Direct measurement of the branching fractions for  $K^0$  and  $K^{*0}$  production in  $J/\psi$  decays

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## Charged Particle Identification

### **Combined particle confidence level**

$$CL(i) = \int_{\chi^2}^{\infty} P(z,n) dz$$

Where i = 1, 2, 3, 4 and 5 denotes e,  $\mu$ ,  $\pi$ , K and proton.

$$\chi^2 = \Delta_{TOF}^2 + \Delta_{dE/dx}^2 + \Delta_{measure}^2$$

$$\Delta_{TOF} = \frac{T_{measure} - T_{expect}}{\sigma_{TOF}}$$

$$\Delta_{dE/dx} = \frac{dE/dx_{measure} - dE/dx_{expect}}{\sigma_{dE/dx}}$$

$$\Delta_{BSC} = \frac{E_{measure} - E_{expect}}{\sigma_E}$$

$$P(z,n) = \frac{1}{2^{\frac{n}{2}}} (z)^{\frac{n}{2}-1} e^{-\frac{n}{2}}$$

For  $\pi$ , CL( $\pi$ )>0.1%, and CL( $\pi$ )>CL(K)

For K, CL(K)>0.1%, and CL(K)>CL(
$$\pi$$
)

For proton, 
$$\frac{CL(p)}{CL(\pi)+CL(K)+CL(p)} > 0.6$$

## Data Sample (April 2001, February 2003, December 2003, January 2004)

### **Integrated luminosity**

Processes used :  $e^+e^- \rightarrow e^+e^-$ ,  $e^+e^- \rightarrow J/\psi \rightarrow e^+e^-$ 

$$\mathcal{L} = \frac{N_{e^+e^-}^{obs}}{\sigma_{Bhabha} \cdot \varepsilon_{Bhabha} + \sigma_{J/\psi \to e^+e^-} \cdot \varepsilon_{J/\psi \to e^+e^-} + \sigma_{int} \cdot \varepsilon_{int}}$$

$$\sigma_{had}^{obs} = \sigma_{J/\psi \to e^+e^-} \cdot \varepsilon_{J/\psi \to e^+e^-} + \sigma_{int} \cdot \varepsilon_{int} = \frac{N_{had}^{obs}}{\epsilon_{had} \cdot \mathcal{L}}$$

Integrated luminosity and the observed inclusive hadronic cross section correlate with each other

Through a  $\chi^2$  fit to the observed inclusive cross sections with J/ $\psi$  resonance parameters fixed at PDG value, we can get  $\mathcal{L}$  at each energy point.

Divide all energy points into 16 energy regions (Add  $\mathcal{L}$  with their weights)

### Selection of the candidate events

At least two charged tracks should be well reconstructed  $\theta$  of each charged track must satisfy  $|cos\theta| < 0.80$ 

For  $e^+e^- \rightarrow K_S^0 + X$ , each track from the region  $|cos\theta| < 0.80$ (for ours, 0.93)  $\sqrt{V_x^2 + V_y^2} < 8.0 cm(for ours, 12.0 cm)$  $|V_{\tau}| < 22.0 cm(for ours, 20.0 cm)$ The best fit value of laboratory decay distance in xy plane for  $K_S^0$  is required to be displaced at least 0.004mm from the interaction point. For  $e^+e^- \rightarrow K^{*0} + X$ , each track from the region  $|cos\theta| < 0.80$  $\sqrt{V_x^2 + V_y^2} < 2.0cm$  $|V_{z}| < 20.0 cm$ 

#### Fit to the mass spectra

Reconstruct  $K_S^0$  and  $K^{*0}$  by examining the invariant mass spectra of  $\pi^+\pi^-$  and  $K^-\pi^+$ 

For  $K_S^0$  (the same as ours)

If more than one combination satisfies the above selection criteria in one event, only those with the longest decay are retained.

For *K*\*0

If more than one combination satisfies the above selection criteria in one event, only those with largest  $CL_{K} * CL_{\pi}$  is retained.

Table 3: Fitted results of mass and the resolution of the selected inclusive  $\pi^+\pi^-$ ,  $K^-\pi^+$ , and  $K^+K^-$  events for M.C. samples.

|                   | $K_S^0$         | $K^{\star 0}$   |
|-------------------|-----------------|-----------------|
| Fitted Mass [MeV] | $497.13\pm0.03$ | $889.85\pm0.14$ |
| Resolution [MeV]  | $6.54 \pm 0.29$ | $7.91 \pm 0.38$ |

function for the signals and a third order polynomial for the background, we obtain the fitted masses and the resolutions, which are shown in Table 3.



Figure 1: The distributions for the inclusive (a)  $\pi^+\pi^-,$  and (b)  $K^-\pi^+$  events for M.C. samples.

The invariant masses are denoted by  $M_{\pi^+\pi^-}$  and  $M_{K^-\pi^+}$ , respectively, and the invariant mass spectra is fitted with a Breit-Wigner convoluted with a Gaussian resolution function for the signals and a polynomial function for background. (After this work, we compared the dynamic variables of MC and DATA as well as detection efficiency)

### **Background contamination**

Possible contaminations :  $e^+e^- \rightarrow (\gamma)e^+e^-$ ,  $e^+e^- \rightarrow (\gamma)\mu^+\mu^-$ 

Accounted for the luminosity, we can estimate the number of background events at each point are less than 10<sup>-3</sup>, thus we can neglect the influence of these background processes.

#### **Monte Carlo simulation**

At this step, the article estimates the detection efficiencies of these two processes.( $\gamma^*$  samples, we use 3.050, 3.060, 3.080, 3.090, 3.095, 3.099, 3.1015, 3.112, 3.120 to estimate other points)



Figure 4: The distributions of the energy dependent efficiencies for (a)  $e^+e^- \rightarrow K_S^0 + X$  and (b)  $e^+e^- \rightarrow K^{\star 0} + X$  processes, where the lines indicate the expected efficiencies.

Table 5: The energy dependent detection efficiencies for  $e^+e^- \rightarrow K_S^0 + X$  and  $e^+e^- \rightarrow K^{*0} + X$  processes, including the branching fractions for the decays  $K^0 \rightarrow K_S^0 \rightarrow \pi^+\pi^-$  and  $K^{*0} \rightarrow K^+\pi^-$ , where the errors are statistical error.

| $E_{\rm cm}   [{\rm GeV}]$ | $\epsilon_{K^0_S}$  | $\epsilon_{K^{\star 0}}$ |
|----------------------------|---------------------|--------------------------|
| 3.0814                     | $0.0931 \pm 0.0025$ | $0.1473 \pm 0.0016$      |
| 3.0878                     | $0.0949 \pm 0.0026$ | $0.1486 \pm 0.0016$      |
| 3.0919                     | $0.0920 \pm 0.0024$ | $0.1456 \pm 0.0016$      |
| 3.0934                     | $0.0967 \pm 0.0022$ | $0.1472 \pm 0.0016$      |
| 3.0949                     | $0.0959 \pm 0.0020$ | $0.1437 \pm 0.0016$      |
| 3.0961                     | $0.1008 \pm 0.0018$ | $0.1490 \pm 0.0016$      |
| 3.0967                     | $0.0933 \pm 0.0016$ | $0.1477 \pm 0.0016$      |
| 3.0973                     | $0.0976 \pm 0.0015$ | $0.1478 \pm 0.0016$      |
| 3.0977                     | $0.0962 \pm 0.0014$ | $0.1477 \pm 0.0016$      |
| 3.0985                     | $0.0954 \pm 0.0014$ | $0.1497 \pm 0.0016$      |
| 3.0995                     | $0.0987 \pm 0.0014$ | $0.1559 \pm 0.0016$      |
| 3.1005                     | $0.0965 \pm 0.0014$ | $0.1490 \pm 0.0016$      |
| 3.1027                     | $0.0965 \pm 0.0014$ | $0.1502 \pm 0.0016$      |
| 3.1061                     | $0.0935 \pm 0.0016$ | $0.1471 \pm 0.0016$      |
| 3.1105                     | $0.0979 \pm 0.0018$ | $0.1485 \pm 0.0016$      |
| 3.1171                     | $0.0951 \pm 0.0019$ | $0.1520 \pm 0.0016$      |

#### **Observed cross section**

The observed cross section for  $e^+e^- \rightarrow M + X$ 

$$\sigma_{e^+e^- \to M+X}^{obs} = \frac{N_{e^+e^- \to M+X}^{obs}}{\mathcal{L} * \mathcal{E}_{e^+e^- \to M+X}}$$



Table 6: The observed cross sections for  $e^+e^- \to K^0 + X$  and  $e^+e^- \to K^{*0} + X$ processes at each combined energy point, where the errors are the combined energy dependent and data sets dependent errors.

| $E_{\rm cm}~[{\rm GeV}]$ | $\sigma_{K^0}^{\text{obs}}$ [nb] | $\sigma_{K^{\star 0}}^{\mathrm{obs}}$ [nb] |
|--------------------------|----------------------------------|--|
| 3.0814                   | $4.11 \pm 1.16$                  | $2.00 \pm 1.63$                            |
| 3.0878                   | $1.73 \pm 1.31$                  | $0.33 \pm 2.01$                            |
| 3.0919                   | $4.93 \pm 1.23$                  | $2.57 \pm 1.57$                            |
| 3.0934                   | $6.11 \pm 1.70$                  | $0.15 \pm 1.75$                            |
| 3.0949                   | $59.78 \pm 4.77$                 | $26.93 \pm 5.47$                           |
| 3.0961                   | $423.43\pm21.45$                 | $153.33\pm18.68$                           |
| 3.0967                   | $735.33\pm45.30$                 | $294.02\pm30.74$                           |
| 3.0973                   | $654.74\pm35.06$                 | $236.50\pm28.78$                           |
| 3.0977                   | $510.01\pm32.44$                 | $214.63\pm23.22$                           |
| 3.0985                   | $246.36\pm15.15$                 | $89.50 \pm 12.55$                          |
| 3.0995                   | $84.76 \pm 6.81$                 | $21.18 \pm 6.77$                           |
| 3.1005                   | $59.36 \pm 6.94$                 | $16.88 \pm 7.89$                           |
| 3.1027                   | $33.95 \pm 4.54$                 | $9.28 \pm 5.28$                            |
| 3.1061                   | $17.54 \pm 3.23$                 | $15.84 \pm 4.84$                           |
| 3.1105                   | $19.54 \pm 3.16$                 | $7.02 \pm 3.70$                            |
| 3.1171                   | $13.36 \pm 2.45$                 | $4.66 \pm 3.45$                            |



## Measurement of $\mathcal{B}(J/\psi \rightarrow \mathcal{M} + X)$

#### **Expected cross section**

The expected observed cross section  $\sigma_{e^+e^- \to M+X}^{expect}(E_{cm})$  consists pf two components:  $\sigma_{e^+e^- \to M+X}^{expect}(E_{cm}) = \sigma_{J/\psi \to M+X}^{expect}(E_{cm}) + \sigma_{e^+e^- \to M+X}^{continuum}(E_{cm})$ 

$$\begin{split} \sigma_{J/\psi \to \mathcal{M}+X}^{\text{expect}}(s) &= \int_{0}^{\infty} ds' G(s,s') \cdot \int_{0}^{1} dx \cdot \sigma_{J/\psi \to \mathcal{M}+X}^{\text{B}}(s'(1-x)) F(x,s') & \beta = \frac{2\alpha}{\pi} (\ln \frac{s}{m_{e}^{2}} - 1), \\ G(s,s') &= \frac{1}{\sqrt{2\pi\xi}} e^{-\frac{(\sqrt{s} - \sqrt{s})^{2}}{2\xi^{2}}} \sigma_{J/\psi \to \mathcal{M}+X}^{\text{expect}}(s) &= \int_{0}^{\infty} ds' G(s,s') \cdot \sigma_{J/\psi \to \mathcal{M}+X}^{\text{sim}}(s') & (1+\delta) = 1 + \frac{\alpha}{\pi} (\frac{\pi^{2}}{3} - \frac{1}{2}) + \frac{3}{4} \beta - \frac{\beta^{2}}{24} (\frac{1}{3} \ln \frac{s}{m_{e}^{2}} + 2\pi^{2} - \frac{37}{4}) \\ \sigma_{J/\psi \to \mathcal{M}+X}^{\text{sim}}(s) &= \frac{12\pi \Gamma^{\text{ee}} \Gamma^{\text{tot}} \cdot \mathcal{B}(J/\psi \to \mathcal{M}+X)}{s^{2}}, \\ [(1+\delta)A^{\beta-2}\phi(\cos\theta,\beta) + (-\beta + \frac{3}{4}\beta^{2})A^{\beta-1}\phi(\cos\theta,\beta + 1)], & cos\theta = \frac{1}{4} (\frac{M^{2}}{s} - 1), \\ \phi(\cos\theta,\beta) &= \frac{\pi\beta \sin(\theta(1-\beta))}{\sin\theta \sin(\pi\beta)}. \end{split}$$

 $\sigma_{e^+e^- \to \mathcal{M}+X}^{\text{continuum}}(s) = h(s) \cdot \sigma_{\mu^+\mu^-}^{\text{B}}(s)$ where  $\sigma_{\mu^+\mu^-}^{\text{B}}(s)$  is the Born cross section for  $e^+e^- \to \mu^+\mu^-$  and h(s) is a factor determined from fit.

# Measurement of $\mathcal{B}(J/\psi \rightarrow \mathcal{M} + X)$

#### **Fitted branching fractions**



 $E_{\rm cm}[GeV]$ 





 $E_{\rm cm}[GeV]$ 

Figure 7: The observed cross sections for  $e^+e^- \to K^{\star 0} + X$  versus the nominal c.m. energies.

Table 7: The measured values of the mass of  $J/\psi$  resonance, the energy spread of BEPC machine, the branching fraction for  $J/\psi \to \mathcal{M} + X$  and h, which are obtained by fitting the observed cross sections discribed in text.

|                             | $\chi^2/\text{NDF}$ | M                  | $\Delta^{\text{BEPC}}$ | h             | $\mathcal{B}(J/\psi \rightarrow M + X)$ |
|-----------------------------|---------------------|--------------------|------------------------|---------------|---|
|                             |                     | [MeV]              | [MeV]                  |               | [%]                                     |
| $\mathcal{M} = K^0$         | 15.9/NDF            | $3096.90 \pm 0.03$ | $0.90\pm0.02$          | $0.18\pm0.03$ | $20.4\pm0.5$                            |
| $\mathcal{M} = K^{\star 0}$ | 10.3/NDF            | $3096.90 \pm 0.06$ | $0.88\pm0.05$          | $0.06\pm0.04$ | $7.7\pm0.4$                             |

# Systematic Error Analysis

To estimate the systematic uncertainties of the branching fractions, the article shifts the measured observed cross section,  $\Gamma^{tot}$  and  $\Gamma^{ee}$  of J/ $\psi$  resonance by +/-1 $\sigma$  to measure the change of the branching fractions.

Table 9: The sources of the systematic uncertainties in the measured branching fractions.

| Sources                     | Uncertanty [%] in                 | Uncertanty [%] in                         |    |
|-----------------------------|-----------------------------------|---|----|
|                             | $\mathcal{B}(J/\psi \to K^0 + X)$ | $\mathcal{B}(J/\psi \to K^{\star 0} + X)$ |    |
| $\Delta \sigma^{\rm obs}$   | 5.5                               | 3.8                                       |    |
| $\Gamma^{ee}_{J/\psi}$      | 3.6                               | 3.6                                       |    |
| $\Gamma_{J/\psi}^{\rm tot}$ | 0.0                               | 0.0                                       | _3 |
| Total                       | 6.6                               | 5.3                                       |    |
|                             |                                   |   |    |
|                             |                                   |   |    |

Table 10: The branching fractions obtained from the fit, where the first errors are from the fit and the second are the systematic.

|                        |                              | ${\mathcal B}\ [\%]$       |           |
|------------------------|------------------------------|----------------------------|-----------|
|                        | $J/\psi \to K^0 + X$         | $20.4\pm0.5\pm1.3$         |           |
|                        | $J/\psi \to K^{\star 0} + X$ | $7.7\pm0.4\pm0.4$          |           |
|                        |                              |                            |           |
| $err_{sys}^{tanle1}$   | $e^{0} = err_{sys}^{table}$  | $^{29}(\% \rightarrow dec$ | imal) * B |
| Eg: 0.066 <sup>.</sup> | *20.4≈1.3                    |                            |           |
|                        |                              |                            |           |

Calculating the relative changes of +σ and -1σ,
respectively and choosing the greater one as the systematic error.