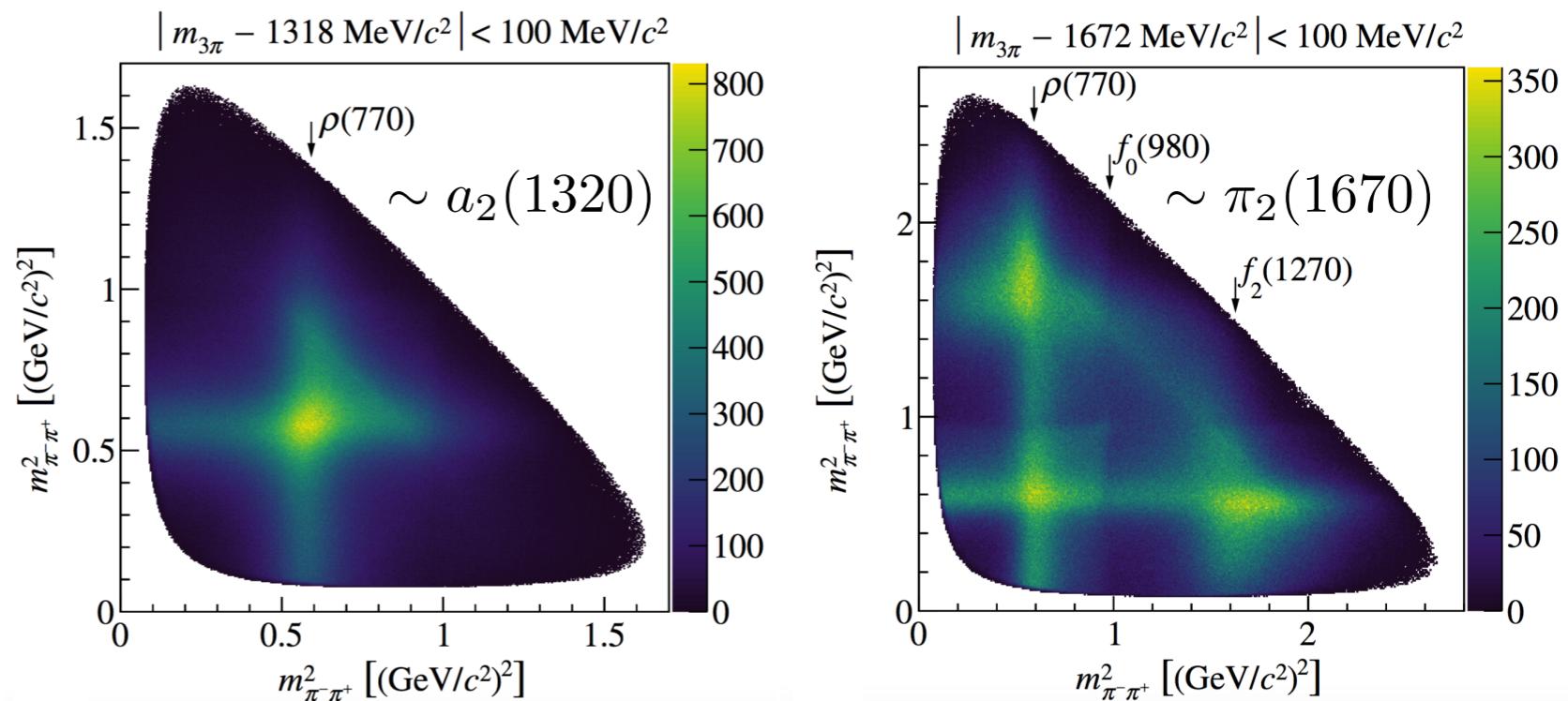


C. Adolph et al. [COMPASS],
Phys. Rev. D 95, no. 3, 032004 (2017)

$3 \rightarrow 3$ Scattering

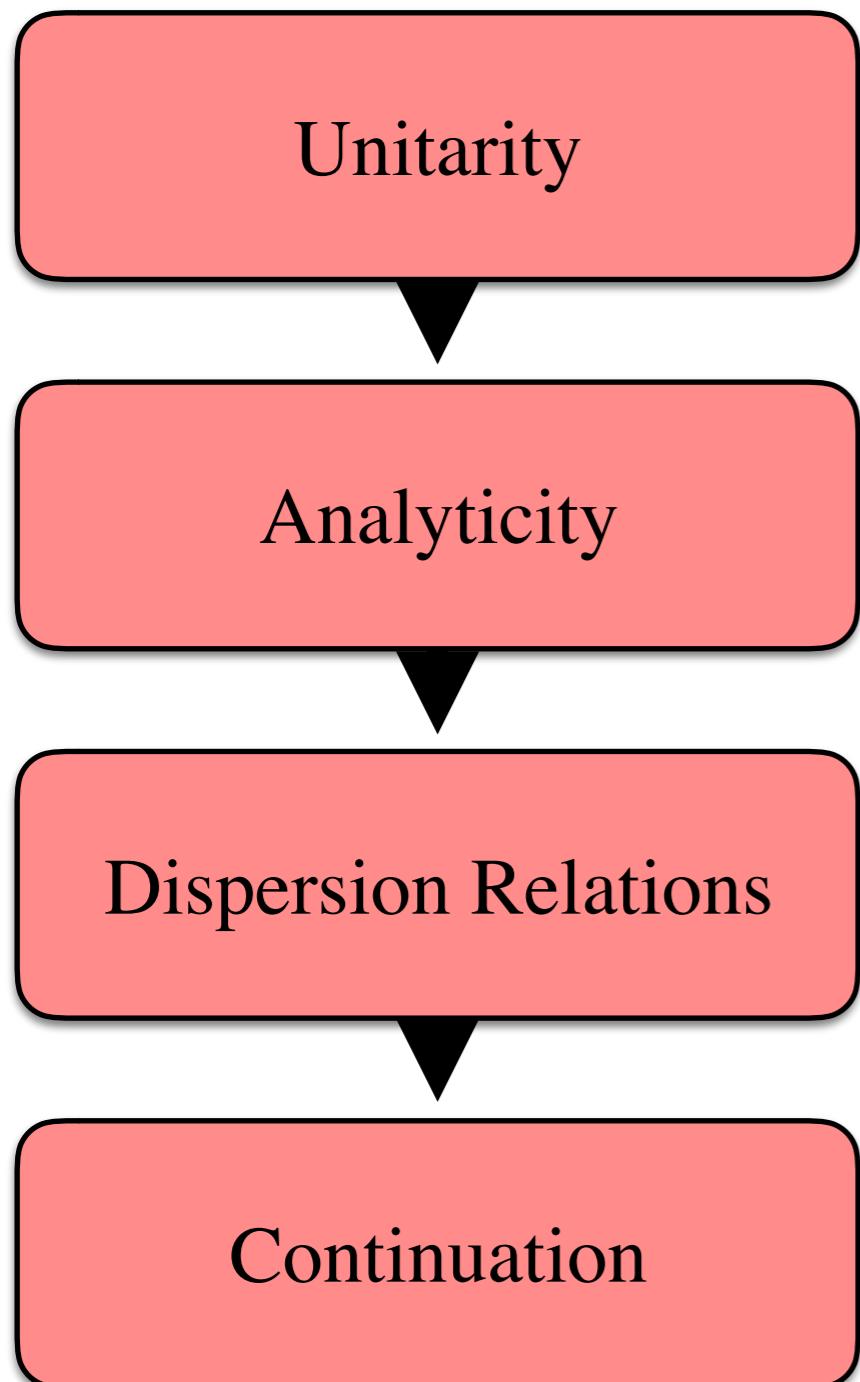


- Consider the elastic scattering $123 \rightarrow 123$ of three massless particles
- Multiple variables → Isobar approximation
 - Interacting two-particle sub-channel (**Isobar**) interacts with third particle (**Spectator**)



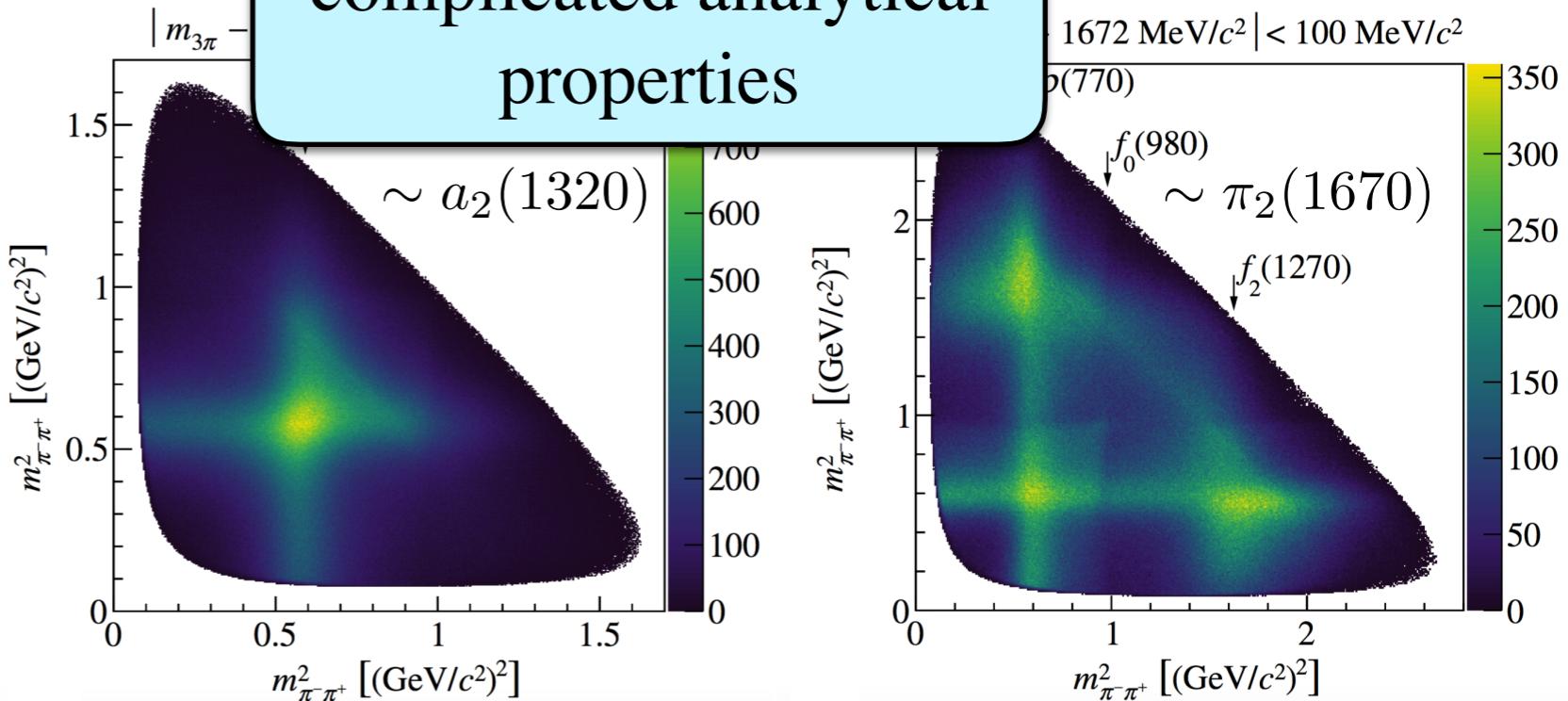
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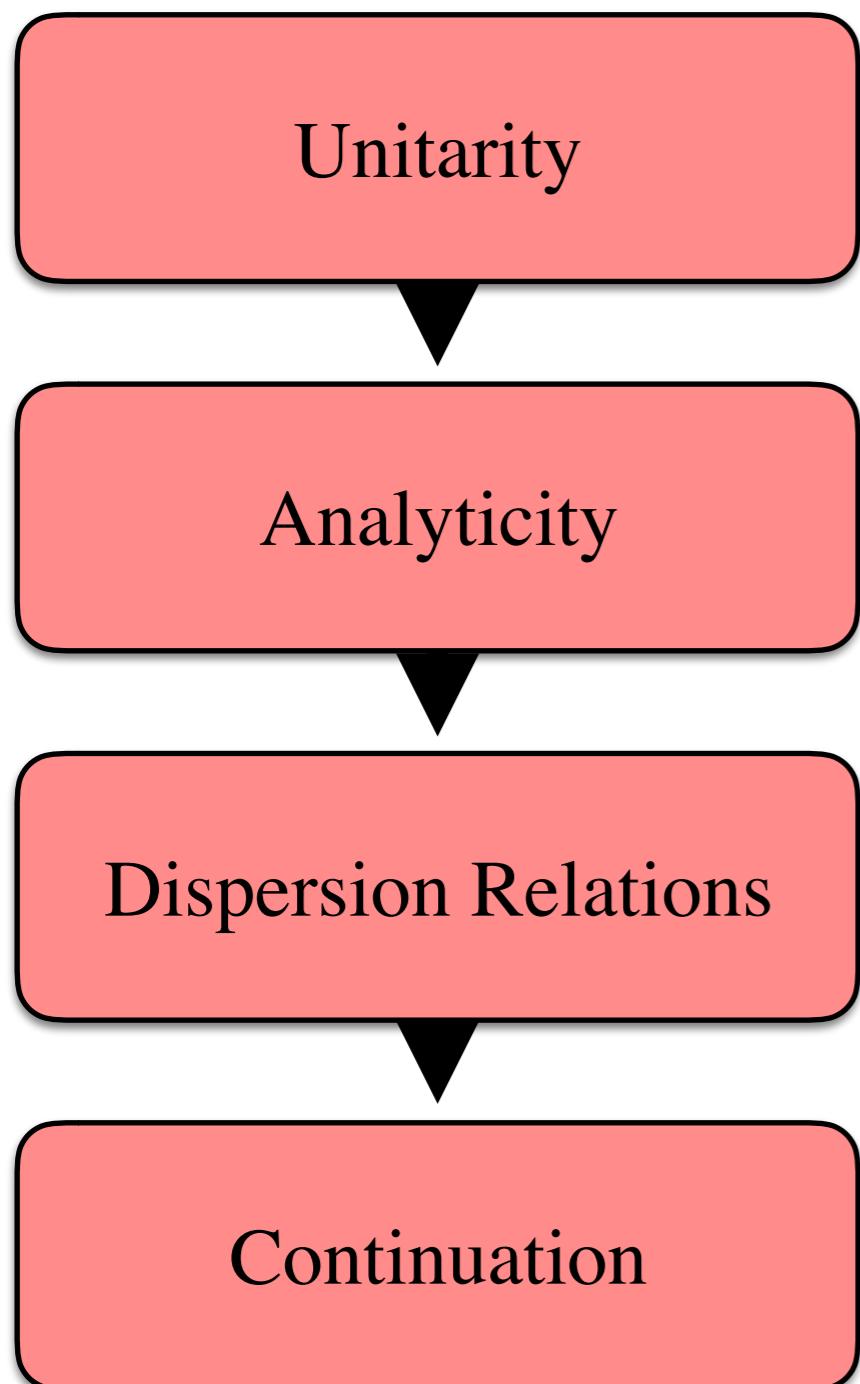
- Consider the elastic scattering $123 \rightarrow 123$ process particles
- Multiple variables → Isobar approximation
 - interacting two-particle sub-channel (**Isobar**) interacting three-particle sub-channel (**Spectator**)

Isobar/Partial Waves → complicated analytical properties

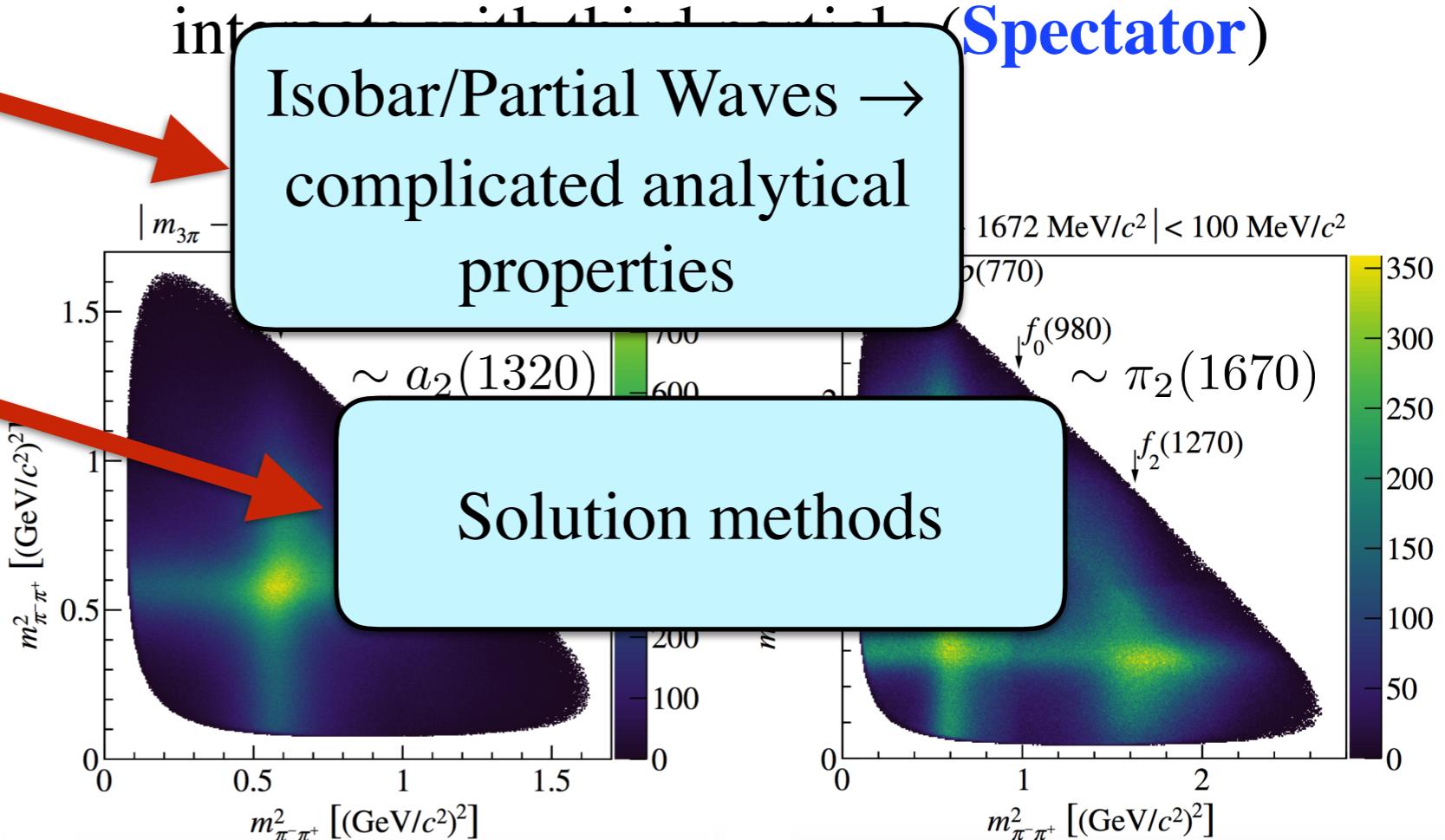


*C. Adolph et al. [COMPASS],
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$3 \rightarrow 3$ Scattering

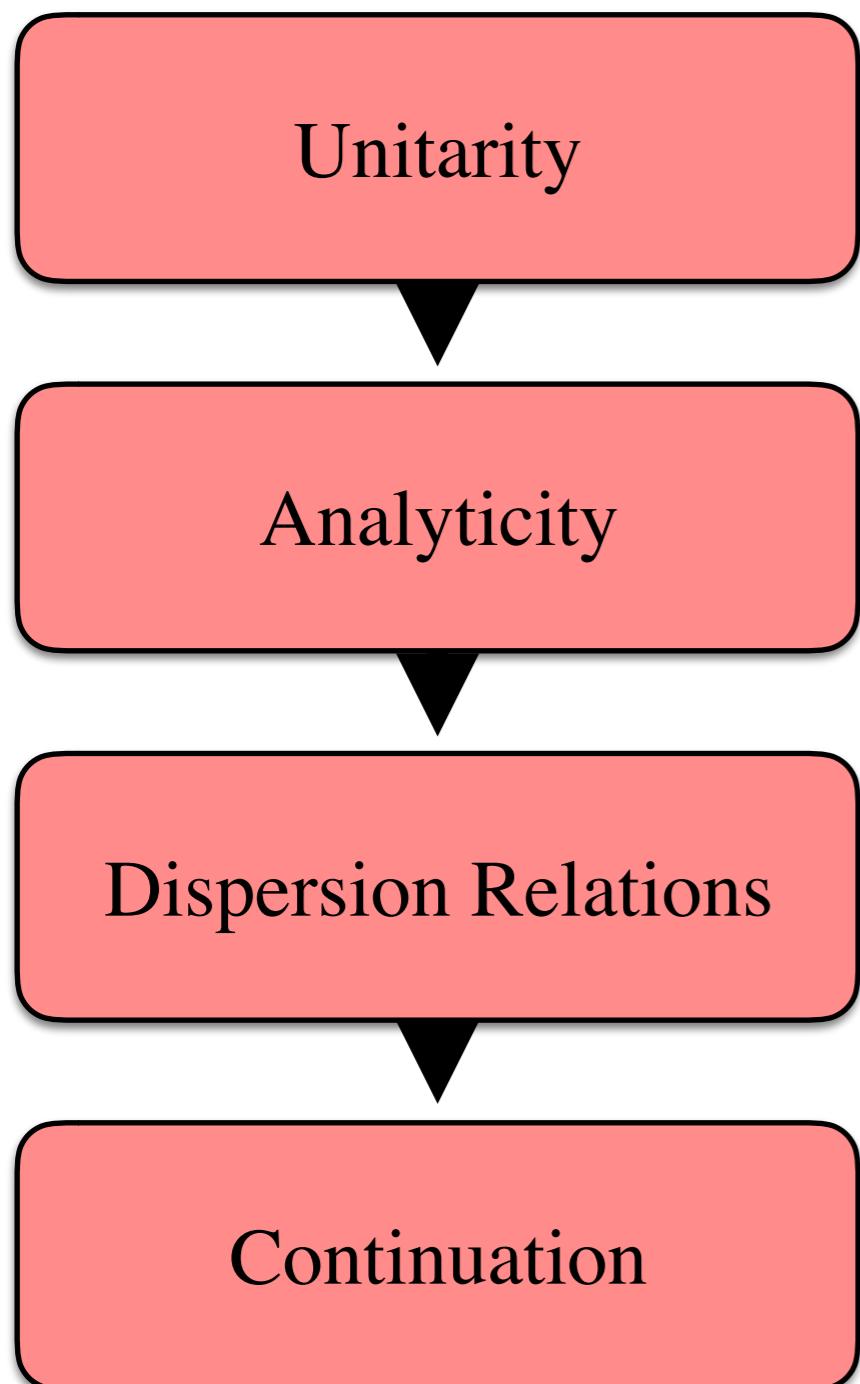


- Consider the elastic scattering $123 \rightarrow 123$
- Massless particles
- Multiple variables → Isobar approximation
 - interacting two-particle sub-channel (**Isobar**)
 - interacting three-particle sub-channel (**Spectator**)

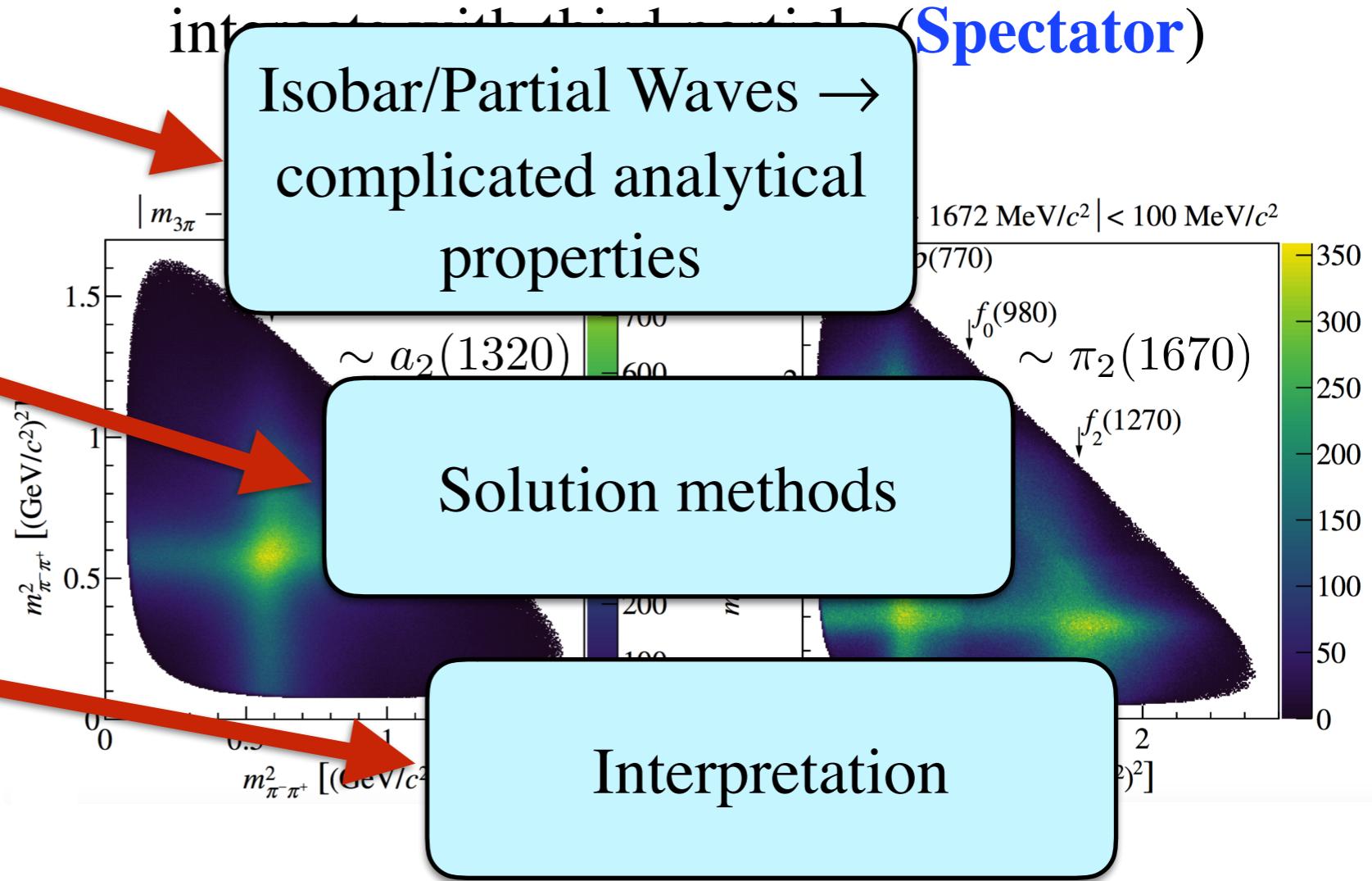


*C. Adolph et al. [COMPASS],
Phys. Rev. D 95, no. 3, 032004 (2017)*

$3 \rightarrow 3$ Scattering



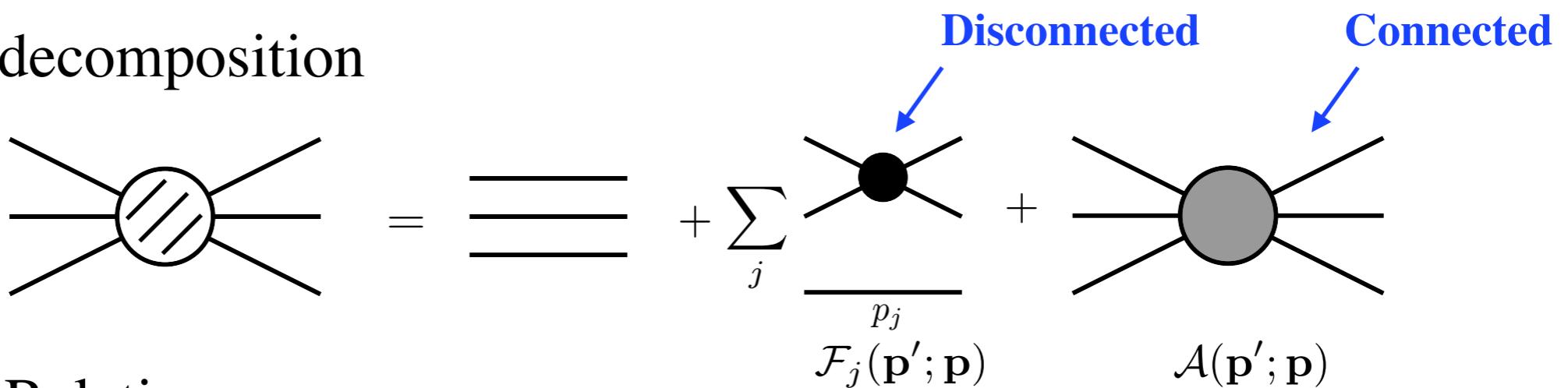
- Consider the elastic scattering $123 \rightarrow 123$ of three massless particles
- Multiple variables → Isobar approximation
 - interacting two-particle sub-channel (**Isobar**) interacting with a three-particle sub-channel (**Spectator**)



Phys. Rev. D 95, no. 3, 032004 (2017)

3 \rightarrow 3 Scattering

- S-matrix decomposition



- Unitarity Relations

Disconnected

$$2 \operatorname{Im} \text{ (disconnected vertex)} = \text{ (two vertices connected by a horizontal line with a red dashed vertical line through it)}$$

Connected

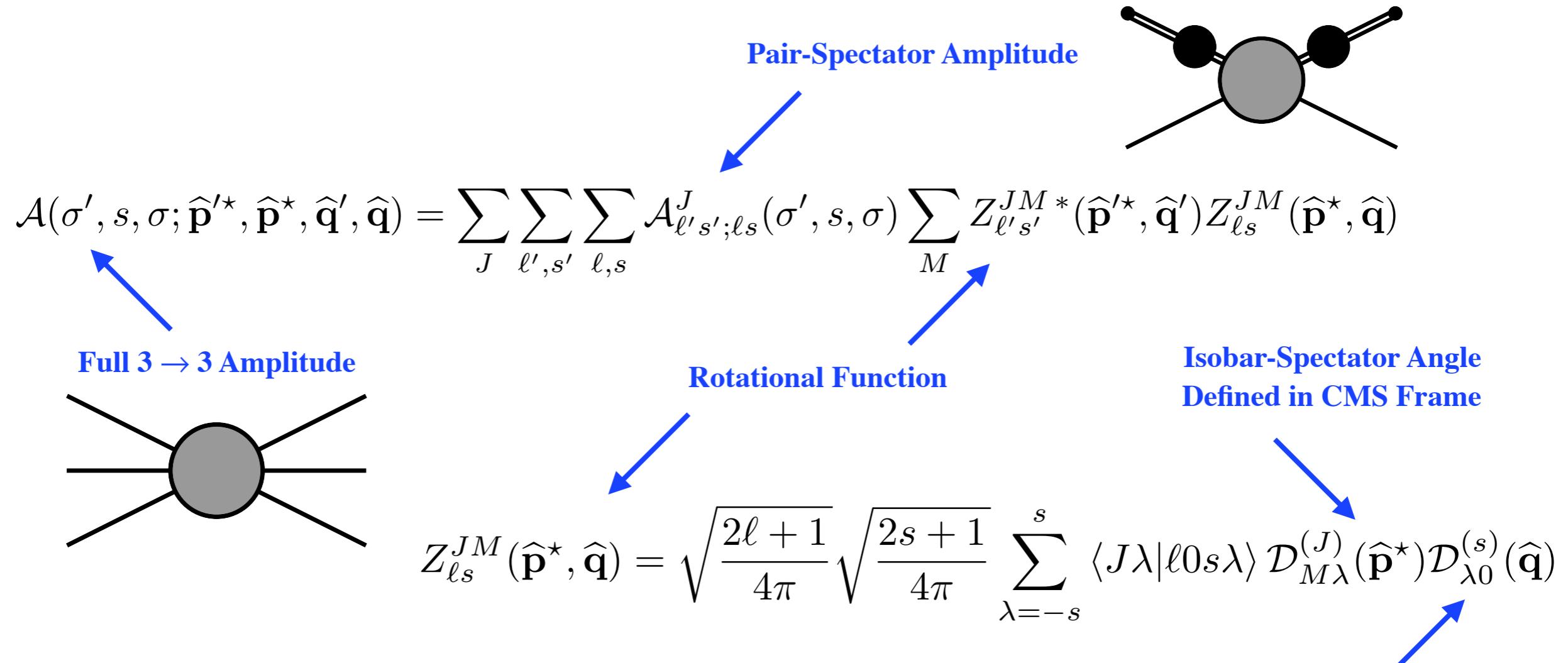
$$2 \operatorname{Im} \text{ (connected vertex)} = \text{ (two connected vertices with a red dashed vertical line through them)}$$

$$\begin{aligned}
 & + \sum_j \text{ (one vertex with one outgoing line labeled } p_j) \\
 & + \sum_k \text{ (one vertex with one outgoing line labeled } p'_k) \\
 & + \sum_{j,k} \text{ (two vertices connected by a horizontal line with a red dashed vertical line through it, with labels } p_j \text{ and } p'_k \text{ on the lines, and } (P - p_j) - p'_k \text{ between the vertices)}
 \end{aligned}$$

*G.N. Fleming,
Phys. Rev. 135, B551 (1964)*

Partial Wave Decomposition

- Partial wave expansion of $3 \rightarrow 3$ elastic amplitude
 - Remove angular dependence \rightarrow 3 Energy variables



- Isobar-Spectator Scattering angle Θ^*

$$\sum_M D_{M\lambda'}^{(J)*}(\hat{\mathbf{p}}'^*) D_{M\lambda}^{(J)}(\hat{\mathbf{p}}^*) = D_{\lambda\lambda'}^{(J)}(\hat{\mathbf{p}}'^* \cdot \hat{\mathbf{p}}^*) \quad \cos \Theta^* = \hat{\mathbf{p}}'^* \cdot \hat{\mathbf{p}}^*$$

Isobar Model

- Assume sum over pair-wise interactions \Rightarrow **Isobar-Spectator amplitudes**
 - Partial wave expansion truncated
 - Leads to complicated sub-channel energy-dependence

$$\text{Diagram: } \text{Three external lines meeting at a central gray circle.} = \sum_{j,k} \text{Diagram: } \text{Three external lines meeting at a central gray circle, which is connected to two black circles, which are further connected to each other. The left black circle is labeled } p'_k \text{ and the right black circle is labeled } p_j.$$

$$\mathcal{A}(\mathbf{p}'; \mathbf{p}) = \sum_{j,k} \mathcal{A}_{kj}(\mathbf{p}'; \mathbf{p})$$

- Partial Wave Projected Isobar-Spectator (PWIS) Amplitudes

$$\tilde{\mathcal{A}}_{\ell' s'; \ell s}^J(\sigma', s, \sigma) = \tilde{\mathcal{A}}_{\ell' s'; \ell s}^J(\sigma', s, \sigma)|_{\text{Iso}} + \sum_{\substack{j, k \\ \text{crossed}}} [X_{\ell' s'; \ell s}^J]_{jk}(\sigma', s, \sigma)$$

Full $3 \rightarrow 3$ PW Amplitude

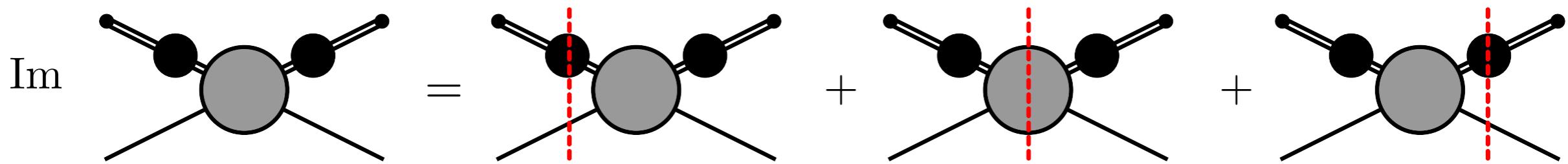
PWIS Amplitude

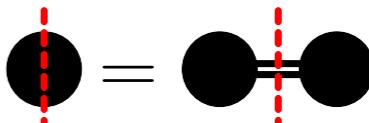
Recoupling PWIS Amplitudes

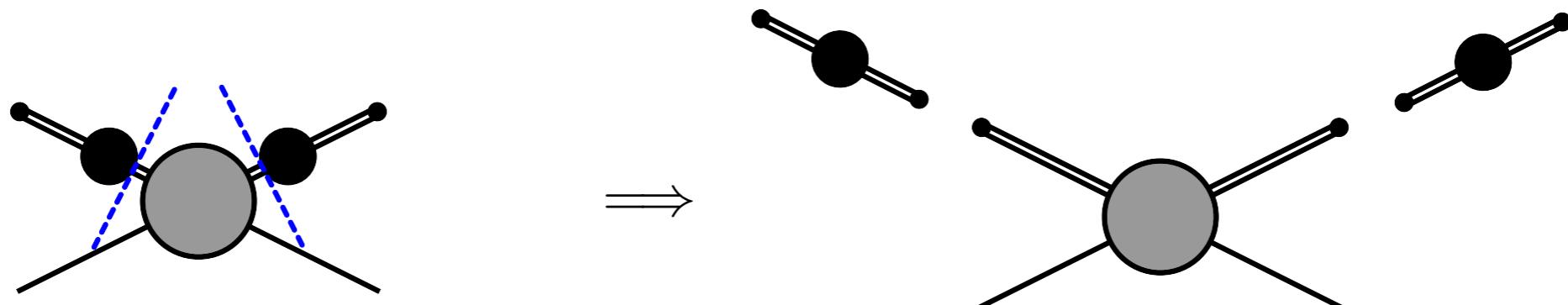
$$\text{Diagram: } \text{Three external lines meeting at a central gray circle.} = \text{Diagram: } \text{Three external lines meeting at a central gray circle, which is connected to two black circles, which are further connected to each other. The left black circle is labeled } 3 \text{ and the right black circle is labeled } 3. + \text{Diagram: } \text{Three external lines meeting at a central gray circle, which is connected to two black circles, which are further connected to each other. The left black circle is labeled } 3 \text{ and the right black circle is labeled } 1. + \text{Diagram: } \text{Three external lines meeting at a central gray circle, which is connected to two black circles, which are further connected to each other. The left black circle is labeled } 1 \text{ and the right black circle is labeled } 3. + \text{Diagram: } \text{Three external lines meeting at a central gray circle, which is connected to two black circles, which are further connected to each other. Both black circles are labeled } 1.$$

Amputation

- Imaginary part of Isobar-Spectator amplitude



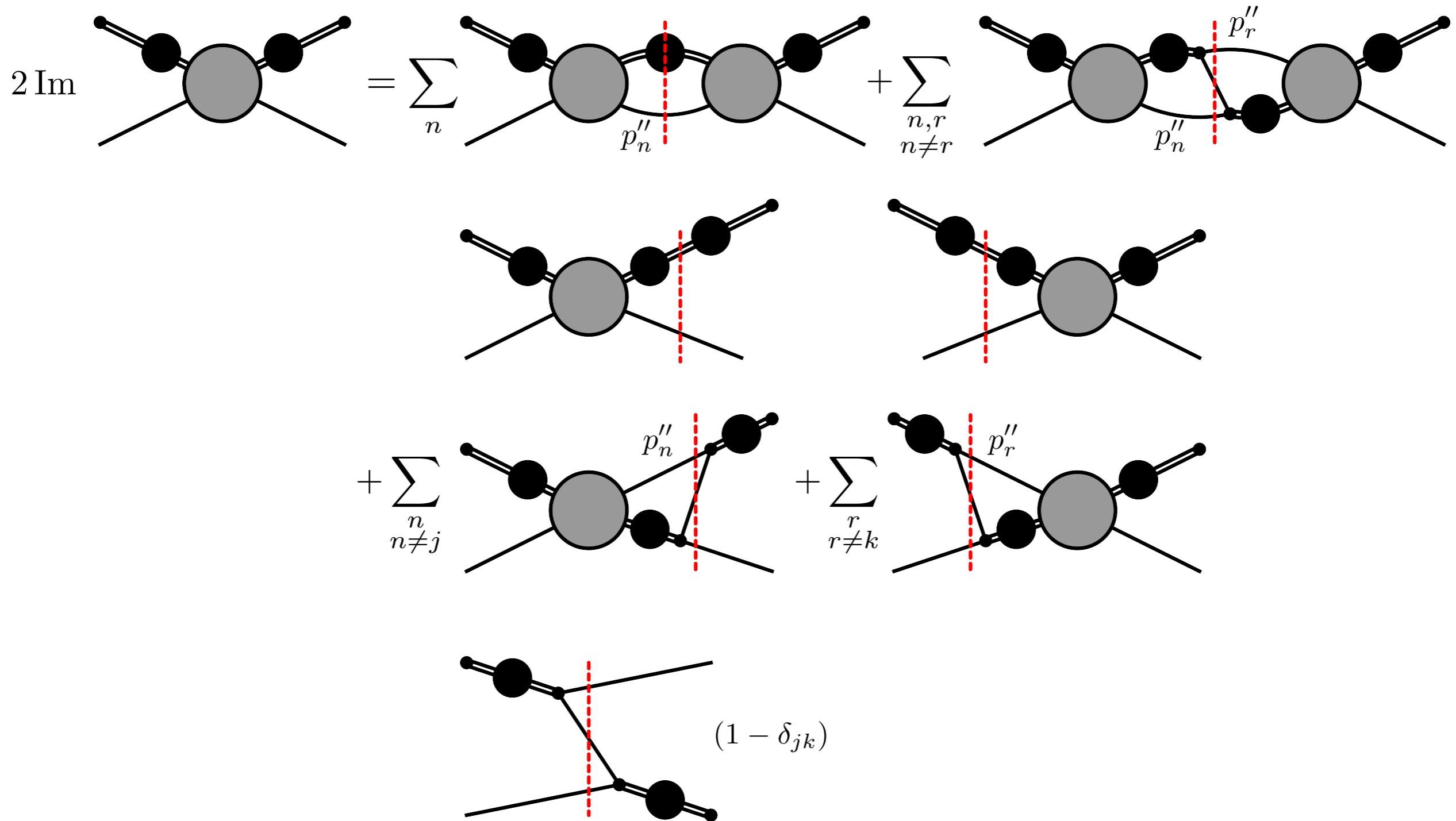
- Use unitarity for partial wave $2 \rightarrow 2$ amplitude 
- Factorize the 2-particle sub-channel denominators



$$\mathcal{A}_{kj} = f_k(\sigma'_k) \tilde{\mathcal{A}}_{kj}(\sigma'_k, s, \sigma_j) f_j(\sigma_j)$$

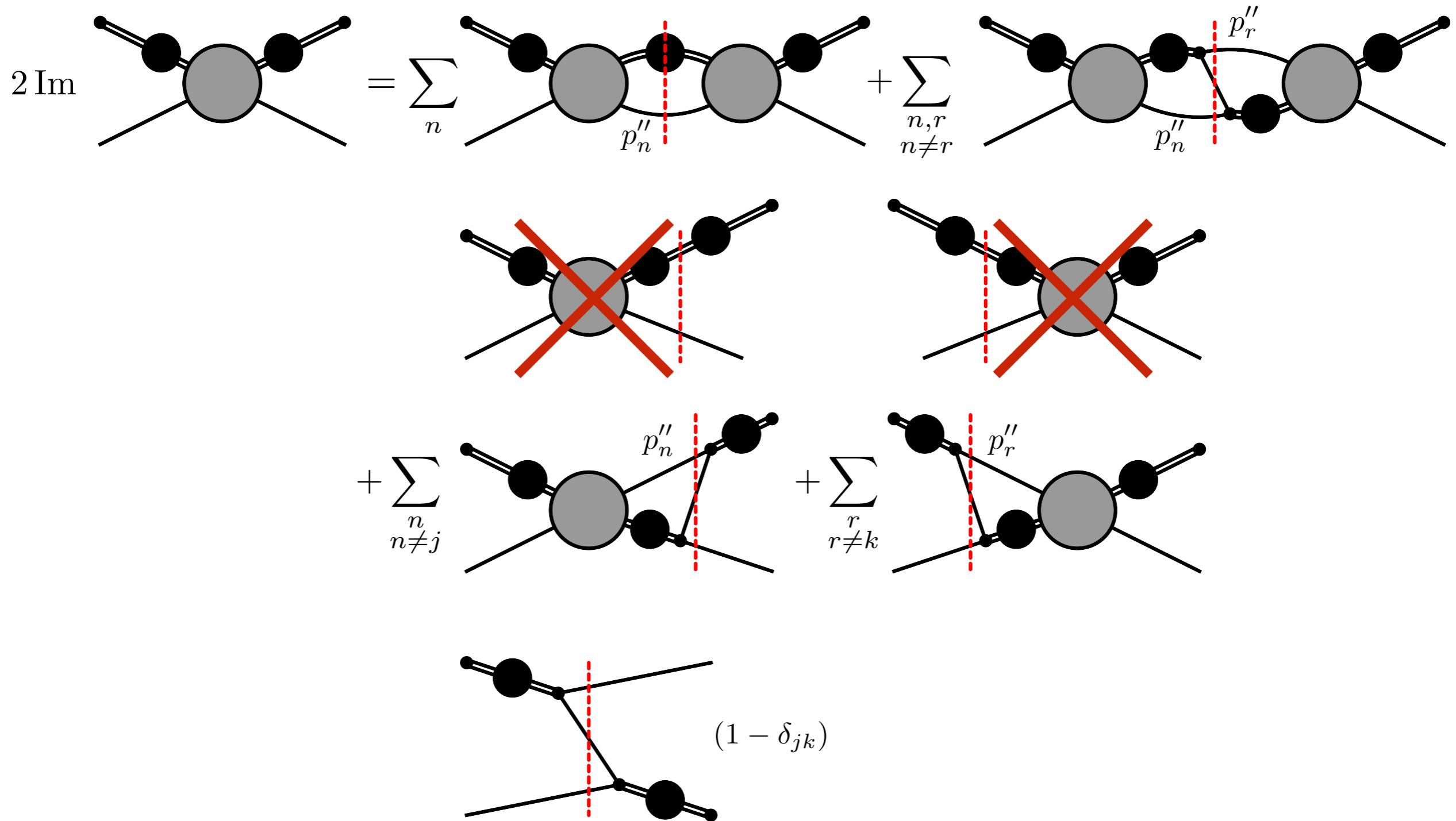
Amputated Isobar-Spectator Unitarity

- PWIS Unitarity Relations



Amputated Isobar-Spectator Unitarity

- PWIS Unitarity Relations



Amputated Isobar-Spectator Unitarity

- PWIS Unitarity Relations

$$2 \operatorname{Im} \text{ (Diagram A)} = \sum_n \text{ (Diagram B)} + \sum_{\substack{n,r \\ n \neq r}} \text{ (Diagram C)} + \sum_{\substack{n \\ n \neq j}} \text{ (Diagram D)} + \sum_{\substack{r \\ r \neq k}} \text{ (Diagram E)}$$

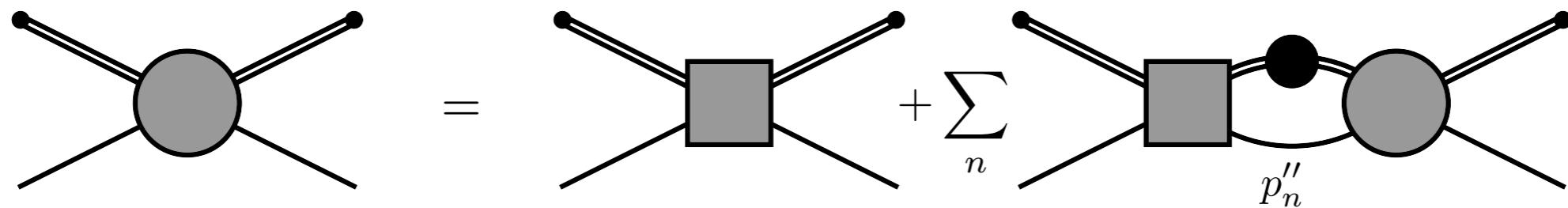
$(1 - \delta_{jk})$

B-Matrix Parameterization

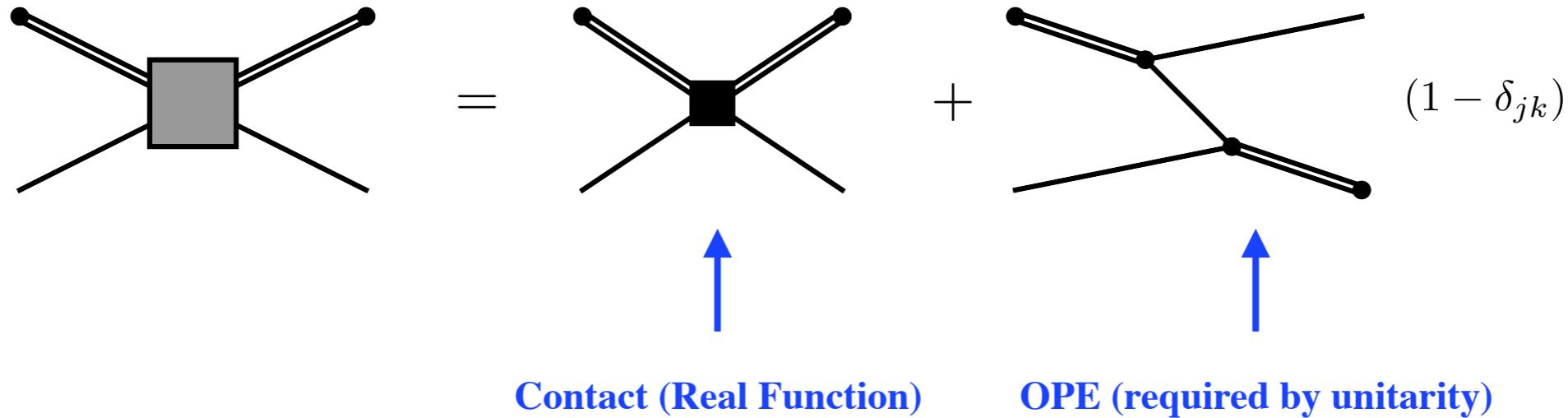
M. Mai et al. [JPAC],
Eur. Phys. J. A53 (2017) no.9, 177

- Unitary parameterization of PWIS amplitude

$$\tilde{\mathcal{A}}_{kj} = \mathcal{B}_{kj} + \sum_n \int \mathcal{B}_{kn} \tau_n \tilde{\mathcal{A}}_{nj}$$



- B-Matrix composed of OPE and Contact

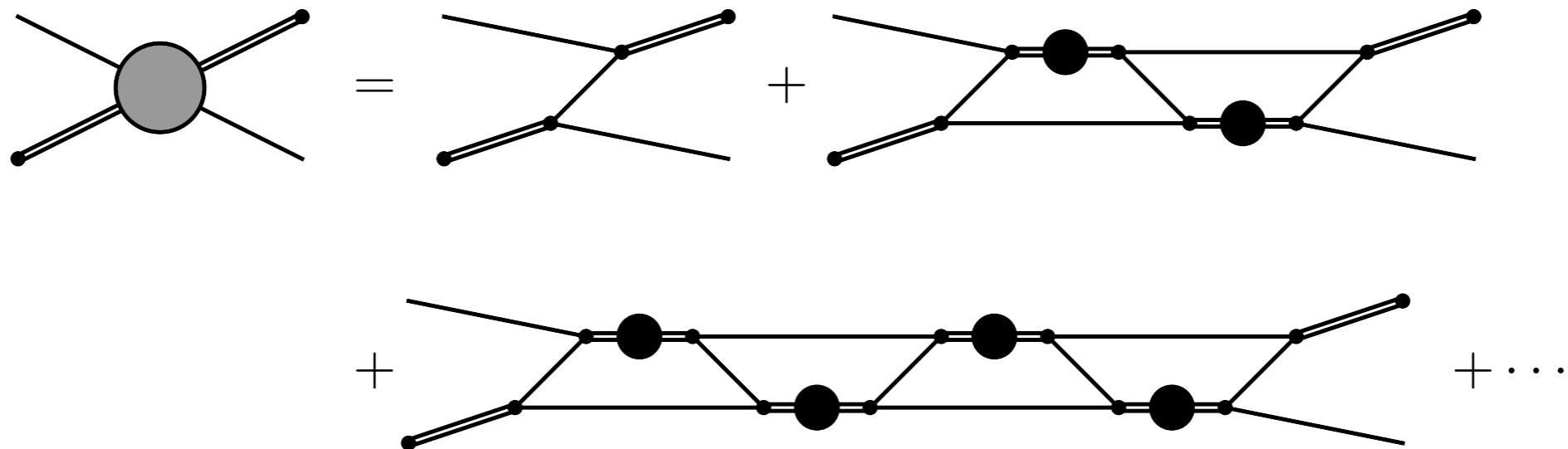


Bubbles, Triangles, and Boxes

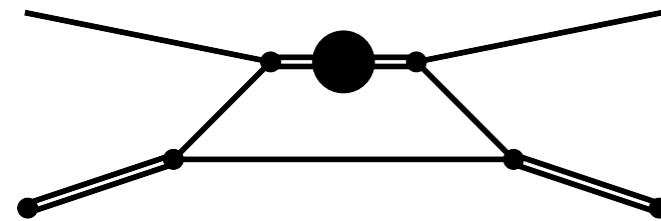
- Ladder summation of exchanges

e.g. OPE driving term

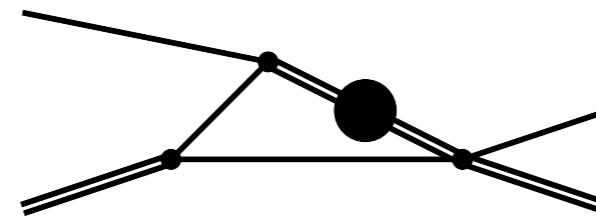
$$\tilde{\mathcal{A}}_{13} = [1 - \mathcal{K}_{11} \cdot \tau_1]^{-1} \tilde{\mathcal{E}}_{13}$$



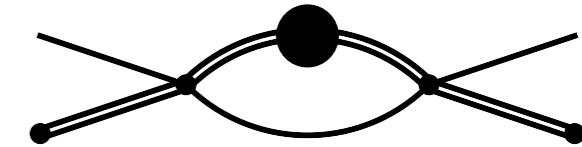
- Analytic properties of OPE, bubble, triangle, box



Box



Triangle



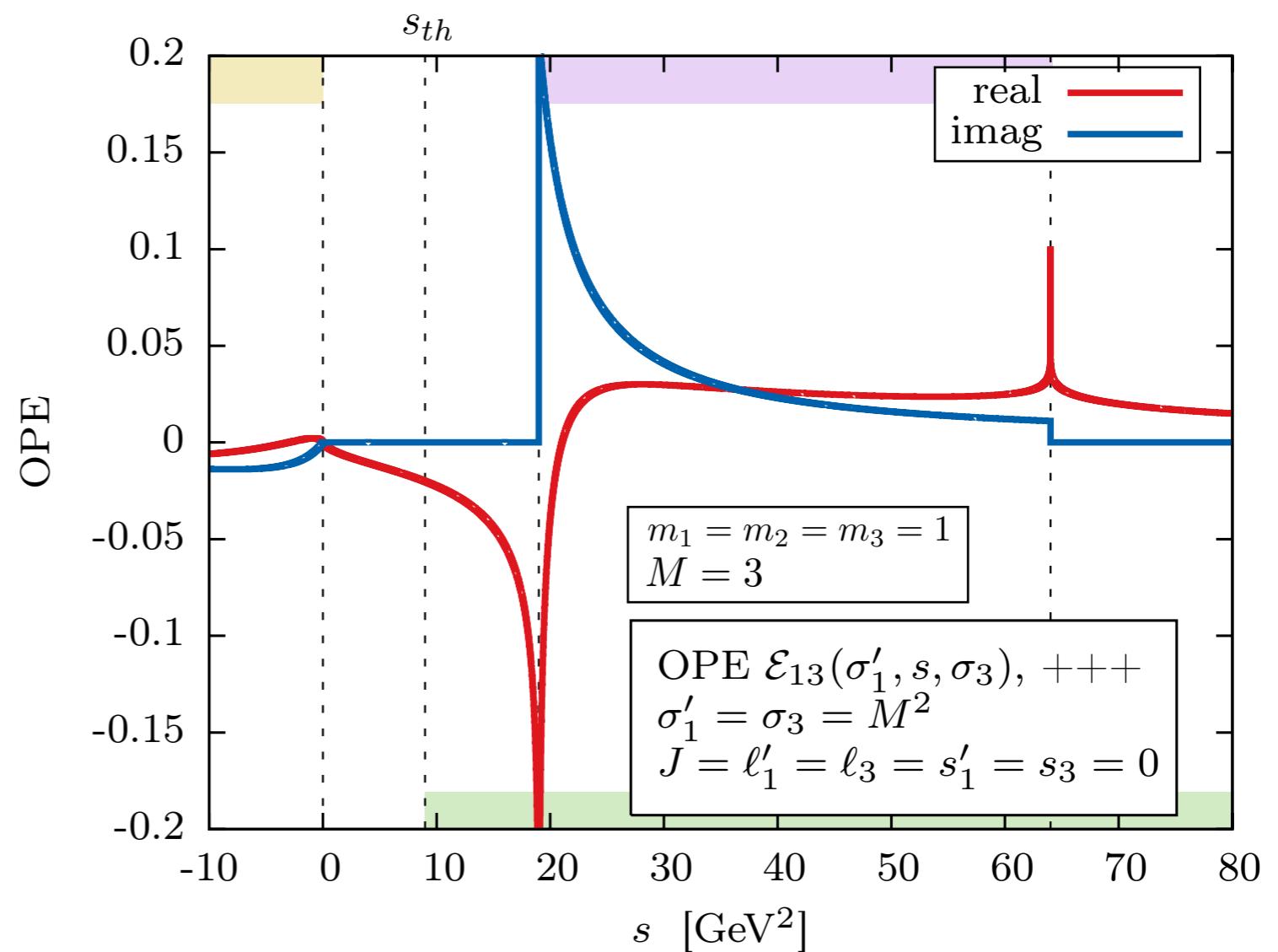
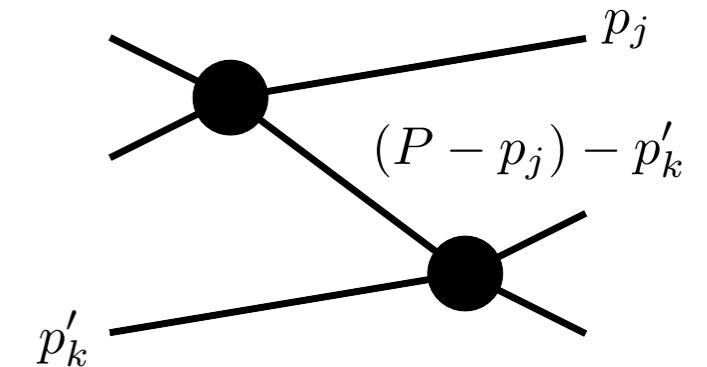
Bubble

One-Particle Exchange

- OPE \Rightarrow Driving term for ladder summation

$$\tilde{\mathcal{E}}_{13}(\sigma'_1, s, \sigma_3) = -\frac{1}{8\pi|\mathbf{p}'_1{}^\star||\mathbf{p}_3{}^\star|} \log \left(\frac{1 + z_{13}^\mu}{1 - z_{13}^\mu} \right) \quad \mathcal{E}_{kj}(\mathbf{p}'; \mathbf{p}) =$$

$$z_{13}^\mu = \frac{\sigma_3 + m_1^2 - m_2^2 - E_3^* \omega_1'{}^*}{|\mathbf{p}'_1{}^\star||\mathbf{p}_3{}^\star|}$$

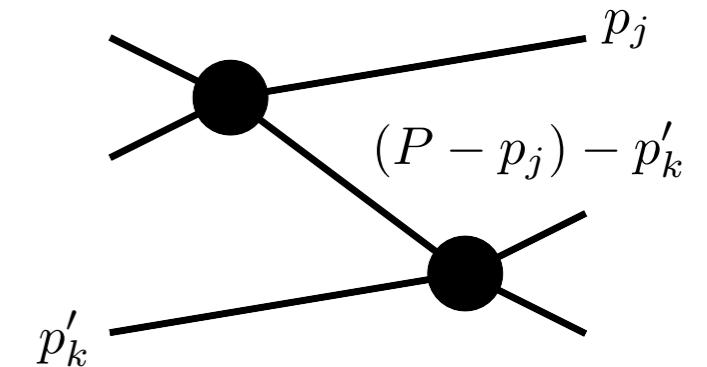


One-Particle Exchange

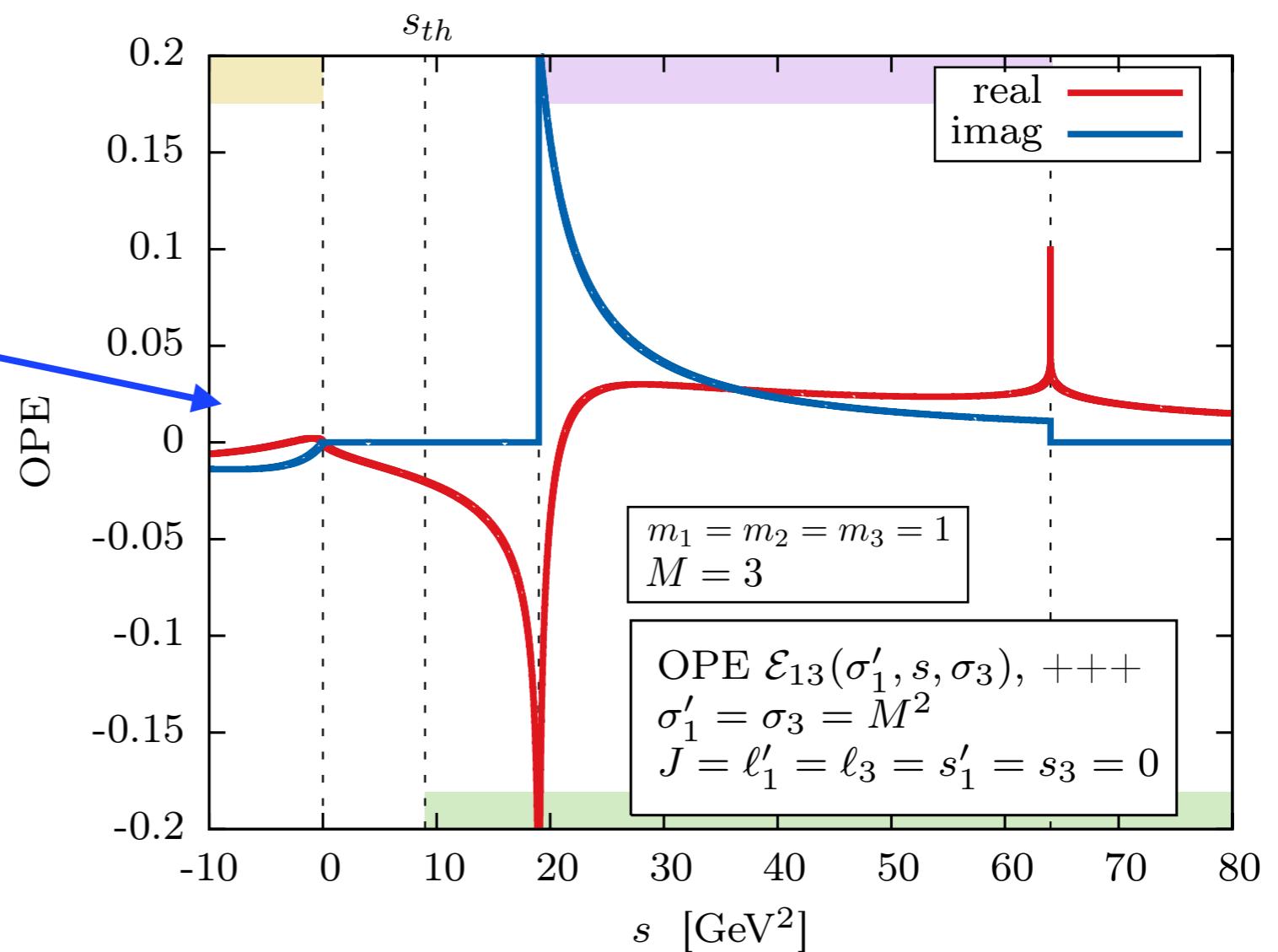
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$$z_{13}^\mu = \frac{\sigma_3 + m_1^2 - m_2^2 - E_3^* \omega_1'{}^*}{|\mathbf{p}'_1{}^\star||\mathbf{p}_3{}^\star|}$$



Virtual Particle Exchange cut

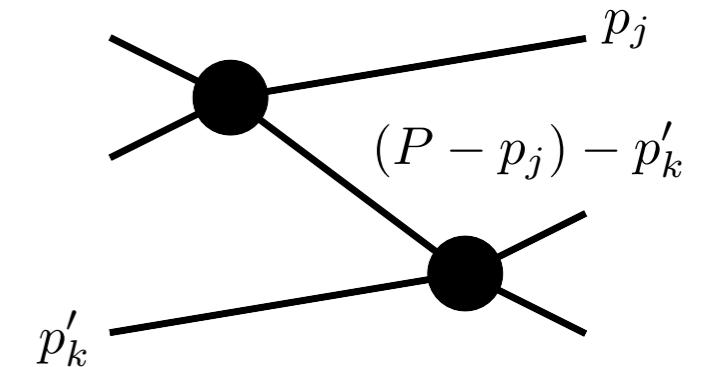


One-Particle Exchange

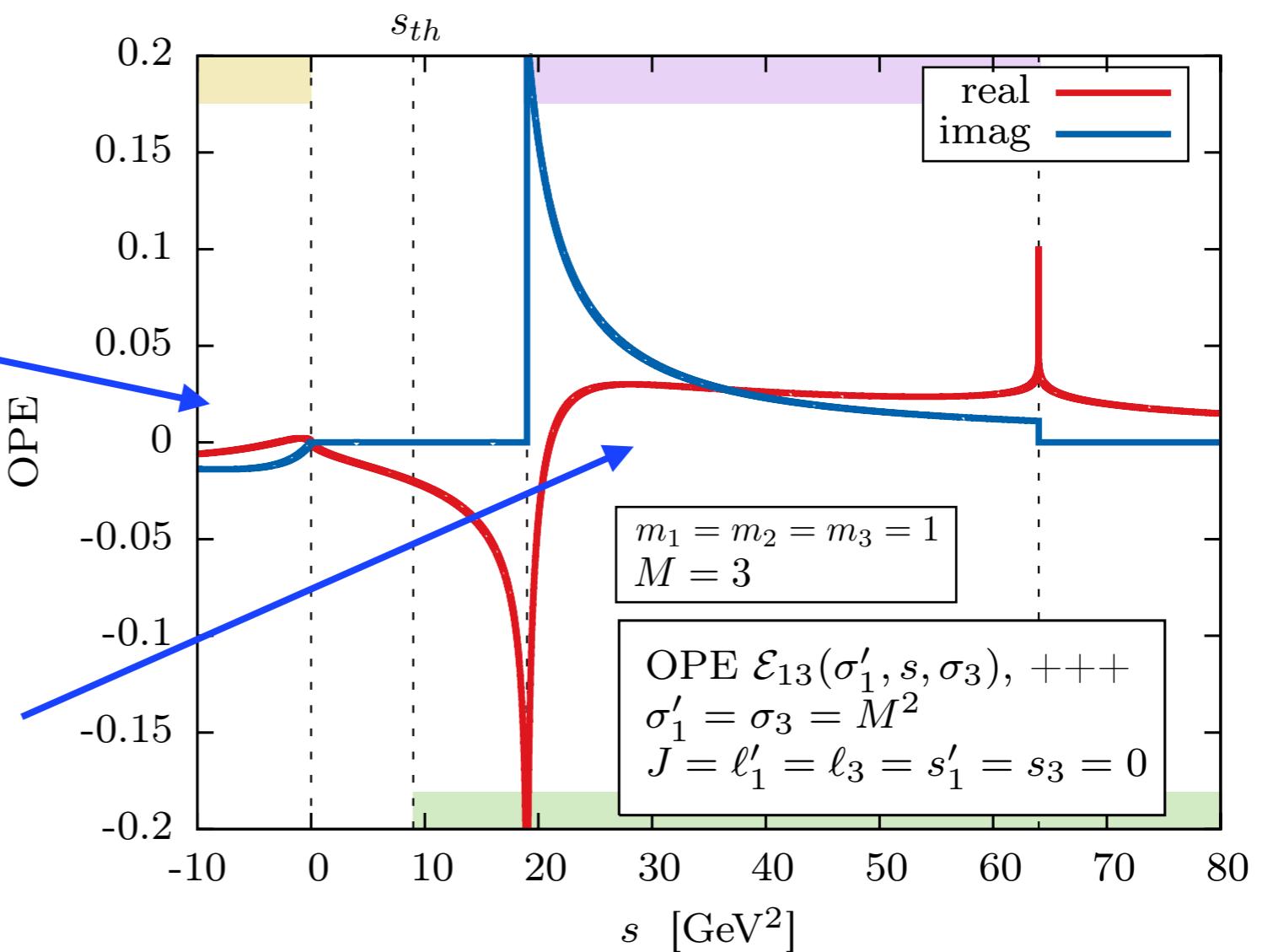
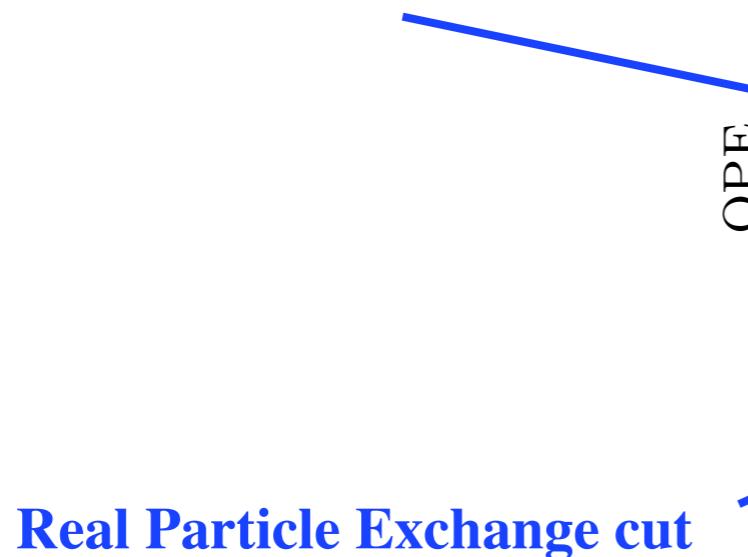
- OPE \Rightarrow Driving term for ladder summation

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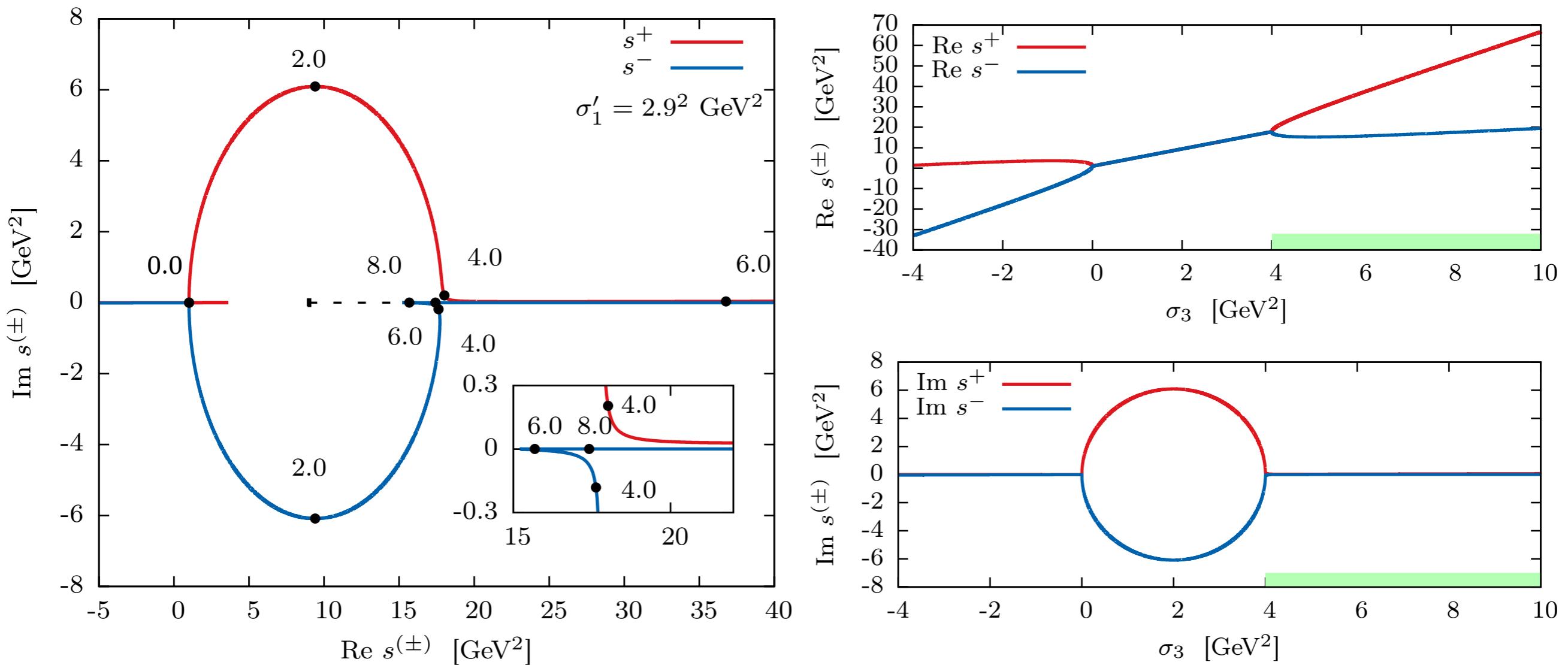
$$z_{13}^\mu = \frac{\sigma_3 + m_1^2 - m_2^2 - E_3^* \omega_1'{}^*}{|\mathbf{p}'_1{}^\star||\mathbf{p}_3{}^\star|}$$



Virtual Particle Exchange cut

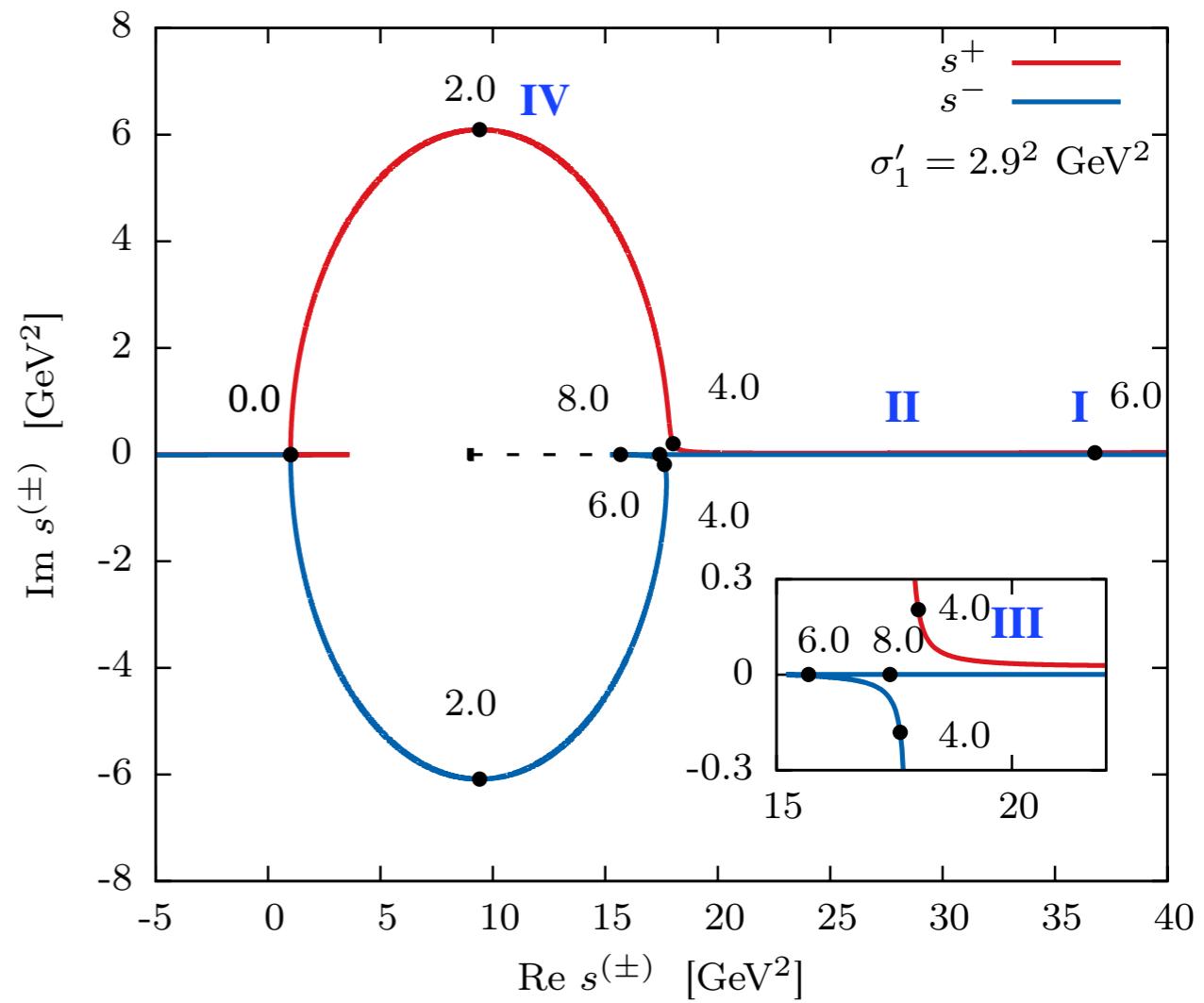
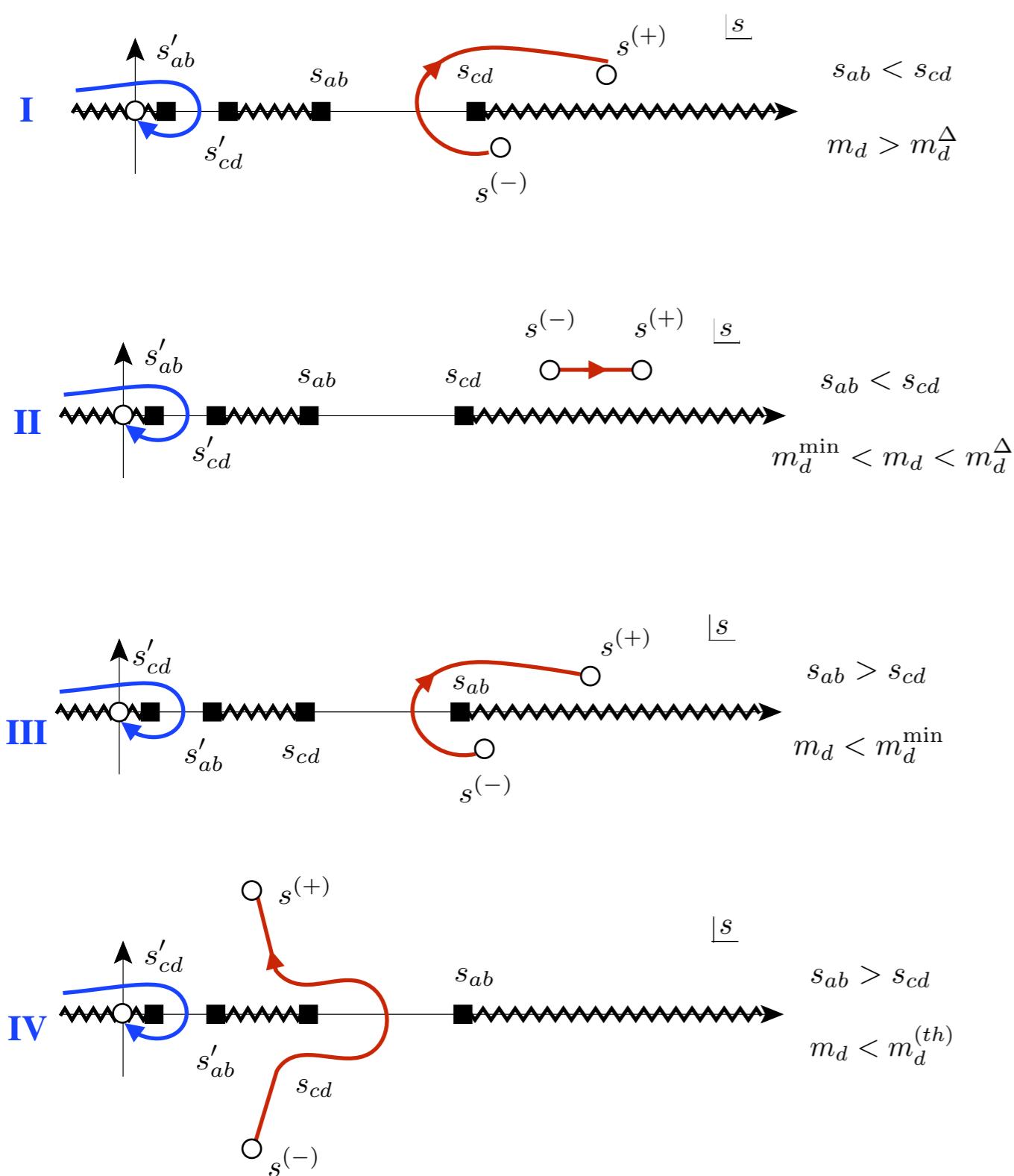


Analytic Structure



$$\tilde{\mathcal{E}}_{13}(\sigma'_1, s, \sigma_3) = \frac{1}{2\pi i} \int_{-\infty}^0 ds' \frac{\Delta_s \tilde{\mathcal{E}}_{13}(\sigma'_1, s', \sigma_3)}{s' - s} + \frac{1}{2\pi i} \int_{s(-)}^{s(+)} ds' \frac{\Delta_s \tilde{\mathcal{E}}_{13}(\sigma'_1, s', \sigma_3)}{s' - s}$$

Analytic Structure



Triangle Diagram

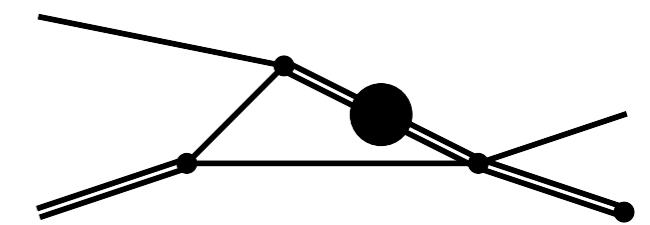
I.J.R. Aitchison and R. Pasquier
Phys. Rev. 152, 4 (1966)

- Triangle diagram - Comparison to off-shell amplitude
 - Imaginary parts equal, difference in the real part

$$\tilde{\mathcal{A}}_{\Delta}(\sigma', s, \sigma) = \int_{-\infty}^{(\sqrt{s}-m)^2} d\sigma'' \tilde{\mathcal{E}}(\sigma', s, \sigma'') \tau(s, \sigma'')$$

Off-shell triangle

$$\mathcal{R}(\sigma', s, \sigma) \rightarrow \text{constant}$$



$$\tilde{\mathcal{A}}_{\Delta}^B(\sigma', s, \sigma) = \int_{\sigma^{(th)}}^{(\sqrt{s}-m)^2} d\sigma'' \tilde{\mathcal{E}}(\sigma', s, \sigma'') \tau(s, \sigma'')$$

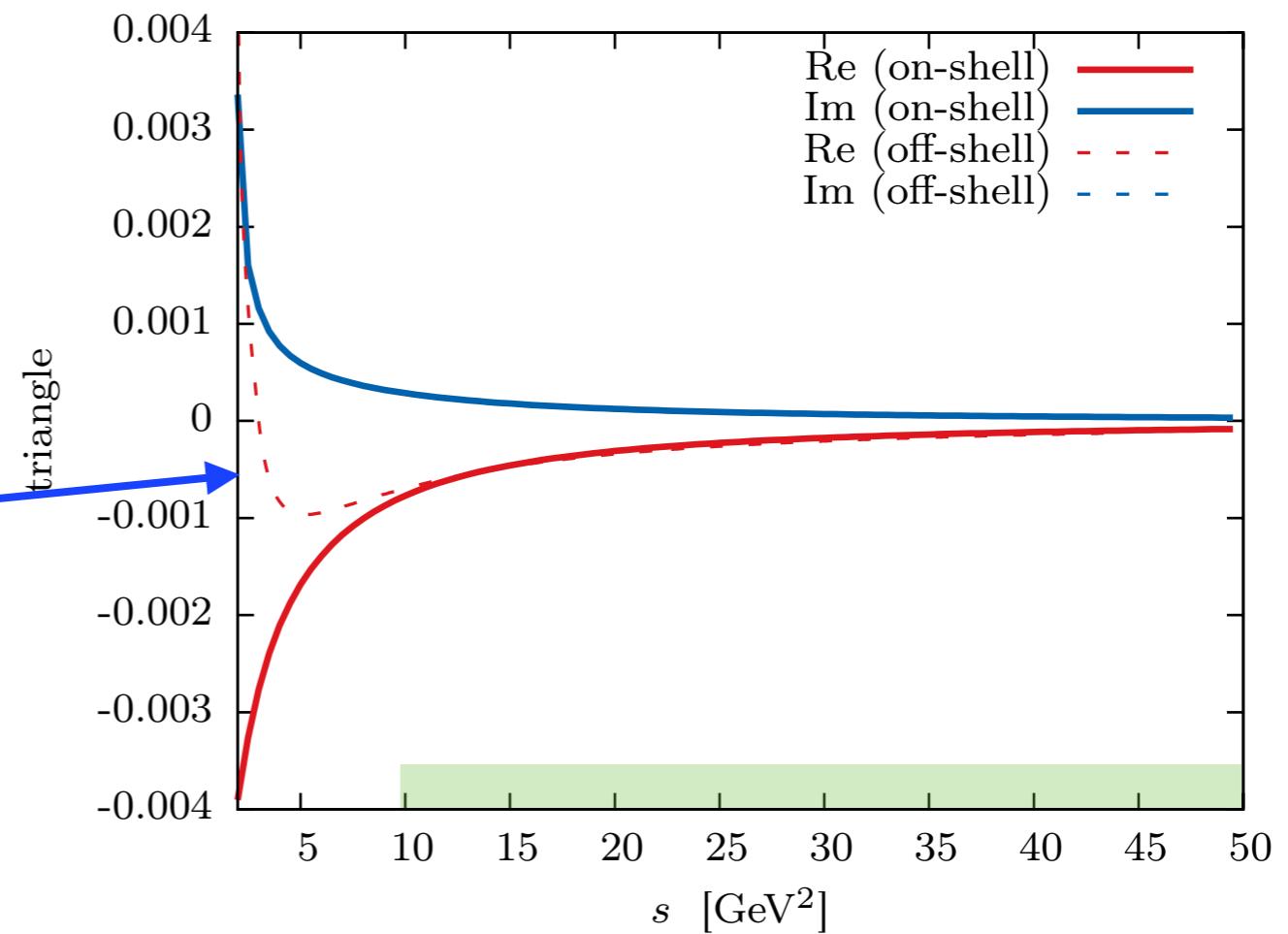
On-shell triangle

$$\tau(s, \sigma) = \frac{1}{\pi} \rho(s, \sigma) f(\sigma)$$

e.g. stable particle

$$\text{Im } \tilde{\mathcal{A}}_{\Delta}(\sigma', s, \sigma) = \text{Re } \tilde{\mathcal{A}}_{\Delta} + i \tilde{\mathcal{E}}(\sigma', s, M^2) \rho(s, M^2)$$

Difference is real part

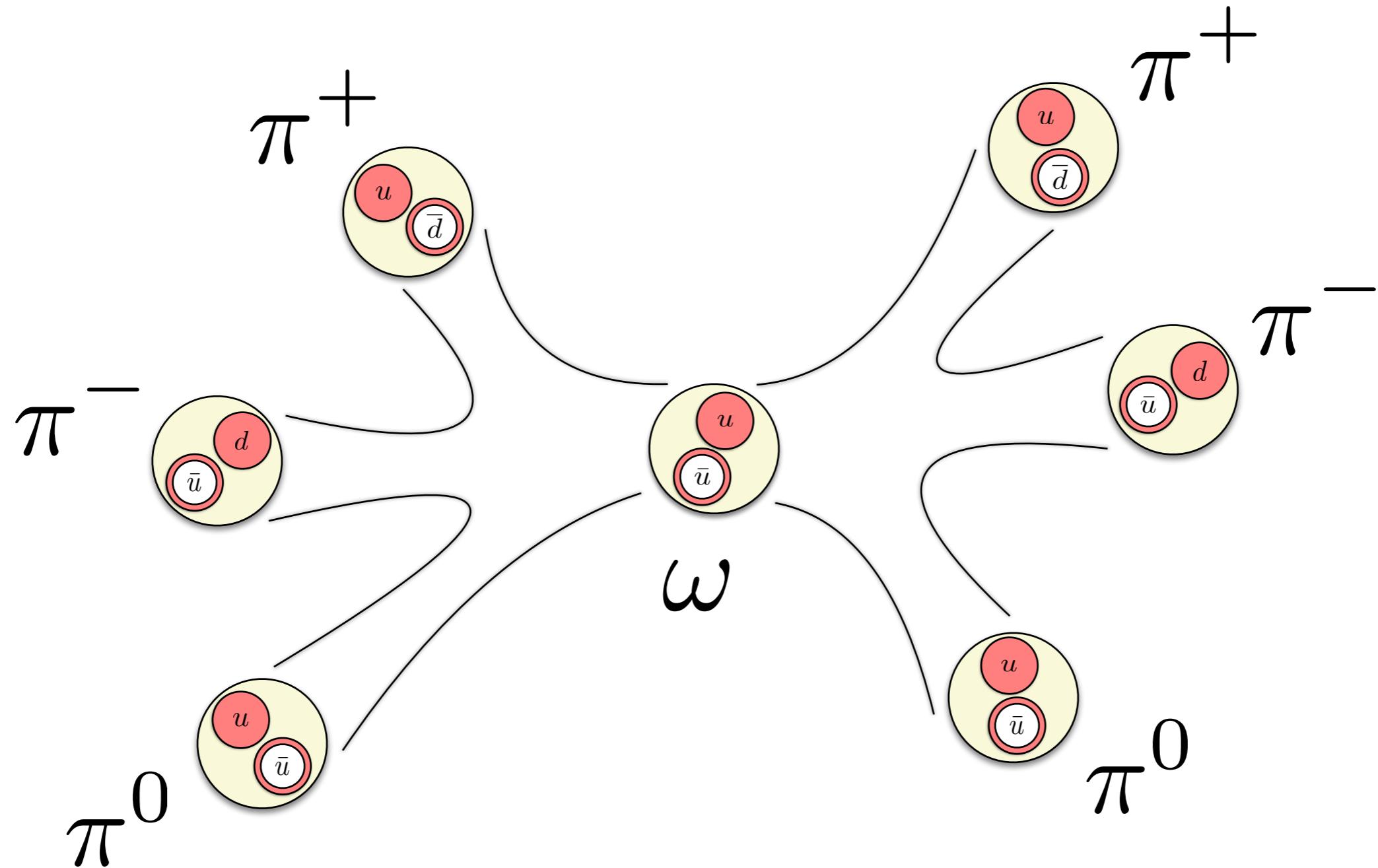


Outlook and Future Directions

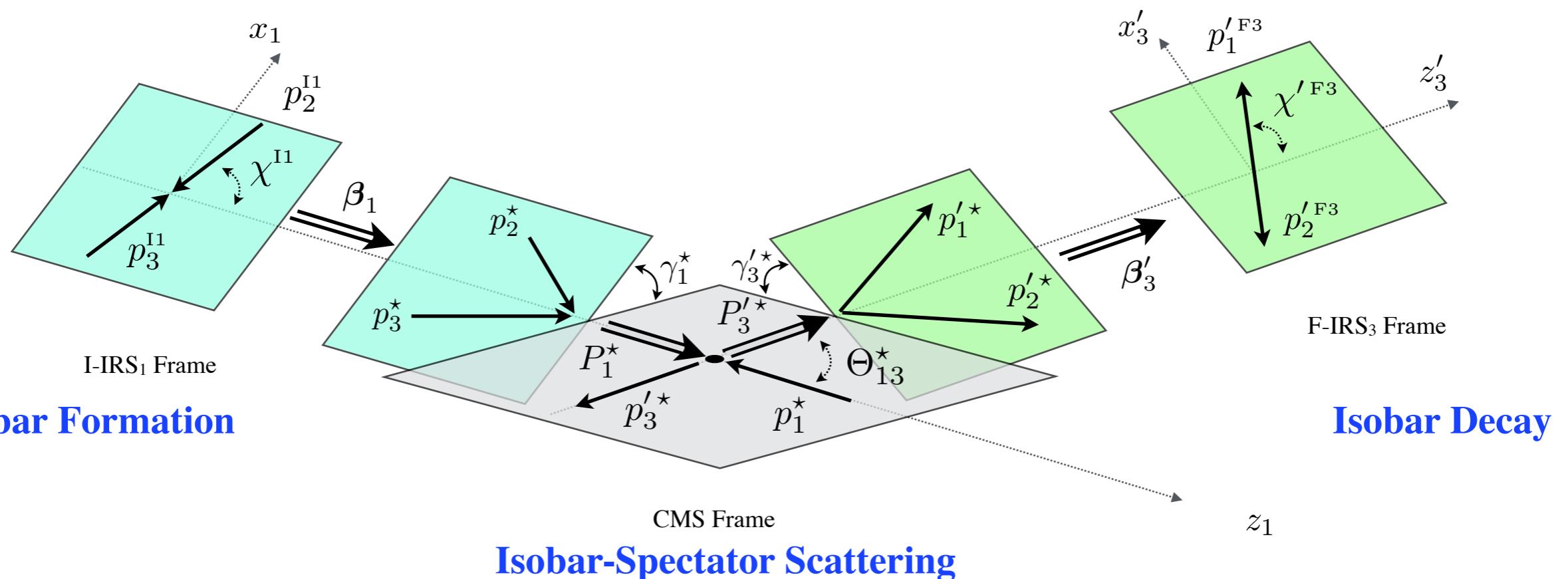
- $3 \rightarrow 3$ unitarity relations within the isobar model
 - Spinless external particles
 - Finite number of isobars with spin
- B-matrix parameterization
 - Satisfies unitarity
 - Numerical test underway to determine feasibility of solution
 - Studying behavior when analytically continued
- Analytic Extensions
 - Can we use the B-matrix to write dispersive representation?
 - Continual Investigation into OPE, Box, Triangle, Bubble

AJ et al [JPAC] - In Preparation

Back-up



Kinematics



Unitarity Relations

Disconnected 2→2 Unitarity Relation

$$\text{Im } \mathcal{F}_j(\mathbf{p}'; \mathbf{p}) = \frac{|\mathbf{q}^{(j)}|}{32\pi^2 \sqrt{\sigma_j}} \int d\Sigma''^{(j'')} \mathcal{F}_j^*(\mathbf{p}''; \mathbf{p}') \mathcal{F}_j(\mathbf{p}''; \mathbf{p}) \Theta(\sigma_j - \sigma_j^{(th)})$$

Connected 3→3 Unitarity Relation

$$\begin{aligned} \text{Im } \mathcal{A}(\mathbf{p}'; \mathbf{p}) &= \frac{1}{2(2\pi)^5} \int \frac{d^3 \mathbf{p}_1''}{2\omega(\mathbf{p}_1'')} \frac{d^3 \mathbf{p}_2''}{2\omega(\mathbf{p}_2'')} \frac{d^3 \mathbf{p}_3''}{2\omega(\mathbf{p}_3'')} \delta^{(4)}(P'' - P) \mathcal{A}^*(\mathbf{p}''; \mathbf{p}') \mathcal{A}(\mathbf{p}''; \mathbf{p}) \\ &+ \sum_k \frac{|\mathbf{q}'^{(k')}|}{32\pi^2 \sqrt{\sigma'_k}} \int d\Sigma''^{(k'')} \mathcal{F}_k^*(\mathbf{p}''; \mathbf{p}') \mathcal{A}(\mathbf{p}''; \mathbf{p})|_{\mathbf{p}_k''=\mathbf{p}'_k} \Theta(\sigma'_k - \sigma_k^{(th)}) \\ &+ \sum_j \frac{|\mathbf{q}^{(j)}|}{32\pi^2 \sqrt{\sigma_j}} \int d\Sigma''^{(j'')} \mathcal{A}^*(\mathbf{p}''; \mathbf{p}')|_{\mathbf{p}_j''=\mathbf{p}_j} \mathcal{F}_j(\mathbf{p}''; \mathbf{p}) \Theta(\sigma_j - \sigma_j^{(th)}) \\ &+ \sum_{\substack{j, k \\ j \neq k}} \pi \delta(u_{jk} - \mu^2) \mathcal{F}_k^*(\mathbf{p}''; \mathbf{p}')|_{\mathbf{p}_j''=\mathbf{p}_j} \mathcal{F}_j(\mathbf{p}''; \mathbf{p})|_{\mathbf{p}_k''=\mathbf{p}'_k}, \end{aligned}$$

Unitarity Relations

Disconnected 2→2 Unitarity Relation

$$2 \operatorname{Im} \text{ (Diagram)} = \text{ (Diagram with red dashed vertical line)}$$

Connected 3→3 Unitarity Relation

$$\begin{aligned} 2 \operatorname{Im} \text{ (Diagram)} &= \text{ (Diagram with red dashed vertical line)} \\ + \sum_j &\quad \text{ (Diagram with red dashed vertical line, p}_j\text{)} + \sum_k \text{ (Diagram with red dashed vertical line, p}'_k\text{)} \\ + \sum_{\substack{j, k \\ j \neq k}} &\quad \text{ (Diagram with red dashed vertical line, p}_j\text{, p}'_k\text{, } (P - p_j) - p'_k\text{)} \end{aligned}$$

Unitarity of B-Matrix

