# From inflation to cosmological EWPT with a complex scalar singlet

#### Speaker: Wei Cheng

Institution: Department of physics, Chongqing University Email: Chengw@cqu.edu.cn

Based on latest work with Ligong Bian

Phys.Rev. D98 (2018), 023524

2018.10 Zhengzhou

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Model

Inflations Results Phase transition Conclusion Dark matter



### Problem and Solving

Horizon and flatness

#### Baryon asymmetry

Both Big bang Nucleosynthesis and measurements of CMB gives:

 $\frac{CMB}{Inflation} \eta_B = \frac{2}{3}$ 







01

Rotation curve of a disc galaxy Gravitational lens
Dark matter

Introduction

The Standard Model should be extension for the three cases!!!

Tommi Tenkanen, Kimmo Tuominen, Ville Vaskonen J. Cosmol.Astropart. Phys. 09 (2016) 037.

> Real scalar singlet +Fermion

Complex scalar singlet Model

The direct detection bounds from XENON1T yield null exclusions!!!

$$\begin{aligned} & \text{Scalar potential} \\ V_0(H,\mathcal{S}) &= \begin{bmatrix} \text{Standard model} \\ -\mu_h^2 |H|^2 + \lambda_h |H|^4 \\ \end{bmatrix} \\ & \text{Scalar partical} \\ \mu_s^2 |S|^2 - \begin{bmatrix} 1 \\ 2\mu_b^2 S^2 + h.c. \end{bmatrix} + \lambda_s |S|^4 \\ +\lambda_s |S|^4 \\ +\lambda_h s |H|^2 |S|^2 \\ \end{bmatrix} \\ & \text{Prescription} \\ & \text{Prescription} \\ H^T &= (0, h)/\sqrt{2} \quad S &= (s + IA)/\sqrt{2} \\ & V_0(h, s, A) &= \frac{\lambda_h h^4}{4} + \frac{1}{4} \lambda_{hs} h^2 A^2 - \frac{\mu_h^2 h^2}{2} + \frac{1}{4} \lambda_{hs} h^2 s^2 + \frac{\lambda_s A^4}{4} - \frac{\mu_s^2 A^2}{2} + \frac{\mu_b^2 A^2}{2} + \frac{\lambda_s s^4}{4} + \frac{1}{2} \lambda_s s^2 A^2 - \frac{\mu_b^2 s^2}{2} - \frac{\mu_b^2 s^2}{2} \\ & \text{DM: } m_A &= \sqrt{2} \mu_b \end{aligned} \end{aligned}$$

$$\frac{dV_0(h, s, A)}{dh}|_{h=v} = 0,$$
  
$$\frac{dV_0(h, s, A)}{ds}|_{s=v_s} = 0.$$
  
$$\mu_h^2 = \lambda_h v^2 + \lambda_{hs} v_s^2/2,$$
  
$$\mu_s^2 = -\mu_b^2 + \lambda_{hs} v^2/2 + \lambda_s v_s^2.$$

## Free parameters



$$\mathcal{M}^{2} = \begin{pmatrix} 2v^{2}\lambda_{h} & vv_{s}\lambda_{hs} \\ vv_{s}\lambda_{hs} & 2v_{s}^{2}\lambda_{s} \end{pmatrix}$$

$$R = ((\cos\theta, \sin\theta), (-\sin\theta, \cos\theta))$$

$$\tan 2\theta = -\lambda_{hs}vv_{s}/(\lambda_{h}v^{2} - \lambda_{s}v_{s}^{2})$$

$$\begin{pmatrix} h_{1} \\ h_{2} \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

$$m_{h_{1},h_{2}}^{2} = \lambda_{h}v^{2} + \lambda_{s}v_{s}^{2} \mp \frac{\lambda_{s}v_{s}^{2} - \lambda_{h}v^{2}}{\cos 2\theta}$$

$$k_{h} = \frac{\cos(2\theta)(m_{h_{1}}^{2} - m_{h_{1}}^{2}) + m_{h_{1}}^{2} + m_{h_{2}}^{2}}{4v^{2}}$$

$$Parameters: v_{s}, m_{h_{2}}^{2}, \mu_{b}$$

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#### Inflations



Slow-roll inflation

$$rac{1}{2}\dot{\phi}^2 \ll V(\phi) \qquad |\ddot{\phi}| \ll \mathcal{H}|\dot{\phi}|$$

Inflations Phase transition Dark matter The friedman Equation:  $3M_P^2\mathcal{H}^2 = \frac{1}{2}\dot{\phi}^2 + V(\phi),$  $\mathcal{H} \equiv \dot{a}/a$ The motion equation of field:  $\ddot{\phi} + 3\mathcal{H}\dot{\phi} + \frac{\partial V(\phi)}{\partial \phi} = 0.$ Slow-roll parameters:

$$\epsilon(\chi) = \frac{M_{\rm p}^2}{2} \left( \frac{dU/d\chi}{U(\chi)} \right)^2 \ll 1 \quad \eta(\chi) = M_{\rm p}^2 \left( \frac{d^2 U/d\chi^2}{U(\chi)} \right) \ll 1$$

*e*-folding number: 
$$N_{\rm e} = \int_{\chi_{\rm end}}^{\chi_{\rm in}} d\chi \frac{1}{M_{\rm p}\sqrt{2\epsilon}} = 60$$

#### Inflations

0.70

0.65

0.50

03

0.20







#### The finite temperature effective potential:



#### **Phase transition**



03

Inflations

#### **Dark matter**

The tree-level direct detection scattering amplitude:



03

Inflations

Phase transition Dark matter

#### **Results** (Real scalar singlet case)



Results Conclusion





The direct detection bounds from XENON1T yield null exclusions ! !

# **Conclusion**

Results Conclusion



	Inflation	SFOEWPT	DM
The complex singlet scalar with the global U(1) being broken	<u>.</u>		
SM +Real singlet scalar	<u></u>		
	2		