

Precision Frontier at the LHC

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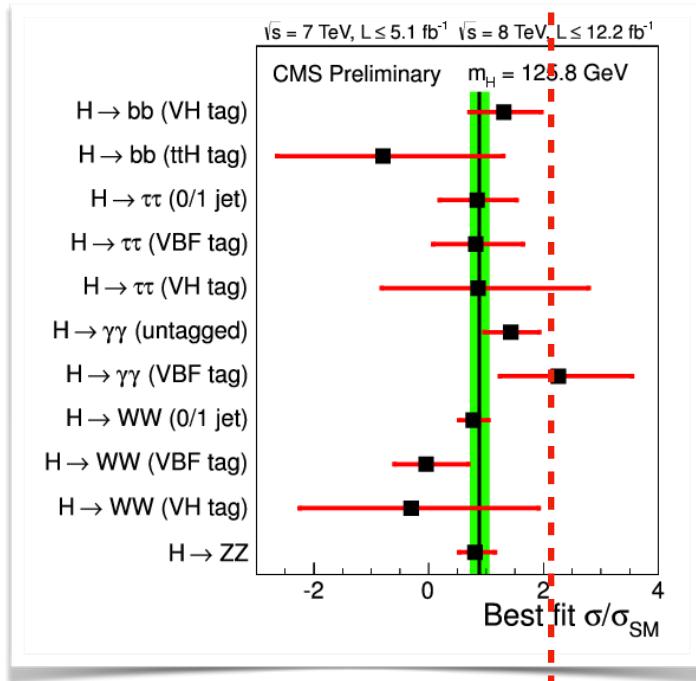


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Outlines

- Why precision?
- Very very very brief overview of IR subtraction/slicing schemes
- Phenomenology - fixed order and resummation

Why Precision?



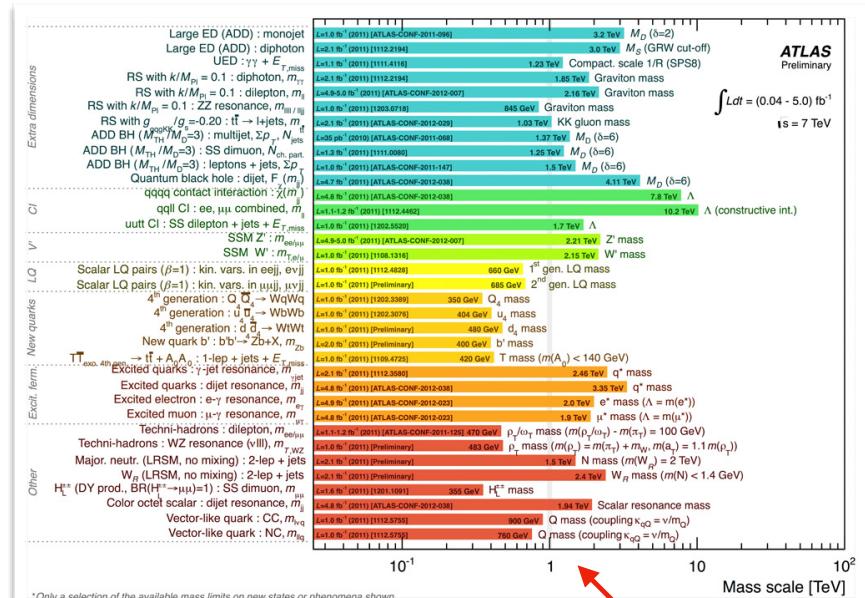
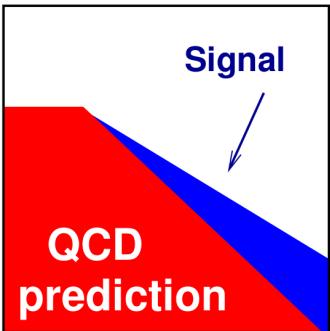
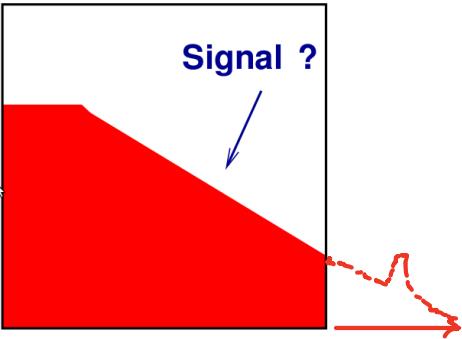
LO @ 8 TeV: $9.6 +/- 25\% \text{ pb}$

N3LO @ 8 TeV: $19.47 + .32\% - 2.99\% \text{ pb}$

Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15

If no higher order prediction is known

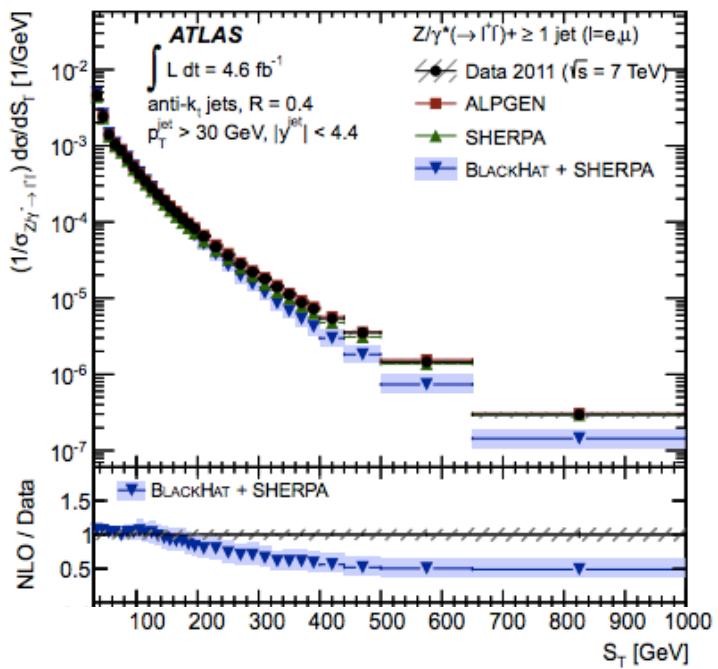
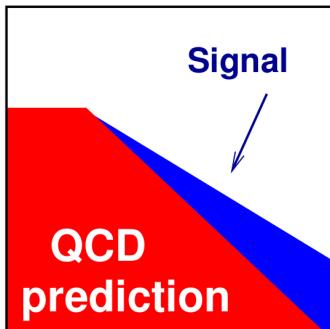
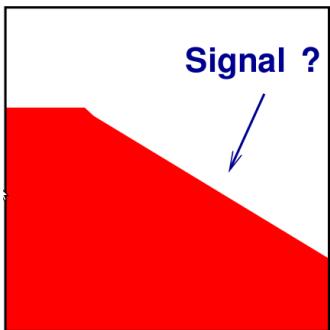
Why Precision?



- Lower limits above 1 TeV now.
- New physics could be out of reach for current LHC
- Predicting the shape correctly will be crucial

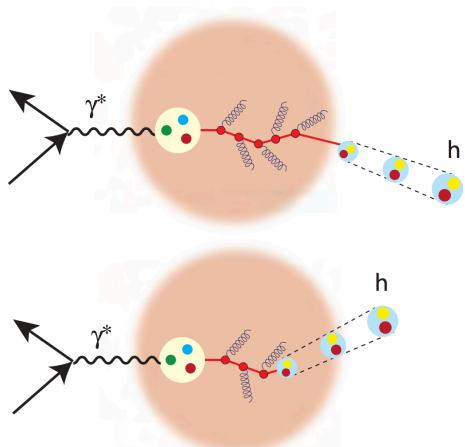
New Physics = Data - SM predictions!!

Why Precision?



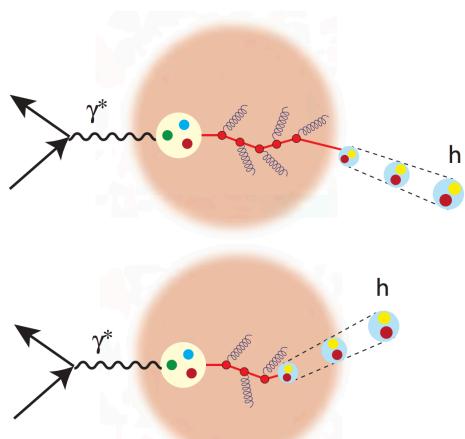
50% difference at high S_T
New Physics? Missing Higher order?

Why Precision?



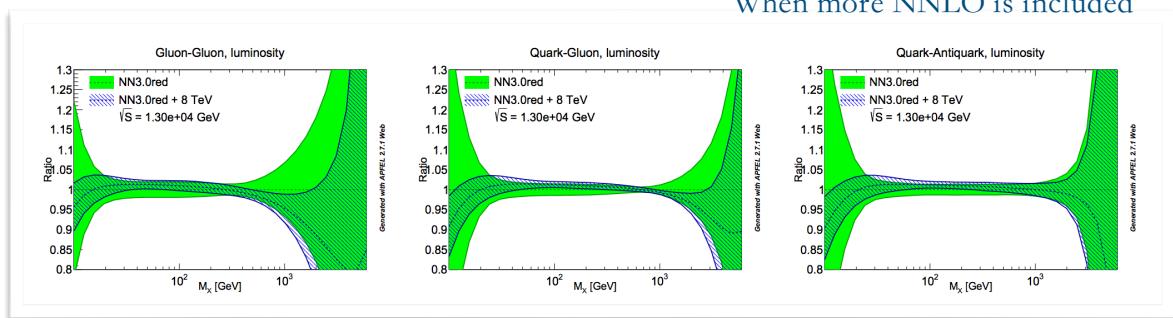
- QCD itself is interesting, Heavy Ion, EIC ...
- Hard probe (perturbative) of nuclei internal structure (non-perturbative)
- To extract the non-perturbative information, we should know the hard part as accurate as possible

Why Precision?



- QCD itself is interesting, Heavy Ion, EIC ...
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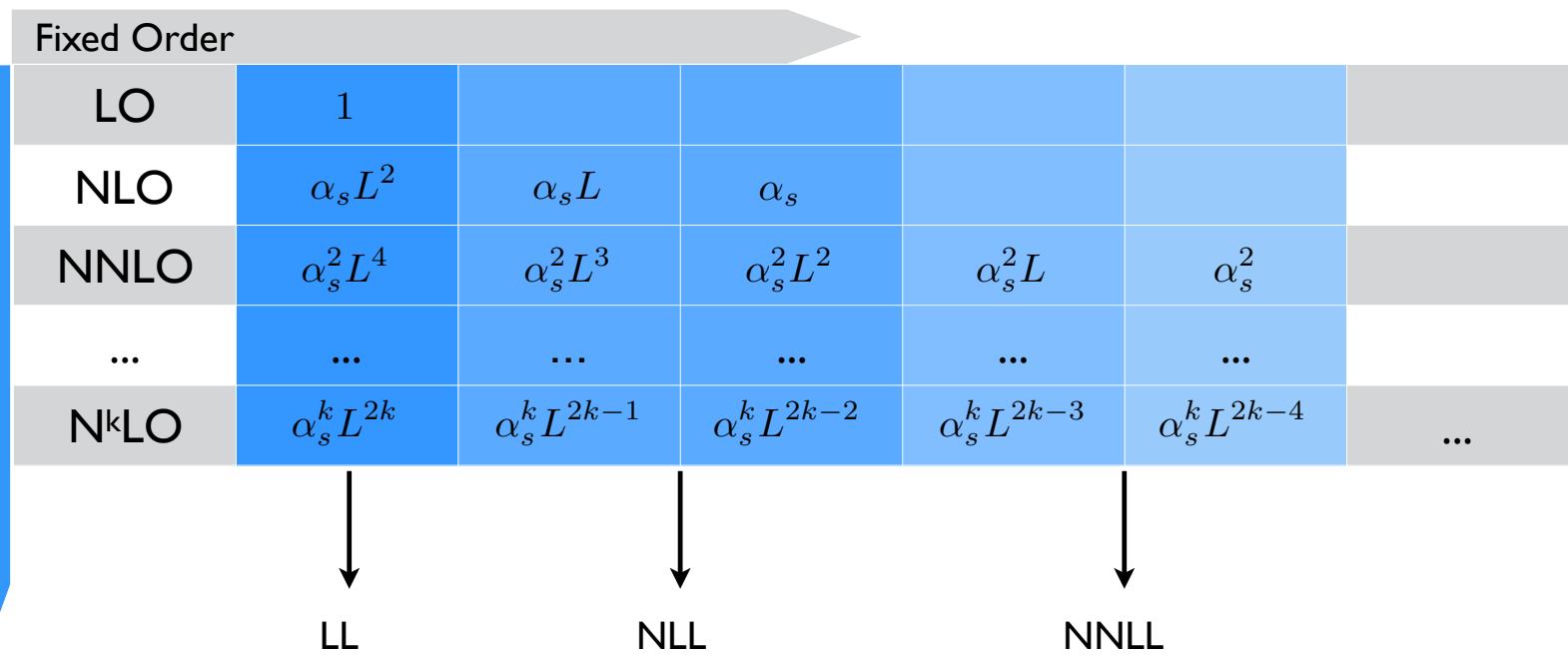
NNPDF30 extractions
When more NNLO is included



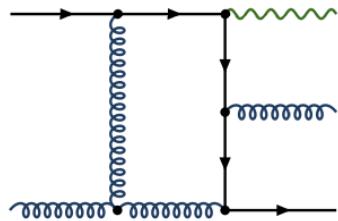
Will see later with more details

Precision (NLO, NNLO, N3LO ... + Res.) is crucial

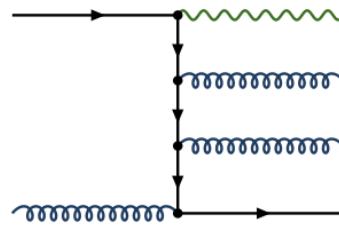
Resummation



Not Just Loops



$$\int d\Phi_3 \sum_{i=0}^2 \frac{vr_i}{\epsilon^i}$$

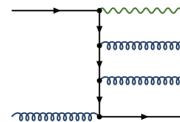


$$\int d\Phi_4 rr$$

- We are mostly interested in fully differential cross section
- IR poles occur in real emissions
- IR poles fully show up for degenerate (soft/collinear) states
ONLY after integrating over phase space, with all kinds of
exp. cuts, jet algorithms ...
- How to isolate for numerical evaluation?

Not Just Loops

Subtraction


$$= \int dz \frac{f(z)}{z^{1+a\epsilon}}$$

$$\int dz \frac{f(z) - f(0)}{z} + \int dz z^{-1-a\epsilon} f(0)$$

Construct counter terms point-wise in the phase space

- Antenna subtraction 2 Gehrmann, Glover
- STRIPPER + modifications Czakon + ...
- Projection to Born Cacciari, et. al.

Slicing

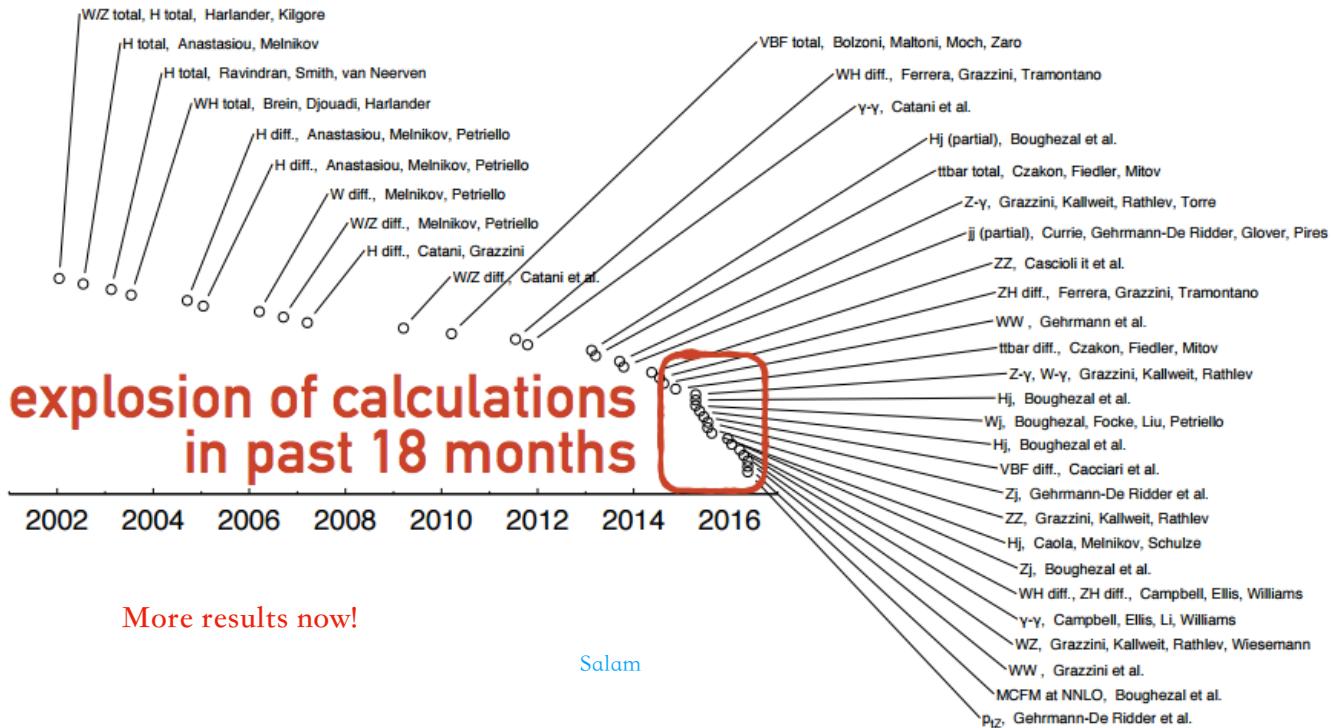
$$\int \frac{f(z)}{z} \theta(z > z_0) - f(0) \frac{z_0^{-a\epsilon}}{a\epsilon} + \dots$$

...

A physical observable (z_0) to regulate all related IR singularities

- qT-subtraction Catani, Grazzini
- Inclusive jet mass Gao, Li, Zhu
- N-jettiness subtraction Boughezal, et. al.

Phenomenology



Phenomenology

process	known	desired
$pp \rightarrow H$	σ $N^3LO_{HEFT} + N^2LO_{QCD}\left(\frac{1}{m_t^6}\right)$	$d\sigma$ $N^3LO_{HEFT} + N^2LO_{QCD}$ + $NLO_{EW} + N^{(1,1)}LO_{QCD \times EW}$
	$d\sigma$ NLO_{EW}	
	$d\sigma$ $N^2LO_{HEFT} + NLO_{QCD} + PS$	
$pp \rightarrow H + j$	$d\sigma$ N^2LO_{HEFT}	$d\sigma$ $N^2LO_{HEFT} + NLO_{QCD} + NLO_{EW}$
	$d\sigma$ NLO_{EW}	
$pp \rightarrow H + 2j$	$d\sigma$ $NLO_{HEFT} + LO_{QCD}$	$d\sigma$ $N^2LO_{HEFT} + LO_{QCD} + NLO_{EW}$ $d\sigma$ $N^2LO_{QCD}(VBF^*)$ $d\sigma$ $NLO_{EW}(VBF)$
	$d\sigma$ $N^2LO_{QCD}(VBF^*)$	
	$d\sigma$ $NLO_{EW}(VBF)$	
$pp \rightarrow H + 3j$	$d\sigma$ NLO_{HEFT}	$d\sigma$ $NLO_{QCD} + LO_{QCD} + NLO_{EW}$
	$d\sigma$ NLO_{EW}	
$pp \rightarrow H + V$	$d\sigma$ N^2LO_{QCD}	$d\sigma$ $N^2LO_{QCD} + NLO_{gg \rightarrow HV} + NLO_{EW}$
$pp \rightarrow HH$	$d\sigma$ N^2LO_{HEFT}	$d\sigma$ $N^2LO_{HEFT} + NLO_{QCD} + NLO_{EW}$
	$d\sigma$ NLO_{QCD}	
$pp \rightarrow H + t\bar{t}$	$d\sigma$ NLO_{QCD}	$d\sigma$ $NLO_{QCD} + NLO_{EW}$
	$d\sigma$ NLO_{EW}	
$pp \rightarrow H + t$	$d\sigma$ NLO_{QCD}	$d\sigma$ $NLO_{QCD} + NLO_{EW}$
$pp \rightarrow H + \bar{t}$		

More results now!

Phenomenology

process	known		desired
$pp \rightarrow V$	σ	$N^3LO_{QCD}(z \rightarrow 0)$	
	$d\sigma$	N^2LO_{QCD}	$d\sigma$
	$d\sigma$	NLO_{EW}	$+ N^{(1,1)}LO_{QCD \times EW} + \text{decays}$
	$d\sigma$	$N^{(1,1)}LO_{QCD \times EW}$	
$pp \rightarrow VV'$	$d\sigma$	$N^2LO_{QCD} + \text{decays}$	$d\sigma$
	$d\sigma$	NLO_{EW}	$+ N^2LO_{QCD} + NLO_{EW} + \text{decays}$
$pp \rightarrow V + j$	$d\sigma$	N^2LO_{QCD}	$d\sigma$
$pp \rightarrow V + 2j$	$d\sigma$	$NLO_{QCD} + \text{decays}$	$d\sigma$
	$d\sigma$	$NLO_{EW} + \text{decays}$	$+ N^2LO_{QCD} + NLO_{EW} + \text{decays}$
$pp \rightarrow VV' + 1, 2j$	$d\sigma$	$NLO_{QCD} + \text{decays}$	$d\sigma$
	$d\sigma$	NLO_{EW}	$+ NLO_{QCD} + NLO_{EW} + \text{decays}$
$pp \rightarrow VV'V''$	$d\sigma$	NLO_{QCD}	$d\sigma$
	$d\sigma$	NLO_{EW}	$+ NLO_{QCD} + NLO_{EW} + \text{decays}$
$pp \rightarrow \gamma\gamma$	$d\sigma$	N^2LO_{QCD}	$d\sigma$
$pp \rightarrow \gamma\gamma + j$	$d\sigma$	NLO_{QCD}	$d\sigma$
			$+ N^2LO_{QCD} + NLO_{EW}$

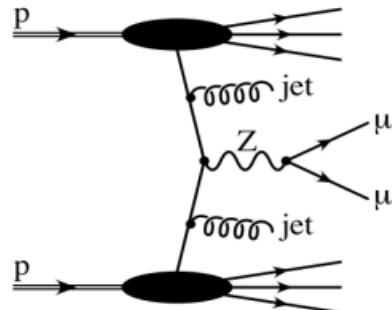
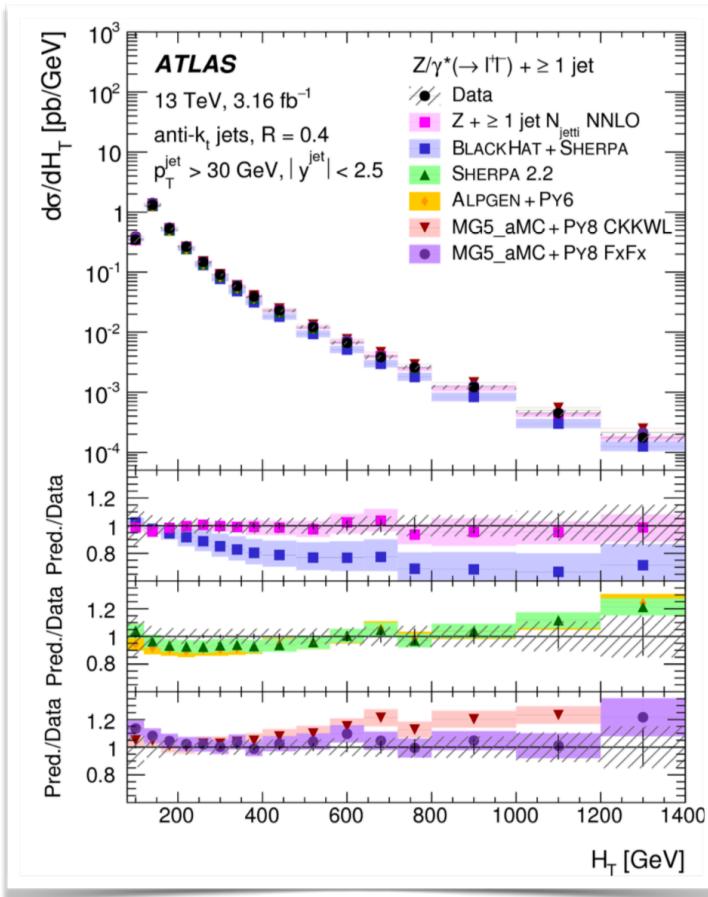
More results now!

Phenomenology

process	known	desired
$pp \rightarrow t\bar{t}$	$d\sigma$ N ² LO _{QCD}	
	$d\sigma$ NLO _{QCD} + decays	$d\sigma$ N ² LO _{QCD} + NLO _{EW} + decays
	$d\sigma$ NLO _{EW}	
$pp \rightarrow t\bar{t} + j$	$d\sigma$ NLO _{QCD} + decays	$d\sigma$ N ² LO _{QCD} + NLO _{EW} + decays
	$d\sigma$ NLO _{EW}	
$pp \rightarrow t\bar{t} + 2j$	$d\sigma$ NLO _{QCD} + on-shell decays	$d\sigma$ NLO _{QCD} + NLO _{EW} + decays
$pp \rightarrow t\bar{t} + V$	$d\sigma$ NLO _{QCD}	$d\sigma$ NLO _{QCD} + NLO _{EW} + decays
	$d\sigma$ NLO _{EW}	
$pp \rightarrow t/\bar{t}$	$d\sigma$ N ² LO _{QCD} (<i>t</i> -channel)	$d\sigma$ N ² LO _{QCD} + NLO _{EW} + decays
$pp \rightarrow 2j$	$d\sigma$ N ² LO _{QCD} (gg,qq)	$d\sigma$ N ² LO _{QCD} + NLO _{EW}
	$d\sigma$ NLO _{EW}	
$pp \rightarrow j + \gamma$	$d\sigma$ NLO _{QCD}	$d\sigma$ N ² LO _{QCD} + NLO _{EW}
	$d\sigma$ NLO _{EW}	
$pp \rightarrow 3j$	$d\sigma$ NLO _{QCD}	$d\sigma$ N ² LO _{QCD} + NLO _{EW}

More results now!

Phenomenology



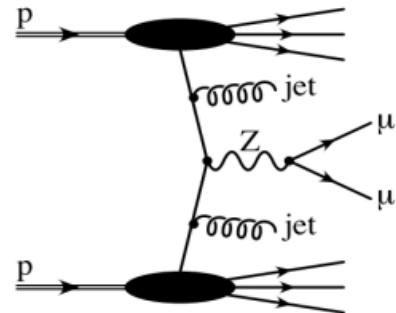
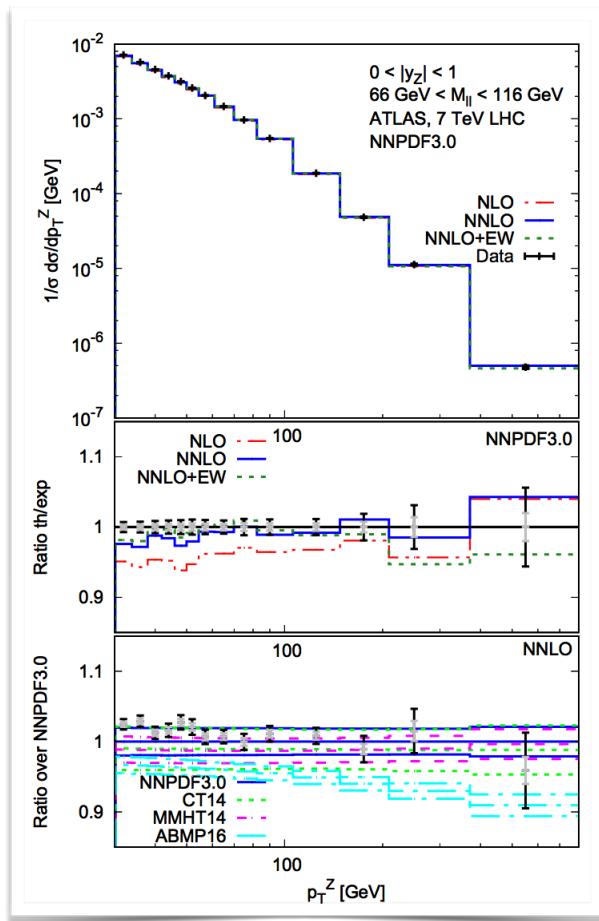
$Z + \text{jets } H_T$ distribution

- common variable in beyond SM searches for heavy particles
- Irreducible background for new physics searches, e.g. dark matter, SUSY

NNLO recovers agreement with data by add on missing high orders.

Boughezal, et. al.

Phenomenology



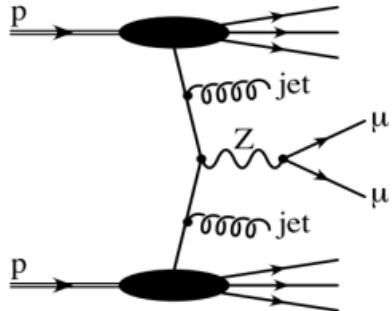
Z pT distribution

- Clean and very small exp. uncertainty $\sim 1\%$
- therefore standard candle at the LHC

NNLO agrees much better than NLO.

Boughezal, et. al.

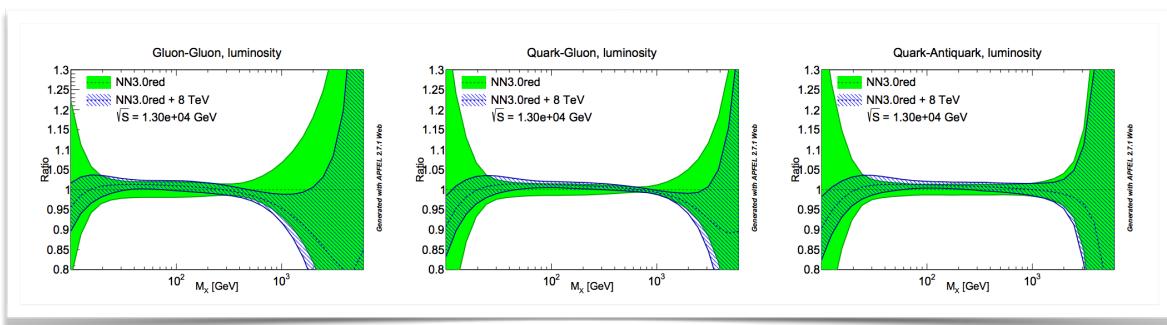
Phenomenology



	Before p_T^Z data	After p_T^Z data
$\sigma_{gg \rightarrow H}$ [pb]	48.22 ± 0.89 (1.8%)	48.61 ± 0.61 (1.3%)
σ_{VBF} [pb]	3.92 ± 0.06 (1.5%)	3.96 ± 0.04 (1.0%)

Z pT distribution

- Clean and very small exp. uncertainty ~ 1%
- therefore standard candle at the LHC



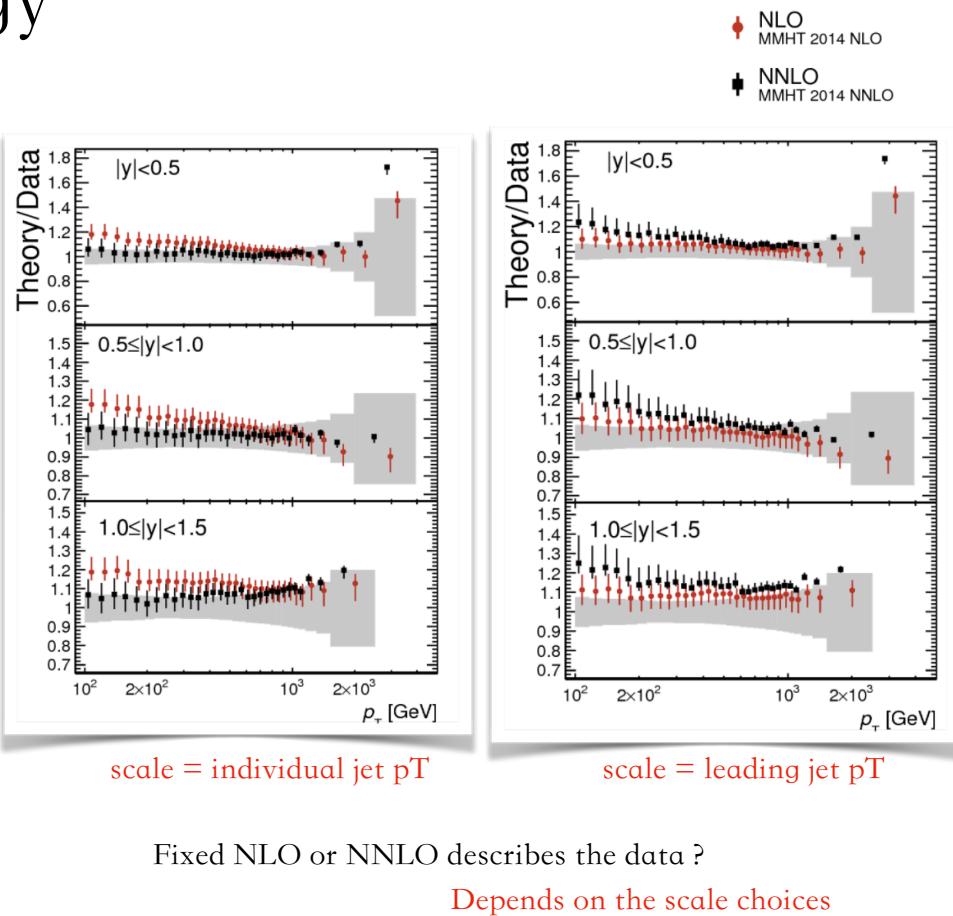
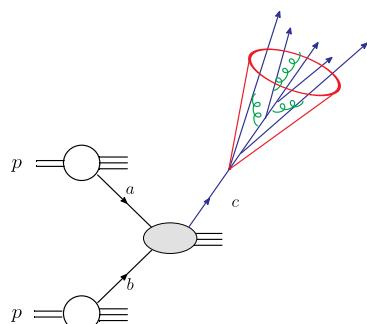
Error reduced by 30% when NNLO Z pT included

Boughezal, et. al.

Phenomenology

Single inclusive jet production

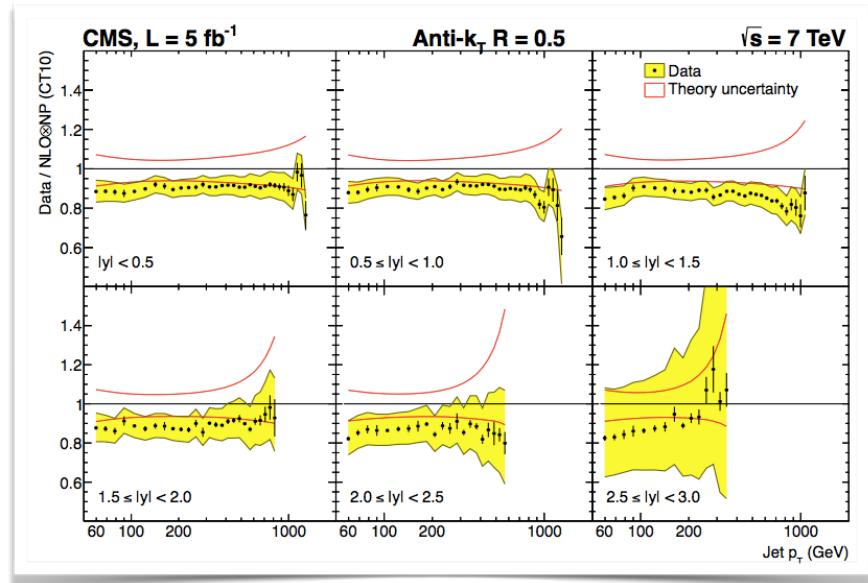
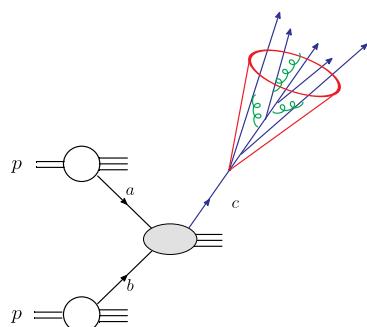
- Another benchmark process at the LHC
- Related to new physics searches, PDF fitting ...



Phenomenology

Single inclusive jet production

- Another benchmark process at the LHC
- Related to new physics searches, PDF fitting ...



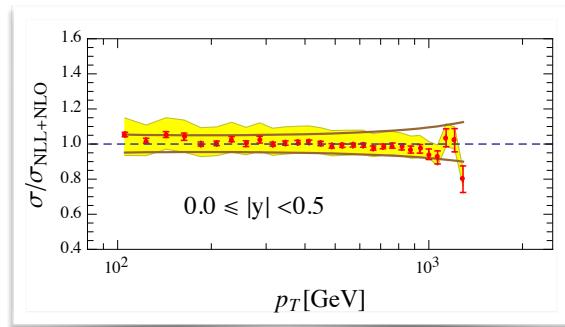
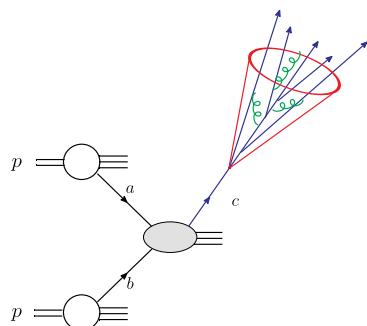
Fixed NLO or NNLO describes the data ?

Systematic discrepancies

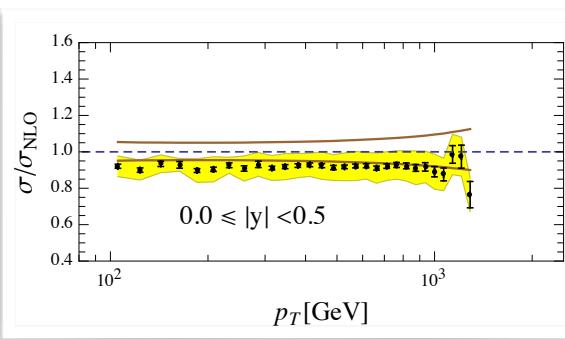
Phenomenology

Single inclusive jet production

- Another benchmark process at the LHC
- Related to new physics searches, PDF fitting ...



XL, Moch, Ringer



Fixed NLO or NNLO describes the data ?

Systematic discrepancies

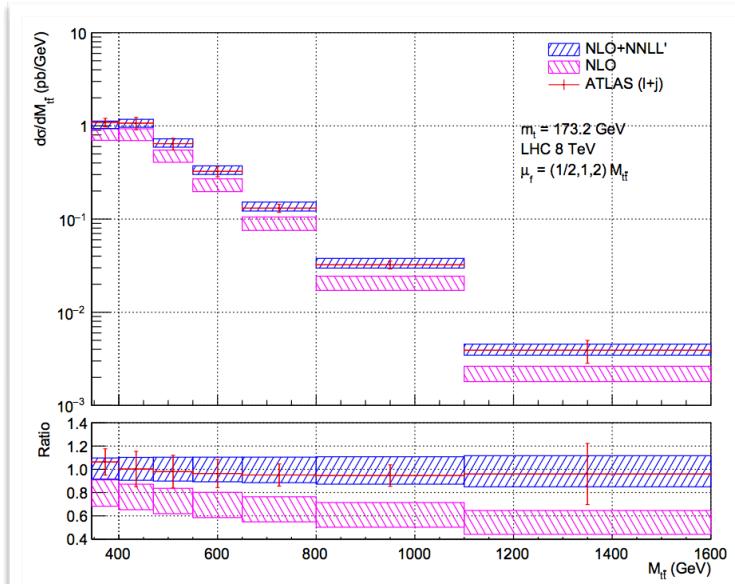
Fixed order is not everything

Resummation (small R) may help here

Phenomenology

t-tbar production

- Another example that resummation could help
- NLO + NNLL' (threshold)

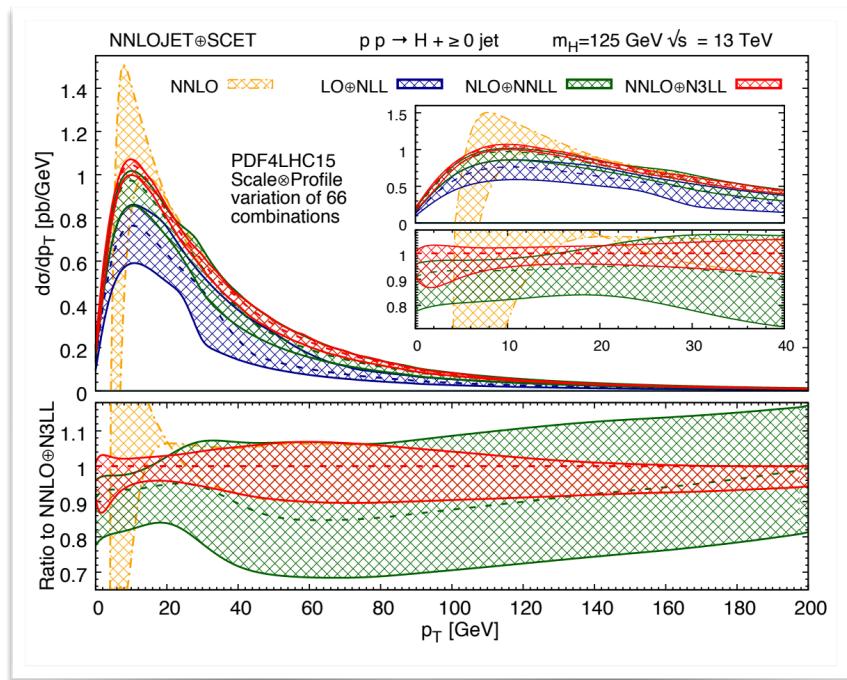


Pecjak, Scott, Wang, Yang

Phenomenology

Higgs production

- First result for NNLO + NNNLL at the LHC
- For small p_T , resummation is crucial



Conclusions

- Precision is important (test QCD + new physics searches)
- We need to deal with both loop corrections and real emissions
- Now the frontier at the LHC is NNLO + beyond