



Weak Decays of Doubly Heavy Baryons

$$\mathcal{B}_{cc} \rightarrow \mathcal{B}_c V$$



姜丽娟

arXiv:1810.00541

Inner Mongolia University

In collaboration with : 贺蓓 李润辉

October 27, 2018 HFCPV

OUTLINE

1 Introduction

2 Framework and Calculations

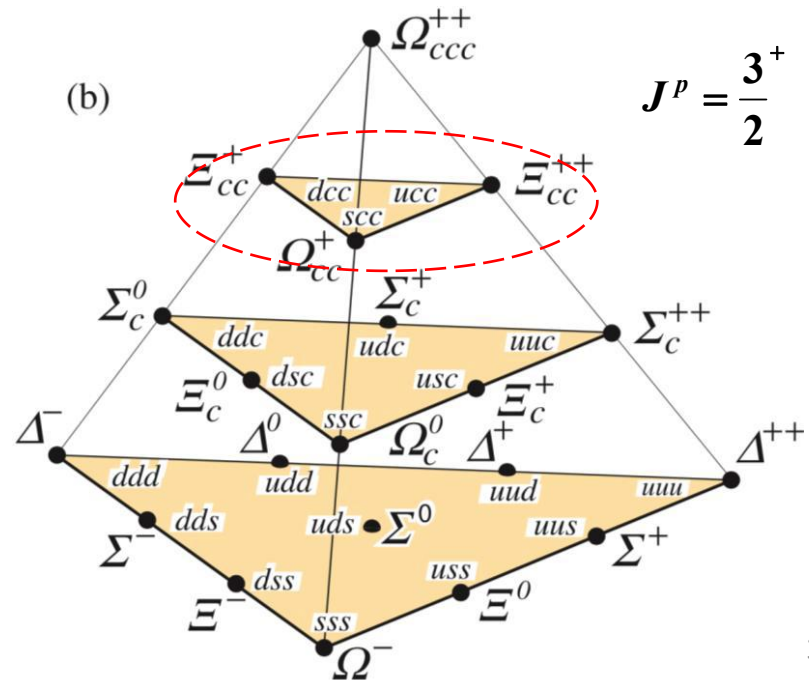
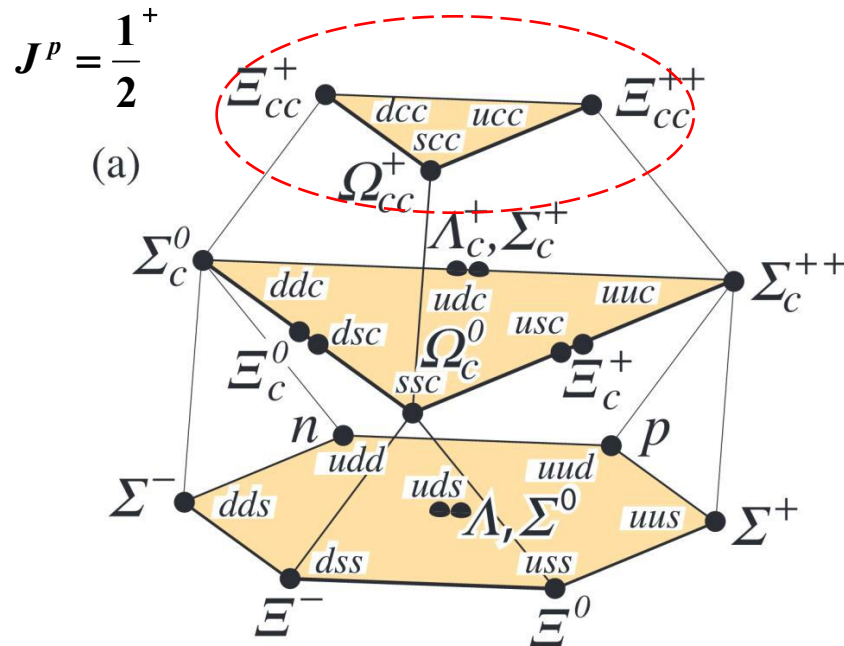
3 Results and Discussions

4 Summary

1 Introduction

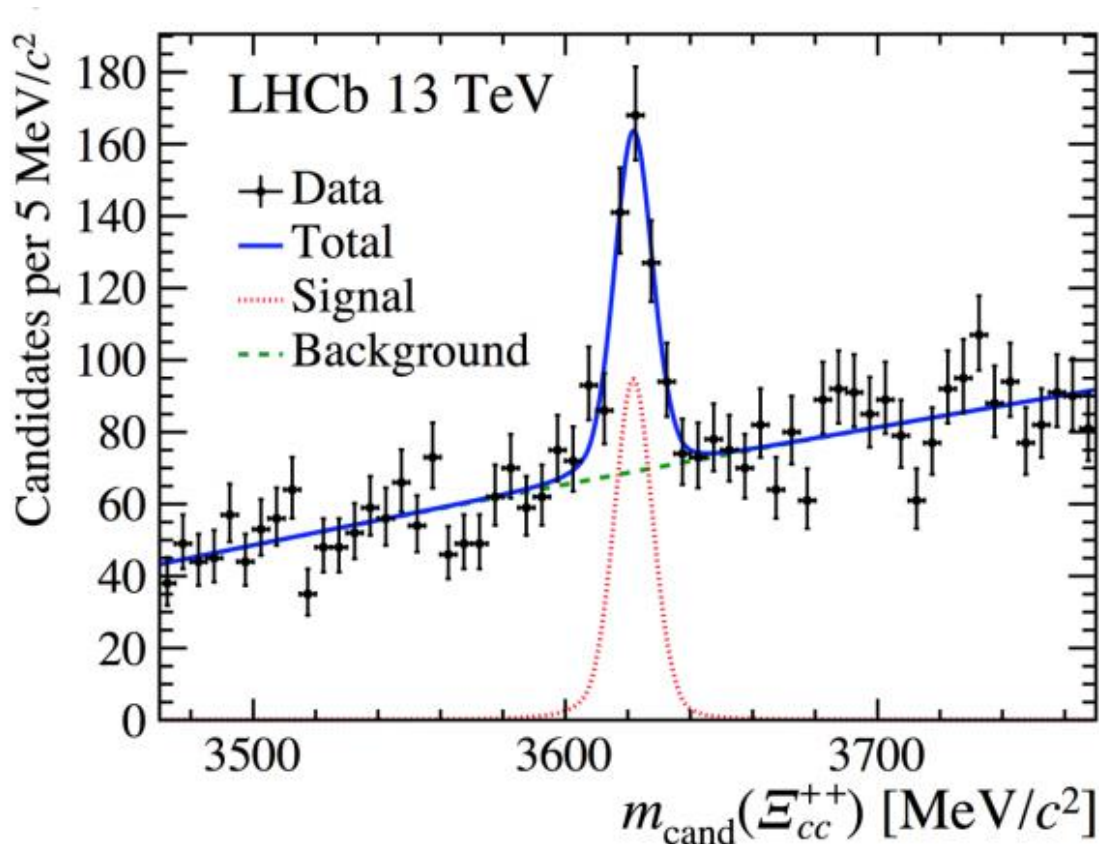
The doubly charm baryons

- ◆ Two SU(4) baryon 20-plets with $J^P = \frac{1}{2}^+$ and $J^P = \frac{3}{2}^+$, each contains a SU(3) triplet with two charm quarks: Ξ_{cc}^{++} (ccu) Ξ_{cc}^+ (ccd) Ω_{cc}^+ (ccs).
- ◆ $J^P = \frac{3}{2}^+$ expected to decay to $J^P = \frac{1}{2}^+$ states via strong/electromagnetic interaction.
- ◆ $J^P = \frac{1}{2}^+$ states decay weakly with a quark transformed to lighter quarks.



1 Introduction

- SELEX observed Ξ_{cc}^+ in $\Xi_{cc}^+ \rightarrow \Lambda_c^+ K^- \pi^+$ and $\Xi_{cc}^+ \rightarrow p D^+ K^-$ decay
- LHCb observed the Ξ_{cc}^{++} state in the $\Lambda_c^+ K^- \pi^+ \pi^+$ decay



PRL119,112011

1 Introduction

- Theoretical work that proposed the discovery channel

Chinese Physics C Vol. 42, No. 5 (2018) 051001

Discovery potentials of doubly charmed baryons *

Fu-Sheng Yu(于福升)^{1,2;1)} Hua-Yu Jiang(蒋华玉)^{1,2} Run-Hui Li(李润辉)³ Cai-Dian Lü(吕才典)^{4,5;2)}
Wei Wang(王伟)^{6;3)} Zhen-Xing Zhao(赵振兴)⁶

¹ School of Nuclear Science and Technology, Lanzhou University, Lanzhou 730000, China

² Research Center for Hadron and CSR Physics, Lanzhou University and Institute of Modern Physics of Chinese Academy of Sciences, Lanzhou 730000, China

³ School of Physical Science and Technology, Inner Mongolia University, Hohhot 010021, China

⁴ Institute of High Energy Physics, Chinese Academy of Sciences, YuQuanLu 19B, Beijing 100049, China

⁵ School of Physics, University of Chinese Academy of Sciences, YuQuanLu 19A, Beijing 100049, China

⁶ INPAC, Shanghai Key Laboratory for Particle Physics and Cosmology, School of Physics and Astronomy, Shanghai Jiao-Tong University, Shanghai 200240, China

Abstract: The existence of doubly heavy flavor baryons has not been well established experimentally so far. In this Letter we systematically investigate the weak decays of the doubly charmed baryons, Ξ_{cc}^{++} and Ξ_{cc}^+ , which should be helpful for experimental searches for these particles. The long-distance contributions are first studied in the doubly heavy baryon decays, and found to be significantly enhanced. Comparing all the processes, $\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$ and $\Xi_c^+ \pi^+$ are the most favorable decay modes for experiments to search for doubly heavy baryons.

- A phenomenological model is employed

Fu-Sheng Yu's report

- Under this framework we calculated the \mathcal{BR} and Γ of $\mathcal{B}_{cc} \rightarrow \mathcal{B}_c V$

2 Framework and Calculations

➤ Effective Hamiltonian

We focus on weak decays induced by the charge current $c \rightarrow s/d$.

The contributing low energy effective Hamiltonian is given by

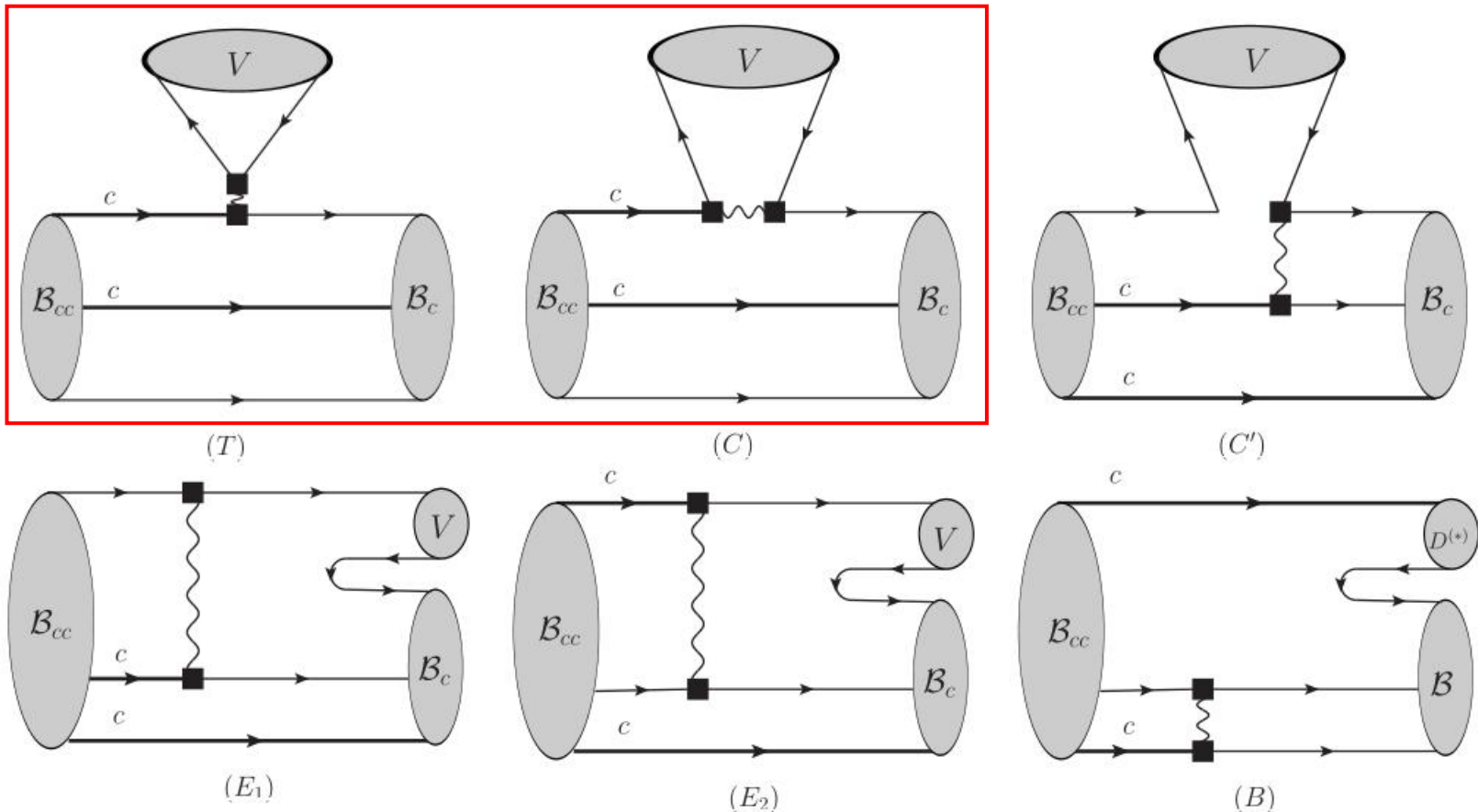
$$H_{eff} = \frac{G_F}{\sqrt{2}} \sum_{q=d,s} V_{cq}^* V_{uD} [C_1(\mu) O_1^q(\mu) + C_2(\mu) O_2^q(\mu)] + h.c.$$

$$O_1^q = (\bar{u}_\alpha D_\beta)_{V-A} (\bar{q}_\beta c_\alpha)_{V-A}, \quad O_2^q = (\bar{u}_\alpha D_\alpha)_{V-A} (\bar{q}_\beta c_\beta)_{V-A}$$

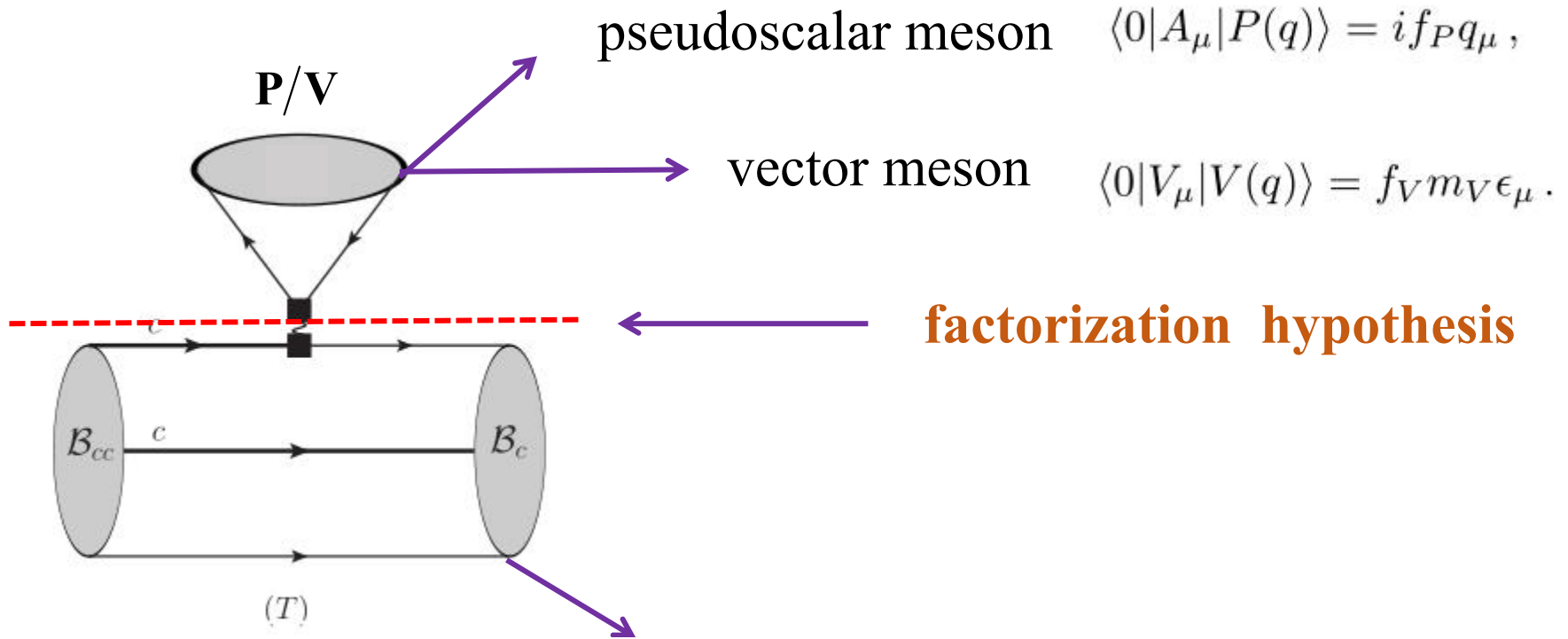
2 Framework and Calculations

➤ Topological Diagrams

Both short distance and long distance contributions are contained.



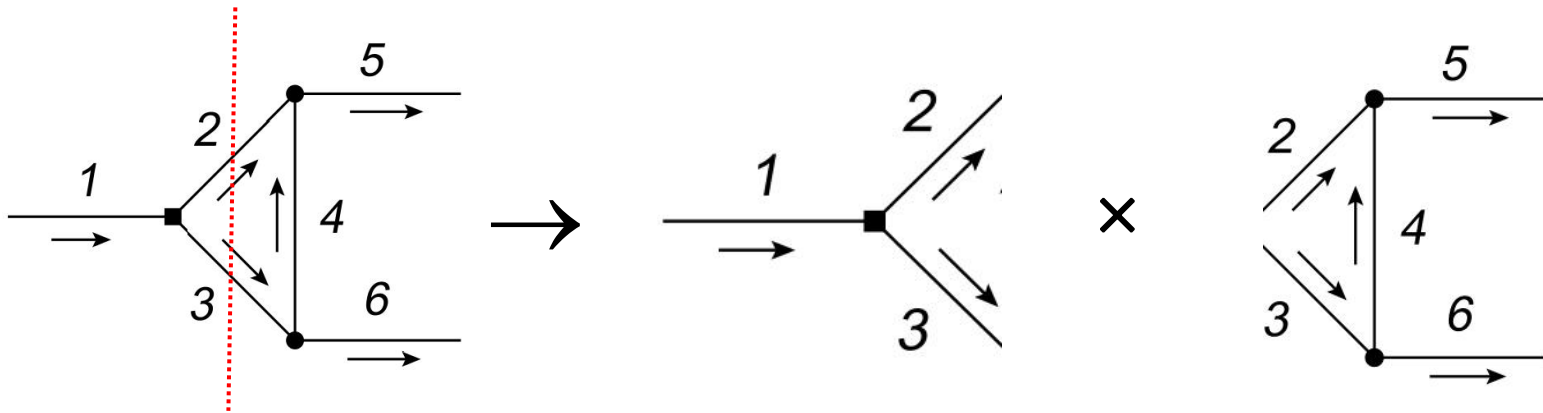
➤ Short Distance Amplitudes



$$\langle \mathcal{B}_c(p', s'_z) | (V - A)_\mu | \mathcal{B}_{cc}(p, s_z) \rangle = \bar{u}(p', s'_z) \left[\gamma_\mu f_1(q^2) + i\sigma_{\mu\nu} \frac{q^\nu}{M} f_2(q^2) + \frac{q^\mu}{M} f_3(q^2) \right] u(p, s_z) - \bar{u}(p', s'_z) \left[\gamma_\mu g_1(q^2) + i\sigma_{\mu\nu} \frac{q^\nu}{M} g_2(q^2) + \frac{q^\mu}{M} g_3(q^2) \right] \gamma_5 u(p, s_z),$$

➤ Long Distance Contributions

The long distance contributions are modeled as final-state interactions (FSIs) and calculated under the one-particle-exchange model



H. Y. Cheng, C. K. Chua and A. Soni, Phys. Rev. D 71, 014030 (2005)

➤ Long Distance Contributions

We adopt Optical theorem to calculate the imaginary part of triangle diagram , treat the intermediate particles 2 and 3 as on-shell.

The absorptive part :

$$\mathcal{A}bs M(p_1 \rightarrow p_5 p_6) = \frac{1}{2} \int \frac{d^3 p_2 d^3 p_3}{(2\pi)^6 4E_2 E_3} (2\pi)^4 \delta^4(p_5 + p_6 - p_2 - p_3) M(p_1 \rightarrow p_2 p_3) T^*(p_5 p_6 \rightarrow p_2 p_3).$$

The dispersive part can be calculated via the dispersion relation:

$$\mathcal{D}is A(m_1^2) = \frac{1}{\pi} \int_s^\infty \frac{\mathcal{A}bs A(s')}{s' - m_1^2} ds',$$

We assume the absorptive part is dominating and neglect the dispersive part.

H. Y. Cheng, C. K. Chua and A. Soni, Phys. Rev. D 71, 014030 (2005)

➤ Long Distance Contributions

To proceed, the Lagrangian for the strong couplings is necessary:

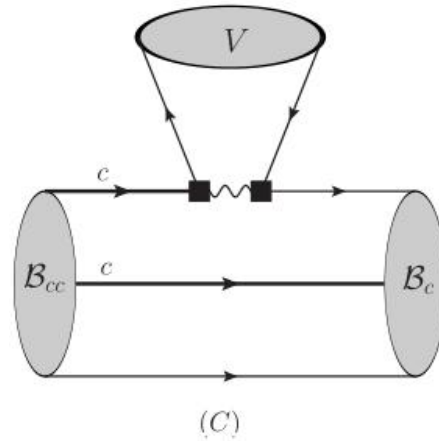
$$\mathcal{L}_{\pi hh} = g_{\pi \mathcal{B}_6 \mathcal{B}_6} \text{Tr}[\bar{\mathcal{B}}_6 i \gamma_5 \Pi \mathcal{B}_6] + g_{\pi \mathcal{B}_3 \mathcal{B}_3} \text{Tr}[\bar{\mathcal{B}}_3 i \gamma_5 \Pi \mathcal{B}_3] + \{g_{\pi \mathcal{B}_6 \mathcal{B}_3} \text{Tr}[\bar{\mathcal{B}}_6 i \gamma_5 \Pi \mathcal{B}_3] + h.c.\},$$

$$\begin{aligned} \mathcal{L}_{\rho hh} = & f_{1\rho \mathcal{B}_6 \mathcal{B}_6} \text{Tr}[\bar{\mathcal{B}}_6 \gamma_\mu V^\mu \mathcal{B}_6] + \frac{f_{2\rho \mathcal{B}_6 \mathcal{B}_6}}{2m_6} \text{Tr}[\bar{\mathcal{B}}_6 \sigma_{\mu\nu} \partial^\mu V^\nu \mathcal{B}_6] \\ & + f_{1\rho \mathcal{B}_3 \mathcal{B}_3} \text{Tr}[\bar{\mathcal{B}}_3 \gamma_\mu V^\mu \mathcal{B}_3] + \frac{f_{2\rho \mathcal{B}_3 \mathcal{B}_3}}{2m_3} \text{Tr}[\bar{\mathcal{B}}_3 \sigma_{\mu\nu} \partial^\mu V^\nu \mathcal{B}_3] \\ & + \{f_{1\rho \mathcal{B}_6 \mathcal{B}_3} \text{Tr}[\bar{\mathcal{B}}_6 \gamma_\mu V^\mu \mathcal{B}_3] + \frac{f_{2\rho \mathcal{B}_6 \mathcal{B}_3}}{m_6 + m_3} \text{Tr}[\bar{\mathcal{B}}_6 \sigma_{\mu\nu} \partial^\mu V^\nu \mathcal{B}_3] + h.c.\}, \end{aligned}$$

$$\mathcal{L}_{\rho\pi\pi} = \frac{ig_{\rho\pi\pi}}{\sqrt{2}} \text{Tr}[V^\mu [\Pi, \partial_\mu \Pi]],$$

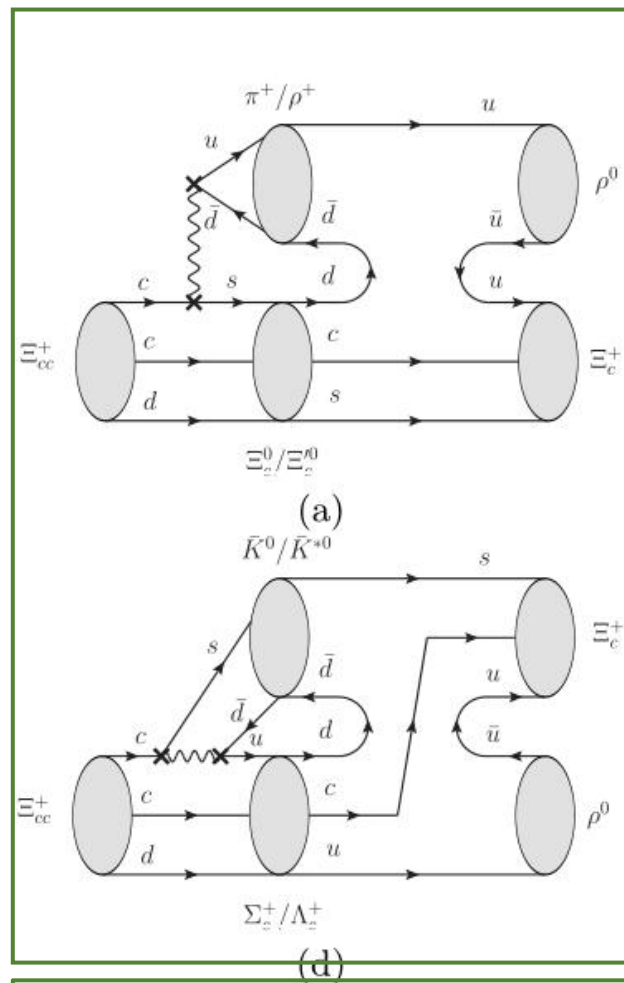
$$\mathcal{L}_{\rho\rho\rho} = \frac{ig_{\rho\rho\rho}}{\sqrt{2}} \text{Tr}[(\partial_\nu V_\mu - \partial_\mu V_\nu) V^\mu V^\nu] = \frac{ig_{\rho\rho\rho}}{\sqrt{2}} \text{Tr}[(\partial_\nu V_\mu V^\mu - V^\mu \partial_\nu V_\mu) V^\nu],$$

➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: Short Distance Contributions

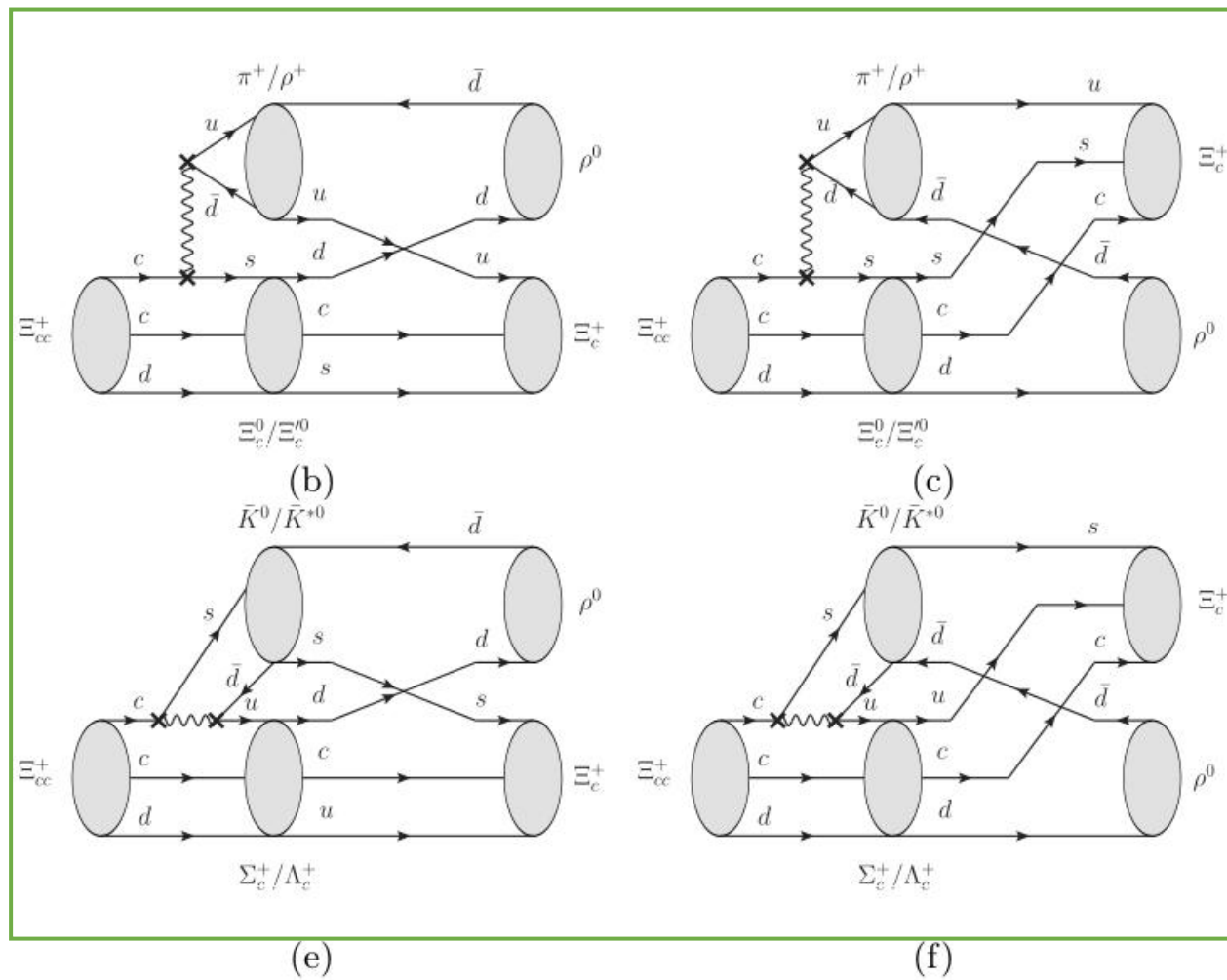


$$\begin{aligned}
 C_{\text{SD}}(\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0) &= \frac{G_F}{\sqrt{2}} V_{cs}^* V_{ud} a_2 f_\rho \epsilon_\mu^* \bar{u}(p_6, s'_z) \left[\left(f_1(m_\rho^2) - \frac{m_{\Xi_{cc}^+} + m_{\Xi_c^+}}{m_{\Xi_{cc}^+}} f_2(m_\rho^2) \right) \gamma^\mu + \frac{2}{m_{\Xi_{cc}^+}} f_2(m_\rho^2) p_6^\mu \right. \\
 &\quad \left. - \left(g_1(m_\rho^2) + \frac{m_{\Xi_{cc}^+} - m_{\Xi_c^+}}{m_{\Xi_{cc}^+}} g_2(m_\rho^2) \right) \gamma^\mu \gamma_5 - \frac{2}{m_{\Xi_{cc}^+}} g_2(m_\rho^2) p_6^\mu \gamma_5 \right] u(p_1, s_z),
 \end{aligned}$$

➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions



Which corresponds to two depictions at hadron level



Which corresponds to one depictions at hadron level

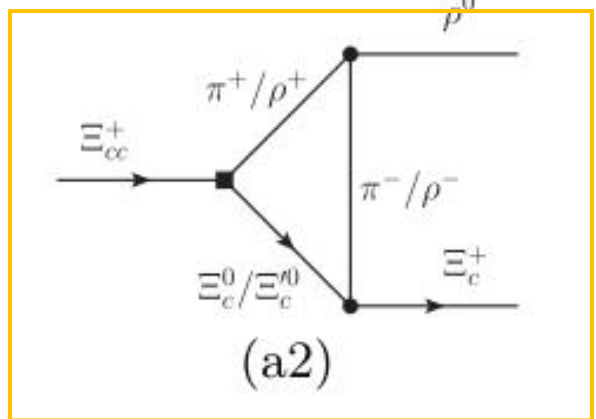
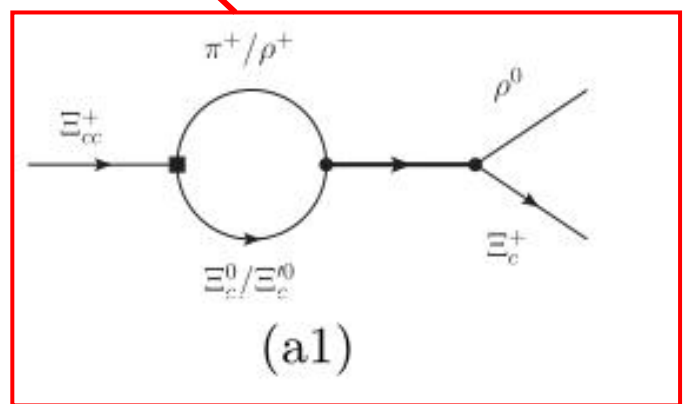
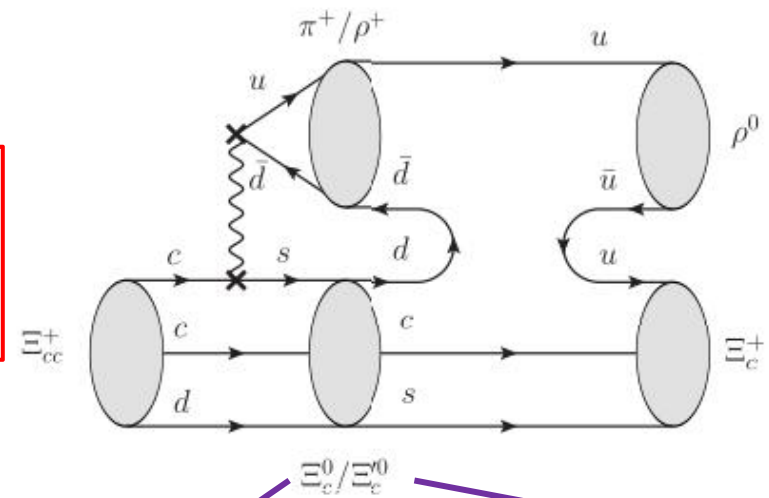
➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions

Ignored

Considered

Which with a resonance-like structure is expected to be highly suppressed

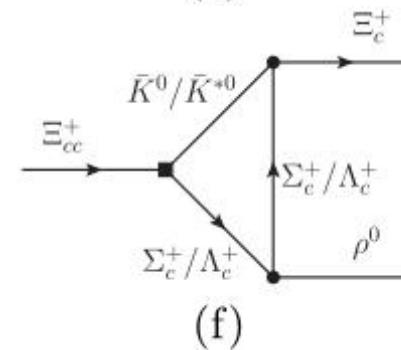
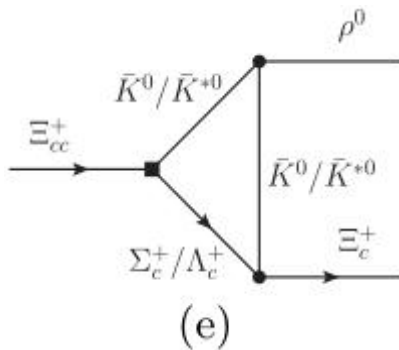
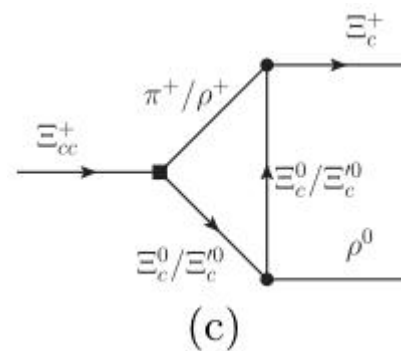
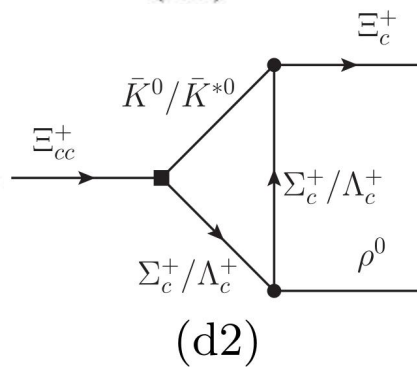
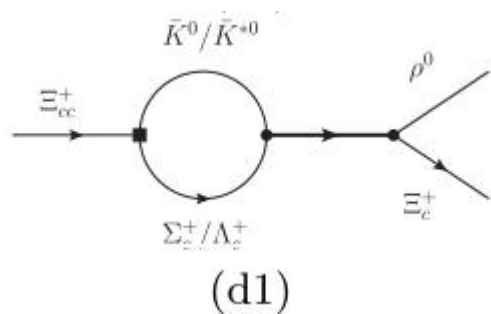
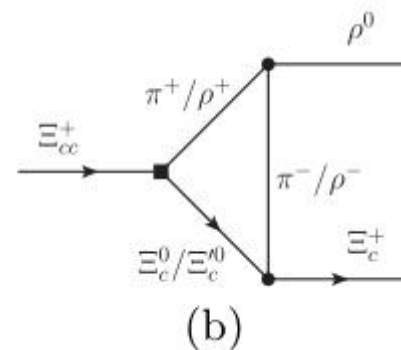
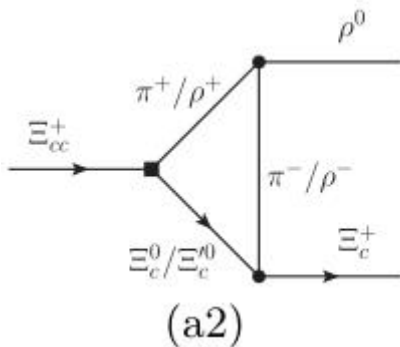
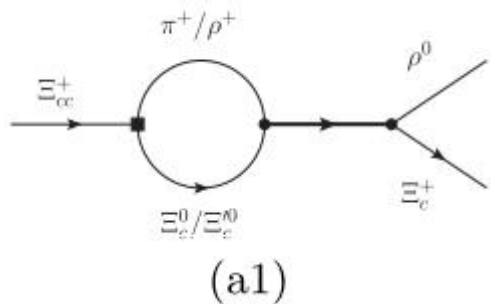
t-channel contribution with a light meson exchange



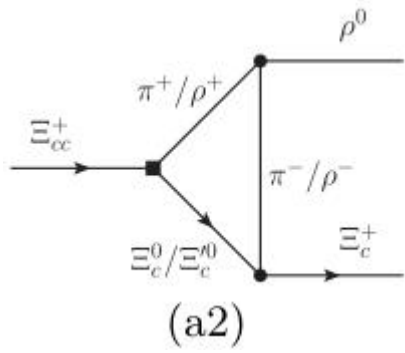
s-channel

t-channel

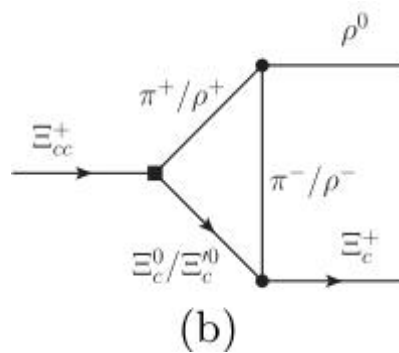
➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions



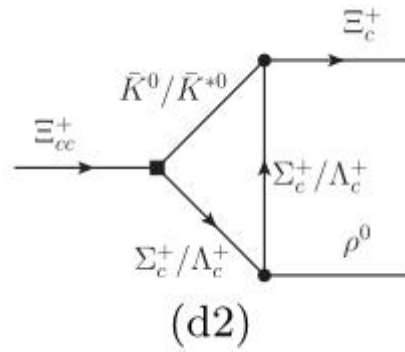
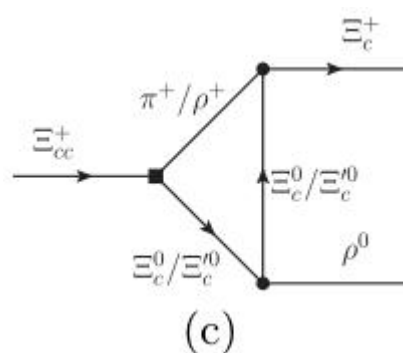
➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions



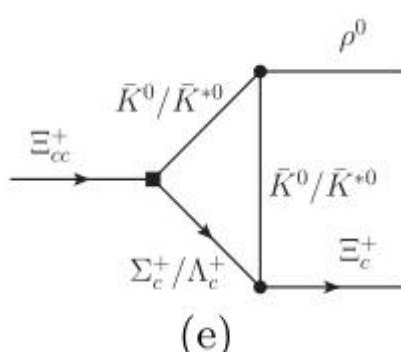
$u\bar{u}$



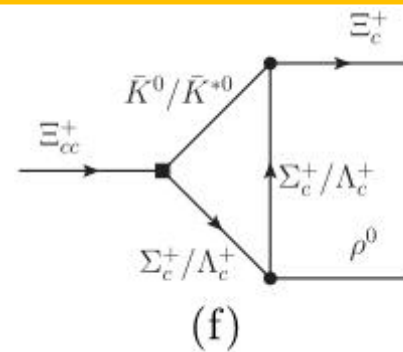
$d\bar{d}$



$u\bar{u}$

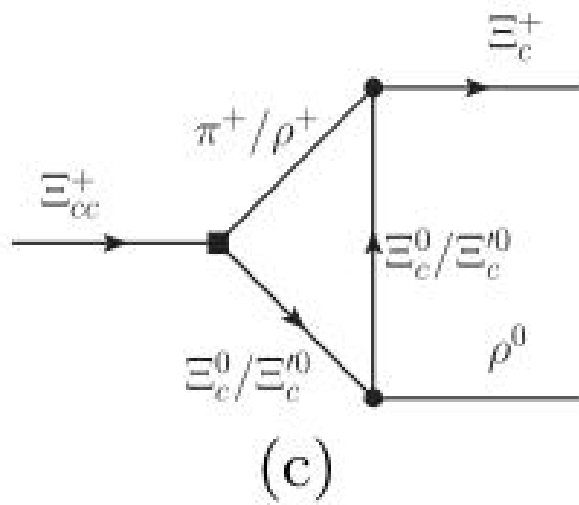


$d\bar{d}$

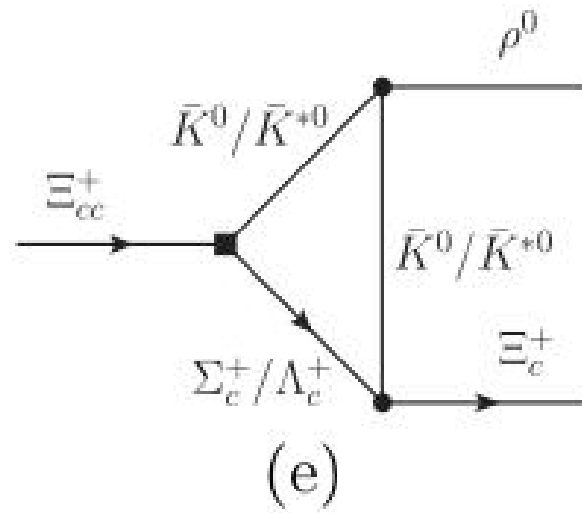


$u\bar{u}$ with an isospin factor $\frac{1}{\sqrt{2}}$, while $d\bar{d}$ with $-\frac{1}{\sqrt{2}}$

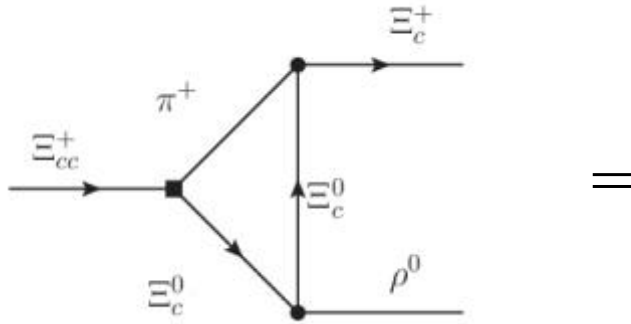
➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions



+



➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$: long Distance Contributions



$$\begin{aligned}
 \text{Abs } M_c(\pi^+; \Xi_c^0; \Xi_c^0) &= -\frac{i}{\sqrt{2}} \int \frac{|\vec{p}_2| \sin\theta d\theta d\varphi}{32\pi^2 m_{\Xi_{cc}^+}} \frac{G_F}{\sqrt{2}} V_{cs}^* V_{ud} a_1 f_\pi g_{\Xi_c^+ \Xi_c^0 \pi^+} \frac{F^2(t, m_{\Xi_c^0})}{t - m_{\Xi_c^0}^2} \epsilon_6^{*\nu} \\
 &\times \bar{u}(p_5, s'_z) \gamma_5 (\not{p}_4 + m_{\Xi_c^0}) \left(f_{1\Xi_c^0 \Xi_c^0 \rho^0} \gamma_\nu + i \frac{f_{2\Xi_c^0 \Xi_c^0 \rho^0}}{2m_{\Xi_c^0}} \sigma_{\mu\nu} p_6^\mu \right) \\
 &(\not{p}_3 + m_{\Xi_c^0}) \left[(m_{\Xi_{cc}^+} - m_{\Xi_c^0}) f_1(m_\pi^2) + (m_{\Xi_{cc}^+} + m_{\Xi_c^0}) g_1(m_\pi^2) \gamma_5 \right] u(p_1, s_z),
 \end{aligned}$$

where

$$F(t, m) = \left(\frac{\Lambda^2 - m^2}{\Lambda^2 - t} \right)^n, \quad \Lambda = m + \eta \Lambda_{\text{QCD}}$$

→ Cheng, Chua, Soni, PRD 71, 014030(2005)

➤ An Example $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$

The amplitude of $\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$ is expressed by

$$\begin{aligned} \mathcal{A}(\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0) = & C_{\text{SD}}(\Xi_{cc} \rightarrow \Xi_c^+ \rho^0) + i[\mathcal{A}bs M_c(\pi^+; \Xi_c^0; \Xi_c^0) + \mathcal{A}bs M_c(\rho^+; \Xi_c^0; \Xi_c^0) \\ & + \mathcal{A}bs M_c(\pi^+; \Xi_c^0; \Xi_c'^0) + \mathcal{A}bs M_c(\rho^+; \Xi_c^0; \Xi_c'^0) + \mathcal{A}bs M_c(\pi^+; \Xi_c'^0; \Xi_c^0) \\ & + \mathcal{A}bs M_c(\rho^+; \Xi_c'^0; \Xi_c^0) + \mathcal{A}bs M_c(\pi^+; \Xi_c'^0; \Xi_c'^0) + \mathcal{A}bs M_c(\rho^+; \Xi_c'^0; \Xi_c'^0) \\ & + \mathcal{A}bs M_e(\bar{K}^0; \Lambda_c^+; K^0) + \mathcal{A}bs M_e(\bar{K}^{*0}; \Lambda_c^+; K^{*0}) + \mathcal{A}bs M_e(\bar{K}^0; \Sigma_c^+; K^0) \\ & + \mathcal{A}bs M_e(\bar{K}^{*0}; \Sigma_c^+; K^{*0})]. \end{aligned}$$

➤ Inputs : Form Factors

Transition form factors with scalar (0^+) and the axial vector (1^+) diquarks.

F	$F(0)$	m_{fit}	δ	F	$F(0)$	m_{fit}	δ
$f_1^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	0.653	1.72	0.27	$f_2^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	-0.738	1.56	0.32
$g_1^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	0.533	2.03	0.38	$g_2^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	-0.053	1.12	1.10
$f_1^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	0.653	1.72	0.27	$f_2^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	-0.738	1.56	0.32
$g_1^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	0.533	2.03	0.38	$g_2^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	-0.053	1.12	1.10
$f_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	0.754	1.84	0.25	$f_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	-0.782	1.67	0.30
$g_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	0.620	2.16	0.35	$g_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	-0.080	1.29	0.52
$f_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	0.754	1.84	0.25	$f_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	-0.782	1.67	0.30
$g_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	0.620	2.16	0.35	$g_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	-0.080	1.29	0.52
$f_1^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	0.653	1.72	0.27	$f_2^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	-0.738	1.56	0.32
$g_1^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	0.533	2.03	0.38	$g_2^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	-0.053	1.12	1.10
$f_1^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	0.754	1.84	0.25	$f_2^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	-0.782	1.67	0.30
$g_1^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	0.620	2.16	0.35	$g_2^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	-0.080	1.29	0.52
$f_1^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.754	1.84	0.25	$f_2^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.782	1.67	0.30
$g_1^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.620	2.16	0.35	$g_2^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.080	1.29	0.52
$f_1^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	0.646	1.68	0.28	$f_2^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	-0.770	1.54	0.33
$g_1^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	0.528	1.99	0.40	$g_2^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	-0.060	1.12	1.02
$f_1^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.646	1.68	0.28	$f_2^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.770	1.54	0.33
$g_1^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.528	1.99	0.40	$g_2^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.060	1.12	1.02
$f_1^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	0.748	1.80	0.27	$f_2^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	-0.819	1.64	0.32
$g_1^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	0.615	2.11	0.36	$g_2^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	-0.088	1.28	0.52

F	$F(0)$	m_{fit}	δ	F	$F(0)$	m_{fit}	δ
$f_1^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	0.637	1.49	0.37	$f_2^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	0.725	1.53	0.32
$g_1^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	-0.167	1.99	0.23	$g_2^{\Xi_{cc}^{++} \rightarrow \Lambda_c^+}$	-0.028*	2.03*	2.62*
$f_1^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	0.637	1.49	0.37	$f_2^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	0.725	1.53	0.32
$g_1^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	-0.167	1.99	0.23	$g_2^{\Xi_{cc}^{++} \rightarrow \Sigma_c^+}$	-0.028*	2.03*	2.62*
$f_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	0.739	1.58	0.36	$f_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	0.801	1.62	0.31
$g_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	-0.198	2.10	0.21	$g_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^+}$	-0.018*	1.62*	1.37*
$f_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	0.739	1.58	0.36	$f_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	0.801	1.62	0.31
$g_1^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	-0.198	2.10	0.21	$g_2^{\Xi_{cc}^{++} \rightarrow \Xi_c^{\prime+}}$	-0.018*	1.62*	1.37*
$f_1^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	0.637	1.49	0.37	$f_2^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	0.725	1.53	0.32
$g_1^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	-0.167	1.99	0.23	$g_2^{\Xi_{cc}^+ \rightarrow \Sigma_c^0}$	-0.028*	2.03*	2.62*
$f_1^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	0.739	1.58	0.36	$f_2^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	0.801	1.62	0.31
$g_1^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	-0.198	2.10	0.21	$g_2^{\Xi_{cc}^+ \rightarrow \Xi_c^0}$	-0.018*	1.62*	1.37*
$f_1^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.739	1.58	0.36	$f_2^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.801	1.62	0.31
$g_1^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.198	2.10	0.21	$g_2^{\Xi_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.018*	1.62*	1.37*
$f_1^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	0.632	1.47	0.38	$f_2^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	0.734	1.52	0.33
$g_1^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	-0.165	1.97	0.27	$g_2^{\Omega_{cc}^+ \rightarrow \Xi_c^0}$	-0.031*	2.32*	3.92*
$f_1^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.632	1.47	0.38	$f_2^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	0.734	1.52	0.33
$g_1^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.165	1.97	0.27	$g_2^{\Omega_{cc}^+ \rightarrow \Xi_c^{\prime0}}$	-0.031*	2.32*	3.92*
$f_1^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	0.735	1.57	0.37	$f_2^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	0.812	1.61	0.32
$g_1^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	-0.196	2.08	0.24	$g_2^{\Omega_{cc}^+ \rightarrow \Omega_c^0}$	-0.021*	1.79*	1.77*

► Inputs : Strong Coupling Constants

TABLE V: Strong coupling constants of VPP and VVV vertices.

vertex	g	vertex	g	vertex	g	vertex	g	vertex	g	vertex	g
$\rho^+ \rightarrow \pi^0 \pi^+$	6.05	$\rho^0 \rightarrow \pi^- \pi^+$	6.05	$\rho^+ \rightarrow K^+ \bar{K}^0$	4.60	$\rho^0 \rightarrow K^0 \bar{K}^0$	-3.25	$\rho^0 \rightarrow K^+ K^-$	3.25	$\omega \rightarrow K^+ K^-$	3.25
$\phi \rightarrow K^- K^+$	4.60	$\bar{K}^{*0} \rightarrow \eta_8 \bar{K}^0$	5.63	$\bar{K}^{*0} \rightarrow K^- \pi^+$	4.60	$\bar{K}^{*0} \rightarrow \bar{K}^0 \pi^0$	-3.25	$K^{*+} \rightarrow K^+ \pi^0$	3.25	$\phi \rightarrow \bar{K}^0 K^0$	4.60
$K^{*+} \rightarrow \eta_8 K^+$	5.63	$K^{*+} \rightarrow \pi^+ K^0$	4.60	$K^{*0} \rightarrow \pi^- K^+$	4.60	$K^{*0} \rightarrow K^0 \eta_8$	5.63	$K^{*0} \rightarrow \pi^0 K^0$	-3.25	$\omega \rightarrow K^0 \bar{K}^0$	3.25
$\rho^+ \rightarrow \rho^0 \rho^+$	7.38	$\rho^0 \rightarrow \rho^- \rho^+$	7.38	$\rho^+ \rightarrow K^{*+} \bar{K}^{*0}$	5.22	$\rho^0 \rightarrow K^{*+} K^{*-}$	3.69	$\omega \rightarrow K^{*+} K^{*-}$	3.69	$\rho^0 \rightarrow K^{*0} \bar{K}^{*0}$	-3.69
$\bar{K}^{*0} \rightarrow \phi \bar{K}^{*0}$	5.22	$\bar{K}^{*0} \rightarrow \bar{K}^{*0} \rho^0$	-3.69	$\bar{K}^{*0} \rightarrow \bar{K}^{*0} \omega$	3.69	$K^{*+} \rightarrow \rho^+ K^{*0}$	5.22	$K^{*+} \rightarrow \phi K^{*+}$	5.22	$K^{*+} \rightarrow K^{*+} \rho^0$	3.69
$K^{*+} \rightarrow \omega K^{*+}$	3.69	$K^{*0} \rightarrow \rho^0 K^{*0}$	-3.69	$K^{*0} \rightarrow \omega K^{*0}$	3.69	$K^{*0} \rightarrow K^{*0} \phi$	5.22	$\phi \rightarrow K^{*-} K^{*+}$	5.22	$\bar{K}^{*0} \rightarrow K^{*-} \rho^+$	5.22
$\omega \rightarrow K^{*0} \bar{K}^{*0}$	3.69	$\phi \rightarrow \bar{K}^{*0} K^{*0}$	5.22								

TABLE VI: Strong coupling constants of $\mathcal{B}_{c3}\mathcal{B}_{c3}P$, $\mathcal{B}_{c3}\mathcal{B}_{c6}P$ and $\mathcal{B}_{c6}\mathcal{B}_{c6}P$ vertices.

vertex	g	vertex	g	vertex	g	vertex	g	vertex	g	vertex	g
$\Xi_c^0 \rightarrow \Xi_c^+ \pi^-$	0.99	$\Xi_c^+ \rightarrow \Xi_c^0 \pi^+$	0.99	$\Xi_c^0 \rightarrow \Xi_c^0 \pi^0$	-0.70	$\Xi_c^0 \rightarrow \Lambda_c^+ K^-$	-0.90	$\Xi_c^+ \rightarrow \Xi_c^+ \pi^0$	0.70	$\Xi_c^0 \rightarrow \Xi_c^0 \eta_8$	-0.70
$\Lambda_c^+ \rightarrow \Lambda_c^+ \eta_1$	0.75	$\Xi_c^+ \rightarrow \Xi_c^+ \eta_1$	0.07	$\Xi_c^0 \rightarrow \Xi_c^0 \eta_1$	0.07	$\Lambda_c^+ \rightarrow \Lambda_c^+ \eta_8$	0.81	$\Xi_c^+ \rightarrow \Xi_c^+ \eta_8$	-0.70	$\Xi_c^+ \rightarrow \Lambda_c^+ \bar{K}^0$	0.90
$\Sigma_c^0 \rightarrow \Sigma_c^+ \pi^-$	8.0	$\Xi_c^0 \rightarrow \Xi_c^0 \pi^0$	-4.0	$\Xi_c^0 \rightarrow \Sigma_c^0 \bar{K}^0$	9.0	$\Sigma_c^+ \rightarrow \Sigma_c^0 \pi^+$	8.0	$\Xi_c^0 \rightarrow \Sigma_c^+ K^-$	6.4	$\Xi_c^0 \rightarrow \Xi_c^+ \pi^-$	5.7
$\Sigma_c^0 \rightarrow \Sigma_c^0 \pi^0$	-8.0	$\Sigma_c^0 \rightarrow \Xi_c^0 K^0$	9.0	$\Xi_c^+ \rightarrow \Xi_c^+ \pi^0$	4.0	$\Xi_c^+ \rightarrow \Sigma_c^+ K^-$	9.0	$\Xi_c^+ \rightarrow \Xi_c^0 \pi^+$	5.7	$\Sigma_c^{++} \rightarrow \Sigma_c^+ \pi^+$	8.0
$\Sigma_c^{++} \rightarrow \Xi_c^+ K^+$	9.0	$\Sigma_c^+ \rightarrow \Xi_c^+ K^0$	6.4	$\Sigma_c^+ \rightarrow \Sigma_c^{++} \pi^-$	8.0	$\Xi_c^+ \rightarrow \Omega_c^0 K^+$	9.0	$\Xi_c^+ \rightarrow \Sigma_c^+ \bar{K}^0$	6.4	$\Sigma_c^+ \rightarrow \Xi_c^0 K^+$	6.4
$\Omega_c^0 \rightarrow \Xi_c^+ K^-$	9.0	$\Xi_c^0 \rightarrow \Omega_c^0 K^0$	9.0	$\Sigma_c^0 \rightarrow \Sigma_c^0 \eta_1$	-2.6	$\Omega_c^0 \rightarrow \Omega_c^0 \eta_1$	-11.0	$\Omega_c^0 \rightarrow \Xi_c^0 \bar{K}^0$	9.0	$\Sigma_c^0 \rightarrow \Sigma_c^0 \eta_8$	4.6
$\Omega_c^0 \rightarrow \Omega_c^0 \eta_8$	-10.4	$\Sigma_c^+ \rightarrow \Sigma_c^+ \eta_1$	-2.6	$\Xi_c^0 \rightarrow \Xi_c^0 \eta_1$	-2.6	$\Xi_c^+ \rightarrow \Xi_c^+ \eta_1$	-2.6	$\Sigma_c^+ \rightarrow \Sigma_c^+ \eta_8$	4.6	$\Xi_c^0 \rightarrow \Xi_c^0 \eta_8$	-2.3
$\Xi_c^0 \rightarrow \Xi_c^+ \pi^-$	4.4	$\Xi_c^+ \rightarrow \Xi_c^+ \pi^0$	3.1	$\Xi_c^+ \rightarrow \Xi_c^0 \pi^+$	4.4	$\Omega_c^0 \rightarrow \Xi_c^+ K^-$	6.5	$\Xi_c^+ \rightarrow \Lambda_c^+ \bar{K}^0$	-4.6	$\Sigma_c^+ \rightarrow \Lambda_c^+ \pi^0$	6.5
$\Xi_c^0 \rightarrow \Omega_c^0 K^0$	6.5	$\Omega_c^0 \rightarrow \Xi_c^0 \bar{K}^0$	6.5	$\Sigma_c^0 \rightarrow \Xi_c^0 K^0$	-7.1	$\Xi_c^+ \rightarrow \Lambda_c^+ K^0$	-4.6	$\Xi_c^+ \rightarrow \Xi_c^0 \pi^+$	4.4	$\Lambda_c^+ \rightarrow \Xi_c^0 K^+$	4.6
$\Xi_c^0 \rightarrow \Sigma_c^+ K^-$	-5.0	$\Xi_c^0 \rightarrow \Xi_c^0 \pi^0$	-3.1	$\Xi_c^+ \rightarrow \Xi_c^0 \pi^+$	4.4	$\Xi_c^0 \rightarrow \Lambda_c^+ K^-$	4.6	$\Xi_c^0 \rightarrow \Xi_c^+ \pi^-$	4.4	$\Xi_c^+ \rightarrow \Omega_c^0 K^+$	6.5
$\Sigma_c^{++} \rightarrow \Xi_c^+ K^+$	-7.1	$\Xi_c^+ \rightarrow \Sigma_c^{++} K^-$	-7.1	$\Sigma_c^{++} \rightarrow \Lambda_c^+ \pi^+$	-6.5	$\Xi_c^0 \rightarrow \Sigma_c^0 \bar{K}^0$	-7.1	$\Lambda_c^+ \rightarrow \Sigma_c^{++} \pi^-$	-6.5	$\Lambda_c^+ \rightarrow \Sigma_c^0 \pi^+$	6.5
$\Sigma_c^+ \rightarrow \Xi_c^0 K^+$	-5.0	$\Sigma_c^0 \rightarrow \Lambda_c^+ \pi^-$	6.5	$\Xi_c^+ \rightarrow \Xi_c^+ \eta_8$	5.4	$\Sigma_c^+ \rightarrow \Xi_c^+ K^0$	-5.0	$\Xi_c^+ \rightarrow \Sigma_c^+ \bar{K}^0$	-5.0	$\Xi_c^0 \rightarrow \Xi_c^0 \eta_8$	5.4
$\Lambda_c^+ \rightarrow \Xi_c^0 K^+$	-0.90	$\Xi_c^+ \rightarrow \Lambda_c^+ K^0$	0.90	$\Xi_c^+ \rightarrow \Xi_c^+ \eta_8$	-2.3						

► Inputs : Strong Coupling Constants

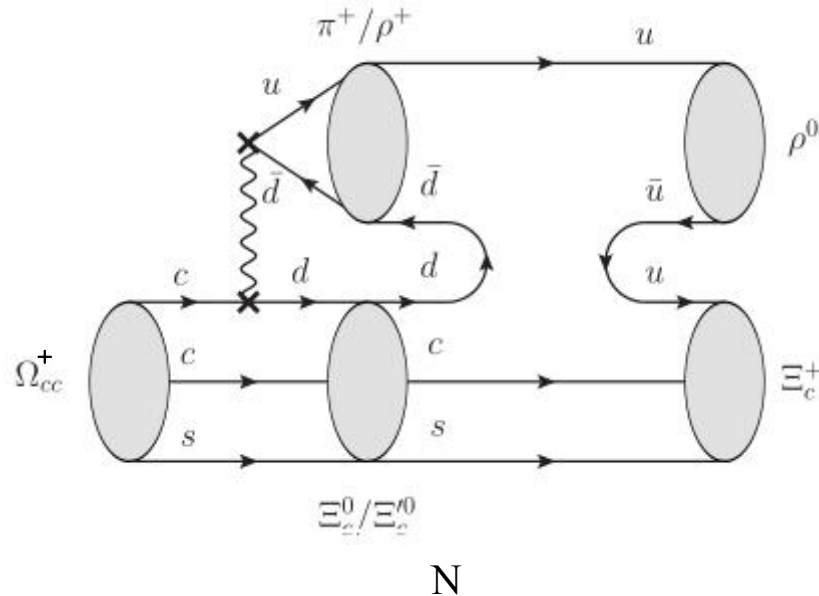
TABLE VII: The strong coupling constants of $\mathcal{B}_{c3}\mathcal{B}_{c3}V$ vertices, each of which owns two $SU(3)$ triplet singly charm baryons and one light vector meson.

vertex	f_1	f_2	vertex	f_1	f_2	vertex	f_1	f_2	vertex	f_1	f_2
$\Xi_c^0 \rightarrow \Lambda_c^+ K^{*-}$	-4.6	-6.0	$\Xi_c^0 \rightarrow \Xi_c^0 \rho^0$	-6.0	-7.5	$\Xi_c^+ \rightarrow \Xi_c^0 \rho^+$	8.5	10.6	$\Lambda_c^+ \rightarrow \Xi_c^0 K^{*+}$	-4.6	-6.0
$\Xi_c^0 \rightarrow \Xi_c^0 \phi$	4.6	6.0	$\Xi_c^0 \rightarrow \Xi_c^0 \omega$	5.5	7.5	$\Lambda_c^+ \rightarrow \Lambda_c^+ \omega$	4.9	6.0	$\Xi_c^0 \rightarrow \Xi_c^+ \rho^-$	8.5	10.6
$\Xi_c^+ \rightarrow \Lambda_c^+ \bar{K}^{*0}$	4.6	6.0	$\Xi_c^+ \rightarrow \Xi_c^+ \rho^0$	6.0	7.5	$\Xi_c^+ \rightarrow \Xi_c^+ \omega$	5.5	7.5	$\Xi_c^+ \rightarrow \Xi_c^+ \phi$	4.6	6.0
$\Sigma_c^0 \rightarrow \Xi_c^{\prime 0} K^{*0}$	5.0	30.0	$\Xi_c^{\prime 0} \rightarrow \Xi_c^{\prime 0} \rho^0$	-2.5	-16.0	$\Xi_c^{\prime 0} \rightarrow \Xi_c^{\prime +} \rho^-$	3.5	22.6	$\Xi_c^{\prime 0} \rightarrow \Xi_c^{\prime 0} \phi$	4.0	21.0
$\Xi_c^{\prime +} \rightarrow \Xi_c^{\prime 0} \rho^+$	3.5	22.6	$\Sigma_c^+ \rightarrow \Xi_c^{\prime 0} K^{*+}$	3.5	21.2	$\Xi_c^{\prime 0} \rightarrow \Xi_c^{\prime 0} \omega$	2.4	15.0	$\Xi_c^{\prime 0} \rightarrow \Sigma_c^+ K^{*-}$	3.5	21.2
$\Xi_c^{\prime 0} \rightarrow \Sigma_c^0 \bar{K}^{*0}$	5.0	30.0	$\Sigma_c^+ \rightarrow \Sigma_c^{++} \rho^-$	4.0	27.0	$\Sigma_c^{++} \rightarrow \Xi_c^{\prime +} K^{*+}$	5.0	30.0	$\Xi_c^{\prime +} \rightarrow \Sigma_c^{++} K^{*-}$	5.0	30.0
$\Xi_c^{\prime +} \rightarrow \Sigma_c^+ \bar{K}^{*0}$	3.5	21.2	$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \rho^0$	2.5	16.0	$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \phi$	4.0	21.0	$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \omega$	2.4	15.0
$\Sigma_c^+ \rightarrow \Xi_c^{\prime +} K^{*0}$	3.5	21.2	$\Sigma_c^{++} \rightarrow \Sigma_c^+ \rho^+$	4.0	27.0	$\Sigma_c^0 \rightarrow \Sigma_c^0 \rho^0$	-4.0	-27.0	$\Sigma_c^0 \rightarrow \Sigma_c^0 \omega$	3.5	24.0
$\Sigma_c^+ \rightarrow \Sigma_c^0 \rho^+$	4.0	27.0	$\Sigma_c^0 \rightarrow \Sigma_c^+ \rho^-$	4.0	27.0	$\Xi_c^{\prime +} \rightarrow \Omega_c^0 K^{*+}$	7.0	35.0	$\Omega_c^0 \rightarrow \Xi_c^{\prime 0} K^{*0}$	7.0	35.0
$\Omega_c^0 \rightarrow \Omega_c^0 \phi$	11.0	52.0	$\Omega_c^0 \rightarrow \Xi_c^{\prime +} K^{*-}$	7.0	35.0	$\Omega_c^0 \rightarrow \Xi_c^{\prime 0} \bar{K}^{*0}$	7.0	35.0	$\Sigma_c^+ \rightarrow \Sigma_c^+ \omega$	3.5	24.0
$\Xi_c^{\prime 0} \rightarrow \Xi_c^0 \rho^0$	-1.5	-11.0	$\Xi_c^{\prime 0} \rightarrow \Xi_c^0 \phi$	-2.1	-13.0	$\Xi_c^{\prime 0} \rightarrow \Xi_c^0 \omega$	1.2	8.0	$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \omega$	1.5	11.0
$\Lambda_c^+ \rightarrow \Xi_c^{\prime 0} K^{*+}$	2.3	14.1	$\Xi_c^{\prime 0} \rightarrow \Lambda_c^+ K^{*-}$	2.3	14.1	$\Xi_c^+ \rightarrow \Xi_c^0 \rho^+$	2.1	15.6	$\Xi_c^{\prime 0} \rightarrow \Xi_c^+ \rho^-$	2.1	15.6
$\Xi_c^{\prime +} \rightarrow \Lambda_c^+ \bar{K}^{*0}$	-2.3	-14.1	$\Sigma_c^+ \rightarrow \Xi_c^0 K^{*+}$	-2.2	-13.0	$\Xi_c^+ \rightarrow \Sigma_c^+ K^{*0}$	-2.2	-13.0	$\Xi_c^{\prime +} \rightarrow \Xi_c^0 \rho^+$	2.1	15.6
$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \rho^0$	1.5	11.0	$\Xi_c^{\prime +} \rightarrow \Xi_c^+ \phi$	-2.1	-13.0	$\Xi_c^0 \rightarrow \Xi_c^+ \rho^-$	2.1	15.6	$\Sigma_c^+ \rightarrow \Xi_c^+ \bar{K}^{*0}$	-2.2	-13.0
$\Sigma_c^+ \rightarrow \Lambda_c^+ \rho^0$	2.6	16.0	$\Xi_c^+ \rightarrow \Omega_c^0 K^{*+}$	3.3	20.0	$\Lambda_c^+ \rightarrow \Xi_c^{\prime +} K^{*0}$	-2.3	-14.1	$\Xi_c^0 \rightarrow \Sigma_c^+ K^{*-}$	-2.2	-13.0
$\Sigma_c^{++} \rightarrow \Xi_c^+ K^{*+}$	-3.1	-18.4	$\Xi_c^+ \rightarrow \Sigma_c^{++} K^{*-}$	-3.1	-18.4	$\Sigma_c^{++} \rightarrow \Lambda_c^+ \rho^+$	-2.6	-16.0	$\Lambda_c^+ \rightarrow \Sigma_c^{++} \rho^-$	-2.6	-16.0
$\Lambda_c^+ \rightarrow \Sigma_c^0 \rho^+$	2.6	16.0	$\Sigma_c^0 \rightarrow \Lambda_c^+ \rho^-$	2.6	16.0	$\Omega_c^0 \rightarrow \Xi_c^+ K^{*-}$	3.3	20.0	$\Xi_c^0 \rightarrow \Sigma_c^0 \bar{K}^{*0}$	-2.2	-13.0
$\Xi_c^0 \rightarrow \Omega_c^0 K^{*0}$	3.3	20.0	$\Sigma_c^0 \rightarrow \Xi_c^0 K^{*0}$	-2.2	-13.0	$\Omega_c^0 \rightarrow \Xi_c^0 \bar{K}^{*0}$	3.3	20.0	$\Lambda_c^+ \rightarrow \Xi_c^+ K^{*0}$	4.6	6.0

- [1] H. Y. Cheng, C. K. Chua and A. Soni, Phys. Rev. D 71, 014030 (2005)
[2] T. M. Aliev, K. Azizi and M. Savci, Phys. Lett. B 696, 220 (2011) doi:10.1016/j.physletb.2010.12.027
[3] T. M. Aliev, K. Azizi and M. Savci, Nucl. Phys. A 852, 141 (2011) doi:10.1016/j.nuclphysa.2011.01.011

➤ Topological FCNC contributions Induced by FSI

An Example:



- The weak decay is induced at quark level by $c \rightarrow u d \bar{d}$. The $d \bar{d}$ pair annihilates subsequently, and the weak transition is equivalent to a flavor changing neutral current $c \rightarrow u \gamma / G$. The same situation exists in $c \rightarrow u s \bar{s}$ inducing decays.
- Topologically, it looks the same as $O_{7\gamma}$ or O_{8G} operators in the low energy effective hamiltonian, which are purely loop effects and suppressed highly at short distance in the standard model.
- However, it is the long distance contributions that are considered here, thus they are kept.

3 Results and Discussions

For lack of experimental data we use $\eta = 1.5$ in this work and range it from 1 to 2 for error estimation.

Channels	$\mathcal{BR}(\%)$	Contributions	CKM	Channels	$\mathcal{BR}(\%)$	Contributions	CKM
$\Xi_{cc}^{++} \rightarrow \Sigma_c^{++} \bar{K}^{*0}$	$5.40^{+5.59}_{-3.66}$	C_{SD}, C	CF	$\Xi_{cc}^{++} \rightarrow \Xi_c^+ \rho^+$	$15.92^{+5.27}_{-3.03}$	T_{SD}, T, C'	CF
$\Xi_{cc}^{++} \rightarrow \Xi_c^+ \rho^+$	$16.50^{+1.21}_{-0.69}$	T_{SD}, T, C'	CF	$\Xi_{cc}^{++} \rightarrow \Sigma_c^+ \rho^+$	$1.57^{+0.63}_{-0.40}$	T_{SD}, T, C', N	SCS
$\Xi_{cc}^{++} \rightarrow \Lambda_c^+ \rho^+$	$2.10^{+1.32}_{-0.82}$	T_{SD}, T, C', N	SCS	$\Xi_{cc}^{++} \rightarrow \Sigma_c^{++} \rho^0$	$0.04^{+0.04}_{-0.03}$	C_{SD}, C, N	SCS
$\Xi_{cc}^{++} \rightarrow \Sigma_c^{++} \omega$	$0.15^{+0.17}_{-0.10}$	C_{SD}, C, N	SCS	$\Xi_{cc}^{++} \rightarrow \Sigma_c^{++} \phi$	$0.09^{+0.08}_{-0.06}$	C_{SD}, C	SCS
$\Xi_{cc}^{++} \rightarrow \Xi_c^+ K^{*+}$	$0.59^{+0.14}_{-0.10}$	T_{SD}, T, N, C'	SCS	$\Xi_{cc}^{++} \rightarrow \Xi_c^+ K^{*+}$	$0.85^{+0.14}_{-0.09}$	T_{SD}, T, N, C'	SCS
$\Xi_{cc}^{++} \rightarrow \Sigma_c^+ K^{*+}$	$0.06^{+0.00}_{-0.01}$	T_{SD}, T, C'	DCS	$\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^{*+}$	$0.05^{+0.00}_{-0.00}$	T_{SD}, T, C'	DCS
$\Xi_{cc}^{++} \rightarrow \Sigma_c^{++} K^{*0}$	$0.02^{+0.02}_{-0.01}$	C_{SD}, C	DCS				

- The \mathcal{BR} of singly CKM suppressed decays range from order of 10^{-3} to 10^{-2} . The doubly CKM suppressed decays have the smallest \mathcal{BR} at the order of 10^{-4} .
- Among decays in the same CKM mode, those with T contributions tend to have largest \mathcal{BR} . The C type decays are about several times smaller than the T type. The other types of decays are suppressed highly.
- Checking the errors, one can find that the T contributions are not sensitive to the variation of η . However, the other types of contributions which increase or decrease rapidly as η changes.
- $\Xi_{cc}^{++} \rightarrow \Xi_c^+ \rho^+$ and $\Xi_{cc}^{++} \rightarrow \Xi_c^+ \rho^+$ have the largest \mathcal{BR} , whose values are around 16% (ρ^+ may have a low reconstruction efficiency in experiments).

3 Results and Discussions

Because there is no experimental data for lifetime of Ξ_{cc}^+ . Instead of branching fractions we present the decay widths in unit of GeV.

Channels	Γ/GeV	Contributions	CKM	Channels	Γ/GeV	Contributions	CKM
$\Xi_{cc}^+ \rightarrow \Sigma_c^+ \bar{K}^{*0}$	$(1.37^{+1.46}_{-0.93}) * 10^{-13}$	C_{SD}, C, E_2	CF	$\Xi_{cc}^+ \rightarrow \Lambda_c^+ \bar{K}^{*0}$	$(1.04^{+1.14}_{-0.72}) * 10^{-13}$	C_{SD}, C, E_2	CF
$\Xi_{cc}^+ \rightarrow \Xi_c^0 \rho^+$	$(3.83^{+0.48}_{-0.37}) * 10^{-13}$	T_{SD}, T, E_1	CF	$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime 0} \rho^+$	$(4.77^{+0.31}_{-0.24}) * 10^{-13}$	T_{SD}, T, E_1	CF
$\Xi_{cc}^+ \rightarrow \Xi_c^+ \rho^0$	$(1.01^{+1.15}_{-0.71}) * 10^{-14}$	C', E_1	CF	$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime +} \rho^0$	$(1.59^{+1.76}_{-1.10}) * 10^{-15}$	C', E_1	CF
$\Xi_{cc}^+ \rightarrow \Xi_c^+ \omega$	$(7.82^{+8.98}_{-5.46}) * 10^{-15}$	C', E_1	CF	$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime +} \omega$	$(1.33^{+1.46}_{-0.92}) * 10^{-15}$	C', E_1	CF
$\Xi_{cc}^+ \rightarrow \Sigma_c^{++} K^{*-}$	$(7.38^{+7.83}_{-5.02}) * 10^{-16}$	E_2	CF	$\Xi_{cc}^+ \rightarrow \Xi_c^+ \phi$	$(5.12^{+5.59}_{-3.52}) * 10^{-15}$	E_2	CF
$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime +} \phi$	$(9.90^{+10.24}_{-6.72}) * 10^{-17}$	E_2	CF	$\Xi_{cc}^+ \rightarrow \Omega_c^0 K^{*+}$	$(2.33^{+2.16}_{-1.54}) * 10^{-14}$	E_1	CF
$\Xi_{cc}^+ \rightarrow \Sigma_c^+ \rho^0$	$(1.26^{+1.43}_{-0.88}) * 10^{-14}$	$C_{SD}, C, C', E_1, E_2, N$	SCS	$\Xi_{cc}^+ \rightarrow \Lambda_c^+ \rho^0$	$(4.97^{+5.93}_{-3.50}) * 10^{-15}$	$C_{SD}, C, C', E_1, E_2, N$	SCS
$\Xi_{cc}^+ \rightarrow \Sigma_c^+ \omega$	$(3.22^{+3.59}_{-2.23}) * 10^{-15}$	$C_{SD}, C, C', E_1, E_2, N$	SCS	$\Xi_{cc}^+ \rightarrow \Lambda_c^+ \omega$	$(1.60^{+1.91}_{-1.13}) * 10^{-15}$	$C_{SD}, C, C', E_1, E_2, N$	SCS
$\Xi_{cc}^+ \rightarrow \Sigma_c^0 \rho^+$	$(9.02^{+3.70}_{-2.45}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS	$\Xi_{cc}^+ \rightarrow \Sigma_c^+ \phi$	$(1.54^{+1.40}_{-1.01}) * 10^{-15}$	C_{SD}, C	SCS
$\Xi_{cc}^+ \rightarrow \Lambda_c^+ \phi$	$(2.61^{+2.67}_{-1.76}) * 10^{-15}$	C_{SD}, C	SCS	$\Xi_{cc}^+ \rightarrow \Xi_c^0 K^{*+}$	$(1.30^{+0.00}_{-0.00}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS
$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime 0} K^{*+}$	$(2.19^{+0.19}_{-0.10}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS	$\Xi_{cc}^+ \rightarrow \Xi_c^+ K^{*0}$	$(1.06^{+0.97}_{-0.69}) * 10^{-15}$	C', E_2, N	SCS
$\Xi_{cc}^+ \rightarrow \Xi_c^{\prime +} K^{*0}$	$(2.64^{+2.66}_{-1.77}) * 10^{-15}$	C', E_2, N	SCS	$\Xi_{cc}^+ \rightarrow \Sigma_c^{++} \rho^-$	$(7.60^{+8.83}_{-5.33}) * 10^{-16}$	E_2, N	SCS
$\Xi_{cc}^+ \rightarrow \Sigma_c^+ K^{*0}$	$(1.96^{+2.03}_{-1.33}) * 10^{-15}$	C_{SD}, C, C'	DCS	$\Xi_{cc}^+ \rightarrow \Lambda_c^+ K^{*0}$	$(9.99^{+11.18}_{-6.95}) * 10^{-16}$	C_{SD}, C, C'	DCS
$\Xi_{cc}^+ \rightarrow \Sigma_c^0 K^{*+}$	$(2.88^{+0.00}_{-0.00}) * 10^{-16}$	T_{SD}, T	DCS				

➤ Estimated with the $\tau_{\Xi_{cc}^+} = 45 \text{ fs}$, the four largest branching fractions are given as

$$BR(\Xi_{cc}^+ \rightarrow \Sigma_c^+ \bar{K}^{*0}) \in [0.3\%, 1.9\%], \quad BR(\Xi_{cc}^+ \rightarrow \Lambda_c^+ \bar{K}^{*0}) \in [0.2\%, 1.5\%],$$

$$BR(\Xi_{cc}^+ \rightarrow \Xi_c^0 \rho^+) \in [2.4\%, 2.9\%], \quad BR(\Xi_{cc}^+ \rightarrow \Xi_c^{\prime 0} \rho^+) \in [3.1\%, 3.5\%].$$

➤ Some pure W exchange decays which are highly suppressed at short distance. These decays are thought to be activated almost by the long distance effects.

3 Results and Discussions

Channels	Γ/GeV	Contributions	CKM	Channels	Γ/GeV	Contributions	CKM
$\Omega_{cc}^+ \rightarrow \Xi_c^+ \bar{K}^{*0}$	$(5.53_{-3.80}^{+5.96}) * 10^{-13}$	C_{SD}, C, C'	CF	$\Omega_{cc}^+ \rightarrow \Xi_c^+ \bar{K}^{*0}$	$(1.06_{-0.72}^{+1.09}) * 10^{-12}$	C_{SD}, C, C'	CF
$\Omega_{cc}^+ \rightarrow \Omega_c^0 \rho^+$	$(8.75_{-0.00}^{+0.00}) * 10^{-13}$	T_{SD}, T	CF	$\Omega_{cc}^+ \rightarrow \Sigma_c^+ \bar{K}^{*0}$	$(3.26_{-2.26}^{+3.62}) * 10^{-15}$	C', E_2, N	SCS
$\Omega_{cc}^+ \rightarrow \Lambda_c^+ \bar{K}^{*0}$	$(2.59_{-1.85}^{+3.21}) * 10^{-16}$	C', E_2, N	SCS	$\Omega_{cc}^+ \rightarrow \Xi_c^+ \rho^0$	$(3.04_{-2.13}^{+3.55}) * 10^{-15}$	C_{SD}, C, E_1, N	SCS
$\Omega_{cc}^+ \rightarrow \Xi_c^+ \rho^0$	$(2.75_{-1.89}^{+2.96}) * 10^{-15}$	C_{SD}, C, E_1, N	SCS	$\Omega_{cc}^+ \rightarrow \Xi_c^+ \omega$	$(2.38_{-1.67}^{+2.78}) * 10^{-15}$	C_{SD}, C, E_1, N	SCS
$\Omega_{cc}^+ \rightarrow \Xi_c^+ \omega$	$(2.53_{-1.70}^{+2.74}) * 10^{-15}$	C_{SD}, C, E_1, N	SCS	$\Omega_{cc}^+ \rightarrow \Xi_c^0 \rho^+$	$(4.45_{-1.57}^{+2.42}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS
$\Omega_{cc}^+ \rightarrow \Xi_c^0 \rho^+$	$(5.59_{-1.98}^{+3.04}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS	$\Omega_{cc}^+ \rightarrow \Xi_c^+ \phi$	$(8.96_{-5.98}^{+8.69}) * 10^{-15}$	$C_{\text{SD}}, C, C', E_2, N$	SCS
$\Omega_{cc}^+ \rightarrow \Xi_c^+ \phi$	$(4.54_{-3.01}^{+4.33}) * 10^{-14}$	$C_{\text{SD}}, C, C', E_2, N$	SCS	$\Omega_{cc}^+ \rightarrow \Omega_c^0 K^{*+}$	$(4.18_{-0.01}^{+0.03}) * 10^{-14}$	T_{SD}, T, E_1, N	SCS
$\Omega_{cc}^+ \rightarrow \Sigma_c^{++} K^{*-}$	$(6.77_{-4.71}^{+7.64}) * 10^{-16}$	E_2, N	SCS	$\Omega_{cc}^+ \rightarrow \Xi_c^+ K^{*0}$	$(8.17_{-5.63}^{+8.97}) * 10^{-16}$	C_{SD}, C, E_2, N	DCS
$\Omega_{cc}^+ \rightarrow \Sigma_c^+ \phi$	$(8.45_{-5.81}^{+9.15}) * 10^{-17}$	C'	DCS	$\Omega_{cc}^+ \rightarrow \Lambda_c^+ \phi$	$(4.25_{-2.93}^{+4.64}) * 10^{-17}$	C'	DCS
$\Omega_{cc}^+ \rightarrow \Xi_c^0 K^{*+}$	$(1.00_{-0.01}^{+0.02}) * 10^{-15}$	T_{SD}, T, E_1	DCS	$\Omega_{cc}^+ \rightarrow \Xi_c^0 K^{*+}$	$(1.47_{-0.03}^{+0.05}) * 10^{-15}$	T_{SD}, T, E_1	DCS
$\Omega_{cc}^+ \rightarrow \Xi_c^+ K^{*0}$	$(8.13_{-5.55}^{+8.64}) * 10^{-16}$	C_{SD}, C, E_2, N	DCS	$\Omega_{cc}^+ \rightarrow \Sigma_c^{++} \rho^-$	$(1.20_{-0.84}^{+1.38}) * 10^{-17}$	E_2, N	DCS
$\Omega_{cc}^+ \rightarrow \Sigma_c^+ \rho^0$	$(9.06_{-6.41}^{+10.75}) * 10^{-18}$	E_1, E_2, N	DCS	$\Omega_{cc}^+ \rightarrow \Lambda_c^+ \rho^0$	$(1.03_{-0.74}^{+1.33}) * 10^{-17}$	E_1, E_2, N	DCS
$\Omega_{cc}^+ \rightarrow \Sigma_c^+ \omega$	$(8.72_{-6.15}^{+10.25}) * 10^{-18}$	E_1, E_2, N	DCS	$\Omega_{cc}^+ \rightarrow \Lambda_c^+ \omega$	$(8.78_{-6.31}^{+11.19}) * 10^{-18}$	E_1, E_2, N	DCS
$\Omega_{cc}^+ \rightarrow \Sigma_c^0 \rho^+$	$(1.39_{-0.97}^{+1.62}) * 10^{-16}$	E_1	DCS				

➤ Estimated with $\tau_{\Omega_{cc}^+} = 75 \text{ fs}$, the branching ratios of the three decays are given as

$$\mathcal{BR}(\Omega_{cc}^+ \rightarrow \Xi_c^+ \bar{K}^{*0}) \in [2.0\%, 13.1\%], \quad \mathcal{BR}(\Omega_{cc}^+ \rightarrow \Xi_c^+ \bar{K}^{*0}) \in [3.9\%, 24.5\%], \quad \mathcal{BR}(\Omega_{cc}^+ \rightarrow \Omega_c^0 \rho^+) \approx 10.0\%.$$



$$\Omega_{cc}^+ \rightarrow pK^+ \pi^- K^- \pi^+$$

4 Summary

- Calculation: **the \mathcal{BR} and Γ of $\mathcal{B}_{cc} \rightarrow \mathcal{B}_c V$**
- Theoretical framework: **The factorization hypothesis**
The final-state interactions (FSIs)
with the one-particle-exchange model
- Results and Discussions:

➤ Results and Discussions:	CF	Large \mathcal{BR}	} T, C
	SCS	Smaller \mathcal{BR}	
	DCS	Smallest \mathcal{BR}	
- The largest \mathcal{BR} s of $\mathcal{B}_{cc} \rightarrow \mathcal{B}_c V$ decays are at the percent level.
- It seems easy to track $\Omega_{cc}^+ \rightarrow pK^+ \pi^- K^- \pi^+$ in experiments.

**Thank you for
your attention!**