# Top quark pair production 

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第十五届粒子物理，核物理和宇宙学交叉学科前沿问题研讨会

## Basic facts about the top

Large mass $m_{t} \approx 173 \mathrm{GeV}$

$=$ Basic facts about the top

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Fermion mass origin
Strong Yukawa coupling $\quad y_{t} \sim 1$
Hierarchy problem
Vacuum stability

## Basic facts about the top

## Large mass $m_{t} \approx 173 \mathrm{GeV}$



Fermion mass origin
Strong Yukawa coupling $\quad y_{t} \sim 1$
Hierarchy problem
Vacuum stability
Short lifetime $\tau \sim 5 \times 10^{-25} \mathrm{~s}$
Decays before hadronization: pQCD dominates!

## Top quark pair production <br> 

A standard candle for the LHC and future colliders

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## Top quark pair production <br> 

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* Test of the SM at the energy frontier
* Possible signals of new physics
* Major background to many searches
* Precise theoretical and experimental results have already enabled us to gain useful information!


## Gluon PDF

Top quark pair production can provide information about the gluon parton distribution functions


Czakon, Mangano, Mitov, Rojo: 1303.7215
Note: only used 7 and 8 TeV data!

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Note: only used 7 and 8 TeV data!
Ongoing: CTEQ analysis with 8 and 13 TeV data

## Deviation?



## Deviation?

## What's going on here?

Top quark mass?
Threshold effect?


## Deviation?



ATLAS-CONF-2018-027


CMS PAS TOP-17-014

In this talk, I'm going to introduce the state-of-the-art QCD prediction for top quark pair production...

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...and some ongoing developments

## NNLO QCD for top pair

## Total cross section

Baernreuther, Czakon, Mitov: 1204.5201; Czakon, Fiedler, Mitov: 1303.6254


## Differential distributions

Czakon, Heymes, Mitov: 1511.00549


## Differential distributions

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Some tension at high energy (boosted kinematics)

## Kinematics

The difficulty for fixed-order calculations: multiplescale process with complicated kinematics!


Many kinematic variables:
top quark mass
$p_{\text {т }}$ of top
$\mathrm{p}_{\text {t }}$ of anti-top
rapidity of top
rapidity of anti-top
Invariant mass $\mathrm{M}_{\mathrm{tt}}$
...

## Kinematics

The difficulty for fixed-order calculations: multiplescale process with complicated kinematics!


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rapidity of top
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Invariant mass $\mathrm{M}_{\mathrm{tt}}$

Which (combination) should be used for the renormalization/factorization scales?

## NNLO with dynamic scale

Czakon, Heymes, Mitov: 1606.03350

## Determine optimal "scale scheme" by minimizing higher order corrections

$$
\begin{aligned}
& \mu_{0} \sim m_{t}, \\
& \mu_{0} \sim m_{T}=\sqrt{m_{t}^{2}+p_{T}^{2}}, \\
& \mu_{0} \sim H_{T}=\sqrt{m_{t}^{2}+p_{T, t}^{2}}+\sqrt{m_{t}^{2}+p_{T, \bar{t}}^{2}},
\end{aligned}
$$

$$
\begin{aligned}
& \mu_{0} \sim E_{T}=\sqrt{\sqrt{m_{t}^{2}+p_{T, t}^{2}} \sqrt{m_{t}^{2}+p_{T, \bar{t}}^{2}}}, \\
& \mu_{0} \sim H_{T, \text { int }}=\sqrt{\left(m_{t} / 2\right)^{2}+p_{T, t}^{2}}+\sqrt{\left(m_{t} / 2\right)^{2}+p_{T, \bar{t}}^{2}}, \\
& \mu_{0} \sim m_{t \bar{t}},
\end{aligned}
$$

## NNLO with dynamic scale

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Vastly different behaviors with different scheme choices (especially in the boosted region)

## Our philosophy

We should study different regions of phase space separately, and combine them to have a good description for all regions!


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Threshold region fixed-order+soft+Coulomb (ongoing)


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Intermediate region
fixed-order+soft
Ahrens, Ferroglia, Neubert,
Pecjak, LLY: 1003.5827

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We should study different regions of phase space separately, and combine them to have a good description for all regions!

Threshold region
fixed-order+soft+Coulomb
(ongoing)

Intermediate region fixed-order+soft

Ahrens, Ferroglia, Neubert, Pecjak, LLY: 1003.5827


Boosted region
fixed-order+soft+quasi-collinear
Pecjak, Scott, Wang, LLY: 1601.07020
Czakon, Ferroglia, Heymes, Mitov, Pecjak,
Scott, Wang, LLY: 1803.07623

## Boosted top quarks

Sensitive to new physics, interesting in its own right!


Actively being probed by LHC experiments

## Producing boosted tops

Hard extra emissions
suppressed

soft gluons


$$
\ln \frac{\hat{s}-M_{t \bar{t}}^{2}}{M_{t \bar{t}}^{2}}
$$

## Producing boosted tops

Hard extra emissions suppressed

soft gluons

$$
\ln \frac{\hat{s}-M_{t \bar{t}}^{2}}{M_{t \bar{t}}^{2}}
$$

Top quark nearly massless

quasi-collinear gluons

$$
\underset{\ln \frac{m_{2}^{2}}{M_{t t_{2}^{2}}}}{\boldsymbol{\nabla}}
$$

## Producing boosted tops

Hard extra emissions suppressed

## Top quark nearly massless


soft gluons

quasi-collinear gluons


$$
\ln \frac{\hat{s}-M_{t \bar{t}}^{2}}{M_{t \bar{t}}^{2}}
$$

Need to resum both!

$$
\ln \frac{m_{t}^{2}}{M_{t \bar{t}}^{2}}
$$

## Soft gluon resummation

Hard function


Kidonakis, Sterman: hep-ph/9705234
Ahrens, Ferroglia, Neubert, Pecjak, LLY: 1003.5827

Evolving from the scale of hard scatterings

## Soft gluon resummation

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Evolving from the scale of hard scatterings
to the scale of soft interactions


## Soft gluon resummation

Hard function


Evolving from the scale of hard scatterings

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Ahrens, Ferroglia, Neubert,
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## IR anomalous dimension

$$
\begin{align*}
\boldsymbol{\Gamma}= & \sum_{(i, j)} \frac{\boldsymbol{T}_{i} \cdot \boldsymbol{T}_{j}}{2} \gamma_{\mathrm{cusp}}\left(\alpha_{s}\right) \ln \frac{\mu^{2}}{-s_{i j}}+\sum_{i} \gamma^{i}\left(\alpha_{s}\right) \\
& -\sum_{(I, J)} \frac{\boldsymbol{T}_{I} \cdot \boldsymbol{T}_{J}}{2} \gamma_{\mathrm{cusp}}\left(\beta_{I J}, \alpha_{s}\right)+\sum_{I} \gamma^{I}\left(\alpha_{s}\right) \\
& +\sum_{I, j} \boldsymbol{T}_{I} \cdot \boldsymbol{T}_{j} \gamma_{\mathrm{cusp}}\left(\alpha_{s}\right) \ln \frac{m_{I} \mu}{-s_{I j}} \\
& +\sum_{(I, J, K)} i f^{a b c} \boldsymbol{T}_{I}^{a} \boldsymbol{T}_{J}^{b} \boldsymbol{T}_{K}^{c} F_{1}\left(\beta_{I J}, \beta_{J K}, \beta_{K I}\right) \\
& +\sum_{(I, J)} \sum_{k} i f^{a b c} \boldsymbol{T}_{I}^{a} \boldsymbol{T}_{J}^{b} \boldsymbol{T}_{k}^{c} f_{2}\left(\beta_{I J}, \ln \frac{-\sigma_{J k} v_{J} \cdot p_{k}}{-\sigma_{I k} v_{I} \cdot p_{k}}\right)
\end{align*}
$$

Becher, Neubert: 0904.1021
Ferroglia, Neubert, Pecjak, LLY: 0907.4791; 0908.3676

$$
\begin{align*}
& F_{1}\left(\beta_{12}, \beta_{23}, \beta_{31}\right)=\frac{\alpha_{s}^{2}}{12 \pi^{2}} \sum_{i, j, k} \epsilon_{i j k} g\left(\beta_{i j}\right) r\left(\beta_{k i}\right) \\
& r(\beta)=\beta \operatorname{coth} \beta \\
& g(\beta)= \\
& \quad \operatorname{coth} \beta\left[\beta^{2}+2 \beta \ln \left(1-e^{-2 \beta}\right)-\operatorname{Li}_{2}\left(e^{-2 \beta}\right)+\frac{\pi^{2}}{6}\right]  \tag{5}\\
& \quad-\beta^{2}-\frac{\pi^{2}}{6}
\end{align*}
$$

3-parton correlations

## The soft function



Known at NLO<br>Ahrens, Ferroglia, Neubert, Pecjak, LLY: 1003.5827<br>Known at NNLO in the massless limit (except an off-diagonal 3-parton piece)<br>Ferroglia, Pecjak, LLY: 1207.4798

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Recent calculation at NNLO with massive tops
Wang, Xu, LLY, Zhu: 1804.05218

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## Recent calculation at NNLO with massive tops

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## NNLO diagrams

Wang, Xu, LLY, Zhu: 1804.05218




c)






a.





## Solving integrals

Wang, Xu, LLY, Zhu: 1804.05218
~60 master integrals

Differential equations

$$
\partial_{\beta} \vec{f}(\epsilon, \beta, \cos \theta)=\epsilon\left(\frac{A}{\beta-1}+\frac{B}{\beta}+\frac{C}{\beta+1}+\frac{D}{\beta-1 / \cos \theta}+\frac{E}{\beta+1 / \cos \theta}\right) \vec{f}(\epsilon, \beta, \cos \theta)
$$



Solution in terms of generalized polylogarithms
Difficult part: boundary conditions

## The boundary conditions

We choose the boundary to be $\beta \equiv \sqrt{1-\frac{4 m_{t}^{2}}{M_{t \bar{t}}}} \rightarrow 0$


Some virtual-real integrals develop
Coulomb/Glauber-type singularities in this limit

Carefully extract the asymptotic behavior, e.g.

$$
g_{6}^{(4)}(\epsilon, \beta \rightarrow 0, y) \approx \frac{\left(e^{-2 i \pi \epsilon}-1\right) \beta^{2 \epsilon} \Gamma(1-2 \epsilon) \Gamma(1+\epsilon)}{4^{1-2 \epsilon} \Gamma(1-\epsilon)}
$$

## A piece of final result

Wang, Xu, LLY, Zhu: 1804.05218

$$
\begin{aligned}
\tilde{s}_{22}^{q \bar{q},(2)} & \left.(0, \beta, y)\right|_{T_{F} N_{l}}=\frac{16\left(7 \beta^{2}-126 \beta+127\right)}{243 \beta} G_{1}+\frac{8\left(5 \beta^{2}+90 \beta+53\right)}{81 \beta}\left(G_{-1,-1}-G_{-1,1}-2 G_{0,-1}\right) \\
& -\frac{16\left(7 \beta^{2}+126 \beta+127\right)}{243 \beta} G_{-1}+\frac{8\left(5 \beta^{2}-90 \beta+53\right)}{81 \beta}\left(G_{1,-1}-G_{1,1}+2 G_{0,1}\right) \\
& +\frac{8\left(\beta^{2}+18 \beta+1\right)}{27 \beta}\left(-G_{-1,-1,-1}+G_{-1,-1,1}+2 G_{-1,0,-1}-2 G_{-1,0,1}-G_{-1,1,-1}+G_{-1,1,1}\right. \\
& \left.+2 G_{0,-1,-1}-2 G_{0,-1,1}-4 G_{0,0,-1}\right)+\frac{8\left(\beta^{2}-18 \beta+1\right)}{27 \beta}\left(4 G_{0,0,1}+2 G_{0,1,-1}-2 G_{0,1,1}\right. \\
& \left.-G_{1,-1,-1}+G_{1,-1,1}+2 G_{1,0,-1}-2 G_{1,0,1}-G_{1,1,-1}+G_{1,1,1}\right) \\
& +\frac{32}{243}\left[28 G_{-1 / y}+98 G_{1 / y}+30\left(2 G_{0,-1 / y}+G_{-1 / y,-1}+G_{-1 / y, 1}-2 G_{-1 / y,-1 / y}\right)\right. \\
& +105\left(2 G_{0,1 / y}+G_{1 / y,-1}+G_{1 / y, 1}-2 G_{1 / y, 1 / y}\right)+18\left(4 G_{0,0,-1 / y}+2 G_{0,-1 / y,-1}+2 G_{0,-1 / y, 1}\right. \\
& -4 G_{0,-1 / y,-1 / y}-G_{-1 / y,-1,-1}+G_{-1 / y,-1,1}+2 G_{-1 / y, 0,-1}+2 G_{-1 / y, 0,1}-4 G_{-1 / y, 0,-1 / y} \\
& \left.+G_{-1 / y, 1,-1}-G_{-1 / y, 1,1}-2 G_{-1 / y,-1 / y,-1}-2 G_{-1 / y,-1 / y, 1}+4 G_{-1 / y,-1 / y,-1 / y}\right) \\
& +63\left(4 G_{0,0,1 / y}+2 G_{0,1 / y,-1}+2 G_{0,1 / y, 1}-4 G_{0,1 / y, 1 / y}-G_{1 / y,-1,-1}+G_{1 / y,-1,1}+2 G_{1 / y, 0,-1}\right. \\
& \left.+2 G_{1 / y, 0,1}-4 G_{1 / y, 0,1 / y}+G_{1 / y, 1,-1}-G_{1 / y, 1,1}-2 G_{1 / y, 1 / y,-1}-2 G_{1 / y, 1 / y, 1}+4 G_{1 / y, 1 / y, 1 / y}\right) \\
& \left.-\frac{332}{3}-\frac{5 \pi^{2}}{2}+6 \zeta_{3}\right]
\end{aligned}
$$

It is remarkable that all the results can be written analytically in terms of multiple polylogarithms


Allows fast numerics!

## Threshold limit

It is interesting to check the threshold limit where the top quarks are produced at rest


# Color singlet: same as Drell-Yan and Higgs production Belitsky: hep-ph/9808389 

Color octet Czakon, Fiedler: 1311.2541

Note: singlet-octet mixing terms do NOT vanish in the threshold limit!

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## Boosted limit

In the limit where the top quarks are highly boosted


Factorization Ferroglia, Pecjak, LLY: 1205.3662
$\boldsymbol{S}_{\text {massive }}\left(s, t, m_{t}, N\right) \rightarrow \boldsymbol{S}_{\text {massless }}(s, t, N) S_{D}^{2}\left(m_{t} / N\right)$

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Also obtain the missing
3-parton piece for free

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Ferroglia, Pecjak, LLY: 1207.4798
Also obtain the missing
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Allows to extract the soft fragmentation function

# Soft and small-mass factorization 

Ferroglia, Pecjak, LLY: 1205.3662
In Mellin space: $\quad Q \sim \sqrt{s}, \sqrt{-t} \gg Q / N \gg m_{t} \gg m_{t} / N$

$$
\hat{\sigma}\left(N, \mu_{f}\right) \sim \operatorname{Tr}\left[\boldsymbol{H}\left(L_{h}, \mu_{f}\right) \boldsymbol{S}\left(L_{s}, \mu_{f}\right)\right] C_{D}^{2}\left(L_{c}, \mu_{f}\right) S_{D}^{2}\left(L_{s c}, \mu_{f}\right)
$$

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& \ln \frac{Q^{2}}{\mu_{f}^{2}}
\end{aligned}
$$

hard log

# Soft and small-mass factorization 

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In Mellin space: $\quad Q \sim \sqrt{s}, \sqrt{-t} \gg Q / N \gg m_{t} \gg m_{t} / N$ $\hat{\sigma}\left(N, \mu_{f}\right) \sim \operatorname{Tr}\left[\boldsymbol{H}\left(L_{h}, \mu_{f}\right) \boldsymbol{S}\left(L_{s}, \mu_{f}\right)\right] C_{D}^{2}\left(L_{c}, \mu_{f}\right) S_{D}^{2}\left(L_{s c}, \mu_{f}\right)$
$\quad \ln \frac{Q^{2}}{\mu_{f}^{2}} \quad \ln \frac{Q^{2}}{\bar{N}^{2} \mu_{f}^{2}}$
hard log
soft log

# Soft and small-mass factorization 

Ferroglia, Pecjak, LLY: 1205.3662
In Mellin space: $\quad Q \sim \sqrt{s}, \sqrt{-t} \gg Q / N \gg m_{t} \gg m_{t} / N$

$$
\begin{array}{cc}
\hat{\sigma}\left(N, \mu_{f}\right) \sim \operatorname{Tr}\left[\boldsymbol{H}\left(L_{h}, \mu_{f}\right)\right. & \left.\boldsymbol{S}\left(L_{s}, \mu_{f}\right)\right] C_{D}^{2}\left(L_{c}, \mu_{f}\right) S_{D}^{2}\left(L_{s c}, \mu_{f}\right) \\
\ln \frac{Q^{2}}{\mu_{f}^{2}} & \ln \frac{Q^{2}}{\bar{N}^{2} \mu_{f}^{2}}
\end{array} \quad \begin{aligned}
& \ln \frac{m_{t}^{2}}{\mu_{f}^{2}} \\
& \text { hard log } \\
& \text { soft log }
\end{aligned} \begin{aligned}
& \text { collinear log } \\
& \text { (small-mass) }
\end{aligned}
$$

# Soft and small-mass factorization 

Ferroglia, Pecjak, LLY: 1205.3662
In Mellin space: $\quad Q \sim \sqrt{s}, \sqrt{-t} \gg Q / N \gg m_{t} \gg m_{t} / N$ $\begin{array}{cc}\hat{\sigma}\left(N, \mu_{f}\right) \sim \operatorname{Tr}\left[\boldsymbol{H}\left(L_{h}, \mu_{f}\right) \boldsymbol{S}\left(L_{s}, \mu_{f}\right)\right] C_{D}^{2}\left(L_{c}, \mu_{f}\right) S_{D}^{2}\left(L_{s c}, \mu_{f}\right) \\ \ln \frac{Q^{2}}{\mu_{f}^{2}} & \ln \frac{Q^{2}}{\bar{N}^{2} \mu_{f}^{2}} \\ \text { hard log } & \ln \frac{m_{t}^{2}}{\mu_{f}^{2}}\end{array} \begin{gathered}\ln \frac{m_{t}^{2}}{\bar{N}^{2} \mu_{f}^{2}} \\ \text { soft log } \\ \begin{array}{c}\text { collinear log } \\ \text { (small-mass) }\end{array}\end{gathered}$

## Soft and small-mass resummation

Massless hard function

Massless soft function
$\boldsymbol{S}\left(L_{s}, \mu_{s} \sim Q / \bar{N}\right)$


All ingredients known at NNLO (for NNLL' resummation)

## NLO+NNLL'

Pecjak, Scott, Wang, LLY: 1601.07020


## NNLO+NNLL'

A joint effort of the NNLO group Czakon, Ferroglia, Heymes, Mitov, Pecjak, and the resummation group Scott, Wang, LLY: 1803.07623

$$
d \sigma^{(\mathrm{N}) \mathrm{NLO}+\mathrm{NNLL}^{\prime}}=d \sigma^{\mathrm{NNLL}_{b+m}^{\prime}}+\left(d \sigma^{(\mathrm{N}) \mathrm{NLO}}-\left.d \sigma^{\mathrm{NNLL}_{b+m}^{\prime}}\right|_{\substack{(\mathrm{N}) \mathrm{NLO} \\ \text { expansion }}}\right)
$$

## NNLO+NNLL'

A joint effort of the NNLO group Czakon, Ferroglia, Heymes, Mitov, Pecjak, and the resummation group Scott, Wang, LLY: 1803.07623

$$
\begin{aligned}
& d \sigma^{(\mathrm{N}) \mathrm{NLO}+\mathrm{NNLLL}^{\prime}}=d \sigma^{\mathrm{NNLL}_{b+m}^{\prime}}+\left(d \sigma^{(\mathrm{N}) \mathrm{NLO}}-\left.d \sigma^{\mathrm{NNLL}_{b+m}^{\prime}}\right|_{\text {expansion }} ^{(\mathrm{N}) \mathrm{NLO}}\right) \\
& d \sigma^{\mathrm{NNLL}_{b}^{\prime}}+\left(d \sigma^{\mathrm{NNLL}_{m}}-\left.d \sigma^{\mathrm{NNLL}_{m}}\right|_{m_{t} \rightarrow 0}\right) \\
& \text { mall mass } \\
& \text { imation } \\
& \text { match to soft } \\
& \text { resummation }
\end{aligned}
$$

soft \& small mass resummation

## NNLO+NNLL'

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## NNLO+NNLL'

A joint effort of the NNLO group Czakon, Ferroglia, Heymes, Mitov, Pecjak, and the resummation group

soft \& small mass resummation
match to soft resummation

Careful to avoid
double-counting!

## NNLO+NNLL'

## State Czakon, Ferroglia, Heymes, Mitov, Pecjak, <br> State-of-the-art QCD prediction Scott, Wang, tur: 1803.07623




Resummation
softens the spectrum

## NNLO+NNLL'

Czakon, Ferroglia, Heymes, Mitov, Pecjak, Scott, Wang, LLY: 1803.07623


Matched result insensitive to scale scheme choices

## Ongoing: rapidity distributions

Pecjak, Scott, Wang, LLY: to appear


Sensitive to gluon PDF at large $x$

# Ongoing: combination with electroweak corrections 

Both EW and resummation effects soften the $\mathrm{p}_{\mathrm{T}}$ spectrum



Czakon et al.: 1705.04105

# Ongoing: combination with electroweak corrections 



NNLO+NNLL'+EW should be better consistent with data!

Stay tuned

## Ongoing: near-threshold



Threshold region sensitive to Coulomb gluons

## Ongoing: near-threshold



Historically, Coulomb gluons have been studied only for the total cross section

Moch, Uwer: 0804.1476
Beneke, Falgari, Klein, Schwinn: 1109.1536

Threshold region sensitive to Coulomb gluons

## Ongoing: near-threshold



Threshold region sensitive to Coulomb gluons

Historically, Coulomb gluons have been studied only for the total cross section

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## Requires new framework to study the $\mathrm{M}_{\mathrm{tt}}$ distribution!

## Summary and outlook

* Top quark production is important
* The most precise QCD calculation: NNLO+NNLL'
* Ongoing:
* Rapidity distributions
* Combination with NLO electroweak corrections
* Near-threshold production


## Thank you!

