

Chiral kinetic theory and magnetic effect

Yoshimasa Hidaka
(RIKEN)

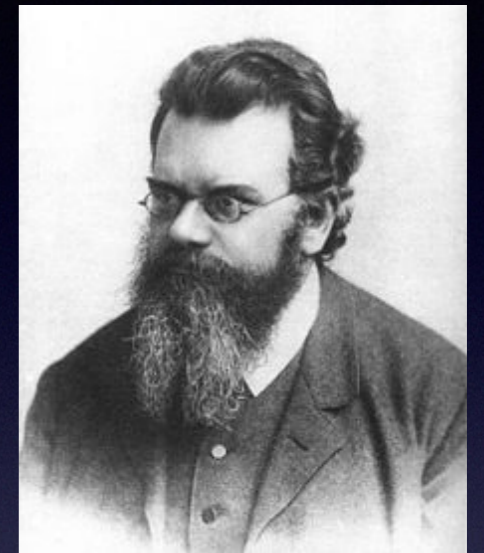
What is chiral kinetic theory?

Relativistic Boltzmann equation

$$(v^\mu \partial_\mu + v^\mu F_{\nu\mu} \partial_{p_\nu}) f = C[f]$$

widely used in plasma physics

Transport coefficient: shear viscosity, etc..



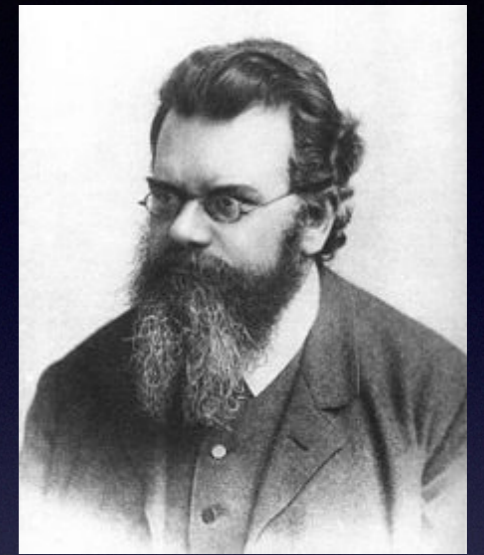
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Relativistic Boltzmann equation with quantum anomaly?

Anomaly matching and effective theory

If UV theory has an anomaly, 't Hooft ('80)
IR theory has the same anomaly.

Anomaly matching and effective theory

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UV theory

QCD

Chiral anomaly: $\partial_\mu j_5^{3\mu} = CE \cdot B$

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Effective theory

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Effective theory

Vacuum

Chiral perturbation
theory

Wess-Zumino term

$$\pi^0 \rightarrow 2\gamma$$

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Manybody systems

Hydrodynamics

Kinetic theory

Anomalous transport

Berry curvature

Chiral magnetic effect (CME)
Chiral vortical effect (CVE)

Equilibrium: CME, CVE
Nonequilibrium: ?

Chiral Kinetic theory

Son, Yamamoto ('12) Stephanov, Yin ('12)

cf. Chang and Niu ('95)

Hamiltonian $H = \boldsymbol{\sigma} \cdot \boldsymbol{p}$

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 $Hu_{\pm} = \pm |\boldsymbol{p}| u_{\pm}$

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 $Hu_{\pm} = \pm |\boldsymbol{p}| u_{\pm}$

When \boldsymbol{p} is time independent:

Wave function

$$\psi(t) = e^{-i|\boldsymbol{p}|t} u_{+}$$

Chiral Kinetic theory

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Hamiltonian $H = \boldsymbol{\sigma} \cdot \boldsymbol{p}$

➔ $Hu_{\pm} = \pm |\boldsymbol{p}| u_{\pm}$

When \boldsymbol{p} is weakly time dependent:

Wave function

$$\psi(t) \approx e^{-i|\boldsymbol{p}|t - i \int^t dt' \boldsymbol{a} \cdot \dot{\boldsymbol{p}}} u_{+}$$

Berry connection: $\boldsymbol{a} := -u_{+}^{\dagger} i \nabla_{\boldsymbol{p}} u_{+}$

Berry('84)

Chiral Kinetic theory

Son, Yamamoto ('12) Stephanov, Yin ('12)

cf. Chang and Niu ('95)

$$S = \int dt (\dot{\mathbf{x}} \cdot \mathbf{p} + \dot{\mathbf{x}} \cdot \mathbf{A} - |\mathbf{p}| - A_0 - \dot{\mathbf{p}} \cdot \mathbf{a})$$

Action with Berry connection: $\mathbf{a} := -u_+^\dagger i \nabla_p u_+$

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Classical EOM $\dot{\mathbf{x}} = \hat{\mathbf{p}} + \dot{\mathbf{p}} \times \boldsymbol{\Omega}$

$$\dot{\mathbf{p}} = \dot{\mathbf{x}} \times \mathbf{B} + \mathbf{E}$$

$$\boldsymbol{\Omega} = \nabla_p \times \mathbf{a} = \frac{\hat{\mathbf{p}}}{2p^2}$$

Chiral Kinetic theory

Son, Yamamoto ('12) Stephanov, Yin ('12)

Chiral kinetic equation (CKE)

$$(\partial_t + \dot{\mathbf{x}} \cdot \nabla_x + \dot{\mathbf{p}} \cdot \nabla_p) f = 0$$

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Current

$$\mathbf{j} = \int_p f \hat{\mathbf{p}} + \mathbf{E} \times \int_p f \boldsymbol{\Omega} + B \int_p f \hat{\mathbf{p}} \cdot \boldsymbol{\Omega}$$

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Hall effect Chiral magnetic effect

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Hall effect Chiral magnetic effect

Anomaly

$$\partial_\mu j^\mu = \frac{1}{4\pi^2} \mathbf{E} \cdot \mathbf{B}$$

Lorentz covariance?

$$S = \int dt (\dot{\mathbf{x}} \cdot \mathbf{p} + \dot{\mathbf{x}} \cdot \mathbf{A} - |\mathbf{p}| - A_0 - \dot{\mathbf{p}} \cdot \mathbf{a})$$

The action looks not Lorentz invariant.

Lorentz covariance?

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The action looks not Lorentz invariant.

1) Energy shift: $|\mathbf{p}| \rightarrow |\mathbf{p}| (1 - \boldsymbol{\Omega} \cdot \mathbf{B})$

Son, Yamamoto ('13)

Lorentz covariance?

$$S = \int dt (\dot{\mathbf{x}} \cdot \mathbf{p} + \dot{\mathbf{x}} \cdot \mathbf{A} - |\mathbf{p}| - A_0 - \dot{\mathbf{p}} \cdot \mathbf{a})$$

The action looks not Lorentz invariant.

1) Energy shift: $|\mathbf{p}| \rightarrow |\mathbf{p}| (1 - \boldsymbol{\Omega} \cdot \mathbf{B})$

Son, Yamamoto ('13)

2) Modified Lorentz transform

Chen, Son, Stephanov, Yee, Yin ('14)

$$\delta \mathbf{x} = \beta t + |\mathbf{p}| \boldsymbol{\beta} \times \boldsymbol{\Omega} \quad \delta \mathbf{p} = \beta \epsilon + |\mathbf{p}| (\boldsymbol{\beta} \times \boldsymbol{\Omega}) \times \mathbf{B}$$

Several approaches

Hamiltonian formalism

Son, Yamamoto ('12)

Semi-classical path integral

Stephanov, Yin ('12)

World line formalism: Mueller, Venugopalan ('17) ('18)

Field theoretical approach

High density effective theory: Son, Yamamoto ('13)

On-shell effective theory: Manuel, Torres-Rincon ('13) ('14)

Carignano, Manuel, Torres-Rincon ('18)

Wigner function: Gao, Liang, Pu, Wang, Wang ('12), Chen, Pu, Wang, Wang ('13)

Wu, Hou, Ren ('17)

Huang, Shi, Jiang, Liao, Zhuang ('18)

CVE: Gao, Pang, Wang ('18)

Kadanoff-Baym: YH, Shi Pu, Yang ('16) ('17), YH, Yang ('18)

QFT approach

Propagator (Wigner function)

$$S^{<}(p, X) = \int d^4s e^{is \cdot p} \langle \psi^\dagger(y) \psi(x) \rangle U(x, y)$$

$$S^{>}(p, X) = \int d^4s e^{is \cdot p} \langle \psi(x) \psi^\dagger(y) \rangle U(x, y)$$

$$\text{where } X = \frac{x+y}{2} \quad s = x - y$$

QFT approach

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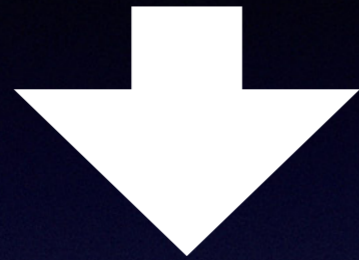
EOM (Schwinger-Dyson equation)



$$\sigma^\mu (p_\mu - A_\mu) \star S^< = \frac{-i\hbar}{2} (\Sigma^< \star S^> - \Sigma^> \star S^<)$$

EOM Up to order \hbar

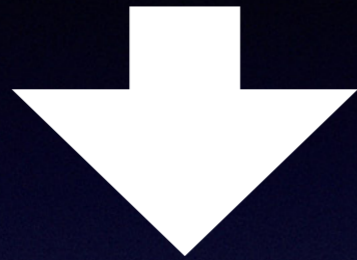
$$\sigma^\mu \left(p_\mu + \frac{i\hbar}{2} \Delta_\mu \right) S^< = \frac{-i\hbar}{2} (\Sigma^< S^> - \Sigma^> S^<)$$



where $\Delta_\mu = \partial_\mu + F_{\nu\mu} \frac{\partial}{\partial p_\nu}$

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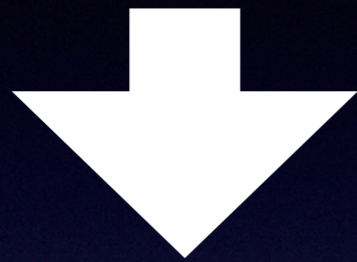
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Chiral kinetic equation (CKE)

$$\Delta_\mu S^<\mu = \Sigma_\mu^< S^>\mu - \Sigma_\mu^> S^<\mu \quad \text{where } S^<\mu = \frac{1}{2} \text{tr} \sigma^\mu S^<$$

EOM Up to order \hbar

$$\sigma^\mu \left(p_\mu + \frac{i\hbar}{2} \Delta_\mu \right) S^< = \frac{-i\hbar}{2} (\Sigma^< S^> - \Sigma^> S^<)$$



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$$S^<\mu = 2\pi\epsilon(p \cdot n) \left[\delta(p^2) (p^\mu + \hbar S_n^{\mu\nu} \mathcal{D}_\nu) + \hbar p_\nu \tilde{F}^{\mu\nu} \delta'(p^2) \right] f$$

YH, Shi Pu, Yang ('16) ('17)

spin: $S_n^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} \frac{p_\alpha n_\beta}{p \cdot n}$ $\mathcal{D}_\mu f = \Delta_\mu f + \Sigma_\mu^< f - \Sigma_\mu^> \bar{f}$

Talk by Amping Huang (Parallel II.3)

Lorentz invariance

Chen, Son, Stephanov ('15), YH. Pu, Yang ('16)

Talk by Jian-Hua Gao (Parallel II.1)

$$S^{<\mu} = 2\pi\epsilon(p \cdot n) \left[\delta(p^2)(p^\mu + \hbar S_n^{\mu\nu} \mathcal{D}_\nu) + \hbar p_\nu \tilde{F}_{\alpha\beta}^{\mu\nu} \delta'(p^2) \right] f$$

$S^{<\mu}$ is Lorentz covariant

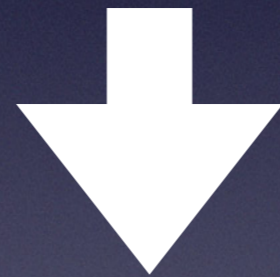
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$S^{<\mu}$ is Lorentz covariant



f is not Lorentz scalar

$$f \rightarrow f + \hbar \frac{\epsilon^{\nu\mu\alpha\beta} p_\alpha n'_\beta n_\mu}{2(p \cdot n)(p \cdot n')} \mathcal{D}_\nu f$$

Application

(local) Equilibrium

$$\text{Current: } J^\mu = 2 \int \frac{d^4 p}{(2\pi)^4} S^{<\mu}(p, X)$$

$$\Rightarrow J = nu + \underbrace{\sigma_B B}_{\text{CME}} + \underbrace{\sigma_\omega \omega}_{\text{CVE}}$$

Dissipative current

CKE with relaxation time approximation

Gorbar, Shovkovy, Vilchinskii, Rudenok, Boyarsky, Ruchayskiy ('16)

Chen, Ishii, Pu, Yamamoto ('16)

YH, Pu, Yang ('17)

$\nabla\mu, \nabla T$ correction

$$\delta J = C_1 \mathbf{E} \times \nabla\mu + C_2 \mathbf{E} \times \nabla T + C_3 \nabla\mu \times \nabla T$$

$$C_i \sim \tau_R$$

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$\odot \mathbf{E}$

low-T

high-T

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Dissipative current

YH, Yang ('18)

Shear and bulk correction

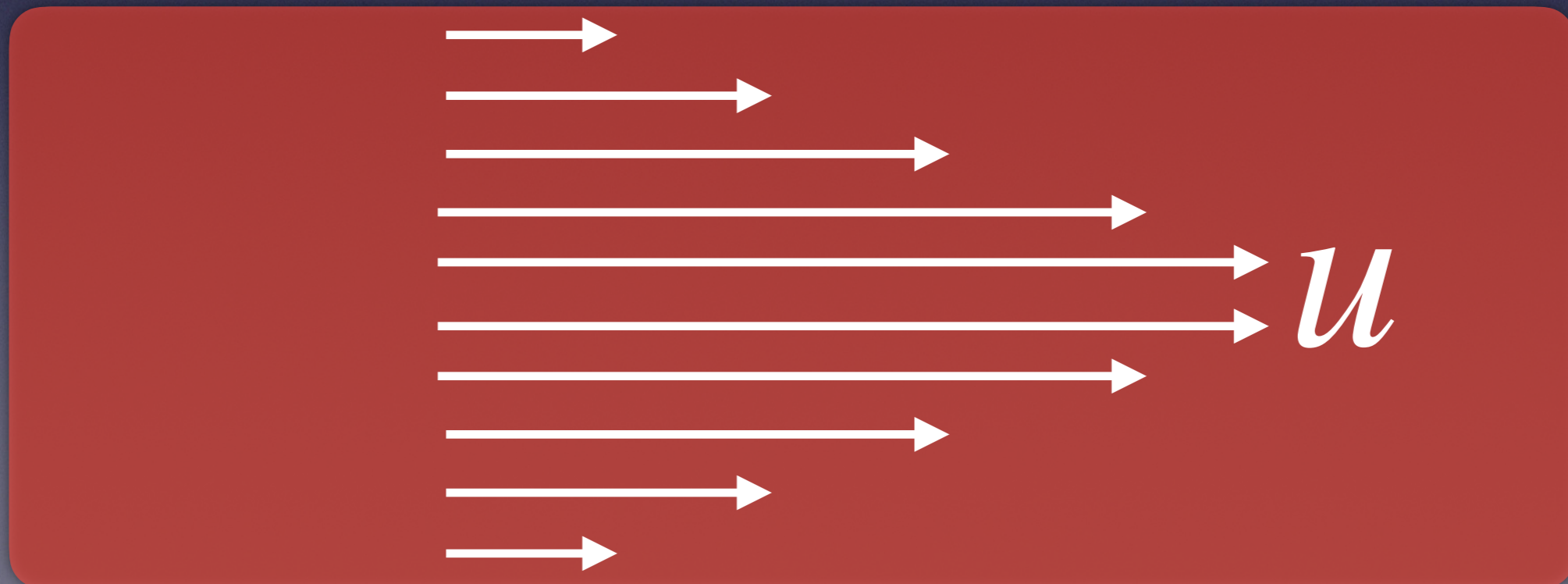
$$\begin{aligned}\delta J^i &= C_4 \pi^{ij} B_j + C_5 \pi^{ij} \omega_j \\ &\quad + C_6 (\nabla \cdot \mathbf{u}) B^i + C_7 (\nabla \cdot \mathbf{u}) \omega^i\end{aligned}$$

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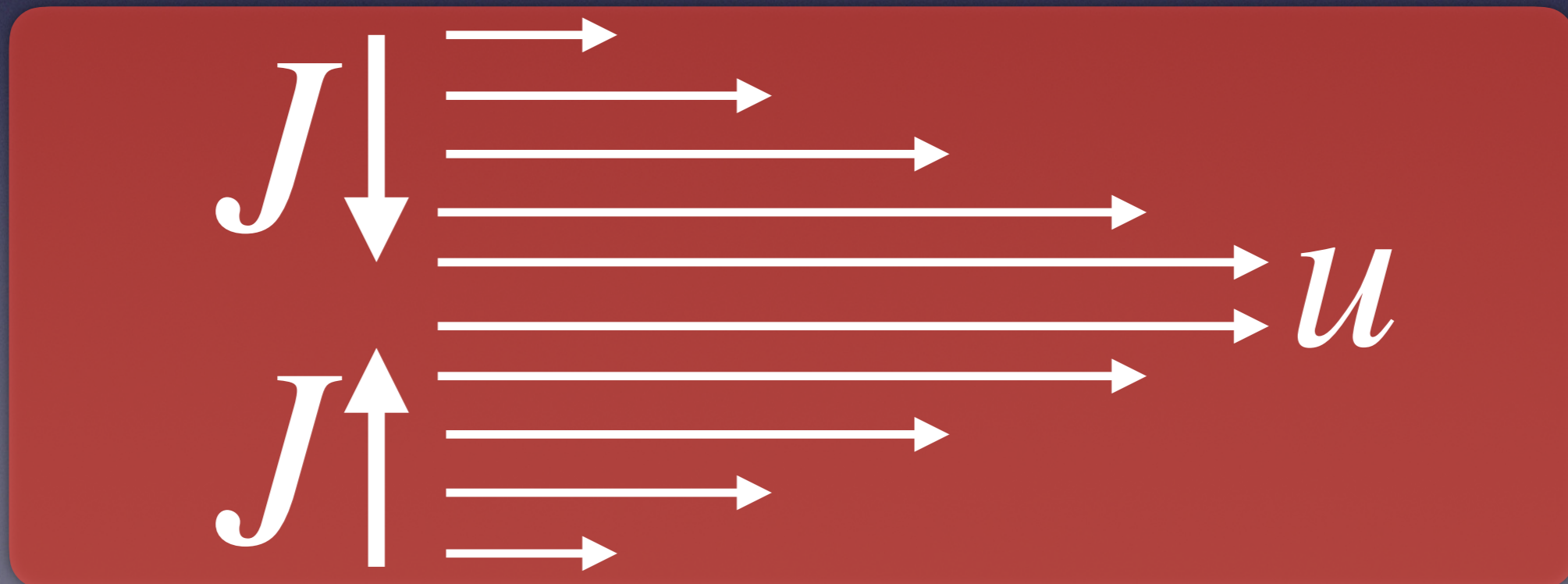


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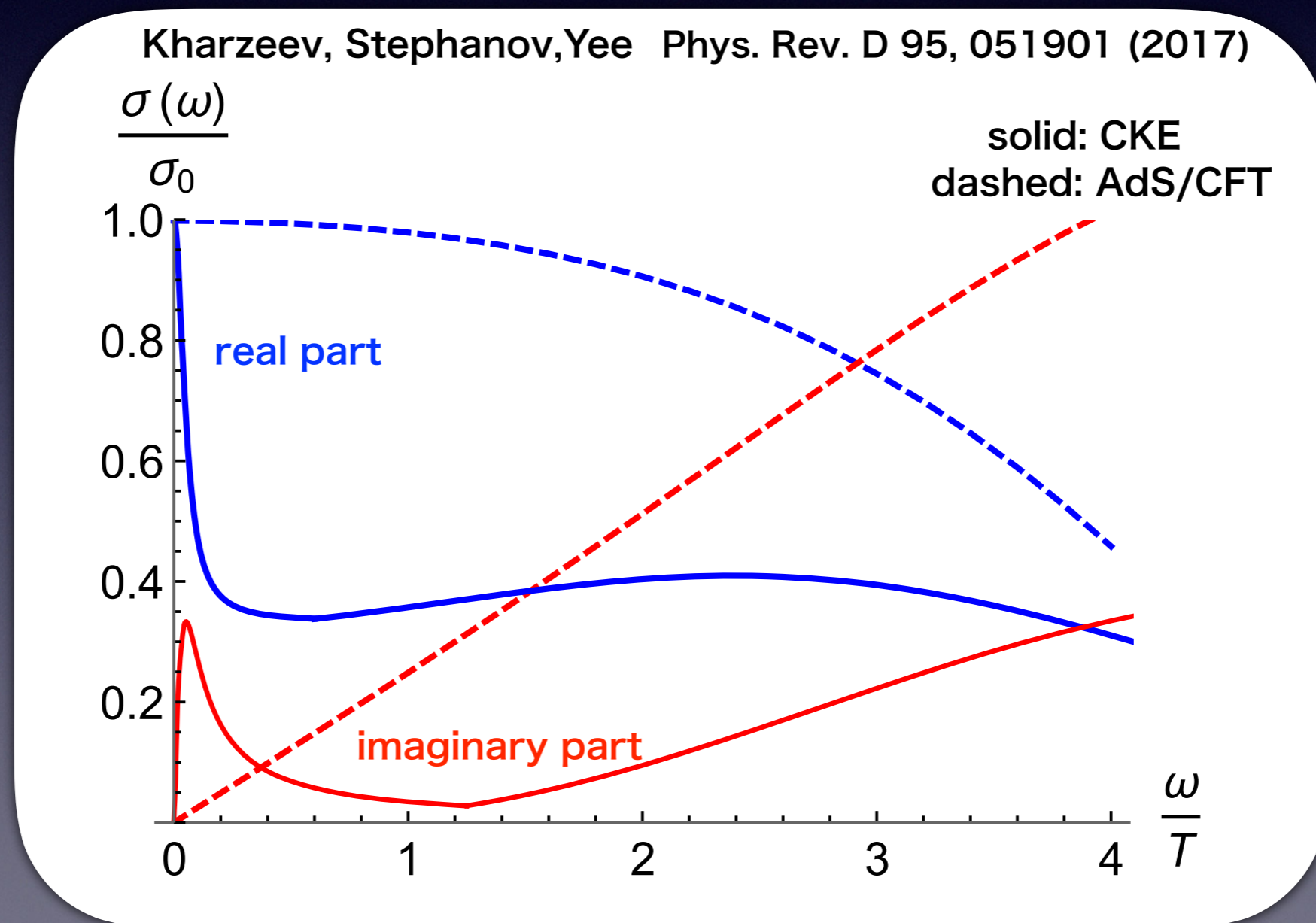


Out of equilibrium transport

Kharzeev, Stephanov, Yee ('17)

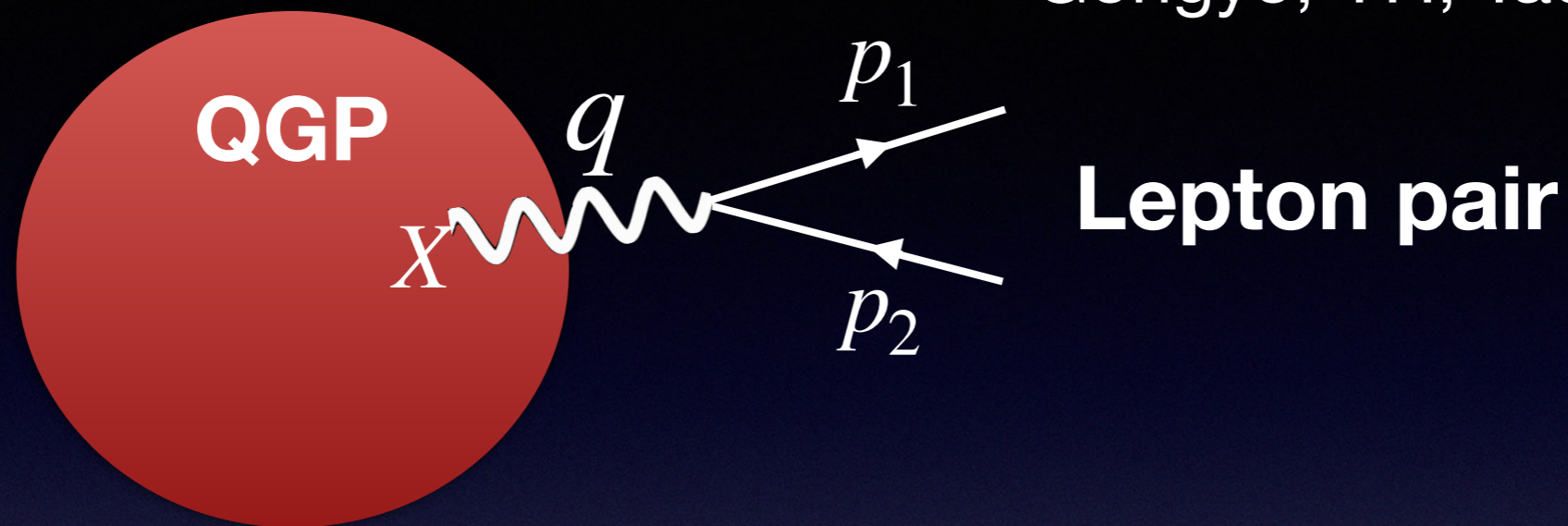
CKE with relaxation time approximation

ω dependence of CME $\sigma(\omega) = \sigma_0 \left(1 - \frac{2}{3} \frac{\omega}{\omega + i\tau_R^{-1}} \right)$



Dilepton production

Gongyo, YH, Tachibana ('18)



Photon polarization function

$$\Pi^{<\mu\nu}(X, q) = \int d^4s e^{iq \cdot s} \langle j^\nu(X - s/2) j^\mu(X + s/2) \rangle$$

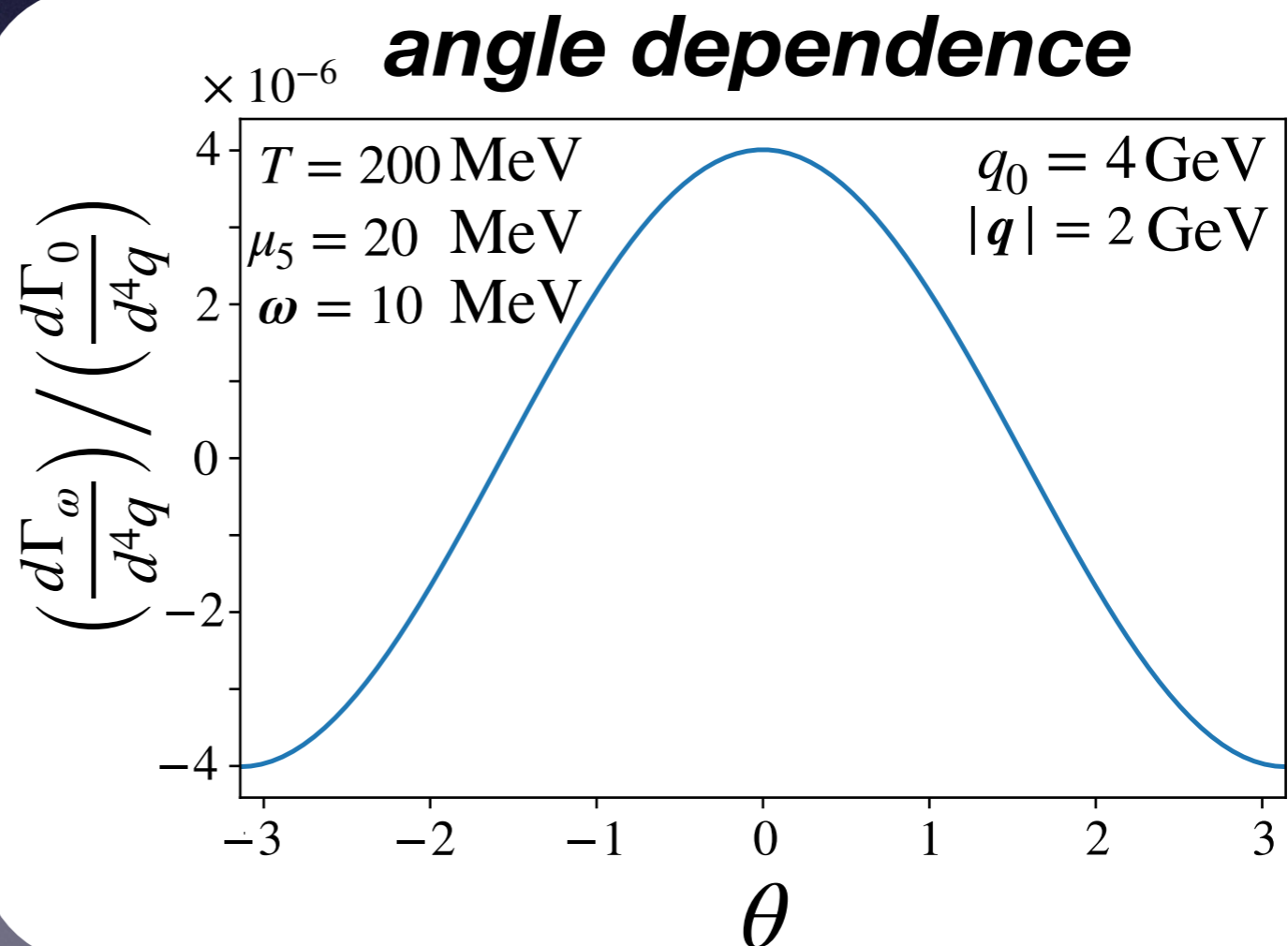
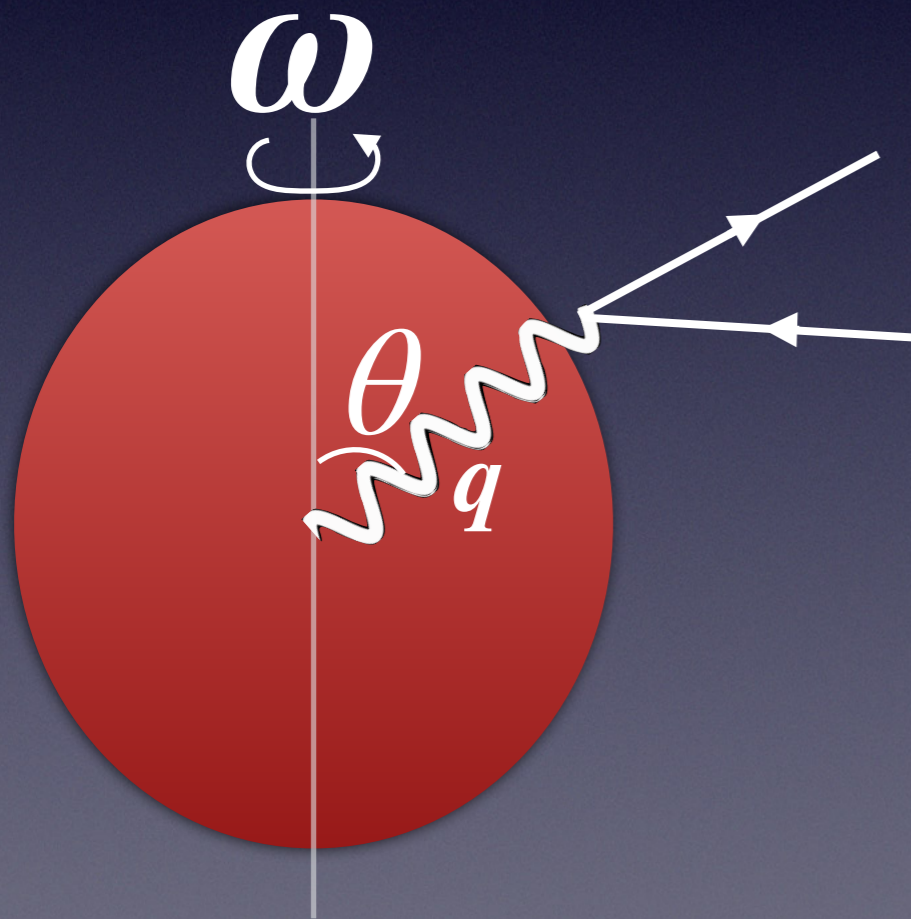
Dilepton production rate

$$\frac{d\Gamma}{d^4q} = - \frac{\alpha}{24\pi^4} \Pi^{<\mu}_{\mu}(q, X)$$

Di-lepton production

Gongyo, YH, Tachibana('18)

$$\frac{d\Gamma}{d^4q} = \frac{d\Gamma_0}{d^4q} + \frac{d\Gamma_\omega}{d^4q} \quad \text{with} \quad \frac{d\Gamma_\omega}{d^4q} = (\boldsymbol{\Omega}_\gamma \cdot \boldsymbol{\omega}) C(q)$$
$$\boldsymbol{\Omega}_\gamma = \frac{\hat{\boldsymbol{q}}}{|q|^2}$$



Puzzle?

CKE from on-shell effective theory

Carignano, Manuel, Torres-Rincon ('18)

$$\left(\Delta_0 + (1 + \mathbf{B} \cdot \boldsymbol{\Omega}) \hat{q}^i \Delta_i + \left(\frac{1}{2} \epsilon^{ijk} E^j \Omega^k - \frac{1}{4} B_{\perp}^i \right) \Delta_i \right) f = 0$$

reproduces consistent anomaly $\partial_{\mu} j_5^{\mu} = \frac{1}{3} \frac{1}{2\pi^2} E \cdot B$

CKE by Son, Yamamoto ('12), YH, Pu, Yang ('17)

$$\left(\Delta_0 + (1 + \mathbf{B} \cdot \boldsymbol{\Omega}) \hat{q}^i \Delta_i + \epsilon^{ijk} E^j \Omega^k \Delta_i \right) f = 0 \quad \text{where } \Delta_{\mu} = \partial_{\mu} + F_{\nu\mu} \frac{\partial}{\partial p_{\nu}}$$

reproduces covariant anomaly

$$\partial_{\mu} j_5^{\mu} = \frac{1}{2\pi^2} E \cdot B$$

Summary

Chiral kinetic theory:

Effective theory

reproducing chiral anomaly

Novel dissipative anomalous transports are found.

Application to HIC and cond-mat

Quarks have mass. What is mass correction to CKE?

Mass correction to CVE cf. Flachi, Fukushima ('17), Lin, Yang ('18)

Talk by Lixin Yang (Parallel II.1)

$$j_5^\mu = \left(\frac{T^2}{6} - \frac{m^2}{4\pi^2} \right) \omega^\mu$$

Analysis with collisions without relaxation time approximation