

Thrust distribution in Higgs decays at the next-to-leading order and beyond

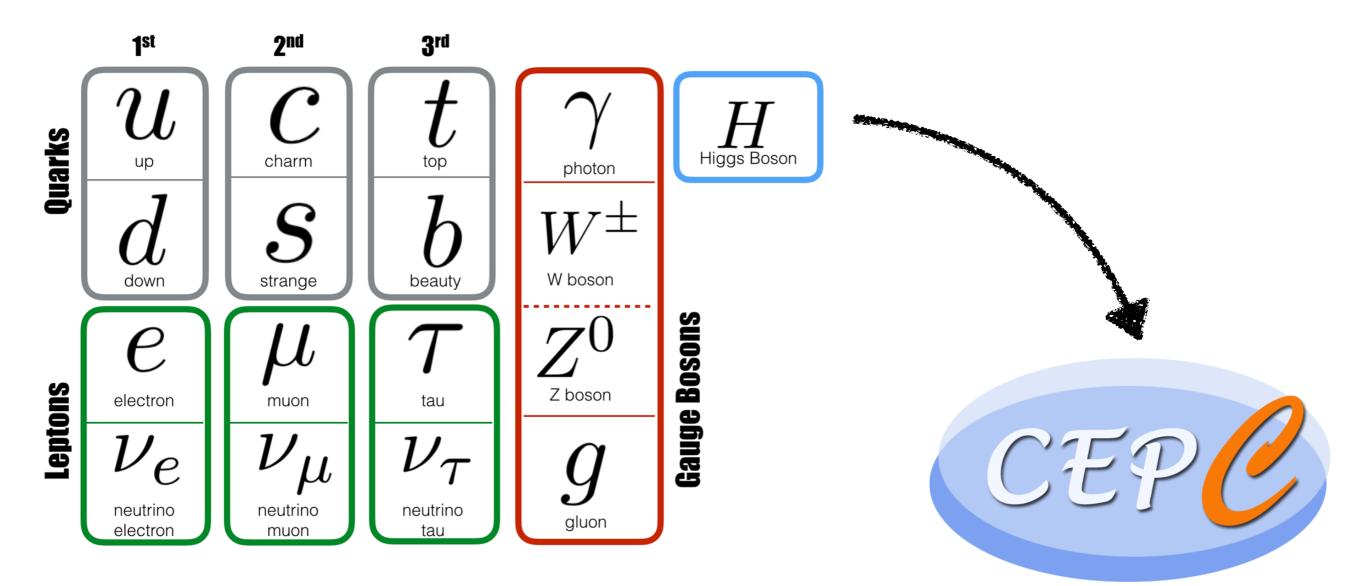
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2018 International Workshop on the High Energy Circular Electron Positron Collider

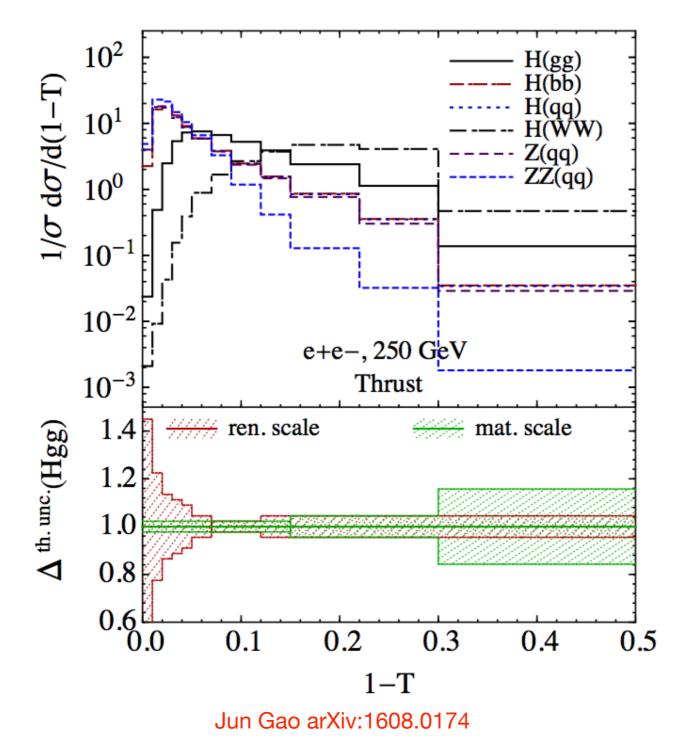
Outline

- Introduction
- Fix Order
- Factorization and NNLO-S
- Summary and Outlook

Motivation: Precision Measurements of Higgs Boson

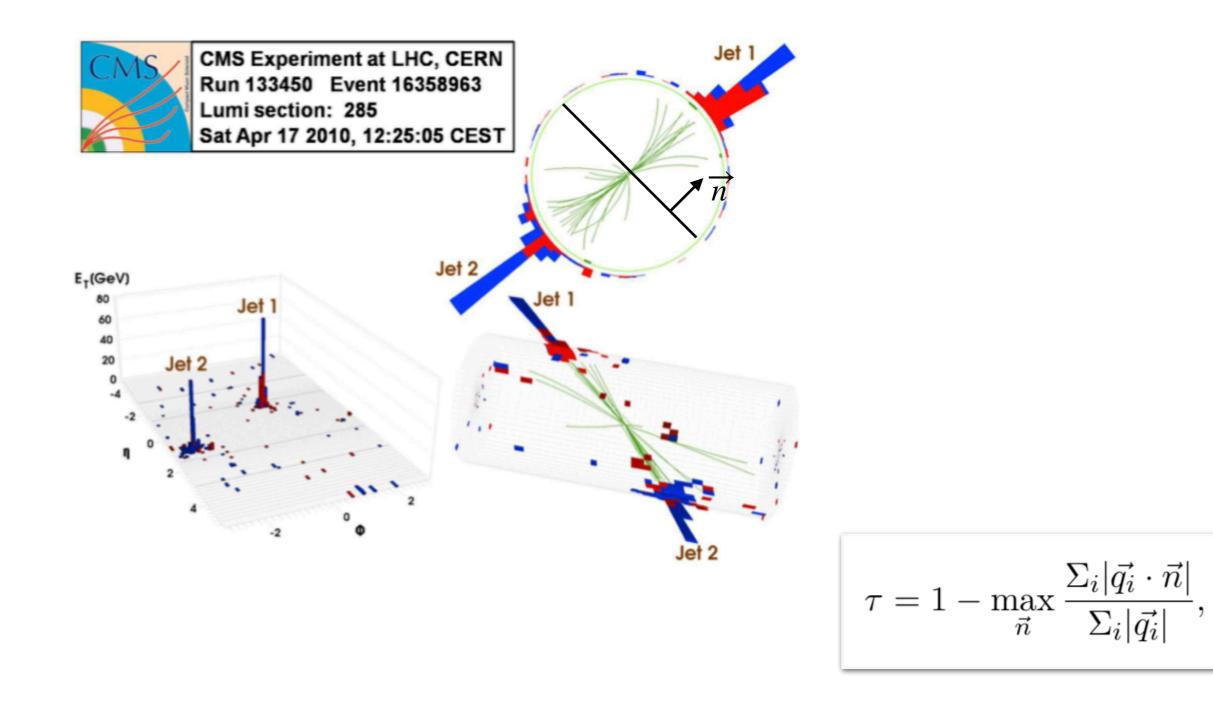


Motivation: Constrains to Light Quark Yukawa Coupling



- The quark signal and The gluon signal has different peak position.
- Helpful to select the signal and improve constrain to the light quark Yukawa coupling.
- Requires high precision theoretical prediction.

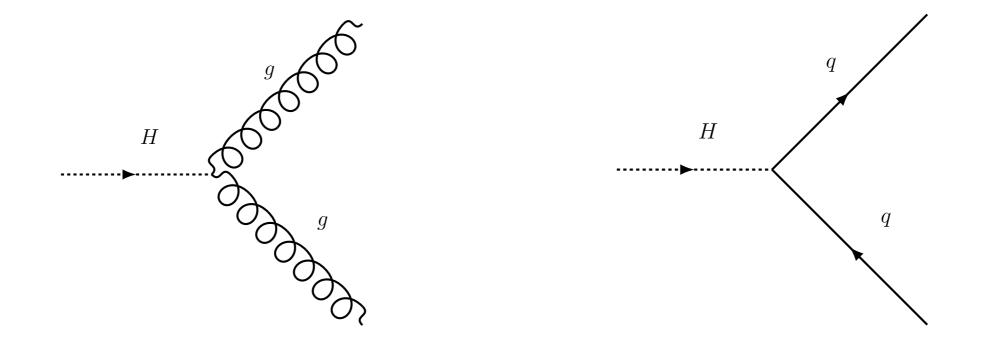
Introduction:Thrust



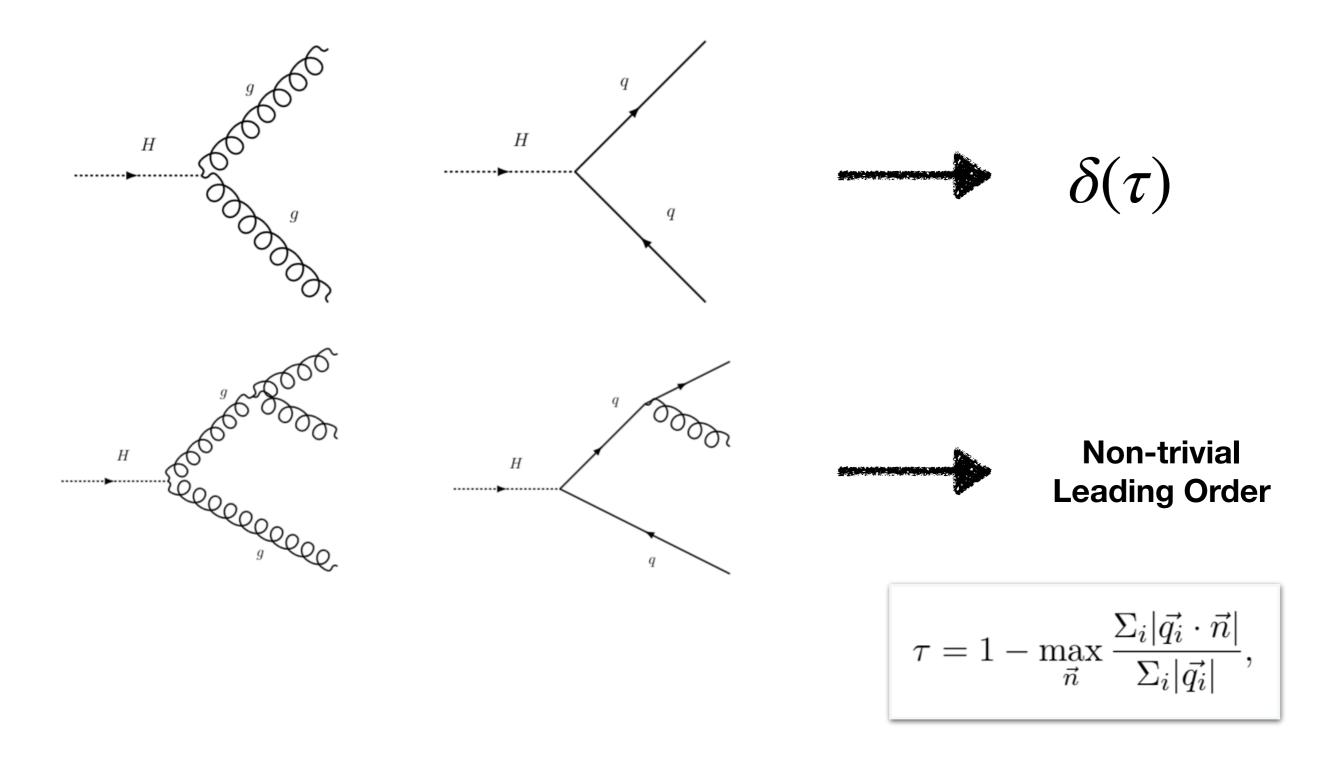
Fix Order

Fix Order: Effective Lagrangian

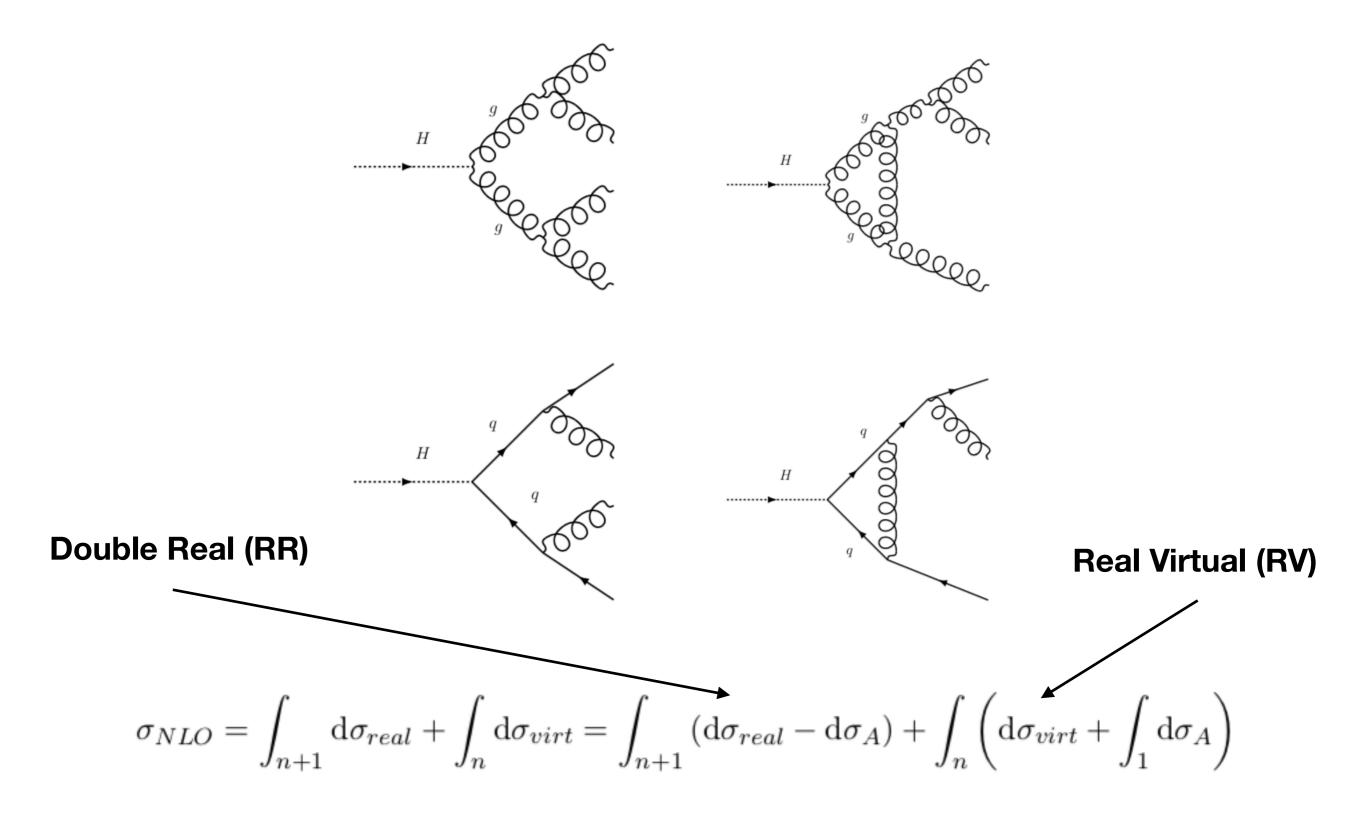
$$\mathcal{L}_{\text{eff}} = \frac{\alpha_s(\mu)C_t(m_t,\mu)}{12\pi v} HG^{\mu\nu,\alpha}G^{\alpha}_{\mu\nu} + \sum_q \frac{y_q(\mu)}{\sqrt{2}} H\bar{\psi}_q\psi_q,$$



Fix Order: Leading Order



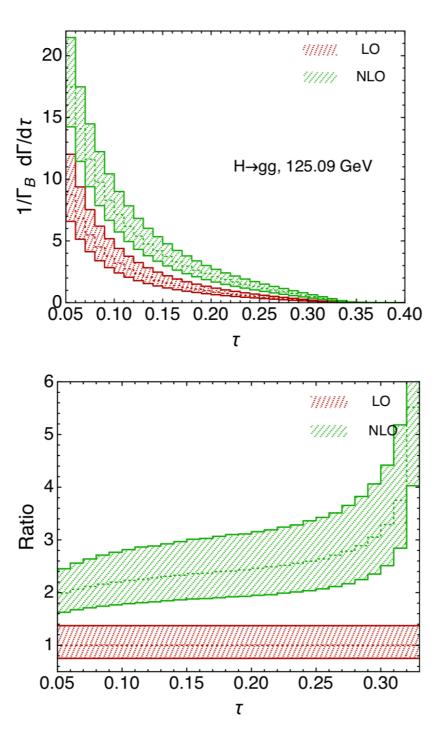
Fix Order: NLO and IR Subtraction



Fix Order: NLO and Numerical Implementations

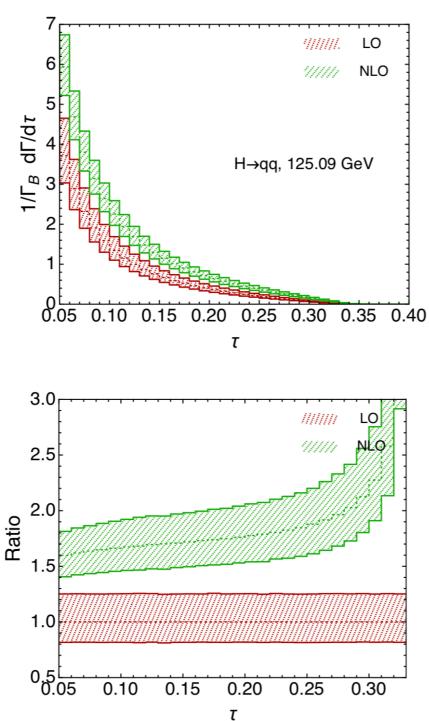
- Tree level matrix elements: HELAS
- One-loop matrix elements: OpenLoops
- IR subtraction: dipole formalism

Fix Order: Numeric Result Gluon channel



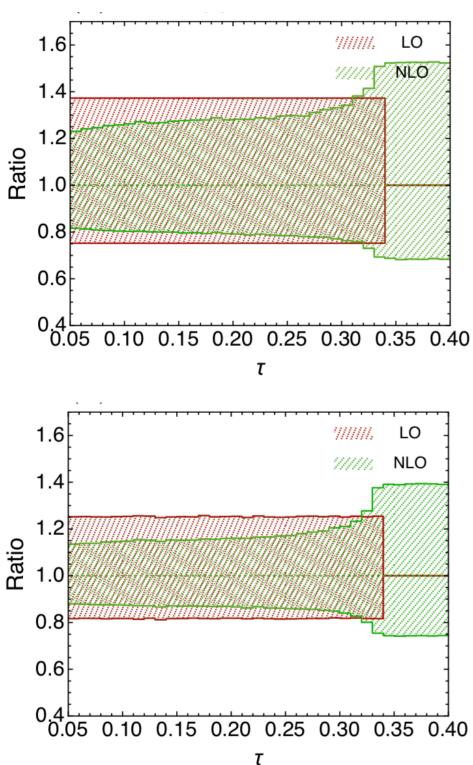
- the NLO correction is relatively large
- The scale dependence is not improved by the next to leading order correction.
- The scale dependence at leading order underestimate the higher order correction

Fix Order: Numeric Result Quark channel



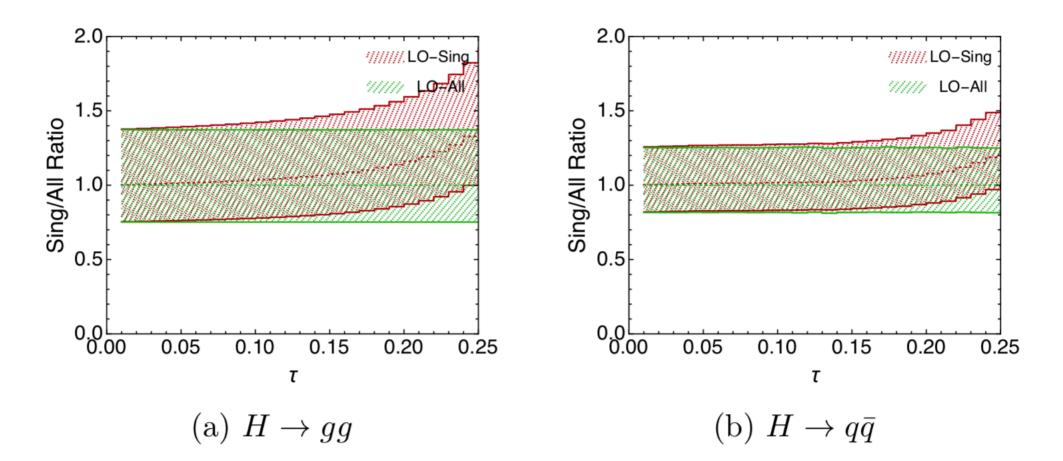
- the NLO correction is relatively large
- The scale dependence is not improved by the next to leading order correction.
- The scale dependence at leading order underestimate the higher order correction

Fix Order: Numeric Result Ratio to central value



- The leading order distribution vanished at tau greater than one third
- The next to leading order correction is the leading contribution.
- The scale dependence is thus larger than the small tau region

Fix Order: Singular Part



$$\frac{1}{\Gamma_B^{(q,g)}} \frac{\mathrm{d}\Gamma^{(q,g)}}{\mathrm{d}\tau} = \sum_{n=1}^{\infty} \left(\frac{\alpha_s}{4\pi}\right)^n \left(\sum_{i=0}^{n-1} A_{n,i} \frac{\ln^i \tau}{\tau} + R_n\right)$$

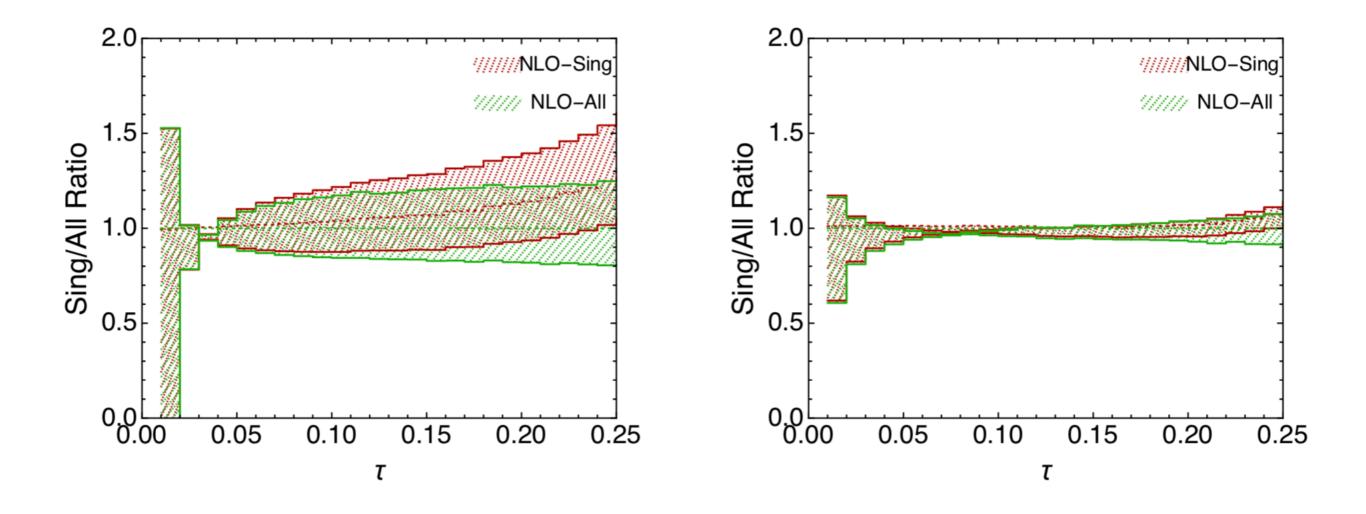
Factorization and NNLO-S

Factorization

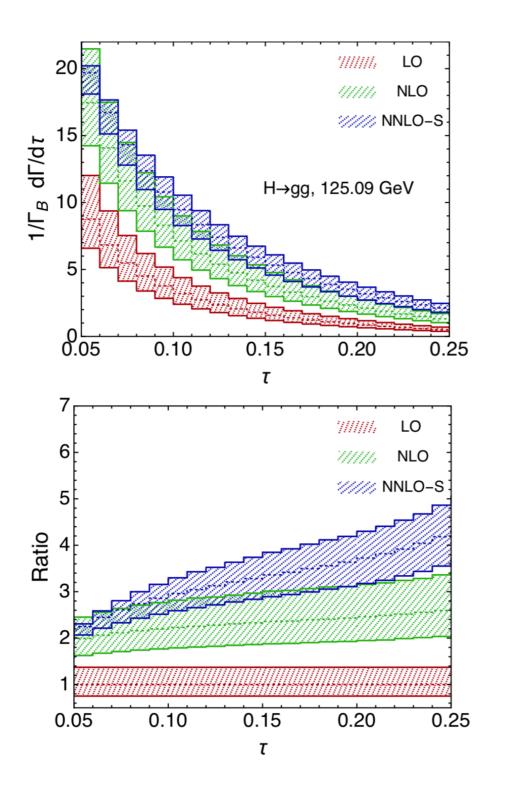
$$\mathcal{L}_{\text{SCET}} = \frac{\alpha_s C_t C_S^g}{12\pi v} H \mathcal{A}_{\perp \bar{n}}^{\mu\nu,\alpha} \mathcal{A}_{\perp n,\mu\nu,\alpha} + \frac{y_q}{\sqrt{2}} C_S^q H \bar{\chi}_{\bar{n}} \chi_n,$$

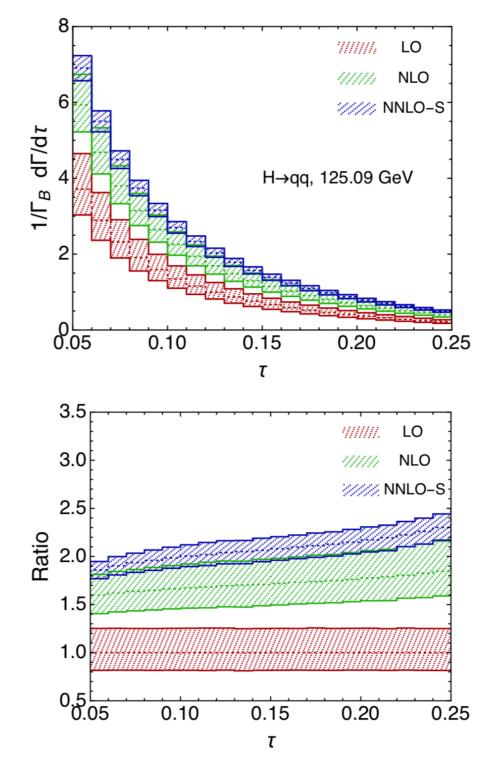
$$\begin{split} \frac{\mathrm{d}\Gamma_{q}}{\mathrm{d}\tau} =& \Gamma_{B}^{q} |C_{S}^{q}(M_{h},\mu)|^{2} \int \mathrm{d}M_{n}^{2} \mathrm{d}M_{\bar{n}}^{2} \mathrm{d}k \,\,\delta\left(\tau - \frac{M_{n}^{2} + M_{\bar{n}}^{2}}{M_{h}^{2}} - \frac{k}{M_{h}}\right) J_{n}^{q}(M_{n}^{2},\mu) J_{\bar{n}}^{q}(M_{\bar{n}}^{2},\mu) \\ & \times S^{q}(k,\mu), \\ \frac{\mathrm{d}\Gamma_{g}}{\mathrm{d}\tau} =& \Gamma_{B}^{g} |C_{t}(m_{t},\mu)|^{2} |C_{S}^{q}(M_{h},\mu)|^{2} \int \mathrm{d}M_{n}^{2} \mathrm{d}M_{\bar{n}}^{2} \mathrm{d}k \,\,\delta\left(\tau - \frac{M_{n}^{2} + M_{\bar{n}}^{2}}{M_{h}^{2}} - \frac{k}{M_{h}}\right) J_{n}^{g}(M_{n}^{2},\mu) \\ & \times J_{\bar{n}}^{g}(M_{\bar{n}}^{2},\mu) \,\,S^{g}(k,\mu), \end{split}$$

Singular Part at NLO

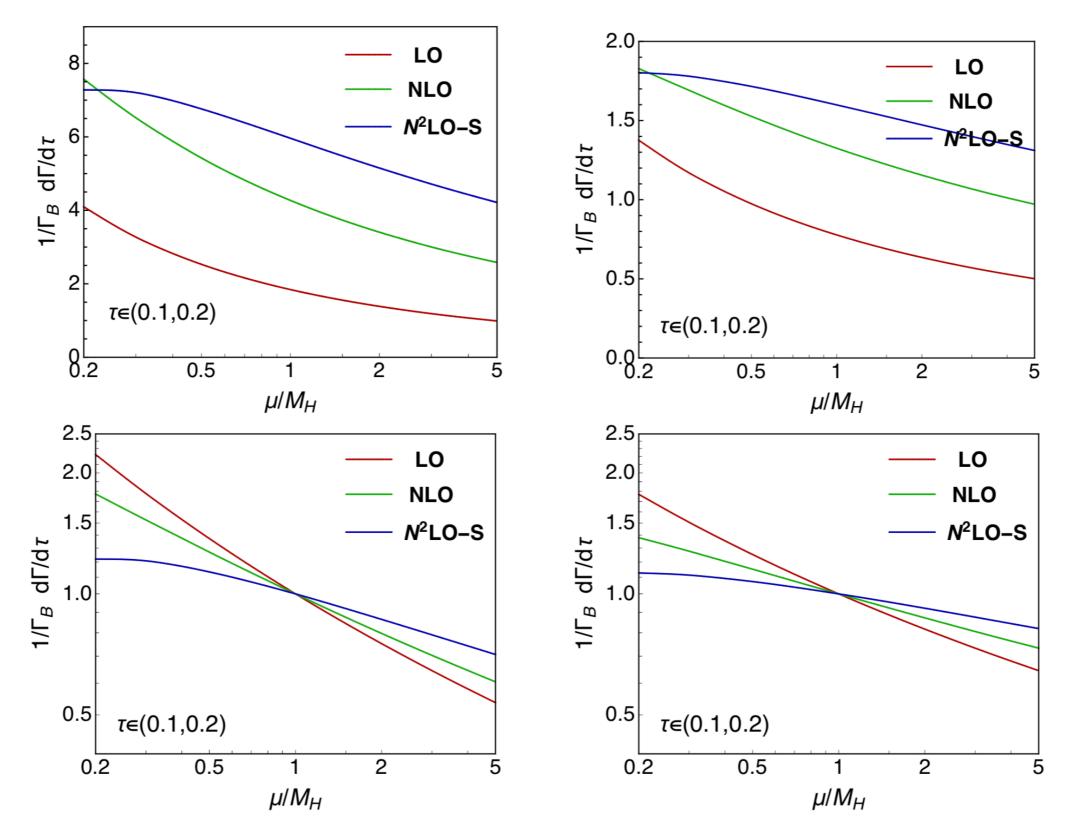


Numeric Result at NNLO-S





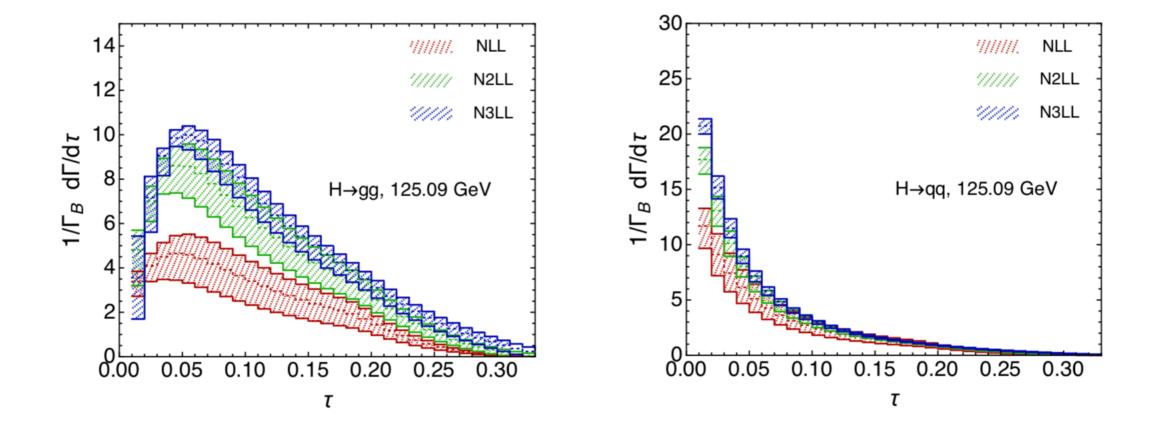
Scale dependence at NNLO-S



Summary and Outlook

- The NLO correction and the NNLO-S correction based on the factorization theorem.
- The perturbative expansion is badly converged and we may need higher order corrections.
- The scale dependence at LO and NLO respectively underestimate the higher order correction.
- To improve our result:
 - Small τ region: N3LL resummation
 - Large τ region: Regular terms of NNLO correction
 - Medium area: Higher order correction

Preliminary Result at N3LL+NLO



Thanks