

Limiting Top-Higgs Interaction and Higgs Width from multi-Top Productions

Qing-Hong Cao

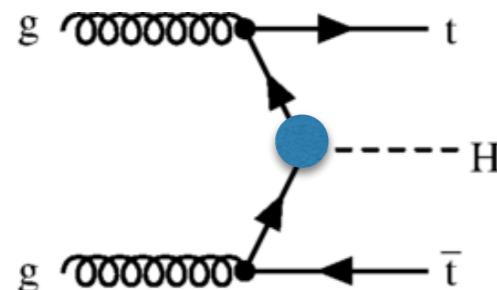
Institute of Theoretical Physics, School of Physics, PKU
Center of High Energy Physics, PKU



北京大学



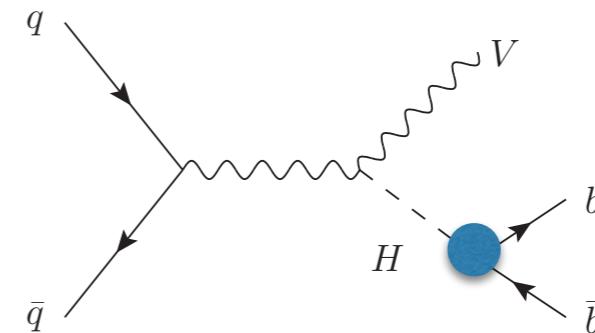
Observation of Higgs-Top and Higgs-Bottom couplings



$$\mu_{ttH} = 0.96^{+0.34}_{-0.31}$$

CMS-PAS-HIG-18-019

Size, CP-phase



$$\mu_{VH} = 1.01 \pm 0.12^{+0.16}_{-0.15}$$

ATLAS, 1808.08238

Size, CP-phase

Higgs-boson width

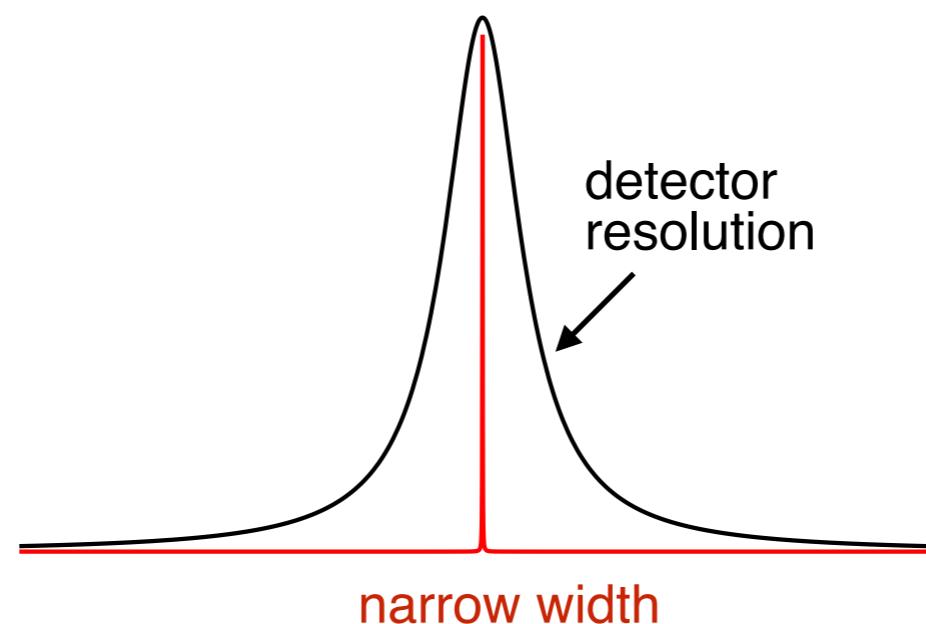
$$\begin{cases} \Gamma_H > \Gamma_H^{\text{SM}} \\ \Gamma_H < \Gamma_H^{\text{SM}} \end{cases}$$

Invisible decay, scalar extension,

Scalar mixing, composite Higgs ...

$$\Gamma_H^{\text{SM}} = 4 \text{ MeV}$$

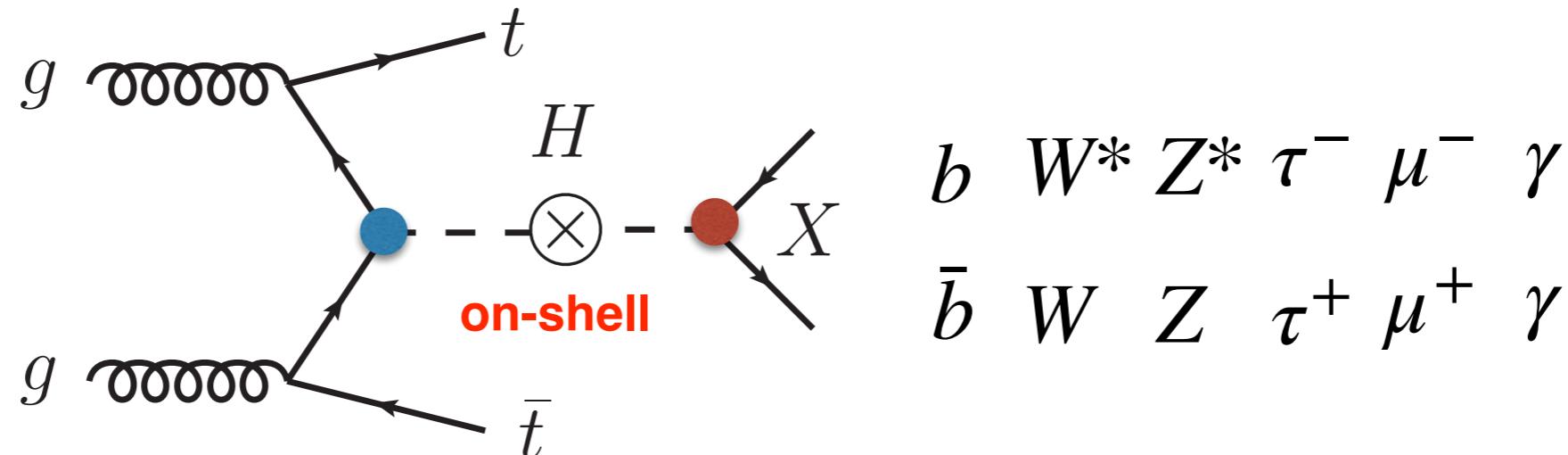
$$\frac{\Gamma_H^{\text{SM}}}{m_H} = 0.000032$$



detector
resolution

One can only obtain
a range but not precise
value of Higgs with

Top-quark Yukawa coupling



Narrow width approximation

$$\hat{\sigma} = \sigma(gg \rightarrow t\bar{t}H) \times \frac{\Gamma(H \rightarrow X)}{\Gamma_H} = \frac{\kappa_t^2 \kappa_X^2}{\Gamma_H / \Gamma_H^{\text{SM}}} \hat{\sigma}_{\text{SM}}$$

$$\kappa_t = \frac{y_t}{y_t^{\text{SM}}}$$

$$\kappa_X = \frac{y_X}{y_X^{\text{SM}}}$$

$$\text{signal strength} \quad \mu \equiv \frac{\hat{\sigma}}{\hat{\sigma}_{\text{SM}}} = \frac{\kappa_t^2 \kappa_X^2}{\Gamma_H / \Gamma_H^{\text{SM}}} \xrightarrow[\kappa_X = 1]{\Gamma_H = \Gamma_H^{\text{SM}}} \mu = \kappa_t^2$$

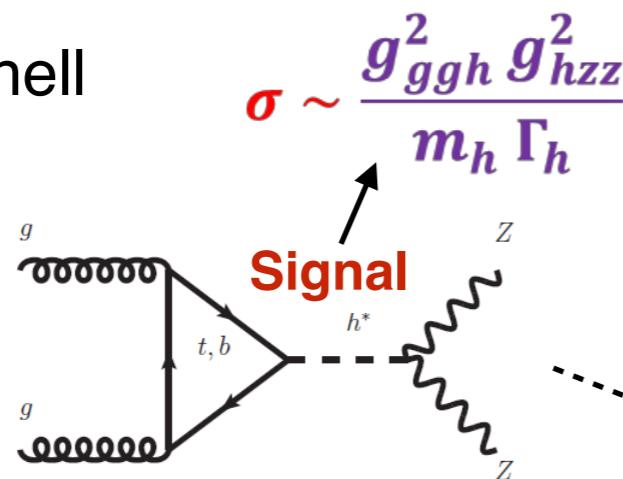
Higgs width affects all the Higgs measurements.

Q1: can we determine κ_t without those assumptions?

Higgs boson width measurement

Caola, Melnikov (2013)

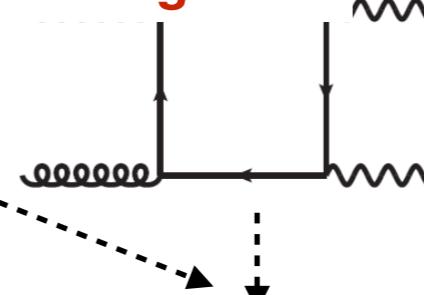
Higgs on-shell production



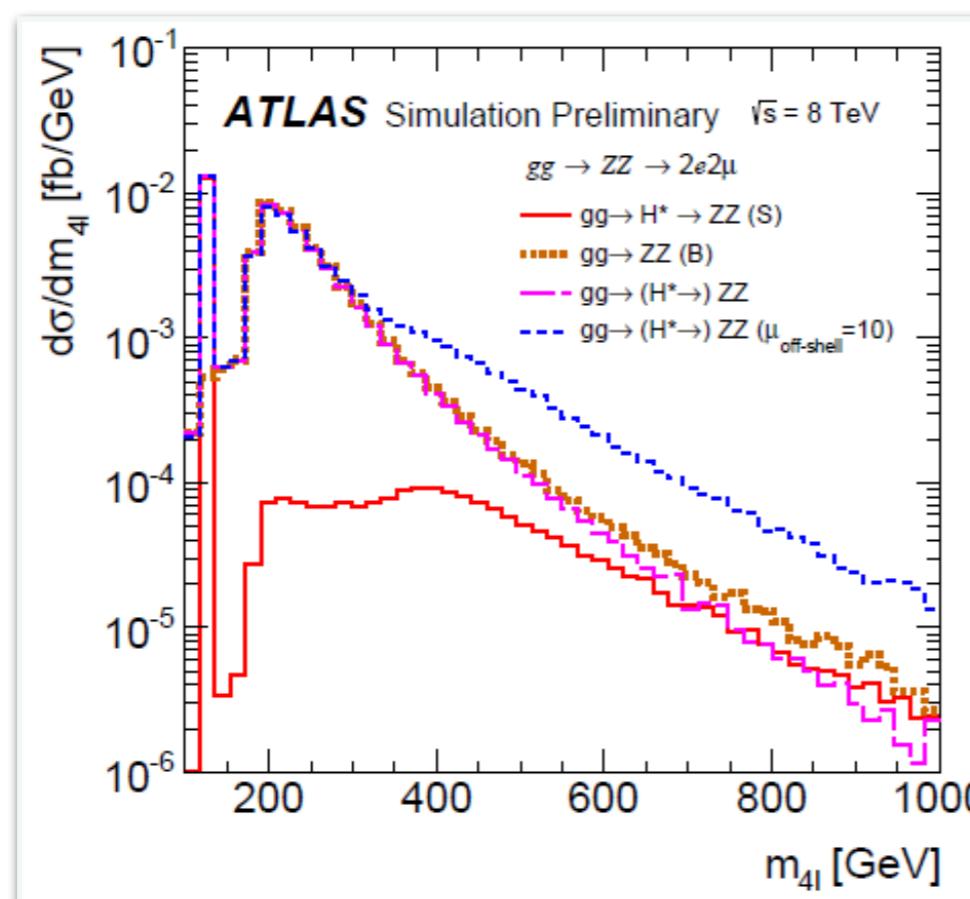
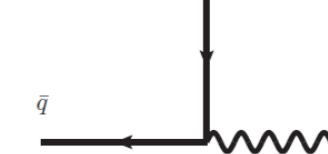
Higgs off-shell production

$$\frac{d\sigma}{dM_{ZZ}^2} \sim \frac{g_{ggh}^2 g_{hzz}^2}{(M_{ZZ}^2 - m_h^2)^2} |M_1|^2 + \frac{g_{ggh} g_{hzz}}{(M_{ZZ}^2 - m_h^2) M_{ZZ}^2} |M_1 M_2^*|$$

Background



Dominant background



On-shell versus off-shell measurements of signal strength



a significant sensitivity to Higgs width

Recent measurements of Higgs width

Direct measurements:

CMS, 1706.09936
(13TeV, 35.9fb⁻¹)

$$gg \rightarrow H \rightarrow ZZ^* \rightarrow 4\ell$$

$$\Gamma_H < 1.10 \text{ GeV}$$

CMS, 1605.02329
(7+8TeV)

$$gg \rightarrow H \rightarrow WW^* \rightarrow 2\ell + \text{MET}$$

$$\Gamma_H < 13 \text{ MeV}$$

ATLAS, 1808.01191
(13TeV, 36.1fb⁻¹)

$$\Gamma_H < 14.4 \text{ MeV}$$

Indirect measurement through invisible decay:

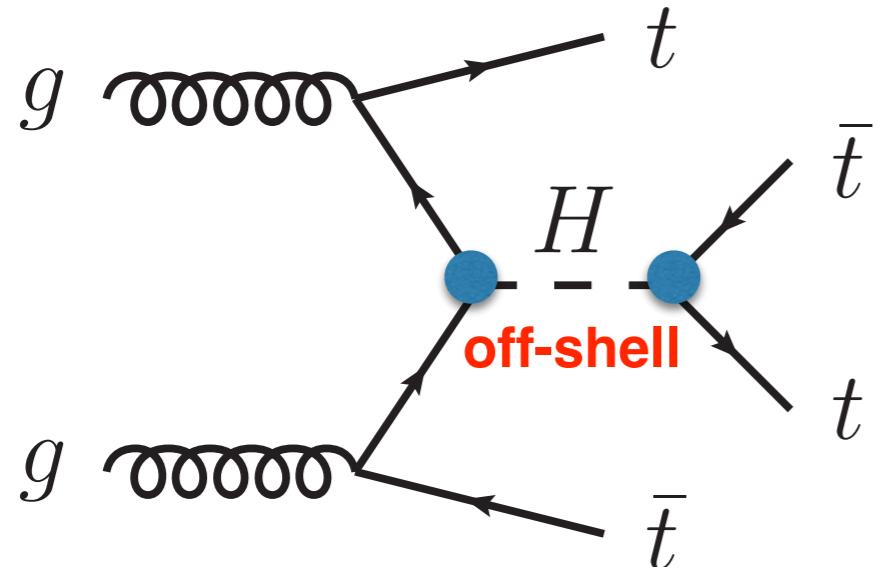
ATLAS, 1809.06682 (13TeV, 36.1fb⁻¹)

$$\text{Br}(H \rightarrow \text{inv}) \leq 0.37 \equiv \text{Br}_{\text{inv}}^{\text{Max}} \xrightarrow{\Gamma_H = \Gamma_H^{\text{known}} + \Gamma_H^{\text{inv}}} \Gamma_H \leq \frac{\Gamma_H^{\text{Known}}}{1 - \text{Br}_{\text{inv}}^{\text{Max}}} \downarrow \Gamma_H^{\text{Known}} = \Gamma_H^{\text{SM}}$$
$$\Gamma_H \leq \frac{\Gamma_H^{\text{SM}}}{1 - \text{Br}_{\text{inv}}^{\text{Max}}} = 1.6\Gamma_H^{\text{SM}} = 6.4 \text{ MeV}$$

Q2: alternative way to measure Higgs width?

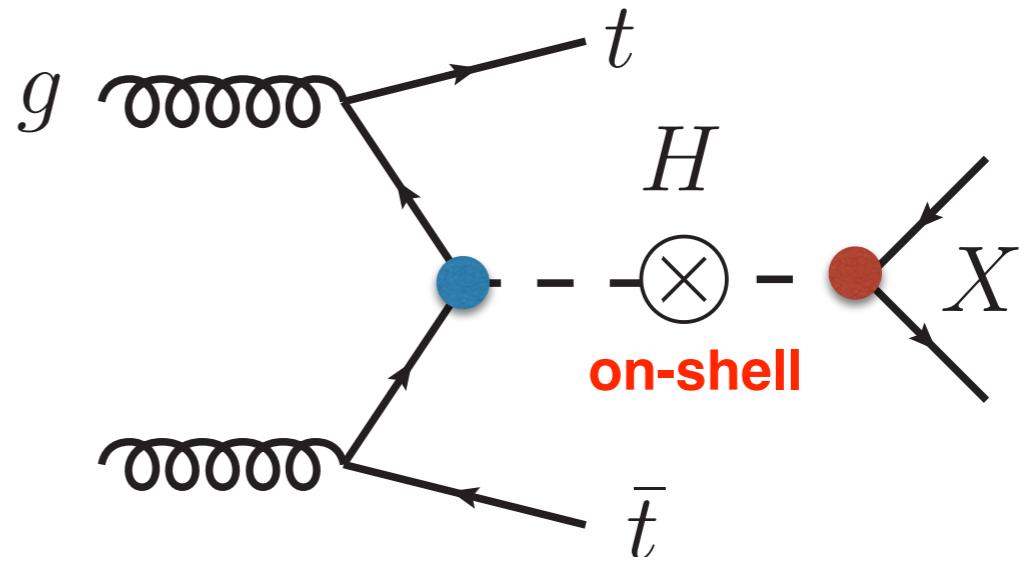
Four top-quark production

QHC, Chen, Liu, 1602.01934



$$\hat{\sigma}(pp \rightarrow t\bar{t}H \rightarrow t\bar{t}t\bar{t})$$

$$\propto \kappa_t^4 \times \hat{\sigma}(pp \rightarrow t\bar{t}H \rightarrow t\bar{t}t\bar{t})_{\text{SM}}$$

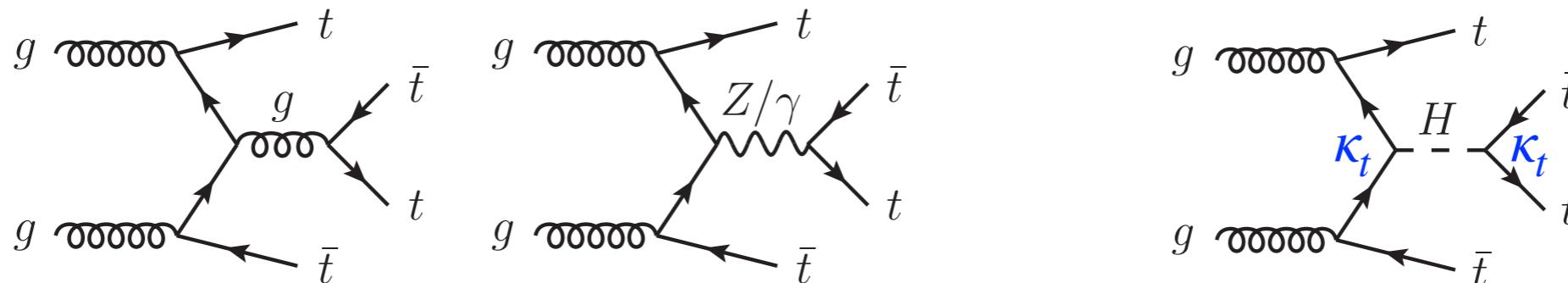


$$\hat{\sigma}(pp \rightarrow t\bar{t}H \rightarrow t\bar{t}X)$$

$$\propto \underbrace{\frac{\kappa_t^2 \kappa_X^2}{\Gamma_H/\Gamma_H^{\text{SM}}}}_{\mu_{t\bar{t}H}^X} \times \hat{\sigma}(pp \rightarrow t\bar{t}H \rightarrow t\bar{t}X)_{\text{SM}}$$

$$\mu_{t\bar{t}H}^X \equiv \frac{\kappa_t^2 \kappa_X^2}{\Gamma_H/\Gamma_H^{\text{SM}}} = \frac{\kappa_t^2 \kappa_X^2}{R_\Gamma} \xrightarrow{\kappa_X \simeq 1} \frac{\kappa_t^2}{R_\Gamma} = \mu_{t\bar{t}H}^X$$

Measuring κ_t from four top-quark production



$$\sigma(t\bar{t}t\bar{t}) = \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g/Z/\gamma} + \kappa_t^2 \sigma_{\text{int}}^{\text{SM}} + \kappa_t^4 \sigma^{\text{SM}}(t\bar{t}t\bar{t})_H$$

LO	8 TeV	1.344	-0.224	0.171	in unit of fb
	13 TeV	9.997	-1.547	1.108	

NLO corrections:

Bevilacqua, Worek (2012)

Alwall et al (2014)

Frederix, Pagani, Zaro (2017)

$K_{\text{factor}} = 1.58$

$$\begin{aligned}\sigma_g^{\text{SM}} &= 1.216 \\ \sigma_{Z/\gamma}^{\text{SM}} &= 0.412\end{aligned}$$

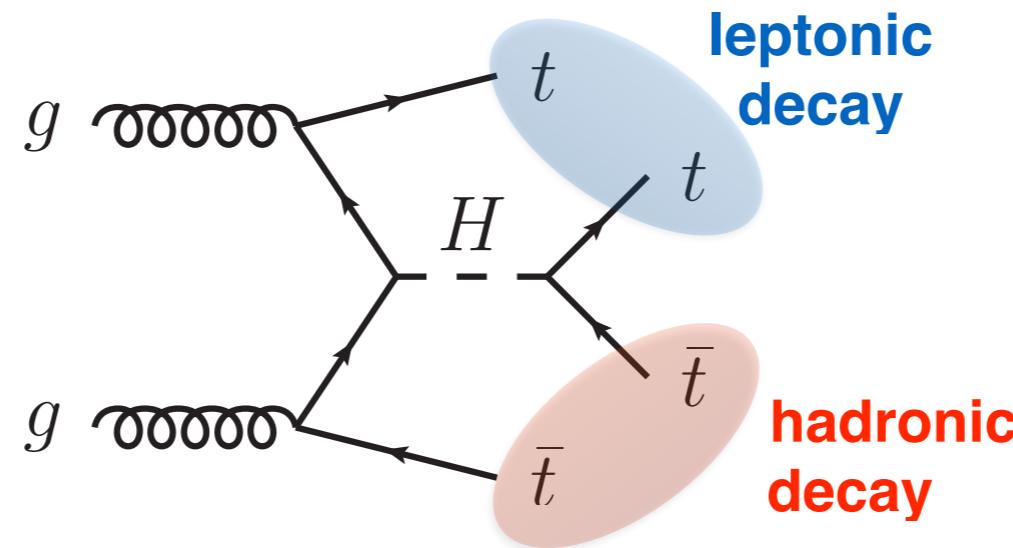
$$\sigma_{g+Z/\gamma}^{\text{SM}} = -0.262$$

significant cancellations between strong and weak interactions

→ **large scale dependence**

Collider simulation

Event Topology: **same-sign charged leptons plus multi-jet (b-jet)**



Event selections:

$$p_T^{j,l} \geq 20\text{GeV}$$

$$|\eta^{j,l}| < 2.5$$

$$N_{l^\pm} = 2$$

$$N_{b-\text{jets}} \geq 3$$

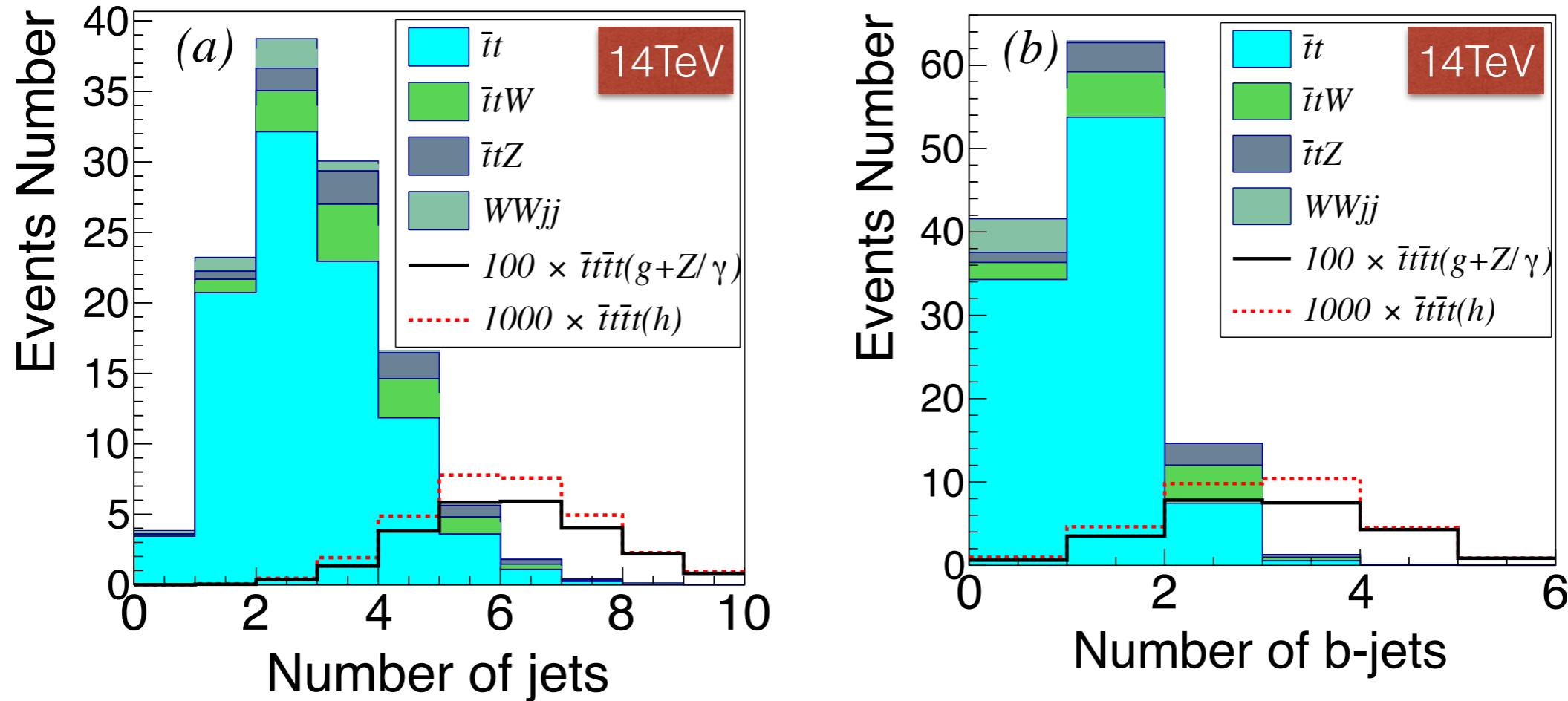
	13-14TeV	27TeV	100TeV
N_{jets}	≥ 5	≥ 6	≥ 6
\cancel{E}_T	$\geq 100\text{GeV}$	150GeV	150GeV
M_T	$\geq 100\text{GeV}$		
H_T	$\geq 700\text{GeV}$	700GeV	800GeV

Backgrounds: $t\bar{t}Z$, $t\bar{t}W^\pm$, $W^\pm W^\pm jj$, $t\bar{t}$

$$K_F^{t\bar{t}W^+} = 1.22 \quad K_F^{t\bar{t}W^-} = 1.27 \quad K_F^{t\bar{t}Z} = 1.49 \quad K_F^{W^+W^+jj} = 0.9 \quad K_F^{t\bar{t}} = 1.4 \quad @14\text{TeV}$$

Collider simulation

Event Topology: same-sign charged leptons plus multi-jet (b-jet)

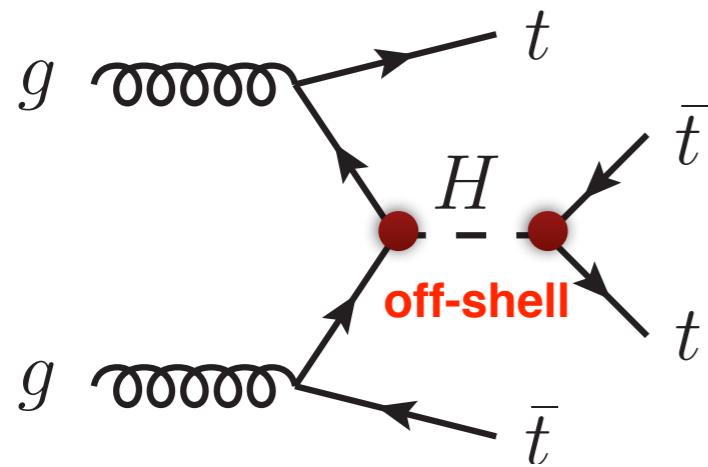


Not adequate to claim a discovery of κ_t at LHC but could set a bound

$$\kappa_t \leq 1.34 \text{ @ 13 TeV LHC (300fb}^{-1}\text{)}$$

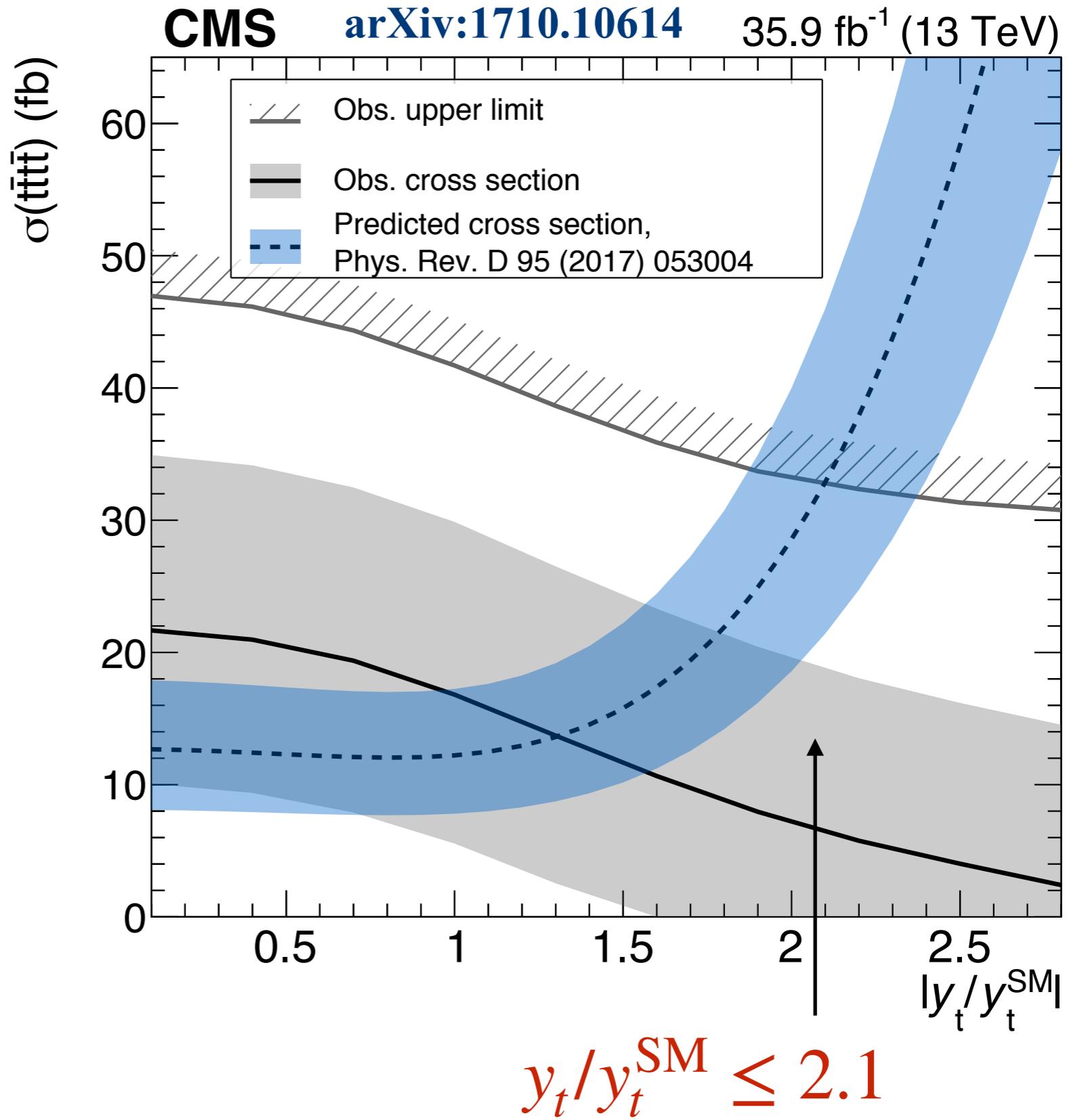
Sizing up top-quark interaction with Higgs

Four-top production



QHC, Chen, Liu
PRD95 (2017) 053004

No assumptions
made on how
Higgs-boson decays



$$\frac{\kappa_t^2}{R_\Gamma} = \mu_{t\bar{t}H}$$

assume $\kappa_X \simeq 1$ and invisible width

**13TeV LHC
(35.9fb⁻¹)**

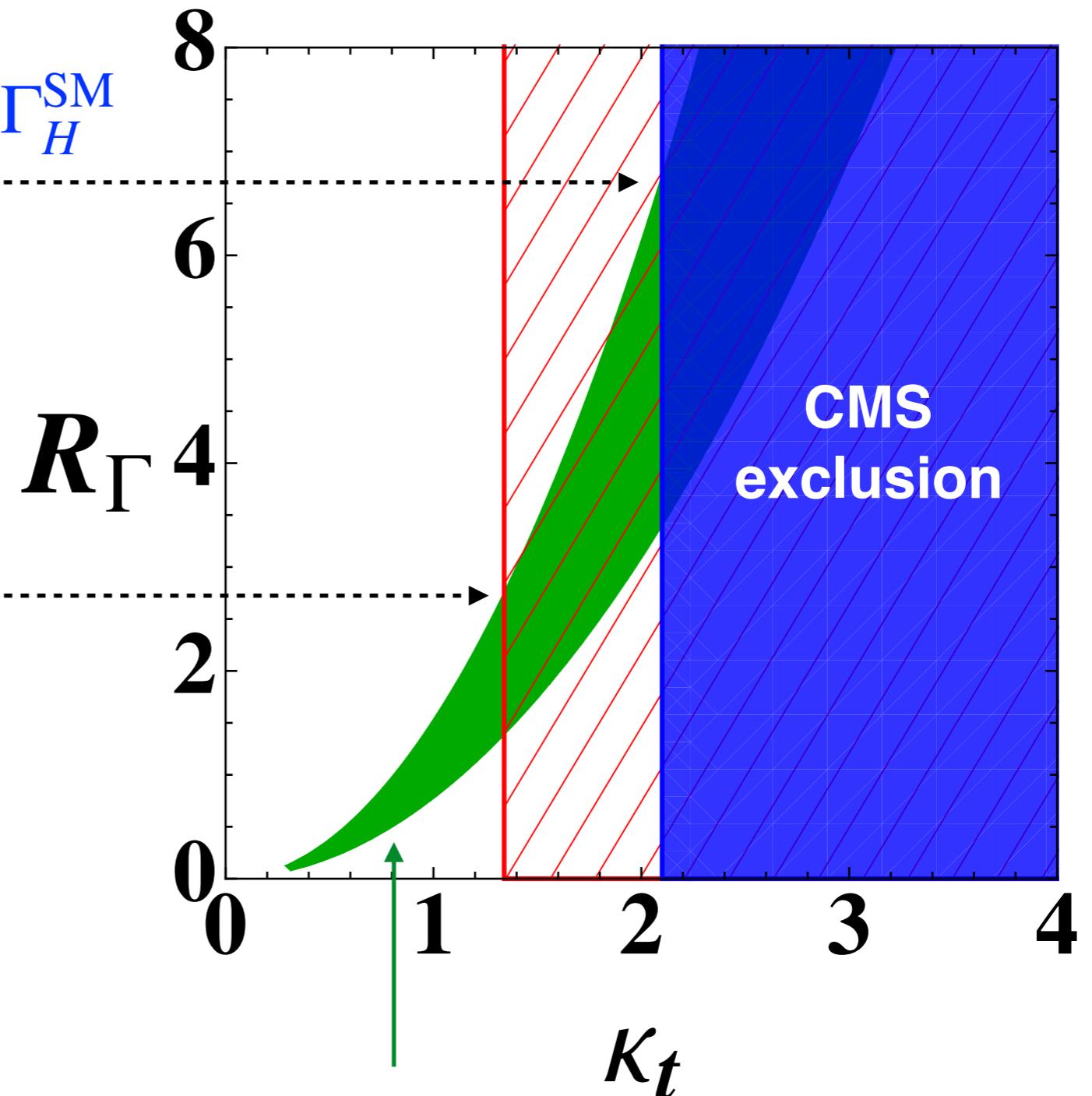
$$\kappa_t \leq 2.1$$

$$\Gamma_H \leq 6.8 \Gamma_H^{\text{SM}}$$

**13TeV LHC
(300fb⁻¹)**

$$\kappa_t \leq 1.34$$

$$\Gamma_H \leq 2.8 \Gamma_H^{\text{SM}} \\ \sim 11\text{MeV}$$



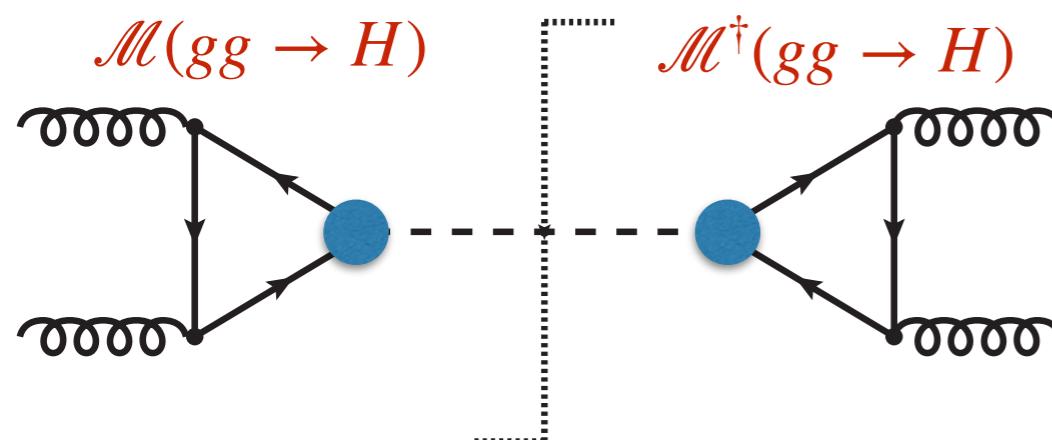
$$\mu_{t\bar{t}H} = 0.96^{+0.34}_{-0.31}$$

CMS-PAS-HIG-18-019

Q3: CP property of top-Higgs interaction

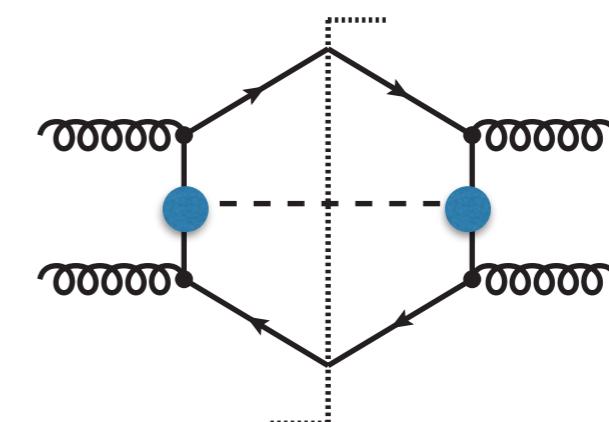
$$\mathcal{L}_{Ht\bar{t}} = -\frac{m_t}{v} H \bar{t}(a_t + i b_t \gamma_5) t$$

Boudjema, Godbole, Guadagnoli, Mohan,
1501.03157



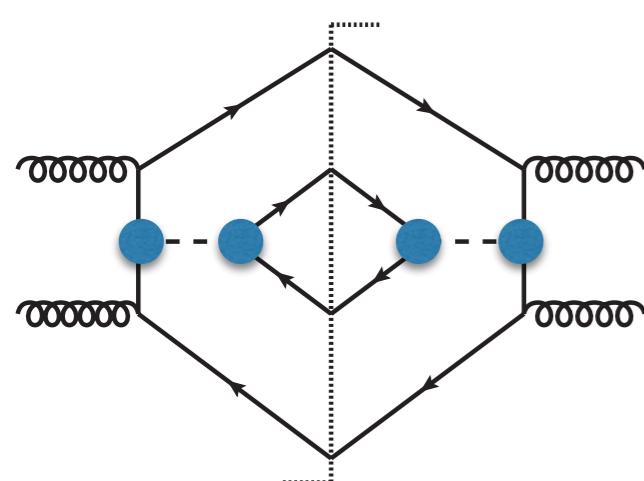
$$\frac{\sigma(gg \rightarrow H)}{\sigma(gg \rightarrow H)_{\text{SM}}} \sim a_t^2 + 2.26 b_t^2$$

pseudo-scalar dominates



$$\frac{\sigma(gg \rightarrow t\bar{t}H)}{\sigma(gg \rightarrow t\bar{t}H)_{\text{SM}}} \sim a_t^2 + 0.46 b_t^2 \quad @13\text{TeV}$$

scalar dominates



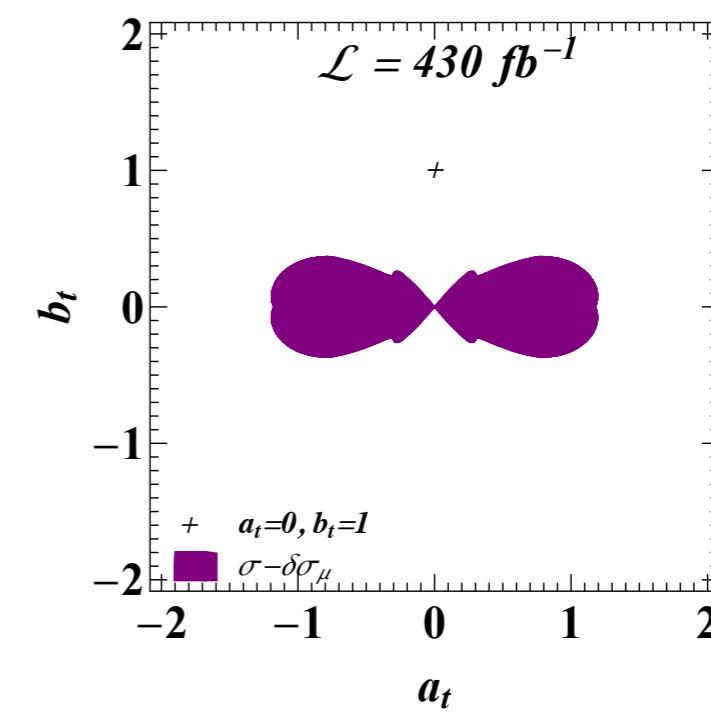
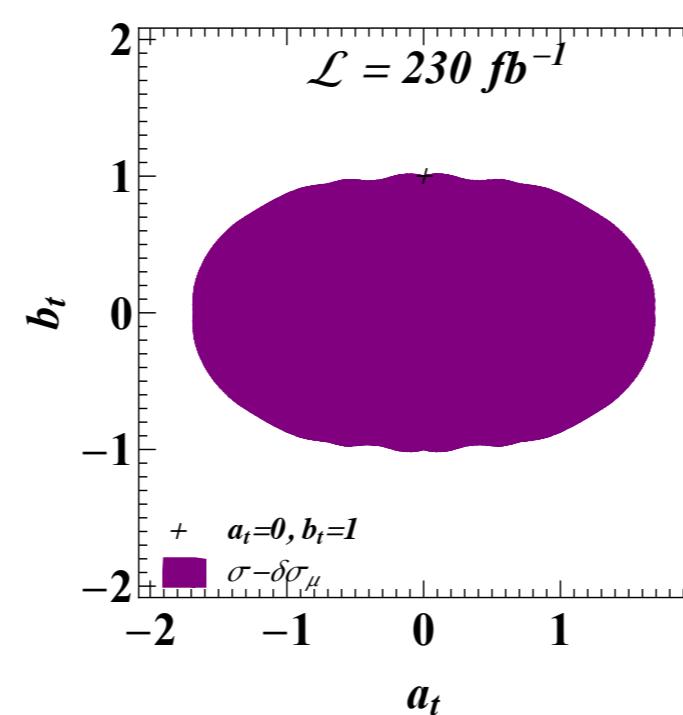
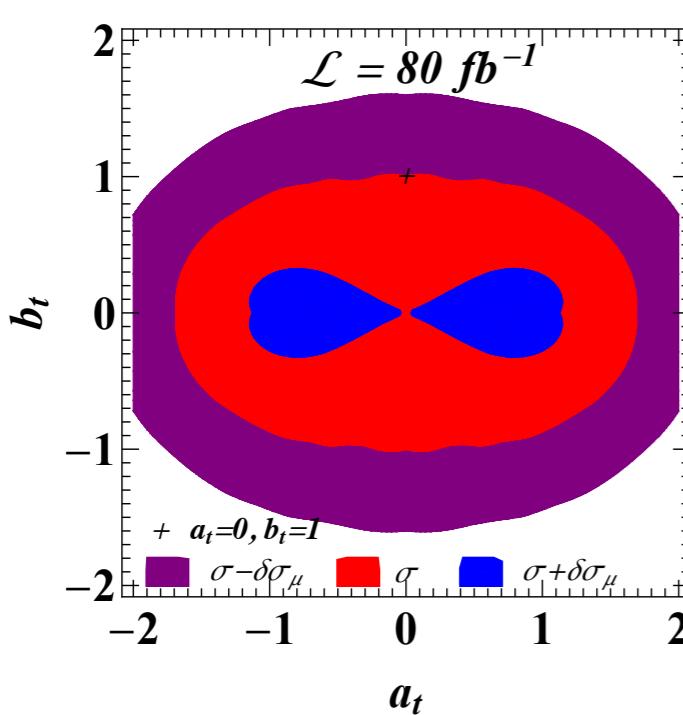
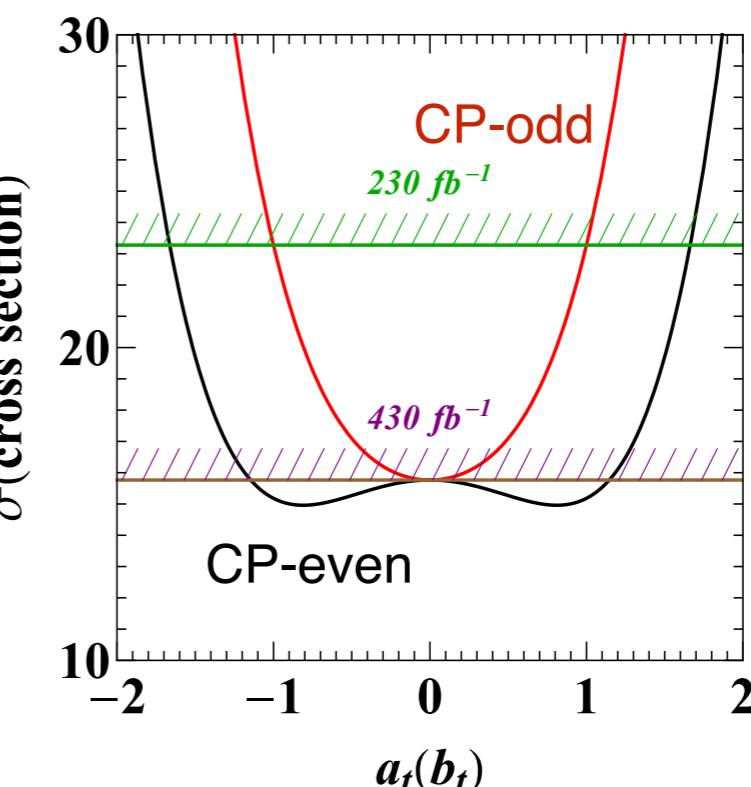
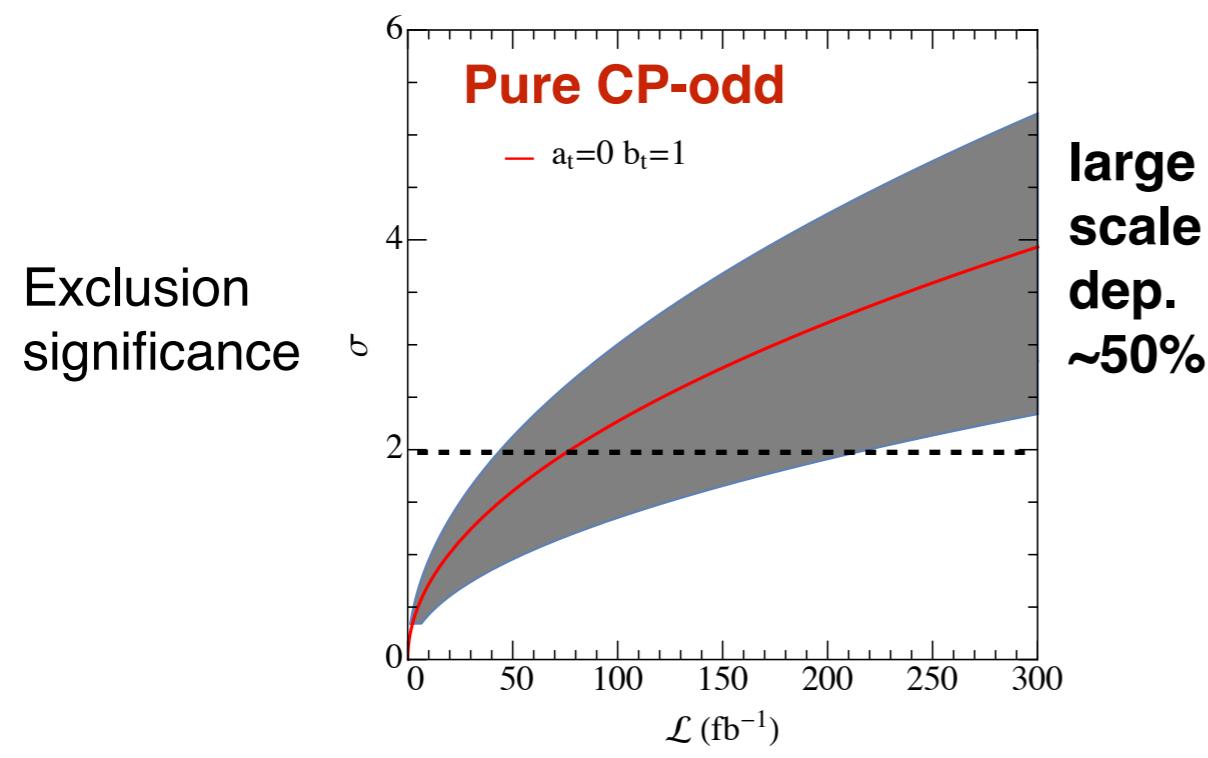
CP-odd
(a=0, b=1)

$$\sigma(t\bar{t}t\bar{t}) = 9.997 + 2.807 \times b_t^2 + 1.788 \times b_t^4$$

CP-even
(a=1, b=0)

$$\sigma(t\bar{t}t\bar{t}) = 9.997 - 1.547 \times a_t^2 + 1.108 \times a_t^4$$

Potential of four-top measurements

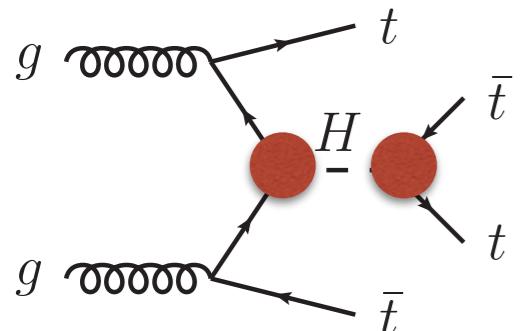


13TeV

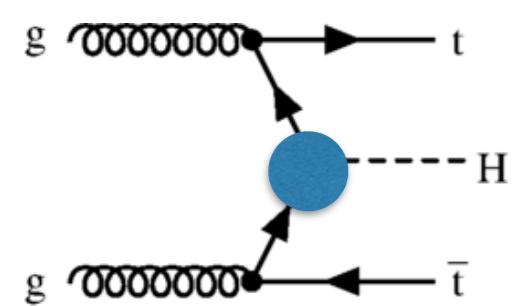
230 fb^{-1} : excludes a pure CP-odd Htt coupling with $b_t=1$

430 fb^{-1} : excludes a pure CP-odd Htt coupling entirely (any b_t)

Global analysis of multi-top productions

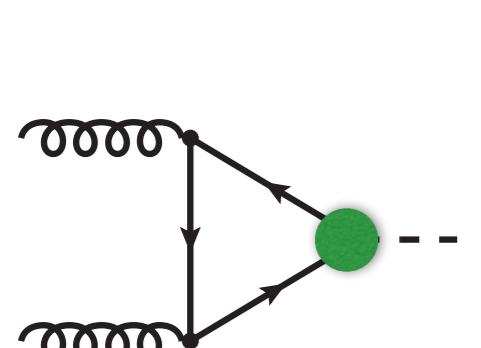


$$\sigma(t\bar{t}t\bar{t}) = 9.998 - 1.522a_t^2 + 2.883b_t^2 + 1.173a_t^4 + 2.713a_t^2b_t^2 + 1.827b_t^4$$



$$\frac{\sigma(gg \rightarrow t\bar{t}H)}{\sigma(gg \rightarrow t\bar{t}H)_{SM}} = (a_t^2 + 0.46b_t^2) \frac{1}{R_\Gamma}$$

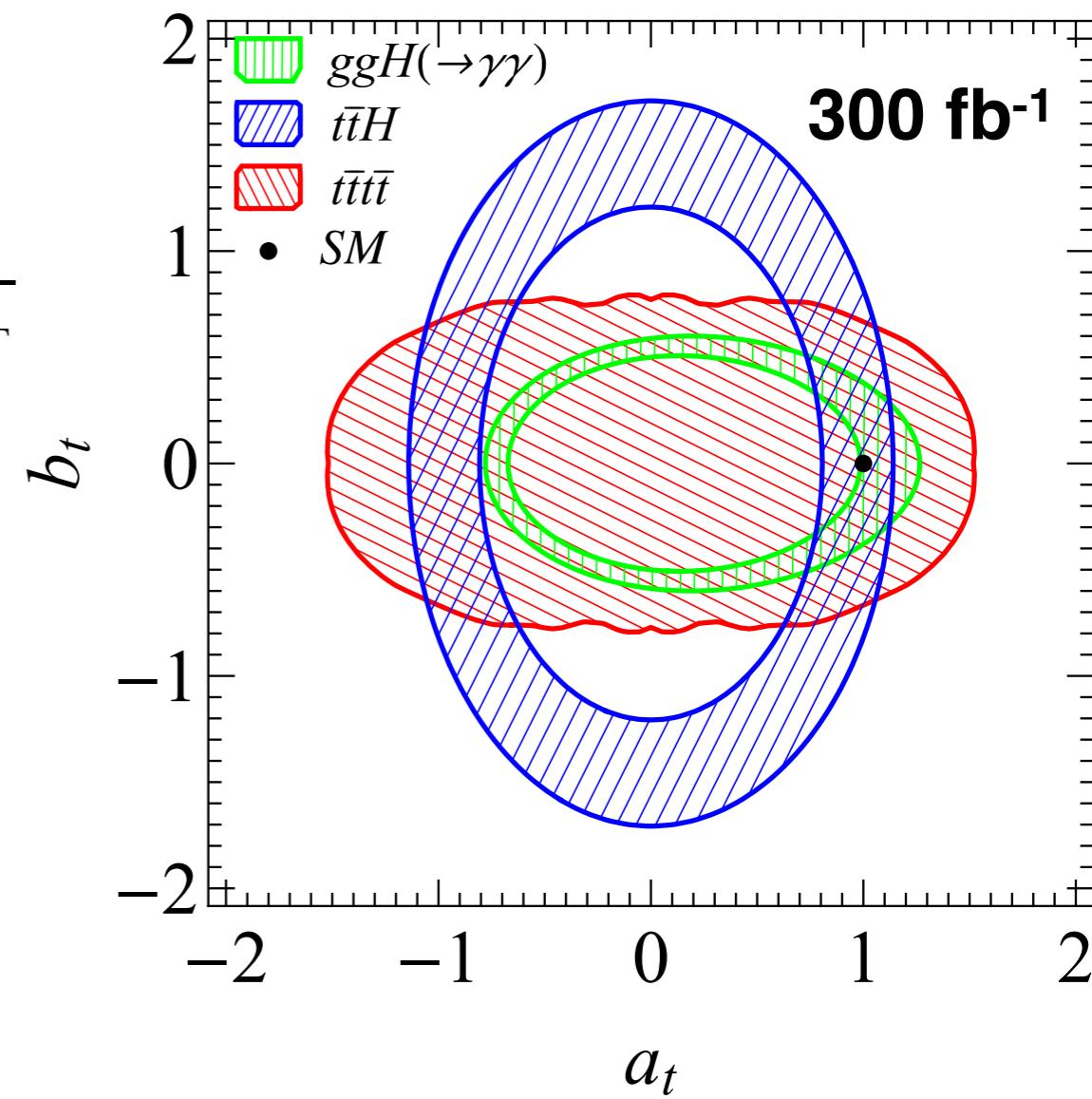
$$\mu_{t\bar{t}H} = 0.96^{+0.34}_{-0.31}$$



$$\begin{aligned} & \frac{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)}{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)_{SM}} \\ &= (a_t^2 + 2.26b_t^2) \\ & \times \frac{0.035(-7.2 + 1.83a_t^2) + 0.27b_t^2}{R_\Gamma} \end{aligned}$$

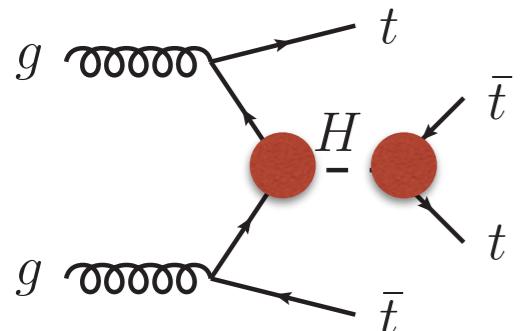
$$\mu_{gg \rightarrow H \rightarrow \gamma\gamma} = 1.16^{+0.21}_{-0.18}$$

13TeV LHC

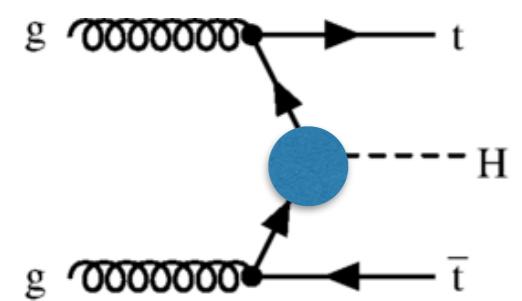


$0.73 < a < 1.14$
 $-0.51 < b < 0.51$

Global analysis of multi-top productions

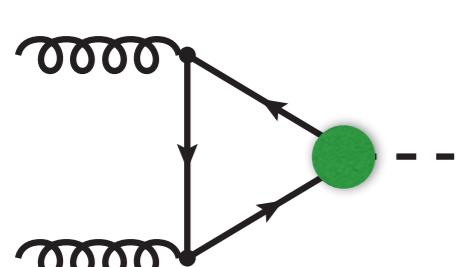


$$\sigma(t\bar{t}t\bar{t}) = 9.998 - 1.522a_t^2 + 2.883b_t^2 + 1.173a_t^4 + 2.713a_t^2b_t^2 + 1.827b_t^4$$



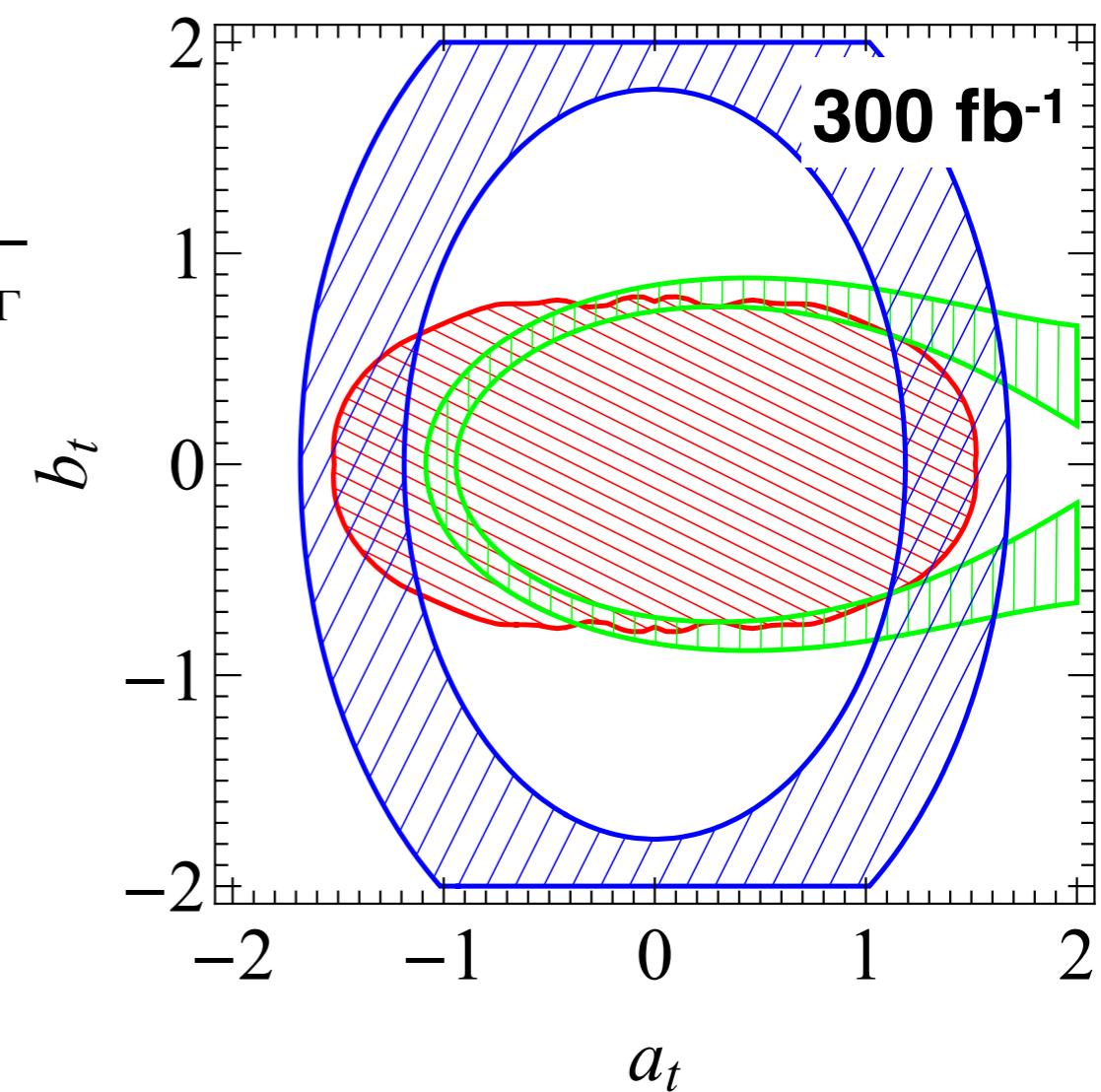
$$\frac{\sigma(gg \rightarrow t\bar{t}H)}{\sigma(gg \rightarrow t\bar{t}H)_{\text{SM}}} = (a_t^2 + 0.46b_t^2) \frac{1}{R_\Gamma}$$

$$\mu_{t\bar{t}H} = 0.96^{+0.34}_{-0.31}$$



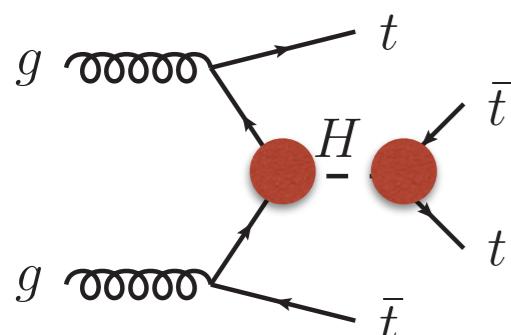
$$\begin{aligned} & \frac{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)}{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)_{\text{SM}}} \\ &= (a_t^2 + 2.26b_t^2) \\ & \times \frac{0.035(-7.2 + 1.83a_t^2) + 0.27b_t^2}{R_\Gamma} \end{aligned}$$

$$\mu_{gg \rightarrow H \rightarrow \gamma\gamma} = 1.16^{+0.21}_{-0.18}$$



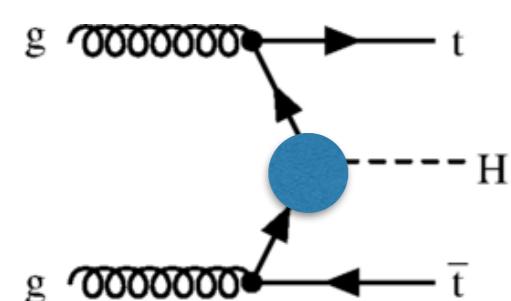
$R_\Gamma < 2.17$

Global analysis of multi-top productions



$$\sigma(t\bar{t}t\bar{t}) = 9.998 - \frac{1.522a_t^2 + 2.883b_t^2 + 1.173a_t^4 + 2.713a_t^2b_t^2 + 1.827b_t^4}{R_\Gamma}$$

assume $\mu_{t\bar{t}t\bar{t}}^H = 1.00^{+0.50}_{-0.50}$



$$\frac{\sigma(gg \rightarrow t\bar{t}H)}{\sigma(gg \rightarrow t\bar{t}H)_{SM}} = (a_t^2 + 0.46b_t^2) \frac{1}{R_\Gamma}$$

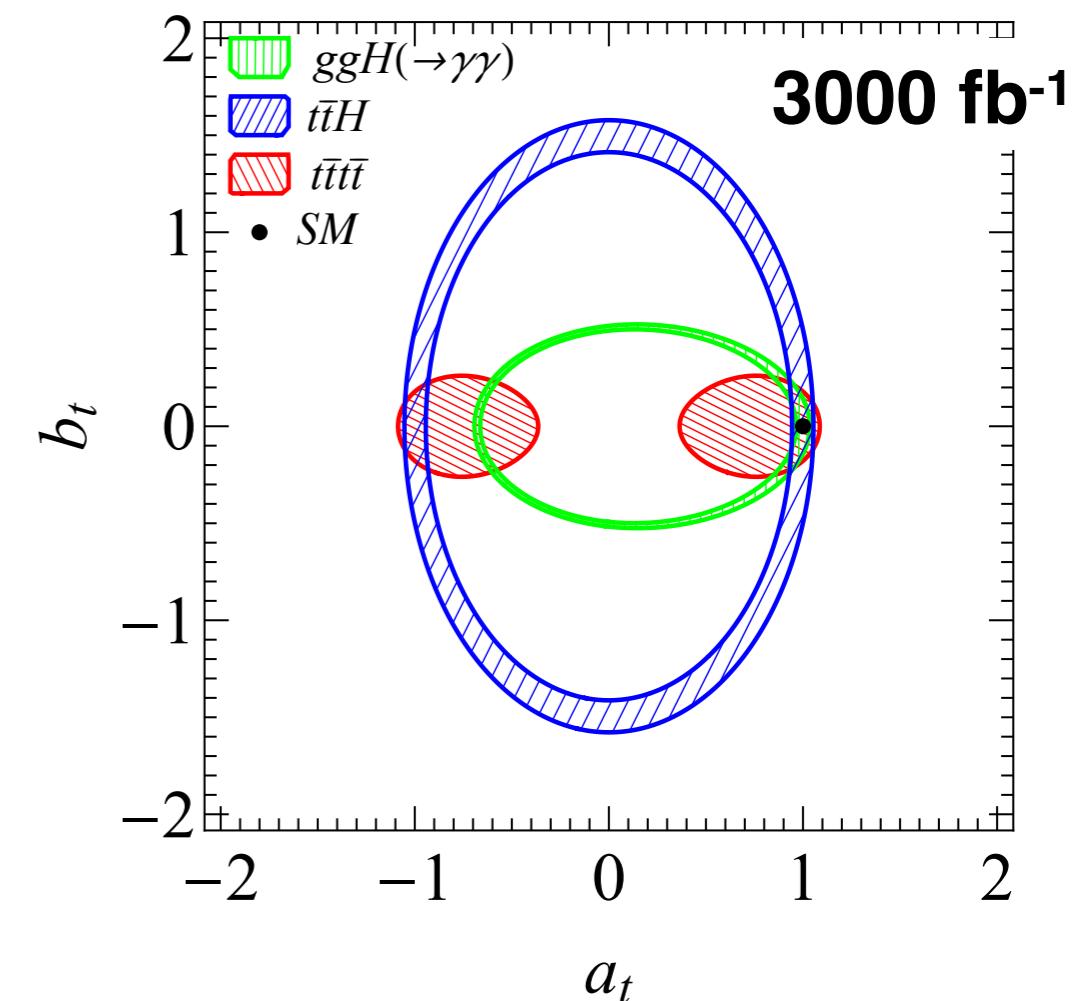
$$\mu_{t\bar{t}H} = 0.96^{+0.34}_{-0.31}$$

$$\frac{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)}{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)_{SM}}$$

$$= (a_t^2 + 2.26b_t^2)$$

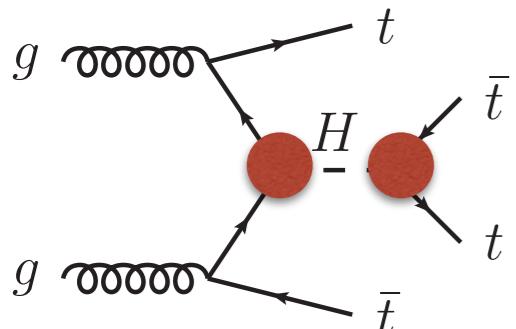
$$\times \frac{0.035(-7.2 + 1.83a_t^2) + 0.27b_t^2}{R_\Gamma}$$

$$\mu_{gg \rightarrow H \rightarrow \gamma\gamma} = 1.00^{+0.05}_{-0.05}$$



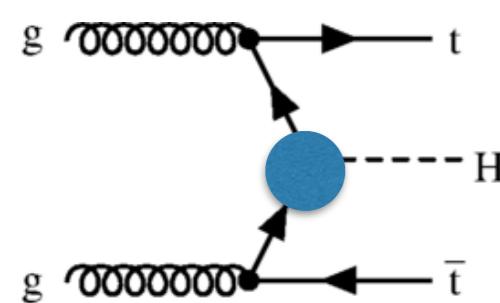
0.93 < a < 1.03
 -0.22 < b < 0.22

Global analysis of multi-top productions



$$\sigma(t\bar{t}t\bar{t}) = 9.998 - \frac{1.522a_t^2 + 2.883b_t^2 + 1.173a_t^4 + 2.713a_t^2b_t^2 + 1.827b_t^4}{R_\Gamma}$$

assume $\mu_{t\bar{t}t\bar{t}}^H = 1.00^{+0.5}_{-0.5}$



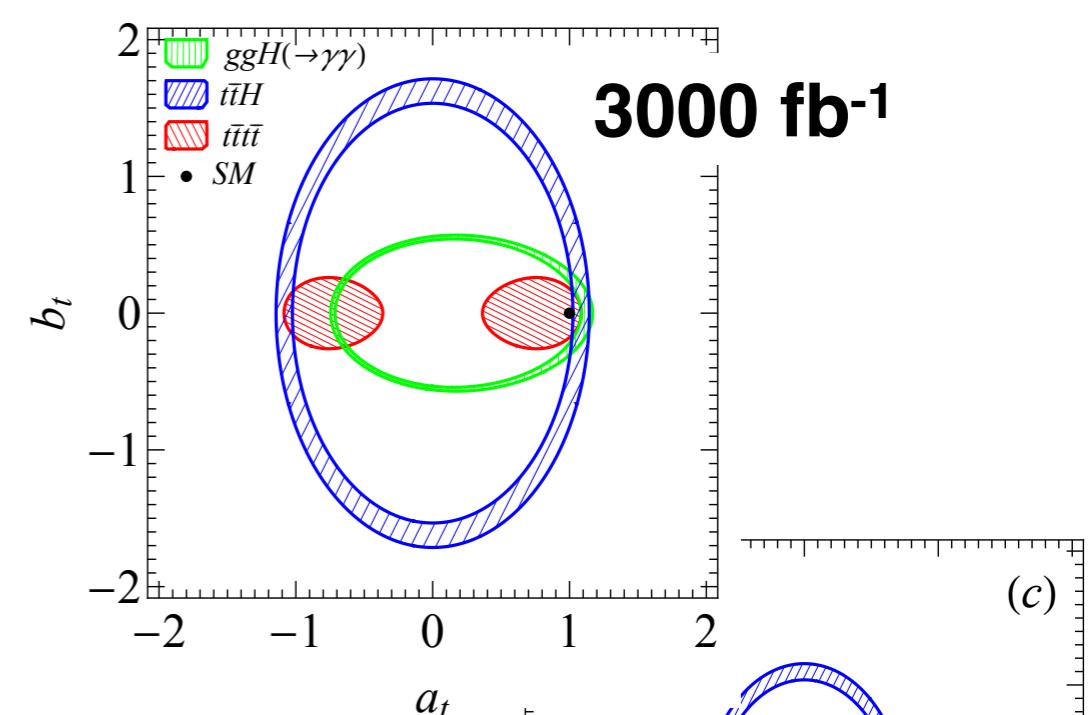
$$\frac{\sigma(gg \rightarrow t\bar{t}H)}{\sigma(gg \rightarrow t\bar{t}H)_{SM}} = (a_t^2 + 0.46b_t^2) \frac{1}{R_\Gamma}$$

$$\mu_{t\bar{t}H} = 0.96^{+0.34}_{-0.31}$$

$$\begin{aligned} & \frac{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)}{\sigma(gg \rightarrow H \rightarrow \gamma\gamma)_{SM}} \\ &= (a_t^2 + 2.26b_t^2) \\ & \times \frac{0.035(-7.2 + 1.83a_t^2) + 0.27b_t^2}{R_\Gamma} \end{aligned}$$

$$\mu_{gg \rightarrow H \rightarrow \gamma\gamma} = 1.00^{+0.05}_{-0.05}$$

@ HL-LHC



$0.54 < R_\Gamma < 1.18$

Summary

The four top-quark production can constrain top-quark Yukawa coupling without assumptions on Higgs boson width or decay branching ratios.

The four top-quark production is sensitive to the CP property of top-Higgs interaction.

$$\mathcal{L}_{Ht\bar{t}} = - \frac{m_t}{v} H \bar{t} (a_t + i b_t \gamma_5) t$$

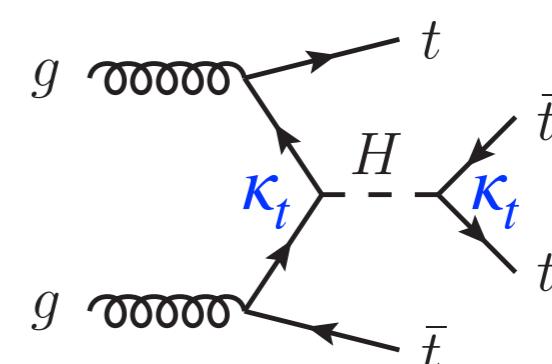
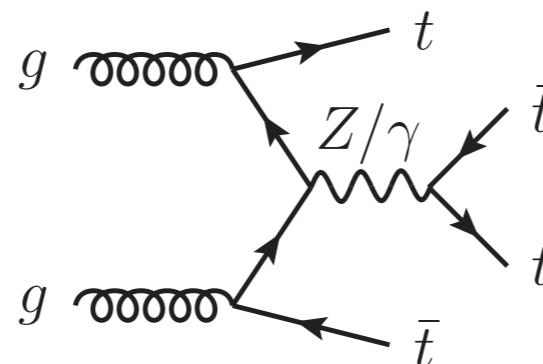
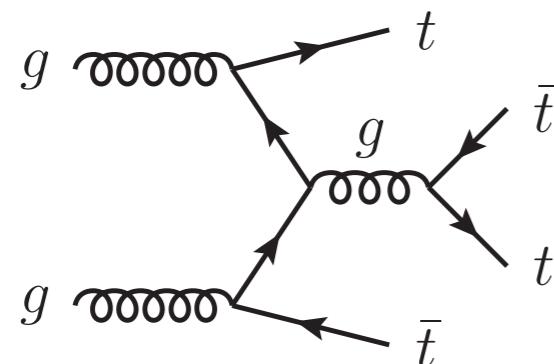
Global fitting of single-Higgs, $t\bar{t}t\bar{t}$ and $t\bar{t}h$ channels yields

$$0.93 < a < 1.13, \quad -0.22 < b < 0.22$$

$$0.54\Gamma_H^{\text{SM}} < \Gamma_H \leq 1.18\Gamma_H^{\text{SM}}$$

Backup Slides

Measuring κ_t from four top-quark production



$$\sigma(t\bar{t}t\bar{t}) = \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g/Z/\gamma} + \kappa_t^2 \sigma_{\text{int}}^{\text{SM}} + \kappa_t^4 \sigma^{\text{SM}}(t\bar{t}t\bar{t})_H$$

LO	8 TeV	1.344	-0.224	0.171	
13 TeV	9.997		-1.547	1.108	in unit of fb
14 TeV	13.14		-2.007	1.515	
27 TeV	115.1		-15.57	11.73	
100 TeV	3276		-356.9	273.1	
Relative ratio	8~12		-1.3	1	
			cancel out around SM $\kappa_t = 1$		

Results for 4-top: inclusive results

- The gg-initiated contributions ($\text{LO}_{1\rightarrow 3}$, $\text{NLO}_{1\rightarrow 4}$) are the dominant ones
- $\text{LO}_{2,3}$ are large and with different sign
- $\text{NLO}_{2,3}$ are sizeable too, and show a very large scale dependence: they are mostly due to QCD-type corrections from $\text{LO}_{2,3}$ respectively
- Remarkably, there is an almost complete cancellation between NLO_2 and NLO_3 which is also quite stable w.r.t. the scale choice
- Such a cancellation will be spoiled by e.g. non-SM Yukawa interactions. A LO-based study may not capture the correct behaviour of the cross section

13 TeV

	$\delta[\%]$	$\mu = H_T/8$	$\mu = H_T/4$	$\mu = H_T/2$
LO ₂		-26.0	-28.3	-30.5
LO ₃		32.6	39.0	45.9
LO ₄		0.2	0.3	0.4
LO ₅		0.02	0.03	0.05
NLO ₁		14.0	62.7	103.5
NLO ₂		8.6	-3.3	-15.1
NLO ₃		-10.3	1.8	16.1
NLO ₄		2.3	2.8	3.6
NLO ₅		0.12	0.16	0.19
NLO ₆		< 0.01	< 0.01	< 0.01
NLO ₂ + NLO ₃		-1.7	-1.6	0.9

R. Frederix, D. Pagani, and M. Zaro,
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