

# Controlling uncertainties in electroweak precision studies

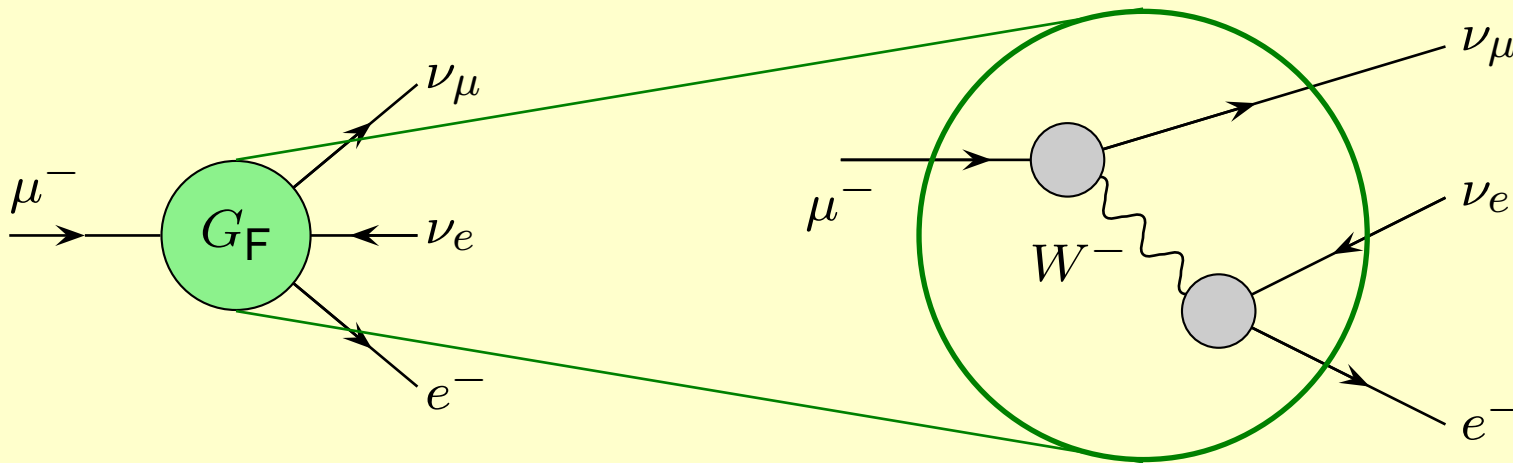
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- “Other” electroweak parameters needed for global fit
- Theory error estimates
- Controlling uncertainties in EW precision tests
- Role of beam polarization

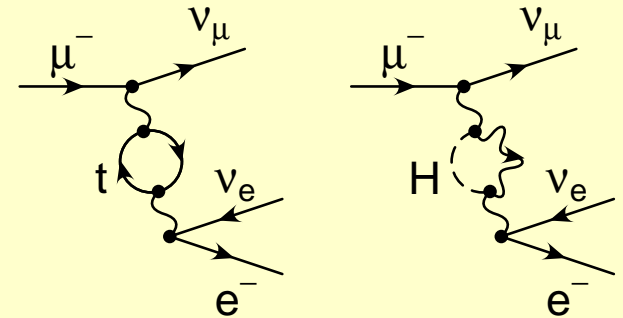
W-boson mass can be calculated from muon decay rate:



$\mu$  decay in Standard Model

$$\frac{G_F^2}{\sqrt{2}} = \frac{e^2}{8s_w^2 M_W^2} (1 + \Delta r)$$

electroweak corrections



- Dependence on  $M_Z, m_t, \Delta\alpha, \dots$
- Comparison to direct measurement of  $M_W$

$e^+e^- \rightarrow f\bar{f}$  for  $E_{\text{CM}} \sim M_Z$ :

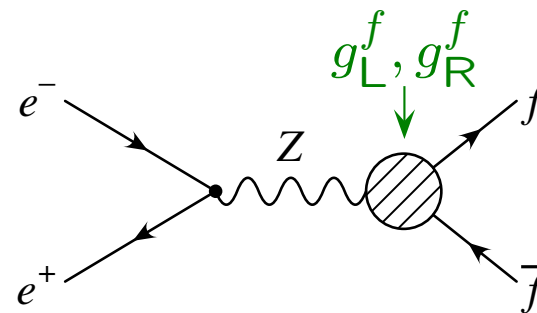
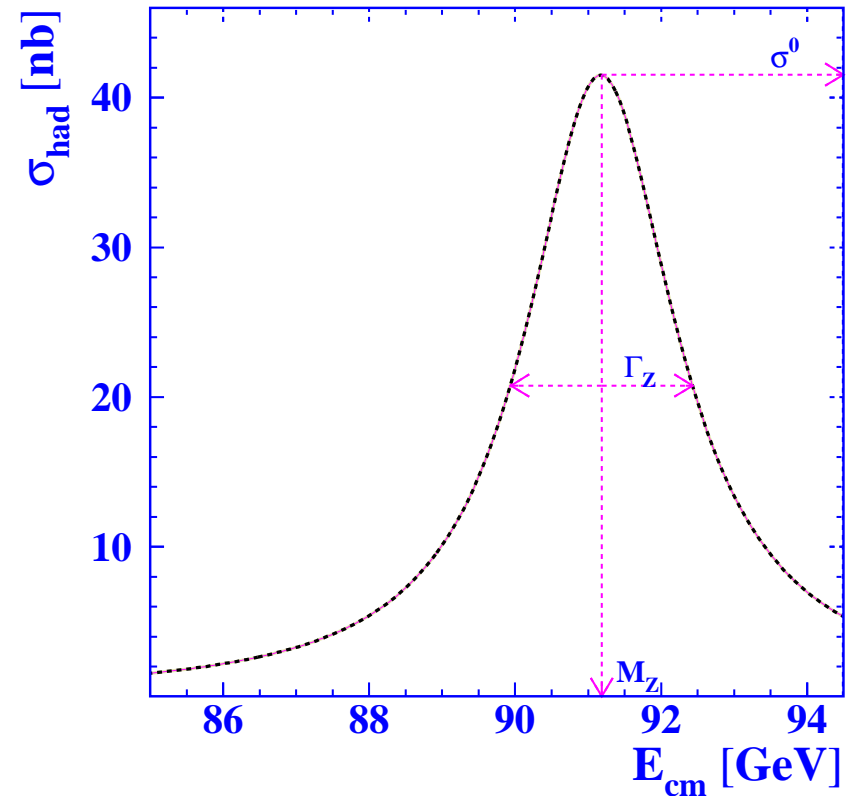
- Mass  $M_Z$
- Width  $\Gamma_Z = \sum_f \Gamma_{ff}$
- Branching ratio  $R_f = \Gamma_{ff}/\Gamma_Z$
- $\sigma^0 \approx \frac{12\pi \Gamma_{ee} \Gamma_{ff}}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2} = \frac{12\pi}{M_Z^2} R_e R_f$

$$\Gamma_{ff} = C [(g_L^f)^2 + (g_R^f)^2]$$

→ Dependence on  $\alpha_s, m_t, M_Z, \dots$

Main exp. systematic uncertainties:

- Beam energy calibration
- Particle ID and acceptance



Forward-backward asymmetry:

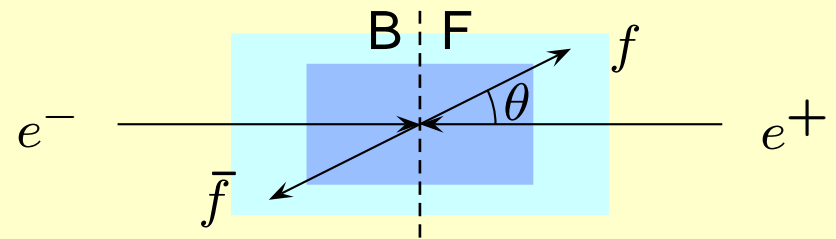
$$A_{FB} \equiv \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f$$

$$\mathcal{A}_f = \frac{2(1 - 4\sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4\sin^2 \theta_{\text{eff}}^f)^2}$$

$$\sin^2 \theta_{\text{eff}}^f = \frac{g_R^f}{2|Q_f|(g_R^f - g_L^f)}$$

Main systematic uncertainties:

- For  $f = b$ : charge tagging, jet clustering
- For  $f = \mu$ : calibration of  $\sqrt{s}$ , muon angle



Polarization asymmetry:

Average  $\tau$  pol. in  $e^+e^- \rightarrow \tau^+\tau^-$ ,  $\langle \mathcal{P}_\tau \rangle = -\mathcal{A}_\tau$

Left-right asymmetry: (using polarization  $e^-$  beams)

$$A_{LR} \equiv \frac{1}{P_{e^-}} \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = \mathcal{A}_e + \Delta A_{\gamma Z} + \Delta A_{\gamma}$$

$$\mathcal{A}_f = \frac{2(1 - 4\sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4\sin^2 \theta_{\text{eff}}^f)^2} \quad \sin^2 \theta_{\text{eff}}^f = \frac{g_R^f}{2|Q_f|(g_R^f - g_L^f)}$$

Limited by systematic uncertainty of  $P_{e^-}$   
0.5% at SLD, 0.1% possible in future

Karl, List '17

Blondel scheme: (if  $e^-$  and  $e^+$  polarization available)

Blondel '88

Four independent measurements for  $P_{e^+}/P_{e^-} = ++, +-, -+, --$

$$A_{LR} = \sqrt{\frac{(\sigma_{++} + \sigma_{-+} - \sigma_{+-} - \sigma_{--})(-\sigma_{++} + \sigma_{-+} - \sigma_{+-} + \sigma_{--})}{(\sigma_{++} + \sigma_{-+} + \sigma_{+-} + \sigma_{--})(-\sigma_{++} + \sigma_{-+} + \sigma_{+-} - \sigma_{--})}}$$

**Note:** No need to know  $|P_{e^\pm}|$  !

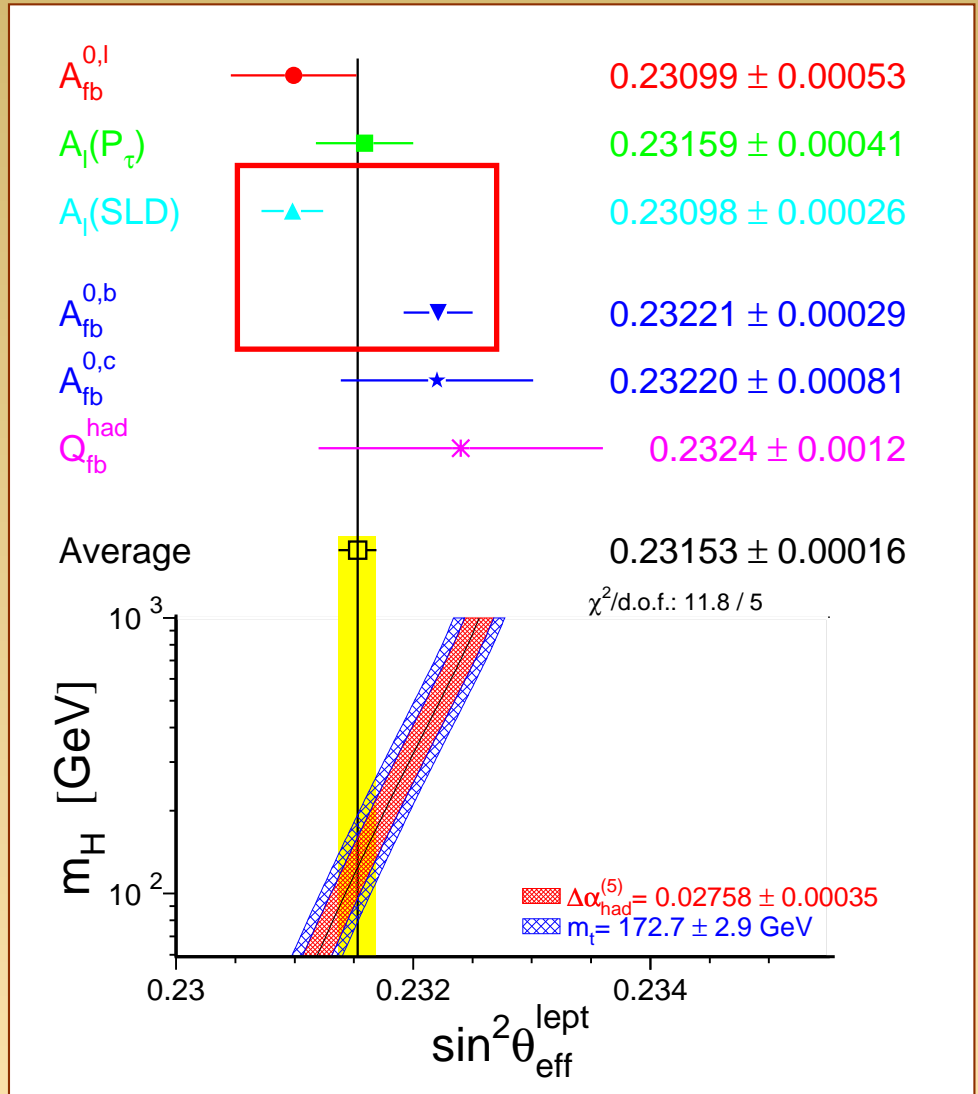
Main systematic uncertainties:

- Difference of  $|P|$  for  $P > 0$  and  $P < 0$
- Difference of  $\mathcal{L}$  for  $P > 0$  and  $P < 0$

$$\delta A_{LR} \approx 10^{-4} \quad \Rightarrow \quad \delta \sin^2 \theta_{\text{eff}}^l \approx 1.3 \times 10^{-5}$$

Mönig, Hawkings '99

- Most precise determination from  $A_{\text{LR}}$  at SLD and  $A_{\text{FB}}^b$  at LEP
- Disagreement by  $\sim 4\sigma$   
→ Underestimated systematics?
- Default at CEPC, FCC-ee:  
 $\sin^2 \theta_{\text{eff}}^l$  from  $A_{\text{FB}}^{\mu\mu}$



	Current exp.	ILC/GigaZ	CEPC	FCC-ee
$M_W$ [MeV]	15	3–4	1	1
$M_Z$ [MeV]	2.3	–	0.5	0.1
$\Gamma_Z$ [MeV]	2.3	0.8	0.5	0.1
$R_\ell = \Gamma_Z^{\text{had}}/\Gamma_Z^\ell$ [ $10^{-3}$ ]	25	10	2	1
$R_b = \Gamma_Z^b/\Gamma_Z^{\text{had}}$ [ $10^{-5}$ ]	66	14	4.3	6
$\sin^2 \theta_{\text{eff}}^\ell$ [ $10^{-5}$ ]	16	1.3	<1	0.5

→ Existing theoretical calculations adequate for LEP/SLC/LHC,  
but not ILC/CEPC/FCC-ee!



- Many seminal works on 1-loop and leading 2-loop corrections

Veltman, Passarino, Sirlin, Marciano, Bardin, Hollik, Riemann, Degrassi, Kniehl, ...

- Full 2-loop results for  $M_W$ ,  $Z$ -pole observables

Freitas, Hollik, Walter, Weiglein '00

Awramik, Czakon, Freitas '06

Awramik, Czakon '02

Hollik, Meier, Uccirati '05,07

Onishchenko, Veretin '02

Awramik, Czakon, Freitas, Kniehl '08

Awramik, Czakon, Freitas, Weiglein '04

Freitas '13,14

Dubovyk, Freitas, Gluza, Riemann, Usovitsch '16,18

- Approximate 3- and 4-loop results (to  $\rho$  parameter)

Chetyrkin, Kühn, Steinhauser '95

Schröder, Steinhauser '05

Faisst, Kühn, Seidensticker, Veretin '03

Chetyrkin et al. '06

Boughezal, Tausk, v. d. Bij '05

Boughezal, Czakon '06

	Experiment	Theory error	Main source
$M_W$	$80.379 \pm 0.012$ MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
$\Gamma_Z$	$2495.2 \pm 2.3$ MeV	0.4 MeV	$\alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
$R_\ell$	$20.767 \pm 0.025$	0.005	$\alpha^3, \alpha^2\alpha_s$
$R_b$	$0.21629 \pm 0.00066$	0.0001	$\alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^\ell$	$0.23153 \pm 0.00016$	$4.5 \times 10^{-5}$	$\alpha^3, \alpha^2\alpha_s$

- Theory error estimate is not well defined, ideally  $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
- Common methods:
  - Count prefactors ( $\alpha, N_c, N_f, \dots$ )
  - Extrapolation of perturbative series
  - Renormalization scale dependence
  - Renormalization scheme dependence

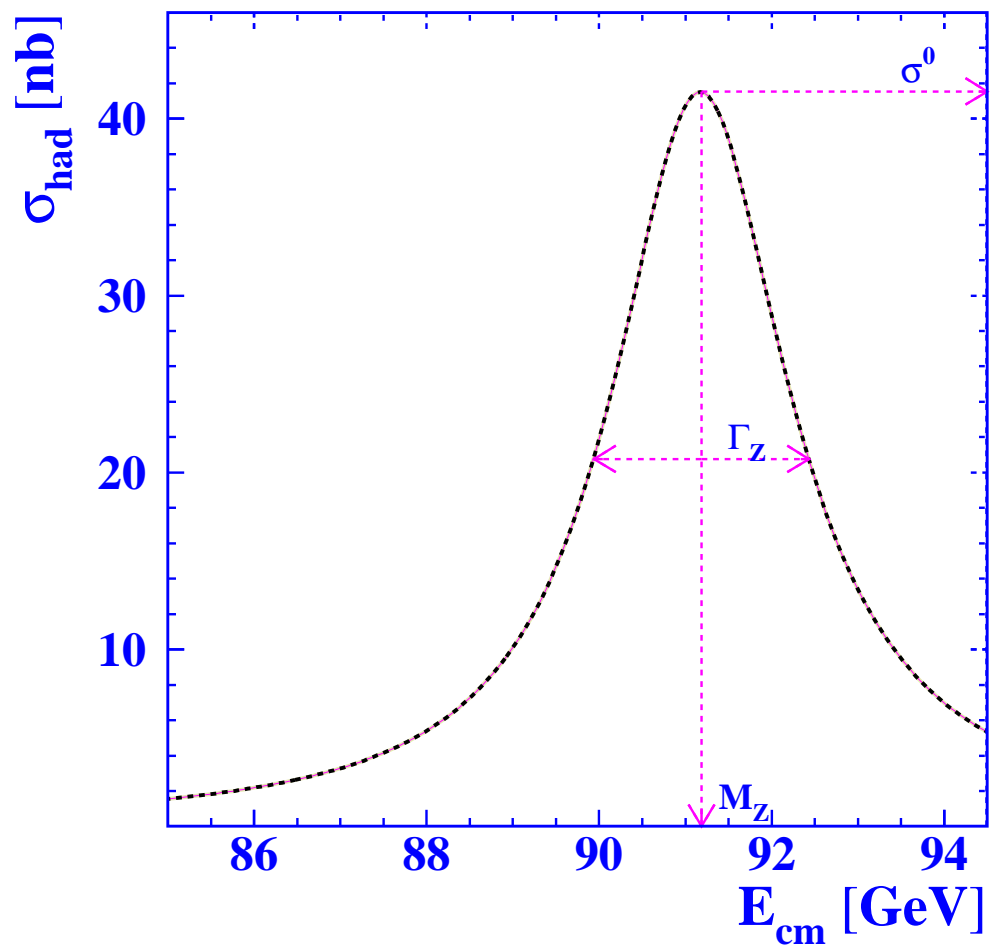
	CEPC	perturb. error with 3-loop <sup>†</sup>	Param. error CEPC**	main source
$M_W$ [MeV]	1	1	2.1	$m_t, \Delta\alpha$
$\Gamma_Z$ [MeV]	0.5	0.15	0.15	$m_t, \alpha_s$
$R_b$ [ $10^{-5}$ ]	4.3	5	$< 1$	
$\sin^2 \theta_{\text{eff}}^l$ [ $10^{-5}$ ]	$< 1$	1.5	2	$m_t, \Delta\alpha$

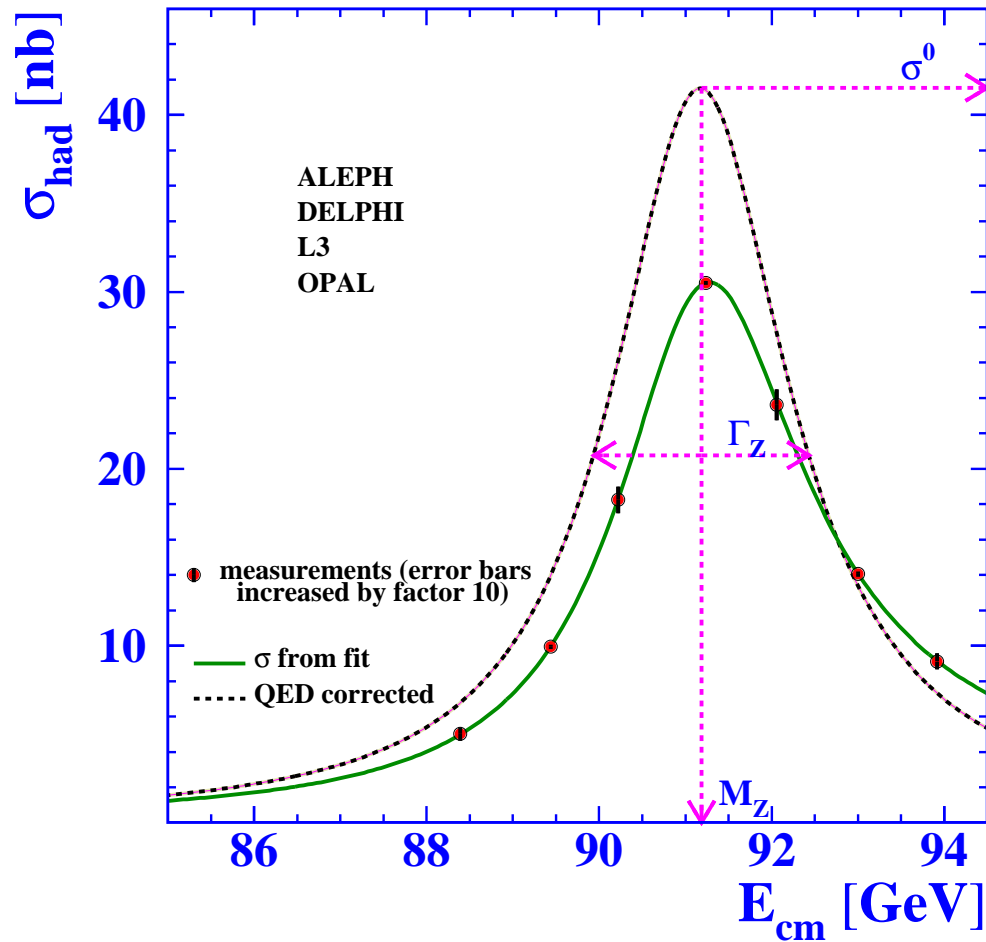
<sup>†</sup> **Theory scenario:**  $\mathcal{O}(\alpha\alpha_s^2)$ ,  $\mathcal{O}(N_f\alpha^2\alpha_s)$ ,  $\mathcal{O}(N_f^2\alpha^2\alpha_s)$   
 ( $N_f^n$  = at least  $n$  closed fermion loops)

Parametric inputs:

**\*\*CEPC:**  $\delta m_t = 600$  MeV,  $\delta\alpha_s = 0.0002$ ,  $\delta M_Z = 0.5$  MeV

also:  $\delta(\Delta\alpha) = 5 \times 10^{-5}$





LEP EWWG '05

- Large effects from initial-state QED radiation
- Theory input necessary to extract relevant EWPOs (“pseudo-observables”)
- talk by F. Piccinini

- $M_Z, \Gamma_Z$ : From  $\sigma(\sqrt{s})$  lineshape
  - Main uncertainties:  $B$ -field calibration, QED
  - $\delta M_Z, \delta \Gamma_Z \sim 0.1$  MeV could be achievable
- $m_t$ : Current status  $\delta m_t \sim 0.4$  GeV at LHC
  - Additional theory uncertainties?

PDG '18

Butenschoen et al. '16

Ferrario Ravasio, Nason, Oleari '18

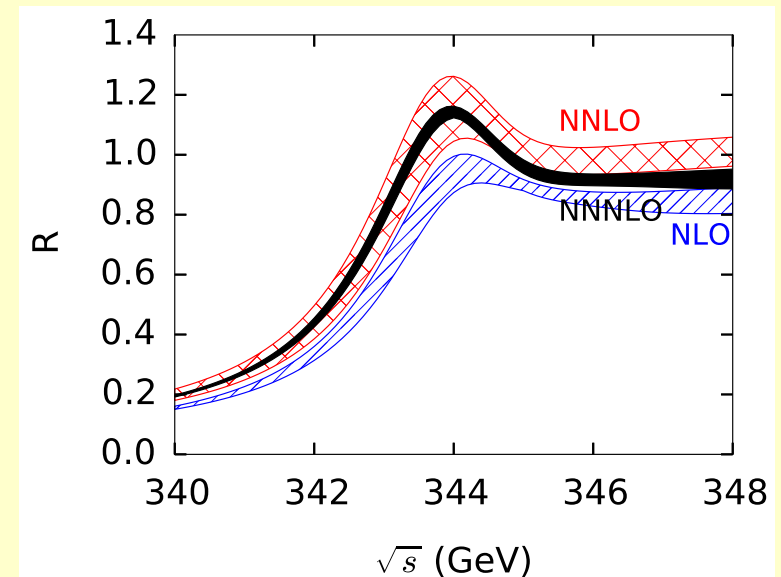
From  $e^+e^- \rightarrow t\bar{t}$  at  $\sqrt{s} \sim 350$  GeV

→ talk by Blondel

$$\delta m_t = [10 \dots 50 \text{ MeV}]_{\text{stat}}$$

- ⊕ [ $< 40$  MeV]<sub>exp. sys.</sub>
- ⊕ [ $\ll 50$  MeV]<sub>QCD</sub>
- ⊕ [10 MeV]<sub>mass def.</sub>
- ⊕ [10...20 MeV] <sub>$\alpha_s$</sub>

$\sim 50 \dots 100$  MeV



Beneke et al. '15

- $\alpha_s$ : d'Enterria, Skands, et al. '15
  - Most precise determination using Lattice QCD from  $\Upsilon$  spectroscopy:  
 $\alpha_s = 0.1184 \pm 0.0006$  HPQCD '10  
→ Difficulty in evaluating systematics
  - $e^+e^-$  event shapes and DIS:  $\alpha_s \sim 0.114$   
Alekhin, Blümlein, Moch '12; Abbate et al. '11; Gehrmann et al. '13  
→ Subject to sizeable non-perturbative power corrections  
→ Systematic uncertainties in power corrections?
  - Hadronic  $\tau$  decays:  $\alpha_s = 0.119 \pm 0.002$  PDG '18  
→ Non-perturbative uncertainties in OPE and from duality violation  
Pich '14; Boito et al. '15,18

- $\alpha_S$ :

- Electroweak precision ( $R_\ell = \Gamma_Z^{\text{had}} / \Gamma_Z^\ell$ ):

$$\alpha_S = 0.120 \pm 0.003$$

PDG '18

→ No (negligible) non-perturbative QCD effects

FCC:  $\delta R_\ell \sim 0.001$

⇒  $\delta \alpha_S < 0.0002$  (subj. to theory error)

**Caviat:**  $R_\ell$  could be affected by new physics

- $R = \frac{\sigma[ee \rightarrow \text{had.}]}{\sigma[ee \rightarrow \mu\mu]}$  at lower  $\sqrt{s}$

e.g. CLEO ( $\sqrt{s} \sim 9 \text{ GeV}$ ):  $\alpha_S = 0.110 \pm 0.015$

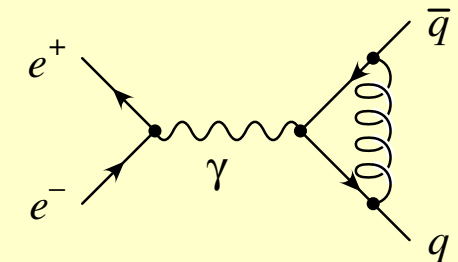
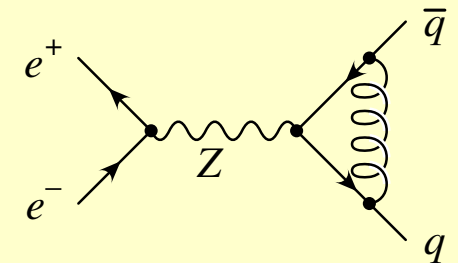
Kühn, Steinhauser, Teubner '07

→ dominated by  $s$ -channel photon, less room for new physics

→ QCD still perturbative

naive scaling to  $50 \text{ ab}^{-1}$  (BELLE-II):  $\delta \alpha_S \sim 0.0001$

d'Enterria, Skands, et al. '15





Effective field theory:  $\mathcal{L} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}(\Lambda^{-3}) \quad (\Lambda \gg M_Z)$

$$\mathcal{O}_{\phi 1} = (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi)$$

$$\alpha \Delta T = -\frac{v^2}{2} \frac{c_{\phi 1}}{\Lambda^2}$$

$$\mathcal{O}_{\text{BW}} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

$$\alpha \Delta S = -e^2 v^2 \frac{c_{\text{BW}}}{\Lambda^2}$$

$$\mathcal{O}_{\text{LL}}^{(3)e} = (\bar{L}_L^e \sigma^a \gamma_\mu L_L^e) (\bar{L}_L^e \sigma^a \gamma^\mu L_L^e)$$

$$\Delta G_F = -\sqrt{2} \frac{c_{\text{LL}}^{(3)e}}{\Lambda^2}$$

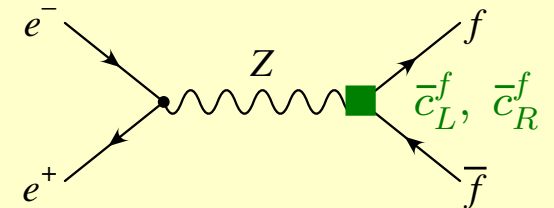
$$\mathcal{O}_R^f = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{f}_R \gamma^\mu f_R)$$

$$f = e, \mu, \tau, b, lq$$

$$\mathcal{O}_L^F = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{F}_L \gamma^\mu F_L)$$

$$F = \begin{pmatrix} \nu_e \\ e \end{pmatrix}, \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}, \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}, \begin{pmatrix} u, c \\ d, s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$$

$$\mathcal{O}_L^{(3)F} = i(\Phi^\dagger \overleftrightarrow{D}_\mu^a \Phi) (\bar{F}_L \sigma_a \gamma^\mu F_L)$$



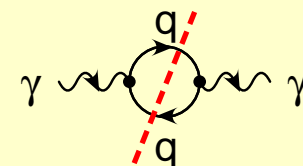
■  $\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 = 0.0315_{\text{lept}} + 0.0276_{\text{had}}$

■ Hadronic effects from  $e^+e^- \rightarrow \text{had. data}$

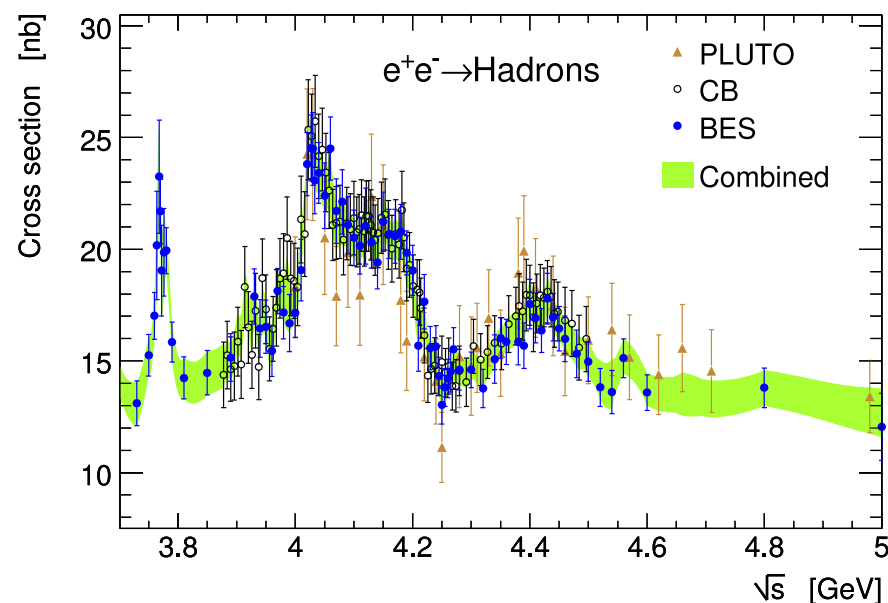
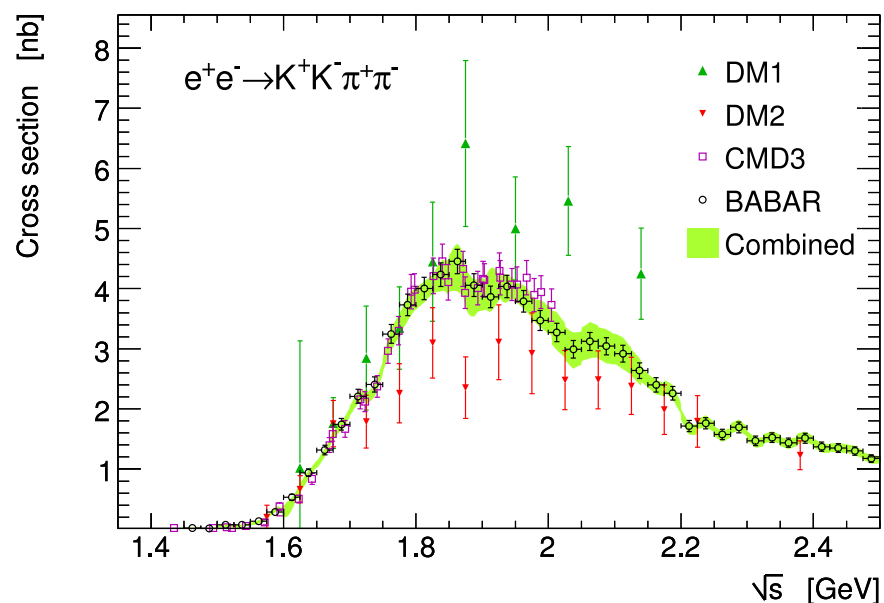
■ Last 5 years: new data from BaBar, VEPP, BES

■ No significant improvement from including  $\tau$  data

■ Robust precision  $\sim 10^{-4}$



Davier et al. '17; Jegerlehner '17  
Keshavarzi, Nomura, Teubner '18

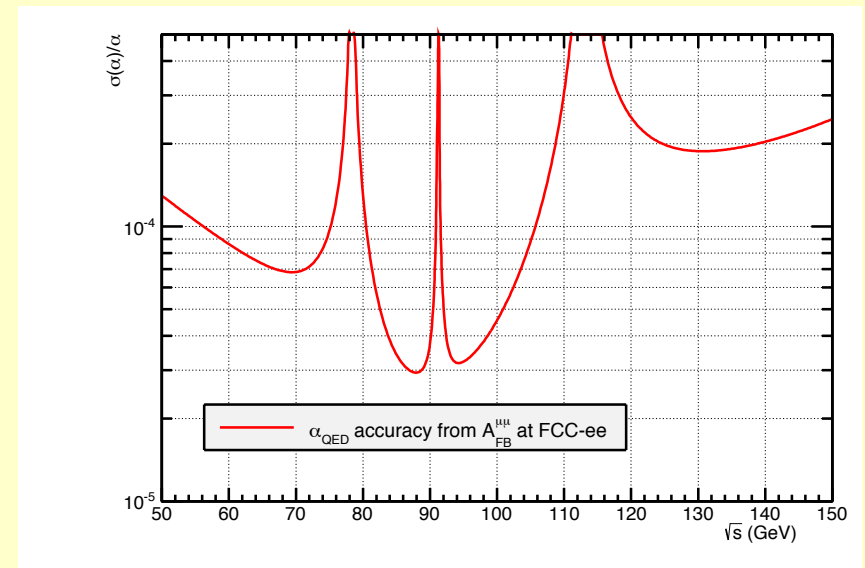


Davier et al. '17

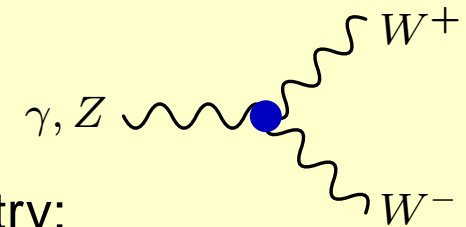
- $\Delta\alpha_{\text{had}}$ : Could be limiting factor
  - a) From  $e^+e^- \rightarrow \text{had.}$  using dispersion relation  
Current:  $\delta(\Delta\alpha_{\text{had}}) \sim 10^{-4}$   
Improvement to  $\delta(\Delta\alpha_{\text{had}}) \sim 5 \times 10^{-5}$  likely
  - b) Direct determination at FCC-ee from  $e^+e^- \rightarrow \mu^+\mu^-$  off the Z peak  
(i.e.  $A_{\text{FB}}^{\mu\mu}$  at  $\sqrt{s} \sim 88$  GeV and  $\sqrt{s} \sim 95$  GeV)  
 $\rightarrow \delta_{\text{th}}(\Delta\alpha_{\text{had}}) \sim 3 \times 10^{-5}$

Janot '15

Requires high-precision theory prediction for  $e^+e^- \rightarrow \mu^+\mu^-$  including 2/3-loop corrections for  $\gamma$ -exchange and box contributions



$\gamma$ WW and ZWW couplings:



General coupling structure consistent with Lorentz symmetry:

$$\begin{aligned} \mathcal{L} = & -ig_V \left[ g_1^V (W_{\mu\nu}^+ W^{-\mu} - W_{\mu\nu}^- W^{+\mu}) V^\nu + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{M_W^2} W_\mu^{+\nu} W_\nu^{-\rho} V_\rho^\mu \right. \\ & + ig_4^V W_\mu^+ W_\nu^- V^{\mu\nu} + ig_5^V \epsilon^{\mu\nu\rho\sigma} W_\mu^+ \overleftrightarrow{\partial}_\rho W_\nu^- V_\sigma \\ & \left. + \tilde{\kappa}_V W_\mu^+ W_\nu^- \tilde{V}^{\mu\nu} + \frac{\tilde{\lambda}_V}{M_W^2} W_\mu^{+\nu} W_\nu^{-\rho} \tilde{V}_\rho^\mu \right] \quad (V = \gamma, Z) \end{aligned}$$

Hagiwara, Peccei, Zeppenfeld, Hikasa '87

EM gauge invariance:  $g_1^\gamma = 1$

SU(2) invariance imposes more constraints

Extension of SM by **higher-dimensional operators**:

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{d=5}^{\infty} \frac{1}{\Lambda^{d-4}} \sum_i c_i \mathcal{O}_i^{(d)}$$

Operators must satisfy SM gauge invariance

**Leading CP-even gauge boson operators ( $d = 6$ ):**

$$\mathcal{O}_{\text{WWWW}} = \text{Tr}[W_{\mu\nu} W^{\nu\rho} W_{\rho}{}^{\mu}]$$

Hagiwara, Ishihara, Szalapski, Zeppenfeld '93

$$\mathcal{O}_{\text{W}} = (D_{\mu}\Phi)^{\dagger} W^{\mu\nu} (D_{\nu}\Phi)$$

$$\mathcal{O}_{\text{B}} = (D_{\mu}\Phi)^{\dagger} B^{\mu\nu} (D_{\nu}\Phi)$$

→ Only 3 indep. parameters

Relations between anom. couplings (aGC) and eff. theory (EFT):

$$g_1^Z = 1 + c_{\text{W}} \frac{M_{\text{Z}}^2}{2\Lambda^2}$$

$$\kappa_{\gamma} = 1 + (c_{\text{W}} + c_{\text{B}}) \frac{M_{\text{W}}^2}{2\Lambda^2}$$

$$\lambda_{\gamma} = \lambda_{\text{Z}} = c_{\text{WWWW}} \frac{3g^2 M_{\text{W}}^2}{2\Lambda^2}$$

$$\kappa_{\text{Z}} = 1 + \left( c_{\text{W}} - c_{\text{B}} \frac{s^2}{c^2} \right) \frac{M_{\text{W}}^2}{2\Lambda^2}$$

- CP-even dim-6 operators can be constrained from final-state distributions
- No significant improvement from beam polarization

$\mathcal{P}_{-80,-20} : \mathcal{P}_{-80,+20} : \mathcal{P}_{+80,-20} : \mathcal{P}_{+80,+20}$	$\Delta g_1^Z$	$\Delta \kappa_\gamma$	$\Delta \lambda_\gamma$	$\Delta \mathcal{P}_{e^-}$	$\Delta \mathcal{P}_{e^+}$
1:1:1:1	1.88	1.73	2.66	0.0009	0.0014
1:4:4:1	1.92	1.68	2.79	0.0010	0.0017
0:1:1:0	1.89	1.69	2.85	0.0015	0.0023
1:4:4:1 and $ \mathcal{P}_e  \neq - \mathcal{P}_e $	1.92	1.68	2.79	0.0016	0.0018

- Strongly coupled models / mixing of Higgs with heavy states:  
Other aGC possible
- Independent constraint of all parameters requires  
 $e^\pm$  **longitudinal and transverse** polarization

**Table 6.** Same as Table 5, but for symmetry (b) and with the L–R-parameterisation. We write again  $\tilde{h}_\pm = \text{Im}(g_1^R \pm \kappa_R)/\sqrt{2}$ . Using this parameterisation, a maximum number of couplings can be measured without transverse beam polarisation. In the  $\gamma$ – $Z$ -parameterisation, the four couplings  $\text{Im } g_1^\gamma$ ,  $\text{Im } g_1^Z$ ,  $\text{Im } \kappa_\gamma$  and  $\text{Im } \kappa_Z$  are not measurable without transverse polarisation

	$\text{Im } g_1^L$	$\text{Im } \kappa_L$	$\text{Im } \lambda_L$	$\text{Im } g_5^L$	$\tilde{h}_-$	$\tilde{h}_+$	$\text{Im } \lambda_R$	$\text{Im } g_5^R$
No polarisation	2.7	1.7	0.48	2.5	11	–	3.1	17
$(P_t^-, P_t^+) = (\mp 80\%, 0)$	2.6	1.2	0.45	2.0	4.5	–	1.4	4.3
$(P_t^-, P_t^+) = (\mp 80\%, \pm 60\%)$	2.1	0.95	0.37	1.6	2.5	–	0.75	2.3
$(P_t^-, P_t^+) = (80\%, 60\%)$	2.6	1.2	0.46	2.0	3.7	3.2	0.98	4.4

$$\sqrt{s} = 500 \text{ GeV}$$

- **Electroweak precision tests** at future  $e^+e^-$  colliders probe physics at beyond the Standard Model at **multi-TeV scale**
- Uncertainties from experimental **systematics**, **perturbative** and **non-perturbative** theory require much work, but can also be mitigated through choice of measurements and analysis
- **Direct determination** of  $\alpha_s$ ,  $m_t$ ,  $\Delta\alpha$  at  $e^+e^-$  colliders is important
- Other lower-energy experiments can provide additional input:  
**BELLE II**, **BES**, ...
- **Beam polarization** can access more information, but non essential for SM program



**Backup slides**

## Example: Error estimation for $\Gamma_Z$

### ■ Geometric perturbative series

$$\alpha_t = \alpha m_t^2$$

$$\mathcal{O}(\alpha^3) - \mathcal{O}(\alpha_t^3) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha^2) \sim 0.26 \text{ MeV}$$

$$\mathcal{O}(\alpha^2 \alpha_s) - \mathcal{O}(\alpha_t^2 \alpha_s) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.30 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.23 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^3) - \mathcal{O}(\alpha_t \alpha_s^3) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s^2) \sim 0.035 \text{ MeV}$$

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$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \mathcal{O}(\alpha_{\text{bos}})^2 \sim 0.1 \text{ MeV}$$

### ■ Parametric prefactors:

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \Gamma_Z \alpha^2 \sim 0.1 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\alpha n_{|q}}{\pi} \alpha_s^2 \sim 0.29 \text{ MeV}$$

**Total:**  $\delta \Gamma_Z \approx 0.5 \text{ MeV}$

## ■ Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

Kureav, Fadin '85

Berends, Burgers, v. Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v. Neerven '89

Skrzypek '92

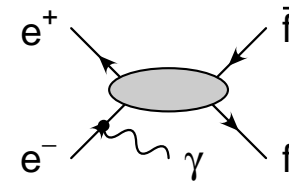
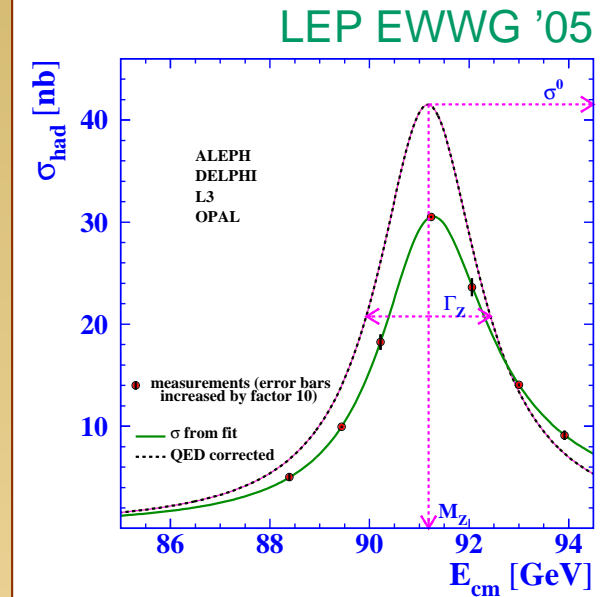
Montagna, Nicosini, Piccinini '97

Soft photons (resummed) + collinear photons

$$\mathcal{R}_{\text{ini}} = \frac{\zeta(1 - s'/s)\zeta^{-1}}{\Gamma(1 - \zeta)} e^{-\gamma_E \zeta + 3\alpha L/2\pi} - \frac{\alpha}{\pi} L \left(1 + \frac{s'}{s}\right) + \alpha^2 L^2 \dots + \alpha^3 L^3 \dots$$

$$\zeta = \frac{2\alpha}{\pi} (L - 1)$$

$$L = \log \frac{s}{m_e^2}$$

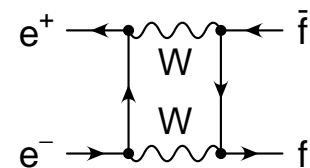
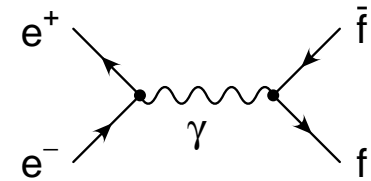
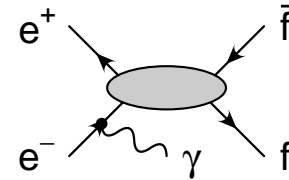
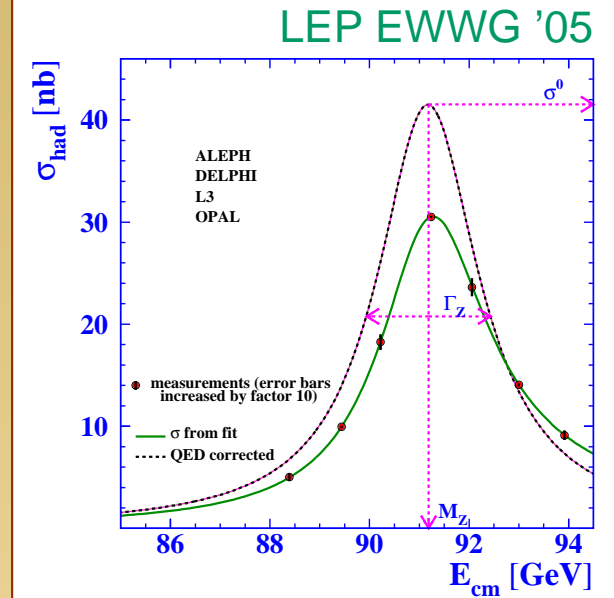


- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{ini}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of  $\gamma$ -exchange,  $\gamma$ -Z interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$



- Deconvolution of initial-state QED radiation:

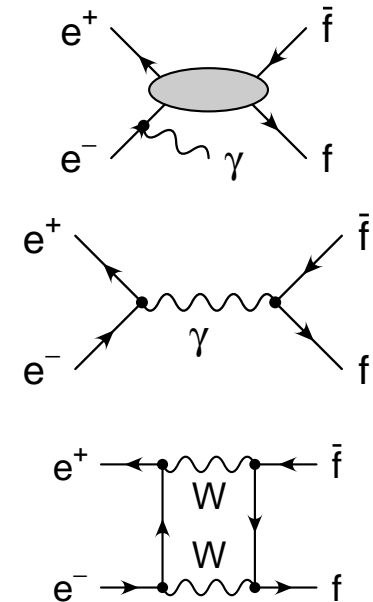
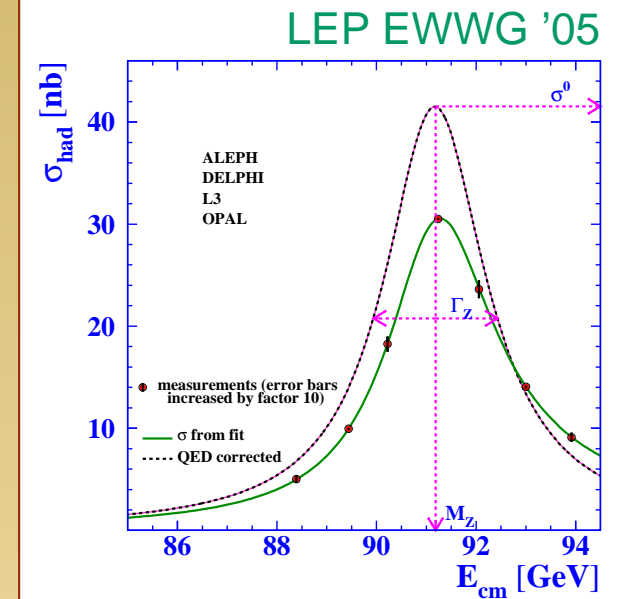
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- Z-pole contribution:

$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2 \overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$



- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

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- Z-pole contribution:

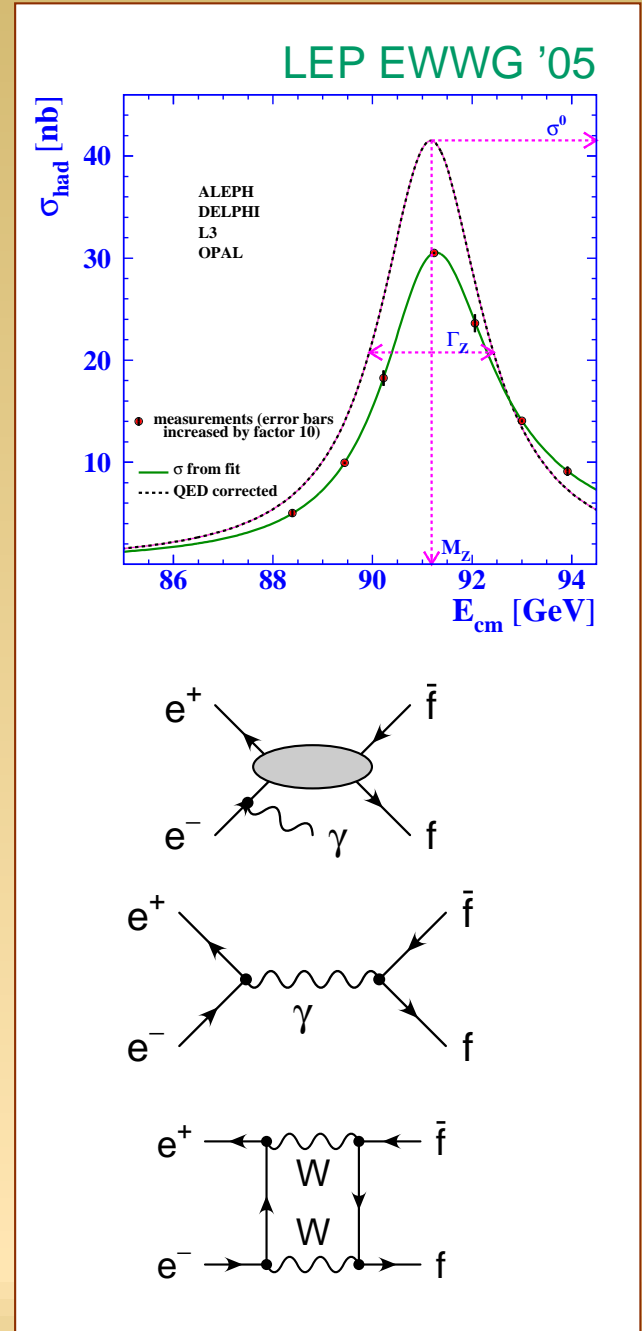
$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2 \overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$

- In experimental analyses:

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2 \Gamma_Z^2 / M_Z^2}$$

$$\overline{M}_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

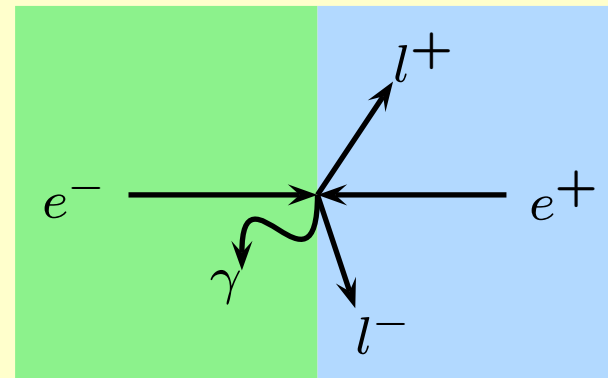
$$\overline{\Gamma}_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$



QED radiation in principle cancels in asymmetries, e.g.  $A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$

Some effects from detector acceptance and cuts

Typical influence  $< 10^{-3}$



Implementation of QED effects:

a) Analytical formulae, e.g. ZFITTER

Arbuzov, Bardin, Christova, Kalinovskaya, Riemann, Riemann, ...

b) Monte Carlo event generator, e.g. KORALZ

Jadach, Ward, ...

- $\alpha_s$ :

Most precise determination using Lattice QCD from  $\Upsilon$  spectroscopy:

$$\alpha_s = 0.1184 \pm 0.0006$$

HPQCD '10

But  $e^+e^-$  event shapes and DIS prefer  $\alpha_s \sim 0.114$

Alekhin, Blümlein, Moch '12; Abbate et al. '11; Gehrmann et al. '13

- Impact on EWPOs:

$$\delta\alpha_s = 0.005 \quad \Rightarrow \quad \delta M_W \approx 3.5 \text{ MeV}$$

$$\delta \sin^2 \theta_{\text{eff}}^l \approx 2 \times 10^{-5}$$

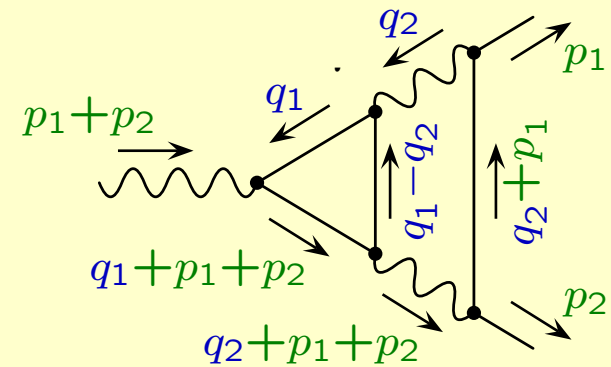
Currently not dominant, but similar magnitude as intrinsic theory error



Experimental precision requires inclusion of **radiative corrections** in theory (1-loop, 2-loop, and partial 3-loop)

Integrals over loop momenta:

$$\int d^4 q_1 d^4 q_2 f(q_1, q_2, p_1, p_2, \dots, m_1, m_2, \dots)$$



Computer algebra tools:

- Generation of diagrams,  $\mathcal{O}(100) - \mathcal{O}(10000)$
- Lorentz and Dirac algebra
- Integral simplification (and expansion)

Evaluation of loop integrals:

- In general not possible analytically
- Numerical methods must be automizable, stable, fastly converging
- Need procedure for isolating divergent pieces

## ■ Subtraction of QED radiation contributions

→ Known to  $\mathcal{O}(\alpha^2)$ ,  $\mathcal{O}(\alpha^3 L^3)$  for **ISR**,  
 $\mathcal{O}(\alpha^2)$  for **FSR** and  $\mathcal{O}(\alpha^2 L^2)$  for **A<sub>FB</sub>**

$$(L = \log \frac{s}{m_e^2})$$

Berends, Burgers, v.Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v.Neerven '89

Skrzypek '92; Montagna, Nicrosini, Piccinini '97

→  $\mathcal{O}(0.1\%)$  uncertainty on  $\sigma_Z$ ,  $A_{FB}$

→ Improvement needed for ILC/FCC-ee

## ■ Subtraction of non-resonant $\gamma$ -exchange, $\gamma$ -Z interf., box contributions, Bhabha scattering

see, e.g., Bardin, Grünewald, Passarino '99

→  $\mathcal{O}(0.01\%)$  uncertainty within SM

(improvements may be needed)

→ Sensitivity to some NP beyond EWPO

