

# From Underlying Event Sensitive To Insensitive

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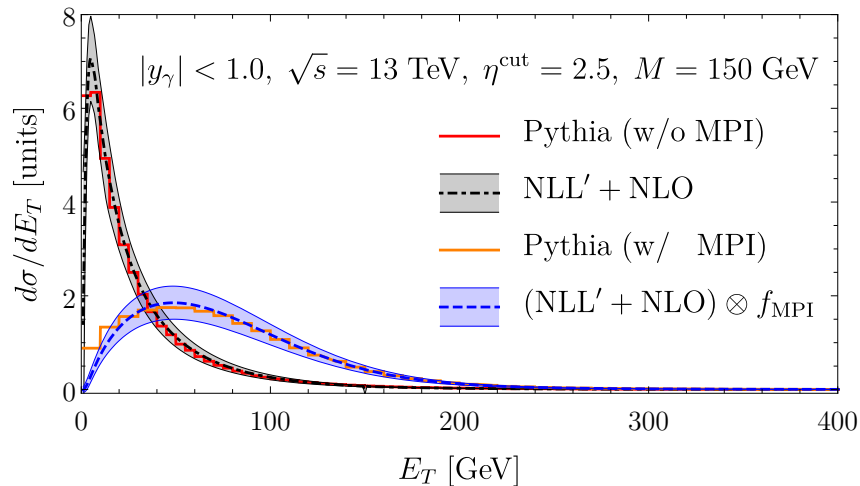
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# Many observables are sensitive to UE/PU

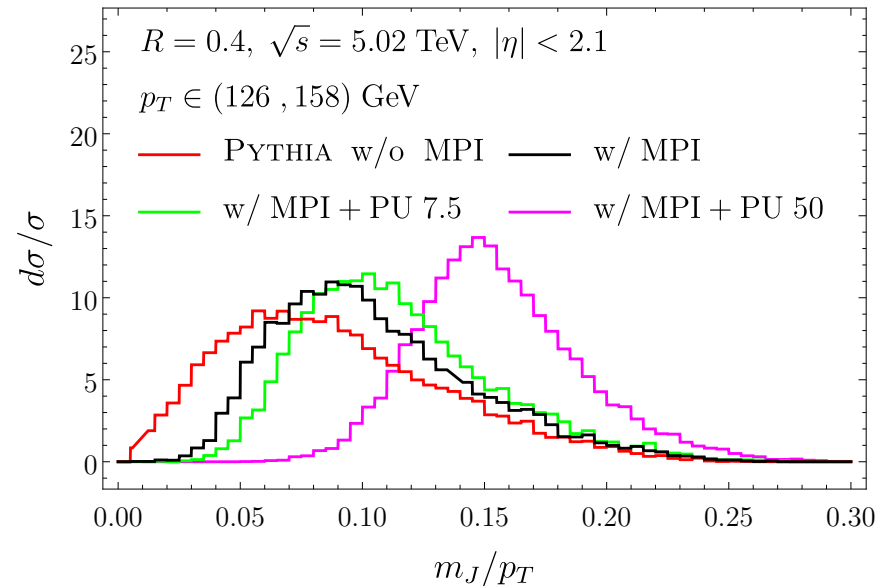
- Drell-Yan: transverse energy (cf pT), beam thrust
- jet substructures: jet mass, angularity, n-subjettiness, ...

Drell-Yan  $E_T$  distribution



$$E_T(\eta_{\text{cut}}) = \sum_i p_T^{(i)} \Theta(\eta_{\text{cut}} - |\eta^{(i)}|)$$

jet mass distribution



# Many observables are sensitive to UE/PU

- ❑ No first principle approach is available for UE.
- ❑ Large source of uncertainties in theory predictions.
- ❑ Several options to handle:
  - ❑ tuned models MPI/soft QCD in Monte Carlo Simulation
  - ❑ useful robust observables:  $p_T$  (over ET), small jet radius
  - ❑ removing soft particles from jet using grooming algorithms
  - ❑ *new statistical approach: subtracted cumulants*

# A lesson with a simple model in Drell-Yan

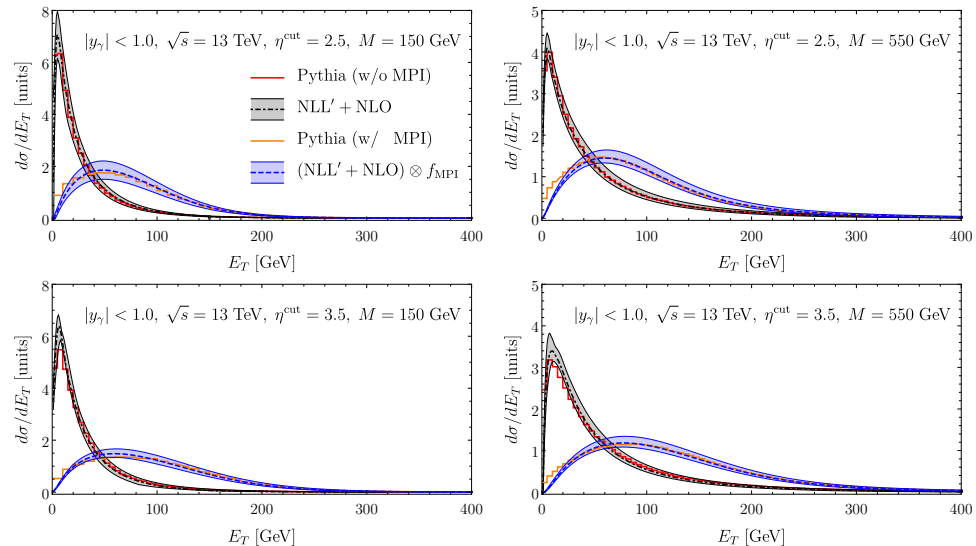
- An excellent agreement between a simple model and Pythia MPI.

$$f_{\text{MPI}}(E_T, \eta_{\text{cut}}) = \mathcal{N} \exp \left[ - \left( \frac{E_T}{\alpha(\eta_{\text{cut}}) \sqrt{\pi}} \right)^2 \right]$$

- A single parameter is constant in wide range of hard scales and rapidities.

$$\alpha(\eta_{\text{cut}}) = A \eta_{\text{cut}}$$

$$A = 22.7 \text{ GeV}$$



- Similar observations in jet mass: inclusive jet, H/Z + jet process, top-quark jet



# Subtraction of moments

- The model is un-correlated to hard scale  $Q$ , while the perturbative contribution is correlated.

$$f_{\text{MPI}}(E_T, \eta_{\text{cut}}) = \mathcal{N} \exp \left[ - \left( \frac{E_T}{\alpha(\eta_{\text{cut}}) \sqrt{\pi}} \right)^2 \right]$$

- The model UE contribution is additive in the moments

$$\langle E_T(Q) \rangle = \langle E_T(Q) \rangle_{\text{pert}} + \langle E_T \rangle_{\text{MPI}}$$

- Uncorrelated model contributions are cancelled in the subtraction between moments at  $Q$  and  $Q_0$

$$\begin{aligned} \Delta E_T(Q, Q_0) &\equiv \langle E_T(Q) \rangle - \langle E_T(Q_0) \rangle \\ &= \Delta E_{T,\text{pert}}(Q, Q_0) \end{aligned}$$

# Subtracted Cumulants

- The cumulant is excellent quantity to generalize this idea
  - defined by a generating function
  - one-to-one with the moments for an observable  $e$ :

$$\kappa_1(e) = \langle e \rangle \quad \kappa_2(e) = \langle e^2 \rangle - \langle e \rangle^2 \quad \kappa_3(e) = \langle e^3 \rangle - 3\langle e^2 \rangle \langle e \rangle + 2\langle e \rangle^3$$

- **Additivity** if  $e_S$  and  $e_B$  are independent

$$\kappa_n(e_S + e_B) = \kappa_n(e_S) + \kappa_n(e_B)$$

- At any  $n$ -th order the subtracted cumulants is insensitive to the uncorrelated contributions ( $e_B = \text{UE, PU}$ )
- What observables work?

Individual particle contributions are added in total.

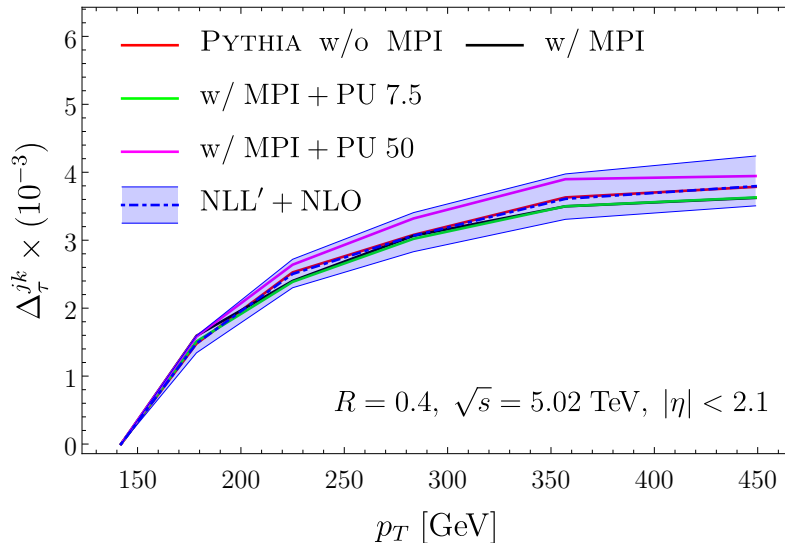
transverse energy, beam thrust, jet angularity,  $\approx$  jet mass,  $\cdots$

# Example with first cumulant of jet mass

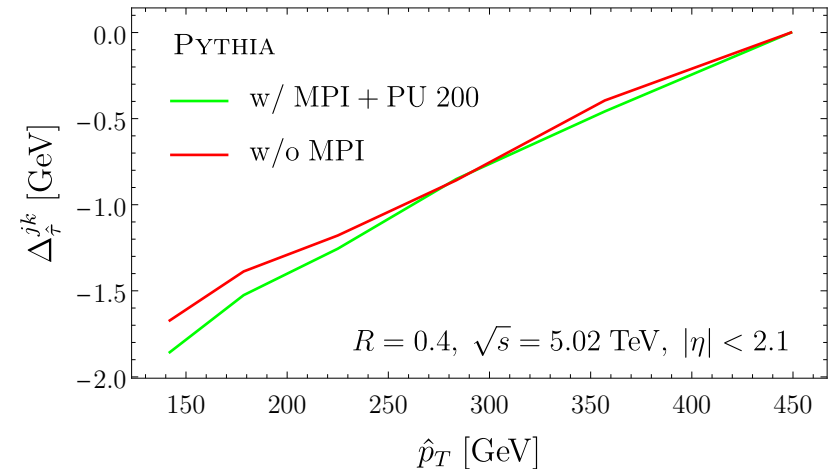
□ definitions:  $\hat{\tau} = 2 \cosh(\eta) \sum_{i \in \text{jet}} p_i^+ = \frac{m_J^2}{p_T} \left[ 1 + \mathcal{O}\left(\frac{m_J^2}{p_T^2}\right) \right]$

$$\Delta_{\tau}^{jk} = \langle \tau \rangle^{[j]} - \langle \tau \rangle^{[k]} \frac{\langle p_T^{-1} \rangle^{[j]}}{\langle p_T^{-1} \rangle^{[k]}}$$

□ Robust under Pythia UE and PU



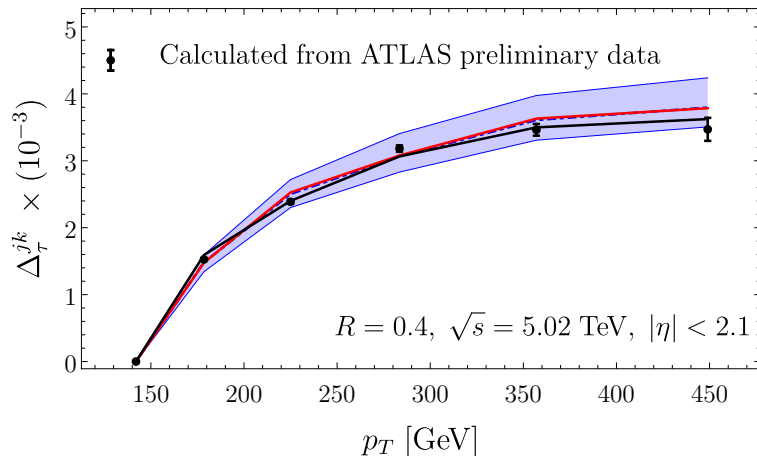
high luminosity (200 PU)



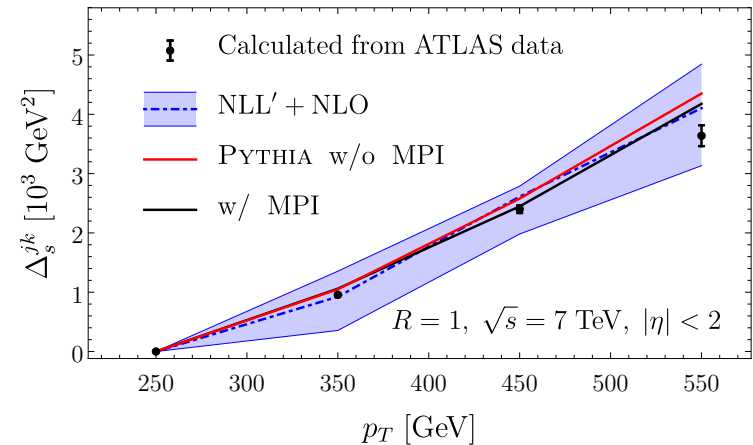
# Comparison to ATLAS measurement

- We use ATLAS data for mass distributions at 5 TeV and at 7 TeV and compute subtracted cumulants.

**5 TeV**



**7 TeV**



- two curves are different:  $\tau = m_J^2/p_T^2$  (left)  $s = m_J^2$  (right)
- *Error bars only include statistical uncertainties.*

# Summary and outlook

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- ❑ excellent performance of a simple UE model observed.
- ❑ subtracted cumulants for DY and for jet substructure:  
mitigating large background such as UE, PU  
a simple way to compare theory calculations to measurement
- ❑ example: jet mass shows a good agreement with ATLAS data
  
- ❑ higher-order cumulants, more jet substructure observables
- ❑ precision predictions beyond NLL'+NLO are available for several jet substructure and can be tested using our method.
- ❑ application to heavy-ion collisions with hard scattering

# Backup

# Sensitivity to quark/gluon jet fraction

- Medium effect in heavy-ion collisions changes q/g jet fraction.

$$\langle \tau \rangle^{[j]} = f_g^{[j]} \langle \tau \rangle_g^{[j]} + (1 - f_g^{[j]}) \langle \tau \rangle_q^{[j]}$$

- visible change in subtracted cumulant

