

*Hydrogen 21cm **absorption line** limits on dark matter and primordial black holes*

& highlights

高宇 Yu Gao

IHEP, CAS



中国科学院高能物理研究所

Institute of High Energy Physics Chinese Academy of Sciences

CMB 21cm absorption signals on passing through intergalactic gas: **neutral Hydrogen presence & cooler than CMB**



philosophy-of-cosmology.ox.ac.uk

Gas temperature decouples from CMB $z \sim 200$

Early reionization window (first discovery claim by EDGES)

[Bowman, et.al. Nature 555, 67 \(2018\).](#)

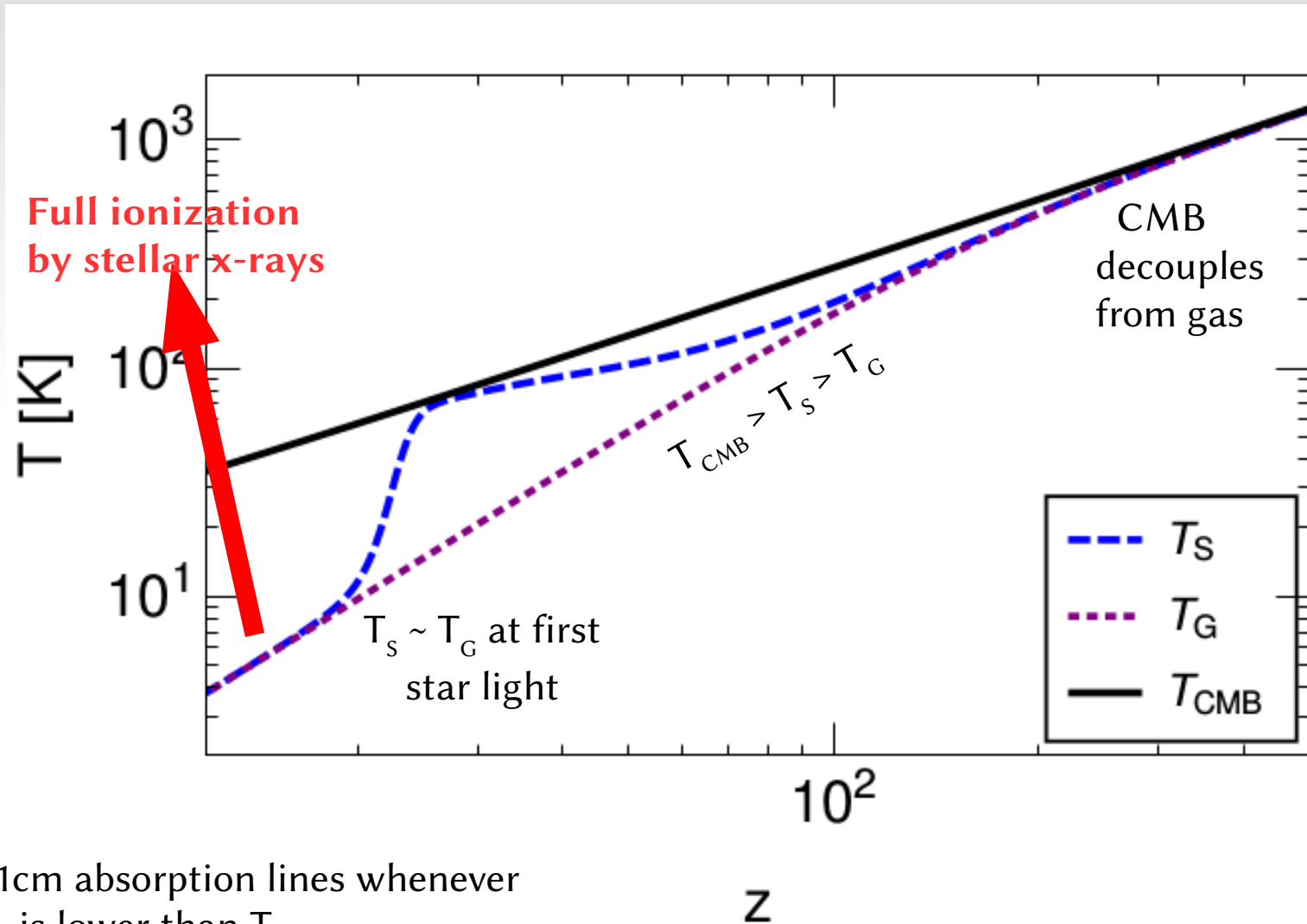
What 21cm data may tell us ...

- The cosmological reionization history:
Ionization fraction x_e , mean temperature T_G
- Distribution of neutral Hydrogen gas
temperature map & power spectrum $P(k)$

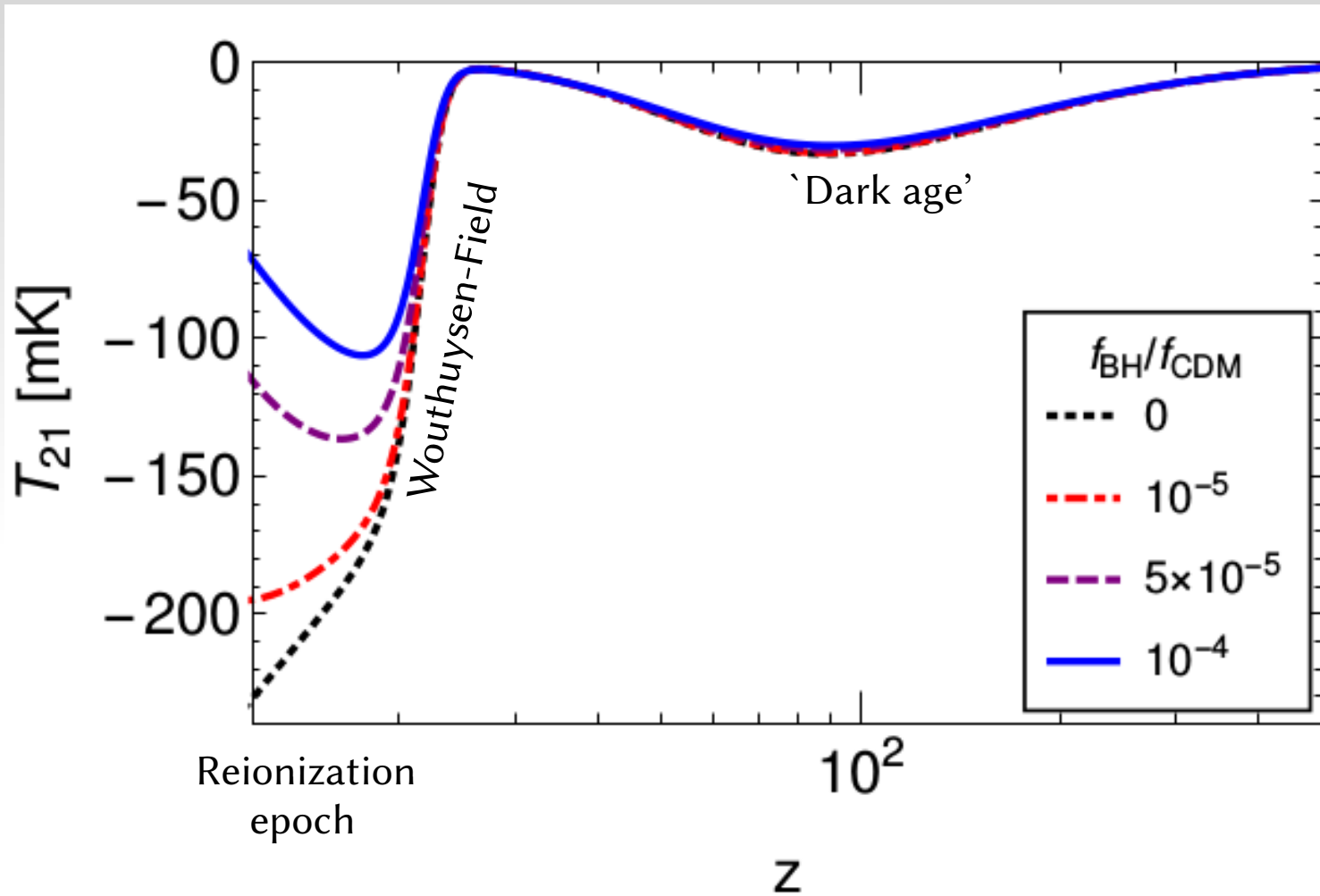
at [a] reionization epoch, $10 < z < 30$

[b] $z < 1$ galactic gas emissions (late time)

Reionization in the 'standard' astrophysics



21cm absorption lines whenever T_{S} is lower than T_{CMB} .



The average 'brightness temperature',
ignoring over-density and comoving velocity gradients

$$T_{21} \approx 0.023\text{K} \cdot x_{\text{H}_\text{I}}(z) \left(\frac{0.15}{\Omega_{\text{m}}} \cdot \frac{1+z}{10} \right)^{\frac{1}{2}} \frac{\Omega_{\text{b}} h}{0.02} \left(1 - \frac{T_{\text{CMB}}}{T_{\text{S}}} \right)$$

First claim of observation: The EDGES 21cm result

J. D. Bowman, A. E. E. Rogers, R. A. Monsalve,
T. J. Mozdzen, and N. Mahesh, *Nature* 555, 67 (2018).

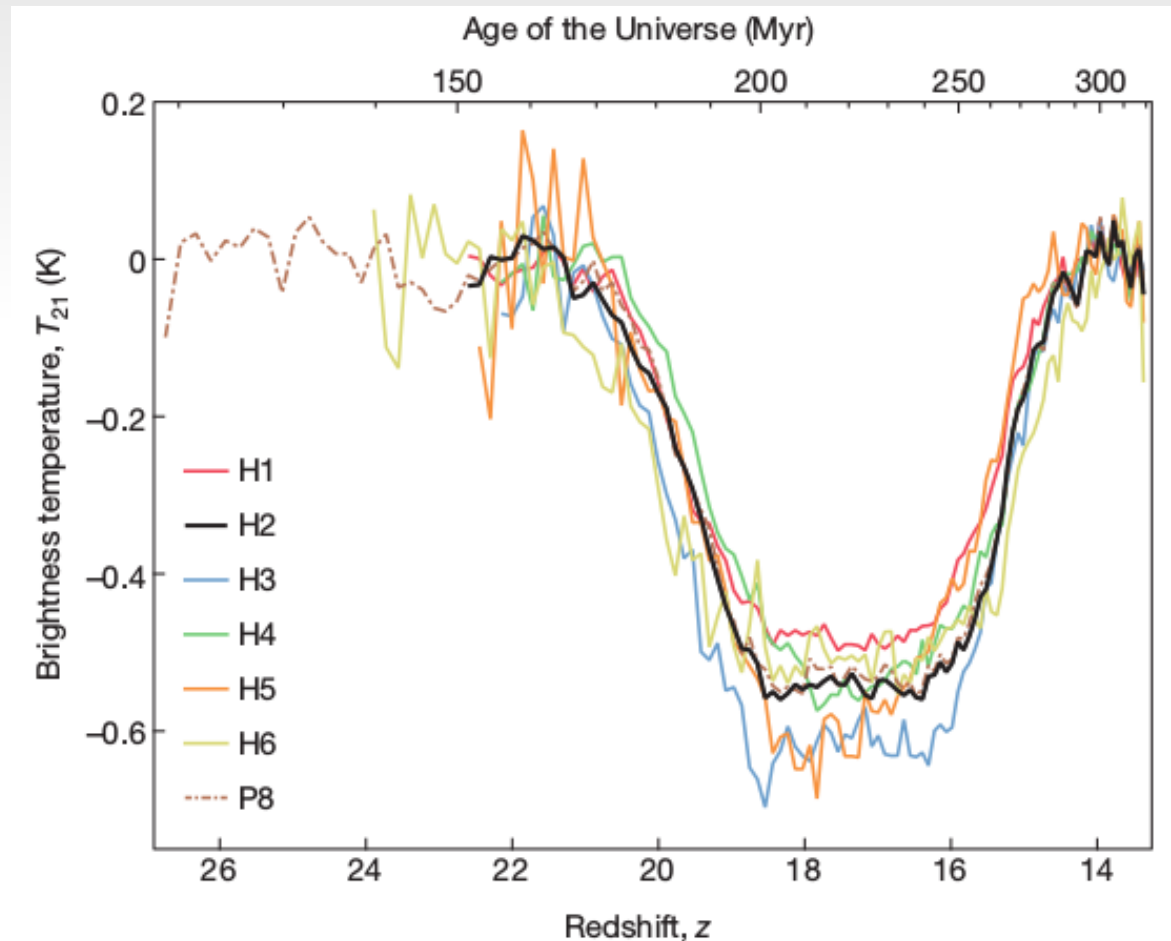
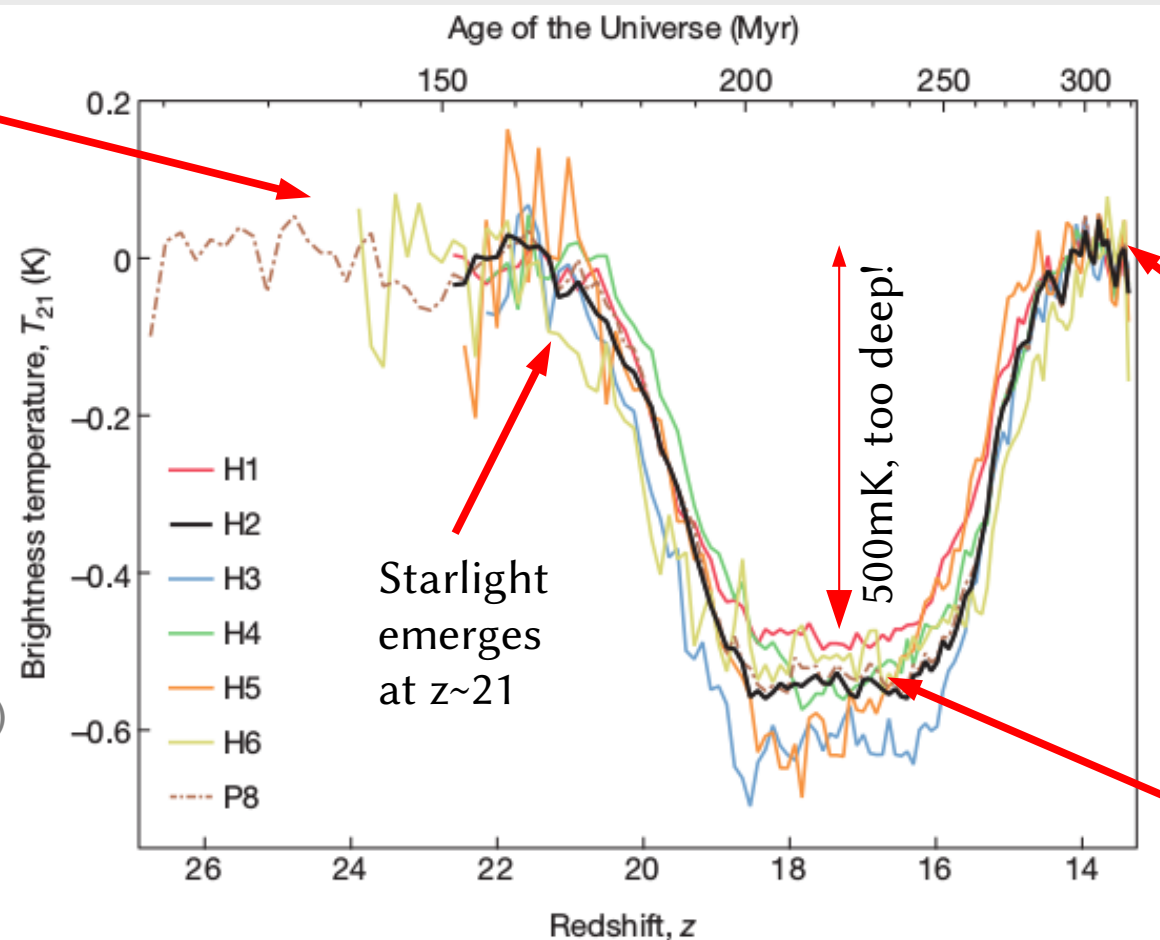


Figure 2 | Best-fitting 21-cm absorption profiles for each hardware case.

First claim of observation: The EDGES 21cm result

J. D. Bowman, A. E. E. Rogers, R. A. Monsalve,
T. J. Mozdzen, and N. Mahesh, Nature 555, 67 (2018).

$T_S \sim T_{\text{CMB}}$
no signal
for $z > 21$



For Λ CDM (non-) interpretations, see [S.Witte, et.al 1804.03888](#)

Figure 2 | Best-fitting 21-cm absorption profiles for each hardware case.

How about (particle) dark matter?

Yeah

DM cooling

Lower gas temperature via collisions: more 21cm signal

Explains the EDGES data
*needs large scattering xsec

DM heating

Raises gas temperature by energy injection: **reduces** 21cm signal

Most stringent bounds on DM annihilation, decays & other energy injections

Nay

CMB uncertainties

Large uncertainty at low frequency; new physics at radio-frequencies

Non-standard cosmology:

Modified Friedmann Eq.

The dark matter connections

- Dark matter can cool down T_G via collision
 - Dark matter can affect T_S via SD scattering
 - Non-CDM affects the 21cm temperature & power spectrum
 - Dark matter, or black holes, can indirectly ruin the 21cm signal if they have emissions, and ...
 - enhance ionized fraction of gas, affect CMB propagation
- Barkana, et.al. :*
A millicharged DM
with $\sigma \sim v^{-4}$
- M. Sitwell, et.al, MNRAS, 438, 2664*
- G. D'Amico, et.al., 1803.03629 (DM annihilation)*
S. Clark, et.al. 1803.09390 (DM decay & PBH rad.)
K. Cheung, et.al, 1803.09398 (DM annihilation)
A. Hektor, et.al, 1803.09697 (PBH accretion)
T. Slatyer, C.-L. Wu, 1803.09734 (DM ann. & decay)
- T. Slatyer, 1506.03811*
T. Slatyer and C.-L. Wu, 1610.06933.
S. Clark, et.al. 1612.07738

DM cooling theories: enhancing the 21cm signal

The scattering (H, He cooling, etc)
mediator must be screened by
SM particles

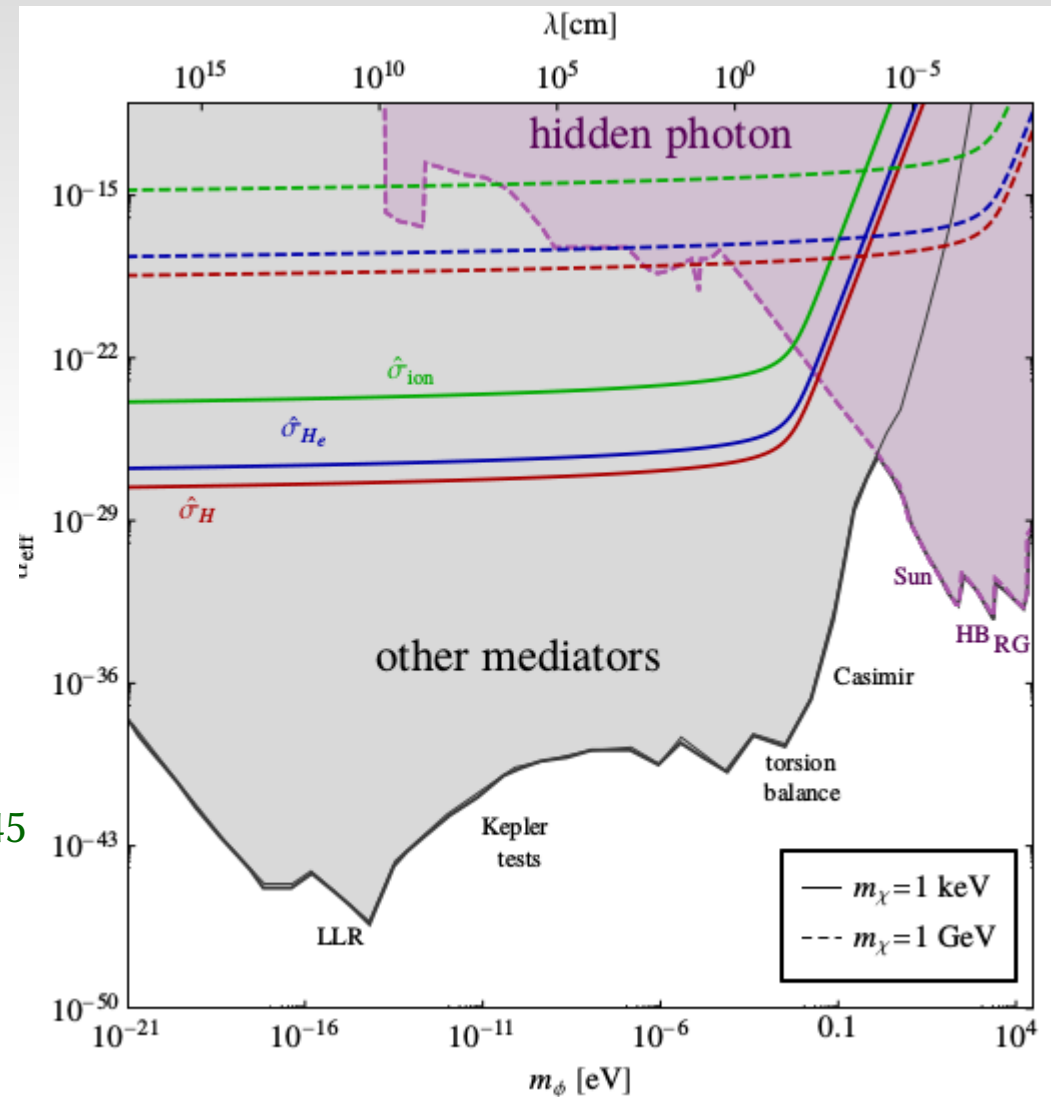
viable candidates: **photon & friends**

R. Barkana, et.al. 1803.03091

A few more DM cases,

S. Fraser et.al 1803.03245

Z. Kang, 1803.04928

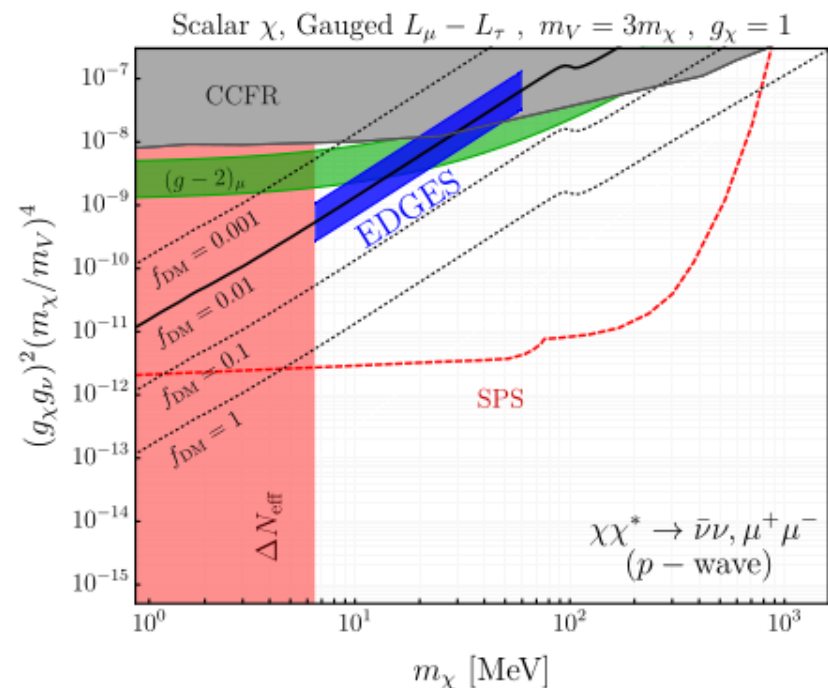
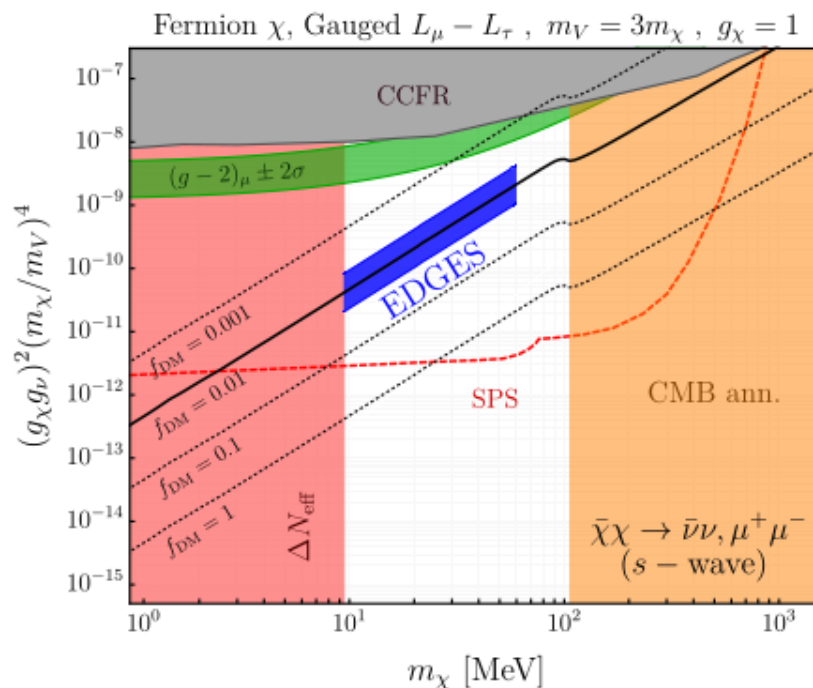


Dark matter as an explanation to the EDGES data

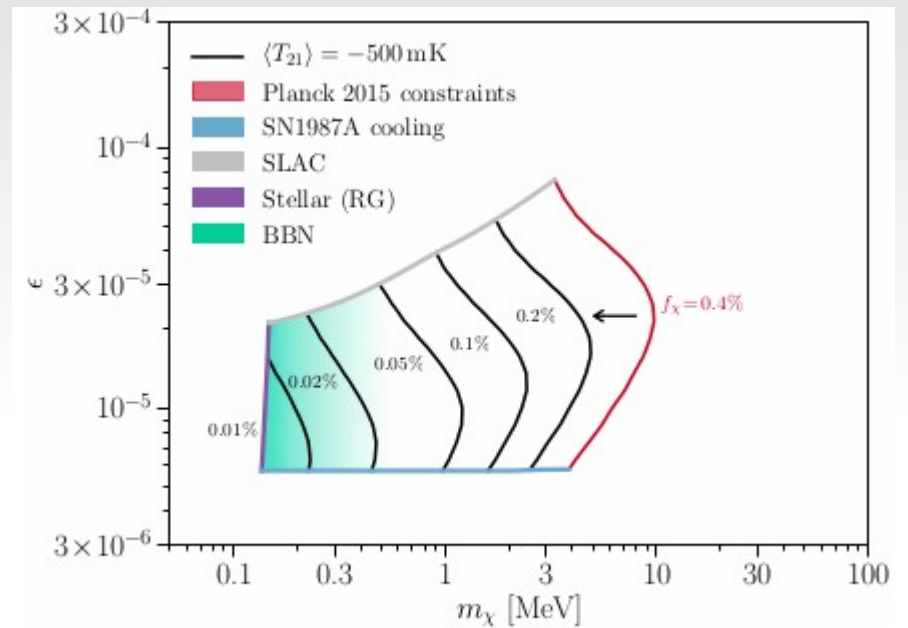
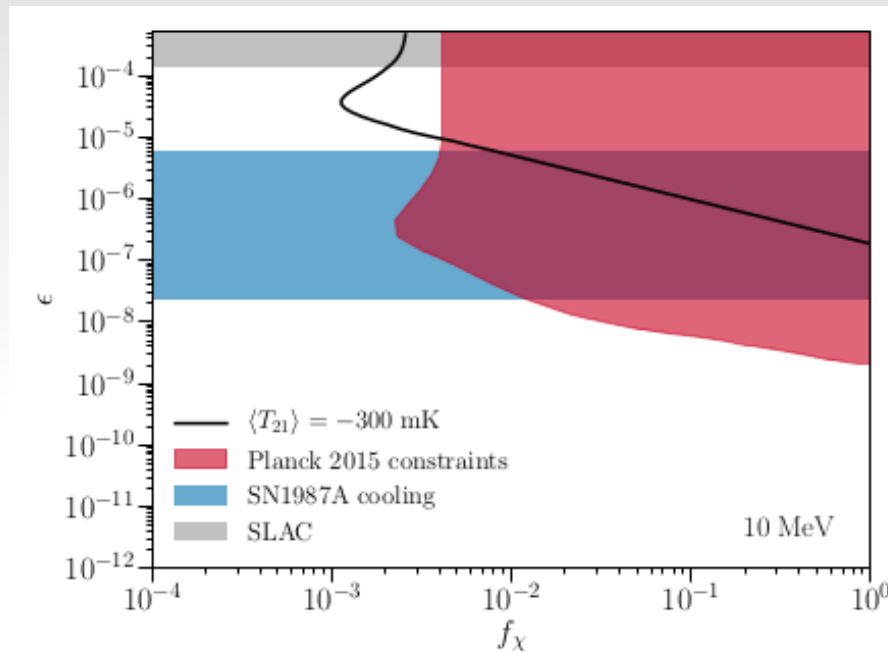
Cooling H_I gas needs a very large scattering cross-section need to avoid constraint from direct detection

→ velocity dependence?

R. Barkana, et.al. 1803.03091



Millicharged DM need to be at percent level abundance, and relic density requires alternative annihilation / depletion process A.Berlin, et.al. 1803.02804



Milli-charged DM constrained to lighter than 85 MeV
and tiny fractions of relic density
for 100% relic density, σ_0 constrained to $< 10^{-42} \text{cm}^2$.

E.D.Kovertz, et.al. 1807.11482

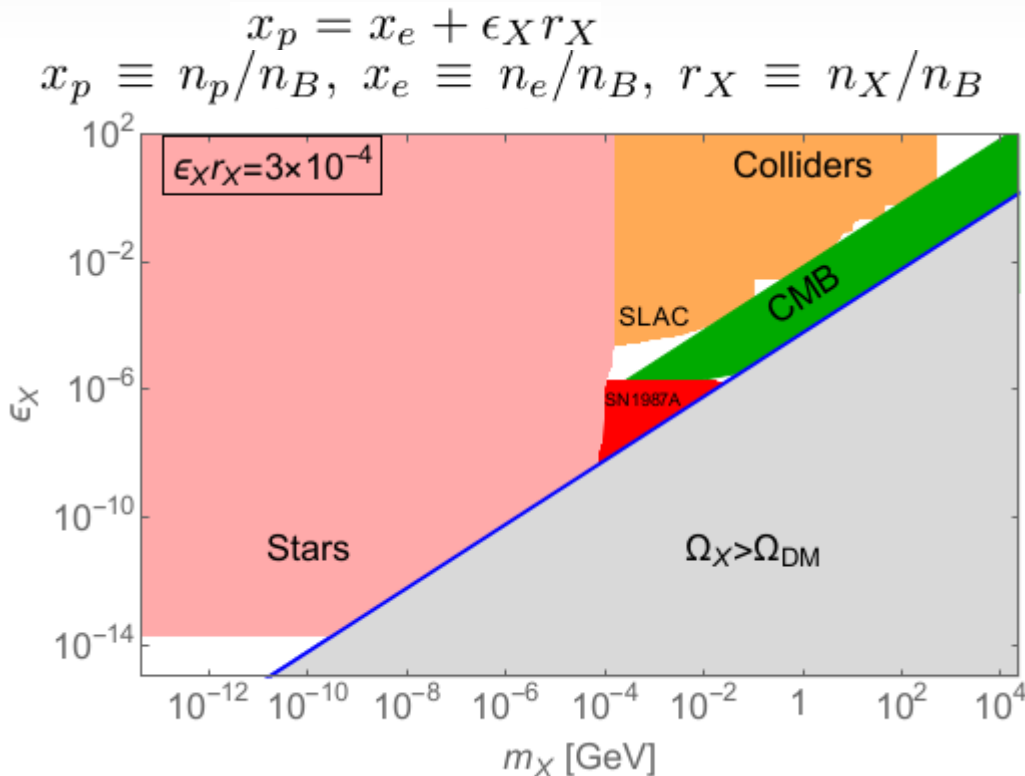
A few non-DM scenarios: Earlier $T_G - T_{\text{CMB}}$ decoupling

Charge sequestration:

A. Falkowski, K. Petraki, 1803.10096

A negatively charged ($-\epsilon$), stable particle to replace some electrons in the Universe

- * reduces x_e during recombination
- * fasten CMB – HI decoupling

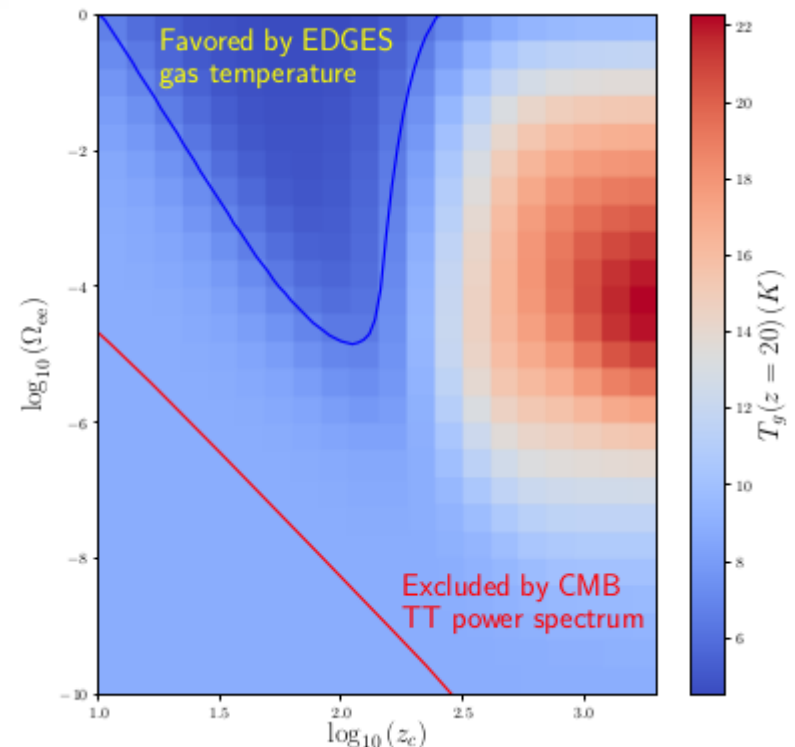


'Early' dark energy?

J. Colin Hill, E. Baxter. 1803.07555

Addition dark energy component with $w=-1$ and it decays away by $z_c \sim 20-1000$.

- * faster expansion rate 'early on'
- * Earlier CMB – HI decoupling



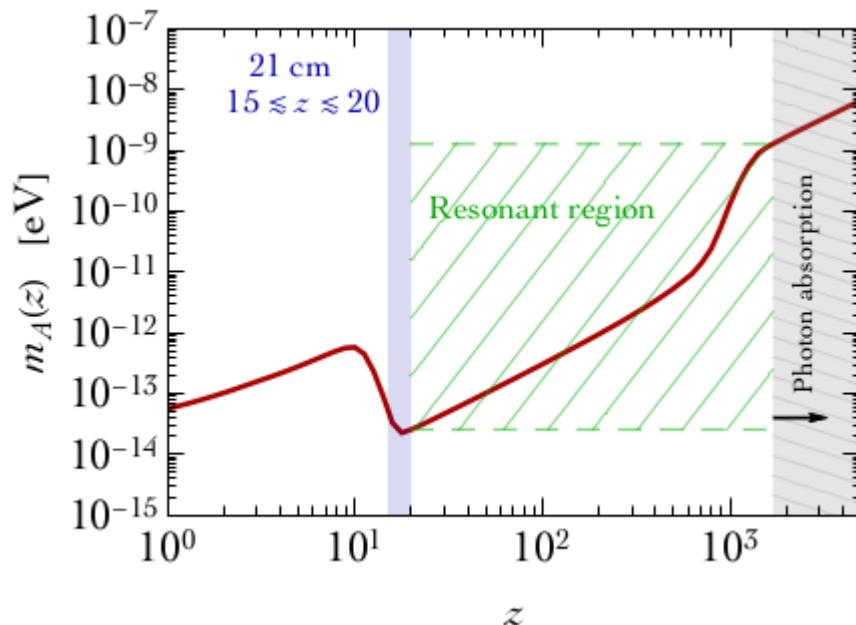
A few non-DM scenarios: low-freq. uncertainties

Modification in CMB:

M. Pospelov, et.al. 1803.07048

Order~1 increase in the Raileigh-Jeans tail of CMB: e.g. osc. with a very light dark photon A' via mixing $\epsilon F'_{\mu\nu} F_{\mu\nu}$

Can be of interest in precision 21cm tests, when A' - A oscillate resonantly at the effective A mass (in plasma) $m_A \sim m_{A'}$

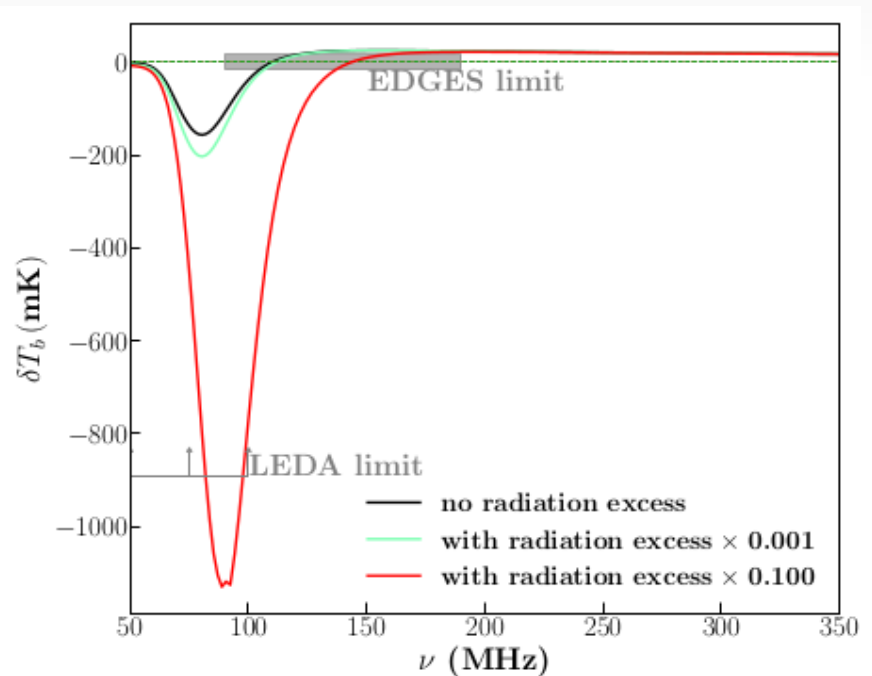


Radio-wavelength backgrounds

C. Feng, G. Holder. 1801.05396

other than CMB is detected by ARCADE

- non black body CMB, enhancing the signal
- EDGES can place a bound on early rad. fields



Other exotics..

- New physics contribution to CMB uncertainty

ALP conversion to CMB: [T.Moroi, K.Nakayama, Y.Tong, 1804.10378](#)

Mirror sector radiation to CMB: [D.A.Sierra, C.S.Fong, 1805.02685](#)

- Modified cosmology / Friedmann equations

Vacuum Energy – Dark Matter interaction: [Y.Wang G-B. Zhao, 1805.11210](#)

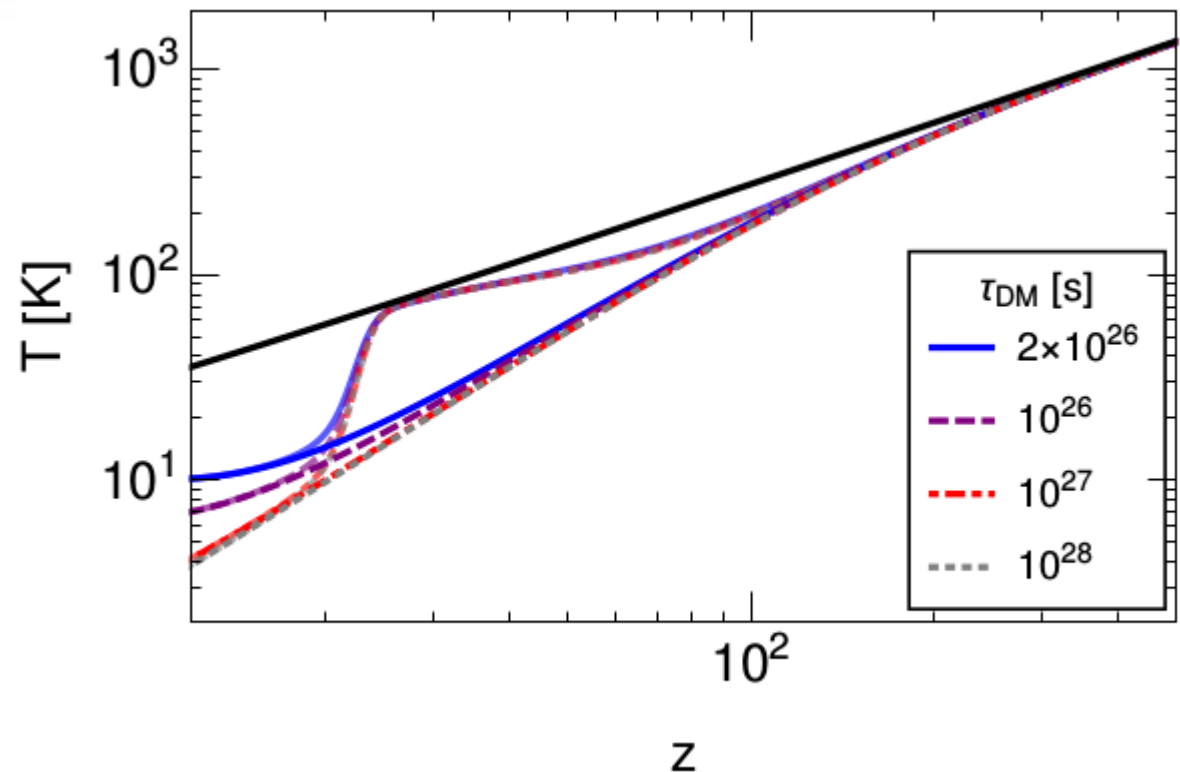
Dark Energy – Dark Matter interaction: [B.Yue, Y.Xu, B.Wang, 1807.05541](#)

DM heating story:
use the 21cm signal as a constraint

- Injection of energy into intergalactic gas suppresses 21cm during reionization epoch
- The observation of (any) 21cm signal means an upper limit
- EDGES may have a lot systematics, future 21cm experiments awaits.

Energy injection effects

- (Historic, cumulative) high-energy injection of electrons/photons
- Increased x_e
- Later CMB – gas decoupling
- Higher T_G
- May reduce or wipe out the 21cm signal



Injection & absorption

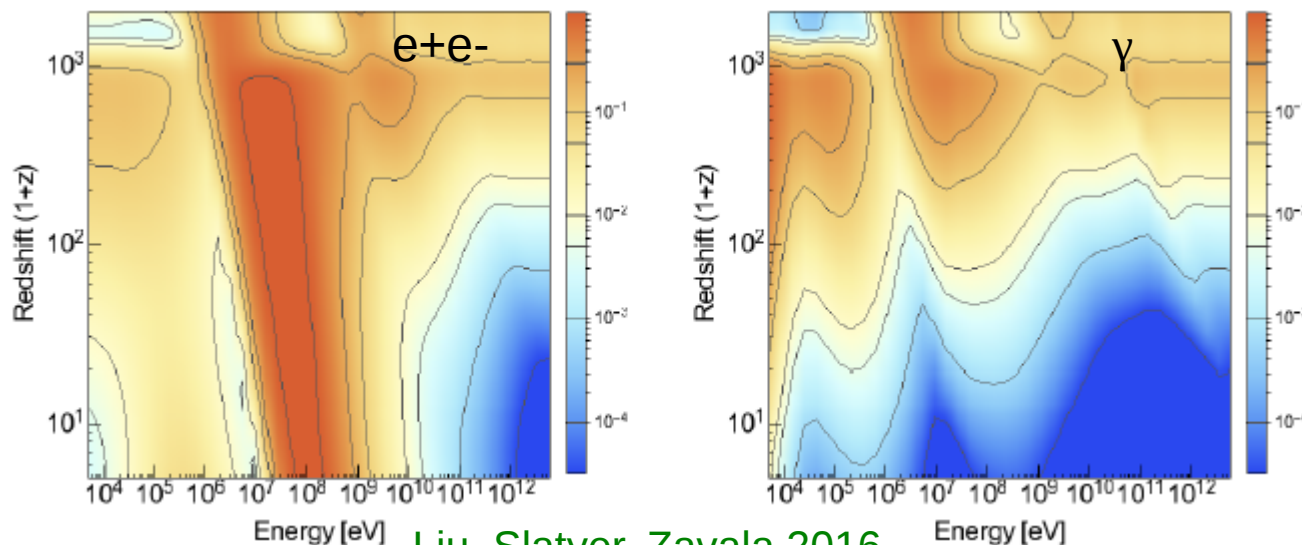
- Injected high-energy particles lose energy by scattering, ionization, excitations, etc...

Not all energy is immediately deposited into the environment (gas, CMB, etc) if particles are too energetic:

- * accumulative over earlier injection
- * efficiency reduces at later time

Numerical calculation

Energy “fraction” into ionization (of H)

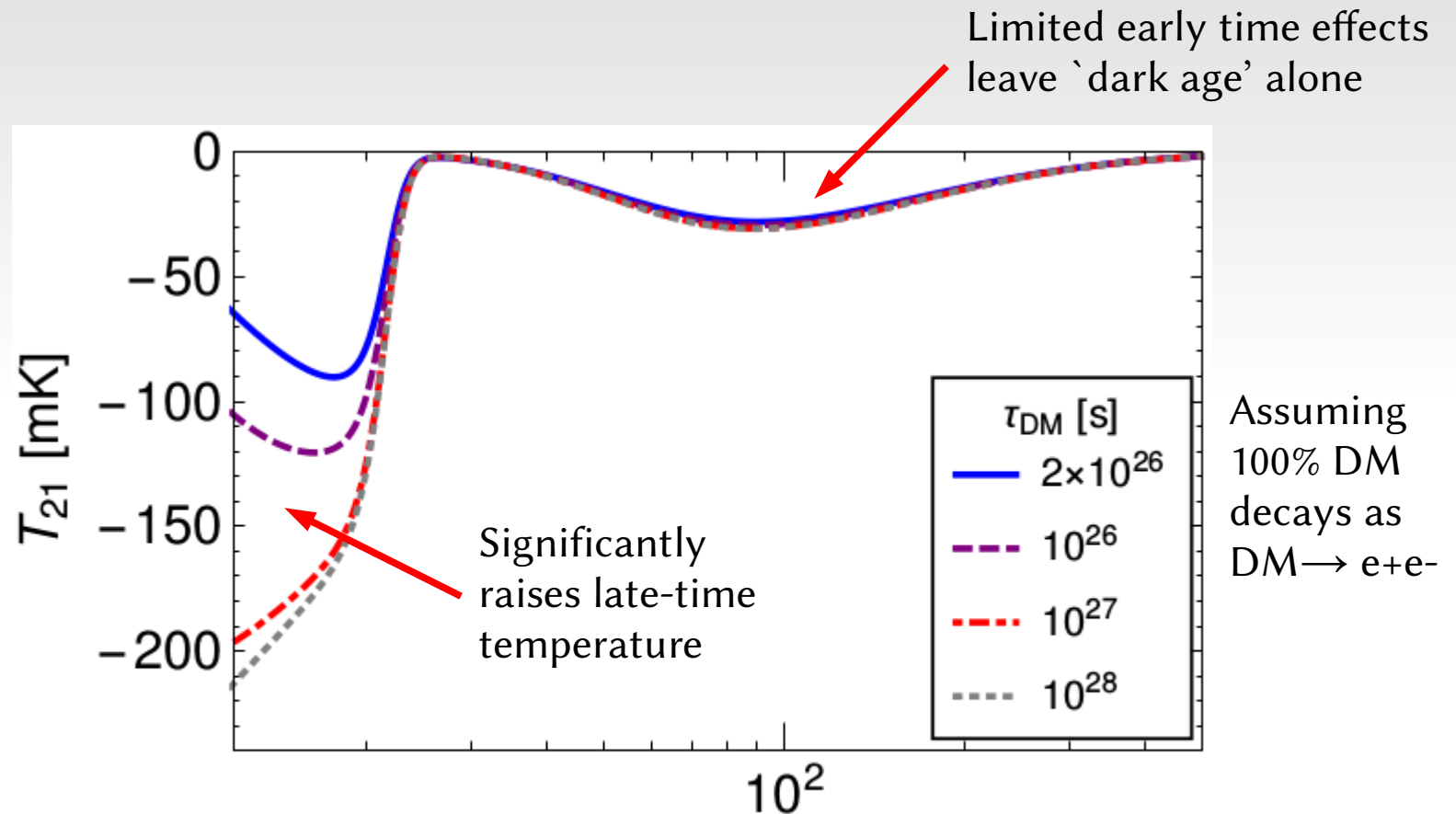


Liu, Slatyer, Zavala 2016

Implementation into
CAMB, HyRec codes:
new physics excitation,
scattering terms, Lyman- α photons, etc.

Also see:
Belotsky, Kirillov 2015

21cm signal suppression by injection



A 100~200 mK temperature enhancement will likely erase the expected 21cm signal in standard astrophysics

EDGES T_{21} uncertainty:
+200 mK by 95% up-fluctuation

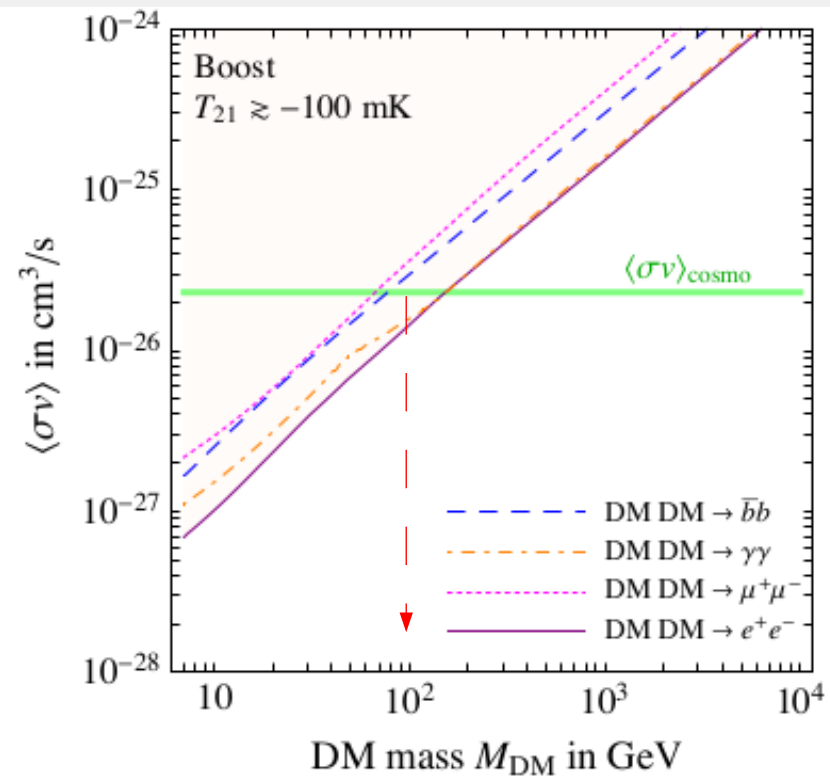
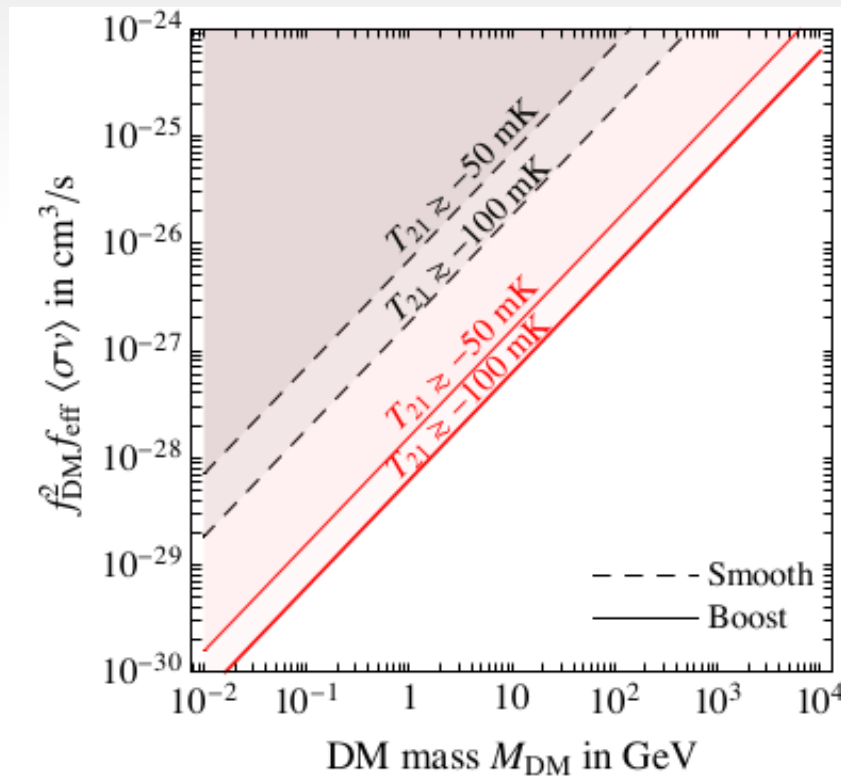
The observation of 21cm signal as an upper limit

On DM annihilation rates:

by requiring injection induced G. D'Amico, P. Panci, A. Strumia 1803.03629

$\Delta T_{21} < +100$ or $+150$ mK

Excluding vanilla thermal wimp below 10^2 GeV?

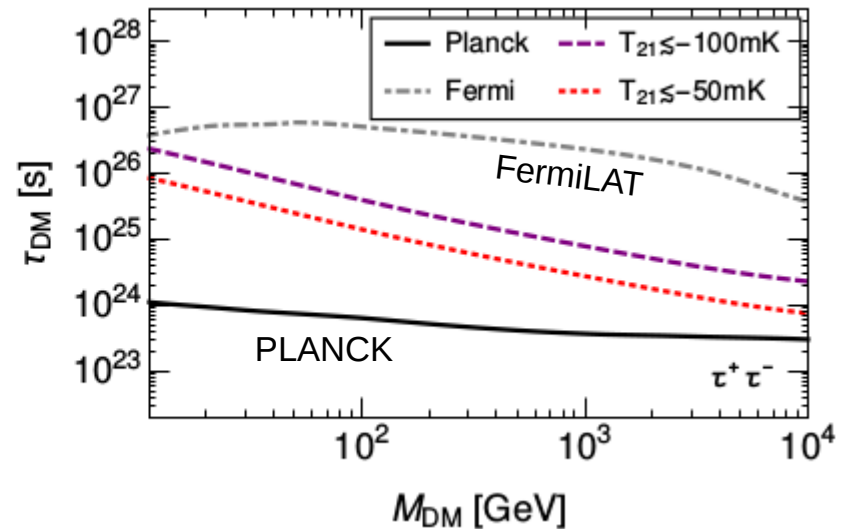
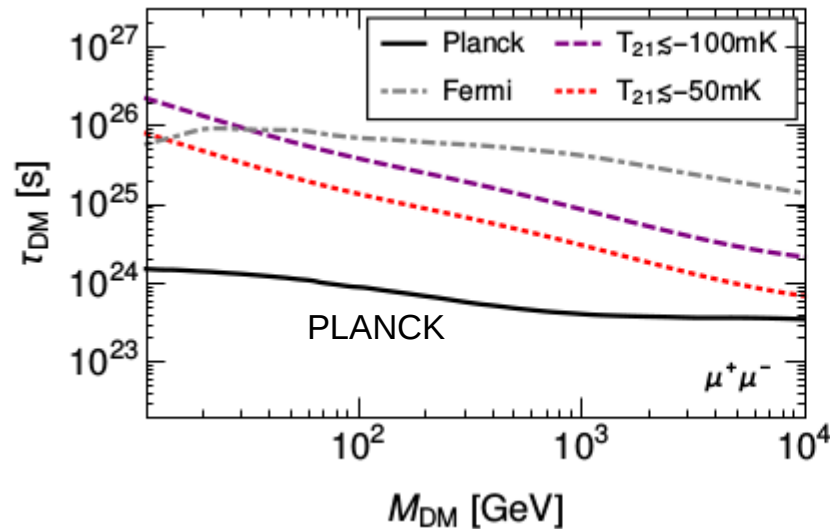
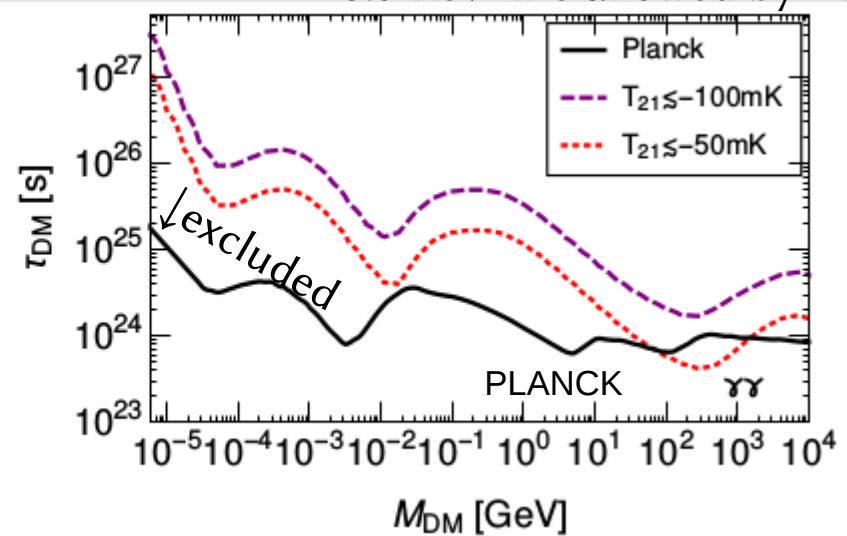
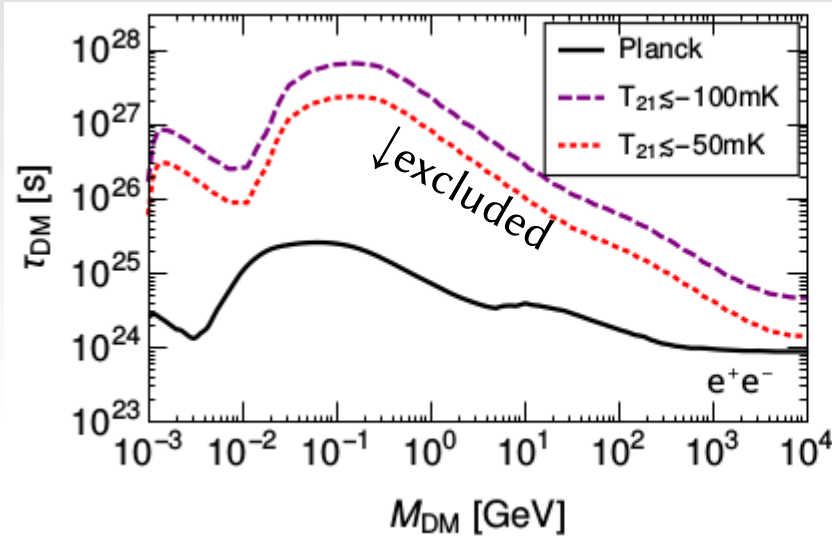


Annihilation: over-density boost B~217 at $z \sim 20$

Lower limit on DM decay lifetime

$$\Delta T_{21} < +100 \text{ or } +150 \text{ mK at } z=17$$

* \leftarrow 3.5 KeV line allowed by 21cm



Lower limit on DM decay lifetime

$\Delta T_{21} < +100$ or $+150$ mK at $z=17$

Gives 1-2 orders of magnitude better bounds on DM injection in comparison to PLANCK TT+TE+EE+lowP data

Lower sensitivity at large DM mass due to poorer effective energy loss efficiency

O(0.1-1) correction from cosmological parameter variations (PLANCK)

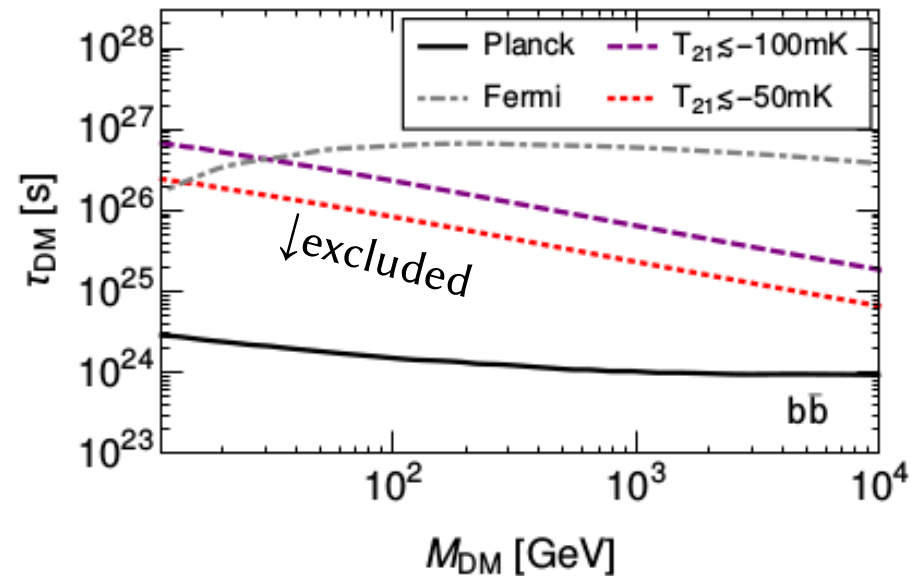
S. Clark, B. Dutta, Y. Gao, Y.-Z. Ma, L.E.Strigari, 1803.09390

Photon line signals:

Very good limit in KeV range. Testing the 3.5 KeV line needs O(mK) T_{21} sensitivity

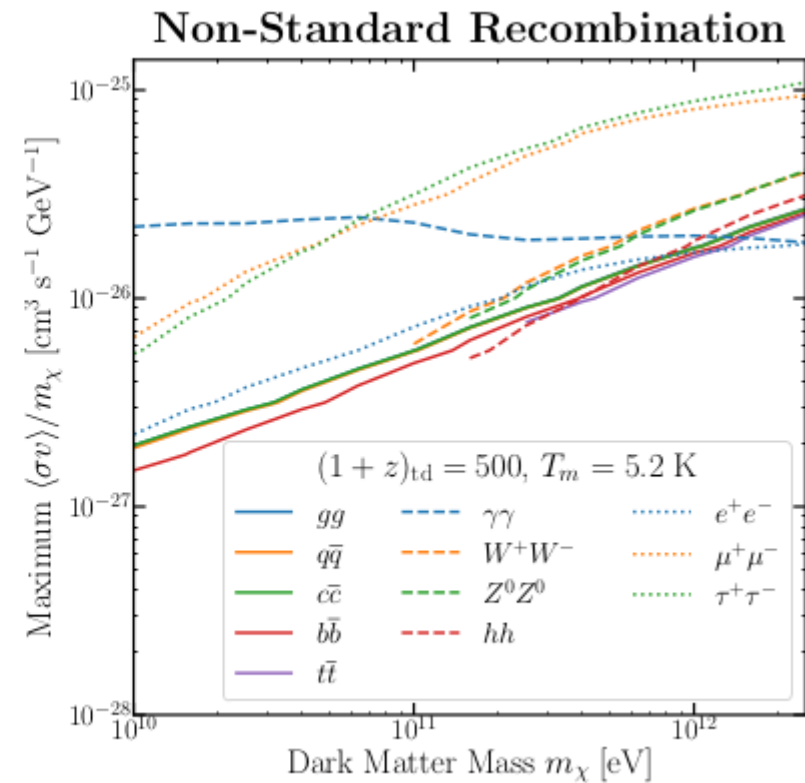
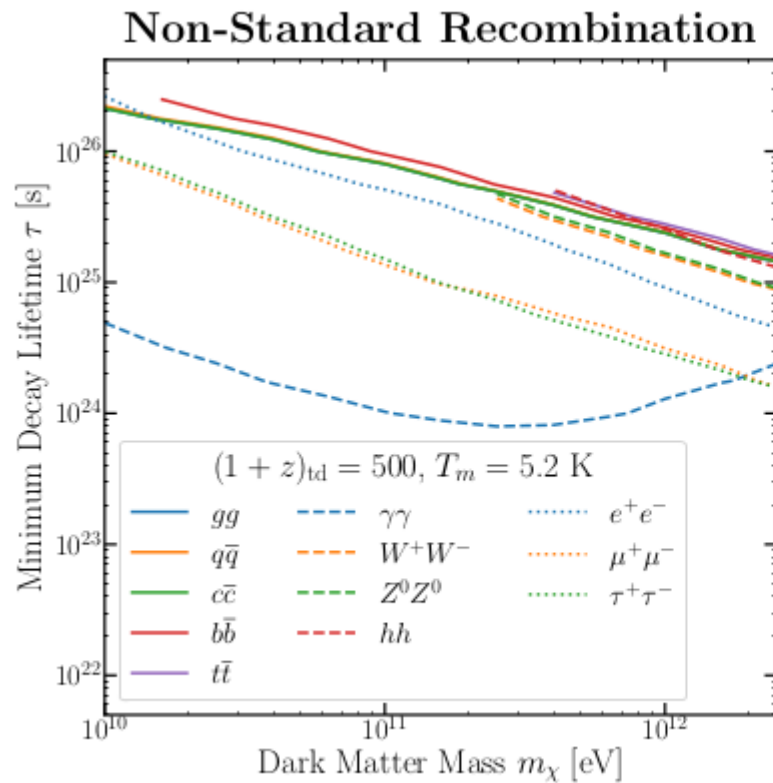
Comparable to X-ray line search $\sim 10^{26}$ s in (sub) MeV range.

Less than Fermi-LAT's pass8 $\sim 10^{29}$ s in GeV range.



Limits on DM annihilation & decays

$T_G(z=17.2) = 5.2\text{K}$ limit, no cooling

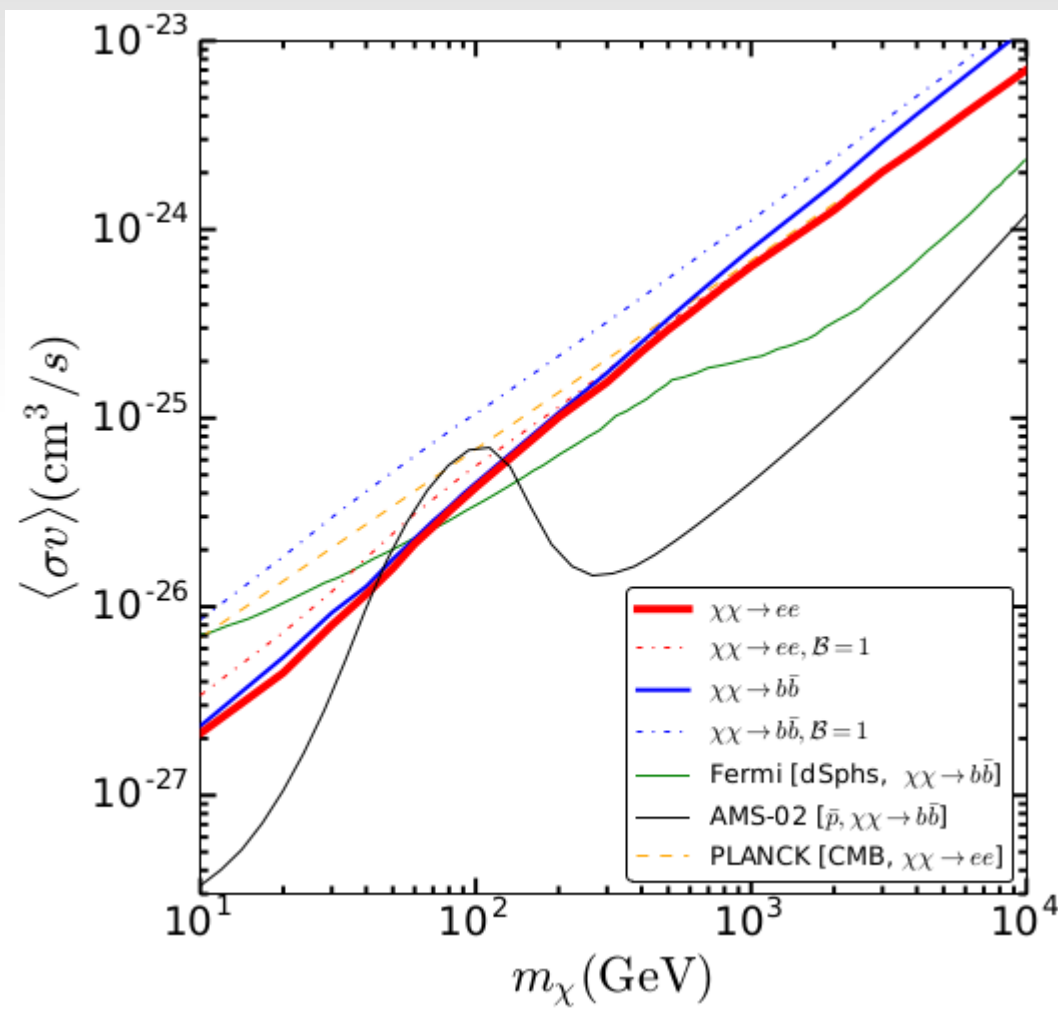
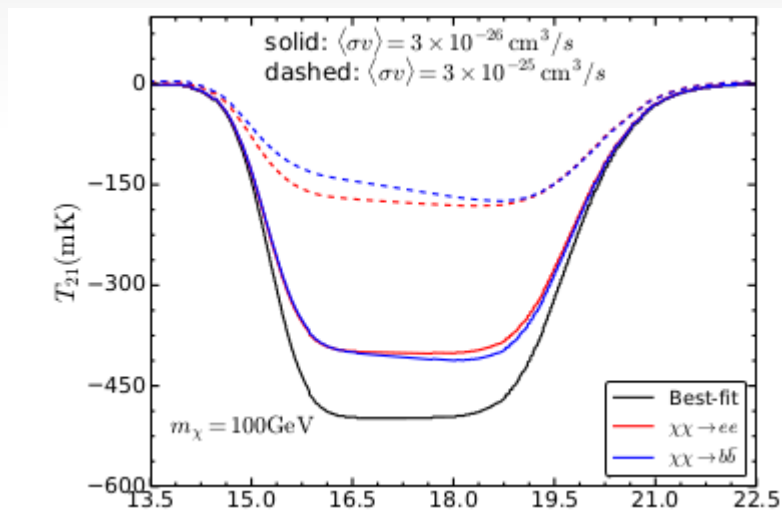


H. Liu, T. R. Slatyer, 1803.09739

Upper limit on DM annihilation

Fitting to the EDGES T_{21} history..

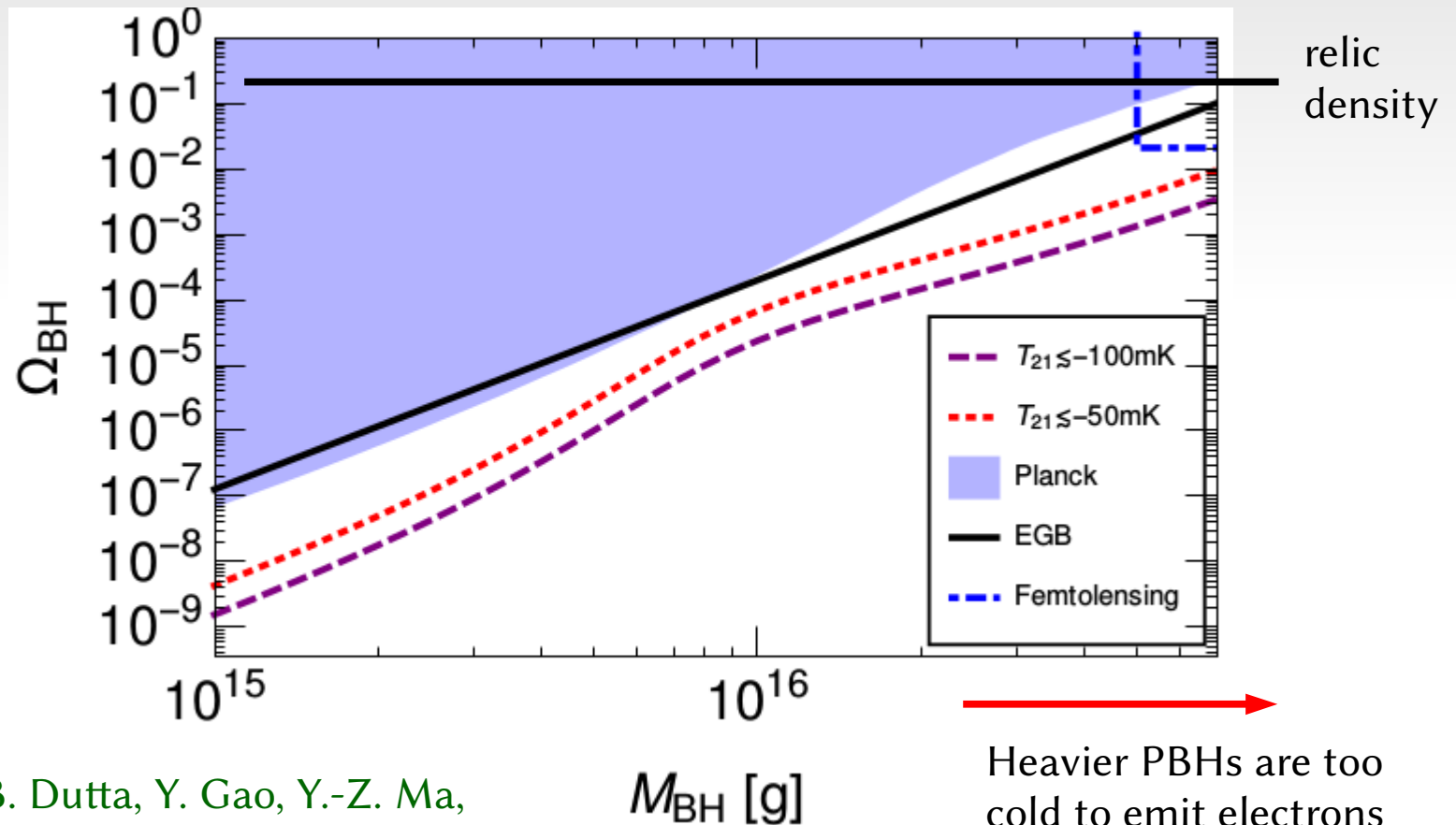
K. Cheung, J-L Kuo, K-W
Ng, Y-L S Tsai, 1803.09398



Upper limit on primordial BH's Hawking radiation

$\Delta T_{21} < +100$ or $+150$ mK at $z=17$

Applicable to long-lived ($m > 10^{15}$ g), evaporating black holes



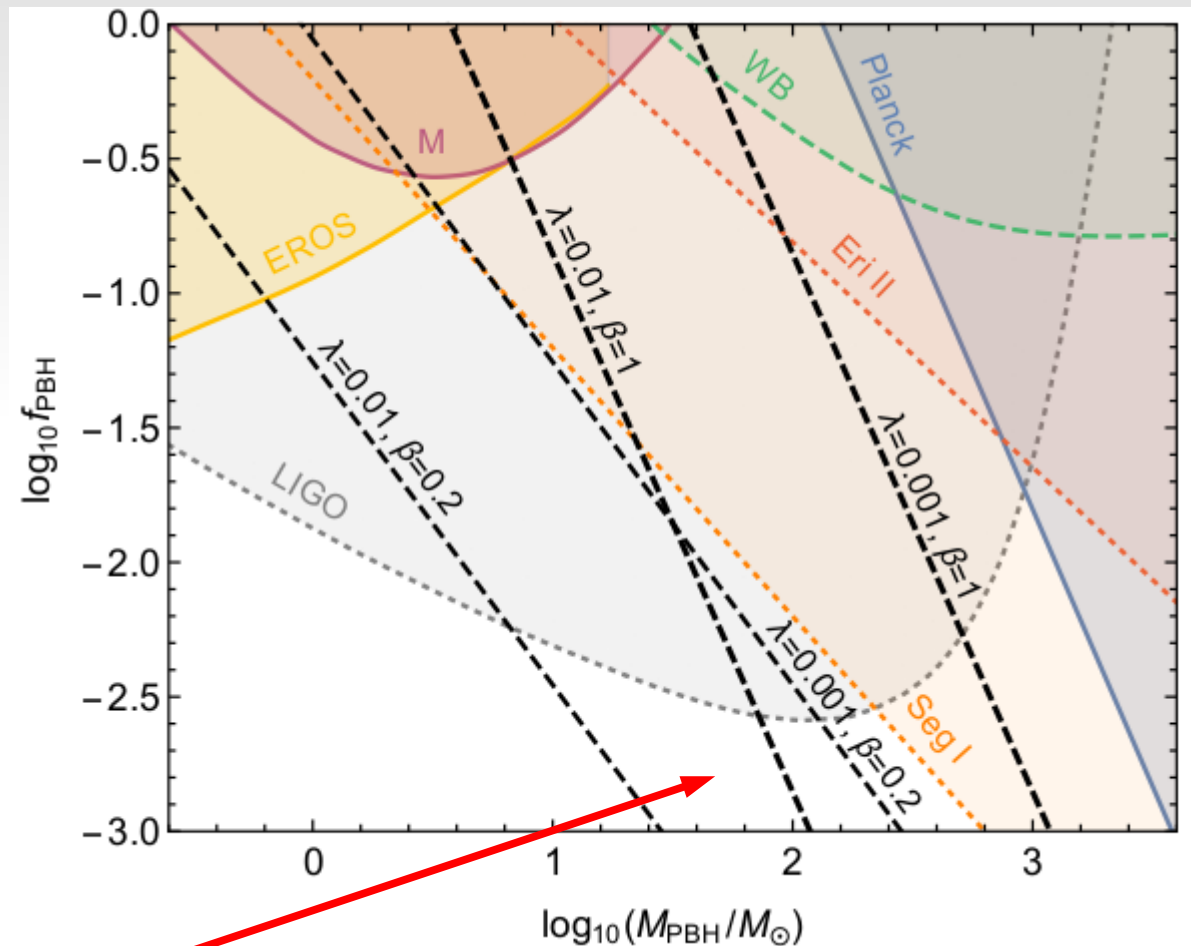
S. Clark, B. Dutta, Y. Gao, Y.-Z. Ma,
L.E.Strigari, 1803.09390

Upper limit on massive PBH's accretion

X-ray emission from BH accretion, immediate energy deposition

Limit set by $T_c(z \sim 17.2) = 8K$
Improvement over PLANCK by 1-2 orders of mag.

A. Hektor, G. Hütsi, L. Marzola,
M. Raidal, V. Vaskonen,
H. Veermäe, 1803.09697



Complementing the LIGO

Summary

- EDGES results can be interpreted in many ways
- Existence of reionization epoch 21cm signal imposes a strong limit on historic energy injection
- Like CMB, injection bounds extend to low energy/mass
- 21cm bounds on DM, PBH are very powerful, even at the current proof-of-principle estimates
- Future 21cm experiments await.

Backup: Injection corrections

Extra contributions to ionization & heating

- Reduce neutral hydrogen fraction

$$\frac{dx_e}{dz} = \left(\frac{dx_e}{dz} \right)_{\text{orig}} - \frac{1}{(1+z)H(z)} (I_{Xi}(z) + I_{X\alpha}(z))$$

Ionization: $I_{Xi}(z) = f_i(E, z) \frac{dE/dV dt}{n_H(z) E_i}$

- IGM temperature

Lyman- α excitation $I_{X\alpha}(z) = f_\alpha(E, z)(1 - C) \frac{dE/dV dt}{n_H(z) E_\alpha}$

$$\frac{dT_{\text{IGM}}}{dz} = \left(\frac{dT_{\text{IGM}}}{dz} \right)_{\text{orig}} - \frac{2}{3k_B(1+z)H(z)} \frac{K_h}{1 + f_{\text{He}} + x_e}$$

- Wouthuysen-Field

$$K_h(z) = f_h(E, z) \frac{dE/dV dt}{n_H(z)}$$

$$T_S = \frac{T_{\text{CMB}} + y_c T_G + y_{\text{Ly}\alpha} T_{\text{Ly}\alpha}}{1 + y_c + y_{\text{Ly}\alpha}},$$

$$y_c = \frac{C_{10}}{A_{10}} \frac{T_\star}{T_G},$$

$$y_{\text{Ly}\alpha} = \frac{P_{10}}{A_{10}} \frac{T_\star}{T_{\text{Ly}\alpha}},$$

For Lyman-alpha during reionization,
See B. Ciardi and P. Madau, astro-ph/0303249

- Energy deposit rate (ionization, excitations, heating)

$$I_{X_i}(z) = \frac{f_i(E, z)}{H_H(z)E_i} \frac{dE}{dVdt},$$

$$I_{X_\alpha}(z) = (1 - C) \frac{f_\alpha(E, z)}{n_H(z)E_\alpha} \frac{dE}{dVdt},$$

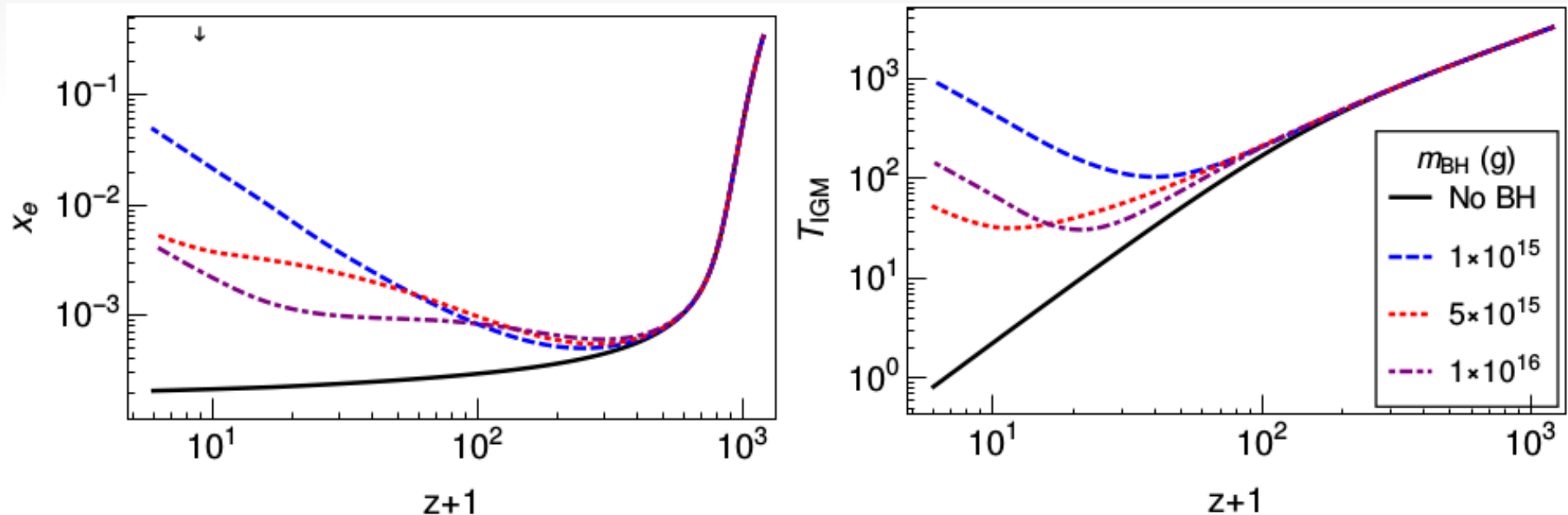
$$K_h(z) = \frac{f_h(E, z)}{n_H(z)} \frac{dE}{dVdt}$$

$$C = \frac{1 + K\Lambda_{2s,1s}n_H(1 + x_e)}{1 + K\Lambda_{2s,1s}n_H(1 - x_e) + K\beta_B n_H(1 - x_e)}.$$

- The ‘effective’ rate $f(E, z)$ is cumulative of historic injection
- Electrons are more effective in energy deposit
- f is averaged over injection spectra and species

$$f_c(m_{DM}, z) = \frac{\sum_s \int f_c(E, z, s) E (dN/dE)_s dE}{\sum_s \int E (dN/dE)_s dE},$$

- Ionization history can be a powerful test (at $z > 7$) on 'prolonged' electron & gamma ray injection from new physics



Num. computation by HyRec

Backup: PBHs (CMB)

Hawking radiation, 'lifetime' and BH mass

- BH evaporates at a temperature

Hawking 75'

$$T_{PBH} = \frac{1}{8\pi GM} = 1.06 \text{TeV} \times \frac{10^{10} \text{g}}{M_{PBH}}$$

- with a peak energy of radiation

$$E_\gamma = 5.71 T_{PBH}, \quad E_\nu = 4.22 T_{PBH}, \quad E_{e^\pm} = 4.18 T_{PBH}$$

Mass loss rate:

$$\dot{M}_{10} = -5.34 \times 10^{-5} \left(\sum_i f_i \right) M_{10}^{-2} \text{ s}^{-1}$$

lifetime:

$$\tau(M) \sim \frac{G^2 M^3}{\hbar c^4} \sim 10^{64} \left(\frac{M}{M_\odot} \right)^3 \text{ yr}$$

- BH evaporation can be a good source of cosmic rays, injection particle species determined by BH mass

A steady radiation injection below ~ 100 MeV

- Relevant for PBH mass above 10^{15} g, or peak radiation energy below muon mass
- Hawking evaporation after recombination yields (mostly) e^+e^- and gamma rays
- For $M \gg 10^{15}$ g, mass loss negligible during the age of the Universe
- A steady and long-lasting injection of radiation that scales as $(1+z)^3$

Extra-galactic source

Of Hawking evaporation rate

$$\dot{M}_{10} = -5.34 \times 10^{-5} \left(\sum_i f_i \right) M_{10}^{-2} \text{ s}^{-1}$$

into light (massless) species,

$$\begin{aligned} f_0 &= 0.267, & f_1^\gamma &= 0.06, & f_{3/2} &= 0.02, \\ f_2^g &= 0.007, & f_{1/2}^\nu &= 0.147, & f_{1/2}^{e^\pm} &= 0.142 \end{aligned}$$

J.MacGibbon, PRD, 1991

and photons & electrons affects
the environment
with unit volume
injection rate,
+redshift

$$\begin{aligned} \frac{dE}{dV dt} &= \dot{M}_{PBH} \eta(E_i, z) n_{PBH} \\ &= \frac{\dot{M}_{10}}{M_{10}} \rho_{cr}(z) \Omega_{PBH}(z) \eta_i(E, z) \end{aligned}$$

$$\frac{dE}{dV dt} \Big|_{BH} \neq \frac{\dot{M}_{10}}{M_{10}} \rho_{cr}(z) \Omega_{PBH}(z) \eta(E_{PBH}, z)$$

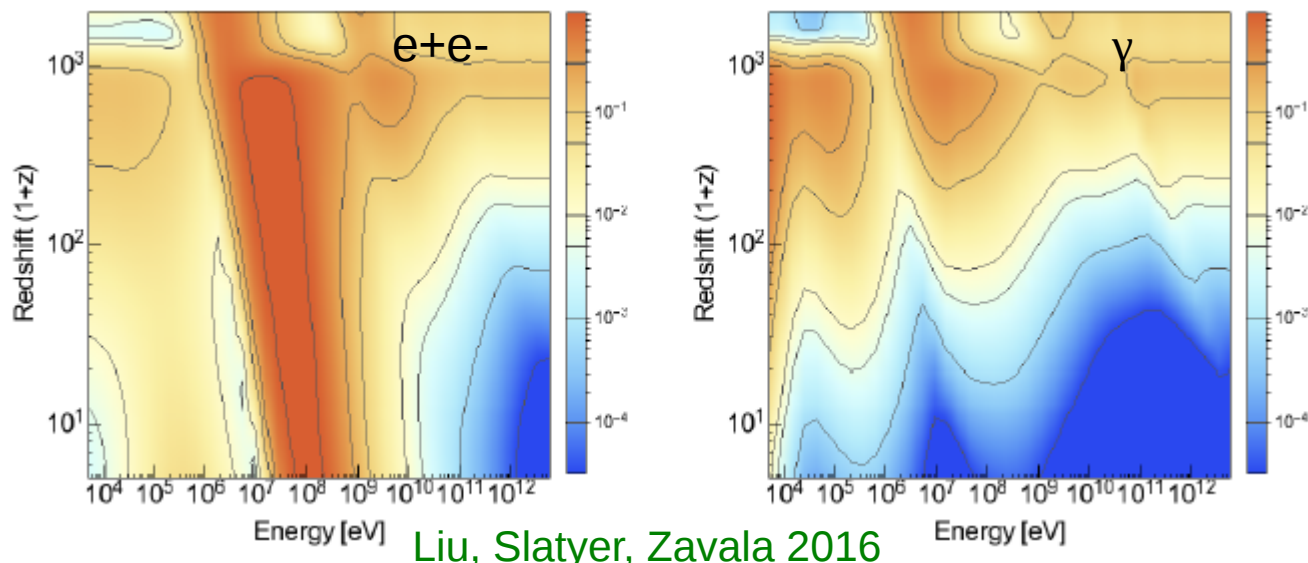
$\Gamma^{-1} \sim M/M$

Injection vs absorption

- photons interact via Compton scattering & absorptions
- electrons lose energy by inverse C. scattering & **ionization**

Not all energy is efficiently absorbed by the environment (gas) esp. if particles are too energetic

Energy “fraction” into ionization (of H)

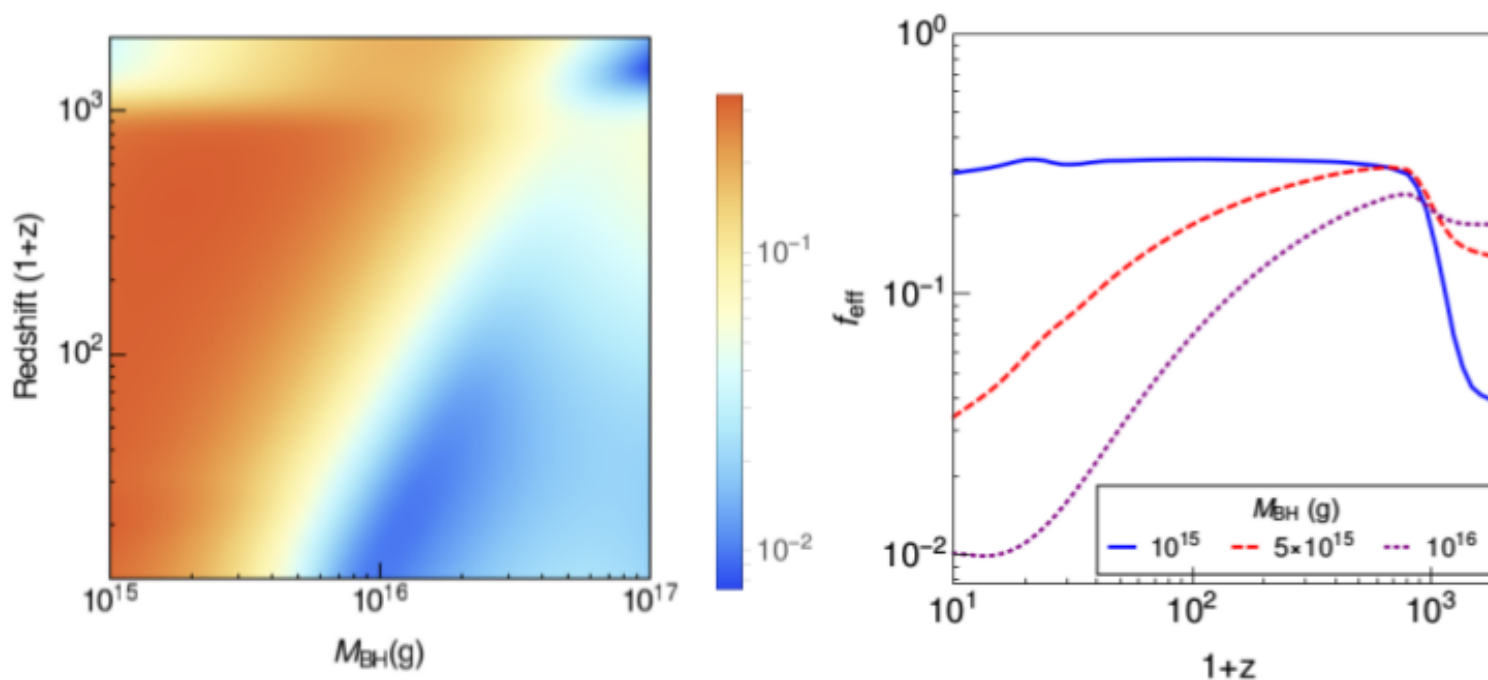


Liu, Slatyer, Zavala 2016

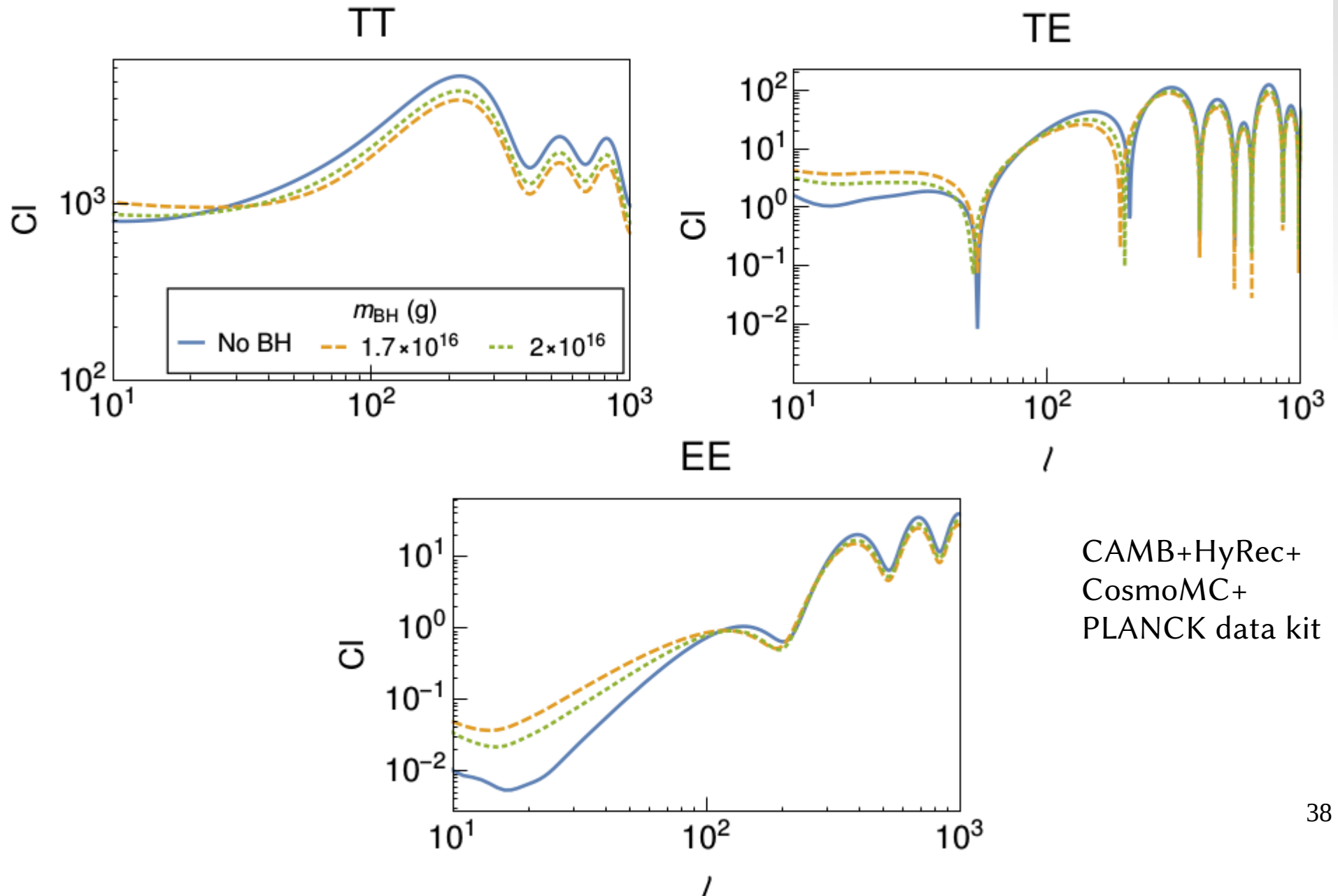
Also see:
Belotsky, Kirillov 2015

BH Effective absorption efficiencies

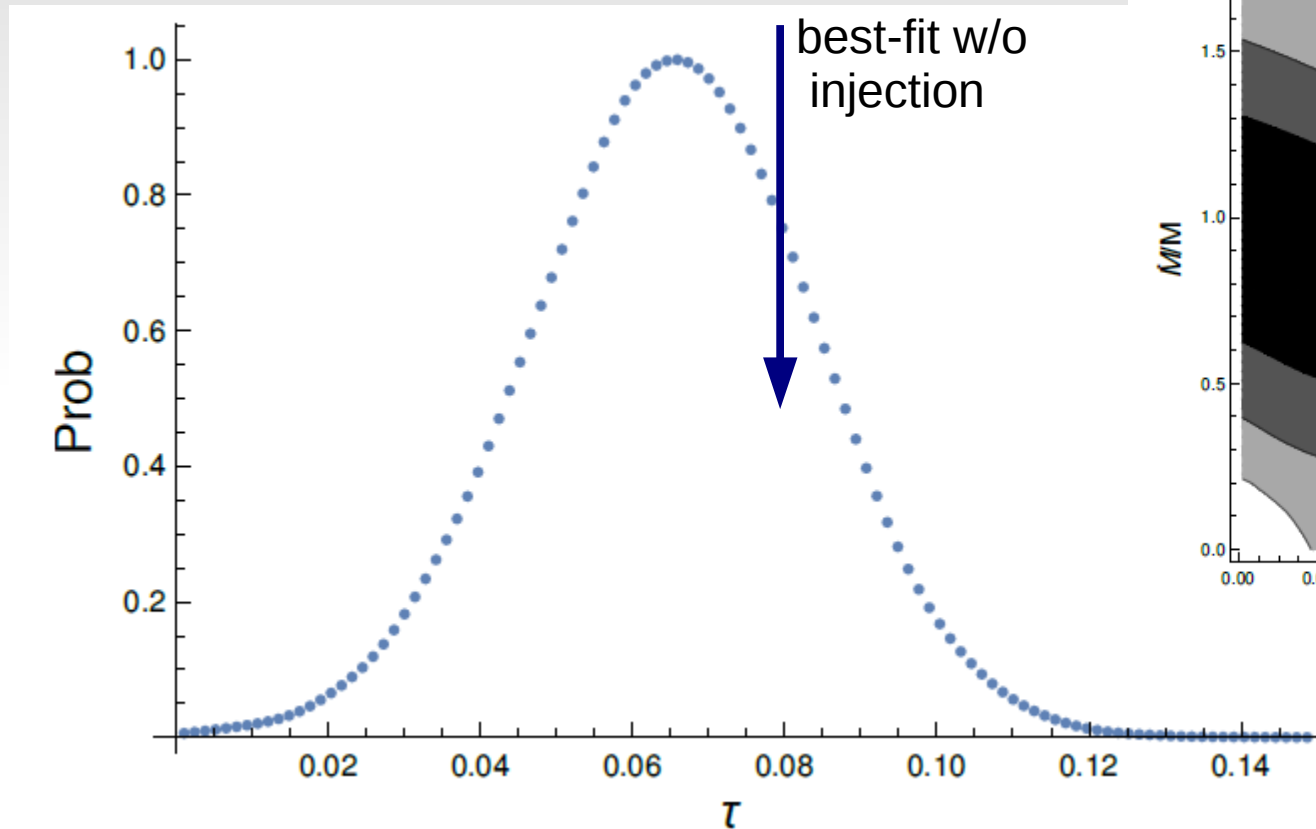
- Average over injection (BH radiation) energy spectrum and particle species
- Delayed injection: integrated over earlier z (up to CMB)



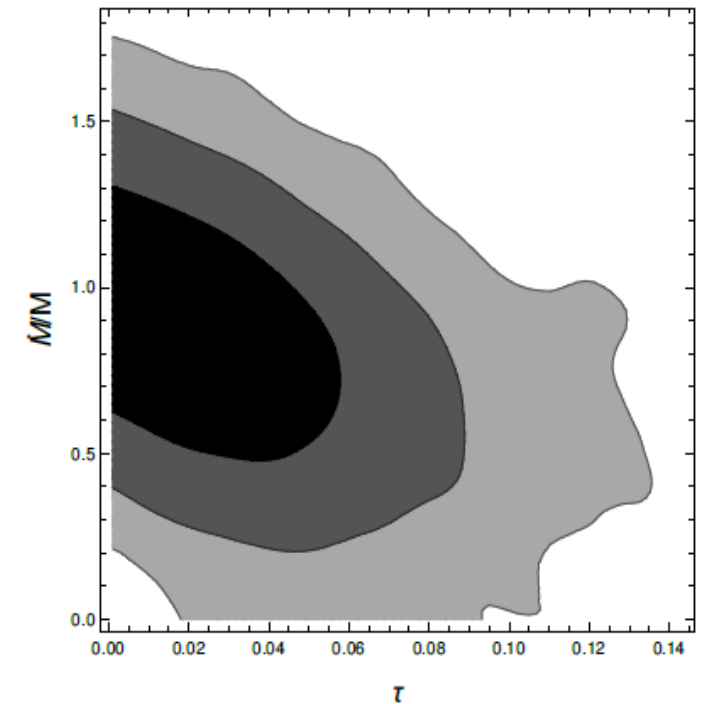
Impact on the CMB Cls



Reduced optical depth



PLANCK: temperature + polarization



Consistent with PLANCK
no injection fit within 2σ

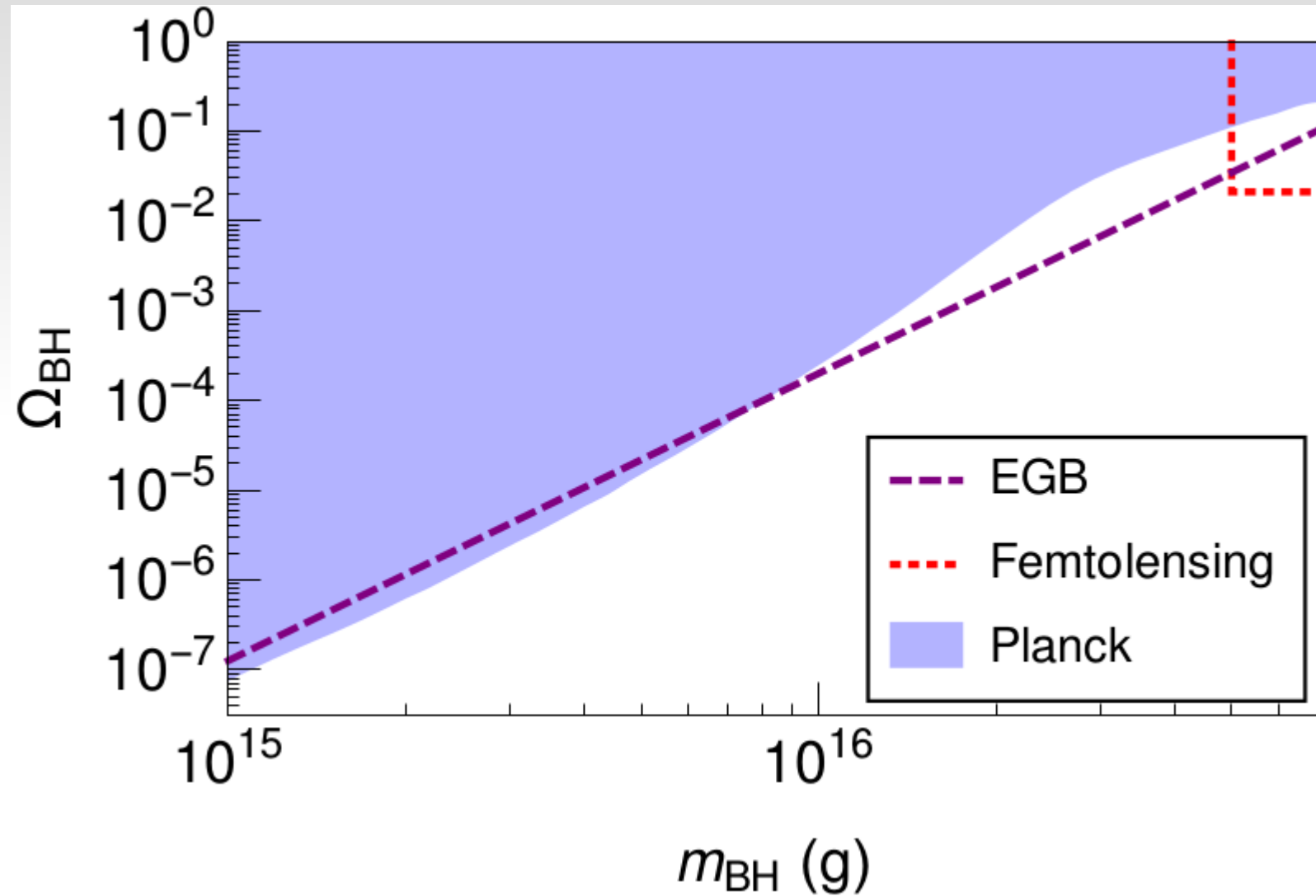
Consistent with reionization
 $z \leq 6$

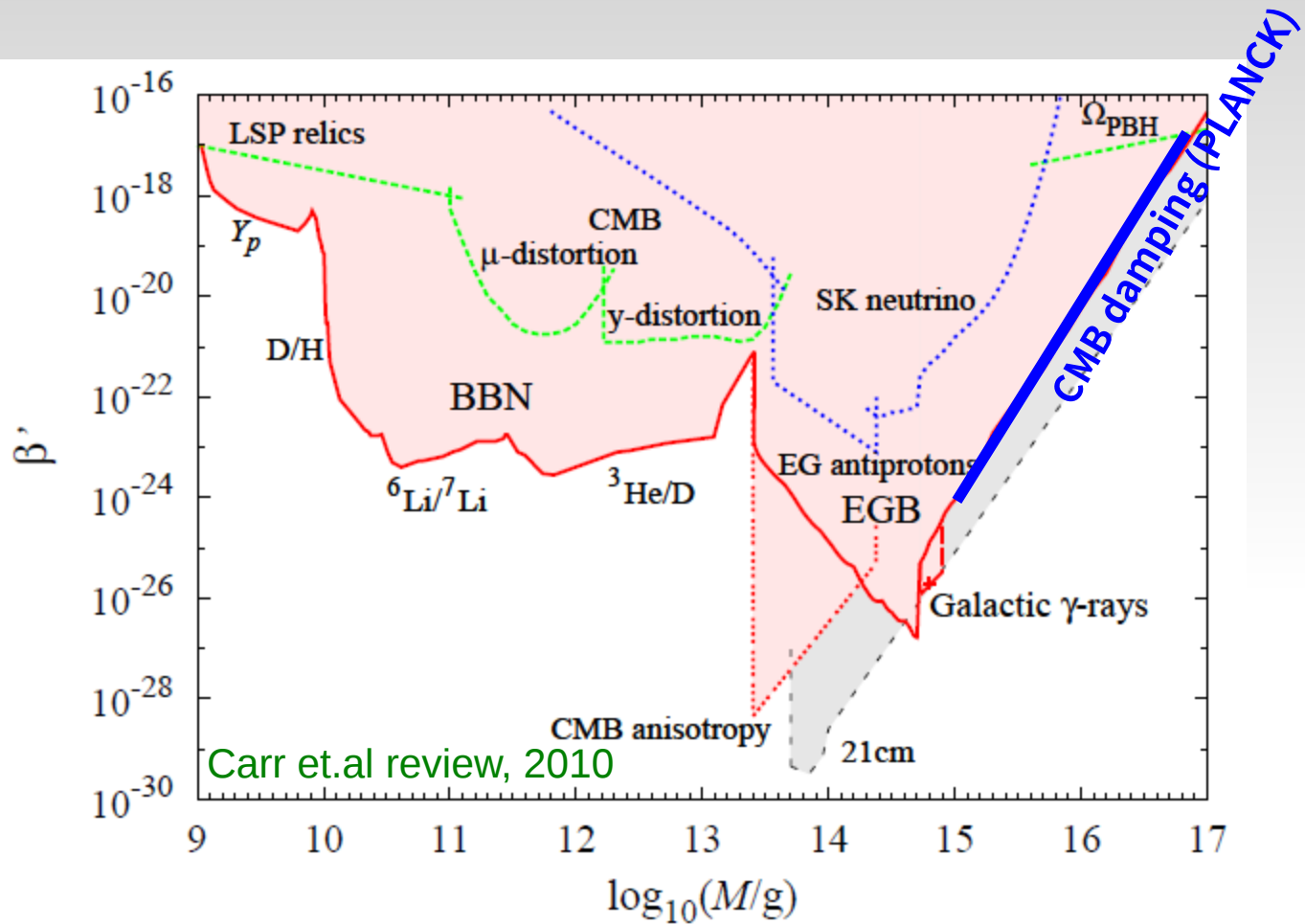
DM Decay: $\tau > 10^{24}$ s

Liu, Slatyer, Zavala 2016

Need polarization data to
break degeneracy

PLANCK constraint on PBH as 'relic' abundance



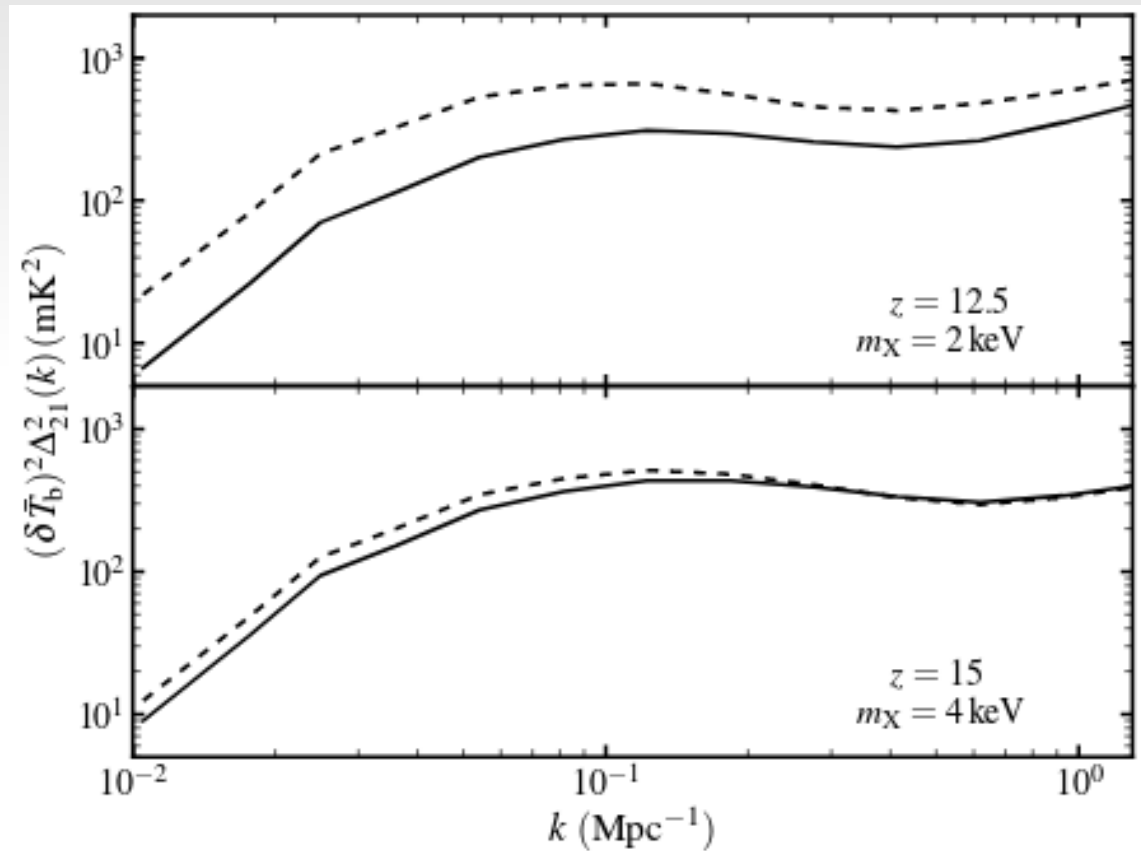


Simple scaling, assume no entropy production after PBHs and $\Omega_r \sim 10^{-4}$ at CMB time, also $\beta' \sim \beta$:

$$\Omega_{\text{PBH}} \simeq \beta \Omega_r (1+z) \sim 10^6 \beta \left(\frac{t}{1\text{s}} \right)^{-1/2} \sim 10^{18} \beta \left(\frac{M}{10^{15}\text{g}} \right)^{-1/2} \quad (M > 10^{15}\text{g}),$$

Backup: 21cm prospects

SKA: Reionization epoch power spectrum



SKA white paper

Warm versus Cold DM at 21cm

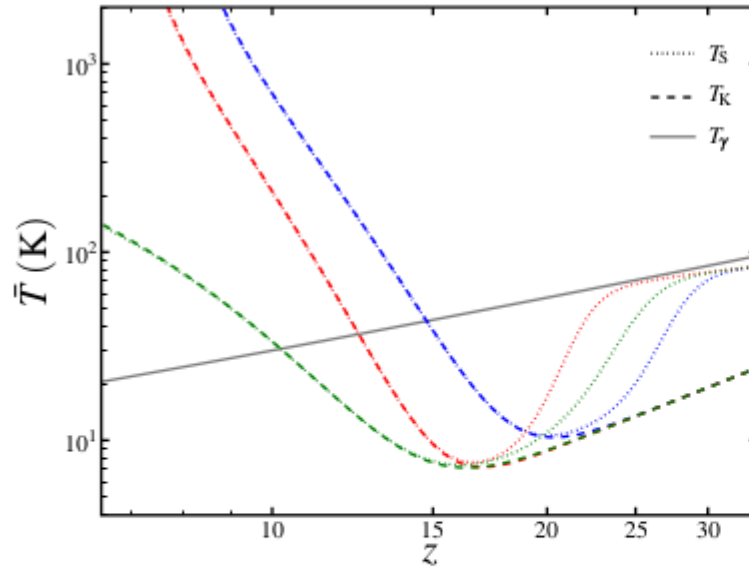


Figure 3. Mean spin temperatures \bar{T}_S for CDM and WDM models. The dotted curves show \bar{T}_S for our fiducial CDM model (blue), WDM with $m_\chi = 3$ keV (red), and CDM with $f_*/f_{*\text{fid}} = 0.1$ (green). In addition, the mean kinetic temperature \bar{T}_K of each model is plotted with a dashed curve in the same colour used for \bar{T}_S . The grey solid line is the CMB temperature. Figure taken from (Sitwell et al., 2014).

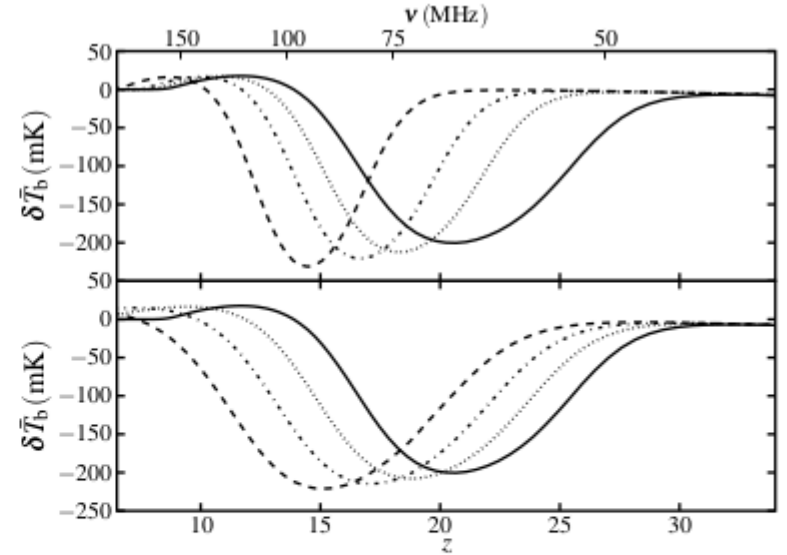
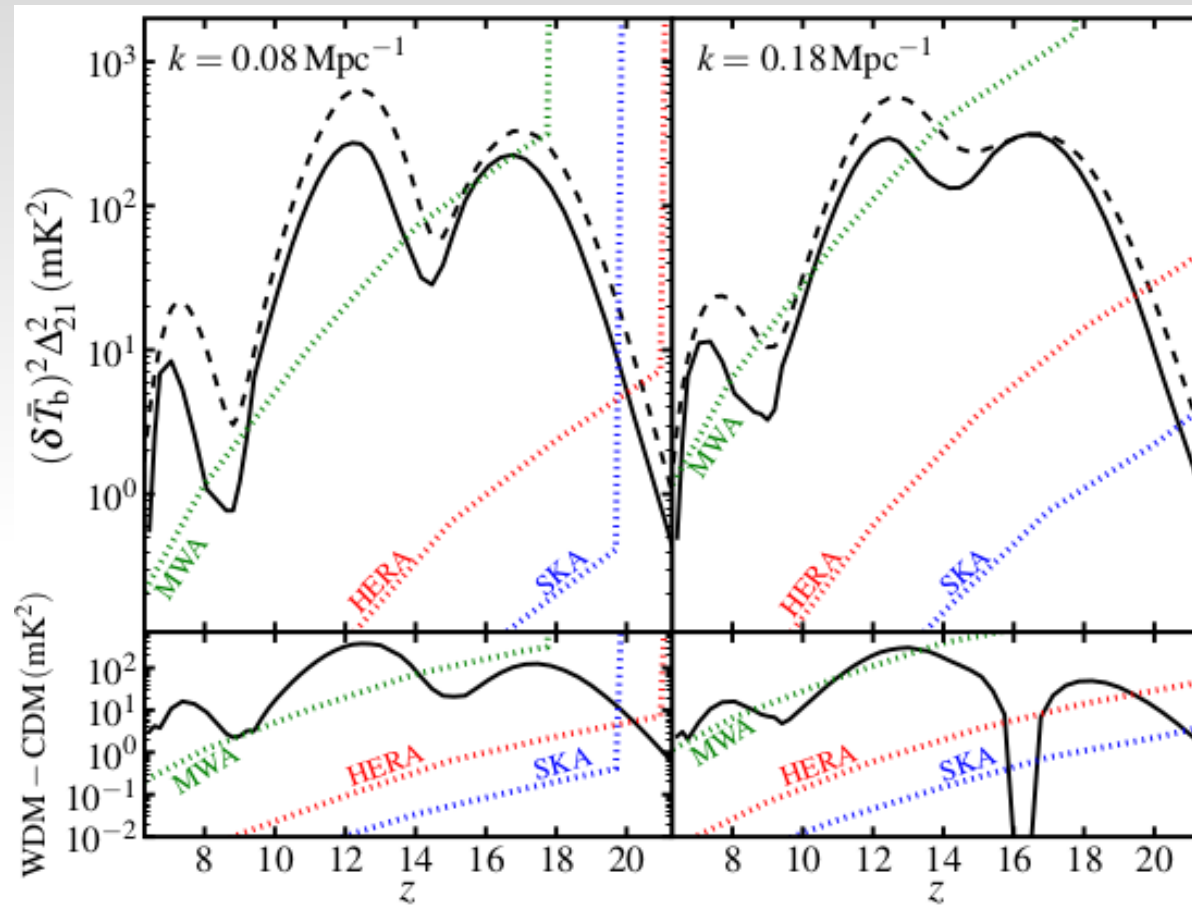


Figure 4. Mean 21 cm brightness temperature $\delta\bar{T}_b$. The solid curve is the fiducial CDM model. The upper plot shows the results of WDM runs where the dashed, dotted-dashed, and dotted curves are for $m_\chi = 2, 3, 4$ keV, respectively. The lower plot shows CDM runs where the dashed, dotted-dashed, and dotted curves are for CDM models with $f_*/f_{*\text{fid}} = 0.03, 0.1, 0.5$, respectively. Figure taken from Sitwell et al. (2014).



M. Sitwell, A. Mesinger, Y.-Z. Ma,
K. Sigurdson, 2014 MNRAS, 438, 2664