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Study of the light baryon states in Λ_c^+ decays

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Outline

Introduction

Possible $\Sigma_{1/2^-}(1380)$ state and a new Λ^ state in $\Lambda_c^+ \rightarrow \eta\pi^+\Lambda$ decay*

Study of $N^(1535)$ in the $\Lambda_c^+ \rightarrow \bar{K}^0\eta p$ decay*

Possible ϕp state in $\Lambda_c^+ \rightarrow \pi^0\phi p$ decay

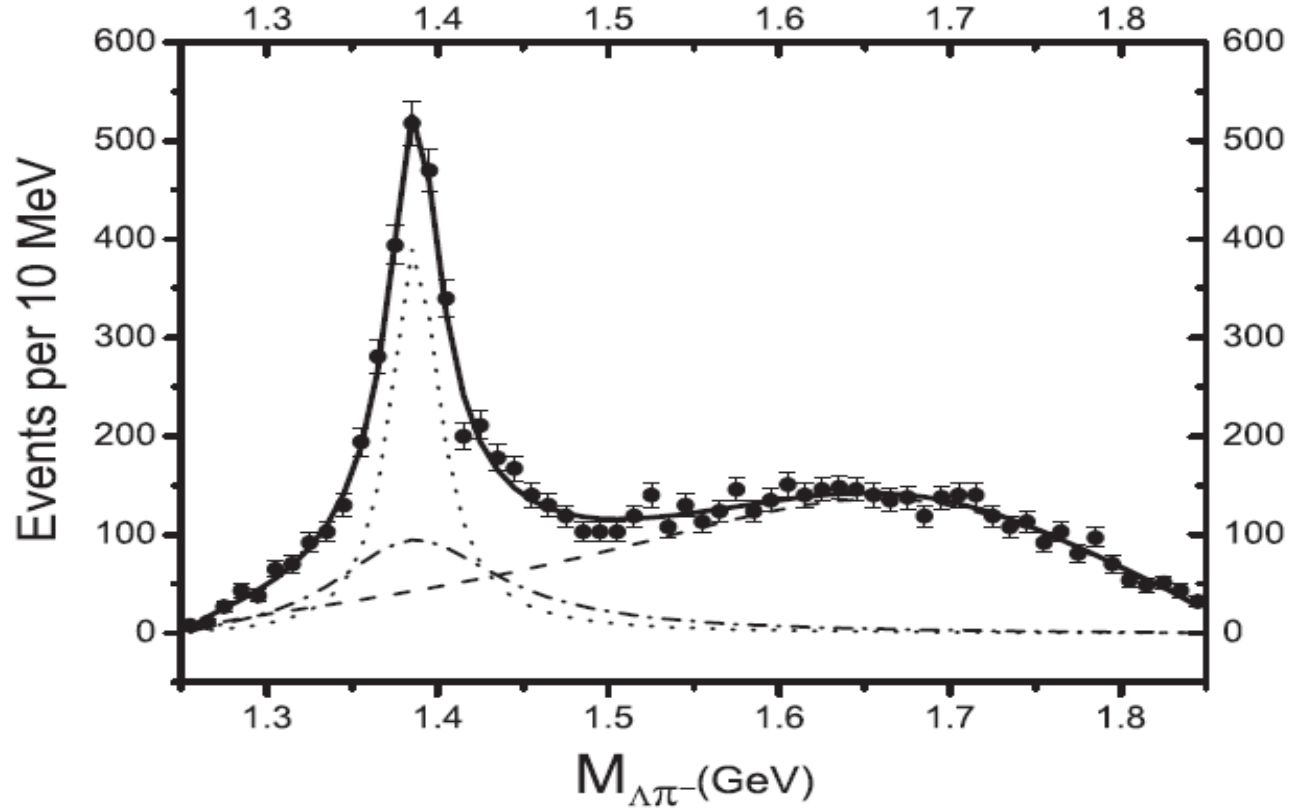
Summary

A possible $\Sigma(1380)$ state with spin-parity $J^P = \frac{1}{2}^-$

Flavor wave functions and masses of the $\frac{1}{2}^-$ pentaquark octet and singlet.

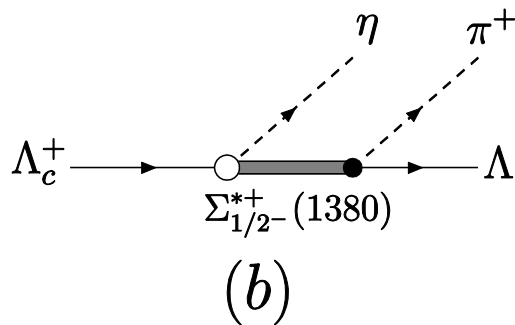
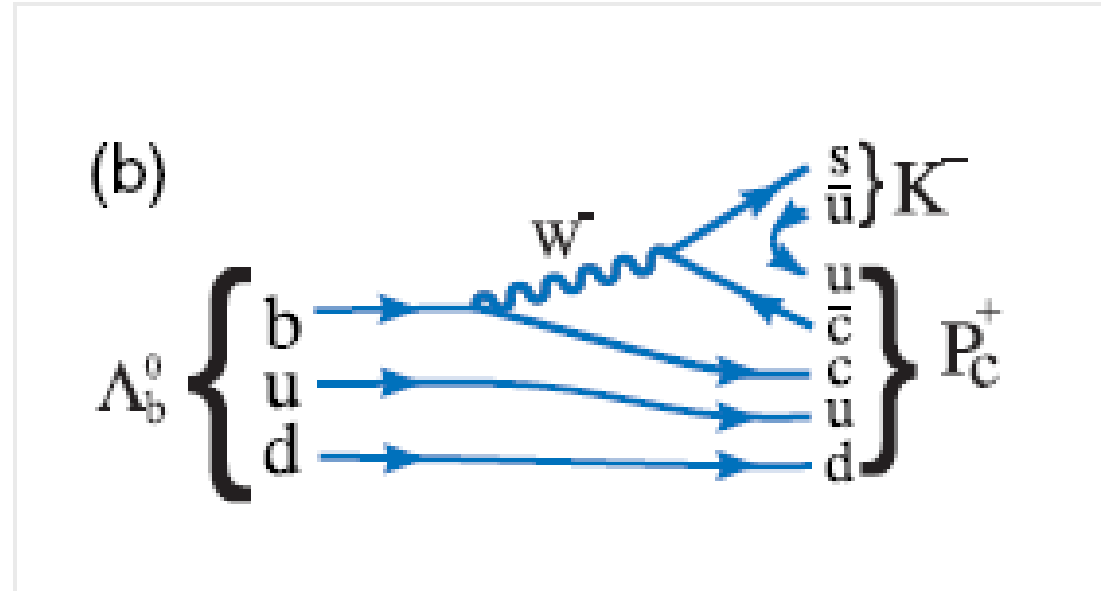
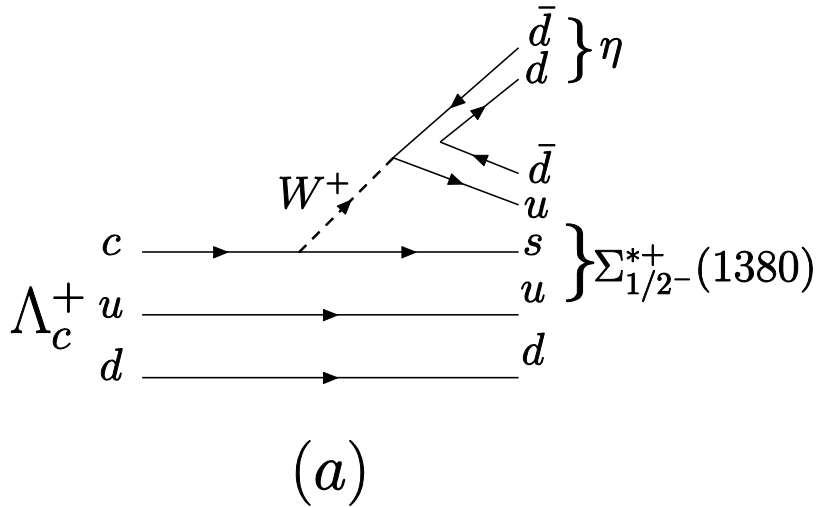
| | (Y, I) | I_3 | flavor wave functions | masses (MeV) |
|--------------|----------|-------|---|--------------|
| Σ_8^+ | $(0, 1)$ | 1 | $[su][ud]_-\bar{d}$ | 1360 |
| Σ_8^0 | | 0 | $\frac{1}{\sqrt{2}}([su][ud]_-\bar{u} + [ds][ud]_-\bar{d})$ | 1360 |
| Σ_8^- | | -1 | $[ds][ud]_-\bar{u}$ | 1360 |

Ao Zhang, Y. R. Liu, P.Z. Huang, W.Z. Deng, X.L. chen and S.L. Zhu,
High Energy Phys. Nucl. Phys. 29, 250 (2005).

Evidence for a new Σ^* resonance with $J^P = 1/2^-$ in the old data of the $K^- p \rightarrow \Lambda \pi^+ \pi^-$ reactionJia-Jun Wu,¹ S. Dulat,^{2,3} and B. S. Zou^{1,3}

| | $M_{\Sigma^{*(3/2)}}$ | $\Gamma_{\Sigma^{*(3/2)}}$ | $M_{\Sigma^{*(1/2)}}$ | $\Gamma_{\Sigma^{*(1/2)}}$ | χ^2/ndf (Fig. 1) |
|------|------------------------|----------------------------|------------------------|----------------------------|------------------------------|
| Fit1 | 1385.3 ± 0.7 | 46.9 ± 2.5 | | | 68.5/54 |
| Fit2 | $1386.1^{+1.1}_{-0.9}$ | $34.9^{+5.1}_{-4.9}$ | $1381.3^{+4.9}_{-8.3}$ | $118.6^{+55.2}_{-35.1}$ | 58.0/51 |

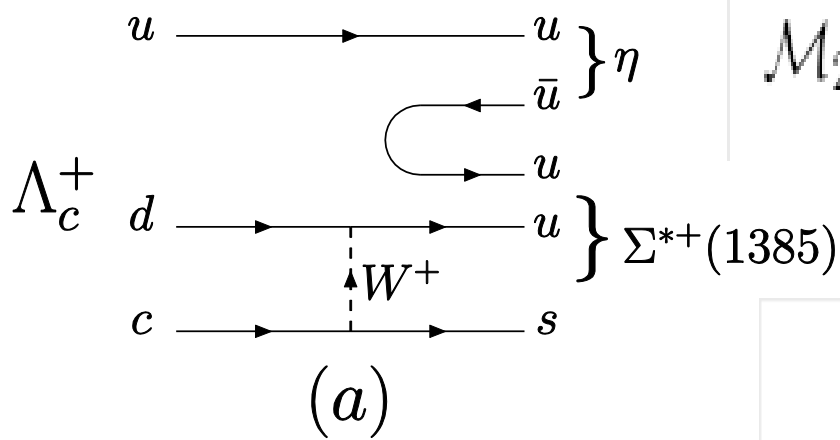
$\Sigma(1380)$ in $\Lambda_c^+ \rightarrow \eta \pi^+ \Lambda$ decay



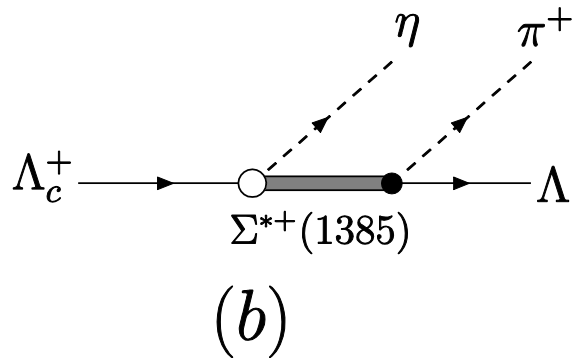
$$\mathcal{M}_1 = ig_{\pi\Lambda\Sigma_1^*} \bar{u}(p_3) G^{\Sigma_1^*}(q) (A_1 + B_1 \gamma_5) u(p),$$

$$G^{\Sigma_1^*}(q) = i \frac{\not{q} + M_{\Sigma_1^*}}{q^2 - M_{\Sigma_1^*}^2 + iM_{\Sigma_1^*}\Gamma_{\Sigma_1^*}},$$

$\Sigma(1385)$ in $\Lambda_c^+ \rightarrow \eta \pi^+ \Lambda$ decay



$$\mathcal{M}_2 = \frac{i g_{\pi \Lambda \Sigma_2^*}}{m_\eta m_\pi} \bar{u}(p_3) p_2^\mu G_{\mu\nu}^{\Sigma_2^*}(q) p_1^\nu (A_2 + B_2 \gamma_5) u(p),$$



$$G_{\mu\nu}^{\Sigma_2^*}(q) = i \frac{\not{q} + M_{\Sigma_2^*}}{q^2 - M_{\Sigma_2^*}^2 + i M_{\Sigma_2^*} \Gamma_{\Sigma_2^*}} P_{\mu\nu},$$

with

$$P^{\mu\nu} = -g^{\mu\nu} + \frac{1}{3} \gamma^\mu \gamma^\nu + \frac{2q^\mu q^\nu}{3M_{\Sigma_2^*}^2} + \frac{\gamma^\mu q^\nu - \gamma^\nu q^\mu}{3M_{\Sigma_2^*}}.$$

Invariant mass, decay angle and energy distributions

$$\frac{d\Gamma}{dM_{\pi+\Lambda}} = \frac{m_{\Lambda}}{32\pi^3 M_{\Lambda_c^+}} \int \sum |\mathcal{M}|^2 |\vec{p}_1| |\vec{p}^*| d\cos\theta^*$$

$$\frac{d\Gamma}{d\cos\theta^*} = \frac{m_{\Lambda}}{32\pi^3 M_{\Lambda_c^+}} \int \sum |\mathcal{M}|^2 |\vec{p}_1| |\vec{p}^*| dM_{\pi+\Lambda}.$$

$$\frac{d\Gamma}{dE_{\pi^+}} = \frac{m_{\Lambda}}{32\pi^3} \int \sum |\mathcal{M}|^2 dE_{\Lambda};$$

Invariant mass distributions

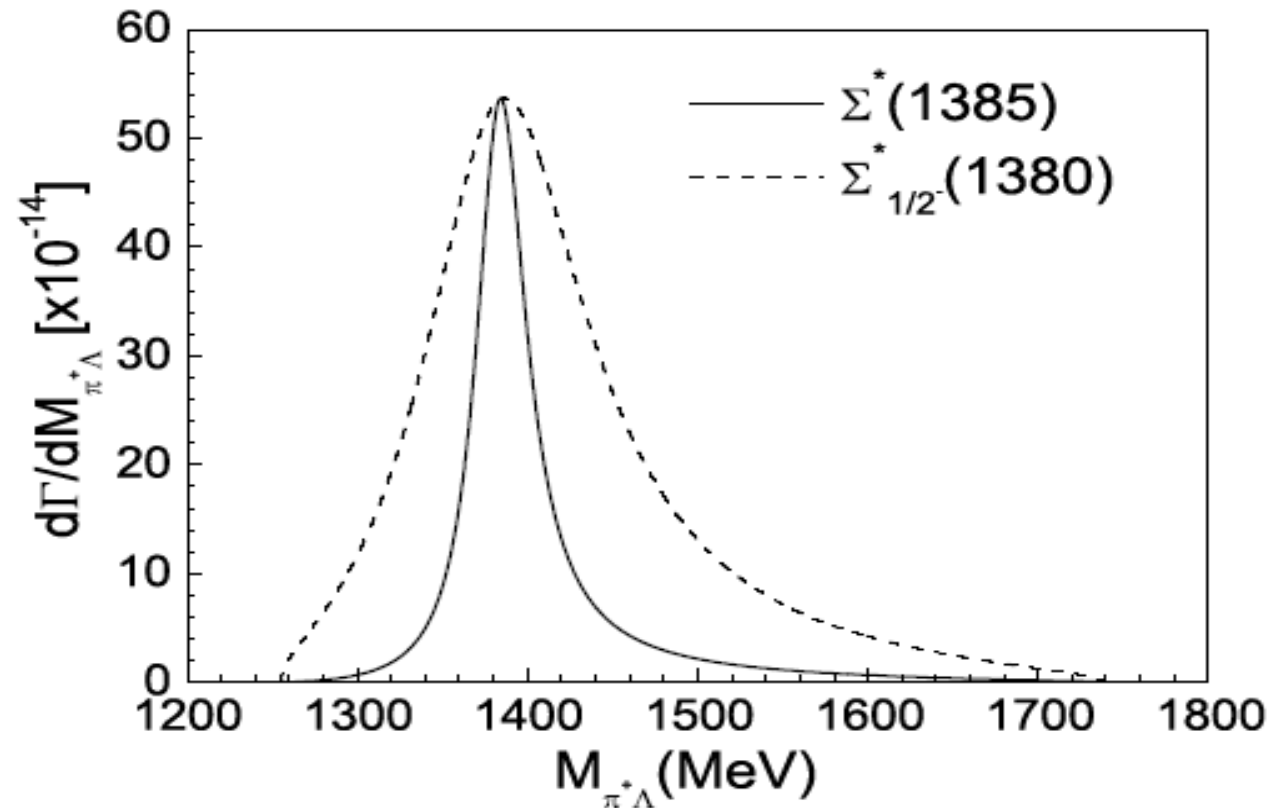


FIG. 5: Invariant mass distributions $d\Gamma/dM_{\pi^+\Lambda}$ as a function of $M_{\pi^+\Lambda}$.

Decay angle and energy distributions

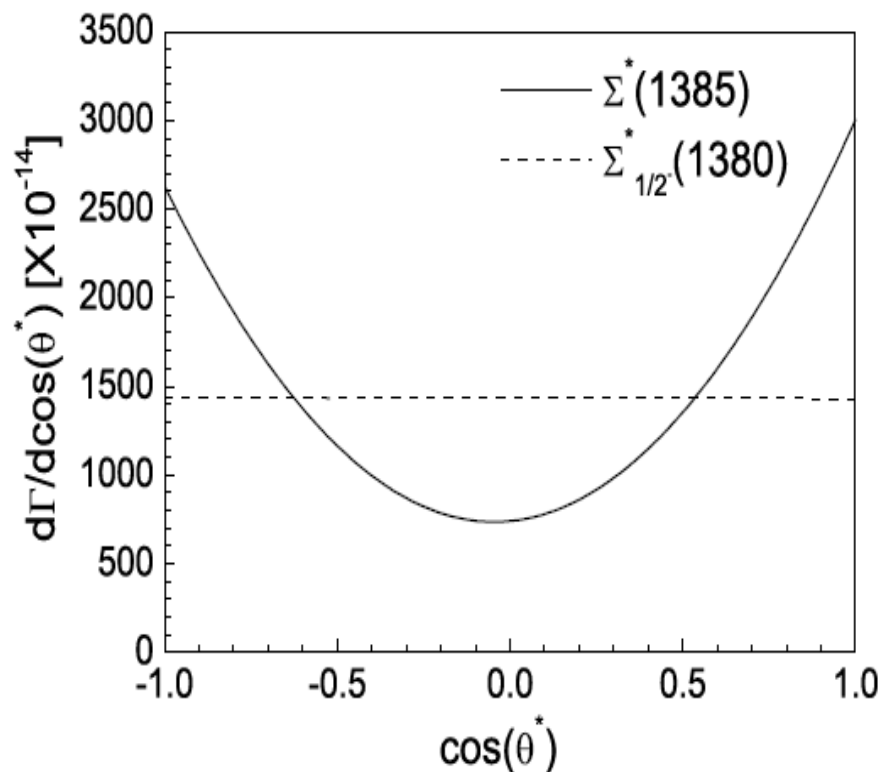


FIG. 6: Angle distributions $d\Gamma/d\cos\theta^*$ in the c.m. frame of $\pi^+\Lambda$ system as a function of $\cos\theta^*$.

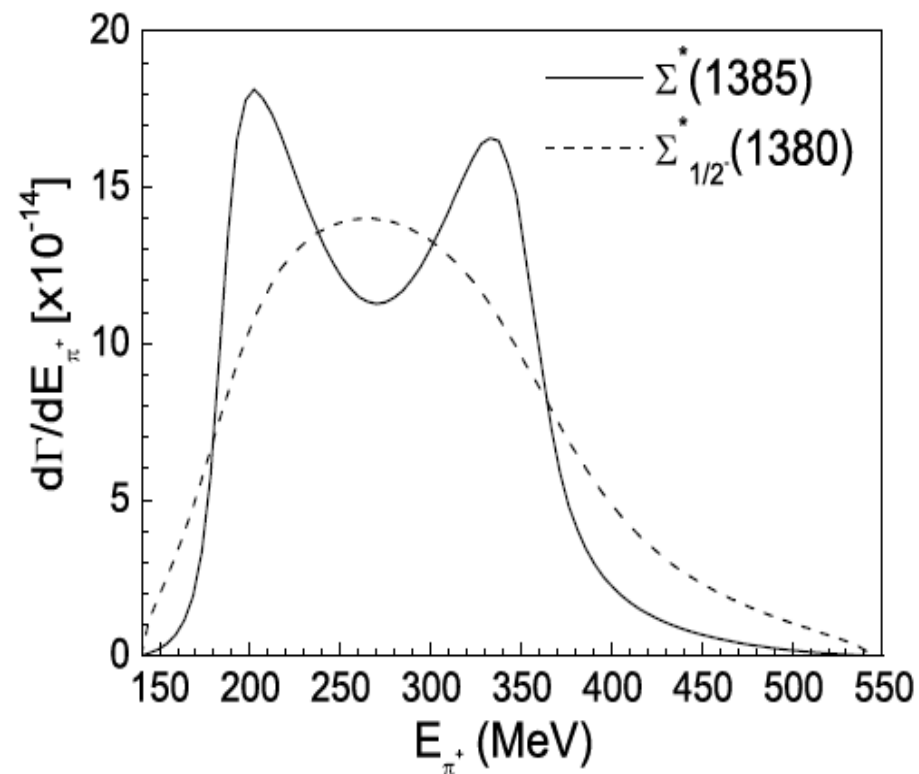
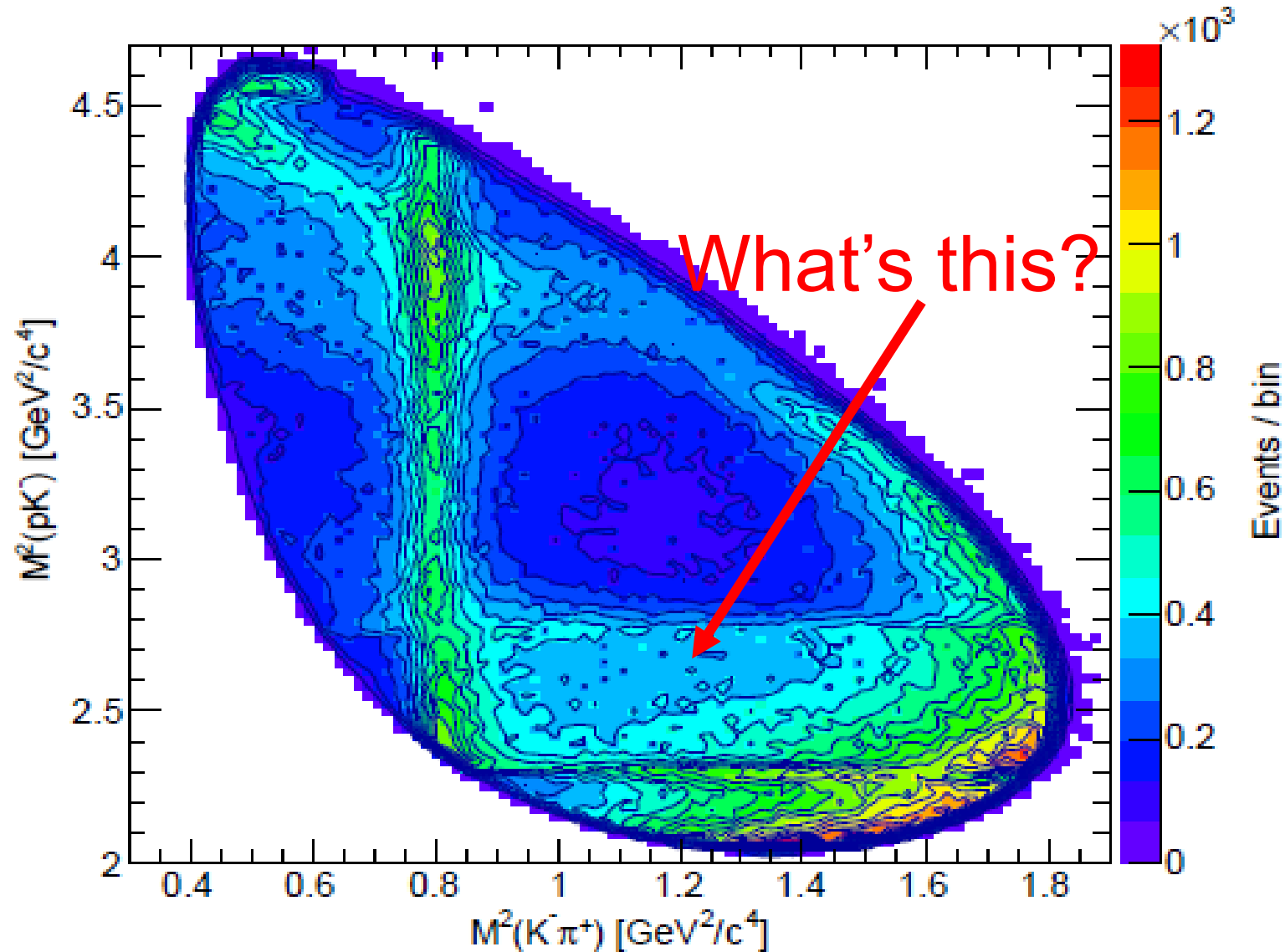


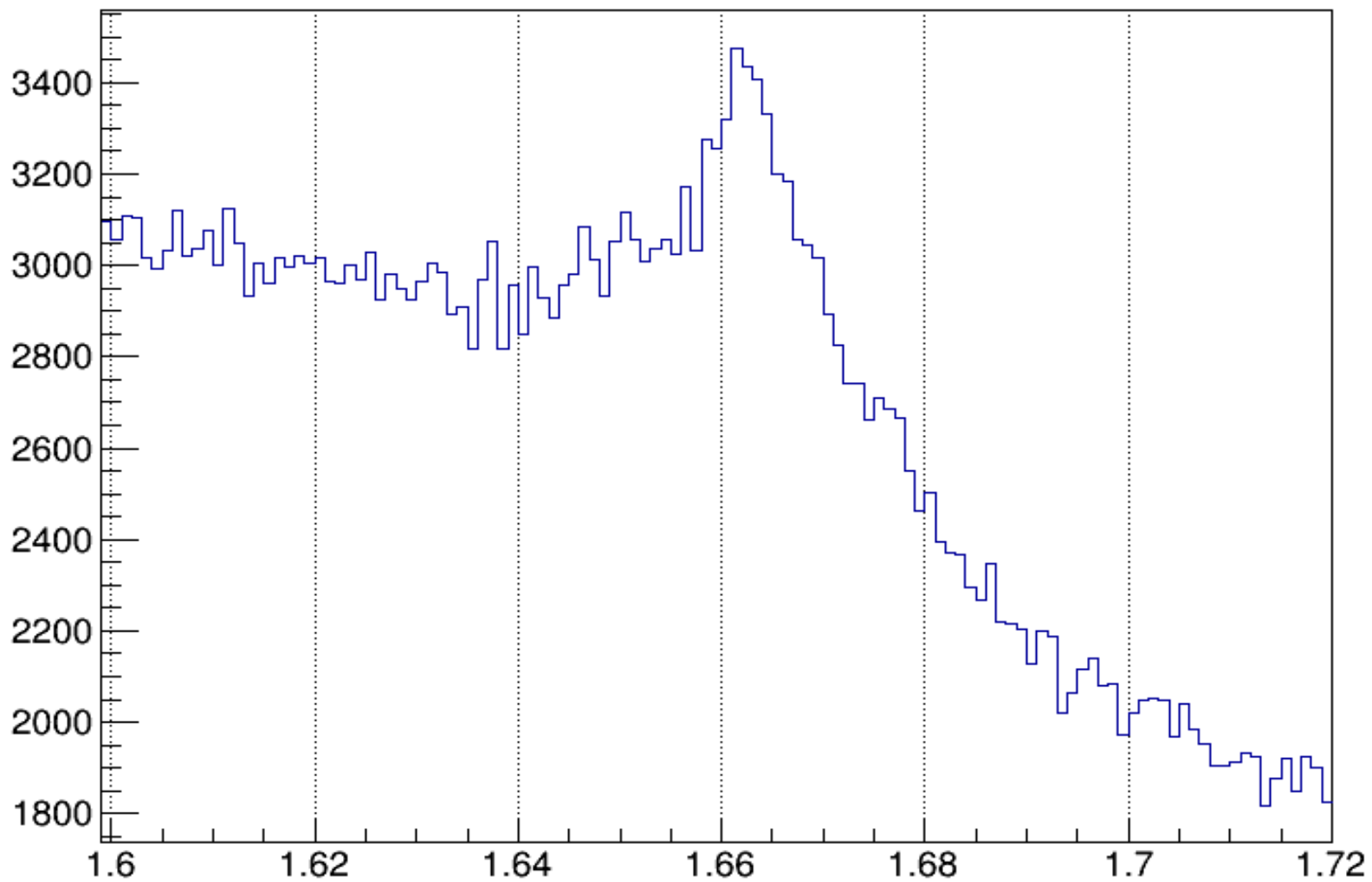
FIG. 7: Energy distributions $d\Gamma/dE_{\pi^+}$ in the rest frame of Λ_c^+ as a function of E_{π^+} .

Dalitz Plot: $\Lambda_c^+ \rightarrow pK^-\pi^+$ decay



Belle Collaboration, PRL 117, 011801 (2016).

■ 1D projection -- $M(pK^-)$



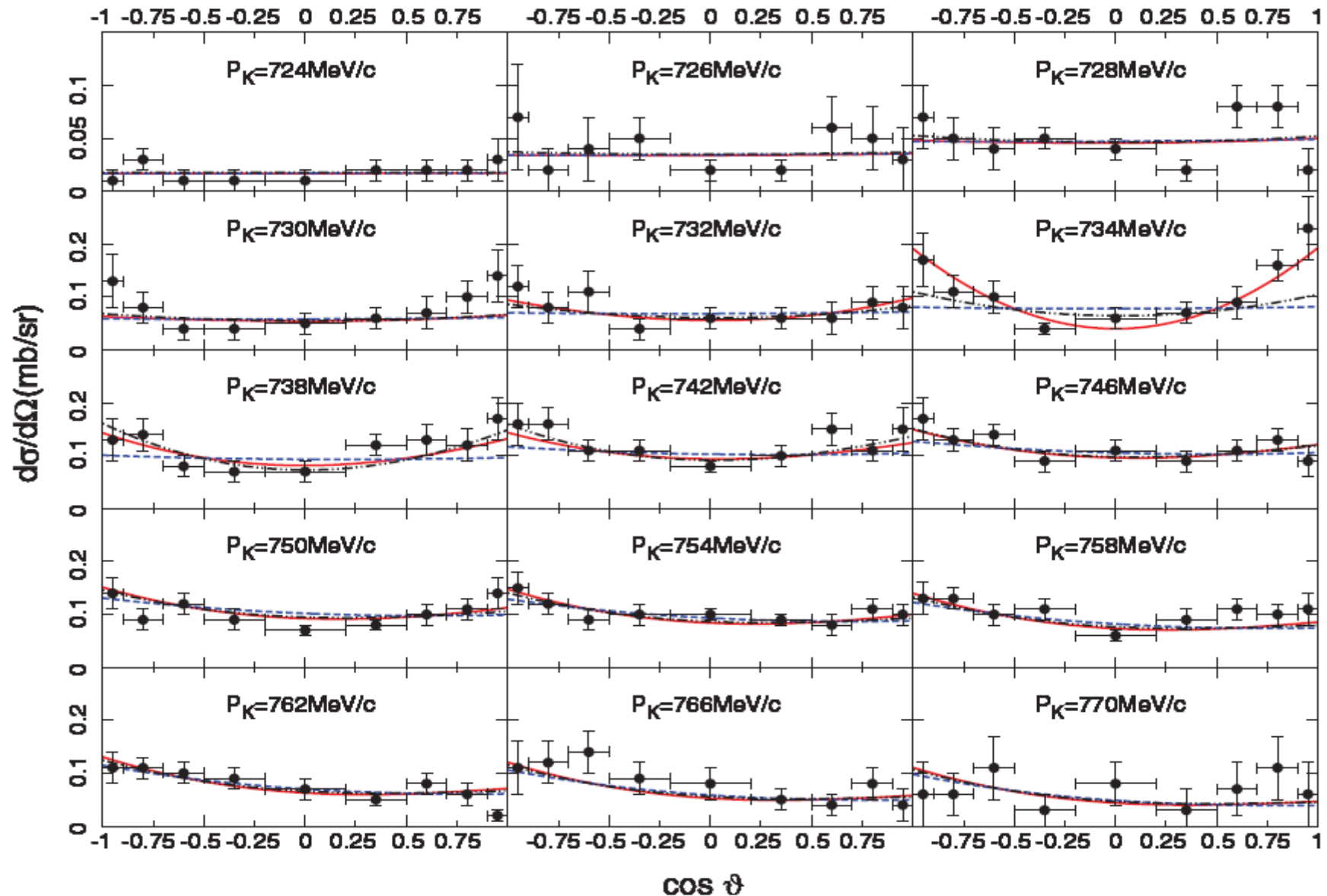
What's that?

- The peak position is ~ 1663 MeV, near the $\eta\Lambda$ threshold (1663.5 MeV)
- Width is ~ 10 MeV, significantly narrower than Λ^* resonances in this region
 - $\Lambda(1670)$: 25-50 MeV
 - $\Lambda(1690)$: ~ 60 MeV
- Maybe due to interference with BG amplitude?

A new narrow Λ^* resonance

- Predictions before: a new narrow Λ^* resonance at this energy with $J=3/2$
 - Kamano et al. [PRC 90, 065204 (2014); PRC 92, 025205 (2015)]
 $J^P=3/2^+$ (P_{03}), $M=1671^{+2}_{-8}$ MeV, $\Gamma=10^{+22}_{-4}$ MeV
 - 刘伯超 & 谢聚军 [PRC 85, 038201 (2012); PRC 86, 055202 (2012)]
 $J^P=3/2^-$ (D_{03}), $M=1668.5 \pm 0.5$ MeV, $\Gamma=1.5 \pm 0.5$ MeV
Couples strongly to $\Lambda\eta$ channel

Angular distributions for $K^-p \rightarrow \eta\Lambda$ reaction



Other contributions

$$\Lambda^*(1670), J^P = 1/2^-, (M, \Gamma) = (1670, 35)\text{MeV} \rightarrow \Lambda\eta \sim 10 - 25\%$$

$$\Lambda^*(1690), J^P = 3/2^-, (M, \Gamma) = (1690, 60)\text{MeV} \rightarrow \Lambda\eta \quad ??$$

$$\Sigma^* \text{ resonances}, \rightarrow \pi^+ \Lambda$$

.....

$$a_0(980), \rightarrow \pi^+ \eta$$

$N^*(1535)$

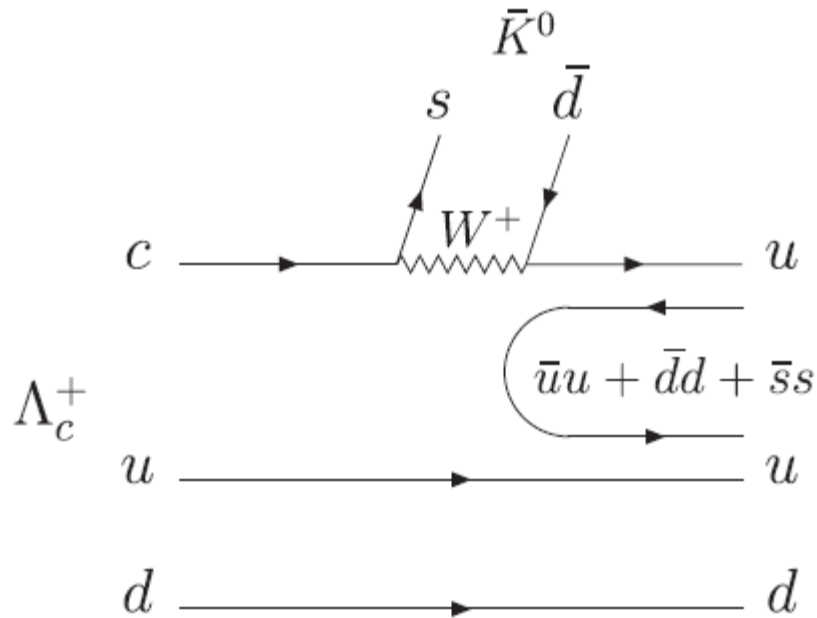
- Couples strongly to $N\eta$ channel
- Pole position $z_R = [(1490 \sim 1530) - i(45 \sim 125)]\text{MeV}$

$$(M_R, \Gamma_R) = (\simeq 1510, \simeq 170)\text{MeV}$$

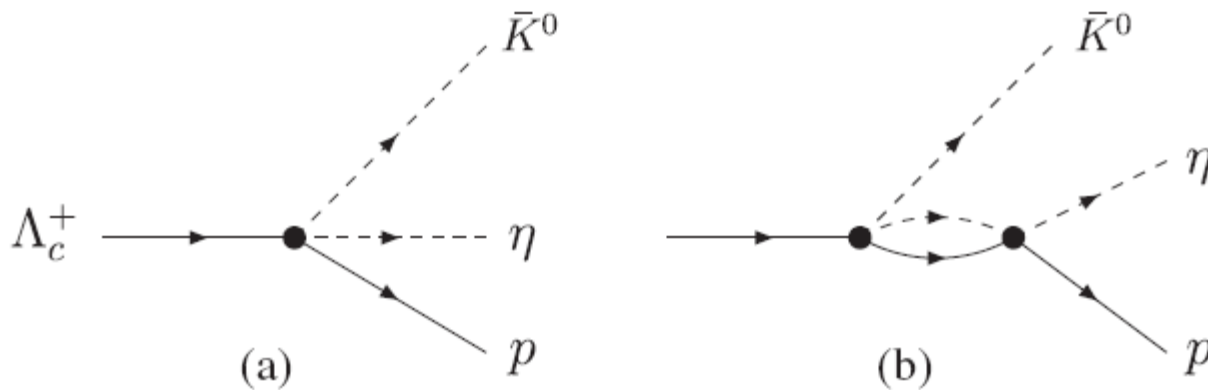
- Breit-Wigner parameterization

$$(M_R, \Gamma_R) = (1525 \sim 1545, 125 \sim 175)\text{MeV} = (\simeq 1535, \simeq 150)\text{MeV}$$

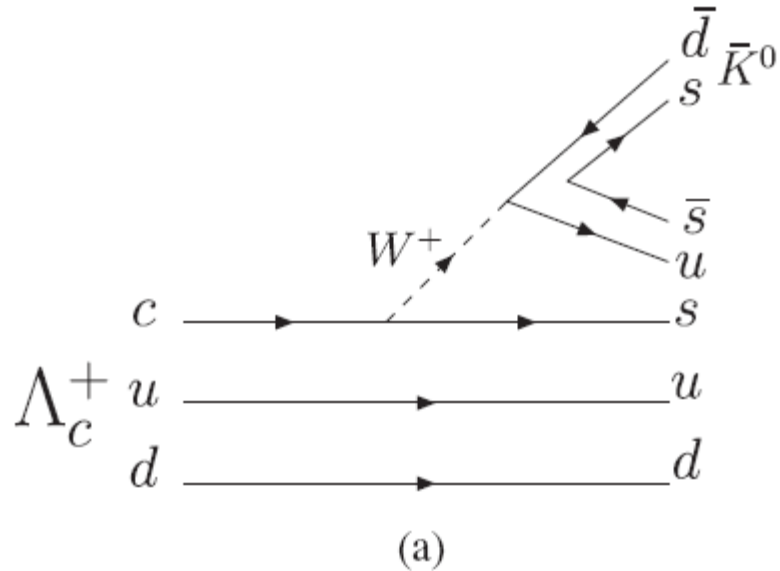
The $N^*(1535)$ as a dynamically generated state



$$T^{MB} = V_P \left(\frac{\sqrt{3}}{3} + \frac{\sqrt{3}}{3} G_{\eta p}(M_{\eta p}) t_{\eta p \rightarrow \eta p}(M_{\eta p}) \right. \\
+ \frac{\sqrt{2}}{2} G_{\pi^0 p}(M_{\eta p}) t_{\pi^0 p \rightarrow \eta p}(M_{\eta p}) \\
+ G_{\pi^+ n}(M_{\eta p}) t_{\pi^+ n \rightarrow \eta p}(M_{\eta p}) \\
\left. - \frac{\sqrt{6}}{3} G_{K^+ \Lambda}(M_{\eta p}) t_{K^+ \Lambda \rightarrow \eta p}(M_{\eta p}) \right),$$



Effective Lagrangian approach and the $N^*(1535)$ resonance as a Breit-Wigner resonance

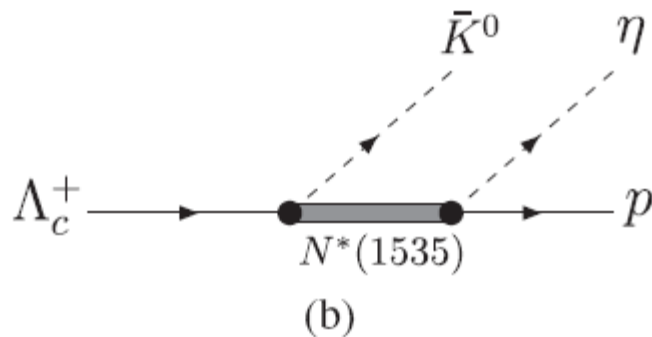


$$T^{N^*} = ig_{N^*N\eta} \bar{u}(p_3, s_p) G_{N^*}(q) (A + B\gamma_5) u(p, s_{\Lambda_c^+}),$$

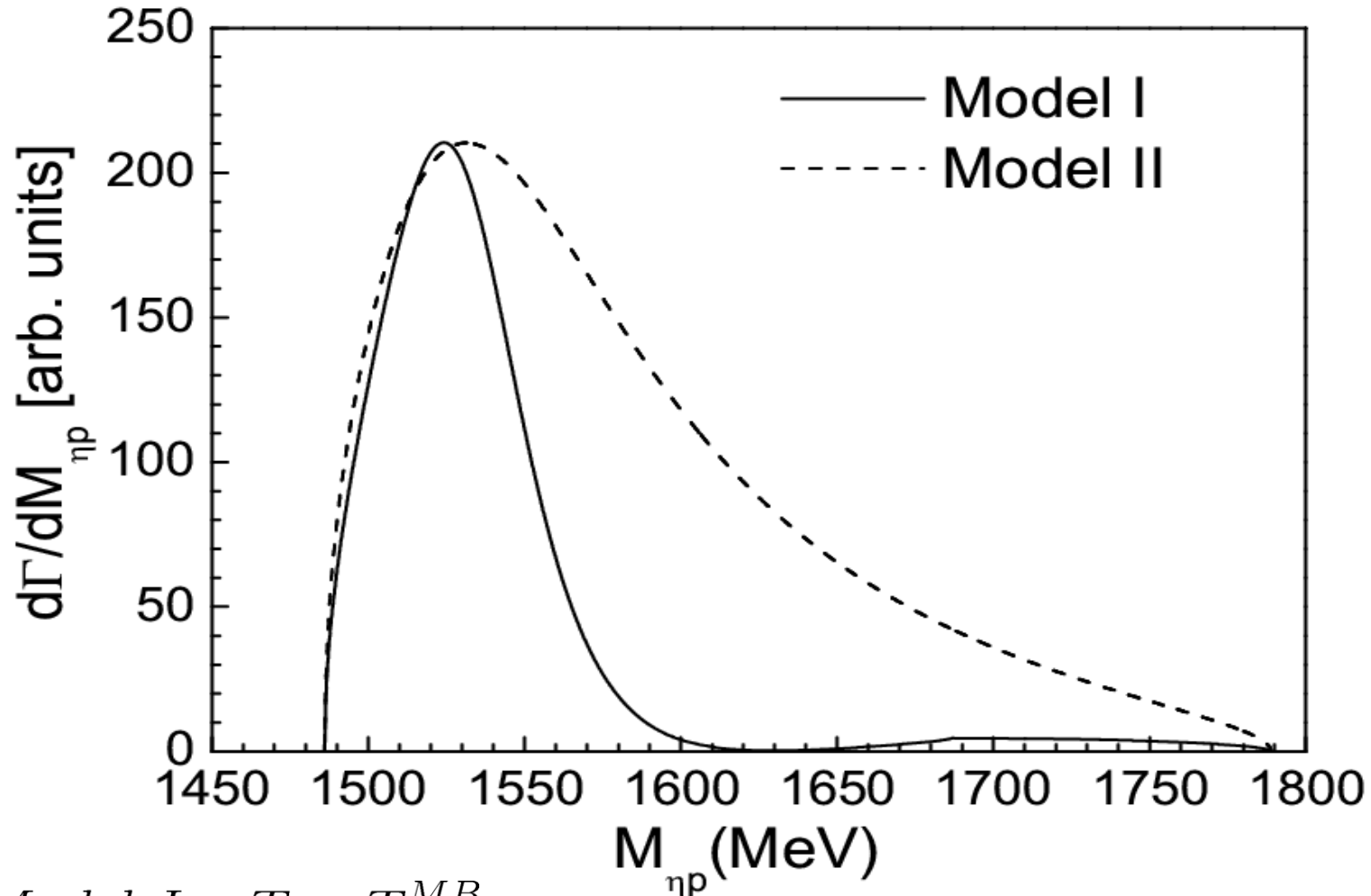
$$G_{N^*}(q) = i \frac{\not{q} + M_{N^*}}{q^2 - M_{N^*}^2 + iM_{N^*}\Gamma_{N^*}(q^2)},$$

$$\Gamma_{N^*}(q^2) = \Gamma_{N^* \rightarrow \pi N}(q^2) + \Gamma_{N^* \rightarrow \eta N}(q^2) + \Gamma_0,$$

$$\Gamma_0 = 19.5 \text{ MeV} \quad \text{for} \quad \Gamma_{N^*}(\sqrt{q^2} = 1535 \text{ MeV}) = 150 \text{ MeV}.$$



Invariant ηp mass distributions



Model I : $T = T^{MB}$

Model II : $T = T^{N^*}$, $M_{N^*} = 1535$ MeV, $\Gamma_{N^*} = \Gamma_{N^*}(q^2)$

Other contributions

$N^*(1650)$

Σ^* resonances

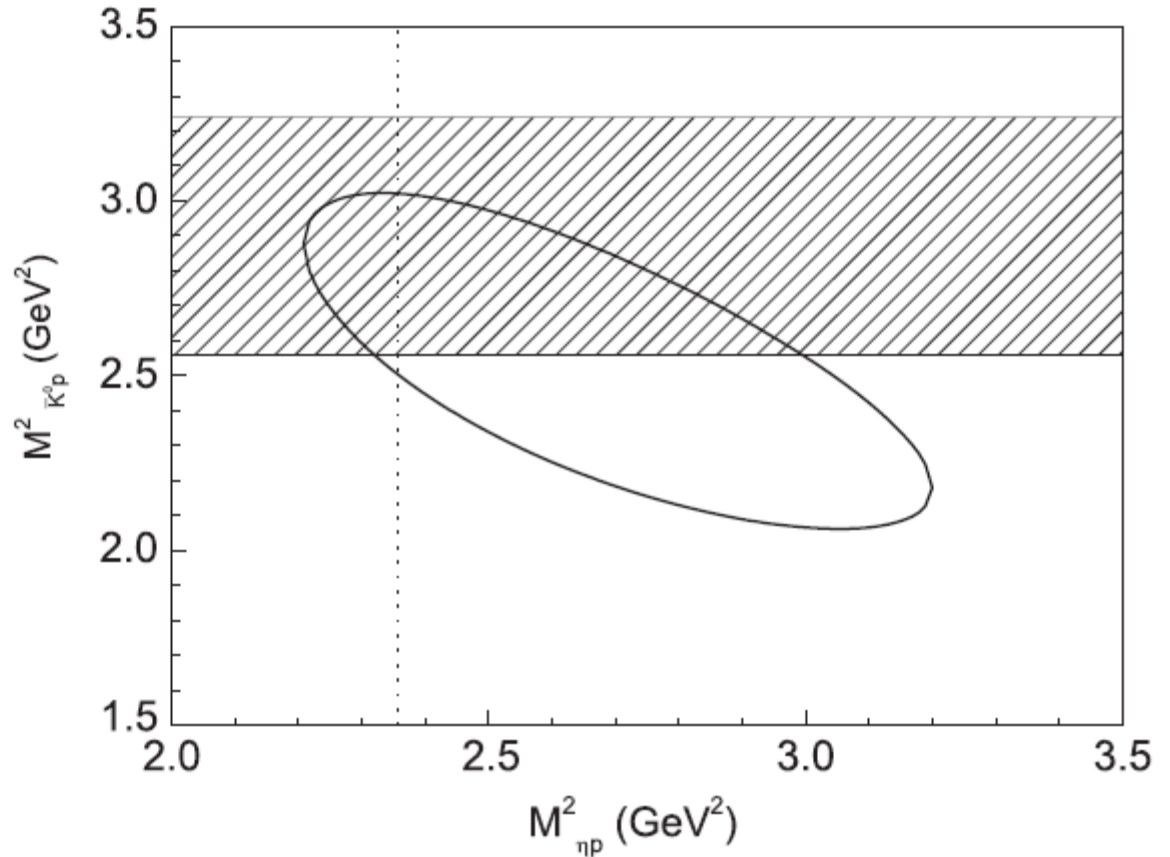


FIG. 7. Dalitz plot for $M_{\eta p}^2$ and $M_{\bar{K}^0 p}^2$ in the $\Lambda_c^+ \rightarrow \bar{K}^0 \eta p$ decay. The $N^*(1535)$ energy is shown by the vertical dotted line, and the horizontal band represents the masses of Σ^* states from 1600 to 1800 MeV.

Possible ϕp state in $\Lambda_c^+ \rightarrow \pi^0 p \phi$ decay

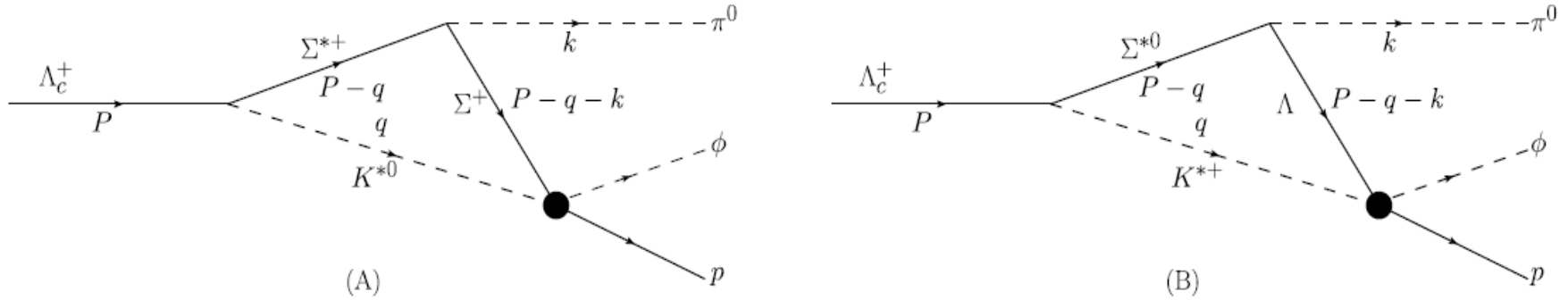


Fig. 1. Triangle diagrams for the $\Lambda_c^+ \rightarrow \pi^0 p \phi$ decay. (A): Σ^+ -exchange. (B): Λ -exchange. The definitions of the kinematical variables (P, q, k) are also shown.

$$\begin{aligned}
 t &= \frac{g_{\Lambda_c \Sigma^* K^*} g_{\vec{\epsilon}_\phi \cdot \vec{k}}}{m_\pi} \sum_{i=\Sigma, \Lambda} c_i \int \frac{d^4 q}{(2\pi)^4} \\
 &\times \frac{i 2m_{\Sigma^*}}{(P-q)^2 - m_{\Sigma^*}^2 + i m_{\Sigma^*} \Gamma_{\Sigma^*}} \frac{i}{q^2 - m_{K^*}^2 + i m_{K^*} \Gamma_{K^*}} \\
 &\times \frac{i 2m_i}{(P-q-k)^2 - m_i^2 + i\epsilon}, \tag{4}
 \end{aligned}$$

where we have defined $c_\Sigma = \frac{\sqrt{6}}{3} t_{K^{*0} \Sigma^+ \rightarrow \phi p}$ and $c_\Lambda = -t_{K^{*+} \Lambda \rightarrow \phi p}$,

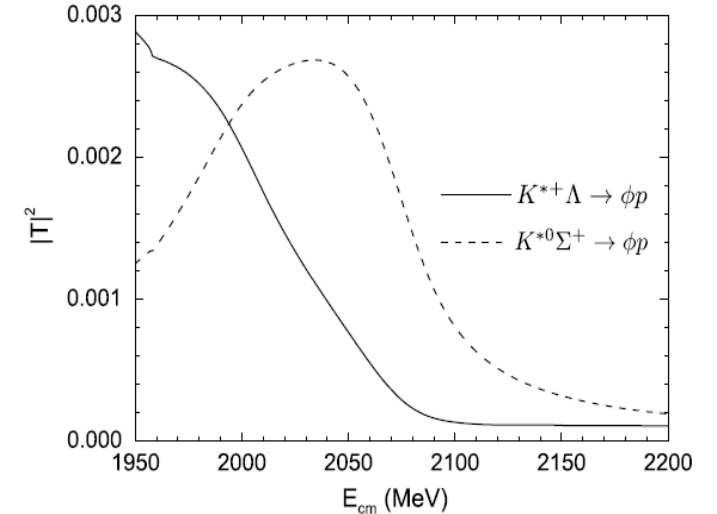
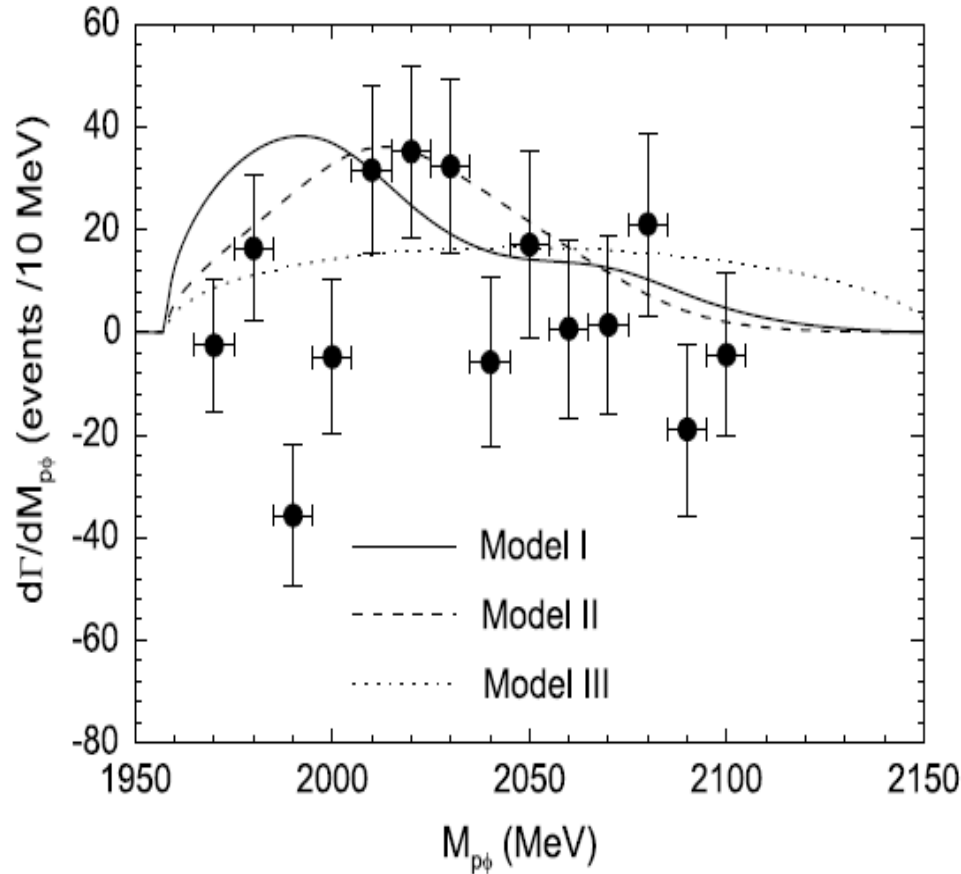


Fig. 3. The squared norm of the T -matrix elements for $K^{*+} \Lambda \rightarrow \phi p$ and $K^{*0} \Sigma^+ \rightarrow \phi p$ as a function of the meson-baryon invariant mass E_{cm} in the model of Ref. [72].



$$t_{\text{I}} = t,$$

$$t_{\text{II}} = t, \text{ but with } t_{K^{*+}\Lambda \rightarrow \phi p} = \frac{\sqrt{6}}{2} t_{K^{*0}\Sigma^+ \rightarrow \phi p}$$

$$= c_1 \frac{\sqrt{6}}{2} \frac{E_{K^*} + E_\phi}{4F_\pi^2},$$

$$t_{\text{III}} = c_2,$$

Fig. 2. Invariant mass distribution of the $\Lambda_c^+ \rightarrow \pi^0 p \phi$ decay. The experimental data are taken from Ref. [47].

[47] B. Pal, et al., Belle Collaboration, arXiv:1707.00089 [hep-ex].

Summary

The $\Lambda_c^+ \rightarrow \eta\pi^+\Lambda$ decay can be used to study the Σ^ and Λ^* resonances*

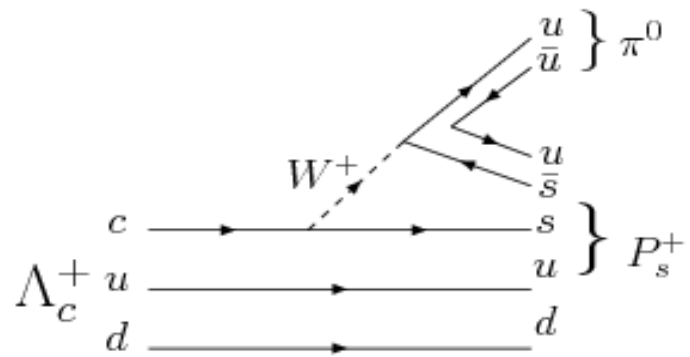
The $\Lambda_c^+ \rightarrow \bar{K}^0\eta p$ decay can be used to study the $N^(1535)$ resonance*

Possible ϕp state, P_s , in the $\Lambda_c^+ \rightarrow \pi^0\phi p$ decay

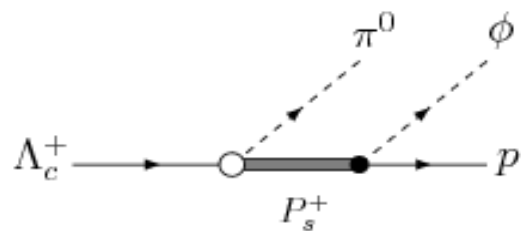
But,

We need more efforts, both on theoretical and experimental sides.

Thank you very much for your attention!



(a)



(b)