



Recent developments in the CTEQ-TEA global analysis

Tie-Jiun Hou

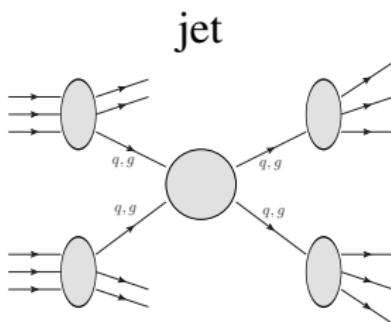
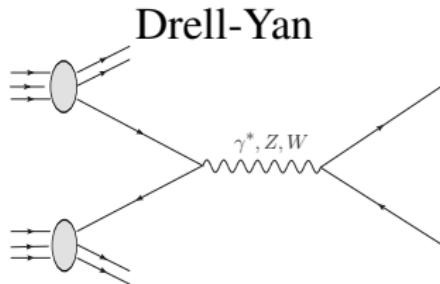
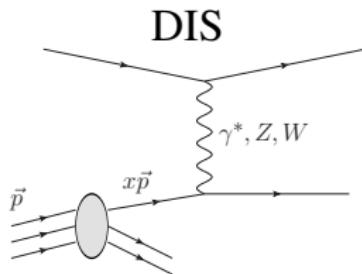
CTP-XJU

August 16, 2018 at SI2018

CTEQ-TEA group

- CTEQ – Tung et al. (TEA)
in memory of Prof. Wu-Ki Tung, who established
CTEQ Collaboration in early 90's.
- Current members:
Sayipjamal Dulat, Tie-Jiun Hou, Ibrahim Sitiwaldi
(Xinjiang U.)
Jun Gao (Shanghai Jiaotong U.)
Marco Guzzi (Kennesaw State U.)
Pavel Nadolsky, Timothy Hobbs, Keping Xie, Boting
Wang (Southern Methodist U.)
Joey Huston, Jon Pumplin, Carl Schmidt, Dan Stump,
Jan Winter, C.-P. Yuan (Michigan State U.)

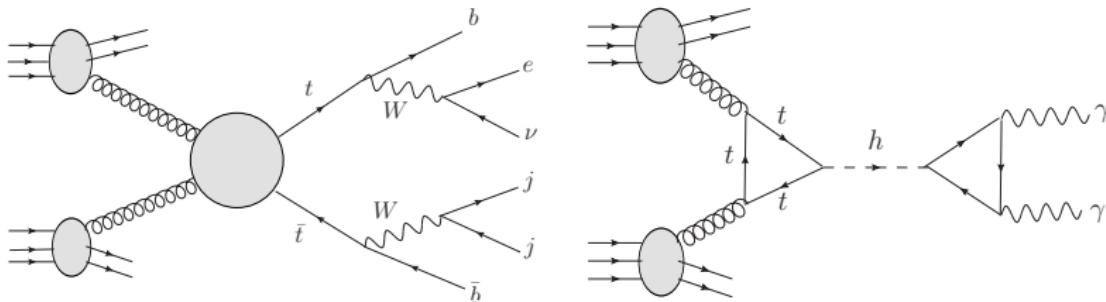
PDF $f(x, Q)$ is universal



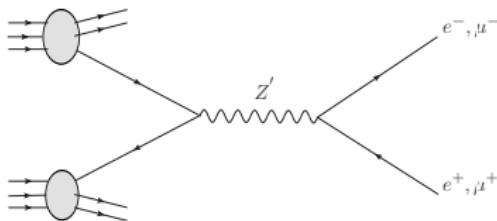
PDF(Parton distribution function) tell us the probability to find out a parton in a proton with particular momentum fraction x and energy Q .

$t\bar{t}$

Higgs



To determine the EW parameters via precision measurements requires knowing PDFs.



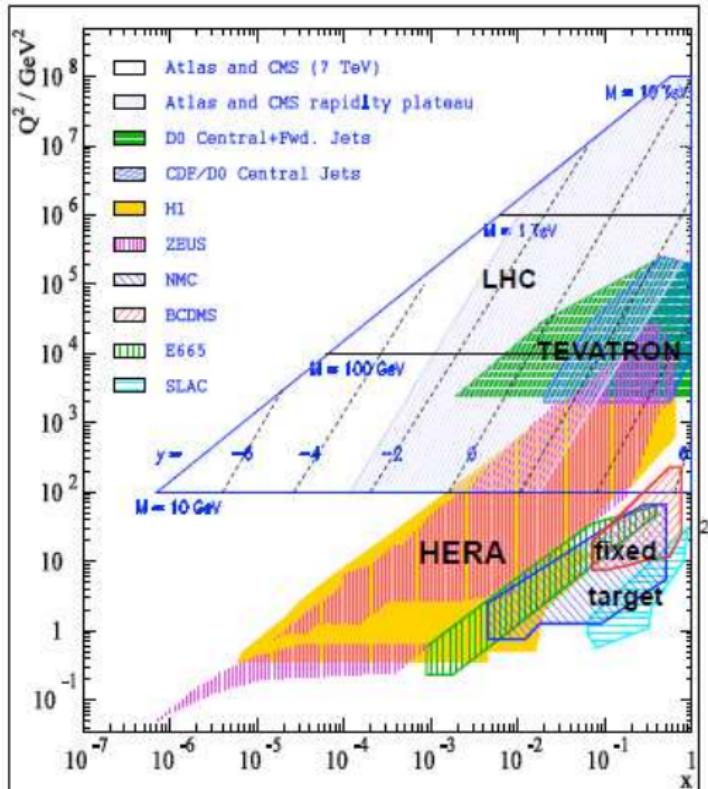
Probing New Physics requires knowing PDFs

PDF evolve

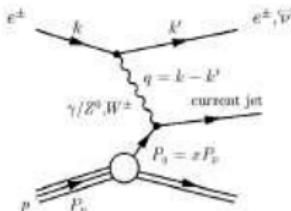
$$\begin{aligned}\frac{\partial q_i(x, \mu^2)}{\partial \ln \mu^2} &= P_{qq}^v \otimes q_i + P_{q\bar{q}}^v \otimes \bar{q}_i + P_{qq}^s \otimes \sum_k^{N_f} q_k + P_{q\bar{q}}^s \otimes \sum_k^{N_f} \bar{q}_k + P_{qg} \otimes g \\ \frac{\partial \bar{q}_i(x, \mu^2)}{\partial \ln \mu^2} &= P_{q\bar{q}}^v \otimes q_i + P_{qq}^v \otimes \bar{q}_i + P_{q\bar{q}}^s \otimes \sum_k^{N_f} q_k + P_{qq}^s \otimes \sum_k^{N_f} \bar{q}_k + P_{qg} \otimes g \\ \frac{\partial g(x, \mu^2)}{\partial \ln \mu^2} &= P_{gq} \otimes \sum_k^{N_f} (q_k + \bar{q}_k) + P_{gg} \otimes g\end{aligned}$$

DGLAP tells us how to evolve from the low energy scale, the input energy scale, to the high energy scale, the energy scale of interaction. But it does not tell us about its x-dependence. We need to use data to determine $f(x, Q)$ as a function of x at Q_0 scale.

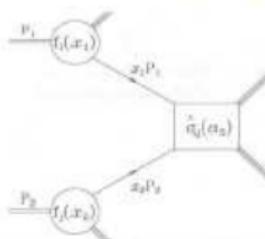
Experimental access to the proton structure



HERA: low and medium x



LHC: important constraints on $g(x)$, flavour separation



Fixed Target: high x , nuclear PDFs

CT global analysis takes $Q_0 = 1.3\text{GeV} \gg \Lambda_{QCD}$, and assume

$$xf_a(x, Q_0, \{a_1, a_2, \dots\}) = x^{a_1}(1-x)^{a_2}P_a(x)$$

- $x \rightarrow 0$: $f \propto x^{a_1}$, Regge-like behavior
- $x \rightarrow 1$: $f \propto (1-x)^{a_2}$, quark counting rules
- $P(x; a_3, a_4, \dots)$: affects intermediate x ; In CT14, Bernstein polynomial is applied.

For every proton flavor:

$$g, \quad u_v, \quad d_v, \quad s, \quad \bar{u}_s = u_s, \quad \bar{d}_s = d_s$$

Where $u = u_v + u_s$, and $d = d_v + d_s$.

In total, there are about 28 shape parameters used in CT14.

Requirements for PDF parametrization

- Valence quark number sum rule

$$\int_0^1 [u(x) - \bar{u}(x)] dx = 2, \int_0^1 [d(x) - \bar{d}(x)] dx = 1$$

$$\int_0^1 [s(x) - \bar{s}(x)] dx = 0$$

Where $u = u_v + \bar{u}$, $d = d_v + \bar{d}$. $(s(x) - \bar{s}(x))$ can be non-zero.

- Momentum sum rule

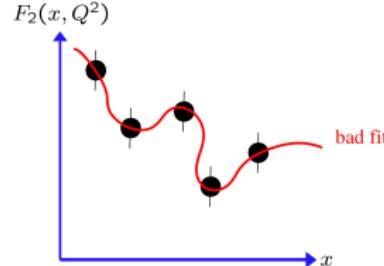
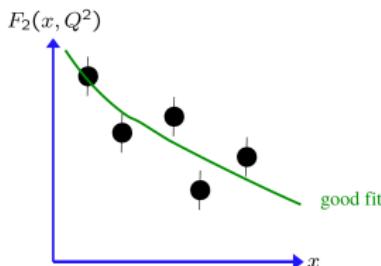
$$\sum_{a=q,\bar{q},g} \int_0^1 x f_{a/p}(x, Q) dx = 1$$

Where

$$\int_0^1 x g(x, Q) dx \sim 0.45$$

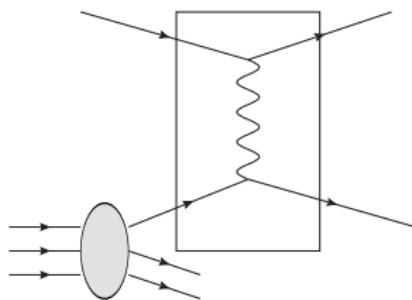
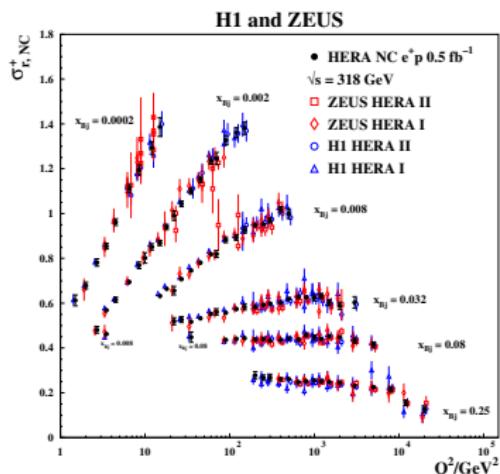
Requirements for PDF parametrization

- A valid PDF set must not produce unphysical predictions for observables
 - Any conceivable hadron cross section σ must be non-negative: $\sigma > 0$. This is typically realized by requiring $f_{a/p}(x, Q) > 0$.
 - Any cross section asymmetry A must lie in the range $-1 \leq A \leq 1$. This constrains the range of allowed PDF parametrizations.
- PDF parametrization for $f_{i/p}(x, Q)$ must be "flexible just enough" to reach agreement with the data, without reproducing random fluctuation.



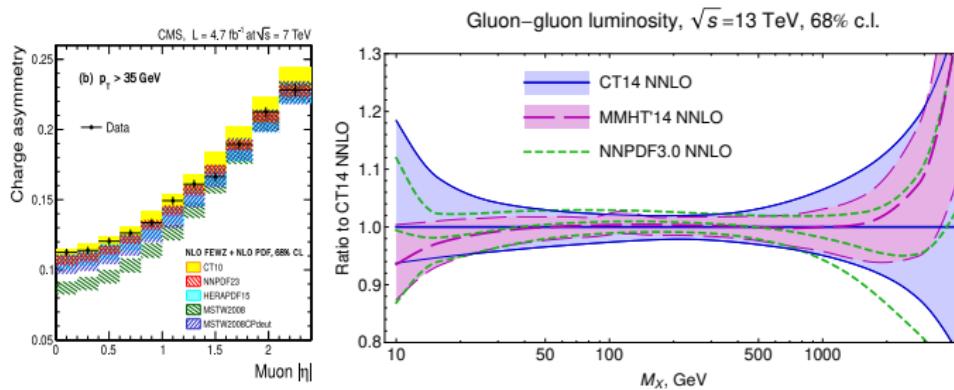
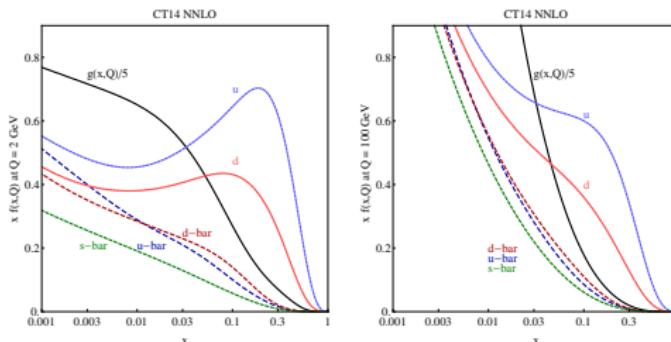
PDF is determined by comparing data and hard cross section

$$\sigma = f(x, Q^2, \{a\}) \otimes \hat{\sigma}$$



PDF	Hard part
LO	LO
NLO	NLO
NNLO	NNLO

PDF will then contribute to the precision measurement and search for new physics.



Criteria for determine PDFs

Define

$$\chi^2_{global} = \sum_i \frac{[D_i - \sum_k \lambda_k \beta_{ki} - T_i(\{a\})]^2}{\sigma_i^2} + \sum_k \lambda_k^2.$$

Where

D_i is the central value of data,

$T_i(\{a\})$ is the theoretical prediction of the data,

σ_i^2 is the quadratic sum of the statistical error and uncorrelated error,

β_{ki} is the matrix for correlated error,

and λ_k are the nuisance parameters.

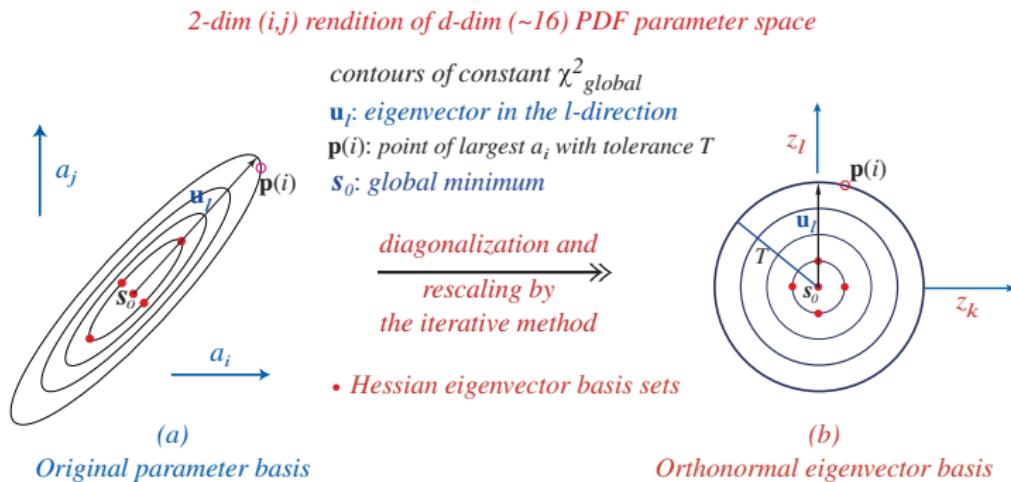
The PDF is obtained by minimize the global χ^2 function respect to $\{a\}$ and $\{\lambda\}$.

Probe the uncertainty of PDFs

The uncertainty of the PDF is estimated by the Hessian method.

$$\chi^2 = \chi_0^2 + \sum_{i,j} H_{ij} y_i y_j, \quad H_{ij} = \frac{1}{2} \left(\frac{\partial^2 \chi^2}{\partial y_i \partial y_j} \right)_0,$$

Where $y_i = a_i - a_i^0$ with a_i^0 to be the parameters at minimal χ_0^2 .



Let $X = X(\{a_i\})$ to be the observable as a function of fitting parameter. Using the linear approximation of parameter $\{z_i\}$, the symmetry uncertainty of X is,

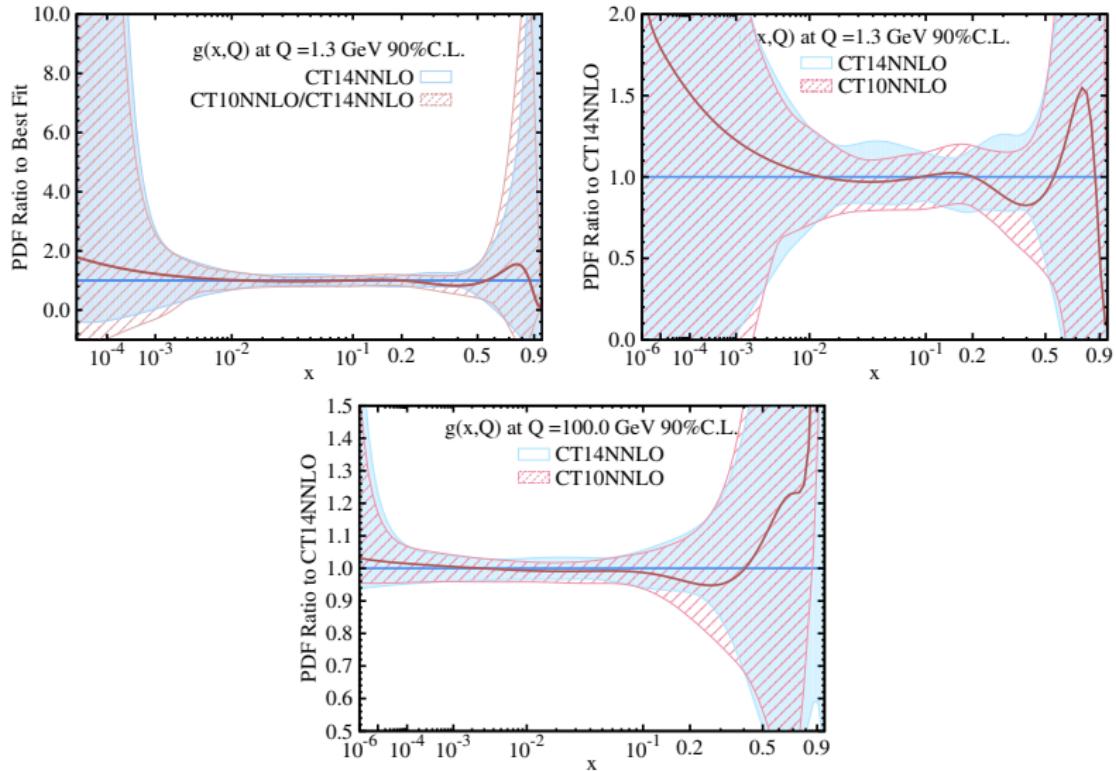
$$\Delta X = \frac{1}{2} \left(\sum_{i=1}^{N_p} [X(\{z_i^+\}) - X(\{z_i^-\})]^2 \right)^{1/2},$$

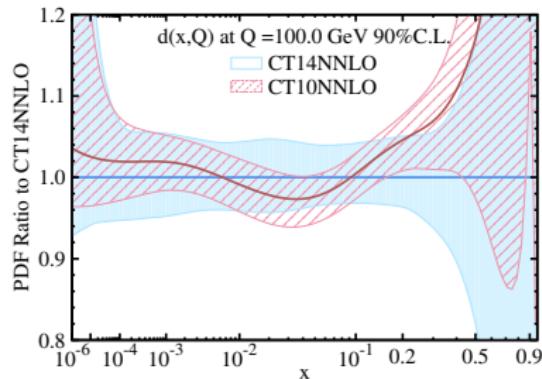
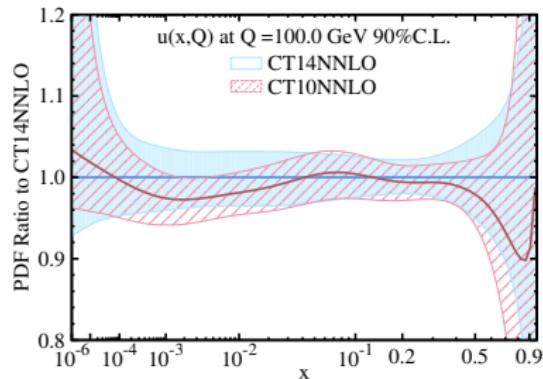
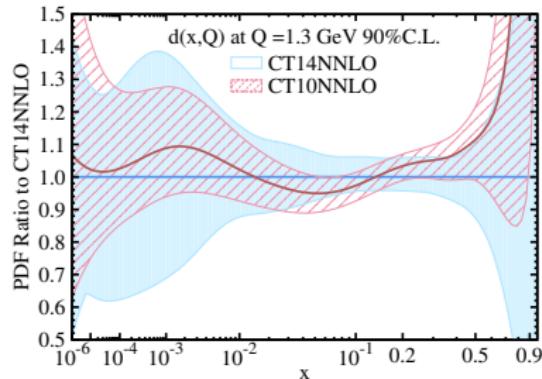
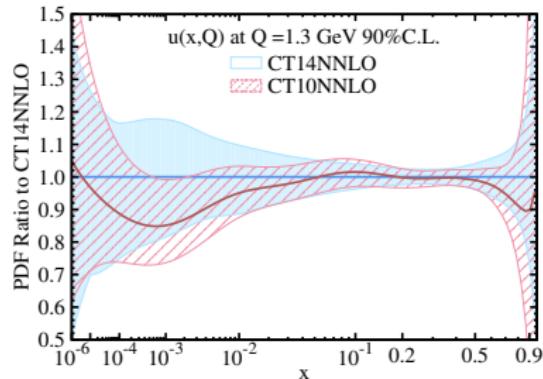
Where $\{z_1^\pm\} = \{\pm T, 0, \dots\}$, $\{z_2^\pm\} = \{0, \pm T, 0, \dots\}$ and so on.
The asymmetry uncertainty of X is,

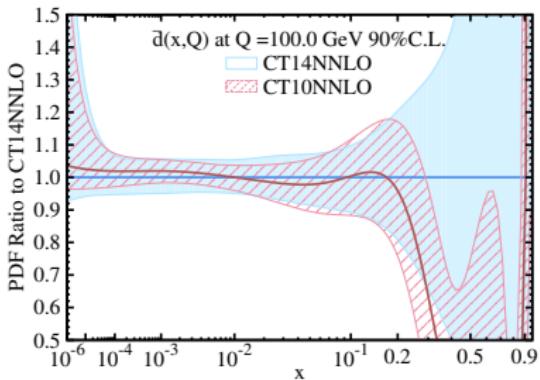
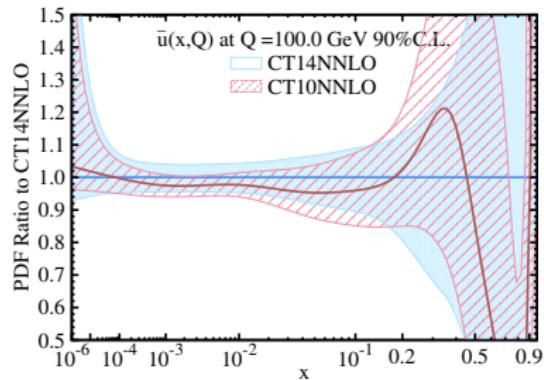
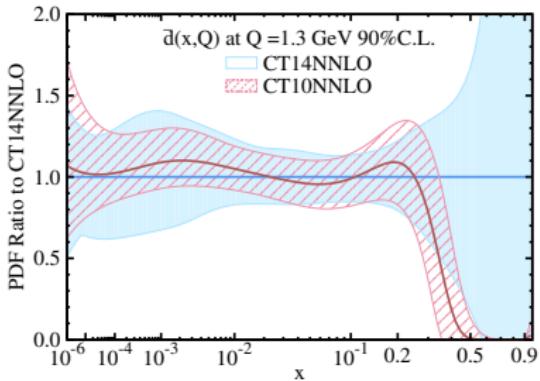
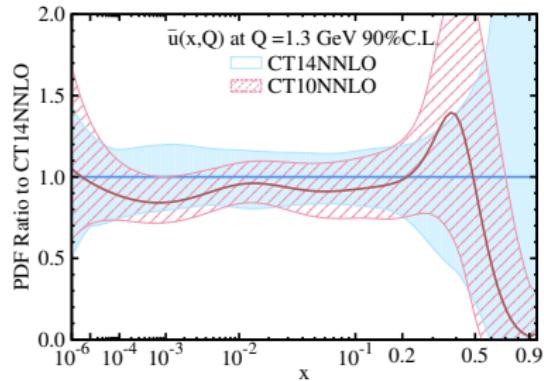
$$\delta^+ X = \sqrt{\sum_{i=1}^{N_a} \left[\max(X_i^{(+)} - X_0, X_i^{(-)} - X_0, 0) \right]^2},$$

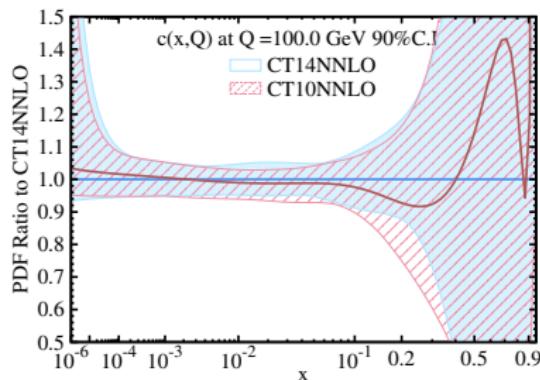
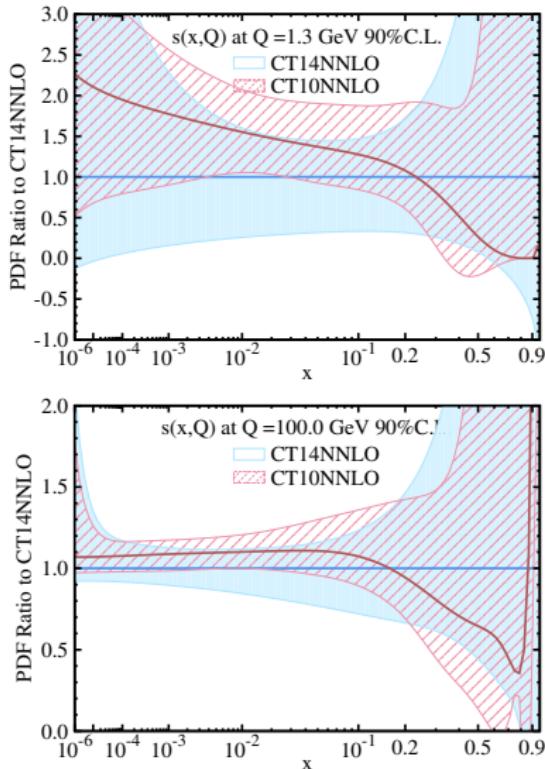
$$\delta^- X = \sqrt{\sum_{i=1}^{N_a} \left[\max(X_0 - X_i^{(+)}, X_0 - X_i^{(-)}, 0) \right]^2},$$

$$g(x, Q)$$

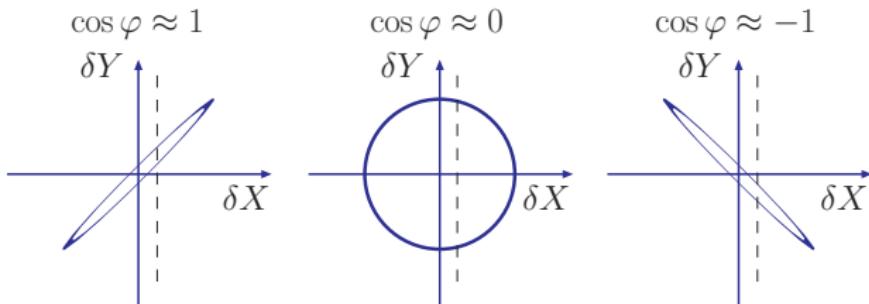


$u(x, Q)$ $d(x, Q)$ 

$\bar{u}(x, Q)$ $\bar{d}(x, Q)$ 

$S(x, Q)$ $c(x, Q)$ 

Is there a way to see the potential impact to PDF
before real global analysis?



In the framework of the Hessian, the correlation between two variables X and Y can be worked out as.

$$\cos \varphi = \frac{\vec{\nabla}X \cdot \vec{\nabla}Y}{\Delta X \Delta Y} = \frac{1}{4\Delta X \Delta Y} \sum_{\alpha=1}^N \left(X_{\alpha}^{(+)} - X_{\alpha}^{(-)} \right) \left(Y_{\alpha}^{(+)} - Y_{\alpha}^{(-)} \right)$$

where the ΔX and ΔY are their symmetric uncertainties. By this correlation angle φ , the tolerance ellipse is defined by

$$X = X_0 + \Delta X \cos \theta, \quad Y = Y_0 + \Delta Y \cos(\theta + \varphi),$$

- The correlation cosine between PDF $f(x, \mu)$ and theoretical prediction T_i contains no information of the experimental uncertainty.
- The correlation cosine $C_f(x_i, \mu_i)$ between PDF $f(x, \mu)$ and residual r_i contains no information of the experimental uncertainty in practice

$$C_f(x_i, \mu_i) = \frac{\vec{\nabla}f(x_i, \mu_i) \cdot \vec{\nabla}r_i}{\Delta f(x_i, \mu_i) \Delta r_i}, \quad \text{where}$$

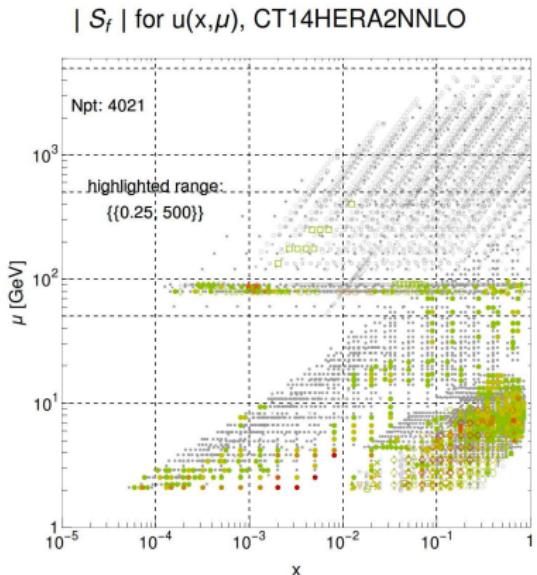
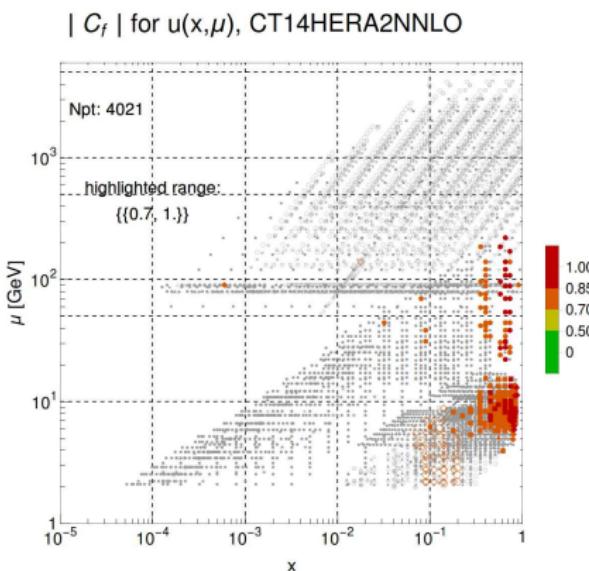
$$\chi^2 = \sum_i^N r_i^2 + \sum_k \lambda_k^2, \quad r_i(\vec{a}) = \frac{D_i - \sum_k \lambda_k \beta_{ki} - T_i(\{a\})}{\sigma_i}$$

- Instead, we concern the "sensitivity" $S_f(x_i, \mu_i)$

$$S_f(x_i, \mu_i) = C_f(x_i, \mu_i) \frac{\Delta r_i}{\sqrt{\frac{\sum_i^N r_i^2}{N}}}$$

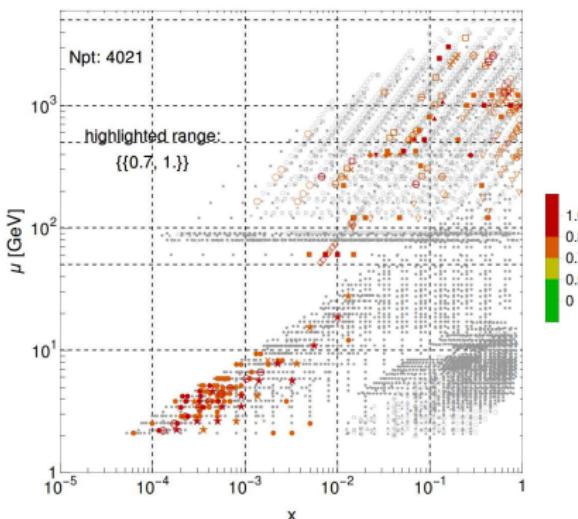
The sensitivity $S_f(x_i, \mu_i)$ help us to visualize the potential impact on PDF in $x - Q$ plane.

$$u(x, \mu)$$

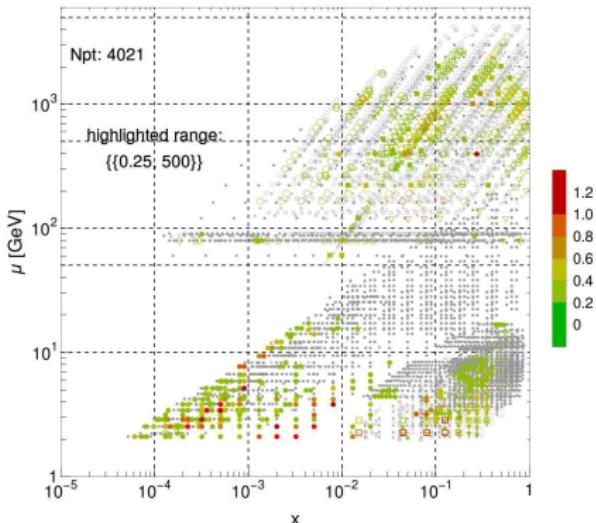


$$g(x, \mu)$$

| C_f | for $g(x,\mu)$, CT14HERA2NNLO



| S_f | for $g(x, \mu)$, CT14HERA2NNLO



Hessian Updating

- Updated Chi-square function :

$$\Delta\chi^2(Z) = \Delta\chi^2_{old}(Z) + (X_i^E - X_i(Z))C_{ij}^{-1}(X_j^E - X_j(Z))$$

- Hessian approximation :

$$\Delta\chi^2_{old}(Z) = T^2 Z^2 \quad (\text{T} = \text{tolerance parameter})$$

$$X_i(Z) = X_i(0) + \Delta X_i \cdot Z \quad \text{with} \quad \Delta X_i^\alpha = \frac{1}{2}(X_i(+e)^\alpha - X_i(-e)^\alpha)$$

- Minimize to find new best fit:

$$Z_{new}^2 = (1 + M)^{-1} A \quad \text{with}$$

$$A^\alpha = \frac{1}{T^2}(X_i^E - X_i(0))C_{ij}^{-1}\Delta X_j^\alpha,$$

$$M^{\alpha\beta} = \frac{1}{T^2}\Delta X_i^\alpha C_{ij}^{-1}\Delta X_j^\beta$$

Updated PDF set

- New best-fit PDF :

$$f_{new}^0 = f^0 + \Delta f \cdot Z$$

- New error PDFs :

$$f^{\pm(r)} = f_{new}^0 \pm \frac{1}{\sqrt{1 + \lambda^{(r)}}} \Delta f \cdot U^{(r)}$$

where $\lambda^{(r)}$ and $U^{(r)}$ are the eigenvalues and eigenvectors of matrix M .

- Direct update of other observables:

$$Y_{new}^0 = Y^0 + \Delta Y \cdot Z$$

$$|\Delta Y| = \Delta Y \cdot (1 + M)^{-1} \Delta Y$$

ePump

Error PDF Updating Method Package

```
FullCT14HERA2.in -- Edited --
++ N(EV pairs)          N(Data Sets)  PDFtype(C/L/N)  Dyn_Tol? (Y/N)  Tol_squared
27                      33           C                  Y                   100.0
++ ObservableFile
N(Observables) Data?(Y/N) Error_type Weight
1120          Y             1.0
CT14HERA2ex/tabs/E160.If1363
337          Y             1.0
CT14HERA2ex/tabs/E181.If1363
258          Y             1.0
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15           Y             1.0
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184          Y             1.0
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".in" file

How to use ePump

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*** N(EV pairs)          * DR515 electron charge asymmetry from W decays from D0 Run-2 9.7 fb-1 (1412.2862)
    27                   * Easy for electron Et>25 GeV and neutrino Et>25 GeV; sqrt(s)=1968 GeV, uncorrelated
*** ObservableFile        * MG15 NLO and NNLO ratios K(W-)/K(W+) for CT14 NNLO. normalized to CT-package LO: + th
CT14HERA2ex/tabs/E160..If1363   3 : NormErr, # of corr_err, Ecm, M_W, METin
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CT14HERA2ex/tabs/E230..If1363
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CT14HERA2ex/tabs/E271..If1363
CT14HERA2ex/tabs/E280..If1363
CT14HERA2ex/tabs/E281..If1363 13   Y   1   1.0
CT14HERA2ex/tabs/E282..If1363 11   Y   1   1.0
CT14HERA2ex/tabs/E285..If1363 72   Y   1   1.0
CT14HERA2ex/tabs/E514..If1363 118  Y   1   1.0
CT14HERA2ex/tabs/E535..If1363 90   Y   1   1.0
CT14HERA2ex/tabs/E538..If1363 133  Y   1   1.0
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PDFs/CT14HERA2ex/If1363  CT14HERA2ex/PDFtmp/If1363

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“.in” file

How to use ePump

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*** N(EV pairs) 27
*** ObservableFile
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CT14HERA2ex/tabs/E101_If1363
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CT14HERA2ex/tabs/E111_If1363
CT14HERA2ex/tabs/E120_If1363
CT14HERA2ex/tabs/E125_If1363
CT14HERA2ex/tabs/E127_If1363
CT14HERA2ex/tabs/E147_If1363
CT14HERA2ex/tabs/E145_If1363
CT14HERA2ex/tabs/E169_If1363
CT14HERA2ex/tabs/E201_If1363
CT14HERA2ex/tabs/E203_If1363
CT14HERA2ex/tabs/E204_If1363
CT14HERA2ex/tabs/E225_If1363
CT14HERA2ex/tabs/E227_If1363
CT14HERA2ex/tabs/E234_If1363
CT14HERA2ex/tabs/E260_If1363
CT14HERA2ex/tabs/E261_If1363
CT14HERA2ex/tabs/E267_If1363
CT14HERA2ex/tabs/E268_If1363
CT14HERA2ex/tabs/E241_If1363
CT14HERA2ex/tabs/E281_If1363
CT14HERA2ex/tabs/E266_If1363
CT14HERA2ex/tabs/E584_If1363
CT14HERA2ex/tabs/E514_If1363
CT14HERA2ex/tabs/E535_If1363
CT14HERA2ex/tabs/E538_If1363

D0S15 electron charge asymmetry from W decays from D0 Run-2 9.7 fb-1 (1412.2862)
  Easy for electron Et>20 GeV and neutrino Et>25 GeV; sqrt(S)=1960 GeV, uncorrelated
  MG15 NLO & NNLO ratios K(W-)/K(W+) for CT14 NNLO, normalized to CT-package LO: + th
  # : NormErr, # of corr_err, Ecm, M_W, Method
  0.8 : 6 1960. 88.380E-25d0
  # of corr_err, Data Col., Stat Err Column, UncSys Column, corr_err Col
  0.8 4 5 7 8 9 10 11 12 13
  ymid pTeffMIN pTeffMAX Easy StatErr TotSys UncSys lob e01% e04% e05% e06%
  0.1 25.0 9.80E+02 0.821 0.8012 0.8011 0.8006 2 0.29 0.14 0.19 1.33
  0.3 25.0 9.80E+02 0.853 0.8011 0.8011 0.8004 3 0.29 0.14 0.19 1.33
  0.5 25.0 9.80E+02 0.874 0.8011 0.8011 0.8004 4 0.29 0.14 0.19 1.33
  0.7 25.0 9.80E+02 0.894 DATA SET 281 ; Norm Fac = 1.00000 ; # of pts = 13 ;
  0.9 25.0 9.80E+02 0.914 R2, r(k) = 4.934 0.110 0.081 0.069 -2.286 0.
  1.1 25.0 9.80E+02 0.934 Y Q Rs Exp Th./Norm
  1.39 25.0 9.80E+02 0.954 Theory Column
  1.7 25.0 9.80E+02 0.974 Data IIf1363.00.dts
  1.9 25.0 9.80E+02 0.984 Data IIf1363.01.dts
  2.1 25.0 9.80E+02 -0.1.00E+00 0.839E+01 1.968E+03 2.18000E-02 1.93596E-02
  2.3 25.0 9.80E+02 -0.1.00E+00 0.839E+01 1.968E+03 2.23000E-02 5.53459E-02
  2.54 25.0 9.80E+02 -0.1.00E+00 0.839E+01 1.968E+03 9.16000E-02 9.12121E-02
  2.92 25.0 9.80E+02 -0.1.00E+00 0.839E+01 1.968E+03 1.19700E-01 1.22005E-01
  9.00E+00 0.839E+01 1.968E+03 1.45200E-01 1.47812E-01
  1.100E+00 0.839E+01 1.968E+03 1.55900E-01 1.62549E-01
  1.390E+00 0.839E+01 1.968E+03 1.57300E-01 1.62798E-01
  1.700E+00 0.839E+01 1.968E+03 1.18000E-01 1.26935E-01
  1.900E+00 0.839E+01 1.968E+03 6.66000E-02 5.79517E-02
  2.100E+00 0.839E+01 1.968E+03 -0.55000E-02 2.17415E-03
  2.300E+00 0.839E+01 1.968E+03 -0.97000E-02 -0.28367E-02
  2.540E+00 0.839E+01 1.968E+03 -0.91000E-01 -0.23848E-01
  2.920E+00 0.839E+01 1.968E+03 -0.3.99700E-01 -0.4.31176E-01
  Data : IIf1363.01.dts
  1.00E-00 0.839E+01 1.968E+03 2.18000E-02 1.96800E-02
  3.00E-01 0.839E+01 1.968E+03 5.23000E-02 5.63526E-02
  5.00E-01 0.839E+01 1.968E+03 9.16000E-02 9.28178E-02
  7.00E-01 0.839E+01 1.968E+03 1.45200E-01 1.24136E-01
  9.00E-01 0.839E+01 1.968E+03 1.45200E-01 1.24136E-01
  1.100E+00 0.839E+01 1.968E+03 0.55900E-01 0.65377E-01
  1.390E+00 0.839E+01 1.968E+03 0.53700E-01 0.65749E-01
  1.700E+00 0.839E+01 1.968E+03 1.10000E-01 2.96277E-01
  1.900E+00 0.839E+01 1.968E+03 6.66000E-02 7.82553E-02
  2.100E+00 0.839E+01 1.968E+03 -0.1.55900E-02 0.88933E-03

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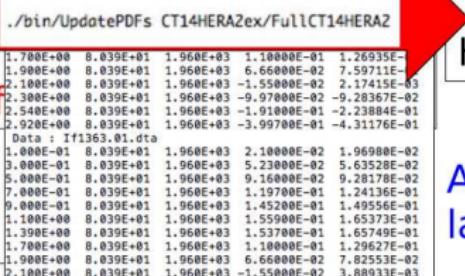
“.theory” file

How to use ePump

```
+++ N(EV pairs)          * DRS15 electron charge asymmetry from W decays from D0 Run-2 9.7 fb^-1 (1412.2862)
    27                   * Easy for electron Et>25 GeV and neutrino Et>25 GeV; sqrt{S}=1960 GeV, uncorrelated
    + ObservablesFile      * MG15 NLO & NNLO ratios K(W-) / K(W+) for CT14 NNLO, normalized to CT-package L0: + th
CT14HERA2ex/tabs/E160.If1363 3 : NormErr, # of corr_err, Ecm, M_W, METmin
CT14HERA2ex/tabs/E161.If1363 0.8   6   1968.  0.38E0  25d0
CT14HERA2ex/tabs/E162.If1363 # of corr_err , Data Column, StatErr Column, UncSys Column, corr_err Col
CT14HERA2ex/tabs/E164.If1363 6
CT14HERA2ex/tabs/E168.If1363 ymid pTMIN pTMAX Easy StatErr TotSys UncSys lob e01% e04% e05% e06%
CT14HERA2ex/tabs/E169.If1363 0.1 25.0 9.88E+02 0.021 0.0011 0.0006 2 0.29 0.14 0.19 1.33
CT14HERA2ex/tabs/E170.If1363 0.3 25.0 9.88E+02 0.023 0.0012 0.0011 3 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E171.If1363 0.5 25.0 9.88E+02 0.025 0.0013 0.0012 5 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E174.If1363 0.7 25.0 9.88E+02 0.027 0.0014 0.0013 7 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E175.If1363 0.9 25.0 9.88E+02 0.029 0.0015 0.0014 9 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E176.If1363 1.1 25.0 9.88E+02 0.031 0.0016 0.0015 11 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E177.If1363 1.39 25.0 9.88E+02 0.033 0.0017 0.0016 13 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E178.If1363 1.7 25.0 9.88E+02 0.035 0.0018 0.0017 15 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E179.If1363 1.9 25.0 9.88E+02 0.037 0.0019 0.0018 17 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E180.If1363 2.1 25.0 9.88E+02 0.039E-01 0.0020 0.0019 19 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E281.If1363 2.3 25.0 9.88E+02 0.041E-01 0.0021 0.0020 21 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E283.If1363 2.54 25.0 9.88E+02 0.043E-01 0.0022 0.0021 23 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E284.If1363 2.92 25.0 9.88E+02 0.045E-01 0.0023 0.0022 25 0.34 0.14 0.19 1.33
CT14HERA2ex/tabs/E227.If1363
CT14HERA2ex/tabs/E234.If1363
CT14HERA2ex/tabs/E260.If1363
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CT14HERA2ex/tabs/E241.If1363
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CT14HERA2ex/tabs/E514.If1363
CT14HERA2ex/tabs/E535.If1363
CT14HERA2ex/tabs/E538.If1363
+++
PDFIn  PDFout
PDFs/CT14HERA2ex/If1363  CT14HERA2ex/PDFtmp/If1363
```

".data" file

".theory" file



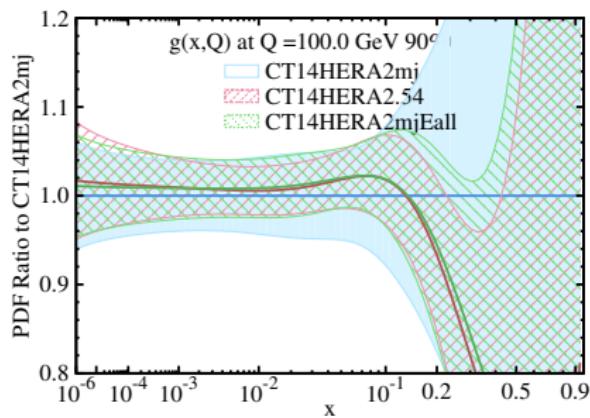
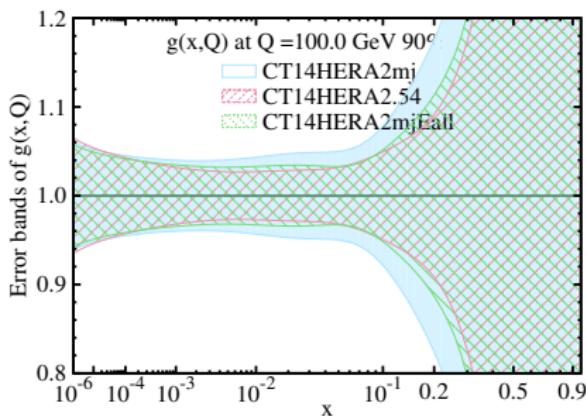
Updated best-fit
and
Hessian Error PDFs

A few seconds
later...

".in" file

Test 1: CT14HERA2 minus Jets

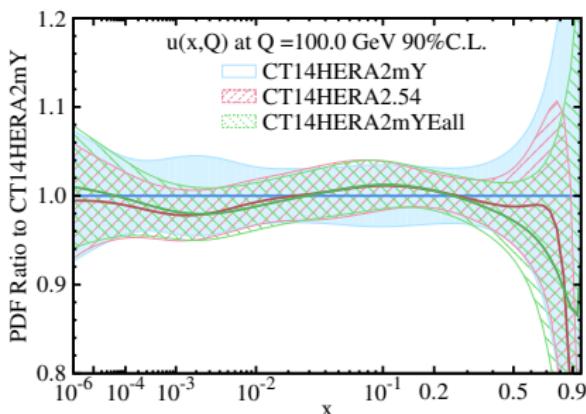
- Remove all CDF, D0, ATLAS 7TeV, CMS TeV jet data from CT14HERA2 and refit → CT14HERA2mj.
- Add back the 4 data sets to CT14HERA2mj by ePump and compare with CT14HERA2.



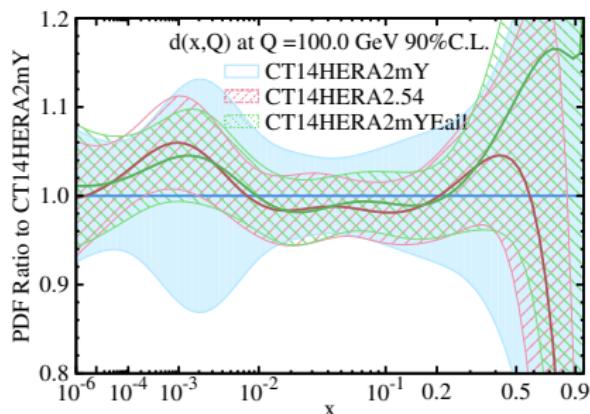
gluon - PDF

Test 2: CT14HERA2 minus Drell-Yan

- Remove 19 Drell-Yan data sets (fixed target exp., Tevatron and LHC) from CT14HERA2 and refit \rightarrow CT14HERA2mY.
- Add back the 19 data sets to CT14HERA2mY by ePump and compare with CT14HERA2.



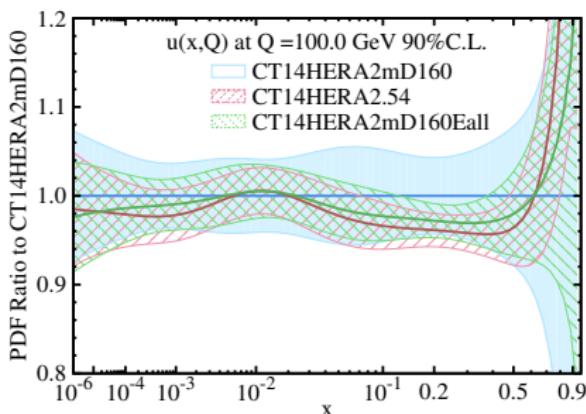
u-PDF



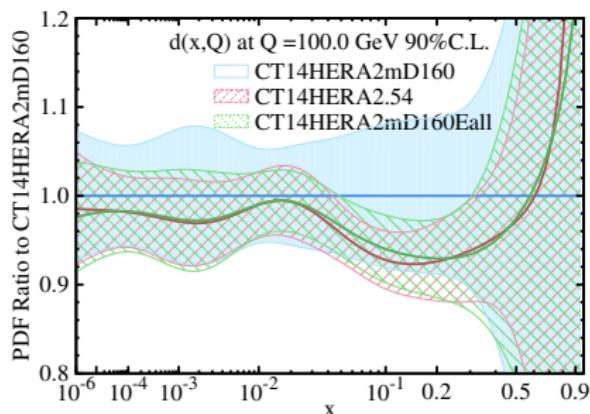
d-PDF

Test 3: CT14HERA2 minus DIS

- Remove 14 DIS data sets (excluding HERA1+2) from CT14HERA2 and refit → CT14HERA2mD160.
- Add back the 14 data sets to CT14HEA2mD160 by ePump and compare with CT14HERA2.



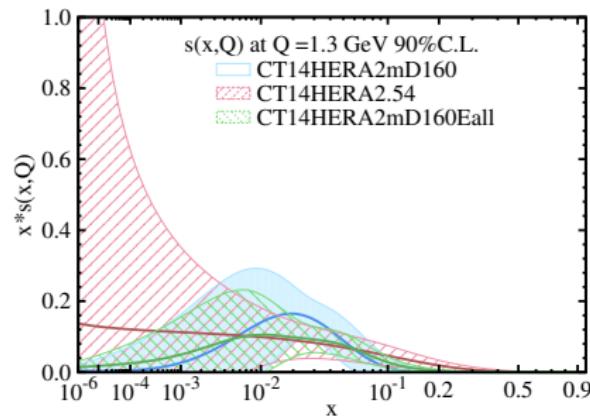
u-PDF



d-PDF

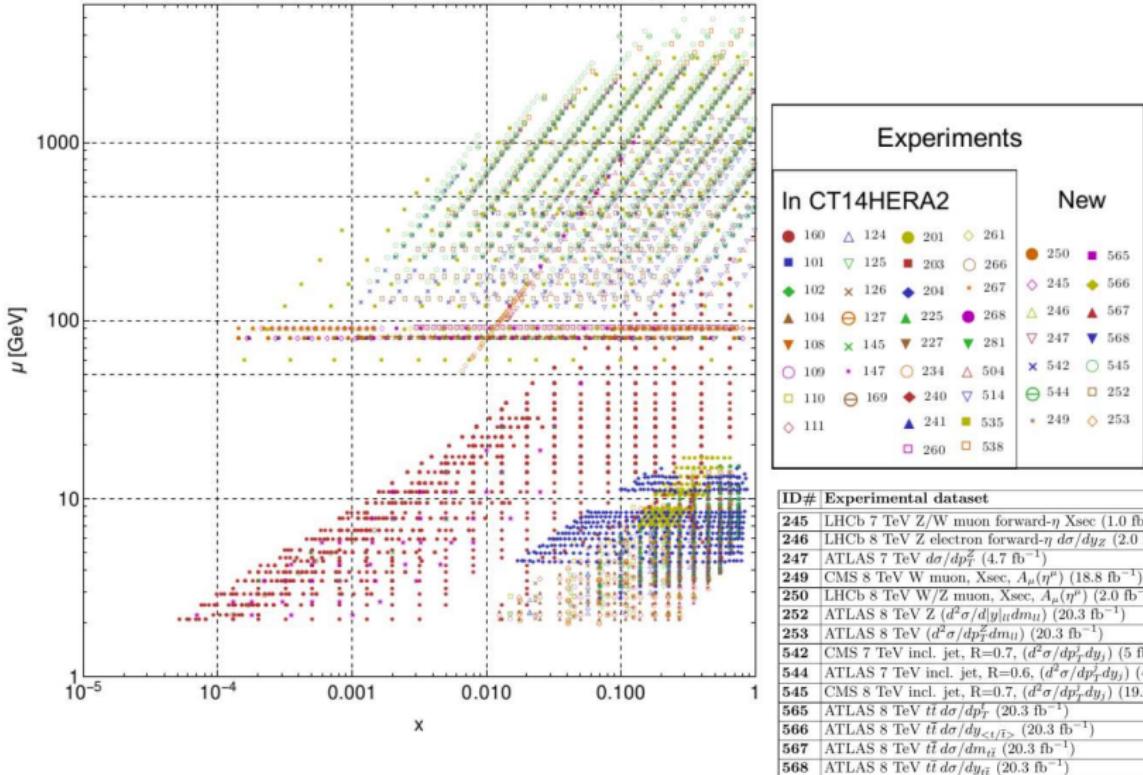
Test 3: CT14HERA2 minus DIS (2)

- Almost all of constraints on s-PDF are from NuTeV and CCFR neutrino dimuon data.



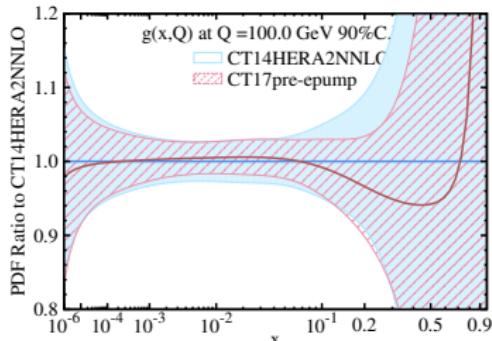
ePump work well in the x region constrained by data.

Experimental data in CTEQ-TEA PDF analysis

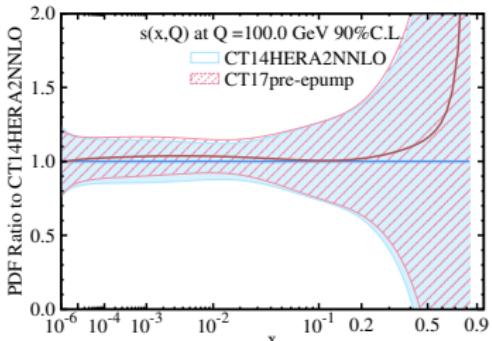


Adding New Data with ePump

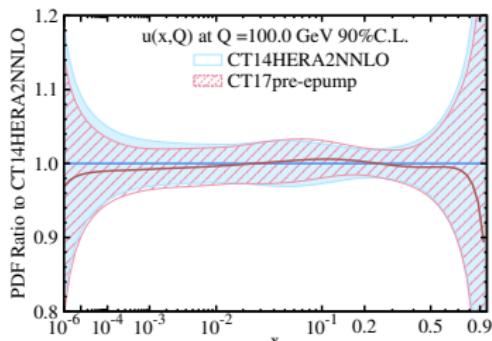
$$g(x, Q)$$



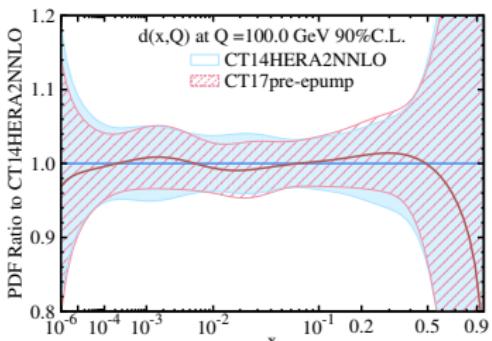
$$s(x, Q)$$



$$u(x, Q)$$



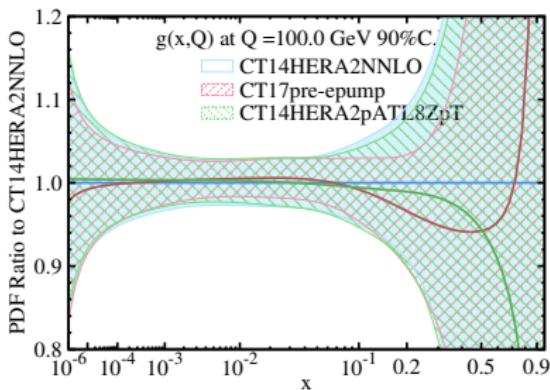
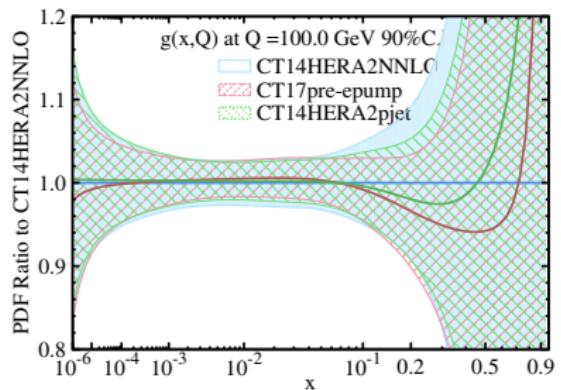
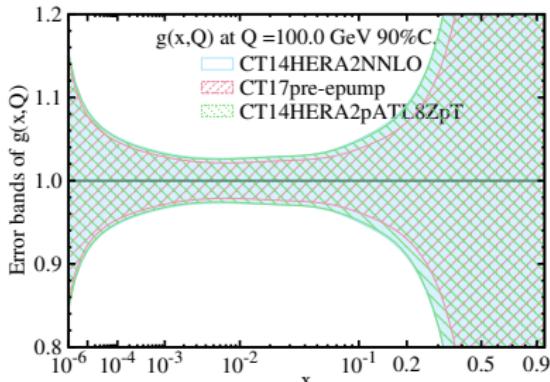
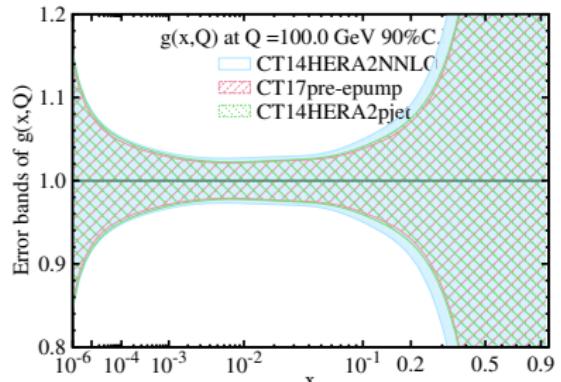
$$d(x, Q)$$



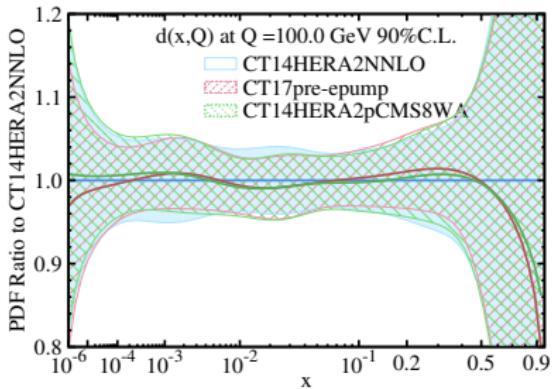
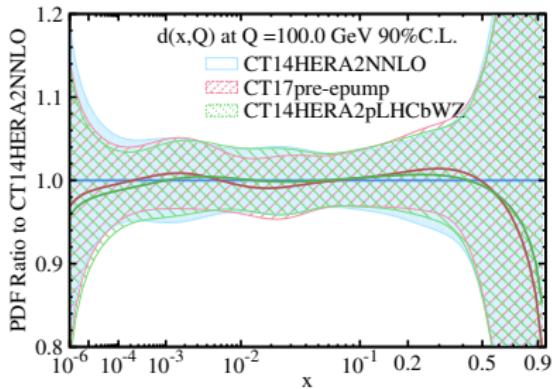
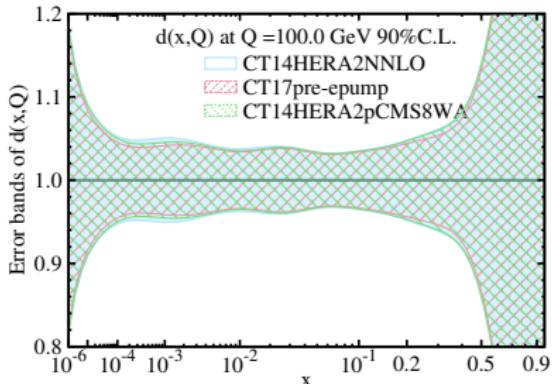
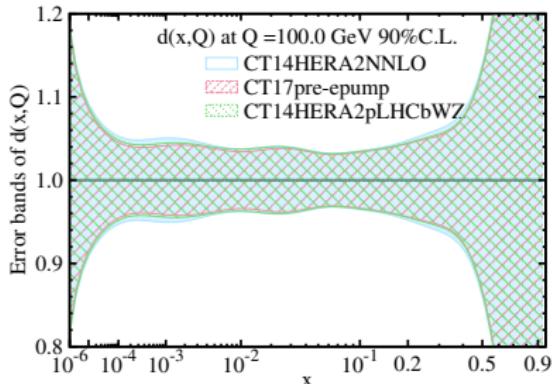
$$u(x, Q)$$

$$d(x, Q)$$

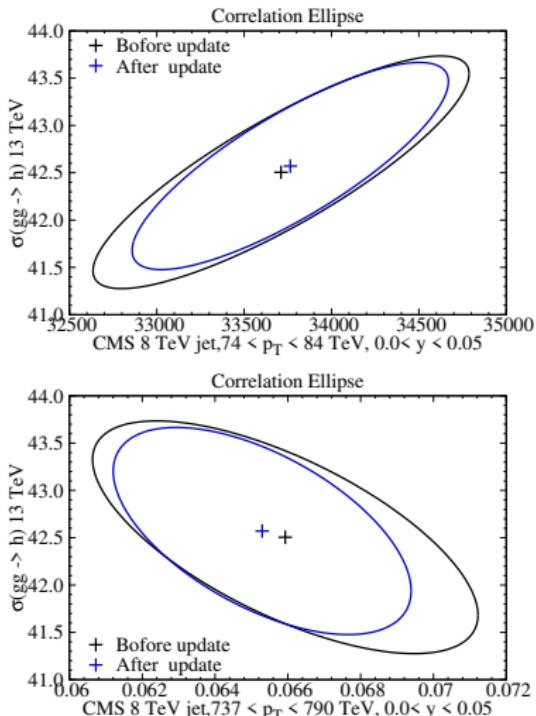
The update on gluon PDF is dominated by jet data.



The update on d quark come from LHCb WZ production and CMS W-asym data.



Effect of Updating on $\sigma(gg \rightarrow H)$



	$\sigma(gg \rightarrow h) 7 \text{ TeV}$
CT14HERA2	14.60 ± 0.44
updated	14.62 ± 0.40
	$\sigma(gg \rightarrow h) 8 \text{ TeV}$
CT14HERA2	18.60 ± 0.55
updated	18.64 ± 0.49
	$\sigma(gg \rightarrow h) 13 \text{ TeV}$
CT14HERA2	42.56 ± 1.23
updated	42.69 ± 1.07
	$\sigma(gg \rightarrow h) 14 \text{ TeV}$
CT14HERA2	47.99 ± 1.39
updated	48.14 ± 1.20

Summary

- The **Sensitivity** $S_f(x_i, \mu_i)$ provide an easy way to visualize the potential impact to PDF in $x - Q$ plane (arxiv:1803.02777).
- The **PDFsense**, which can work out the sensitivity, is public available at
<http://metapdf.hepforge.org/PDFSense/>
- **UpdatePDFs** function in **ePump** package is a fast and efficient method to estimate the effect of new data on the a current set of best-fit and Hessian error PDFs (arxiv:1806.07950).
- Extensively cross-checked against CT14 global fits
- ePump code is public available at
<http://hep.pa.msu.edu/epump/>