Some backgrounds (motivation) to pickup the papers

-- A little exploration to get a hint for new idea to BESIII data analysis (not only these paper)

-- (my point of view) Graviton/Gravitino search has advantage of

-- One of hot topic now since the GW is detected recently

-- Since no new particle is found recently, DM papers assuming the candidate of Gravitino/Graviton is increasing (I feel but no sure !)

-- Graviton/Gravitino can couple with photons that makes the search "relatively" clean as mentioned in the paper and that is common advantage for colliders. Search for new phenomena in high-mass diphoton final states using 37 fb⁻¹ of protonproton collisions collected at sqrt(s) = 13 TeV with the ATLAS detector

ATLAS Collaboration, Physics Letters B 775 (2017) 105-125

Model

Heavy Higgs of spin-0 -- extended Higgs sector



• Randall and Sundrum (RS) model -- exchange of spin-2 particle



• ATLAS (CMS) reported a bump structure around 750 GeV from this diphoton search in 2016.

• ATLAS (CMS) reported a bump structure around 750 GeV





• How we can distinguish the spin ?

... seems to be distinguishable if we can detect with enough statistics

From "Spin-identification of Randall-Sundrum Graviton resonances ... " A. A. Pankov, A.V.Tsytrinov, and I. A. Serenkova



Event Selection

Data Set : total 36.7 fb⁻¹ of pp collision, sqrt(s)=13 TeV , 2015/2016

Event Reconstruction (mainly described for photons):

-- Two photons of highest energies which record more than 40 GeV / 30 GeV is retained.

- -- "Tight" photon
- -- Isolation from tracks

Event Selection :

-- For search of spin-0 signals, the transverse energy $E_T > 0.4m_{\gamma\gamma}$ for the leading photon, and $E_T > 0.3m_{\gamma\gamma}$ for the second photon, where $m_{\gamma\gamma}$ denotes diphoton invariant mass

-- For search of spin-2 resonant and non-resonant signals, $E_T > 55$ GeV (for each photon)

Diphoton invariant mass distribution

-- Estimation of background components is compared with several methods. --



Fig. 1. The diphoton invariant-mass distributions of the data are shown in the upper panels for (a) the spin-0 and (b) the spin-2 selections and their decomposition into contributions from genuine diphoton ($\gamma\gamma$), photon+jet (γj and $j\gamma$) and dijet (jj) events as determined using the 2×2D sideband method. The bottom panels show the purity of diphoton events as determined by the matrix method and the 2×2D sideband method. Each point in the distributions is plotted in the centre of the corresponding bin. The total uncertainties, including statistical and systematic components added in quadrature, are shown as error bars.

paper claimed that estimation of background component from 3 different methods is consistent each other.

Reference : Background channel to the di-photon events



From "Search for Randall-Sundrum Gravitons at the LHC", Evan Wulf

Diphoton invariant mass distribution



Fig. 2. Distributions of the diphoton invariant mass for events passing (a) the spin-0 selection or (b) the spin-2 selection, with the background-only fits superimposed. The data points are plotted at the centre of each bin. The error bars indicate statistical uncertainties. The differences between the data and the fits are shown in the bottom panels. The arrows in the lower panels indicate values outside the range by more than one standard deviation. There is no data event with $m_{\gamma\gamma} > 2700$ GeV.

Significance and C.L. upper limit



Short summary of this paper

-- The data are consistent with the SM background expectation and no excess is observed around 750 GeV .

-- In the combined dataset (2015+2016), the largest local deviation for spin-0 (spin-2) resonance search is 2.6σ (3.0σ), and global significance of this excess is null (0.8σ) for each.

-- RS1 model with κ/M_{PL} = 0.1 is excluded below m_G=4.1 TeV.

Search for supersymmetry in a final state containing two photons and missing transverse momentum in sqrt(s)=13 TeV pp collisions at the LHC using the ATLAS detector

ATLAS Collaboration, Euro Phys. J. C (2016) 517

Model

Model of Gauge-mediated supersymmetry breaking (GGM)

- Lightest supersymmetric particle (LSP) is the gravitino G with a mass of significantly less than 1 GeV.
- Final states consists of di-photon + missing transverse momentum (its magnitude is E_T^{miss})



Fig. 1 Typical production and decay-chain processes for the gluinopair production GGM model for which the NLSP is a bino-like neutralino

ATLAS SUSY Searches* - 95% CL Lower Limits Status: July 2015

| Sta | atus: July 2015 | | | | | | $\sqrt{s} = 7, 8 \text{ TeV}$ |
|---|--|--|---|---|---|--|---|
| | Model | e, μ, τ, γ | Jets | E ^{miss} _T | ∫ <i>L dt</i> [fb | $\sqrt{s} = 7 \text{ TeV} \qquad \sqrt{s} = 8 \text{ TeV}$ | Reference |
| Inclusive Searches | $ \begin{array}{l} \text{MSUGRA/CMSSM} \\ \tilde{q} \tilde{q}, \tilde{q} \rightarrow q \tilde{k}_{1}^{0} \\ \tilde{q} \tilde{q}, \tilde{q} \rightarrow q \tilde{k}_{1}^{0} \text{ (compressed)} \\ \tilde{q} \tilde{q}, \tilde{q} \rightarrow q \ell \ell \ell \ell r / \nu \nu \tilde{k}_{1}^{0} \\ \tilde{q} \tilde{s}, \tilde{s} \rightarrow q \tilde{q} \tilde{k}_{1}^{1} \\ \tilde{g} \tilde{s}, \tilde{s} \rightarrow q \tilde{q} \tilde{k}_{1}^{1} \rightarrow q W^{\pm} \tilde{k}_{1}^{0} \\ \tilde{g} \tilde{s}, \tilde{s} \rightarrow q \tilde{q} \ell \ell \ell r / \nu \nu \tilde{k}_{1}^{0} \\ \tilde{g} \tilde{s}, \tilde{s} \rightarrow q q \ell \ell \ell r / \nu \nu \tilde{k}_{1}^{0} \\ \tilde{g} \tilde{s} \tilde{s} \tilde{s} q \ell \ell \ell r LSP \\ \tilde{g} \tilde{s} \tilde{s} \tilde{s} \tilde{s} q \ell \ell \ell r LSP \\ \tilde{g} \tilde{s} \tilde{s} \tilde{s} \tilde{s} \tilde{s} \tilde{s} \tilde{s} s$ | $\begin{array}{c} 0\text{-3 } e, \mu/1\text{-2 }\tau \\ 0 \\ \text{mono-jet} \\ 2 \ e, \mu \ (\text{off-}Z) \\ 0 \\ 0 \text{-1 } e, \mu \\ 2 \ e, \mu \\ 1\text{-2 }\tau + 0\text{-1 }\ell \\ 2 \ \gamma \\ \gamma \\ 2 \ e, \mu \ (Z) \\ 0 \end{array}$ | 2-10 jets/3 2-6 jets 1-3 jets 2-6 jets 2-6 jets 2-6 jets 0-3 jets 0-2 jets 2 jets 2 jets 2 jets 2 jets 2 jets | b Yes Yes | 20.3 20.3 20.3 20.3 20 20 20 20.3 20.3 2 | $\begin{array}{c c c c c c c c c c c c c c c c c c c $ | 1507.05525 1405.7875 1507.05525 1503.03290 1405.7875 1507.05525 1501.03555 1407.0603 1507.05493 1507.05493 1507.05493 1507.05493 1500.03290 1502.01518 |
| 3 rd gen. ẽ med. | $\begin{array}{l} gg, g \rightarrow b \bar{b} \tilde{\chi}_1^0 \\ gg, g \rightarrow t \tilde{\chi}_1^0 \\ gg, g \rightarrow t \tilde{\chi}_1^0 \\ gg, g \rightarrow t \tilde{\chi}_1^1 \\ gg, g \rightarrow b \bar{p} \tilde{\chi}_1^+ \end{array}$ | 0 0 0-1 <i>e</i> , µ 0-1 <i>e</i> , µ | 3 b 7-10 jets 3 b 3 b | Yes Yes Yes Yes | 20.1 20.3 20.1 20.1 | Image: Second | 1407.0600 1308.1841 1407.0600 1407.0600 |
| 3 rd gen. squarks direct production | $\begin{array}{c} b_{1}b_{1}, b_{1} \rightarrow b\tilde{\chi}_{1}^{0} \\ b_{1}b_{1}, b_{1} \rightarrow \delta\tilde{\chi}_{1}^{1} \\ \tilde{\iota}_{1}\tilde{\iota}_{1}, \tilde{\iota}_{1} \rightarrow b\tilde{\chi}_{1}^{1} \\ \tilde{\iota}_{1}\tilde{\iota}_{1}, \tilde{\iota}_{1} \rightarrow Wb\tilde{\chi}_{1}^{0} \text{ or } \tilde{\chi}_{1}^{0} \\ \tilde{\iota}_{1}\tilde{\iota}_{1}, \tilde{\iota}_{1} \rightarrow \tilde{\chi}_{1}^{0} \\ \tilde{\iota}_{1}\tilde{\iota}_{1}, \tilde{\iota}_{1} \rightarrow \tilde{\chi}_{1}^{0} \\ \tilde{\iota}_{2}\tilde{\iota}_{2}, \tilde{\iota}_{2} \rightarrow \tilde{\iota}_{1} + Z \end{array}$ | 0 2 e, µ (SS) 1-2 e, µ 0-2 e, µ (0 2 e, µ (Z) 3 e, µ (Z) | 2 b 0-3 b 1-2 b 0-2 jets/1-2 nono-jet/c-ta 1 b 1 b | Yes Yes Yes b Yes ag Yes Yes Yes | 20.1 20.3 1.7/20.3 20.3 20.3 20.3 20.3 | $ \begin{array}{c c c c c c c c c c c c c c c c c c c $ | 1308.2631 1404.2500 1209.2102, 1407.0583 1506.08616 1407.0608 1403.5222 1403.5222 |
| EW direct | $ \begin{array}{l} \tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\ell} \nu(\ell \tilde{\nu}) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\nu} \nu(\tau \tilde{\nu}) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{2}^{0} \rightarrow \tilde{\chi}_{1}^{-} \nu \tilde{\ell}_{L} \ell(\tilde{\nu}\nu), \ell \tilde{\nu} \tilde{\ell}_{L} \ell(\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} Z \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{2}^{+} \tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} \Lambda \tilde{\chi}_{1}^{0}, h \rightarrow b \tilde{b} / W W / \tau \\ \tilde{\chi}_{2}^{+} \tilde{\chi}_{3}^{-} \tilde{\chi}_{2}^{-} \tilde{\chi}_{3} \rightarrow \tilde{\ell}_{R} \ell \\ GGM (wino NLSP) weak prod. \end{array} $ | 2 e,μ 2 e,μ 2 τ 3 e,μ 2-3 e,μ τ/γγ e,μ,γ 4 e,μ 1 e,μ + γ | 0 0 | Yes Yes Yes Yes Yes Yes Yes | 20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3 | $\begin{tabular}{ c c c c c c } \hline $$ $$ $$ $$ $$ $$ $$ $$ $$ $$ $$ $$ $$ | 1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 1501.07110 1405.5086 1507.05493 |
| Long-lived particles | $\begin{array}{l} \text{Direct} \tilde{\chi}_1^+ \tilde{\chi}_1^- \text{ prod., long-lived} \tilde{\chi} \\ \text{Direct} \tilde{\chi}_1^+ \tilde{\chi}_1^- \text{ prod., long-lived} \tilde{\chi} \\ \text{Stable, stopped } \tilde{g} \text{ R-hadron} \\ \text{Stable } \tilde{g} \text{ R-hadron} \\ \text{GMSB, stable } \tilde{\tau}, \tilde{\chi}_1^0 {\rightarrow} \tilde{\tau}(\tilde{e}, \tilde{\mu}) {+} \tau \\ \text{GMSB, } \tilde{\chi}_1^0 {\rightarrow} \gamma \tilde{G}, \text{ long-lived} \tilde{\chi}_1^0 \\ \tilde{g}_{\tilde{g}}, \tilde{\chi}_1^0 {\rightarrow} \gamma \tilde{G}, \text{ long-lived} \tilde{\chi}_1^0 \\ \text{GGM } \tilde{g}_{\tilde{g}}, \tilde{\chi}_1^0 {\rightarrow} Z \tilde{G} \end{array}$ | $ \begin{array}{c} \overset{\pm}{\underset{1}{\overset{1}{\overset{1}{\overset{1}{\overset{1}{\overset{1}{\overset{1}{1$ | 1 jet - 1-5 jets - - - τ ts - | Yes Yes - - Yes - - | 20.3 18.4 27.9 19.1 19.1 20.3 20.3 20.3 | $\begin{array}{c c c c c c c c c c c c c c c c c c c $ | 1310.3675 1506.05332 1310.6584 1411.6795 1410.5542 1504.05162 1504.05162 |
| RPV | LFV $pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e\mu/e\tau/\mu\tau$ Bilmear RPV CMSSM $\tilde{x}_{1}^{\dagger}\tilde{x}_{1}^{-}, \tilde{x}_{1}^{\dagger} \rightarrow w\tilde{x}_{0}^{0}, \tilde{x}_{1}^{0} \rightarrow ee\tilde{v}_{\mu}, e\mu\tilde{v}$ $\tilde{x}_{1}^{\dagger}\tilde{x}_{1}^{-}, \tilde{x}_{1}^{\dagger} \rightarrow w\tilde{x}_{0}^{0}, \tilde{x}_{1}^{0} \rightarrow \tau r\tilde{v}_{e}, er\tilde{v}$ $\tilde{g}\tilde{s}, \tilde{s} \rightarrow qq$ $\tilde{g}\tilde{s}, \tilde{s} \rightarrow q\tilde{x}_{1}^{0}, \tilde{x}_{1}^{0} \rightarrow qqq$ $\tilde{g}\tilde{s}, \tilde{s} \rightarrow \tilde{t}_{1}\tilde{t}, \tilde{t}, \tilde{t} \rightarrow bs$ $\tilde{t}_{1}\tilde{t}_{1}, \tilde{t}, \tau \rightarrow bs$ $\tilde{t}_{1}\tilde{t}_{1}, \tilde{t}, \tau \rightarrow bf$ | $\begin{array}{c} & e \mu, e \tau, \mu \tau \\ 2 e, \mu (SS) \\ \tau & 4 e, \mu \\ \tau & 3 e, \mu + \tau \\ 0 \\ 2 e, \mu (SS) \\ 0 \\ 2 e, \mu (SS) \\ 0 \\ 2 e, \mu \end{array}$ | - 0-3 b - 6-7 jets 0-3 b 2 jets + 2 b 2 b | - Yes Yes - - Yes - - | 20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3 | №r 1.7 TeV X ₁₁₁ =0.11, A _{132/133/233} =0.07 ∅, ĝ 1.35 TeV m(𝔅)=m(𝔅), cr _{LSP} <1 mm | 1503.04430 1404.2500 1405.5086 1502.05686 1502.05686 1502.05686 1404.250 ATLAS-CONF-2015-026 ATLAS-CONF-2015-015 |
| Other | Scalar charm, $\tilde{c} \rightarrow c \tilde{\ell}_1^0$ | 0 | 2 c | Yes | 20.3 | ک 490 GeV m(k ⁰ ₁)<200 GeV | 1501.01325 |
| 10 ⁻¹ Mass scale [TeV] | | | | | | | |

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

ATLAS_Preliminary

Event Selection

Data Set : 3.2 fb⁻¹ of pp collision, sqrt(s)=13 TeV, 2015

Event Reconstruction : electron/muon/jet/photon...

Event Selection : two tight photons with $p_T > 75$ GeV

 $H_{\rm T}$: scalar sum of the $P_{\rm T}$ of photons and leptons/jets

meff : $H_T + E_T^{miss} ==$ Total E involved in the GGM process

| SR | $W\gamma\gamma$ CR |
|--|--|
| 2 Tight photons with $p_{\rm T} > 75 {\rm GeV}$ | 2 Tight photons with $p_{\rm T} > 50 {\rm ~GeV}$ |
| | 1 e or μ with $p_{\rm T} > 25 { m GeV}$ |
| $\Delta \phi_{\min}(\text{jet}, p_{\text{T}}^{\text{miss}}) > 0.5$ | $\Delta \phi_{\min}(\text{jet}, p_{\text{T}}^{\text{miss}}) > 0.5$ |
| $E_{\rm T}^{\rm miss} > 175~{\rm GeV}$ | $50 < E_{\mathrm{T}}^{\mathrm{miss}} < 175 \mathrm{~GeV}$ |
| $m_{\rm eff} > 1500 { m ~GeV}$ | N(jets) < 3 |
| | $m_{e_{\mathcal{V}}} \notin 83-97 \text{ GeV}$ |

Reference : "loose" photon and "tight" photon

| Category | Description | Name | loose | tight |
|------------------|---|------------------------|--------------|--------------|
| Acceptance | $ \eta < 2.37$, with $1.37 < \eta < 1.52$ excluded | _ | \checkmark | \checkmark |
| Hadronic leakage | Ratio of $E_{\rm T}$ in the first sampling layer of the hadronic calorimeter to $E_{\rm T}$ of the EM cluster (used over the range $ \eta < 0.8$ or $ \eta > 1.37$) | | ~ | \checkmark |
| | Ratio of $E_{\rm T}$ in the hadronic calorimeter to $E_{\rm T}$ of the EM cluster (used over the range $0.8 < \eta < 1.37$) | R _{had} | ~ | \checkmark |
| EM Middle layer | Ratio of $3 \times 7 \ \eta \times \phi$ to 7×7 cell energies | R_{η} | \checkmark | \checkmark |
| | Lateral width of the shower | w_{η_2} | \checkmark | \checkmark |
| | Ratio of $3 \times 3 \ \eta \times \phi$ to 3×7 cell energies | R_{ϕ} | | \checkmark |
| EM Strip layer | Shower width calculated from three strips around the strip with maximum energy deposit | <i>w</i> _{s3} | | \checkmark |
| | Total lateral shower width | $w_{s tot}$ | | \checkmark |
| | Energy outside the core of the three central strips but within seven strips divided by energy within the three central strips | $F_{\rm side}$ | | \checkmark |
| | Difference between the energy associated with the second maximum in the strip layer and the energy re- constructed in the strip with the minimum value found between the first and second maxima | ΔE | | V |
| | Ratio of the energy difference associated with the largest and second largest energy deposits to the sum of these energies | E _{ratio} | | √ |

Table 1: Discriminating variables used for loose and tight photon identification.

"measurement of the photon identification efficiencies with the ATLAS detector using LHC Run01 data" (2016) arXiv: 1606.01813

E_T^{miss} & m_{eff} distribution of the di-photon sample



Table 1 Requirements defining the signal region (SR) and the $W\gamma\gamma$ CR referred to in Sect. 6

| SR | $W_{\gamma\gamma}$ CR |
|--|--|
| 2 Tight photons with $p_{\rm T} > 75 \text{ GeV}$ | 2 Tight photons with $p_{\rm T} > 50 {\rm ~GeV}$ |
| | 1 <i>e</i> or μ with $p_{\rm T} > 25 { m GeV}$ |
| $\Delta \phi_{\min}(\text{jet}, p_{\text{T}}^{\text{miss}}) > 0.5$ | $\Delta \phi_{\min}(\text{jet}, p_{\text{T}}^{\text{miss}}) > 0.5$ |
| $E_{\rm T}^{\rm miss} > 175~{\rm GeV}$ | $50 < E_{\rm T}^{\rm miss} < 175 {\rm ~GeV}$ |
| $m_{\rm eff} > 1500~{\rm GeV}$ | N(jets) < 3 |
| | $m_{e\gamma} \notin 83-97 \text{ GeV}$ |

Table 2 Summary of background estimates by source, and total combined background, in the signal region. The uncertainties shown

| Source | Number of events |
|--|------------------------|
| QCD $(\gamma\gamma, \gamma j, jj)$ | $0.05^{+0.20}_{-0.05}$ |
| $e \rightarrow \gamma$ fakes | 0.03 ± 0.02 |
| W_{YY} | 0.17 ± 0.08 |
| $Z\gamma\gamma$ | 0.02 ± 0.02 |
| Sum | $0.27_{-0.10}^{+0.22}$ |
| $(m_{\tilde{g}}, m_{\tilde{\chi}_1^0}) = (1500, 100)$ | 7.0 |
| $(m_{\tilde{g}}, m_{\tilde{\chi}_1^0}) = (1500, 1300)$ | 8.0 |

Results and C.L. upper limit

After the selection, number of events in the signal region is 0 where as that of expected SM background event is 0.27 .

By considering uncertainties,
 95% C.L. on the non-SM
 is 3 events.

• 95% C.L. on the cross section 0.93 fb.

with many (widely accepted) assumption on SUSY model , , ,

 Table 4
 Numbers of events observed in the SR after the successive application of the selection requirements, as well as the size of the expected SM background

| Requirement | Number of events |
|--|------------------------|
| Two photons, $p_{\rm T}^{\gamma} > 75$ | 4982 |
| $\Delta \phi_{\min}(\text{jet}, p_{\text{T}}^{\text{miss}}) > 0.5$ | 4724 |
| $m_{\rm eff} > 1500 { m ~GeV}$ | 1 |
| $E_{\rm T}^{\rm miss} > 175~{\rm GeV}$ | 0 |
| Expected SM background | $0.27^{+0.22}_{-0.10}$ |



Short summary of this paper

-- No events in the signal region is observed where the SM background is extimated to be 0.27 events.

-- 95% C.L. upper limit of 3.0 events (0.93 fb) on the number of events (cross section) due to the BSM is obtained.

-- In the context of GGM model (with a bino-like NLSP) it leads to a lower limit of 1650 GeV in the gluino mass.

Reference -- ATLAS-CONF-2017-080

-- Update analysis by using total 36.1 fb⁻¹ pp collision at sqrt(s)=13 TeV

Table 8: Summary of the observed number of events (N_{obs}), and the number of events expected from SM sources (N_{exp}^{SM}), for each of the seven SRs. Also shown are the derived (S_{obs}^{95}) and expected (S_{exp}^{95}) model-independent 95% CL limits on the number of events from non-SM processes, and the observed ($\langle \epsilon \sigma \rangle_{obs}^{95}$) and expected ($\langle \epsilon \sigma \rangle_{exp}^{95}$) 95% CL limits on the visible cross section from non-SM processes. The last column of the table shows the significance Z of the observed excess (if any), and the probability p, capped at 0.5, that a background-only experiment is more signal-like than the observed number of events in the given signal region.

| Signal Region | Nobs | $N_{\mathrm{exp}}^{\mathrm{SM}}$ | $S_{ m obs}^{95}$ | $S_{ m exp}^{95}$ | $\langle \epsilon \sigma \rangle_{\rm obs}^{95}$ [fb] | $\langle \epsilon \sigma \rangle^{95}_{\exp}$ [fb] | $Z\left(p ight)$ |
|---------------------------|------|----------------------------------|-------------------|---------------------|---|--|------------------|
| $SR_{S-L}^{\gamma\gamma}$ | 0 | $0.50^{+0.30}_{-0.26}$ | 3.0 | $3.1^{+1.4}_{-0.2}$ | 0.083 | $0.086^{+0.039}_{-0.003}$ | 0.00 (0.50) |
| $SR_{S-H}^{\gamma\gamma}$ | 0 | $0.48^{+0.30}_{-0.25}$ | 3.0 | $3.1^{+1.3}_{-0.1}$ | 0.083 | $0.086^{+0.036}_{-0.003}$ | 0.00 (0.50) |
| $SR_{W-L}^{\gamma\gamma}$ | 6 | 3.7 ± 1.1 | 8.6 | $5.8^{+2.8}_{-1.6}$ | 0.238 | $0.161^{+0.078}_{-0.044}$ | 1.06 (0.14) |
| $SR_{W-H}^{\gamma\gamma}$ | 1 | $2.05^{+0.65}_{-0.63}$ | 3.7 | $4.4^{+1.9}_{-1.0}$ | 0.103 | $0.122^{+0.053}_{-0.028}$ | 0.00 (0.50) |
| $SR_L^{\gamma j}$ | 4 | $1.33^{+0.54}_{-0.32}$ | 7.6 | $4.7^{+1.6}_{-0.8}$ | 0.210 | $0.130^{+0.044}_{-0.022}$ | 1.81 (0.035) |
| $SR_{L200}^{\gamma j}$ | 8 | $2.68^{+0.64}_{-0.63}$ | 11.5 | $5.4^{+2.2}_{-1.2}$ | 0.318 | $0.151^{+0.060}_{-0.033}$ | 2.36 (0.009) |
| $SR_{H}^{\gamma j}$ | 3 | $1.14^{+0.61}_{-0.36}$ | 6.6 | $5.9^{+1.8}_{-1.1}$ | 0.183 | $0.162^{+0.050}_{-0.030}$ | 1.20 (0.116) |



For BESIII --- Example

$\chi_{c2}(1P)$

$$I^{G}(J^{PC}) = 0^{+}(2^{+})$$

Mass $m = 3556.20 \pm 0.09$ MeV Full width $\Gamma = 1.93 \pm 0.11$ MeV

Chic2->etac(1S) + G(-> $\gamma\gamma$) ?

| $\chi_{c2}(1P)$ DECAY MODES | Fraction (Γ_i/Γ) | Confidence level | р (MeV/c) | |
|-----------------------------------|------------------------------|-----------------------|--------------|-----|
| $J/\psi(1S)\pi^{+}\pi^{-}\pi^{0}$ | < 1.5 | % | 90% | 185 |
| $\pi^0 \eta_c$ | < 3.2 | imes 10 ⁻³ | 90% | 512 |
| $\eta_c(1S)\pi^+\pi^-$ | < 5.4 | imes 10 ⁻³ | 90% | 459 |

Citation: K.A. Olive et al. (Particle Data Group), Chin. Phys. C38, 090001 (2014) (URL: http://pdg.lbl.gov)

LIGHT \tilde{G} (Gravitino) MASS LIMITS FROM COLLIDER EXPERIMENTS

The following are bounds on light (\ll 1 eV) gravitino indirectly inferred from its coupling to matter suppressed by the gravitino decay constant.

Unless otherwise stated, all limits assume that other supersymmetric particles besides the gravitino are too heavy to be produced. The gravitino is assumed to be undetected and to give rise to a missing energy $(\not E)$ signature.

| CL% DOCUMENT ID | | | TECN | COMMENT | | | | | |
|---|--|--|---|--|--|--|--|--|--|
| ullet $ullet$ ullet $ullet$ $ullet$ $ullet$ $ullet$ $ullet$ u | | | | | | | | | |
| 95 | ¹ ABDALLAH | 05 B | DLPH | $e^+e^- \rightarrow \widetilde{G}\widetilde{G}\gamma$ | | | | | |
| 95 | ² ACHARD | 04E | L3 | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| | ³ HEISTER | 03 C | ALEP | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁴ ACOSTA | 02H | CDF | $p \overline{p} \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁵ ABBIENDI,G | 00 D | OPAL | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁶ ABREU | 00Z | DLPH | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁷ AFFOLDER | 00J | CDF | $p \overline{p} \rightarrow \widetilde{G} \widetilde{G} + \text{jet}$ | | | | | |
| 95 | ⁸ ACCIARRI | 99R | L3 | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁹ ACCIARRI | 98V | L3 | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| 95 | ⁹ BARATE | 98J | ALEP | $e^+ e^- \rightarrow \widetilde{G} \widetilde{G} \gamma$ | | | | | |
| | <u>CL%</u> use the fo 95 95 95 95 95 95 95 95 95 | CL%DOCUMENT IDuse the following data for av951952ACHARD3HEISTER954ACOSTA955ABBIENDI,G956ABREU957AFFOLDER959ACCIARRI959BARATE | CL%DOCUMENT IDuse the following data for average951 ABDALLAH952 ACHARD952 ACHARD954 ACOSTA955 ABBIENDI,G956 ABREU957 AFFOLDER959 ACCIARRI959 BARATE959 BARATE | CL%DOCUMENT IDTECNuse the following data for averages, fits, li951 ABDALLAH952 ACHARD952 ACHARD954 ACOSTA955 ABBIENDI,G956 ABREU957 AFFOLDER959 ACCIARRI959 ACCIARRI959 BARATE959 BARATE | | | | | |

e+e->(J/psi) ->
$$G + G + \gamma$$

Quarkonium decay into photon plus graviton: a golden channel to discriminate General Relativity from Massive Gravity?

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(Dated: November 27, 2017)

Abstract

After the recent historical discovery of gravitational wave, it is curious to speculate upon the detection prospect of the quantum graviton in the terrestrial accelerator-based experiment. We carefully investigate the "golden" channels, $J/\psi(\Upsilon) \rightarrow \gamma + \text{graviton}$, which can be pursued at BESIII and Belle 2 experiments, by searching for single-photon plus missing energy events. Within the effective field theory (EFT) framework of General Relativity (GR) together with Nonrelativistic QCD (NRQCD), we are capable of making solid predictions for the corresponding decay rates. It is found that these extremely suppressed decays are completely swamped by the Standard Model background events $J/\psi(\Upsilon) \to \gamma + \nu \bar{\nu}$. Meanwhile, we also study these rare decay processes in the context of massive gravity, and find the respective decay rates in the limit of vanishing graviton mass drastically differ from their counterparts in GR. Counterintuitive as the failure of smoothly recovering GR results may look, our finding is reminiscent of the van Dam-Veltman-Zakharov (vDVZ) discontinuity widely known in classical gravity, which can be traced to the finite contribution of the helicity-zero graviton in the massless limit. Nevertheless, at this stage we are not certain about the fate of the discontinuity encountered in this work, whether it is merely a pathology or not. If it could be endowed with some physical significance, the future observation of these rare decay channels, would, in principle, shed important light on the nature of gravitation, whether the graviton is strictly massless, or bears a very small but nonzero mass.

PACS numbers: 04.60.Bc, 14.40.Pq, 14.70.Kv



FIG. 1: Four LO Feynman diagrams for $c\bar{c}({}^{3}S_{1}^{(1)}) \rightarrow \gamma + \mathcal{G}$.

but I feel, the conclusion of this article is not so much attractive. . .

arXiv:1711.09058v1 [hep-ph] 24 Nov 2017