

5-body contributions to inclusive semileptonic rare B decays

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Outline

- Motivation
- Calculation
 - Amplitudes
 - Phase space integral
 - Results
- Summary and Outlook

Motivation

Motivation

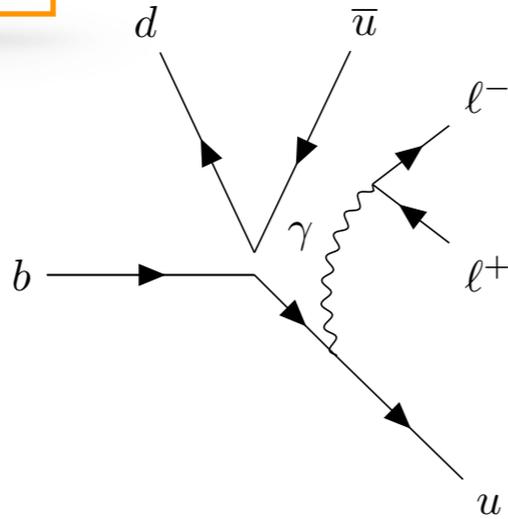
- For $\bar{B} \rightarrow X_s \ell^+ \ell^-$, a thorough phenomenological analysis is available in the market [Huber, Hurth & Lunghi, 15']
- A similar study for $\bar{B} \rightarrow X_d \ell^+ \ell^-$ is in order (for Belle II). What's new?

$P_{1,2}^u = (\bar{d}\gamma^\mu P_L u)(\bar{u}\gamma_\mu P_L b)$ not CKM suppressed for b-d transition \Rightarrow

- new NLO & NNLO QCD corrections [Greub et al, 03; Seidel, 04']
- new power corrections [Ali, Hiller, Handoko & Morozumi 96'; Ligeti & Tackmann 07; Benzke, Turkczyk & Hurth, 17']
- new log-enhanced EM contributions [Huber, Hurth & Lunghi, 07']
- ...
- five-body processes at the quark level (totally new)

Motivation

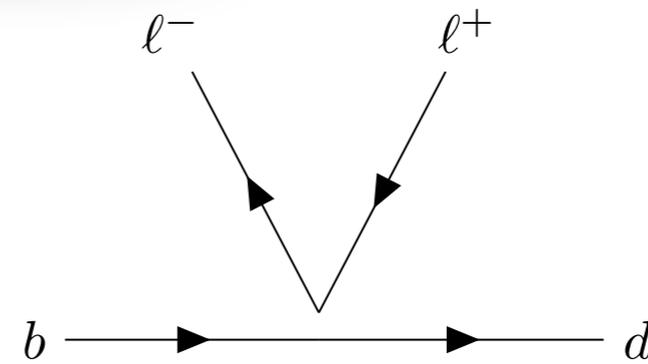
5 body



$$P_{1,2}^u = (\bar{d}\gamma^\mu P_L u)(\bar{u}\gamma_\mu P_L b)$$

$$|V_{ud}^* V_{ub}|^2 \alpha_s^2 \kappa^2 C_{1,2}^2$$

3 body



$$P_9 = (\bar{d}\gamma^\mu P_L b)(\bar{l}\gamma_\mu l)$$

$$|V_{td}^* V_{tb}|^2 C_9^2$$

$$\kappa \equiv \frac{\alpha_e}{\alpha_s}$$

- Formally, $\lambda^4 \alpha_s^2 \kappa^2$ vs $\lambda^4 \kappa^2$ ($C_9 \sim \kappa$), 5-body is necessary to complete NNLO
 - Numerically, $\lambda^4 \alpha_s^2 \kappa^2$ vs $\lambda^4 \alpha_s^2 \kappa^2$ ($C_9 \sim \alpha_s \kappa$), LO!!
- only suppressed by phase space

Calculation

Calculation

$$\Gamma = \frac{1}{2m_b} \int |\mathcal{A}|^2 d\text{PS}$$

- Write down the decay amplitude
- square it
- and make the phase-space integration

Amplitudes

Amplitudes

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QCD} \times \text{QED}}(u, d, s, c, b, e, \mu, \tau)$$

$$- \frac{4G_F}{\sqrt{2}} V_{ud}^* V_{ub} (C_1^u P_1^u + C_2^u P_2^u) + \frac{4G_F}{\sqrt{2}} V_{td}^* V_{tb} \sum_{i=3}^{10} C_i(\mu) P_i,$$

$$P_1^u = (\bar{d}_L \gamma_\mu T^a u_L) (\bar{u}_L \gamma^\mu T^a b_L),$$

$$P_2^u = (\bar{d}_L \gamma_\mu u_L) (\bar{u}_L \gamma^\mu b_L),$$

$$P_3 = (\bar{d}_L \gamma_\mu b_L) \sum_q (\bar{q} \gamma^\mu q),$$

$$P_4 = (\bar{d}_L \gamma_\mu T^a b_L) \sum_q (\bar{q} \gamma^\mu T^a q),$$

$$P_5 = (\bar{d}_L \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} b_L) \sum_q (\bar{q} \gamma^{\mu_1} \gamma^{\mu_2} \gamma^{\mu_3} q),$$

$$P_6 = (\bar{d}_L \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} T^a b_L) \sum_q (\bar{q} \gamma^{\mu_1} \gamma^{\mu_2} \gamma^{\mu_3} T^a q),$$

$$P_7 = \frac{e}{16\pi^2} m_b (\bar{d}_L \sigma^{\mu\nu} b_R) F_{\mu\nu},$$

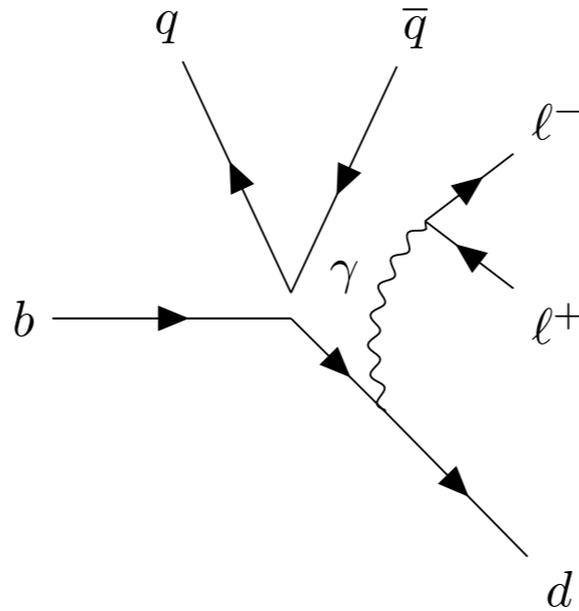
$$P_8 = \frac{g}{16\pi^2} m_b (\bar{d}_L \sigma^{\mu\nu} T^a b_R) G_{\mu\nu}^a,$$

$$P_9 = (\bar{d}_L \gamma_\mu b_L) \sum_l (\bar{l} \gamma^\mu l),$$

$$P_{10} = (\bar{d}_L \gamma_\mu b_L) \sum_l (\bar{l} \gamma^\mu \gamma_5 l).$$

Amplitudes

$$P_{1-6}^{(u)} = (\bar{q}\Gamma_1 q)(\bar{d}\Gamma_2 b)$$



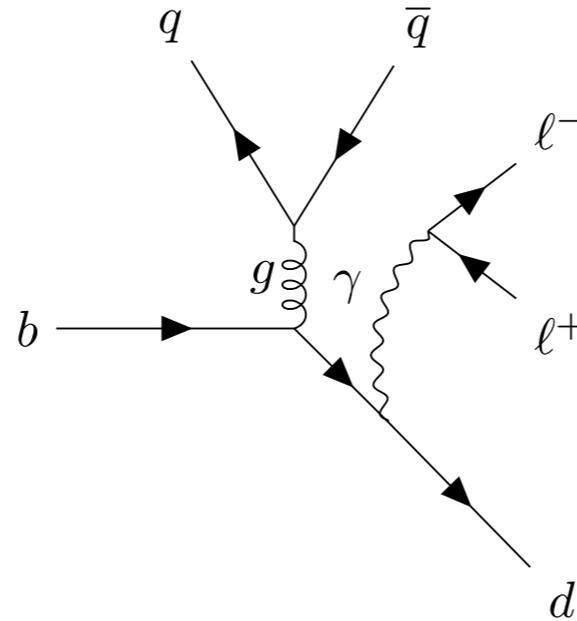
$$4 \times 6$$

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4

Amplitudes

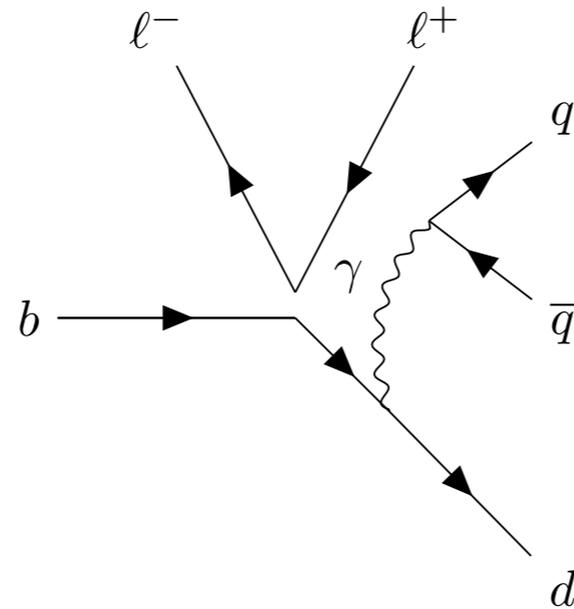
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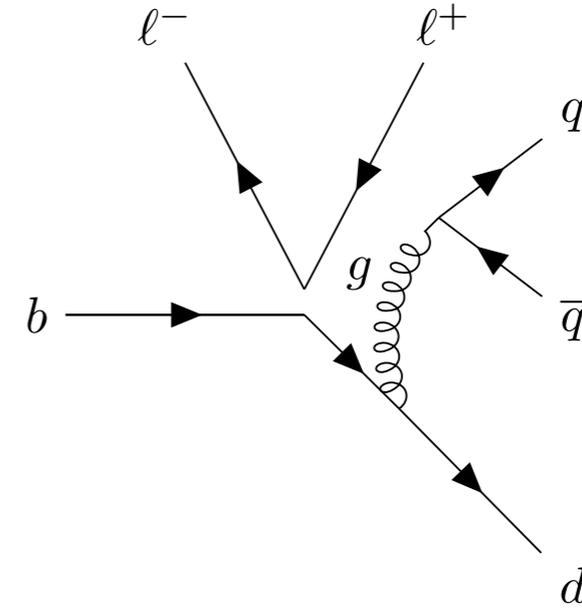
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$$P_{9,10} = (\bar{d}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu (\gamma_5) \ell)$$



$$4 \times 2$$



$$2 \times 2$$

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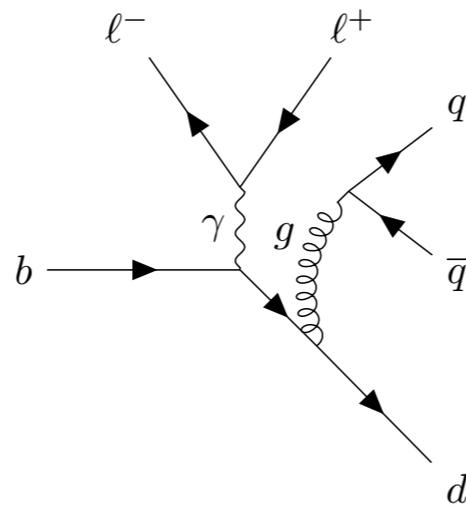
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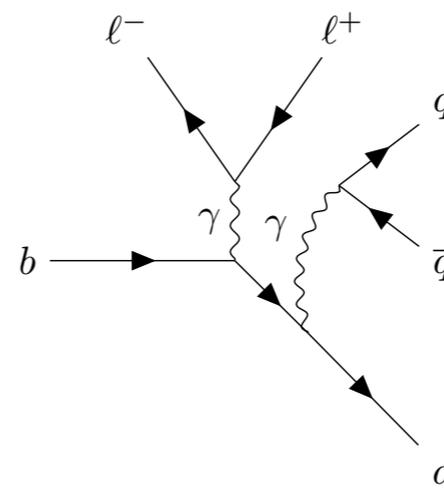
$$P_{9,10} = (\bar{d}_L \gamma_\mu b_L) (\bar{\ell} \gamma^\mu (\gamma_5) \ell)$$

$$4 \times 2 + 2 \times 2$$

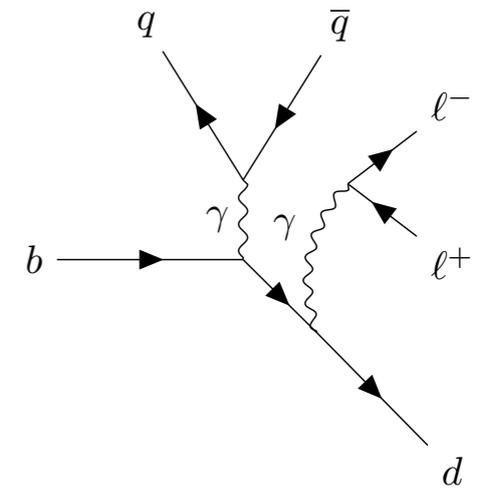
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$$2 + 4 + 4$$

In total, 50

Amplitudes

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- Amplitude square

- 50*50 = 2500 (only for u)
- 42*42 (for s)
- 42*42 (for d)?

In total, 50

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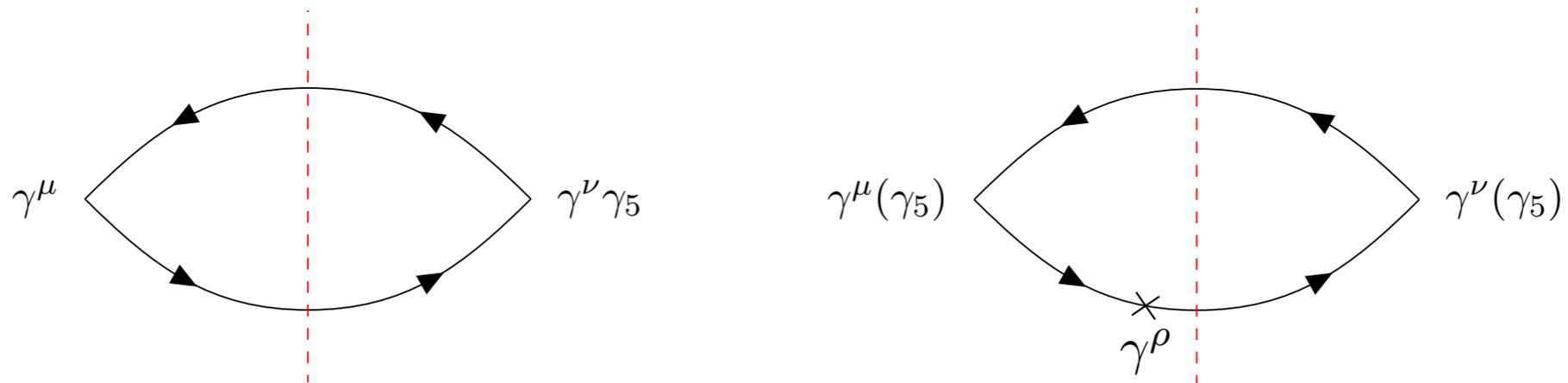
- 50*50 = 2500 (only for u)
- 42*42 (for s)
- 42*42 (for d)?
- 84*84 (for d)

identical particle

In total, 50

Amplitudes

- Rules to reduce terms
 - color conservation
 - what vanishes in the sense of exchanging fermions



- $A_s = A_u(Q_u \rightarrow Q_s)$, $A_d(I) = A_u(Q_u \rightarrow Q_d)$
- $Q_u + Q_d + Q_s = 0$
- power counting up to $\alpha_s^3 \kappa^3$

Amplitudes

- Power counting

\mathcal{P}_{1-6}	\mathcal{P}_7	\mathcal{P}_8	\mathcal{P}_9	\mathcal{P}_{10}
$\alpha_s^2 \kappa^2$	$\begin{pmatrix} \text{QCD } \alpha_s^3 \kappa^2 \\ \text{QED } \alpha_s^3 \kappa^3 \end{pmatrix}$	$\alpha_s^3 \kappa^2$	$\begin{pmatrix} \text{QCD } \alpha_s^3 \kappa^2 \\ \text{QED } \alpha_s^3 \kappa^3 \end{pmatrix}$	$\begin{pmatrix} \text{QCD } \alpha_s^3 \kappa^2 \\ \text{QED } \alpha_s^3 \kappa^3 \end{pmatrix}$
...	×	×	×	×
...	×	×	×	×
...	×	×	×	×
...	×	×	×	×

$$C_{1-8} \sim 1$$

$$C_9 \sim 0.04\kappa + 2\alpha_s\kappa + \dots$$

$(\sim \alpha_s\kappa)$

$$C_{10} \sim 0.005\kappa^2 - 4\alpha_s\kappa + \dots$$

$(\sim \alpha_s\kappa^2)$

Phase space integral

Phase space integral

- Start from

$$\int d\Phi_5 \equiv \int \prod_{i=1}^5 \frac{d^3 p_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta^4(p_b - \sum_{i=1}^5 p_i)$$

- an issue of simplifying the calculation again

$$|\mathcal{A}|^2 = |\mathcal{A}|_I^2 (\{p_i \cdot p_j\}) + \text{Re} (\mathcal{A}_{II}^2 (\{p_i \cdot p_j\}) i\epsilon_{\mu\nu\rho\sigma} p_i^\mu p_j^\nu p_k^\rho p_l^\sigma)$$

Under the phase space int. $\int \prod_{i=1}^5 \frac{d^3 p_i}{(2\pi)^3 2E_i}$

p_i^0 is even, p_i^1-3 are odd

⇒ $\text{Re} (\mathcal{A}_{II}^2 (\{p_i \cdot p_j\}) i\epsilon_{\mu\nu\rho\sigma} p_i^\mu p_j^\nu p_k^\rho p_l^\sigma)$ must be odd

⇒ $\int d\Phi_5 |\mathcal{A}|^2 = \int d\Phi_5 |\mathcal{A}|_I^2 (\{p_i \cdot p_j\})$

Phase space integral

- The integral formula [Kumar, 69']

$$\begin{aligned}
 \int d\Phi_5 &\equiv \int \prod_{i=1}^5 \frac{d^3 p_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta^4(p_b - \sum_{i=1}^5 p_i) \\
 &= \frac{\pi^2 m_b^6}{4(2\pi)^{11}} \int_0^1 ds_3 \int_{s_3}^1 ds_2 \int_{s_2}^1 ds_1 \int_{u_{1-}}^{u_{1+}} du_1 \int_{u_{2-}}^{u_{2+}} du_2 \int_{u_{3-}}^{u_{3+}} du_3 \int_{t_{2-}}^{t_{2+}} dt_2 \int_{t_{3-}}^{t_{3+}} dt_3 \\
 &\quad \times \frac{1}{(1-u_1)(1-u_2)(1-u_3) \sqrt{(1-\eta_2^2)(1-\xi_2^2) - (\omega_2 - \eta_2 \xi_2)^2}} \\
 &\quad \times \frac{1}{\sqrt{(1-\eta_3^2)(1-\xi_3^2) - (\omega_3 - \eta_3 \xi_3)^2} \sqrt{\lambda(1, s_2, -s_1 + s_2 - u_1 + 1)}} \\
 &\quad \times \frac{1}{\sqrt{\lambda(1, t_2, t_2 - u_1 - u_2 + 1)} \sqrt{\lambda(1, s_3, -s_1 + s_3 - u_1 - u_2 + 2)}},
 \end{aligned}$$

$$s_n = \frac{(p_b - \sum_{i=1}^n p_i)^2}{m_b^2}, \quad u_n = \frac{(p_b - p_{n+1})^2}{m_b^2}, \quad t_n = \frac{(p_b - \sum_{i=2}^{n+1} p_i)^2}{m_b^2}.$$

Phase space integral

$$p_i \cdot p_j \Rightarrow f(s, t, u)$$

$$\begin{aligned} \frac{p_1 \cdot p_b}{m_b^2} &\rightarrow \frac{1 - s_1}{2}, & \frac{p_2 \cdot p_b}{m_b^2} &\rightarrow \frac{1 - u_1}{2}, & \frac{p_3 \cdot p_b}{m_b^2} &\rightarrow \frac{1 - u_2}{2}, & \frac{p_4 \cdot p_b}{m_b^2} &\rightarrow \frac{1 - u_3}{2}, \\ \frac{p_1 \cdot p_2}{m_b^2} &\rightarrow \frac{1 - s_1 + s_2 - u_1}{2}, & \frac{p_1 \cdot p_3}{m_b^2} &\rightarrow \frac{u_1 - s_2 + s_3 - t_2}{2}, & \frac{p_1 \cdot p_4}{m_b^2} &\rightarrow \frac{t_2 - s_3 - t_3}{2}, \\ \frac{p_2 \cdot p_3}{m_b^2} &\rightarrow \frac{1 + t_2 - u_1 - u_2}{2}, & p_2 \cdot p_4 &= \frac{1 - t_2 + t_3 - u_3}{2} - p_3 \cdot p_4 \end{aligned}$$

An exercise of counting numbers

$$\left(\sum_{i=1}^5 p_i \right)^2 = m_b^2$$

No. of independent p-products: $4 \cdot 5 / 2 - 1 = 9$

No. of integral variables: $3 \cdot 5 - 4 - 3 = 8$

One p-product redundant?

$$\delta^4 \left(p_b - \sum_{i=1}^5 p_i \right)$$

3 space angles

Phase space integral

One p-product redundant?

The answer is yes.

In the dim-4 space-time, at most 4 independent momenta, so

$$\sum_{i=1}^5 c_i p_i = 0 \quad \Rightarrow \quad \det \begin{pmatrix} p_1 \cdot p_1 & p_1 \cdot p_2 & p_1 \cdot p_3 & p_1 \cdot p_4 & p_1 \cdot p_5 \\ p_2 \cdot p_1 & p_2 \cdot p_2 & p_2 \cdot p_3 & p_2 \cdot p_4 & p_2 \cdot p_5 \\ p_3 \cdot p_1 & p_3 \cdot p_2 & p_3 \cdot p_3 & p_3 \cdot p_4 & p_3 \cdot p_5 \\ p_4 \cdot p_1 & p_4 \cdot p_2 & p_4 \cdot p_3 & p_4 \cdot p_4 & p_4 \cdot p_5 \\ p_5 \cdot p_1 & p_5 \cdot p_2 & p_5 \cdot p_3 & p_5 \cdot p_4 & p_5 \cdot p_5 \end{pmatrix} = 0$$

$$\Rightarrow p_3 \cdot p_4 = [111 \text{ terms}] \pm \sqrt{[2179 \text{ terms}]}$$

The solution is only the first part, according to [Kumar, 69'], and this is wrong!

Phase space integral

Kumar's method

$$\int d^4 p_4 \delta(p_4^2) \delta((p_b - p_{1-4})^2) \delta((p_b - p_{2-4})^2 - t_3) \delta((p_b - p_4)^2 - u_3) p_4^\mu$$
$$= (\alpha p_{1-3} + \beta p_{2-3} + \gamma p_b)^\mu \times I$$

✓

4 delta fix p4 ↪

$$p_4 = \alpha p_{1-3} + \beta p_{2-3} + \gamma p_b$$

[Kumar, 69']

X

The key point: the 4 delta don't fix one p4 point, but two (two solutions).

$$p_4(\text{I}) + p_4(\text{II}) = \alpha p_{1-3} + \beta p_{2-3} + \gamma p_b$$

✓

Results

Results

One example (dominant)

$$|\langle d\ell^+ \ell^- u\bar{u} | P_{1,2}^u | b \rangle|^2$$

First, I show the results before phase space integration (without expression of $p_3^* p_4$ substituted).

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$$|\langle d\ell^+ \ell^- u\bar{u} | P_{1,2}^u | b \rangle|^2$$

First, I show the results before phase space integration (without expression of $p_3^* p_4$ substituted).

Then, let's have a look at the result after the integration, (with only one variable s , the dilepton mass square, unintegrated.)

$$\left| \langle d\ell^+ \ell^- u\bar{u} | P_{1,2}^u | b \rangle \right|^2$$

$$\begin{aligned} & \tilde{\alpha}_e^2 Q_L^2 \left[\frac{Q_d Q_u}{81} \left(-108\sqrt{4-s}\sqrt{s} (4s^2 - 7s - 6) f_1(s) \ln(s) + 108 (4s^3 - 15s^2 + 10) [f_1(s)]^2 \right. \right. \\ & + 216\sqrt{4-s}\sqrt{s} (4s - 1) f_1(s) - 72(s + 1) (4s^2 + 5s - 5) f_2(s) - 9 (12s^3 - 45s^2 - 10) \ln^2(s) \\ & + 432f_3(s) - \frac{(s - 1) (955s^3 - 425s^2 - 749s + 87)}{s} + \left. \frac{12 (84s^3 - 15s^2 - 56s + 3) \ln(s)}{s} \right) \\ & + \frac{Q_d^2}{324} \left(216\sqrt{4-s}\sqrt{s} (s^2 + 2s - 12) f_1(s) \ln(s) - \frac{216 (s^4 - 18s^2 + 16s + 6) [f_1(s)]^2}{s} \right. \\ & + \frac{108 (4s^3 + 13s^2 - 110s - 60) f_1(s)}{\sqrt{4-s}\sqrt{s}} + \frac{18 (3s^4 - 18s^2 + 16s + 6) \ln^2(s)}{s} \\ & + \left. \frac{(s - 1) (259s^3 + 43s^2 + 133s - 1863)}{s} - \frac{6 (36s^3 + 477s^2 - 172s - 120) \ln(s)}{s} \right) \\ & + \frac{Q_u^2}{324} \left(\frac{144(s + 1) (s^3 + 5s^2 + 13s - 3) f_2(s)}{s} + \frac{72 (18s^2 + 10s - 3) \ln^2(s)}{s} + 1728f_3(s) \right. \\ & + \left. \frac{(s - 1) (259s^3 + 715s^2 + 6835s - 2901)}{s} - \frac{12 (24s^3 + 348s^2 + 148s - 111) \ln(s)}{s} \right) \Big] \end{aligned}$$

$$f_1(s) = \frac{\pi}{6} - \arctan\sqrt{s/(4-s)},$$

$$f_2(s) = 2 \text{Li}_2(-s) - \ln^2(s) + 2 \ln(s) \ln(1+s) + \zeta_2,$$

$$f_3(s) = 4 \text{Li}_3(-s) - 2 \text{Li}_2(-s) \ln(s) - \frac{1}{6} \ln^3(s) + \zeta_2 \ln(s) + 3\zeta_3,$$

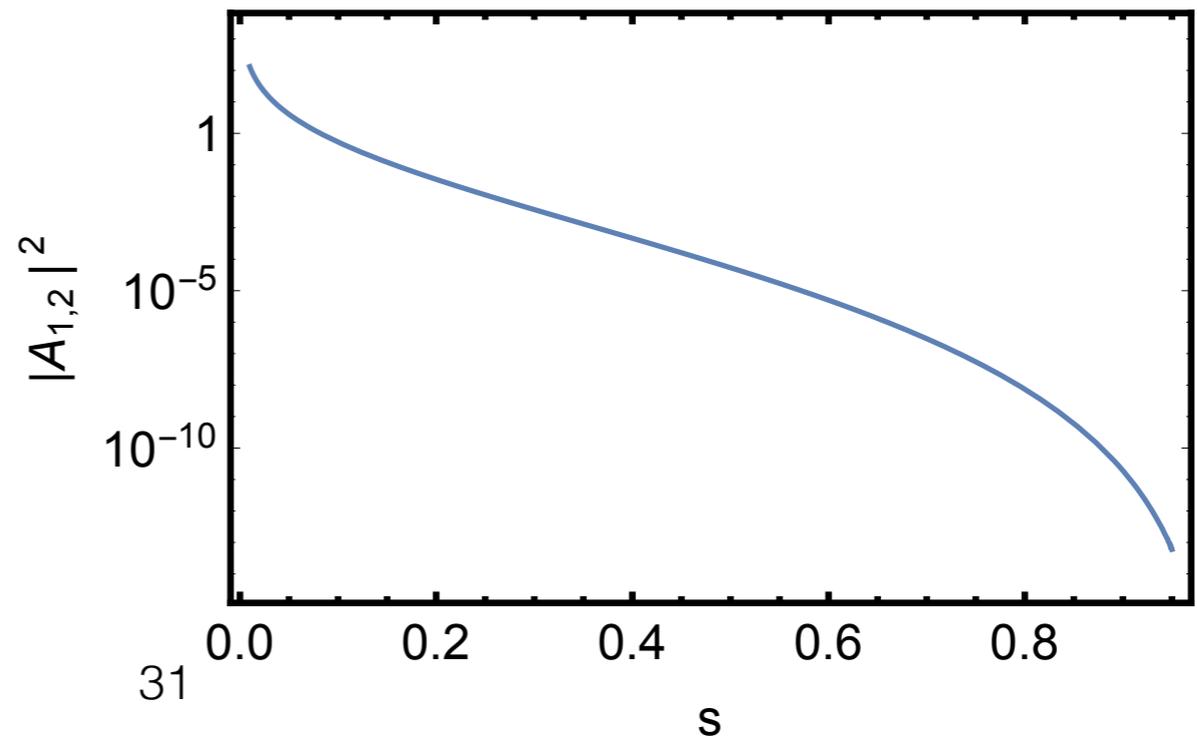
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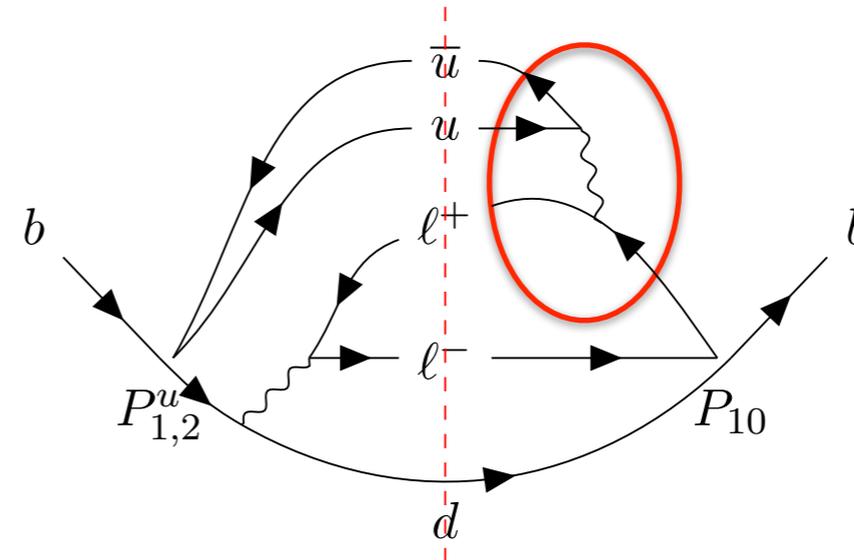
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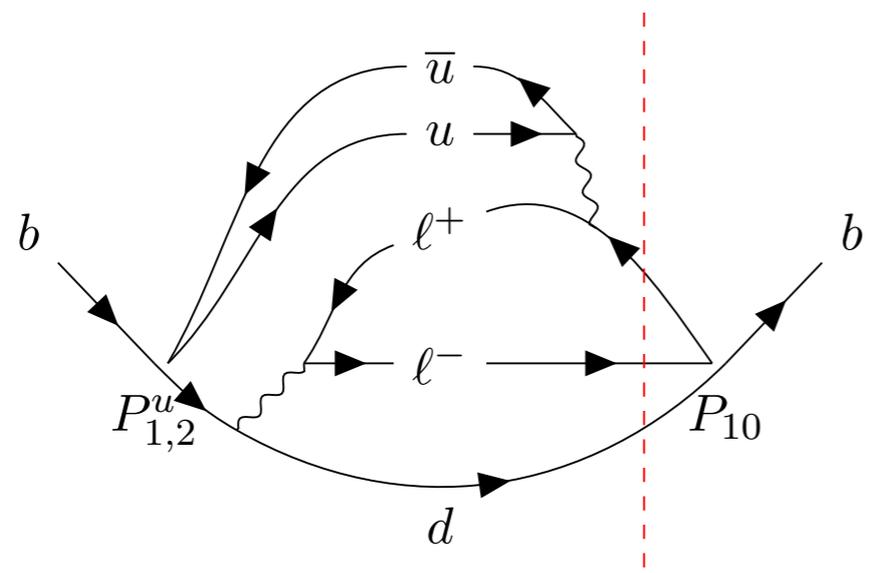
IR divergence

P_{10} interference with others

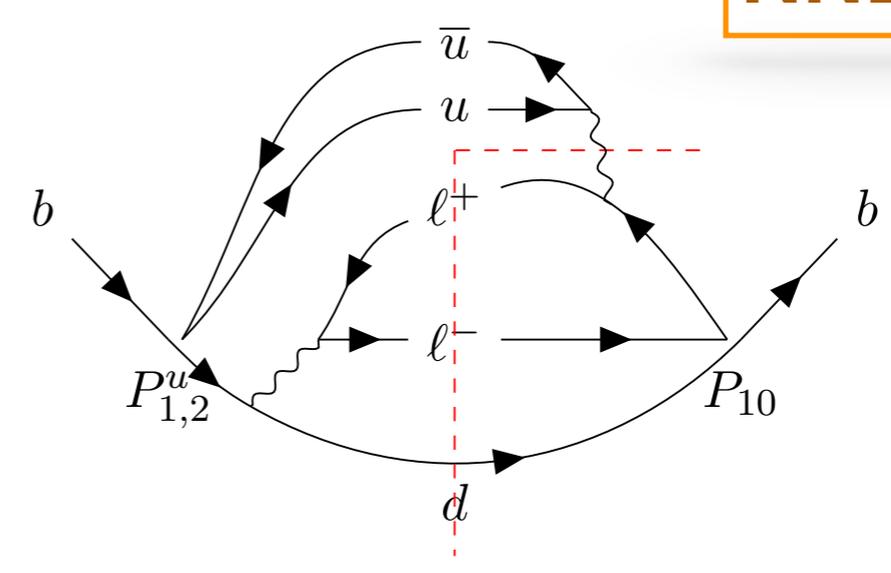


To cancel the divergence

NNLO QED



3 body, 2 loop



4 body, 1 loop

Numerical results (preliminary)

Branching ratio (5 body compared to 3 body)

$M_{\ell\ell}^2$	[1, 6] GeV ²	[1, 3.5] GeV ²
$\mathcal{B}(5)$ [10^{-10}]	9.6	9.3
$\mathcal{B}(3)$ [10^{-8}]	6.8	3.8
$\mathcal{B}(5/3)$	1.4%	2.4%

Percentage of other contributions

- 1/mb corrections 0.2%
- 1/mc corrections -1%
- log-enhanced QED corrections 5% (e), 2% (muon)

Summary and Outlook

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Thank you for your attention!