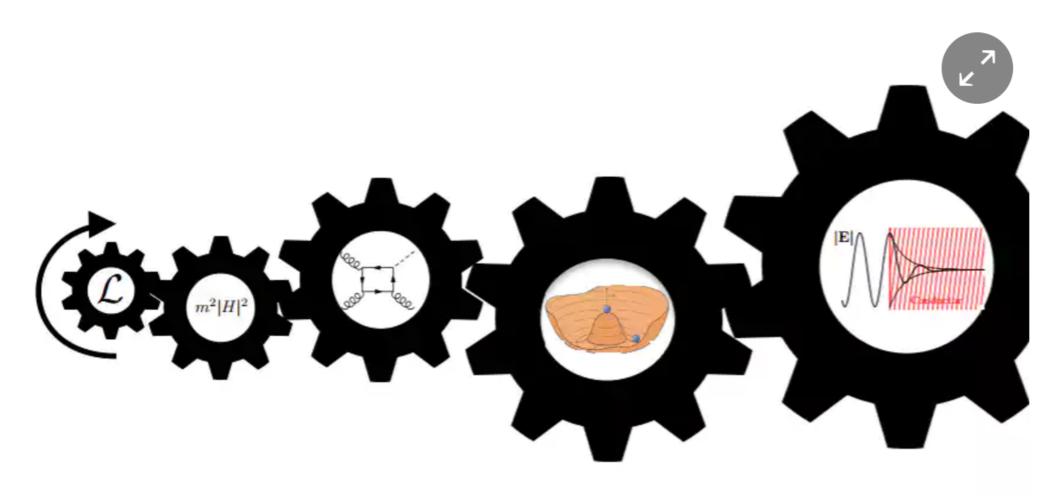
# Searching for a new world New Physics at the LHC and beyond

LianTao Wang U. Chicago

FeynRules/Madgraph School. Nov. 19, 2018. USTC HeFei.

# From gravity to the Higgs we're still waiting for new physics

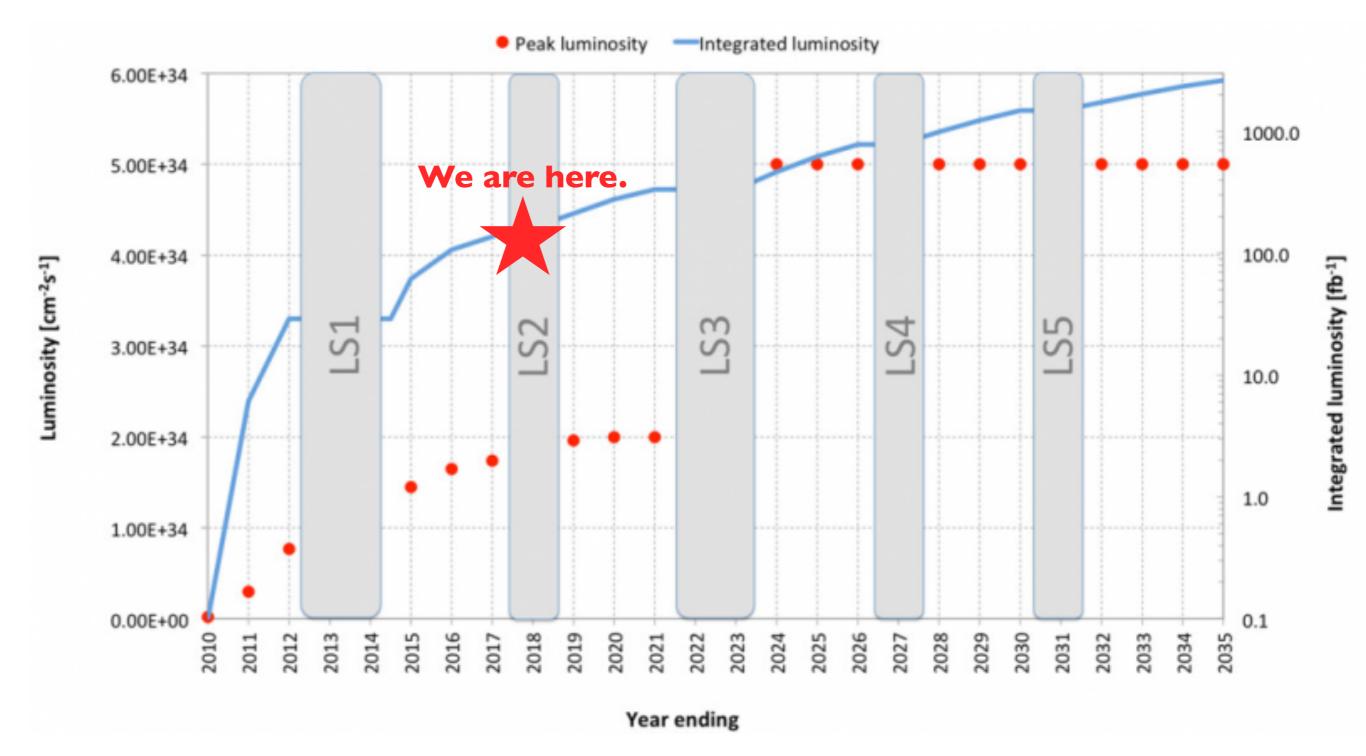
Annual physics jamboree Rencontres de Moriond has a history of revealing exciting results from colliders, and this year new theories and evidence abound



#### SM: complete yet incomplete

- Complete: could be a consistent theory valid up to the Planck scale.
- Incomplete: many open questions
  - Origin of electroweak scale
  - Dark matter
  - Origin of CP, flavor
  - Matter anti-matter asymmetry
  - **>**
- Goal of particle physics: answer these questions.
- Colliders (LHC and beyond) will be crucial.

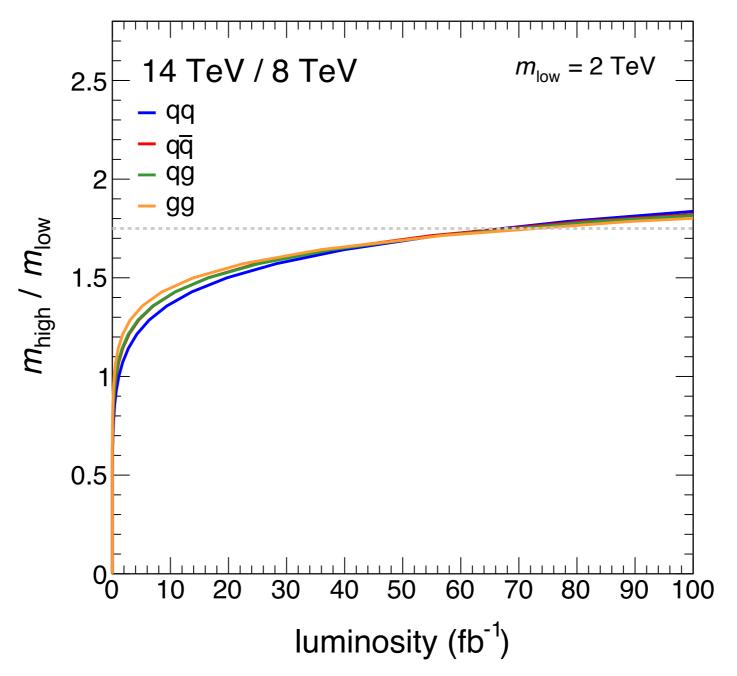
# Road ahead at the LHC



#### LHC is pushing ahead.

Exp. collaborations are pursuing a broad and comprehensive physics program: SUSY, composite H, extra Dim, etc.

#### As data accumulates



Rapid gain initial 10s-100 fb<sup>-1</sup>, slow improvements afterwards.

Progress will become slower, harder

# New directions?

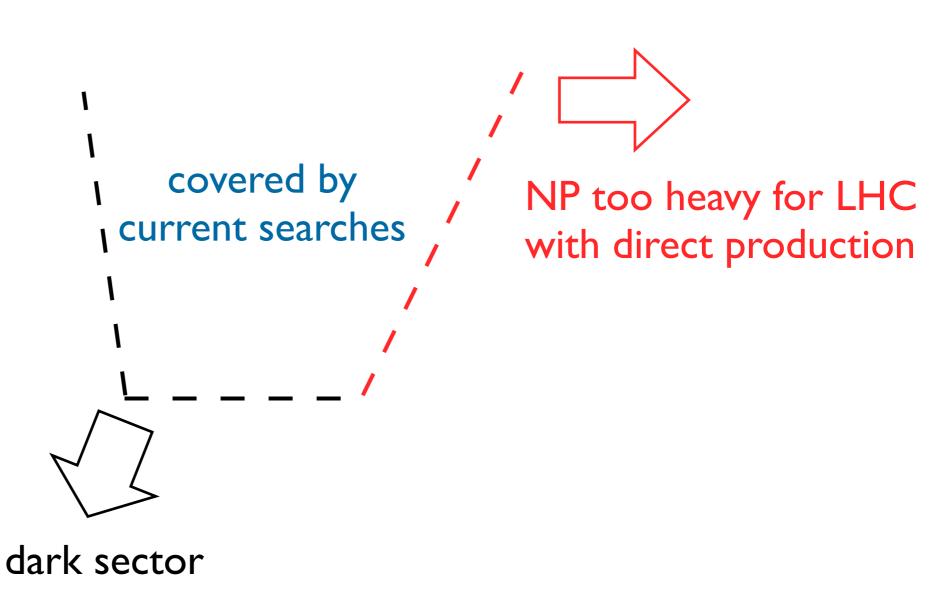
# stronger coupling



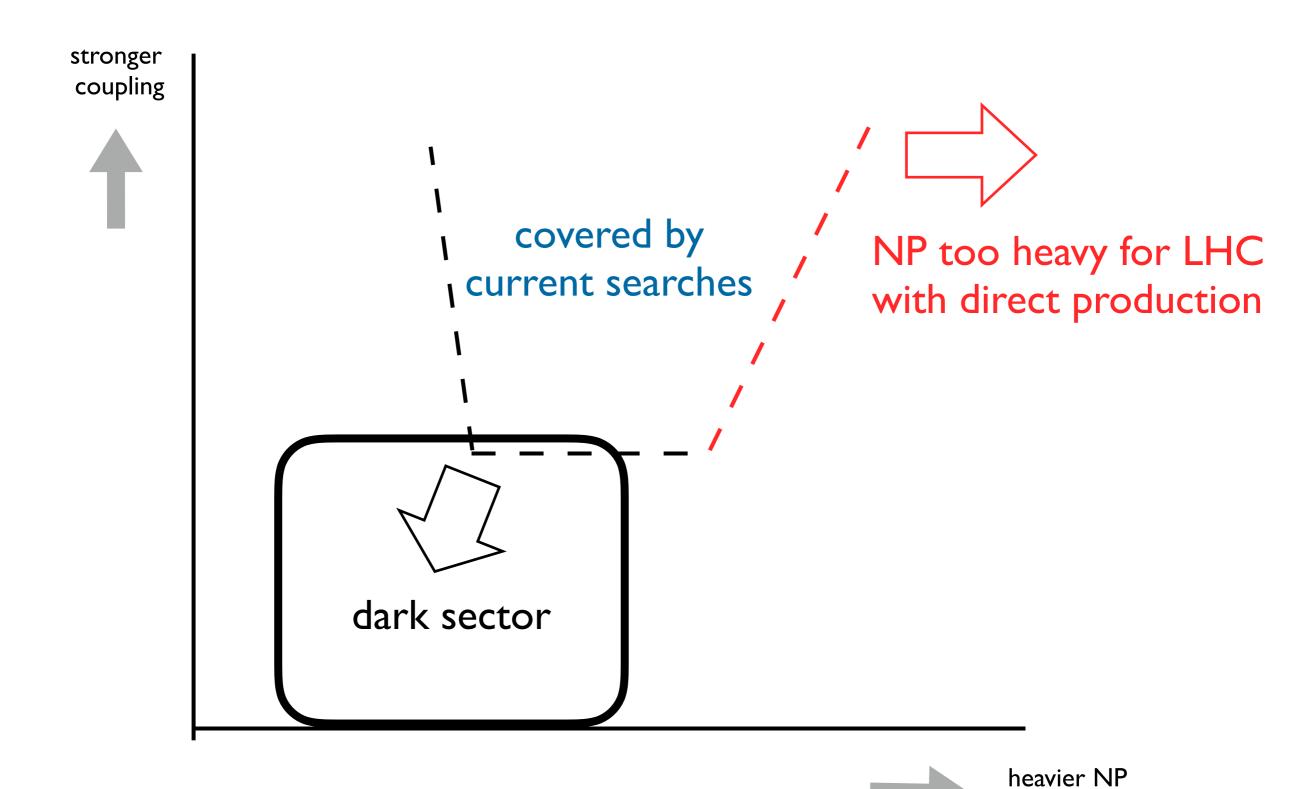
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covered by current searches
```

# stronger coupling





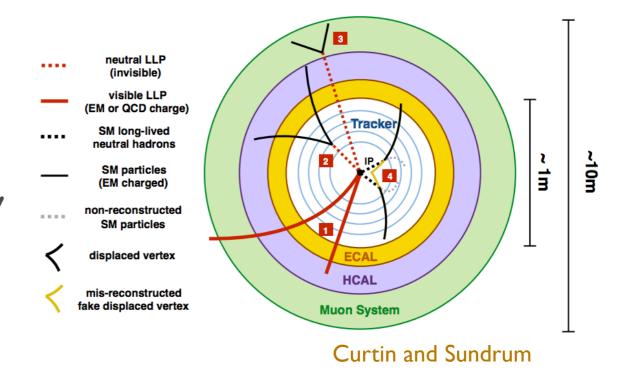
heavier NP particle



particle

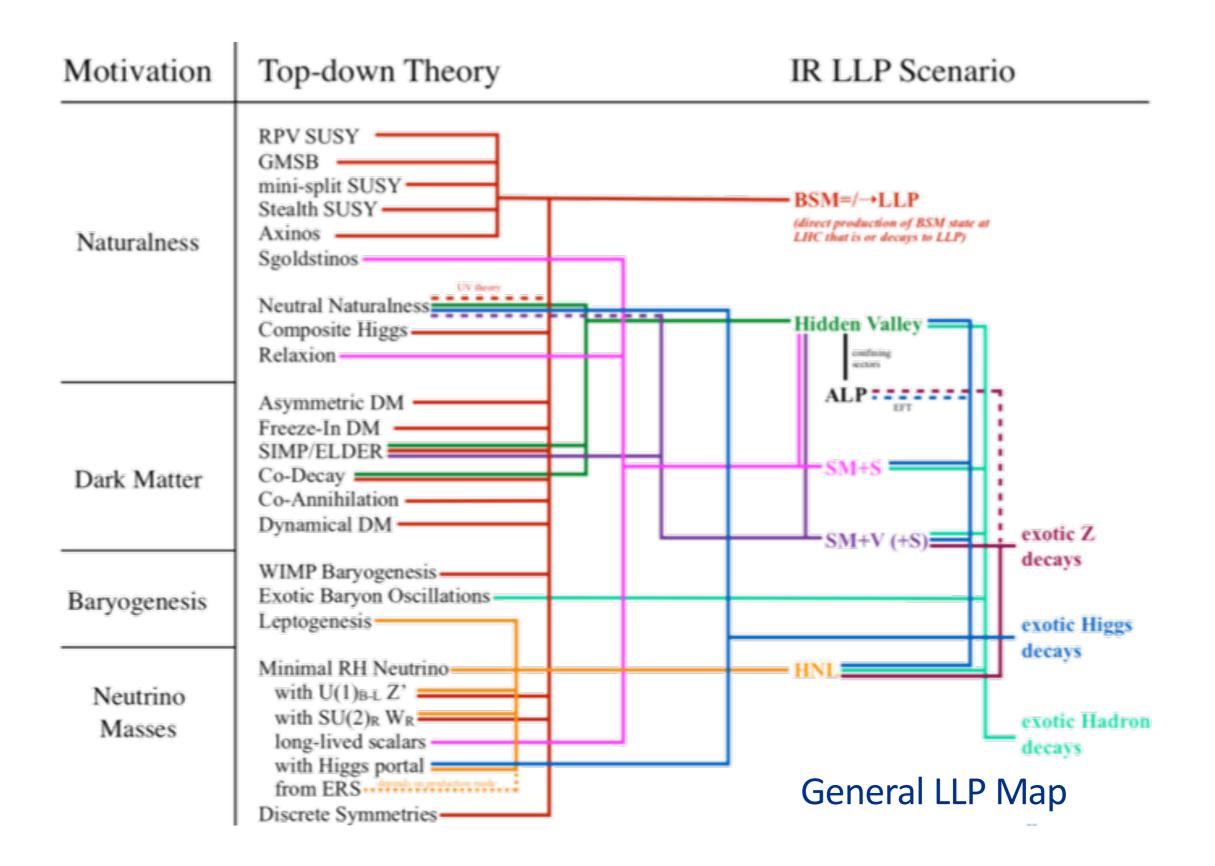
## Example: Long Lived particles (LLP)

- Very weakly coupled to the SM.
  - $\triangleright$  Connection with dark matter, neutrino, etc.  $\tau$
- Displaced-Long lived, soft, kink,
   Covered by LHC searches already.
- Cosmological constraints from BBN:  $\tau < 0.1 \text{ sec } (10^7 \text{ m})$

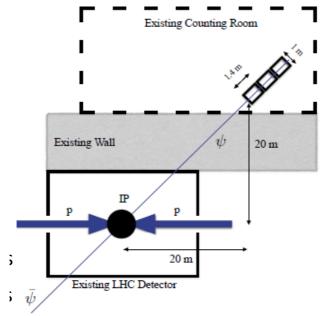


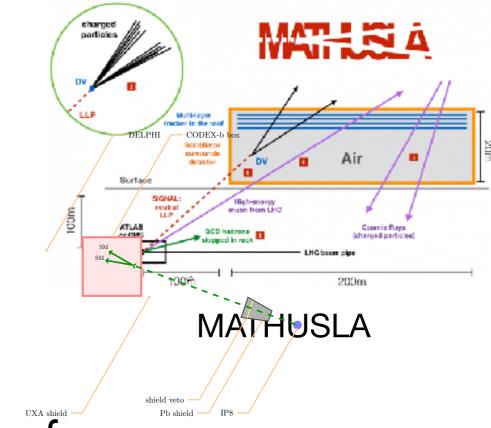
Here, I focus on: decay length >> 10 meters

#### tons of models

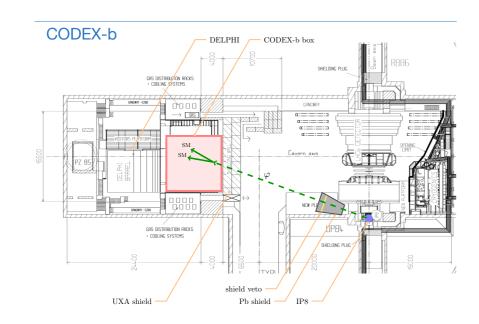


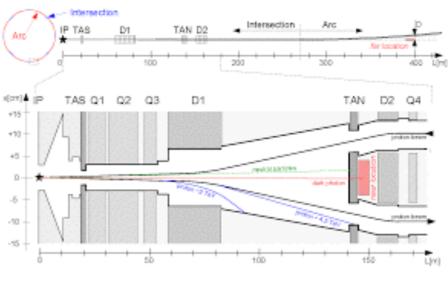
#### Far detectors





new detectors far away from the interaction region

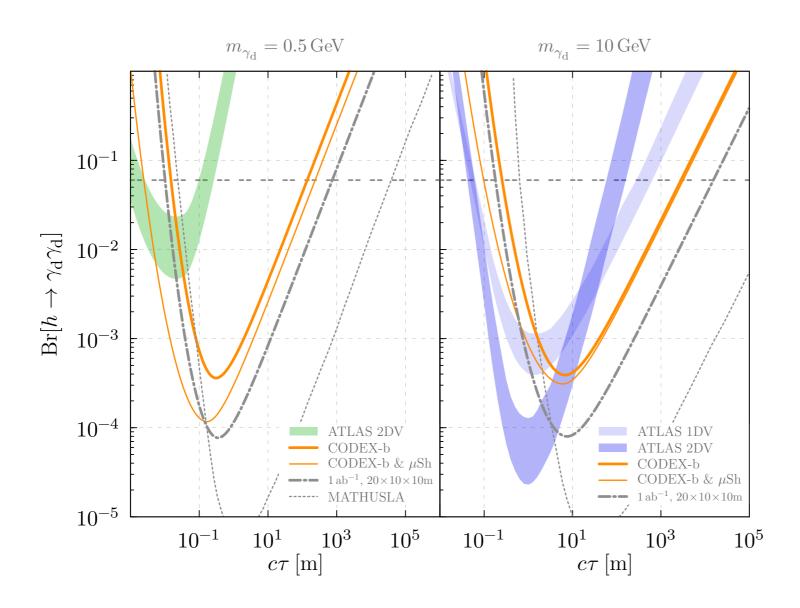


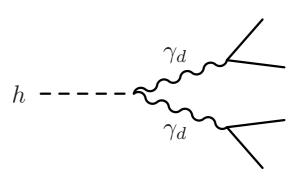


**FASER** 

#### Could reach T≈104-5 m

#### Exotic Higgs decays

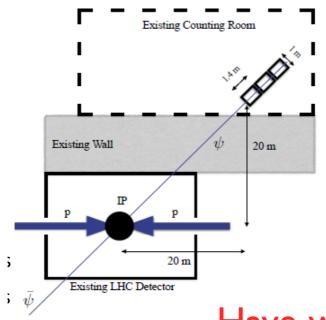


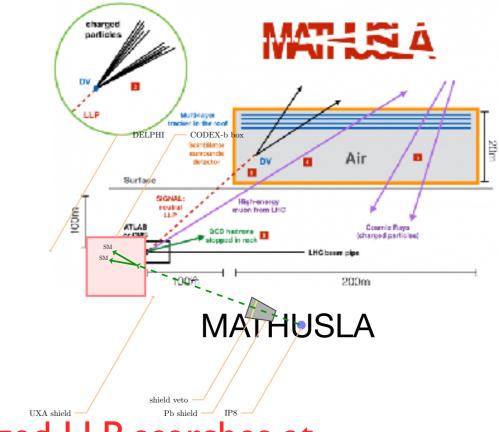


Application:
Neutral Naturalness
(See back-up material)

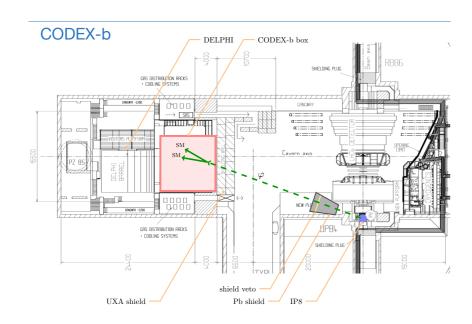
For low masses, ATLAS/CMS are background limited, CODEX-b & MATHUSLA have an edge

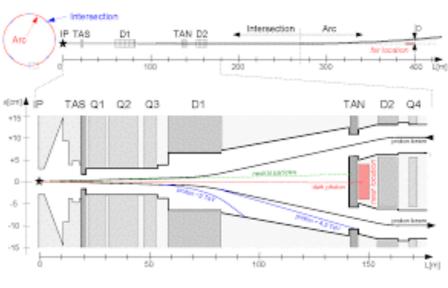
#### Far detectors





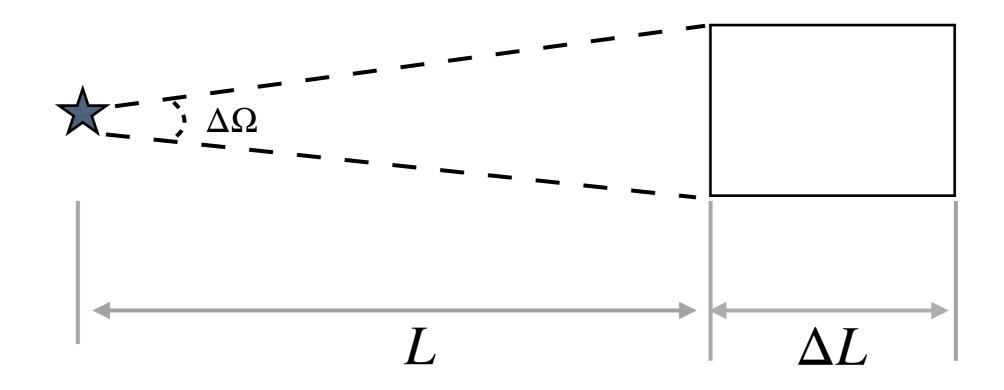
Have we fully optimized LLP searches at the interaction points ATLAS, CMS, LHCb?





**FASER** 

#### Optimal place to catch LLP



Number of particle decayed within detector volume:

$$\#_{\text{in}} \simeq \#_{\text{produced}} \times \frac{\Delta\Omega}{4\pi} \times \frac{\Delta L}{d} e^{-L/d}$$

 $d = \gamma c \tau$  decay length  $d \gg \Delta L, L$ 

Very long lived:  $d \ge 100$ s meters

### Optimal place to catch LLP

Number of particle decayed within detector volume:

$$\#_{\text{in}} \simeq \#_{\text{produced}} \times \frac{\Delta \Omega}{4\pi} \times \frac{\Delta L}{d} e^{-L/d}$$
  $d = \gamma c \tau$ 

ATLAS/CMS (LHCb)

 $\Delta\Omega$ 

 $\sim 4\pi$ 

 $\Delta L$ 

1 - 10 meters

 $\boldsymbol{L}$ 

1 - 10 meters

Far detectors

< 0.1

1 - 10 meters

10 - 100 meters

### Optimal place to catch LLP

$$\#_{\text{in}} \simeq \#_{\text{produced}} \times \frac{\Delta \Omega}{4\pi} \times \frac{\Delta L}{d} e^{-L/d}$$
  $d = \gamma c \tau$ 

ATLAS/CMS (LHCb)

$$\Delta\Omega$$
 ~  $4\pi$ 

$$\Delta L$$
 1 – 10 meters

$$L = 1 - 10$$
 meters

Far detectors

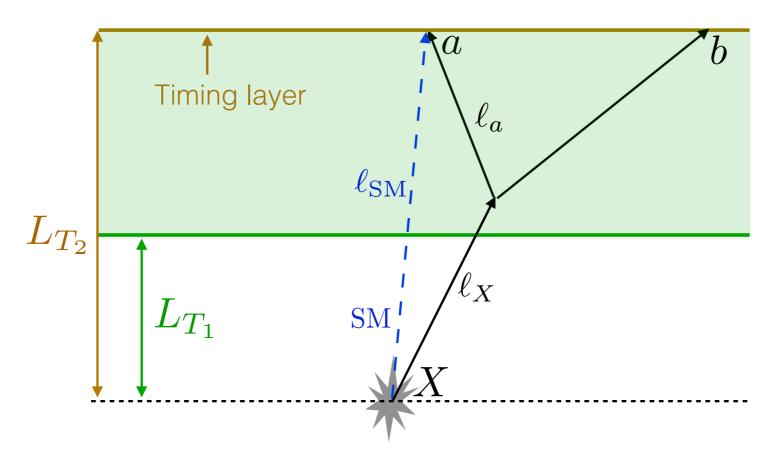
$$1 - 10$$
 meters

$$10 - 100$$
 meters

Advantage of far detector? Far away from interaction point, less background.

New proposal: use timing information Significantly lower background near interaction point.

#### Time delay

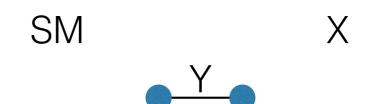


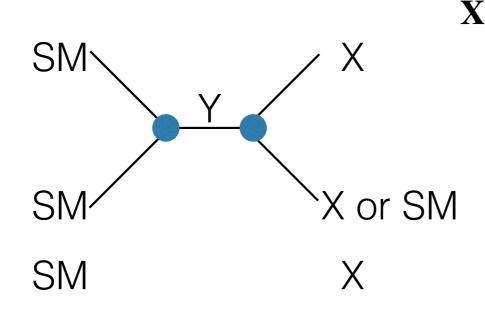
$$\Delta t = \frac{\ell_X}{\beta_X} + \frac{\ell_a}{\beta_a} - \frac{\ell_{\rm SM}}{\beta_{\rm SM}}$$
  $\beta_a \simeq \beta_{\rm SM} \simeq 1$ 

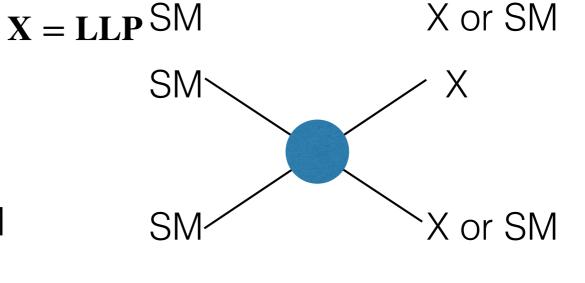
Good for massive LLP produced with small or moderate boost

$$\beta_X < 1$$

# Basic topologies







boost:<sub>SM</sub>

$$\gamma \simeq \frac{m_Y}{2m_X^2}$$
 or SM

challenging for  $m_X \ll m_Y$ 

benchmark: Higgs portal

$$Y = Higgs$$

 $X \rightarrow SM$  Long lived

boost:

$$\gamma \sim 1$$

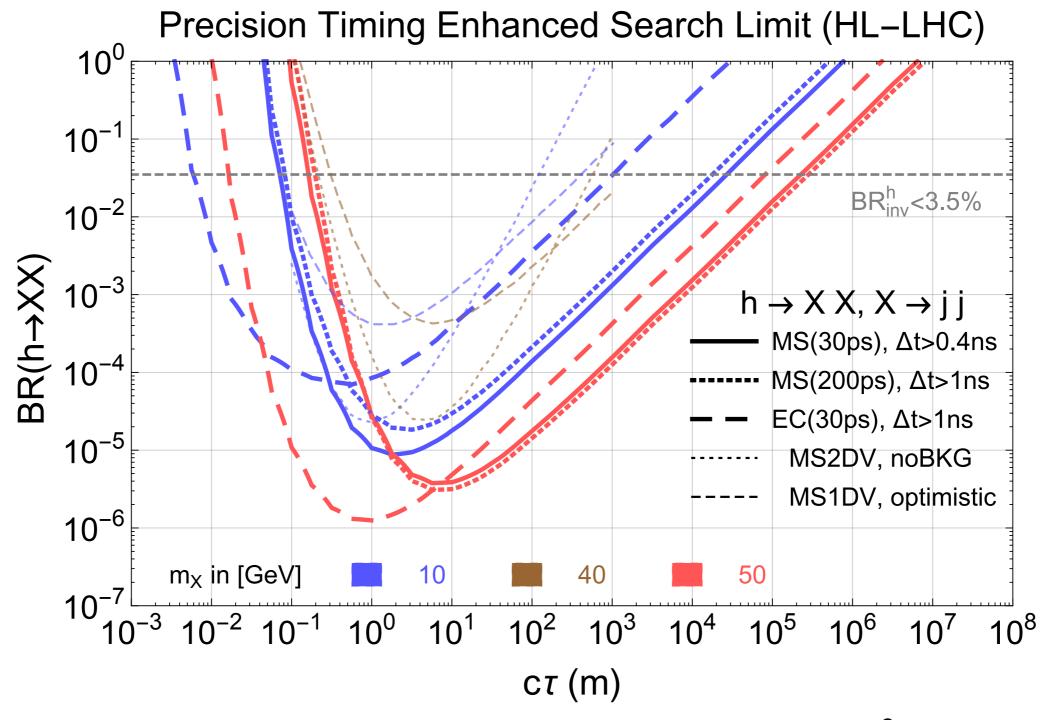
slow moving, sizable  $\Delta t$ 

benchmark: SUSY

$$X = neutralino$$

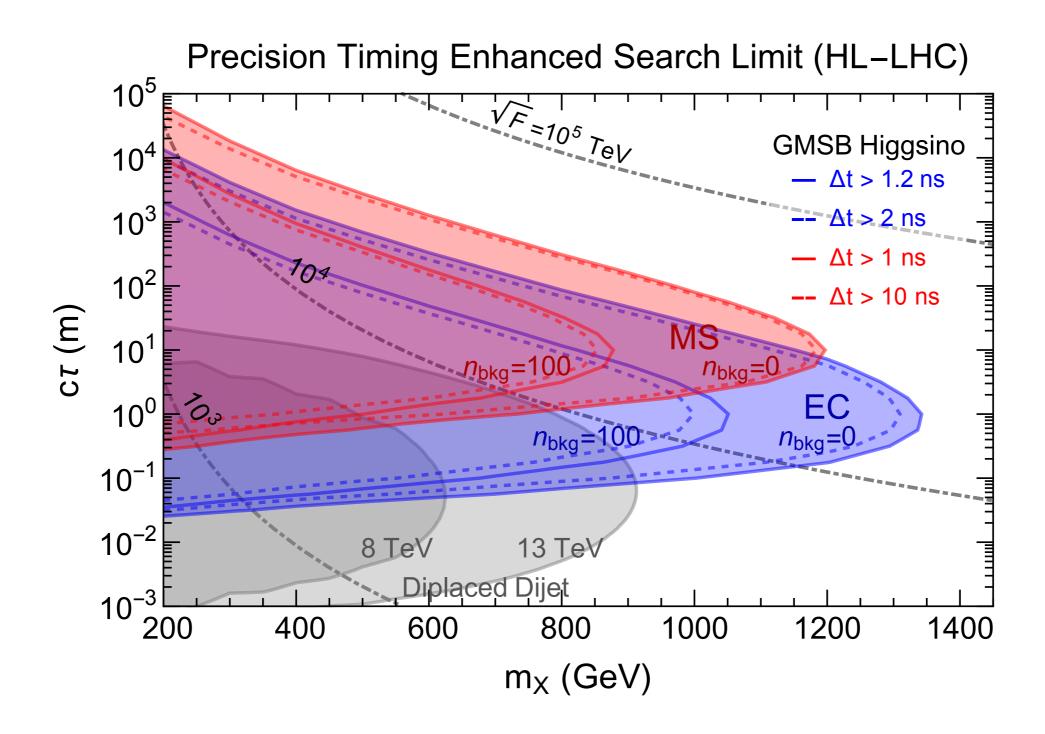
$$\chi_0 \rightarrow \text{gravitino} + \dots$$
 Long lived

## Sensitivity to Higgs portal

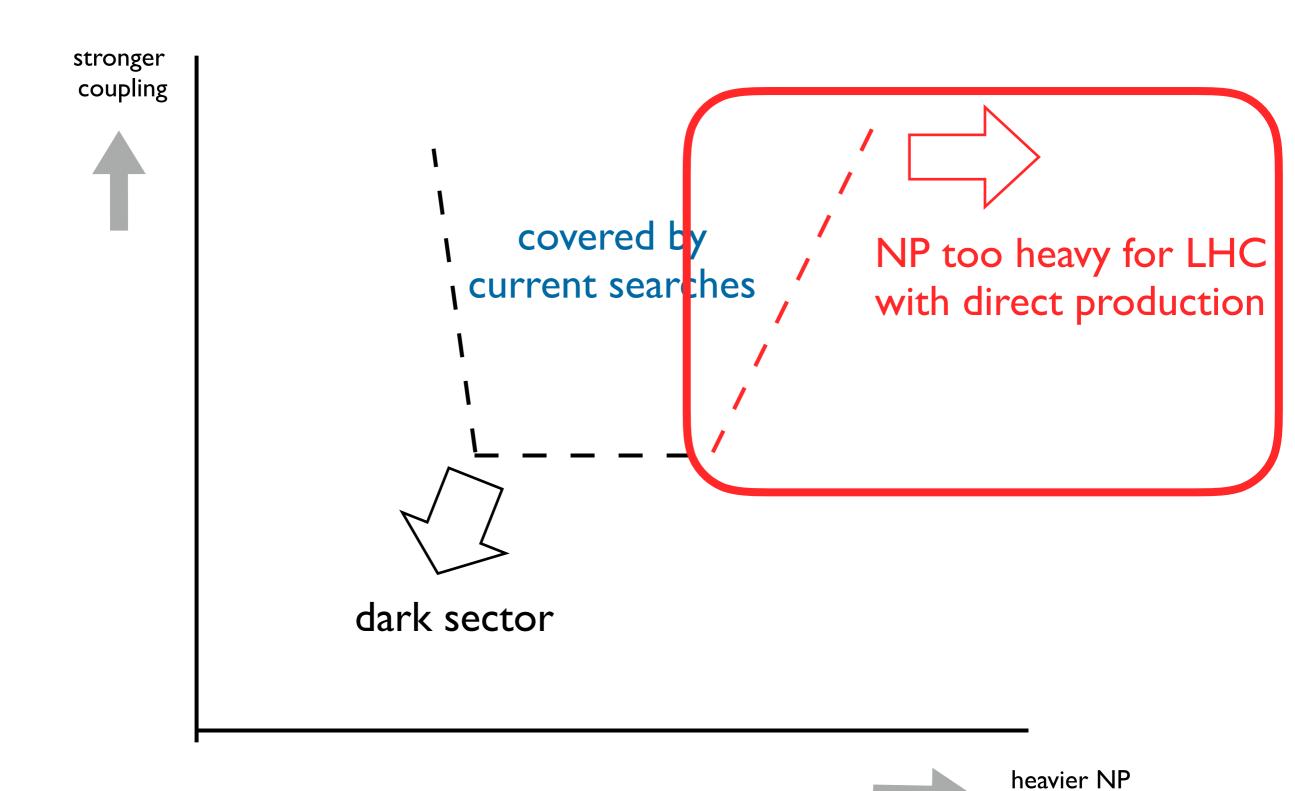


For example, for BR( $h \to XX$ )  $\sim 10^{-3}$  EC(MS) reach can be  $c\tau \sim 10^3 (10^4)$  meters

### Sensitivity to SUSY



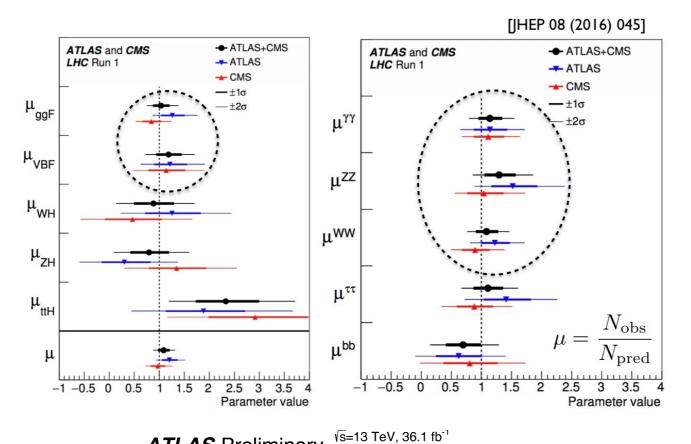
Slower moving LLP, timing cuts can be further relaxed.



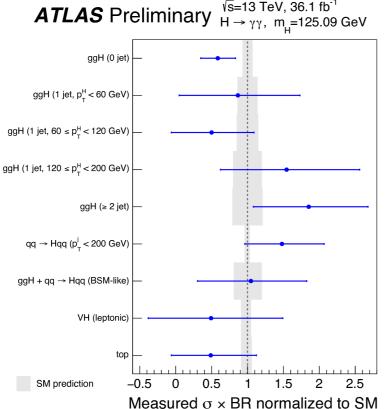
particle

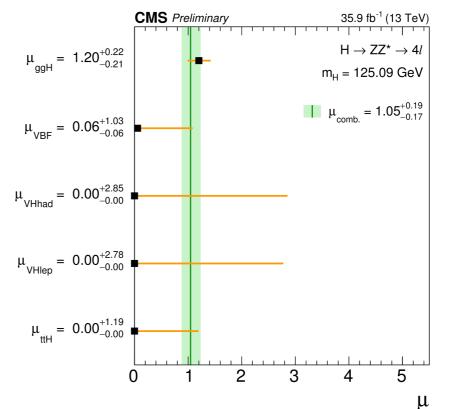
# Revealing trace of new physics with precision measurements

### Higgs Standard Model-like



# Agree to about 10-20%





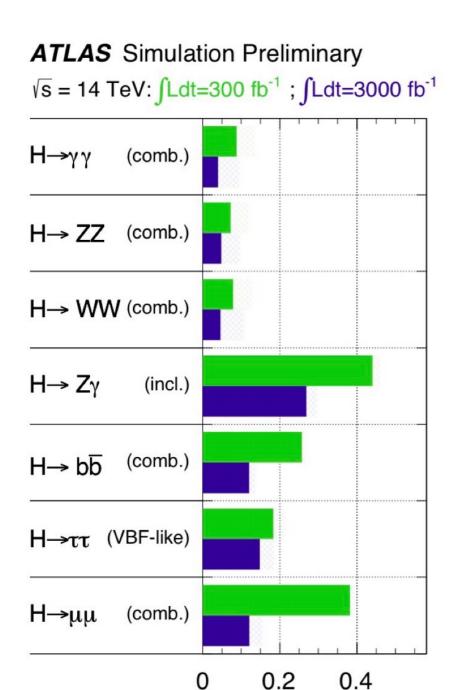
#### Not entirely surprising

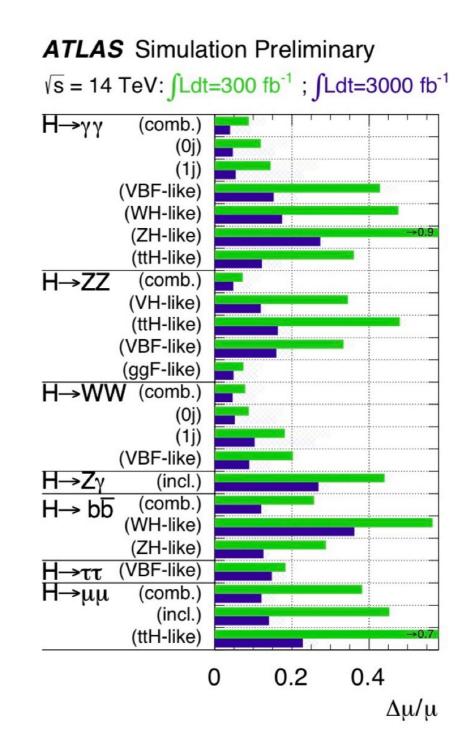
 In general, deviation induced by new physics is of the form

$$\delta \simeq c \frac{v^2}{M_{
m NP}^2}$$
 M<sub>NP</sub>: mass of new physics c: O(1) coefficient

- Current LHC precision: 10%  $\Rightarrow$  sensitive to M<sub>NP</sub> < 500-700 GeV
- At the same time, direct searches constrain new physics below TeV already.
- Unlikely to see O(1) deviation.

# Significant improvement with high lumi

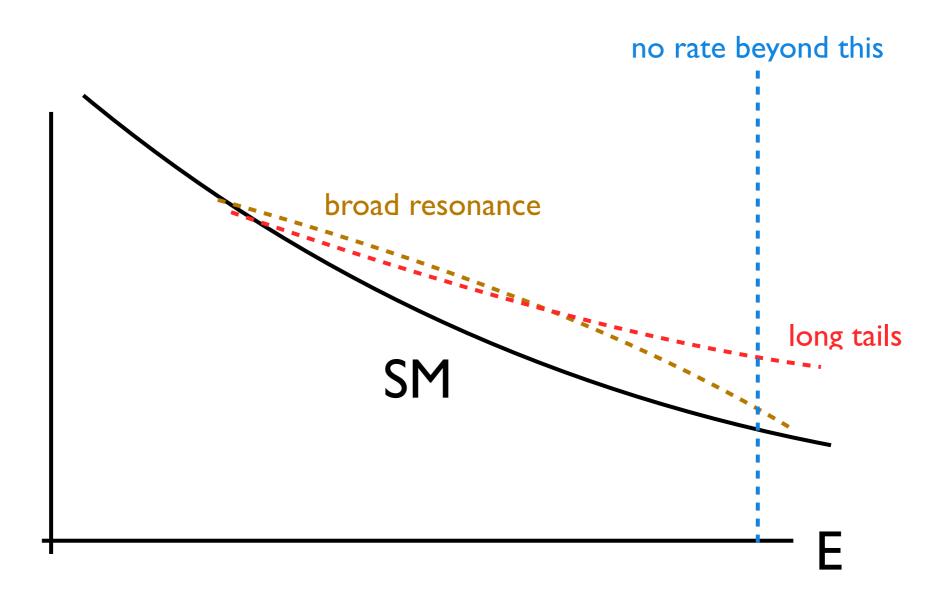






 $\Delta \mu / \mu$ 

# Precision measurement with distribution

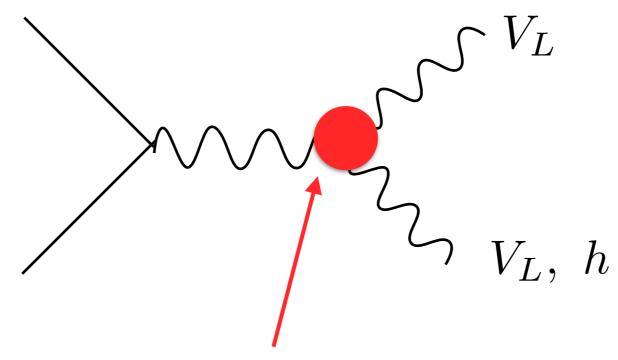


Low S/B, systematic dominated.

Room to improve.

#### Diboson production at the LHC

$$q\bar{q} \rightarrow VV, \quad V = W, Z, h.$$



New physics contribution

New physics effect encoded in the non-renormalizable operators:

$$\frac{1}{\Lambda^2}\mathcal{O}$$

 $\Lambda$ : new physics scale

#### Precision measurement at the LHC possible?

LEP precision tests probe NP about 2 TeV

$$\frac{\delta\sigma}{\sigma_{\rm SM}} \sim \frac{m_W^2}{\Lambda^2} \sim 2 \times 10^{-3} \quad \to \Lambda \ge 2 \text{ TeV}$$

At LHC, new physics effect grows with energy

$$\frac{\delta\sigma}{\sigma_{\rm SM}} \sim \frac{E^2}{\Lambda^2} \sim 0.25$$
  $E \sim 1 \text{ TeV}, \ \Lambda \sim 2 \text{ TeV}$ 

LHC needs to make a 20% measurement to beat LEP LHC has potential.

#### Picking final state important

#### At LHC, interference with SM crucial

Signal-SM interference

$$\frac{\delta\sigma}{\sigma_{\rm SM}} \sim \frac{E^2}{\Lambda^2} \sim 0.25$$

Without interference

$$\frac{\delta\sigma}{\sigma_{\rm SM}} \sim \frac{E^4}{\Lambda^4} \sim 0.05$$

- I. WZ final states, only longitudinal mode useful
- 2. W/Z+h

# Will be challenging

SM WW, WZ processes are dominated by transverse modes

$$\sigma_{SM}^{total}/\sigma_{SM}^{LL}\sim 15-50$$

New technique such as polarization tagging of W/Z crucial

Wh/Zh(bb) channels have large reducible background

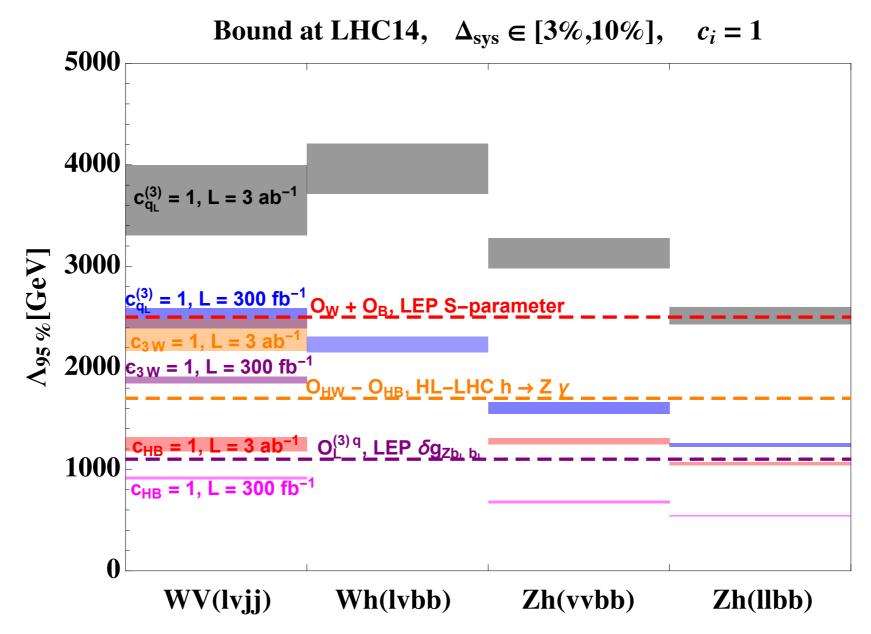
LHC @ 8 TeV : 
$$\sigma_b^{red}/\sigma_{SM}^{Wh}\sim 200-10$$

Difficult measurement. Large improvement needed. Room for developing new techniques

# Operators: d=6

name	structure	coefficient (power counting)
$\mathcal{O}_H$	$rac{1}{2}(\partial_{\mu} H ^2)^2$	$c_H/f^2$
$\mathcal{O}_y$	$yQ_LHu_R H ^2$	$c_y/f^2$
$\mathcal{O}_W$	$ig\left(H^{\dagger}\sigma^{a}\overleftrightarrow{D}^{\mu}H\right)D^{\nu}W^{a}_{\mu\nu}$	$c_W/m_*^2$
$\mathcal{O}_B$	$ig' ig( H^\dagger \overleftrightarrow{D}^\mu H ig) D^ u B_{\mu u}$	$c_B/m_*^2$
$\mathcal{O}_{HW}$	$ig(\mathring{D^{\mu}}H)^{\dagger}\sigma^{a}(\mathring{D^{\nu}}H)\mathring{W}^{a}_{\mu\nu}$	$c_{HW}/m_*^2 \times (g_*/4\pi)^2$
$\mathcal{O}_{HB}$	$ig'(D^{\mu}H)^{\dagger}(D^{\nu}H)B_{\mu\nu}$	$c_{HB}/m_*^2 \times (g_*/4\pi)^2$
$O_L^q$	$ig^2ig(H^\dagger \overleftrightarrow{D}_\mu Hig)ar{Q}_L \gamma^\mu Q_L$	$c_q/m_*^2  imes \epsilon_q^2$
$O_L^{q,3}$	$ig^2ig(H^\dagger\sigma^a\overleftrightarrow{D}_\mu Hig)ar{Q}_L\sigma^a\gamma^\mu Q_L$	$c_{q,3}/m_*^2 \times \epsilon_q^2$
$O_R^u$	$ig^2ig(H^\dagger \overleftrightarrow{D}_\mu Hig)ar{u}_R \gamma^\mu u_R$	$c_u/m_*^2  imes \epsilon_u^2$
$O_R^u$	$ig^2ig(H^\dagger \overleftrightarrow{D}_\mu Hig)ar{d}_R \gamma^\mu d_R$	$c_d/m_*^2  imes \epsilon_d^2$
$O_T$	$\left(H^{\dagger} \overrightarrow{D}_{\mu} H\right)^{2}$	$c_T/f^2$
$\mathcal{O}_6$	$ H ^6$	$\lambda_3/f^2$

#### Projections



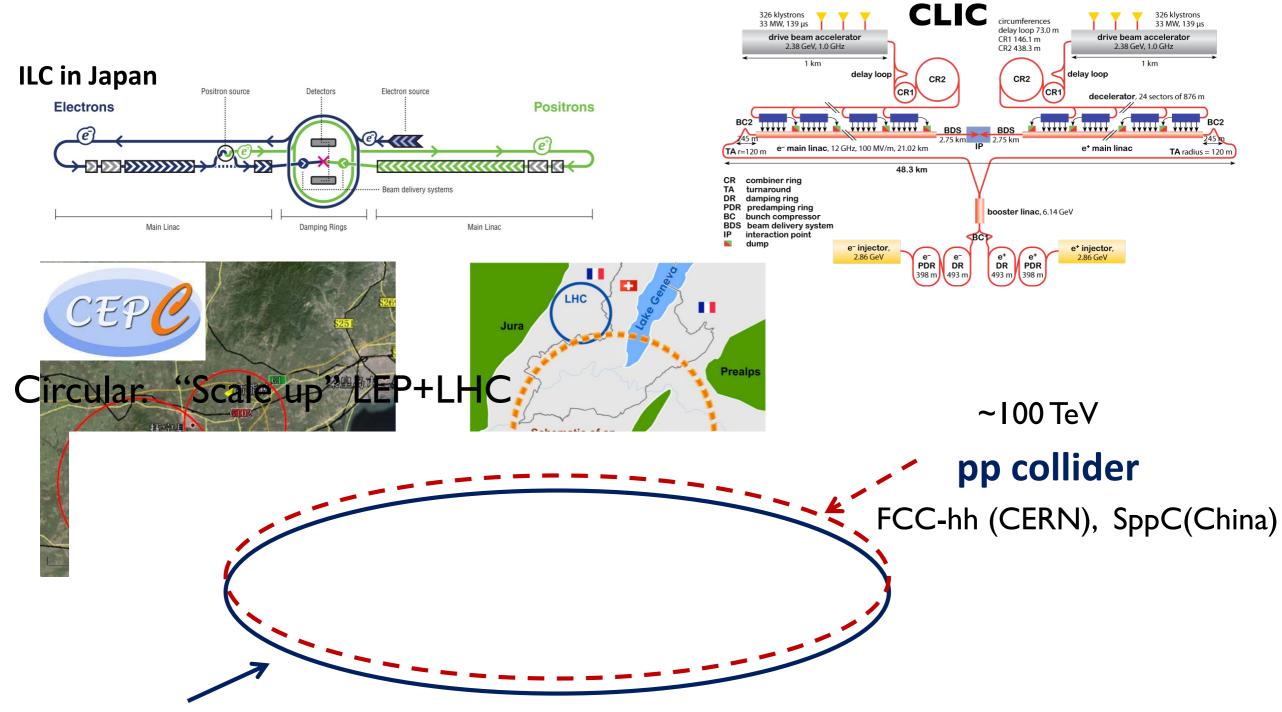
Possible to reach 4 TeV.

Better than LEP, and many LHC direct searches

See also: Alioli, Farina, Pappadopulo, Ruderman, Franceschini, Panico, Pomarol, Riva, Wulzer, Azatov, Elias-Miro, Regimuaji, Venturini

# Beyond LHC

#### Future Colliders



250 GeV e<sup>-</sup>e<sup>+</sup> Higgs Factory

FCC-ee (CERN), CEPC(China)

## Ambitious program

#### FCC-ee:

FCC-ee possible operation model						
working point	luminosity/IP [10 <sup>34</sup> cm <sup>-2</sup> s <sup>-1</sup> ]	total luminosity (2 IPs)/ yr	physics goal	run time [years]		
Z first 2 years	100	26 ab <sup>-1</sup> /year	150 ab <sup>-1</sup>	4		
Z later	200	52 ab <sup>-1</sup> /year				
W	32	8.3 ab <sup>-1</sup> /year	10 ab <sup>-1</sup>	1		
Н	7.0	1.8 ab <sup>-1</sup> /year	5 ab <sup>-1</sup>	3		
machine modification for RF installation & rearrangement: 1 year						
top 1st year (350 GeV)	0.8	0.2 ab <sup>-1</sup> /year	0.2 ab <sup>-1</sup>	1		
top later (365 GeV)	1.5	0.38 ab <sup>-1</sup> /year	1.5 ab <sup>-1</sup>	4		

 $\sim 10^6$  Higgses,  $\sim 10^{13}$  Zs,...

13 yr run plan: Higgs=3, Z=4, top=5, W=1

#### **CEPC Operation Plan**

Partide type	Energy (c.m.) (GeV)	Luminosity per IP (10 <sup>34</sup> cm <sup>-2</sup> s <sup>-1</sup> )	Luminosity per year (ab <sup>-1</sup> , 2 IPs)	Years	Total luminosity (ab <sup>-1</sup> , 2 IPs)	Total number of particles
Н	240	3	0.8	7	5.6	1 x 10 <sup>6</sup>
Z	91	32	8	2	16	7 x 10 <sup>11</sup>
W	160	10	2.6	1	2.6	8 x 10 <sup>6</sup>

#### **CEPC** yearly run time assumption:

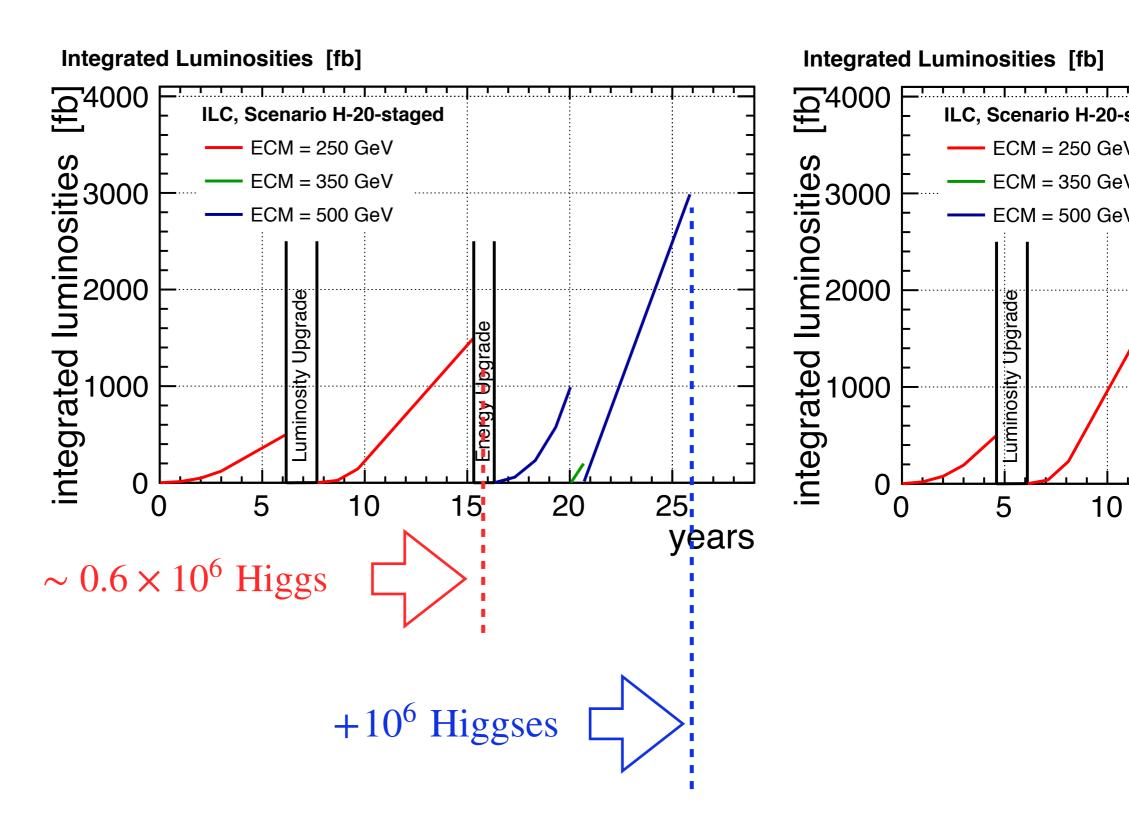
- Operation 8 months, or 250 days, or 6,000 hrs
- Physics (60%) 5 months, or 150 days, or 3,600 hrs, or 1.3 Snowmass Unit.

Currently, no plan to scan the ttbar threshold.

## **CEPC**

staging scheme	physics focus
7 year at Higgs ~1M events	H
240 GeV (initial stage)	indir. BSM
2 years at Z upto 10 <sup>12</sup> events	Z, W
1 year at WW ~20M events	EW Physics

## ILC run plan



No Z-pole or WW run planned

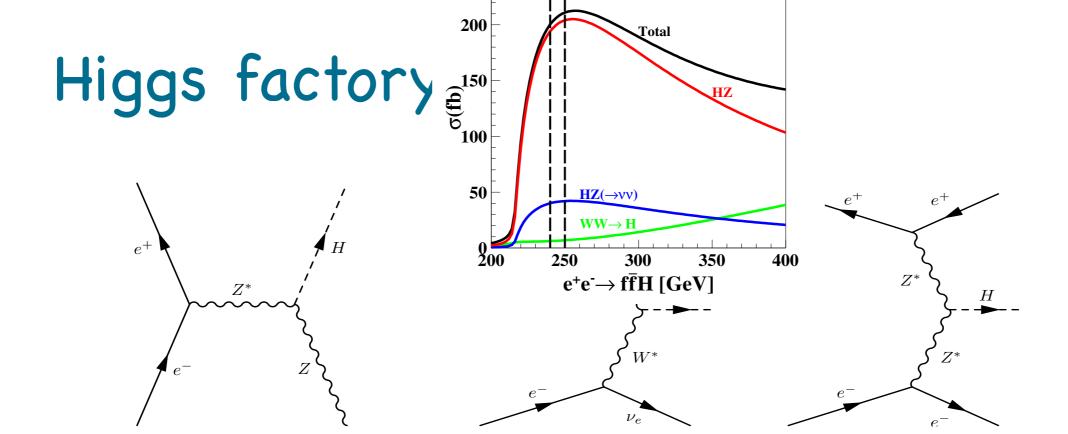
## What precision should we aim for?

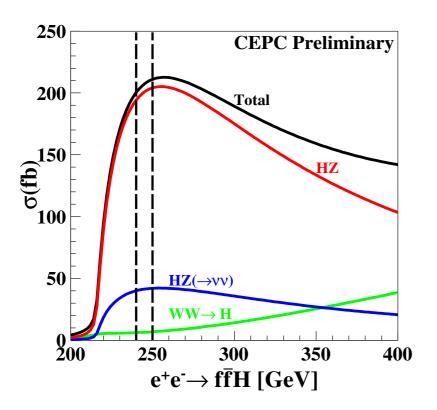
- In precision measurements, we are going after deviations of the form

$$\delta \simeq c \frac{v^2}{M_{
m NP}^2}$$
 M<sub>NP</sub>: mass of new physics c: O(1) coefficient

- Take the Higgs coupling.
  - ▶ LHC precision: 5-10% ⇒ sensitive to  $M_{NP}$  < TeV

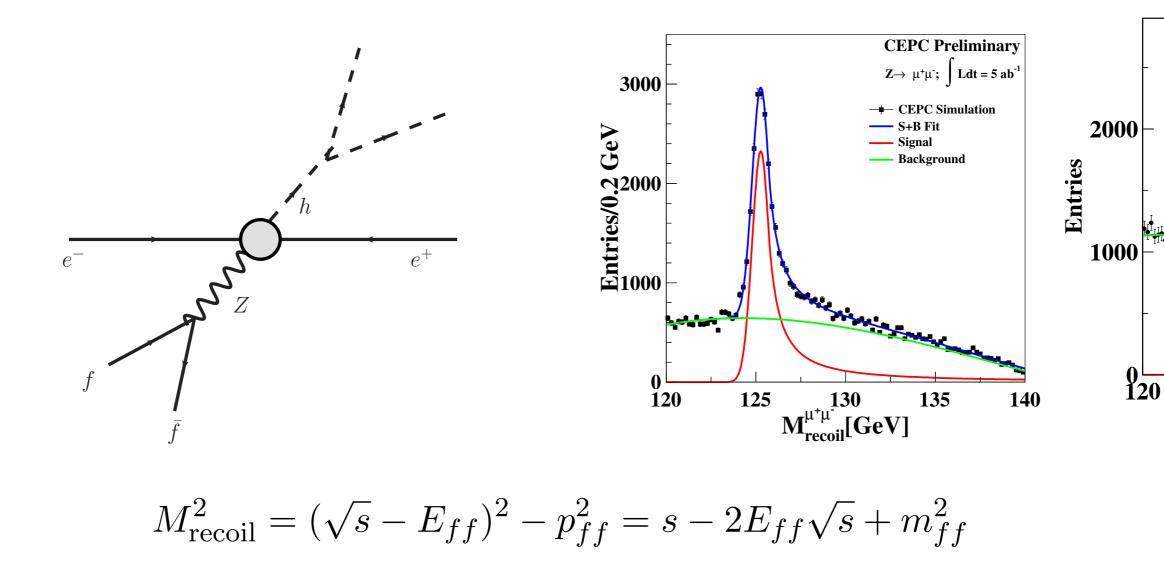
• To go beyond the LHC, need 1% or less precision.





Process	Cross section	Nevents in 5 $ab^{-1}$		
Higgs boson production, cross section in fb				
$e^+e^- \to ZH$	212	$1.06 \times 10^{6}$		
$e^+e^- \to \nu\nu H$	6.72	$3.36 \times 10^4$		
$e^+e^- \to eeH$	0.63	$3.15 \times 10^{3}$		
Total	219	$1.10 \times 10^{6}$		

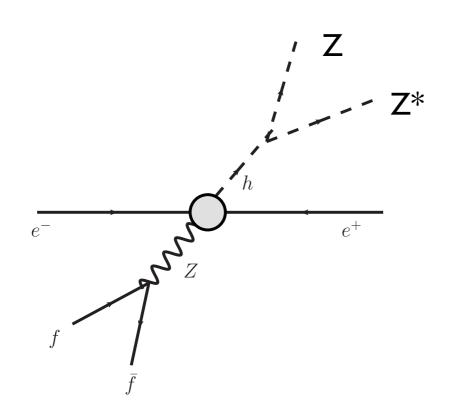
#### Zh cross section



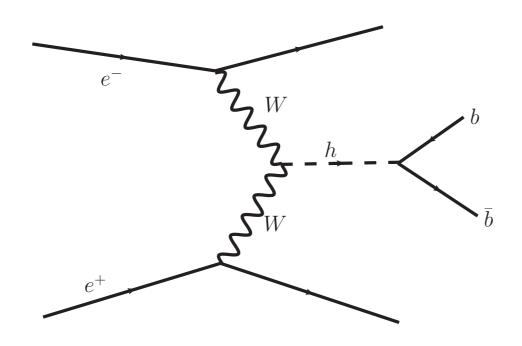
Can use recoil mass to identify Zh process, independent of Higgs decay

⇒ inclusive measurement of Zh cross section

# Higgs width. Unique capability of lepton colliders.

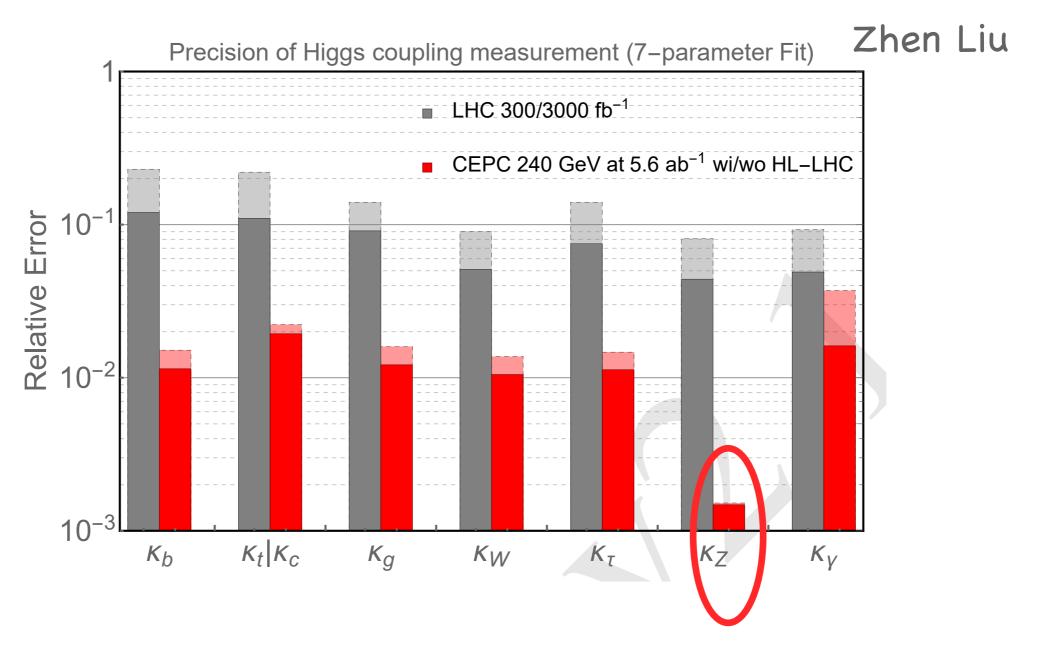


$$\Gamma_H \propto \frac{\Gamma(H \to ZZ^*)}{\mathrm{BR}(H \to ZZ^*)} \propto \frac{\sigma(ZH)}{\mathrm{BR}(H \to ZZ^*)}$$



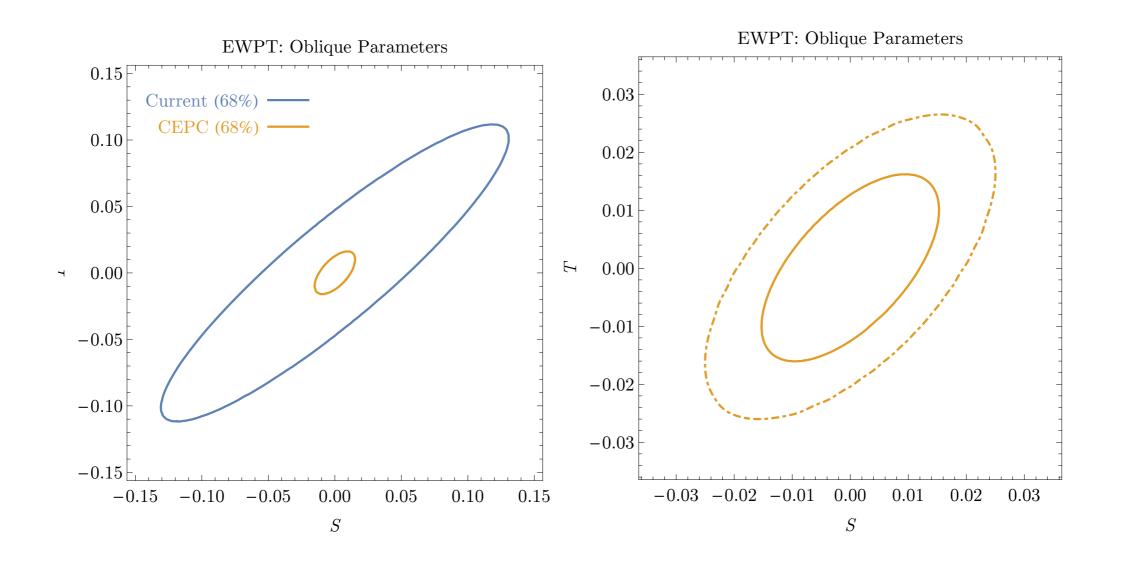
$$\Gamma_H \propto \frac{\Gamma(H \to bb)}{\mathrm{BR}(H \to bb)} \propto \frac{\sigma(\nu\nu H \to \nu\nu bb)}{\mathrm{BR}(H \to bb) \cdot \mathrm{BR}(H \to WW^*)}$$

### CEPC can do very well



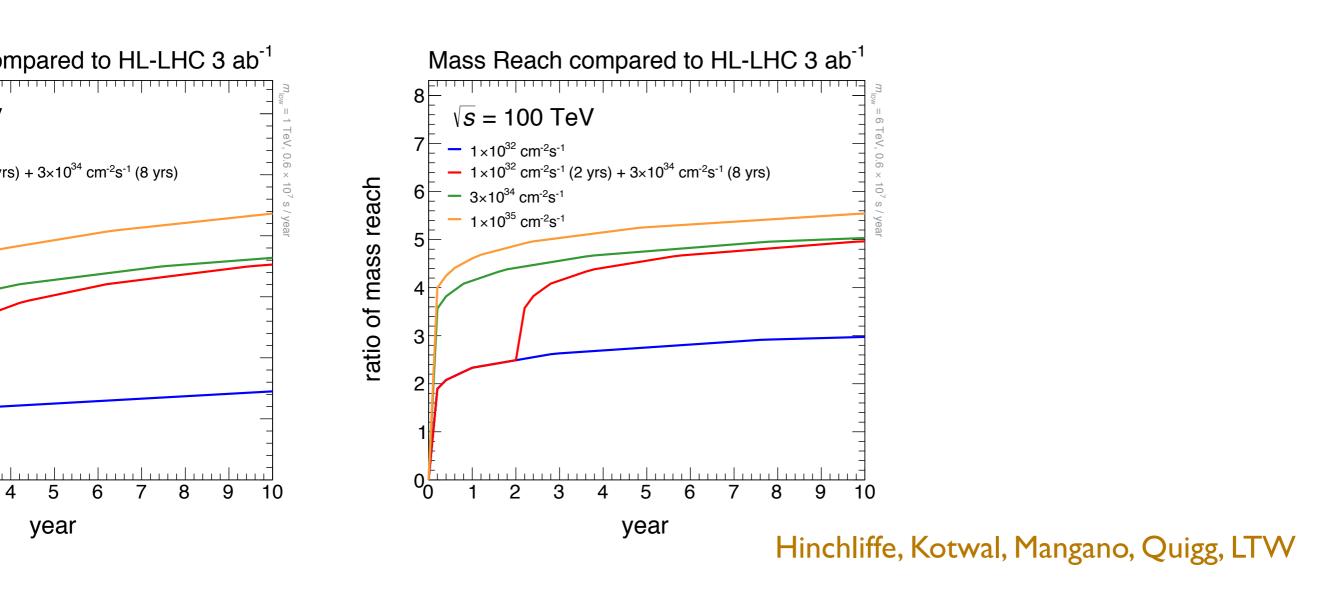
Up to sub percent precision, reach to new physics at multi-TeV scale. Far beyond the reach of LHC.

## Electroweak precision



FCC can do even better (by a factor of a few)

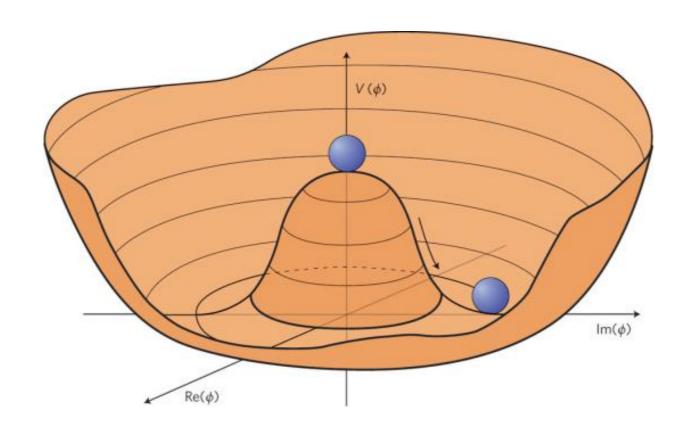
## 100-ish TeV pp collider



A factor of at least 5 increase in reach beyond the LHC, with modest luminosity

# Electroweak symmetry breaking

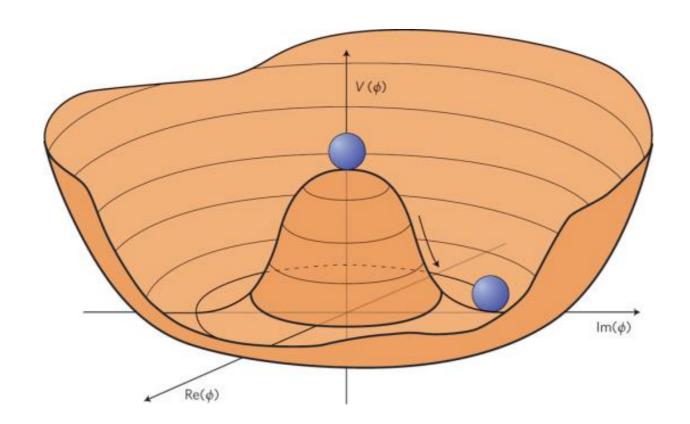
# "Simple" picture: Mexican hat



$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4$$
$$\langle h \rangle \equiv v \neq 0 \rightarrow m_W = g_W \frac{v}{2}$$

Similar to, and motivated by Landau-Ginzburg theory of superconductivity.

# "Simple" picture: Mexican hat



$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4$$

$$\langle h \rangle \equiv v \neq 0 \quad \rightarrow \quad m_W = g_W \frac{v}{2}$$

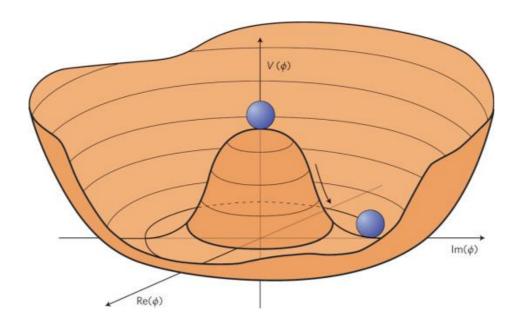
Similar to, and motivated by Landau-Ginzburg theory of superconductivity.

However, this simplicity is deceiving.

Parameters not predicted by theory. Can not be the complete picture.

5

# Mysteries of the electroweak scale.



## Mysteries

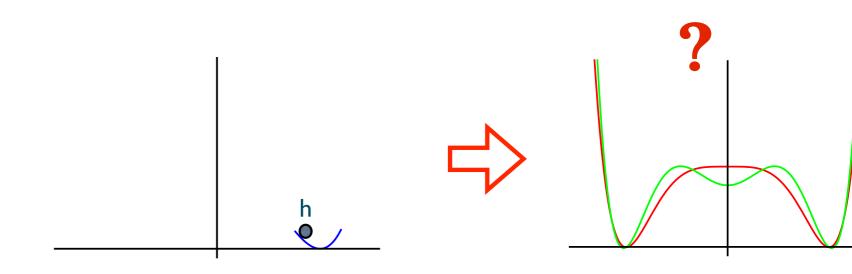


Figure 8: Question of the nature of the electroweak phase transition.

Understanding this physics is also directly relevant to one of the m damental questions we can ask about any symmetry breaking phenomenant which is what is the order of the associated phase transition. How experimentally decide whether the electroweak phase transition in the universe was second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order? This question is a should not be a second order or first order?

Tuesday electroweak symmetry, we must now undertake an experimental pro-

Wednesday, August 13, 19 robe how electroweak symmetry is restored at high energies.

Tuesday, January 20, 15

A first-order phase transition is also strongly motivated by the poof electroweak baryogenesis [18]. While the origin of the baryon asymment of the most fascinating questions in physics, it is frustratingly stronger to build models for baryogenesis at ultra-high energy scale and direct experimental consequences. However, we aren't forced to detect the description of the property of the poor of the most fascinating questions.

## Mysteries

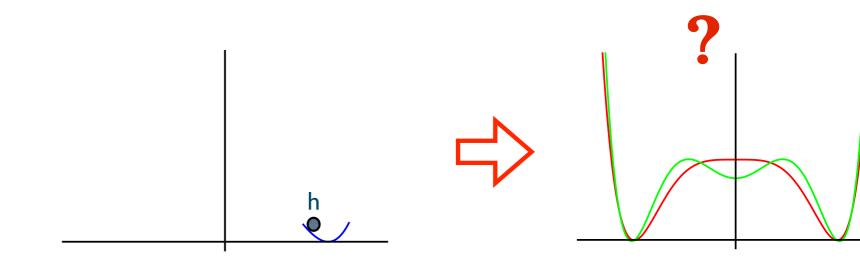


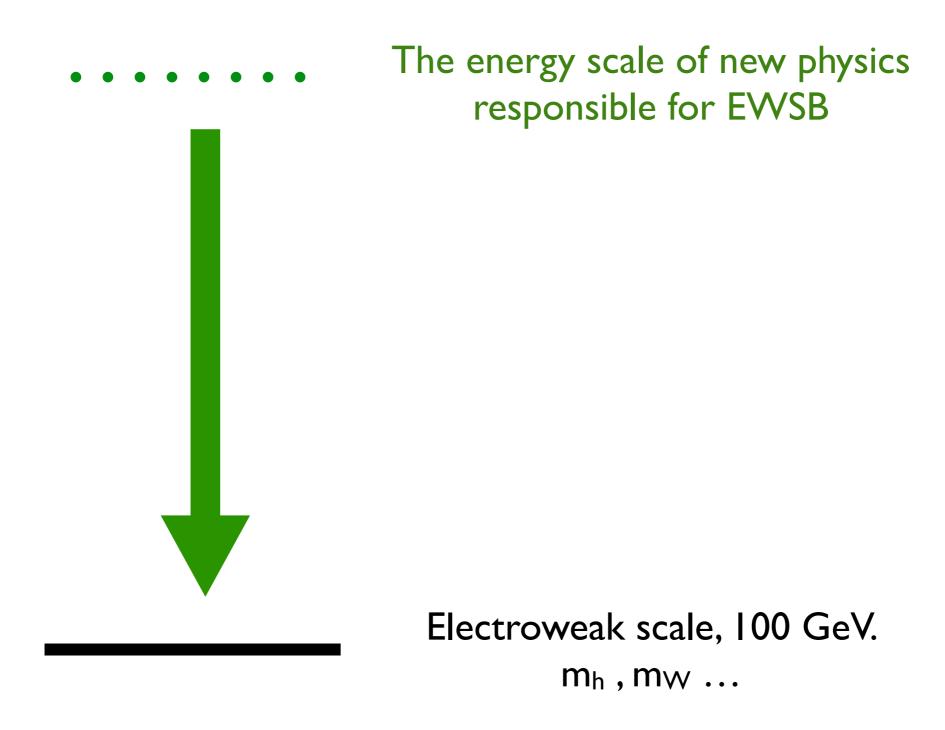
Figure 8: Question of the nature of the electroweak phase transition.

- Understanding this physics is also directly relevant to one of the m damental questions we can ask about any symmetry breaking phenowhich is what is the order of the associated phase transition. How experimentally decide whether the electroweak phase transition in the universe was second order or first order? This question is a should not be a second order or first order?
- Tuesday electroweak symmetry, we must now undertake an experimental progression, of stellar by electroweak symmetry it is the color of the color of
  - A first-order phase transition is also strongly motivated by the poof electroweak baryogenesis [18]. While the origin of the baryon asymment of the most fascinating questions in physics, it is frustratingly stronger to build models for baryogenesis at ultra-high energy scale and direct experimental consequences. However, we aren't forced to de-

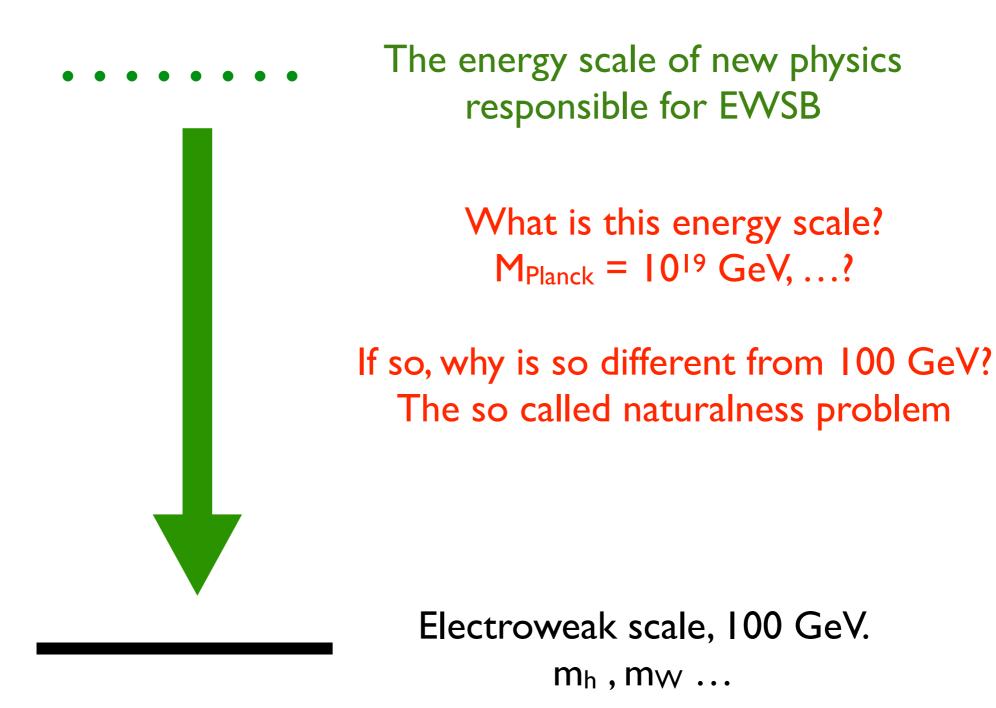
- How to pre
- Full Higgs

Tuesday, January 20, 15

### How to predict Higgs mass?



#### How to predict Higgs mass?



#### Naturalness of electroweak symmetry breaking



The energy scale of new physics responsible for EWSB

TeV new physics.

Naturalness motivated

Many models, ideas.

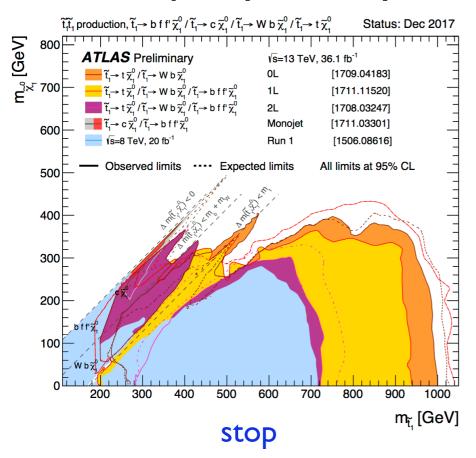
Electroweak scale, 100 GeV.  $m_h$ ,  $m_W$  ...

#### Models

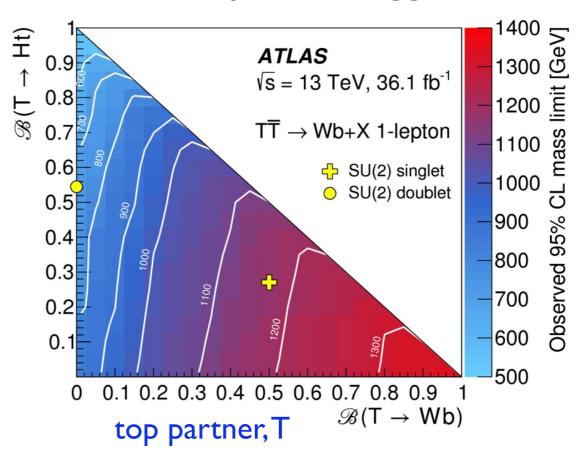
- "conventional"
  - ▶ Supersymmetry (SUSY), Composite Higgs, ...
  - A bit uncomfortable, still viable.

## All eyes on these searches

#### Supersymmetry



#### Composite Higgs



fine-tuning = comparison:

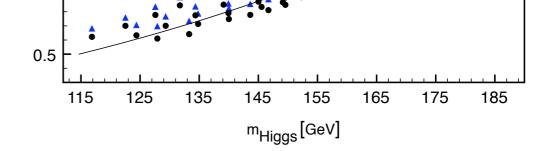
$$\frac{1}{16\pi^2}m_{\rm T}^2$$
  $vs$   $m_h^2 = (125 \text{ GeV})^2$ 

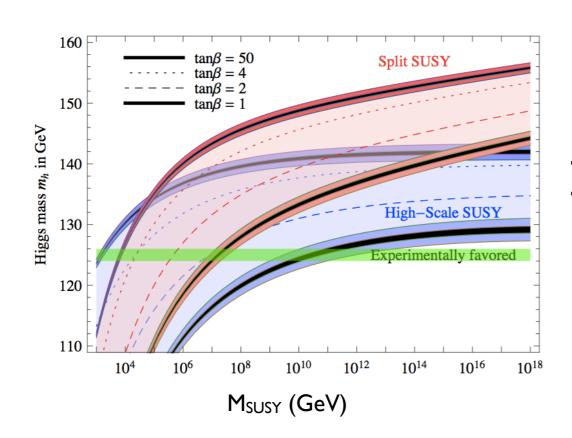
current limit:

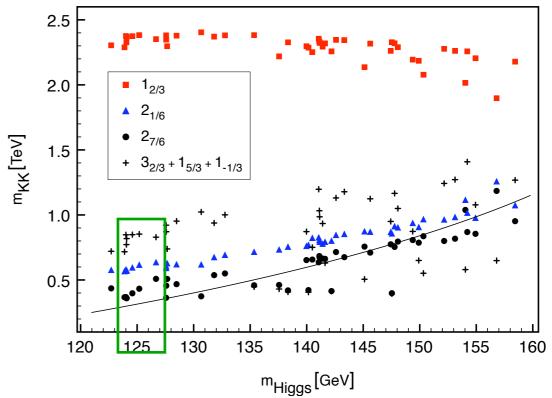
$$m_{\rm T} \sim 1 {\rm \ TeV}$$

My view: not a big problem yet.

## A confusing picture







Supersymmetry
Stop too heavy to be natural

Composite top partner too light, excluded

Such conclusions too simplistic, "work around" available. A bit uncomfortable, yes. Not time to give up just yet.

#### Models

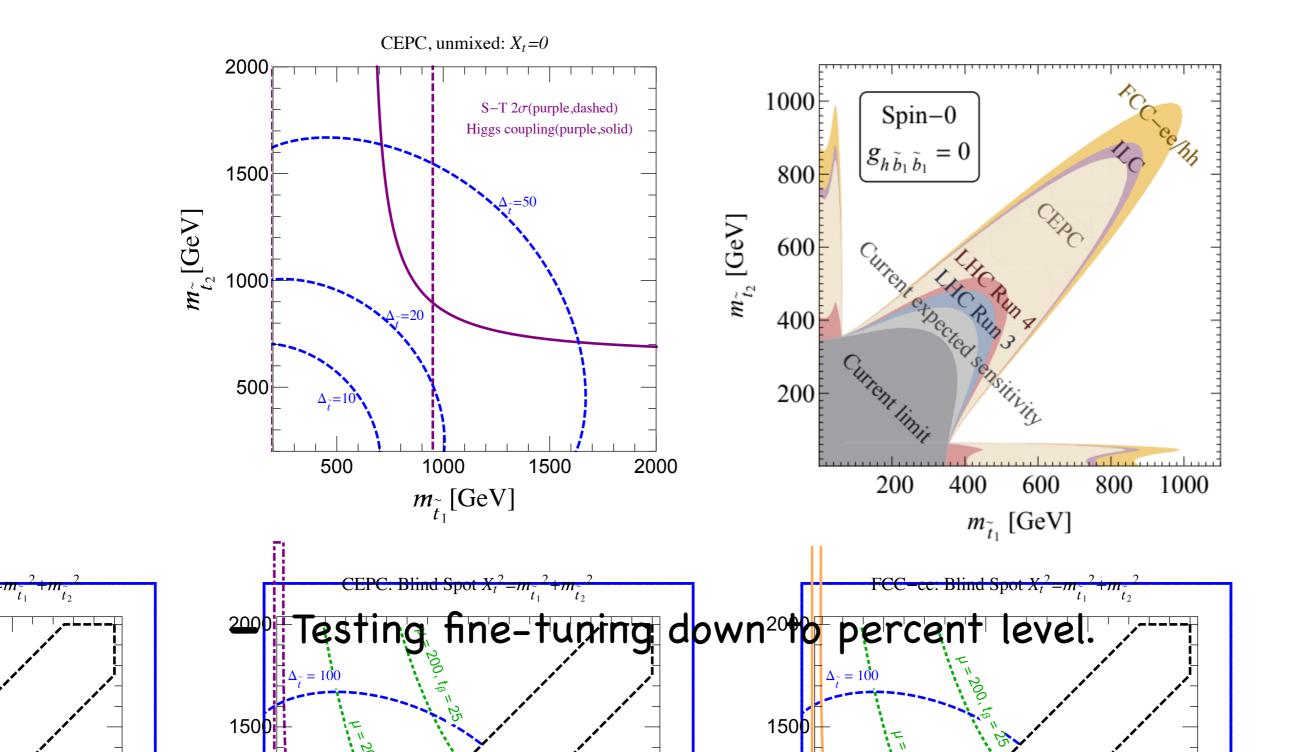
- "conventional"
  - ▶ Supersymmetry (SUSY), Composite Higgs, ...
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#### Models

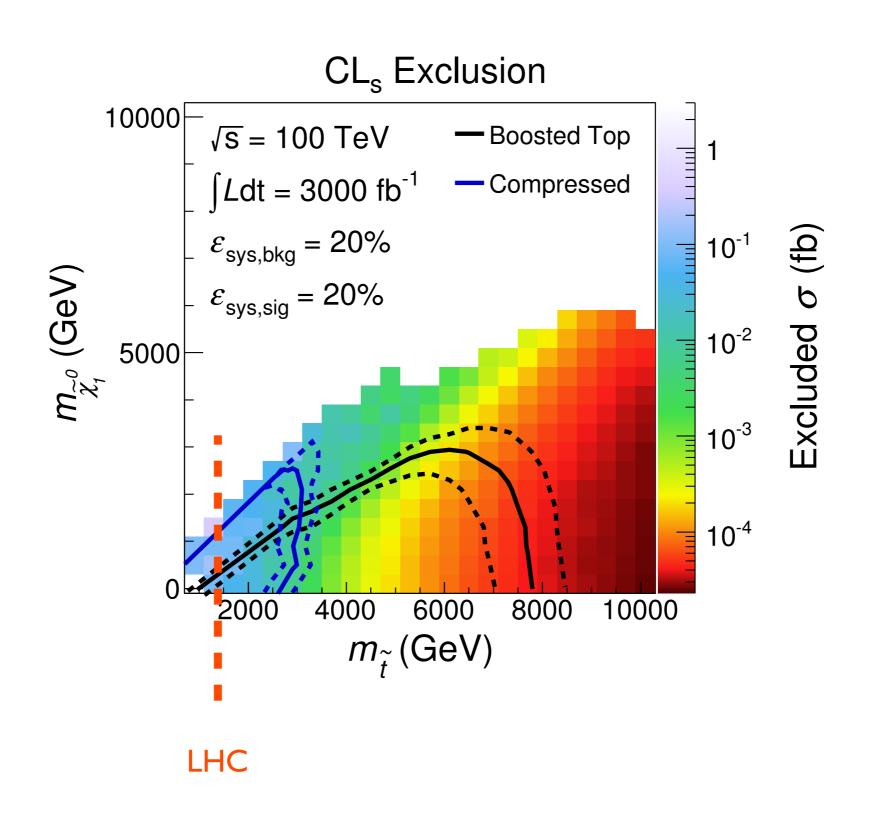
- New attempts
  - ▶ Neutral naturalness, N-naturalness, relaxion...
  - None of these is in terribly good shape. But interesting, could develop into something better.
- Why so many models?
  - Because the situation is confusing.
- Future Colliders crucial in testing them.

#### Naturalness in SUSY

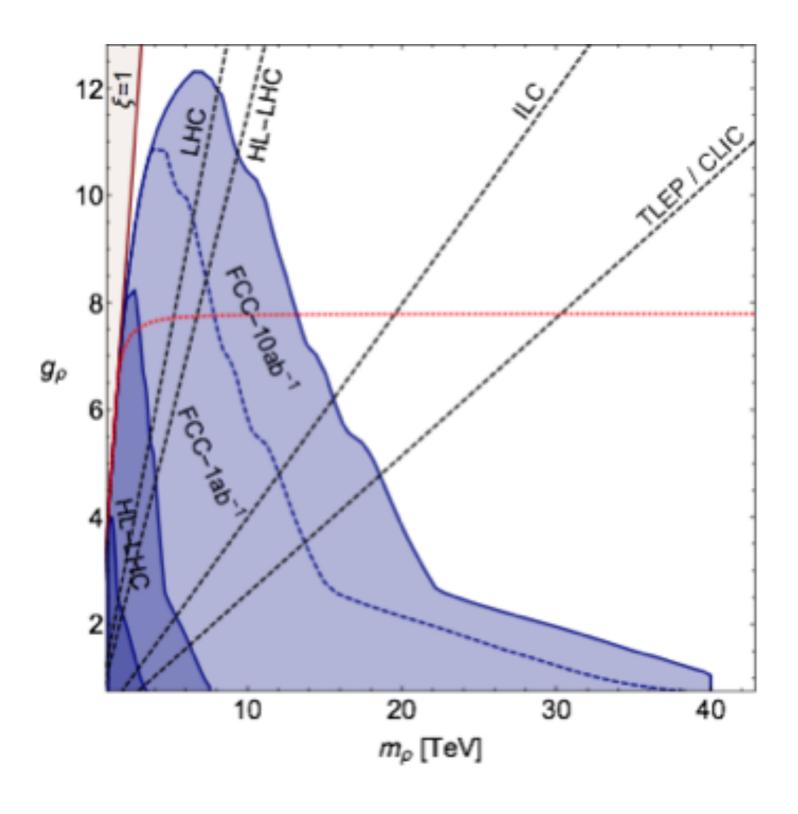
- LHC searches model dependent, many blind spots.



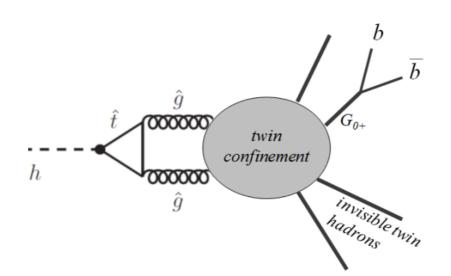
## Testing naturalness: Supersymmetry



# Composite Higgs



## Neutral naturalness: "twin"



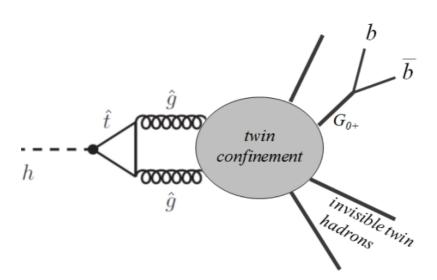
Chacko, Goh, Harnik Craig, Katz, Strassler, Sundrum

UV completion with composite Higgs: Low, Tesi, and LTW

Top partner T not colored.

Higgs decay through hidden world and back

## Neutral naturalness: "twin"

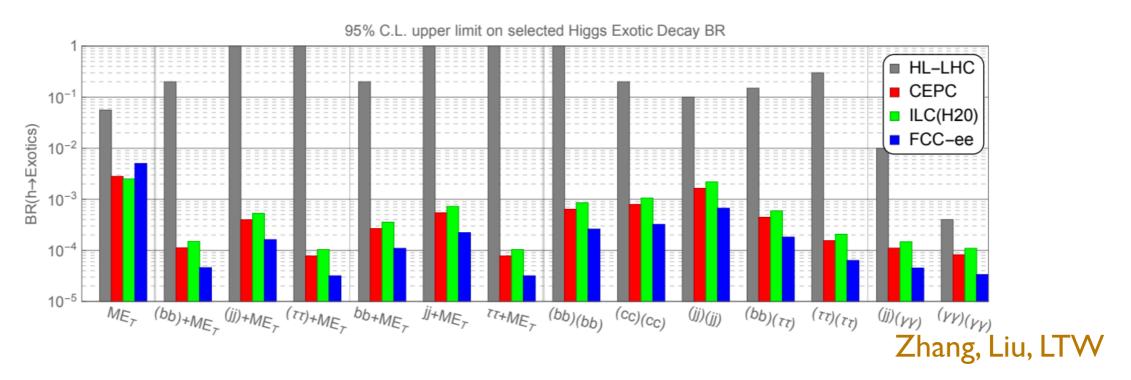


Chacko, Goh, Harnik Craig, Katz, Strassler, Sundrum

UV completion with composite Higgs: Low, Tesi, and LTW

Top partner T not colored.

Higgs decay through hidden world and back



Can be tested at LHC and Higgs factories.

Mysteries of the electroweak scale.

## Mysteries

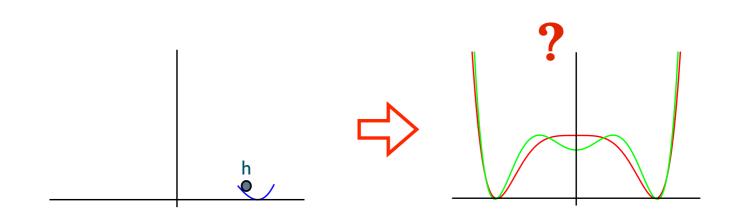


Figure 8: Question of the nature of the electroweak phase transition.

How to pre

Understanding this physics is also directly relevant to one of the most fundamental questions we can ask about any symmetry breaking phenomenon, which is what is the order of the associated phase transition. How can we experimentally decide whether the electroweak phase transition in the early universe was second order or first order? This questions it is a should be taken to the condition of the con

What does ous next step following the Higgs discovery: having understood what breaks the tries of the tries o Wednesday, August 13, the robe how electroweak symmetry is restored at thigh energies.

A first-order phase transition is also strongly motivated by the possibility

of electroweak baryogenesis [18]. While the origin of the baryon asymmetry is one of the most fascinating questions in physics, it is frustratingly straight-Is it connected wards or this empted for any organisis an unitar-tight energy scales, with no direct experimental consequences. However, we aren't forced to defer this physics to the deep ultraviolet: as is well known, the dynamics of electroweak

asymmetry?

symmetry breaking itself provides all the ingredients needed for baryogene-

sis. At temperatures far above the weak scale, where electroweak symmetry

## Mysteries

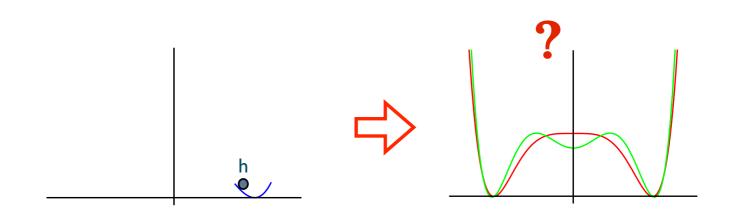


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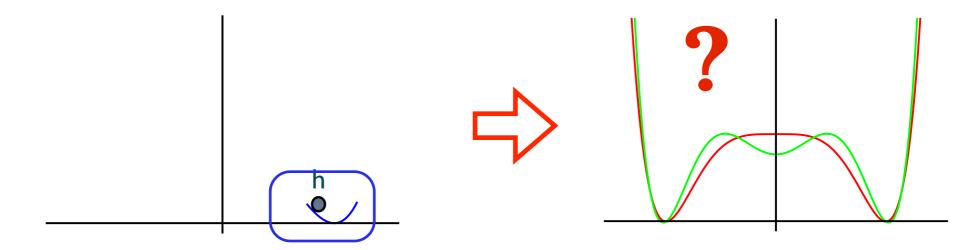
of electroweak baryogenesis [18]. While the origin of the baryon asymmetry is

What does ous next step following the Higgs discovery: having understood what breaks the standard cleent of weak symmetry, we must now indertake an experimental program to Wednesday, August 13, the robe how electroweak symmetry is restored at thigh energies.

A first-order phase transition is also strongly motivated by the possibility

one of the most fascinating questions in physics, it is frustratingly straight-Is it connected and the models for any ordenesis and the trees scales, with no direct experimental consequences. However, we aren't forced to defer this asymmetry? physics to the deep ultraviolet: as is well known, the dynamics of electroweak symmetry breaking itself provides all the ingredients needed for baryogenesis. At temperatures far above the weak scale, where electroweak symmetry

## Nature of EW phase transition



What we know from LHC LHC upgrades won't go much further

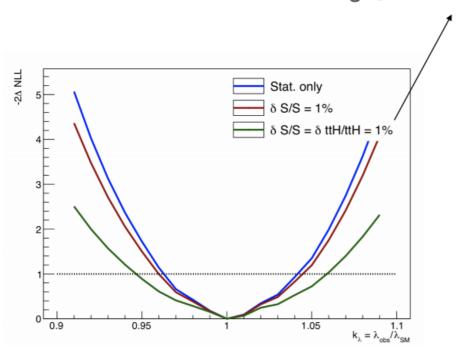
"wiggles" in Higgs potential

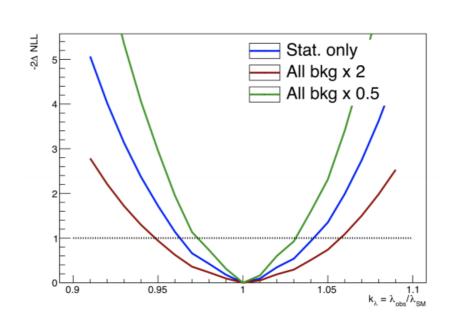
Wednesday, August 13, 14 Big difference in triple Higgs coupling

## Triple Higgs coupling at 100 TeV collider

#### Precision on the self-coupling

assuming QCD can be measured from sidebands





nominal background yields:

$$\delta \kappa_{\lambda}(\text{stat}) \approx 3.5 \%$$
  
 $\delta \kappa_{\lambda}(\text{stat + syst}) \approx 6 \%$ 

varying (0.5x-2x) background yields:

 $\delta \kappa_{\lambda}(\text{stat}) \approx 3 - 5 \%$ 

Talk by Michele Selvaggi at 2nd FCC physics workshop

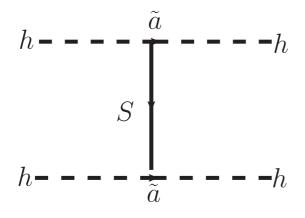
## But, there should be more

$$V(h) = \frac{m^2}{2}h^2 + \lambda h^4 + \frac{1}{\Lambda^2}h^6 + \dots$$

- 1st order EW phase transition means there is new physics close to the weak scale.
- Can be difficult to discover at the LHC.
- Will leave more signature in Higgs coupling.

# For example

$$m^2h^{\dagger}h + \tilde{\lambda}(h^{\dagger}h)^2 + m_S^2S^2 + \tilde{a}Sh^{\dagger}h + \tilde{b}S^3 + \tilde{\kappa}S^2h^{\dagger}h + \tilde{h}S^4$$





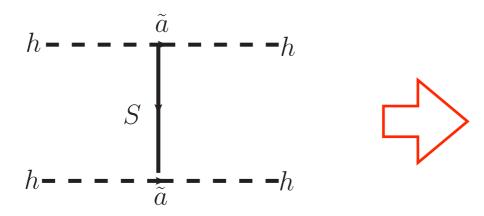
$$\frac{c}{m_S^2} \left( h^\dagger \partial h \right)^2$$

shift in h-Z coupling

$$\delta_{Zh} \sim c \frac{v^2}{m_S^2}$$

# For example

$$m^2h^{\dagger}h + \tilde{\lambda}(h^{\dagger}h)^2 + m_S^2S^2 + \tilde{a}Sh^{\dagger}h + \tilde{b}S^3 + \tilde{\kappa}S^2h^{\dagger}h + \tilde{h}S^4$$



$$\frac{c}{m_S^2} \left( h^\dagger \partial h \right)^2$$

### shift in h-Z coupling

$$\delta_{Zh} \sim c \frac{v^2}{m_S^2}$$

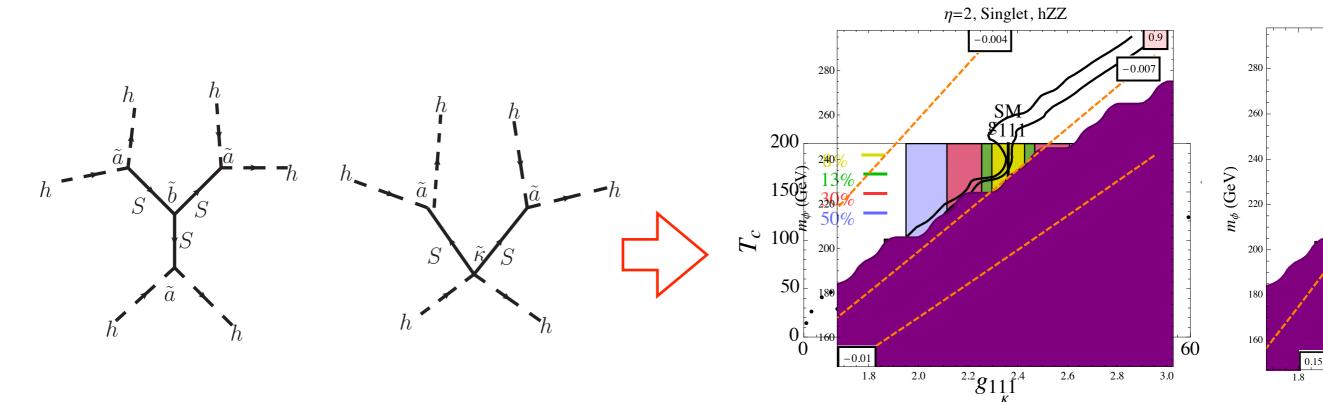
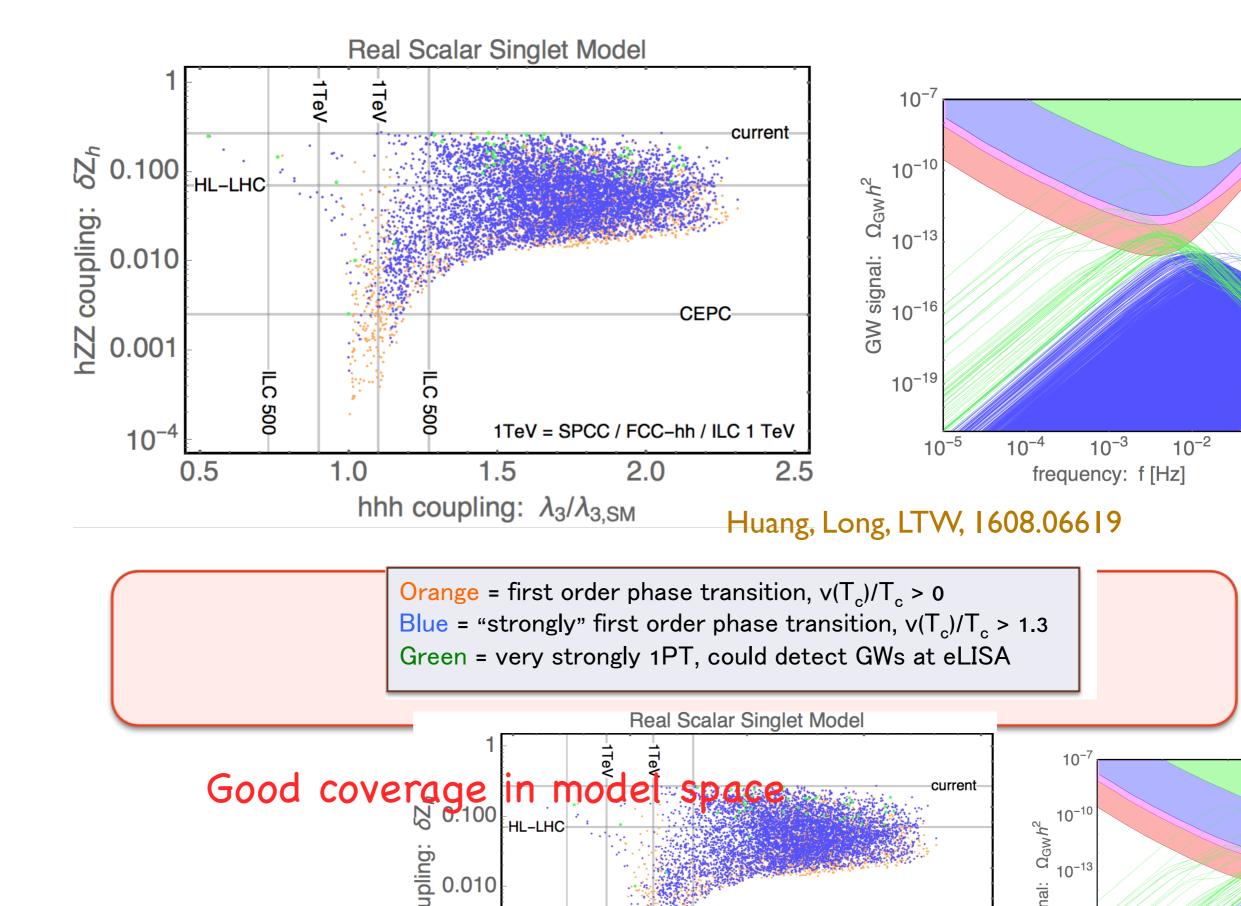
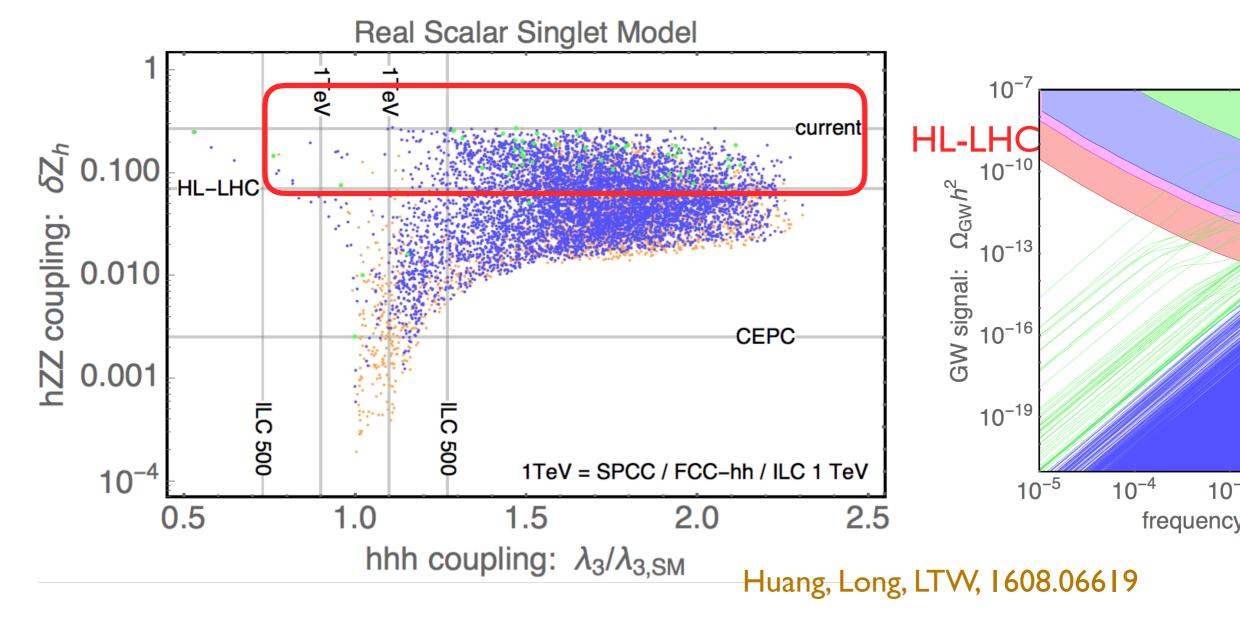


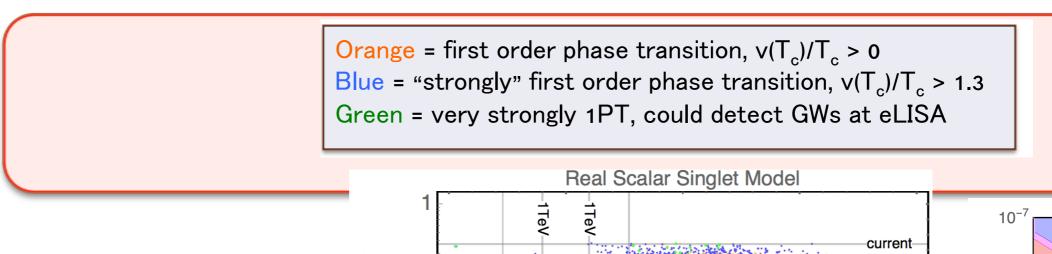
Figure 6. The region exparameter space where a strongly Singlet benchmark model. Also shown are the fractions

# Probing EWSB at higgs factories



# Probing EW phase transition





### Conclusion

- LHC still has a lot to say.
  - ▶ 15+ years of operation, 95+% of data to come.
  - Need to think about how to new searches with this data.
- Beyond the LHC, we need future colliders to address the open questions of the Standard Model.

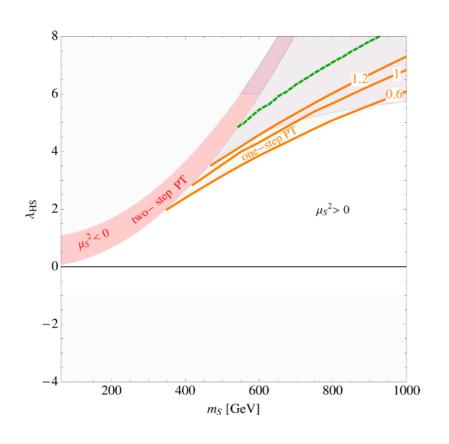
# extra

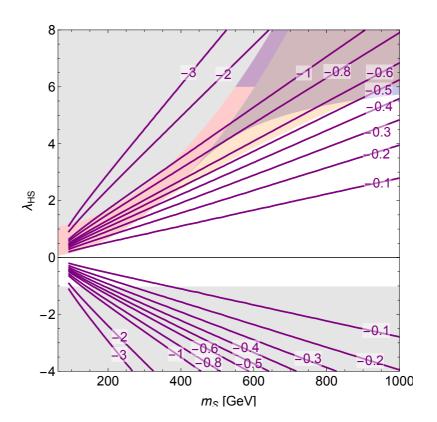
# Nightmare scenario:

Meade et al

Singlet model with a  $Z_2$   $S \rightarrow -S$ 

$$S \rightarrow -S$$





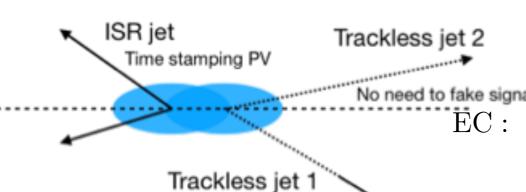
h<sup>6</sup> term generated at 1-loop order Only marginally visible.

### Late comers will be spotted easily:

	$L_{T_2}$	$L_{T_1}$	Trigger	$\epsilon_{ m trig}$	$\epsilon_{ m sig}$	$\epsilon_{ m fake}^{j}$	Ref.
EC	1.17 m	0.2 m	DelayJet	0.5	0.5	$10^{-3}$	[12]
MS	10.6 m	4.2 m	MS RoI	0.25,  0.5	0.25	$5 \times 10^{-9}$	[24]

CMS timing module ATLAS MS LLP search

(without timing



Pile-Up background, time spread 190 ps (beam property)

No need to fake signal 
$$ext{EC}: N_{ ext{bkg}}^{ ext{PU}} = \sigma_{ ext{j}} \mathcal{L}_{ ext{int}} \epsilon_{ ext{trig}}^{ ext{EC}} \left( \bar{n}_{ ext{PU}} rac{\sigma_{ ext{j}}}{\sigma_{ ext{inc}}} \epsilon_{ ext{fake}}^{j, ext{EC}} f_{ ext{nt}}^{j} 
ight) pprox 2 imes 10^{7},$$

MS: 
$$N_{\text{bkg}}^{\text{PU}} = \sigma_{\text{j}} \mathcal{L}_{\text{int}} \epsilon_{\text{trig}}^{\text{MS}} \left( \bar{n}_{\text{PU}} \frac{\sigma_{\text{j}}}{\sigma_{\text{inc}}} \epsilon_{\text{fake}}^{j,\text{MS}} f_{\text{nt}}^{j} \right) \approx 50, (5)$$

Pile-up BKG: intrinsic resolution

Fake displaced obj

~190 ps

EC (30ps) cut:  $\Delta t > 1$  ns

BKG(EC-PU)  $\sim 1.3$ 

MS (30ps) cut:  $\Delta t > 0.4$  ns

BKG(MS-PU) ~ 0.86

The detector time resolution for

MS can be downgraded to

hundreds of ps

**LPC TOTW** 

MS (200ps) cut:  $\Delta t > 1$ ns

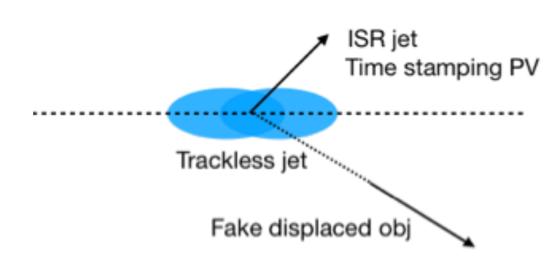
BKG(MS-PU) << 1



### Late comers will be spotted easily:

	$L_{T_2}$	$L_{T_1}$	Trigger	$\epsilon_{ m trig}$	$\epsilon_{ m sig}$	$\epsilon_{ m fake}^{j}$	Ref
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CMS timing module ATLAS MS LLP search



Same-vertex hard scattering background, time spread 30 ps (precision timing)

EC:  $N_{\rm bkg}^{\rm SV} = \sigma_{\rm j} \mathcal{L}_{\rm int} \epsilon_{\rm trig}^{\rm EC} \epsilon_{\rm fake}^{\rm j,EC} \approx 1 \times 10^{11}$ 

MS:  $N_{\rm bkg}^{\rm SV} = \sigma_{\rm j} \mathcal{L}_{\rm int} \epsilon_{\rm trig}^{\rm MS} \epsilon_{\rm fake}^{\rm j,MS} \approx 4 \times 10^5$ ,

resolution ~30 ps

EC (30ps) cut:  $\Delta t > 0.4 \text{ ns}$ 

MS (30ps) cut:  $\Delta t > 1$ ns

BKG(SV) << 1

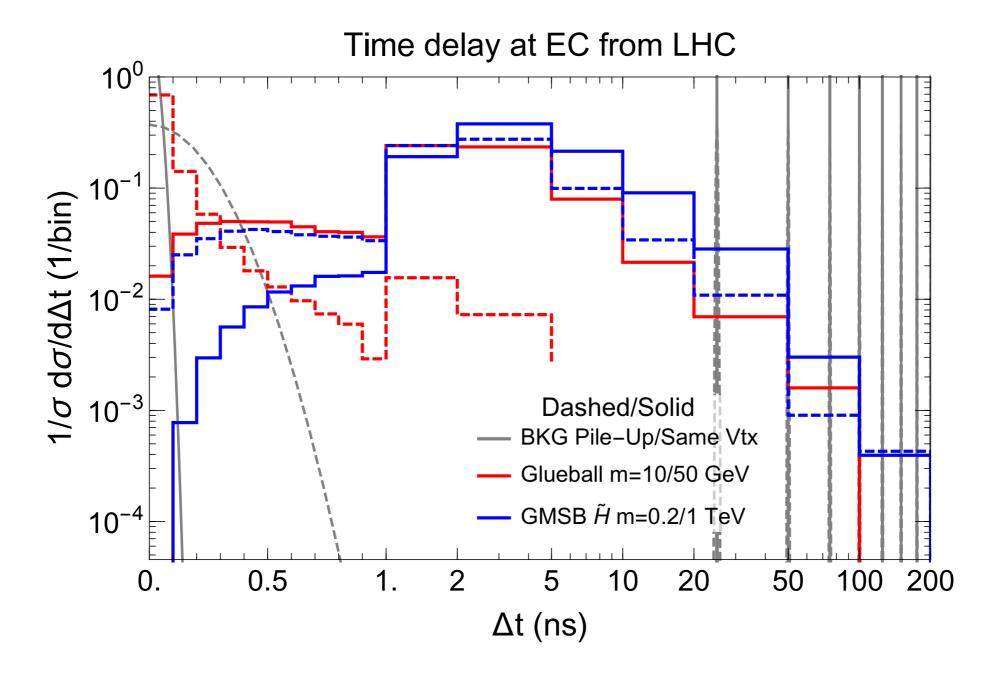
Hard collision BKG: detector time The detector time resolution for MS can be downgraded to hundreds of ps MS (200ps) cut:

 $\Delta t > 1$ ns

BKG(MS-SV)  $\sim 0.11$ 



### Search based on EC

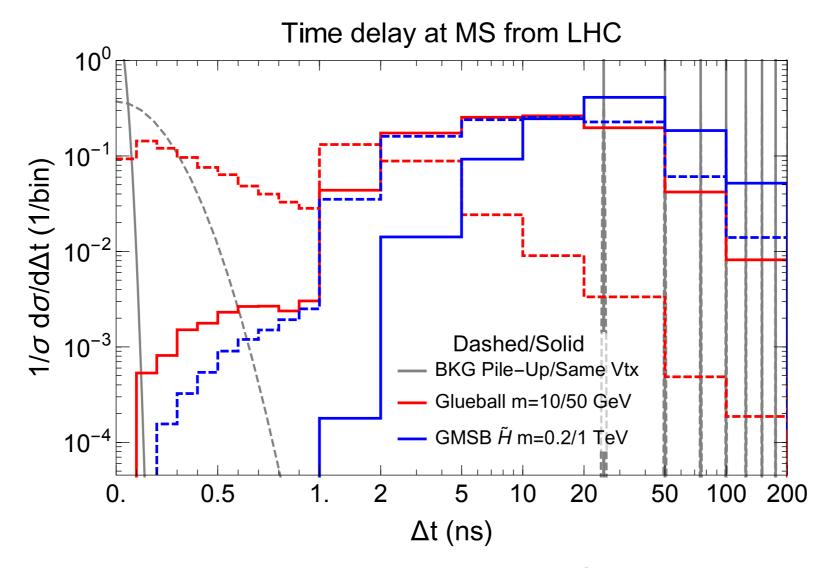


After timing cut:  $\Delta t > 0.8 \text{ ns}$ 

Back ground dominated by pile up

$$\#_{\text{background}} \sim 1$$

### Search based on MS



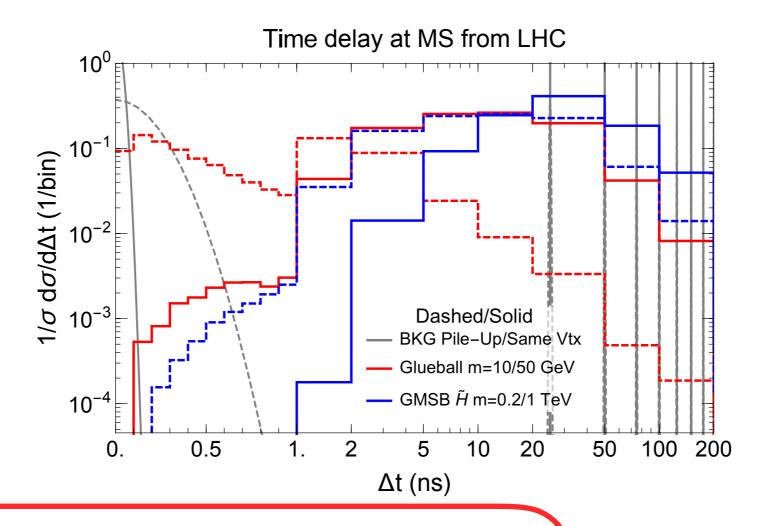
Pile up background smaller, shielded by HCAL etc.

Before timing cut:  $\sim 50$ 

After timing cut:  $\Delta t > 0.4 \text{ ns} \text{ } \#_{\text{background}} \sim 1$ 

Further away, larger  $\Delta t$  for signal.

### Search based on MS



Pile up background smaller, shielded by HCAL etc.

$$\Delta t > 0.4 \text{ ns} \quad \#_{\text{background}} \sim 1$$

Further away, larger  $\Delta t$  for signal.

no need for super good timing resolution

$$\delta t \sim 200 \text{ ps}$$
 will do

# Signal ISR jet (time stamp) SM ISR jet (cinted stamp) SM X X SM X X S

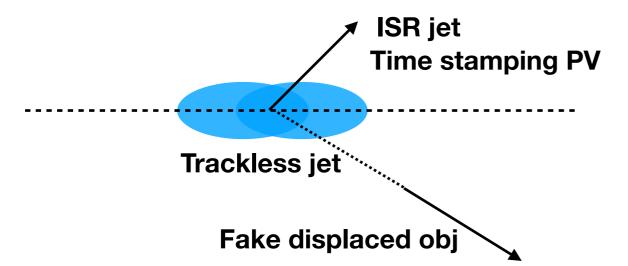
ISINGR jet providers the time for the hard collision

2. LLP decay before reaching timing layer.



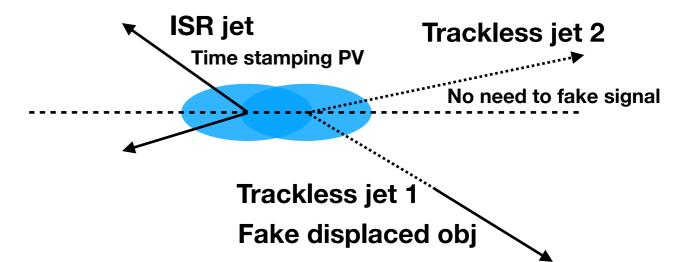
# background

Same hard interaction



Time delay from resolution of timing detector.

Pile up



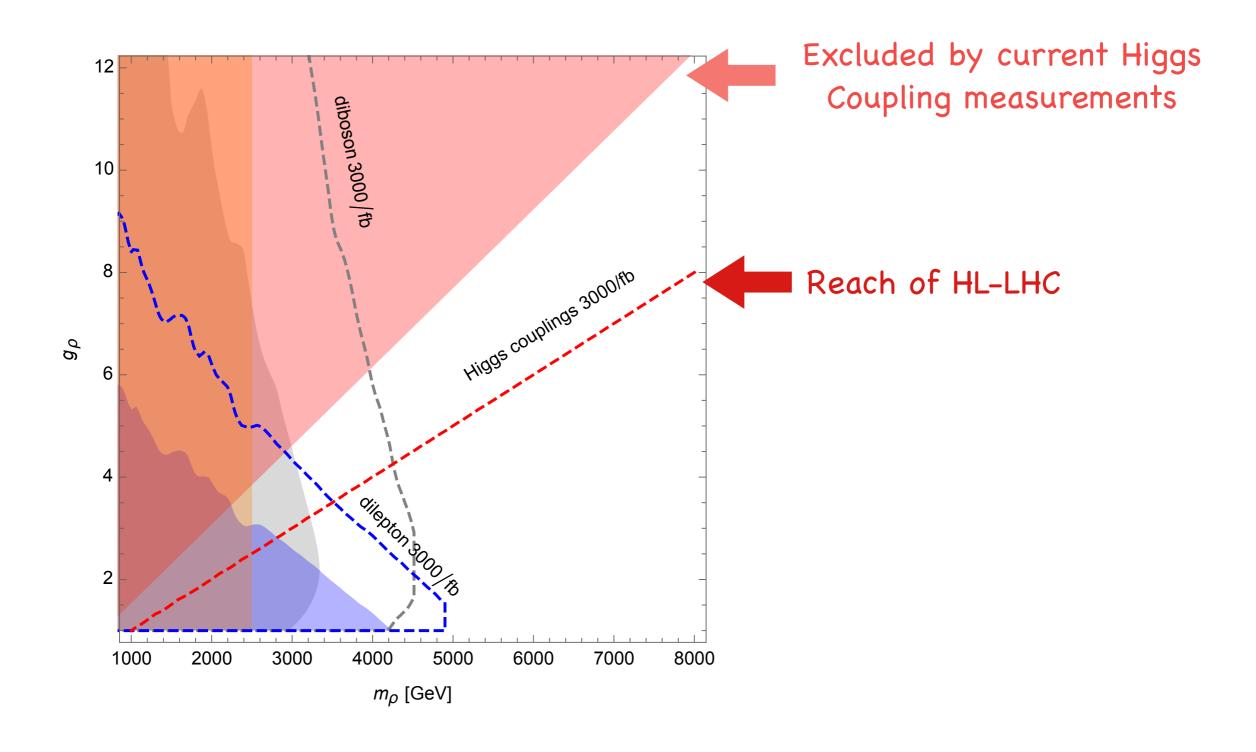
Time delay from spread of the proton bunch

~ 190 ps

# Importance of precision measurement

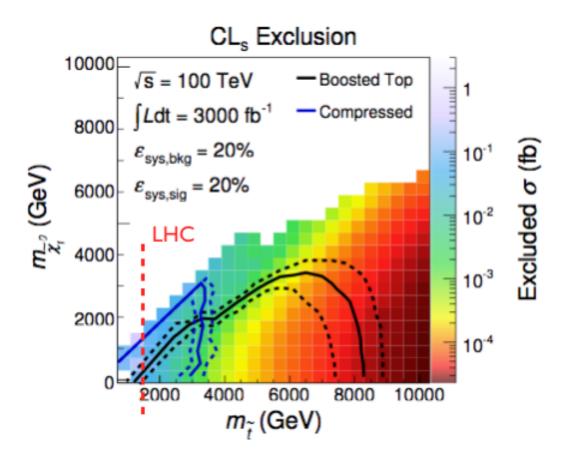
- No clear indication where new physics might be.
  - Precision measurement can give crucial guidance.
- Lots of data still to come
  - Room to improve! Statistics and systematics.
- Will be a important part of the legacy of the LHC.
  - LEP taught us a lot. LHC will do the same.

# Higgs coupling vs direct search

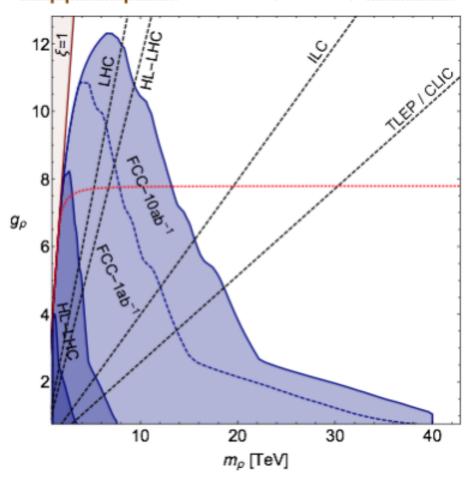


### Testing naturalness at 100 TeV pp collider





### Pappadopulo, Thamm, Torre, Wulzer, 2014



Fine tuning:  $(M_{NP})^{-2}$ 

# Why is Higgs measurement crucial?

- Naturalness is the most pressing question of EWSB.
  - How should we predict the Higgs mass?
- We may not have the right idea. No confirmation of any of the proposed models.
- Need experiment!
- Fortunately, with Higgs, we know where to look.
- And, the clue to any possible way to address naturalness problem must show up in Higgs coupling measurement.