

# Exclusive Quarkonium Electroproduction

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#### Introduction

 $\gamma^*(q)p(p) \to V(q')p(p')$ 

Kinematic region: 
$$z = \frac{q' \cdot p}{q \cdot p} \approx 1$$
 (i.e.,  $t = (p' - p)^2 \approx 0$ ).

Also be termed as "elastic" or "diffractive" process.

 $q^2 = 0$ : photoproduction

 $q^2 \neq 0$  : electroproduction





#### Introduction

Regge theory:

 $J/\psi$  production: exchange Pomeron(c = +1, I = 0);

 $\eta_c$  production: exchange Odderon( c = -1, I = 0).





#### Factorization

1. Partonic process. Hard scale provided by  $M_V$  or Q. Perturbative QCD.

2.  $q\bar{q}$  pair to quarkonium state. NRQCD long-distanse matrix element.

3. Parton distribution.  $k_t$  factorization or collinear factorization.





### $k_t$ fatorization

- Based BFKL method and resummed all leading-log terms.
- Introduce the "unintegrated" pdf:  $f(x, k_T)$

$$xg(x,\mu^2) = \int^{\mu^2} \frac{d^2k_T}{\pi k_T^2} f(x,k_T)$$

- The amplitude<sup>[1]</sup>: Im $\mathcal{M} \propto xg(x, \mu^2)$ where  $\mu^2 \sim (Q^2 + M_{J/\psi}^2)/4$  $x = (Q^2 + M_{J/\psi}^2)/(W^2 + Q^2)$
- Strong dependence on scale.





### $k_t$ fatorization

• Improvement<sup>[1]</sup>:

$$\begin{split} \left[ \frac{\alpha_s(\bar{Q}^2)}{\bar{Q}^4} x g(x, \bar{Q}^2) \right] &\longrightarrow \int_{Q_0^2}^{(W^2 - M_{J/\psi}^2)/4} \frac{\mathrm{d}k_T^2 \, \alpha_s(\mu^2)}{\bar{Q}^2(\bar{Q}^2 + k_T^2)} \frac{\partial \left[ x g(x, k_T^2) \sqrt{T(k_T^2, \mu^2)} \right]}{\partial k_T^2} + \\ & \ln \left( \frac{\bar{Q}^2 + Q_0^2}{\bar{Q}^2} \right) \frac{\alpha_s(\mu_{\mathrm{IR}}^2)}{\bar{Q}^2 Q_0^2} x g(x, Q_0^2) \sqrt{T(Q_0^2, \mu_{IR}^2)} \,. \end{split}$$

- Shortage:
  - 1. Effect beyond leading-log.
  - 2. Difficult to NLO.
  - 3.  $\eta_c$ ?



#### Collinear factorization

Collinear factorization is commonly used in DVCS and exclusive VM production, and lead to the concept of generalized parton distributions (GPDs).





#### GPDs

twist-2 GPDs:  $H^{g}(x,\xi,t), H^{q}(x,\xi,t)...$  with the skewness  $\xi$ 

$$\int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle P' | \overline{\psi(-\lambda n/2)} \gamma^{\mu} \psi(\lambda n/2) | P \rangle = H(x,\xi,\Delta^2) \overline{U(P')} \gamma^{\mu} U(P) + E(x,\xi\Delta^2) \overline{U(P')} \frac{i\sigma^{\mu\nu} \Delta_{\nu}}{2M} U(P) + \cdots,$$

$$\int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle P' | \overline{\psi(-\lambda n/2)} \gamma^{\mu} \gamma_5 \psi(\lambda n/2) | P \rangle = \widetilde{H}(x,\xi,\Delta^2) \overline{U(P')} \gamma^{\mu} \gamma_5 U(P) + \widetilde{E}(x,\xi,\Delta^2) \overline{U(P')} \frac{\gamma_5 \Delta^{\mu}}{2M} U(P) + \cdots,$$



 $|x| > \xi$  : DGLAP region  $|x| < \xi$  : ERBL region



#### GPDs

• Forward limit:

$$H^g(x,\xi,t) \xrightarrow{\xi=0} xg(x) \ H^q(x,\xi,t) \xrightarrow{\xi=0} q(x)$$

- extrapolate from PDFs:
  - 1) foward model for  $x >> \xi$ ,  $H(x,\xi) \approx H(x,0)$ ; for other region ...
  - 2) Shuvaev transform (poor for large  $\xi$ )
  - 3) Double distribution model (singular)

Fraught with uncertainties!

For review on GPDs: M. Diehl, Phys, Rept, 388 41 (2003), A. V. Belitskly, A. V. Radyushkin, Phys, Rept 418, 1 (2005).



### Photoproduction

#### $\Upsilon$ photoproduction at NLO<sup>[1]</sup>:



[1] D. Yu. Ivanov, A. Schafer, L. Szymanowski, G. Krasnikov, Eur. Phys. J. C 34, 297 (2004); 75, 75 (E) (2015).



# Photoproduction

#### $J/\psi$ photoproduction at NLO:



small value of  $m_c$  lead to : 1) large  $\alpha_s(m_c)$ ; 2) small  $\xi$ .



# Photoproduction

Improve the convergence of perturbative series by resumming  $\ln \frac{1}{\xi}$  terms<sup>[1]</sup>.



[1] S. P. Jones, A. D. Martin, M. G. Ryskin, T. Teubner, J. Phys. G: Nucl. Part. Phys. 43 035002 (2016).



hard scale: 
$$Q^2, M_V^2$$
  
skewness:  $\xi = \frac{Q^2 + M_V^2}{2s + Q^2 - M_V^2}$ 

expect a better convergence than photoproduction

$$\mathcal{M} \sim \int_{-1}^{1} dx \left[ C^{(0)}(x,\mu_F) \text{GPD}(x,\mu_F) + C^{(1)}(x,\mu_R,\mu_F) \text{GPD}(x,\mu_F) \right]$$

partonic process:





- Ultraviolet singularities. Removed by  $\overline{\text{MS}}$  renomalization.
- Infrared singularities. Parts cancelled each other, remaining absorbed by GPDs.
- Singularities at  $x = \pm \xi$







$$A^{\lambda\lambda\prime} = c_0^{\lambda\lambda\prime} + \sum_{i=1}^{11} c_i^{\lambda\lambda\prime} f_i \pm \{y_1 \to 1 - y_1 - y_2\},\$$

$$d_1 = (1+y_2)^{-1}, d_2 = (1+2y_2)^{-1}, d_3 = (1-2y_1)^{-1}, d_4 = (2y_1+2y_2-1)^{-1}, d_5 = (2y_1+y_2-1)^{-1}, d_6 = (6y_1+y_2-5)^{-1}, d_7 = (6y_1+5y_2-1)^{-1}.$$

$$\begin{aligned} f_1 &= \sqrt{\frac{-y_1}{1-y_1}} \frac{\ln(\sqrt{1-y_1}+\sqrt{-y_1})}{-y_1}, f_2 = f_1^2, f_3 = \sqrt{\frac{1+y_2}{y_2}} \frac{\ln(\sqrt{1+y_2}+\sqrt{y_2})}{1+y_2}, f_4 = f_3^2, \\ f_5 &= d_1 d_3 d_5 \ln(\frac{1-y_1}{1+y_2}), f_6 = d_1 d_3 d_4 d_5 \ln(2+2y_2), \\ f_7 &= d_1 \ln(1+y_2)^2, f_8 = d_1^2 (\ln(\frac{4m^2}{\mu_F^2}) \ln(\frac{1-y_1}{1+y_2}) + \ln(1-y_1)^2), \\ f_9 &= d_1^2 \text{Li}_2 (-1+2y_1), f_{10} = d_1^2 \text{Li}_2 (-1-2y_2), \\ f_{11} &= d_1^2 d_5 C_0 (-4y_2, -1+2y_1, 1-2y_1 - 2y_2, 1, 1, 0), \end{aligned}$$
(13)

<sup>1</sup> For 
$$f_{11}$$
,  $C_0(-4y_2, -1+2y_1, 1-2y_1-2y_2, 1, 1, 0) = \frac{1}{2(2y_1+y_2-1)} \{ \text{Li}_2[\frac{2(1-y_1)}{1+y_2}] - \text{Li}_2[\frac{(1-2y_1-2y_2)(1+y_2)}{1+y_2-4(1-y_1)(y_1+y_2)}] + \text{Li}_2[\frac{2(1-y_1)(1-2y_1-2y_2)}{1+y_2-4(1-y_1)(y_1+y_2)}] - \text{Li}_2[\frac{2(1-y_1)y_2}{(1+y_2)y_2-\sqrt{(1+y_2)y_2}(2y_1+y_2-1)}] - \text{Li}_2[\frac{2(1-y_1)y_2}{(1+y_2)y_2+\sqrt{(1+y_2)y_2}(2y_1+y_2-1)}] + \{y_1 \rightarrow 1 - y_1 - y_2\}$ 



For  $A_{g}^{++}$  and  $A_{g}^{--}$ ,  $c_{i}^{++} = c_{i}^{--} = c_{i}$ ,  $c_0 = d_1^2 d_2 d_3 d_4 (\frac{13\pi^2 s^4}{4^2 \Omega^{12}} - \frac{31\pi^2 s^4}{2^{12} \Omega^{14}} + \frac{\pi^2 s^4}{8} + \frac{2\pi^2}{9} - \frac{\pi^2 s^3}{12^{12}}$  $+ \tfrac{88a^{3}b}{5} + \tfrac{38}{52}\pi^{2}a^{3}b + \tfrac{183\pi^{2}a^{3}}{495} - \tfrac{82\pi^{2}a^{3}}{95} - \tfrac{33\pi^{3}}{5} + \tfrac{23\pi^{2}b^{2}}{5} + \tfrac{87}{54}\pi^{2}a^{2}b^{2}$  $+ \tfrac{19\pi^2a^2}{576b^2} - \tfrac{133a^2b}{9} - \tfrac{278}{108}\pi^2a^2b - \tfrac{28\pi^2a^2}{2165} + \tfrac{328\pi^2a^2}{432} + \tfrac{11a^2}{3} - \tfrac{11\pi^2a}{1728b^2}$  $+ \frac{11ab}{3} + \frac{29}{144} \pi^2 ab - \frac{29\pi^2 a}{2845} + \frac{34\pi^2 a}{72} - \frac{8a}{9} + \frac{71\pi^2}{1728a} - \frac{319\pi^2}{1728}) - d_1 \frac{g_0}{24} \ln \frac{g_0}{a_1^8},$  $c_1 = y_1^{-1} \left( -\frac{2\pi^2}{104} + \frac{23\pi}{104} - \frac{1}{1044} - \frac{2\pi}{10} - \frac{1}{104} - \frac{37}{1044} + \frac{11}{10} \right),$  $c_2 = y_1^{-1} \left( -\frac{110}{385} + \frac{1}{7165} + \frac{1}{385} - \frac{1}{55} - \frac{1}{5} \right),$  $c_3 = d_1^2 d_{8,02}^2 \left( -\frac{2a^4}{0b} + \frac{11a^3}{12b} - \frac{2a^3}{0} + \frac{4a^2b}{0} + \frac{4a^2}{2b} - \frac{5a^2}{3} - \frac{5ab}{3a} + \frac{4a}{0} \right),$  $c_4 = d_{112}^3 (-\frac{\hbar a^2}{1225^2} - \frac{11a^2}{885} + \frac{\hbar a}{225^2} - \frac{\hbar a}{65} - \frac{a}{86} - \frac{11}{1225} - \frac{47}{122}),$  $c_5 = d_3 d_5 \left( -\frac{4a^4}{2b} + \frac{19a^5}{3b} + \frac{8a^2b}{2} - \frac{73a^4}{2b} - \frac{37a^4}{8} - \frac{67a^2b}{18} + \frac{65a^3}{18b} + \frac{301a^3}{38} \right)$  $-\frac{4a^{2}b^{2}}{b}+\frac{83a^{2}b^{2}}{16}+\frac{61a^{2}b}{36}-\frac{13a^{2}}{32b}-\frac{47a^{2}}{12}+\frac{13ab^{2}}{18}+\frac{b^{2}}{36}-\frac{185ab^{2}}{36}+\frac{b^{2}}{92a}$  $+\frac{37ab}{4a}+\frac{47a}{77}-\frac{13b^3}{3a}+\frac{53b^2}{3a}-\frac{41b}{77}$ For  $A_q^{++}$  and  $A_q^{--}$ ,  $c_i^{++} = c_i^{--} = c_i$ ,  $c_8 = a_2^2 d_3 d_4 d_8 \left(\frac{80a^7b}{6b} + \frac{14a^7}{6b} - \frac{76a^7}{6} + \frac{176a^6b^2}{5} - \frac{292a^6b}{5} - \frac{86a^6}{6b} + \frac{146a^6}{5} + 48a^8b^3\right)$  $-\frac{828a^{5}b^{2}}{3}+228a^{5}b+\frac{77a^{5}}{28}-87a^{5}-\frac{16a^{4}b^{4}}{6}-\frac{748a^{4}b^{3}}{6}+\frac{2218a^{4}b^{2}}{6}-\frac{1083a^{4}b}{64}-\frac{49a^{4}}{64}$  $+ \frac{1281a^4}{18} + \frac{154a^3b^3}{3} - \frac{410a^3b^2}{3} + \frac{881a^3b}{3} + \frac{115a^3}{725} - \frac{85a^3}{5} + \frac{157a^2b^2}{8} - \frac{367a^2b}{12}$  $-\frac{7a^2}{345}+\frac{29a^2}{4}+\frac{27ab}{79}-\frac{a}{7}),$  $c_7 = \frac{4}{55}, c_8 = -\frac{4^2}{55} - \frac{4}{5},$ For  $A_{\mu}^{00}$ ,  $c_{1}^{00} = -\sqrt{y_{2}}c_{1}$ ,  $c_{0} = -\frac{\delta a^{2}}{18812} - \frac{b^{2}}{312} - \frac{\delta a^{2}}{312} + \frac{13b}{3442} + \frac{\delta a}{3412} + \frac{b^{2}}{122} - \frac{47a}{122} - \frac{13b}{122} + \frac{7a}{2} + \frac{11}{34} + \frac{b}{34} + \frac{36}{344} - \frac{49}{124}$  $c_{10} = \frac{5a^2}{3ab} + \frac{5a}{4bb} - \frac{a}{5b} - \frac{1}{bb} + \frac{13}{14a},$  $c_{11} = d_5 d_6^{-1} d_7^{-1} \left( -\frac{a^3}{(24)^2} + \frac{a^2}{240^2} - \frac{5a^2}{245} + \frac{a^2}{6} - \frac{a5}{14} + \frac{a}{365} - \frac{a}{14} + \frac{1}{124} \right),$ (15)For  $\tilde{A}_{e}$ ,  $\tilde{c}_{i}^{00} = 0$ ,  $\tilde{c}_{i}^{++} = -\tilde{c}_{i}^{--} = c_{i}$ ,

 $c_8 = d_8^{-2} \frac{2}{56}, c_8 = d_2 d_8^{-1} y_{2\frac{5}{5}}^{24}.$  (20)

For  $\tilde{A}_{a}, \tilde{c}_{a}^{00} = 0, \tilde{c}_{a}^{++} = -\tilde{c}_{a}^{--} = c_{a}$ .

For  $A_x^{00}$ ,  $c_i^{00} = -\sqrt{y_0}c_i$ ,  $c_0 = d_1^2 d_3 d_4 (\frac{2\pi^2 a^2}{4\pi^2 b^2} + \frac{28\pi^2 b}{b} + \frac{a}{2^2} \pi^2 a^2 b - \frac{28\pi^2 a^2}{54b} + \frac{2\pi^2 a^2}{2^2} - \frac{18\pi^2}{9}$  $-\frac{2\pi^{2}a}{3\pi^{2}a} - \frac{44ab}{3\pi^{2}} + \frac{12}{3\pi}\pi^{2}ab + \frac{31\pi^{2}a}{3\pi^{2}} - \frac{43\pi^{2}a}{3\pi^{2}} + \frac{8a}{3\pi} - \frac{37\pi^{2}}{3\pi^{2}} + \frac{28\pi^{2}}{3\pi^{2}}) - d_{1}\frac{d_{1}}{2\pi}\ln\frac{d_{1}}{2\pi},$  $c_1 = -\frac{7}{6\pi} - \frac{2}{6\pi}, c_2 = y_1^{-1}(-\frac{2\pi}{6\pi} + \frac{7}{76\pi} - \frac{7}{6\pi}),$  $c_3 = d_1^2 d_5^2 \left( \frac{2a^4}{ab} + \frac{8a^3}{ab} - \frac{a^3}{3} + \frac{17a^2b}{a} - \frac{7a^2}{ab} - \frac{14a^2}{a} - \frac{23ab}{a} + \frac{23a}{a} \right),$  $c_4 = d_1^3 y_2 \left( -\frac{2a^2}{6\lambda} - \frac{7a}{6\lambda} - \frac{2a}{6} + \frac{1}{6\lambda} - \frac{1}{8} \right),$  $c_5 = d_3 d_5 \left( \frac{32a^5}{2b} - \frac{28a^4}{2b} - \frac{49a^4}{5} - 2a^3b + \frac{a^3}{3b} + \frac{145a^3}{14} + \frac{19a^2b^2}{5} + \frac{a^2b}{3} \right)$  $+\frac{a^2}{55}-\frac{41a^2}{9}-\frac{5ab^2}{5}+\frac{7b^2}{18a}+\frac{13ab}{9}+\frac{17a}{18}-\frac{7b^2}{9}-\frac{11b}{18}),$  $c_8 = d_2 d_3 d_4 d_8 \left(-\frac{16}{9} a^8 b^2 + 12 a^8 b + \frac{28 a^8}{24} - \frac{98 a^8}{9} - \frac{80 a^4 b^3}{3}\right)$  $+ \frac{488\pi^{4}b^{2}}{6b} - \frac{428\pi^{4}b}{4b} - \frac{14\pi^{4}}{4b} + \frac{194\pi^{4}}{4b} + \frac{384\pi^{3}b^{3}}{4b} - \frac{914\pi^{3}b^{2}}{4b} + \frac{814\pi^{3}b}{6b} + \frac{7\pi^{3}}{4b} - \frac{281\pi^{3}}{18} + \frac{368\pi^{2}b^{2}}{18}$  $= \frac{847\pi^{2}b}{18} = \frac{7\pi^{2}}{188} + \frac{43\pi^{2}}{9} + \frac{32\pi b}{9} = \frac{\pi}{9}),$  $c_7 = \frac{a}{2a}, c_8 = -\frac{a^2}{1b} - \frac{b}{2}, c_9 = -\frac{a^2}{1b} - \frac{7b}{1ba^2} + \frac{a}{4b} + \frac{bb}{1a} + \frac{b}{1a} - \frac{b}{1ba} - \frac{7b}{1b} - \frac{2}{4b} + \frac{b}{4},$  $c_{10} = \frac{a^2}{36} - \frac{a}{36} + \frac{a}{2} + \frac{1}{126} - \frac{2}{6}, c_{11} = d_3 d_6^{-1} d_7^{-1} (\frac{a^2}{16} - \frac{a}{6} - \frac{a}{166} + \frac{2a}{16} - \frac{1}{16}).$ (16)  $c_0 = d_1^2 \frac{\pi^2}{2\pi}, c_3 = d_1^2 d_3 y_2 \frac{\pi}{2}, c_4 = d_1^3 y_2 \frac{\pi}{2\pi},$  $c_8 = -\frac{2a^2}{bb} + \frac{2a}{bb} - \frac{8a}{b} + \frac{2b}{b} + \frac{2}{b}, c_8 = d_2(\frac{8a^3}{b} - 8a^2b + \frac{8a^2}{b} + \frac{8a^2}{b} - \frac{32a}{b} + \frac{4}{b}),$  $c_7 = -\frac{4}{32}, c_8 = c_9 = \frac{4\pi}{32} - \frac{4}{3}, c_{10} = \frac{1}{32} - \frac{4\pi}{32}, c_{11} = -\frac{2\pi^2}{32} + \frac{8\pi^2}{3} - \frac{8\pi^5}{3} - \frac{2\pi}{3},$ (17) $c_3 = d_1^2 d_5 (1 - y_2) \frac{8}{9}, c_6 = \frac{8a^2}{05} - \frac{4a}{05} - \frac{4}{9}, c_6 = \frac{16ab}{9} - \frac{8a}{3} + \frac{8}{9},$  $c_7 = -\frac{4}{65}, c_8 = c_9 = \frac{4\pi}{65} - \frac{4}{9}, c_{10} = -\frac{4\pi}{65} - \frac{4}{9}, c_{11} = \frac{8\pi^2}{9} - \frac{8\pi^3}{3} + \frac{16\pi}{9}$ (18)  $c_1 = -\frac{1}{36a_1}, c_2 = \frac{1}{36a_2}, c_4 = -d_1y_2\frac{1}{36a_2}$  $c_5 = d_5 d_5^{-1} (\frac{19a^2}{18} - \frac{ab}{8} - \frac{7a}{12} + \frac{bb}{38}), c_6 = d_5 d_4 d_5^{-2} (-\frac{2a^2b}{3} + \frac{ba^2}{5} + \frac{14ab}{5} - \frac{28a}{18} + \frac{13}{38}),$  $c_9 = d_1(\frac{1}{24} - \frac{5a}{725}), c_{11} = d_1 d_5^{-2} d_5^{-1} d_7^{-1} \frac{1}{7255}$ (19)



#### At high energy limit:

$$\mathcal{M} \sim \alpha_s F^g(\xi, \xi, \mu_F) + \frac{\alpha_s^2}{\pi} \ln \frac{\bar{m}^2}{\mu_F^2} \ln \frac{1}{\xi} \left[ 3F^g(\xi, \xi, \mu_F) + \frac{4}{3}F^{q(+)}(\xi, \xi, \mu_F) \right],$$

where

$$\bar{m} = \sqrt{(m^2 + \frac{Q^2}{4})}$$



#### Parameter list:

- $\Lambda^3_{\text{QCD}} = 332 \text{MeV}, \ \Lambda^4_{\text{QCD}} = 292 \text{MeV}, \ \Lambda^5_{\text{QCD}} = 210 \text{MeV};$
- $|R_{J/\psi}(0)|^2 = 0.903 \text{GeV}^3$ ,  $|R_{\Upsilon}(0)|^2 = 7.76 \text{GeV}^3$ ;
- $1.4 \le m_c \le 1.6, 4.8 \le m_b \le 5.0;$
- $max\{1, \frac{1}{2}\sqrt{m^2 + Q^2/4}\} \le \mu_F \le 2\sqrt{m^2 + Q^2/4};$
- $\mu_R = \mu_F$ .

The choice of  $\mu_R$  eliminates the contribution from  $\beta_0$  term ( $\beta_0 \ln \frac{\mu_F^2}{\mu_R^2}$ ), which related to the BLM(PMC) method.



HERA上 $J/\psi$ 产生:





HERA上 $J/\psi$ 产生:





HERA上Y产生:





# Conclusion

- The quarkonium exclusive production process can potentially serve to constrain the parton distributions in low  $x_b$  region.
- For  $J/\psi$  electroproduction, the NLO result agrees with data nicely, with theoretical uncertainty largely suppressed at large  $Q^2$  region.
- Precise phenomenological study require more reliable approach for the evaluation of GPDs.



# 谢谢

