



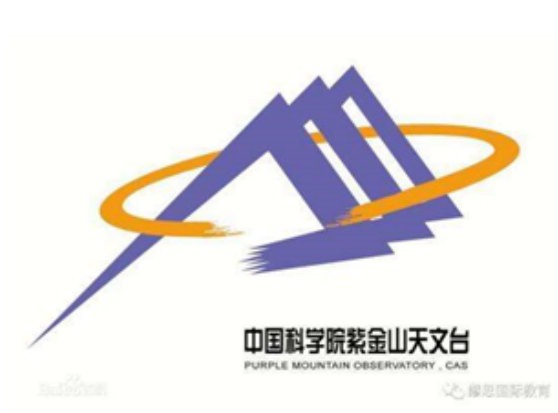
# Particle acceleration in SNRs —Origin of Galactic cosmic rays

Siming Liu, Houdun Zeng, Yiran Zhang, Zhaodong Shi, Yuliang Xin, Qiang Yuan

Purple Mountain Observatory

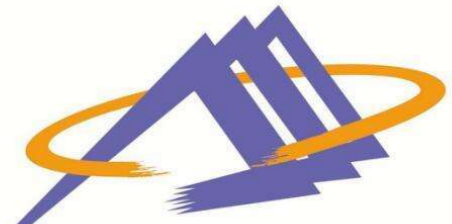
**LHAASO, 林芝**

# Outline



- Background
- Evolution of high-energy particle distribution in SNRs
- Cosmic ray electrons and positrons
- Conclusions

# Background



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Supernova remnants (SNRs) have been proposed as the dominant contributors to galactic cosmic rays (Baade & Zwicky 1934).

1、 SNRs have enough total power — 3 SNRs per century with a  $10^{51}$  erg shock kinetic energy for each SNR, a 10% efficiency leads to the observed CR density ( $1 \text{ eV/cm}^3$ );

2、 Direct evidence:

Radio emission (1948)

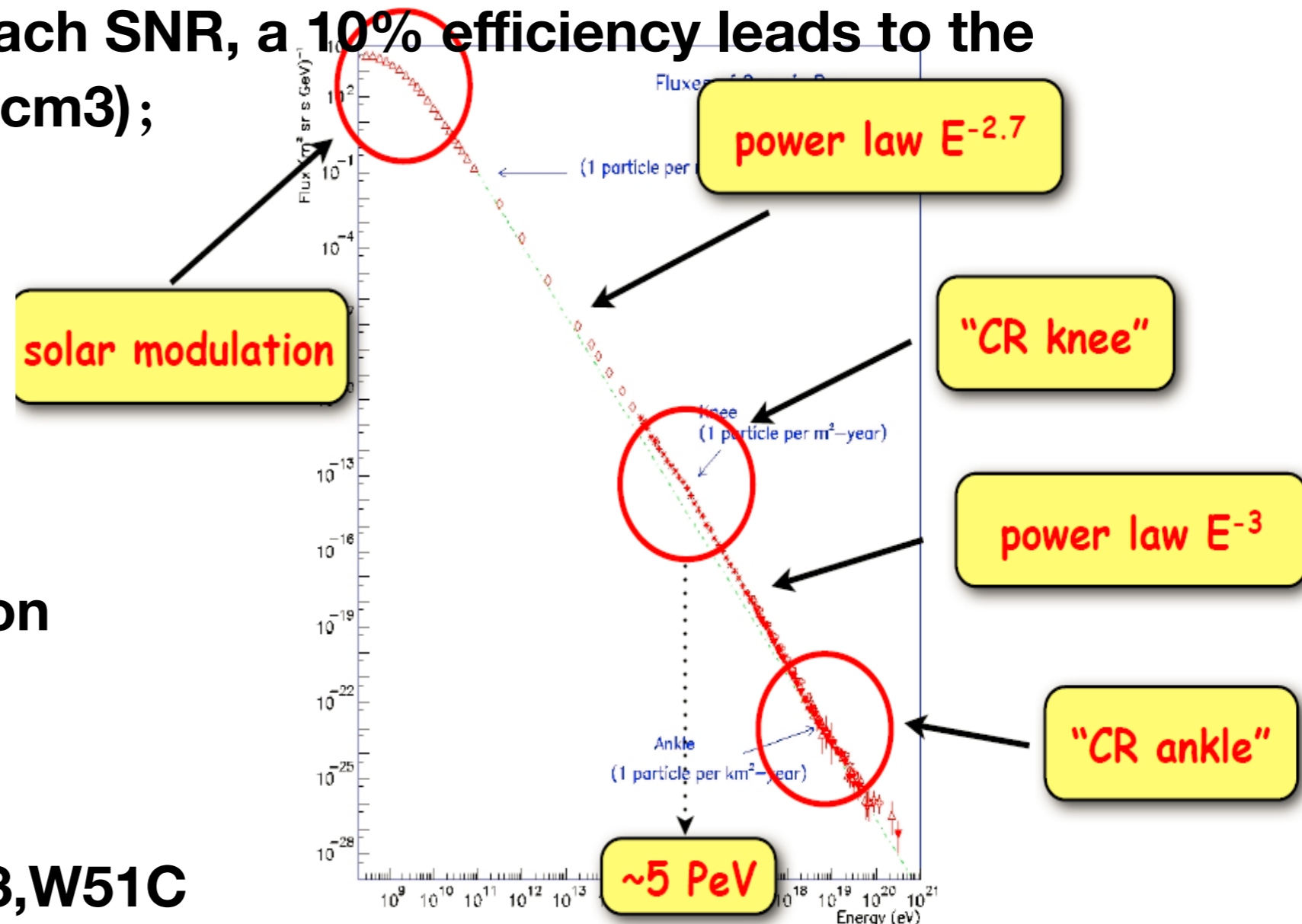
— GeV electrons

Non-thermal X-ray emission (1995)

— TeV electrons

$\pi^0$  bump (2013) W44, IC443, W51C

— GeV protons



# Background

standard theory of shock predicts  $p=2.0$

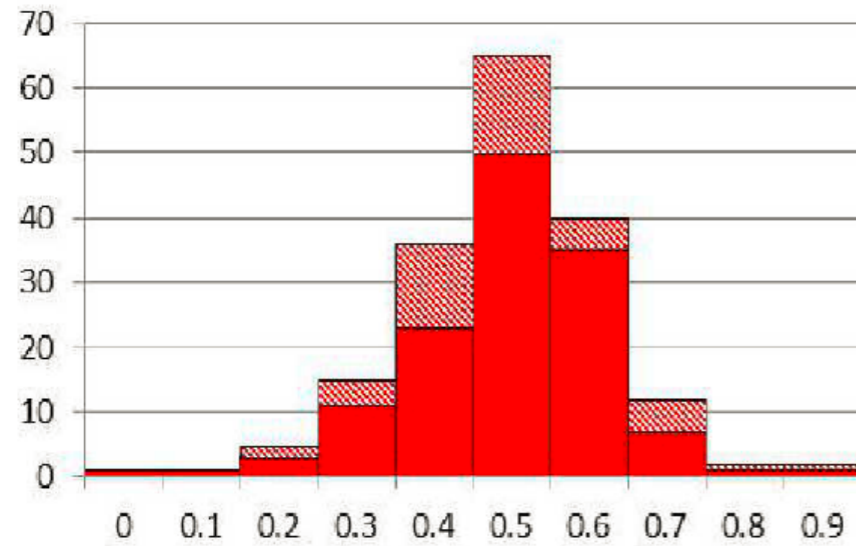
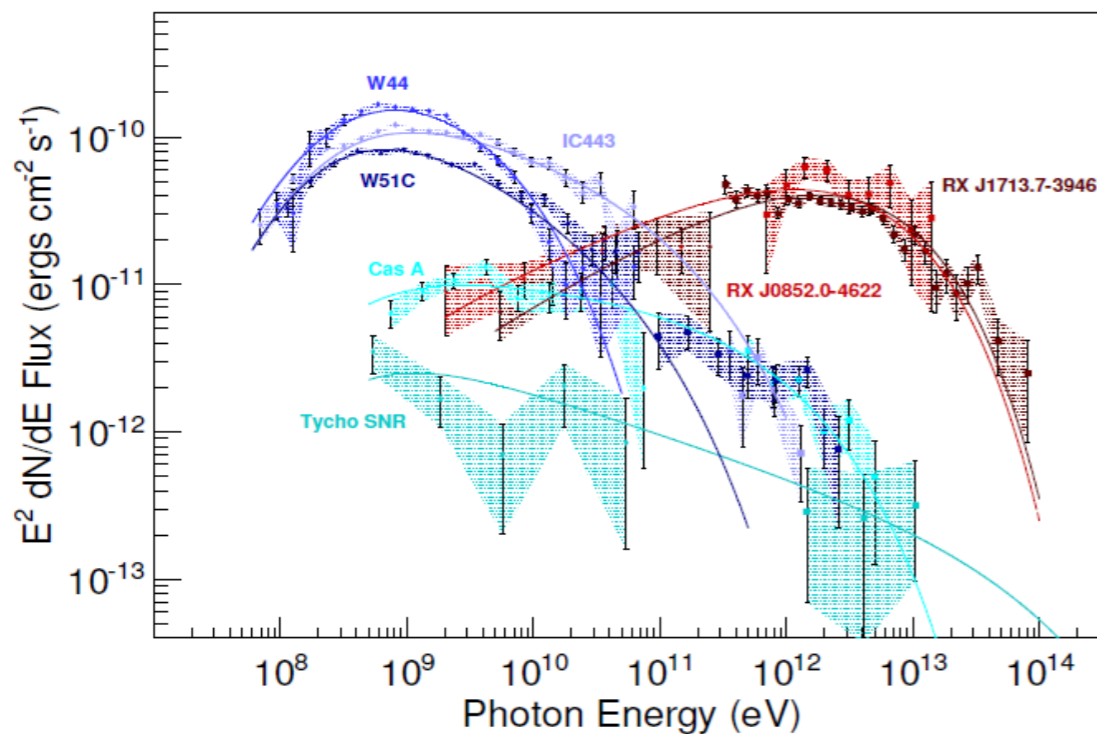
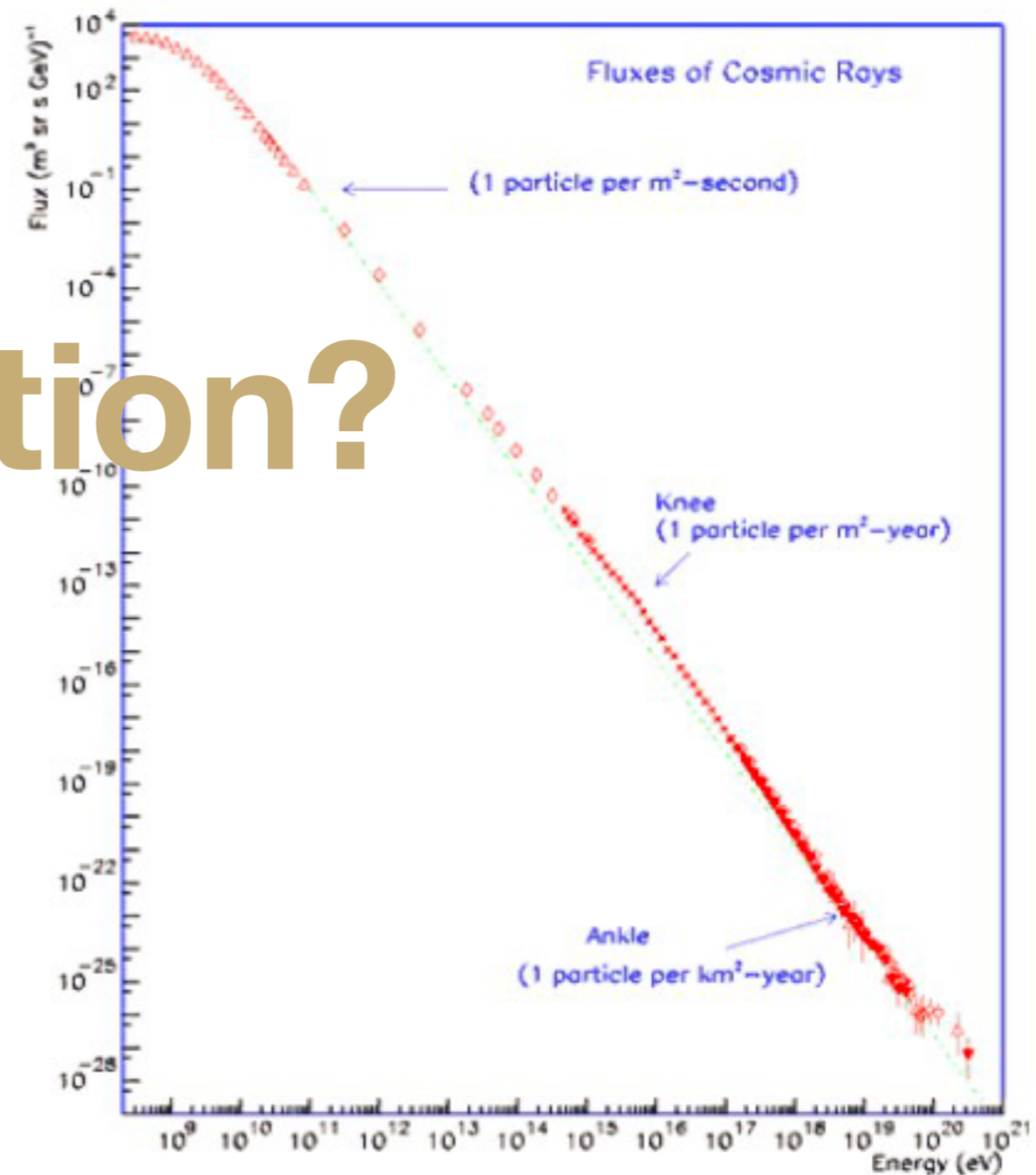


Fig. 6 Histogram of spectral indices of shell-type Galactic SNRs as extracted from Cohen (2014). The solid bars correspond to firm spectral index determination, while the shadowed ones to SNRs with uncertain spectral indices

## Relation?



# Background



Injection Power :  
 Proton ~ 3e48 erg/year  
 Electron ~ 4e46 erg/year

3 SNRs/100 yrs with  
 1e50 erg protons and  
 1e48 erg electrons  
 For each SNR.

10% efficiency for type  
 Ia SNRs with a kinetic  
 energy of 1e51 ergs

$$q^f(R, z) \propto \left(\frac{R}{R_\odot}\right)^\alpha \exp\left[-\frac{\beta(R - R_\odot)}{R_\odot}\right] \exp\left(-\frac{|z|}{z_s}\right),$$

$$q(p) \propto \begin{cases} p^{-\alpha_1}, & p < p_{br}, \\ p^{-\alpha_2}, & p \geq p_{br}, \end{cases}$$

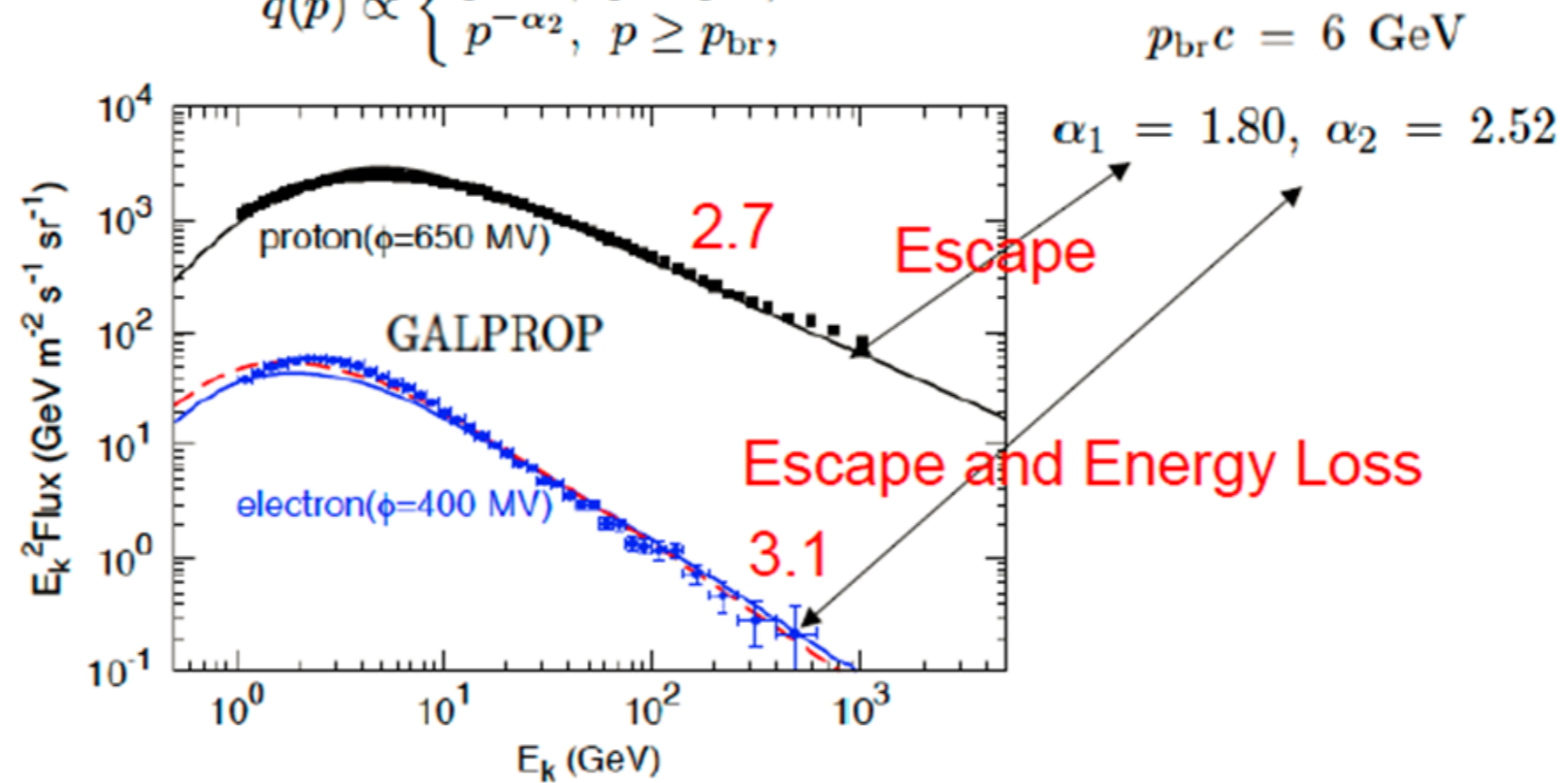


FIG. 1.— The expected fluxes of CR protons and electrons at the Earth, for the same spectral shape of the injected particles, compared with the PAMELA observational data (Adriani et al. 2011a,b). We adopt two parameter settings to calculate the electron spectrum: for solid line the magnetic field is the canonical one adopted in GALPROP and  $K_{ep} \approx 1.3\%$ ; for dashed line the magnetic field is two times larger and  $K_{ep} \approx 1.9\%$ .

# Background



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THE ASTROPHYSICAL JOURNAL LETTERS, 844:L3 (5pp), 2017 July 20

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<https://doi.org/10.3847/2041-8213/aa7de1>

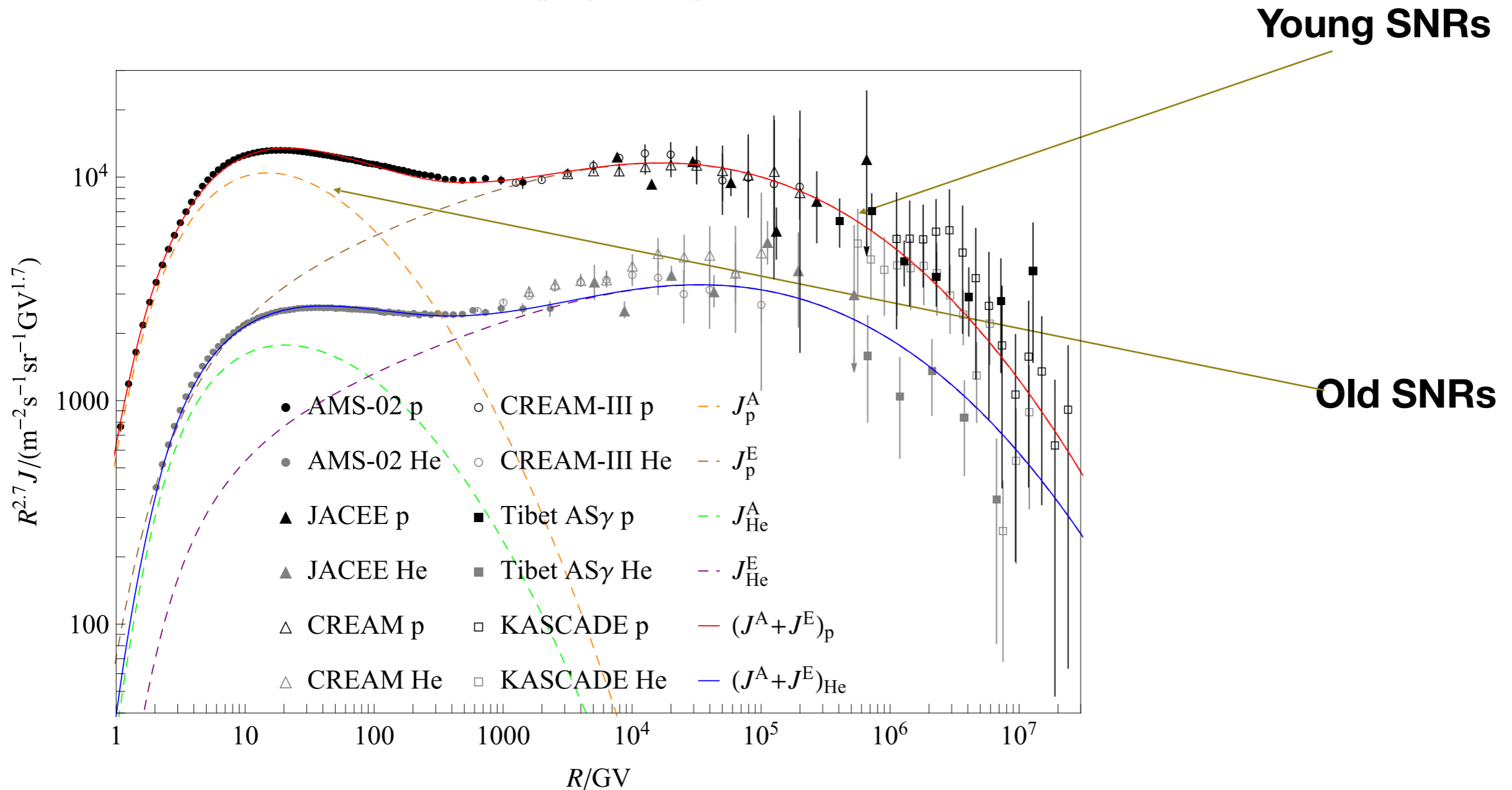


CrossMark

## Anomalous Distributions of Primary Cosmic Rays as Evidence for Time-dependent Particle Acceleration in Supernova Remnants

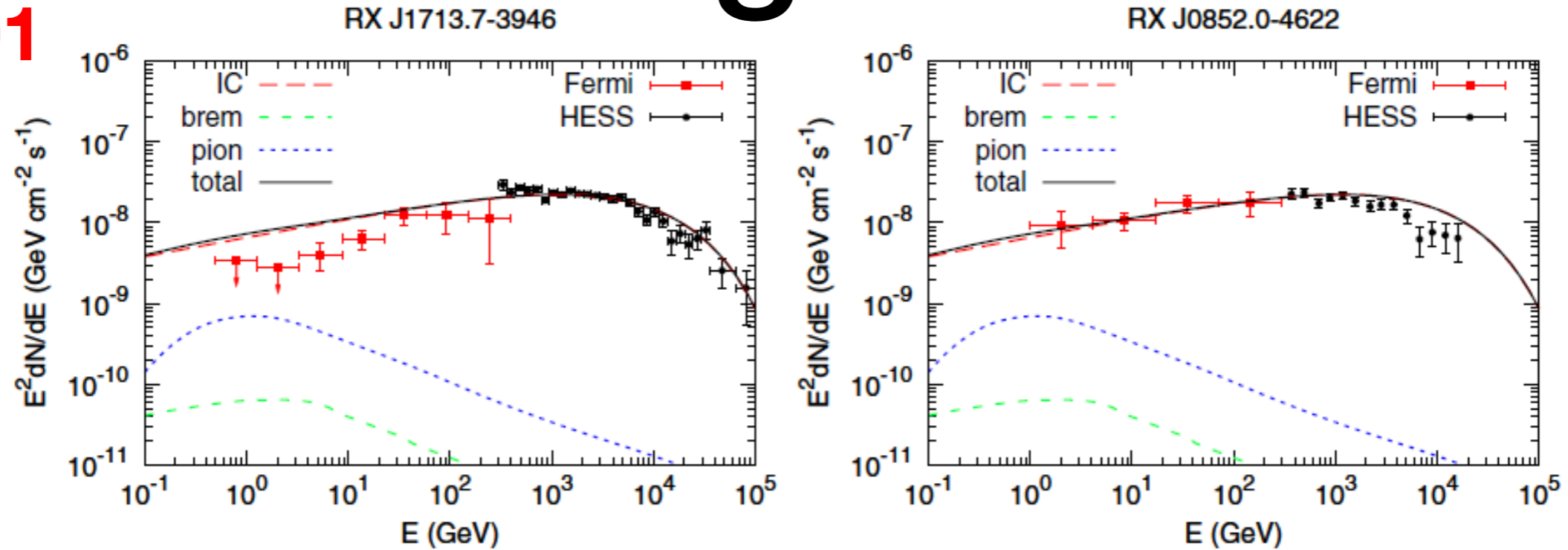
Yiran Zhang<sup>1,2</sup>, Siming Liu<sup>1,2</sup>, and Qiang Yuan<sup>1,2</sup>

<sup>1</sup>Key Laboratory of Dark Matter and Space Astronomy, Purple Mountain Observatory, Chinese Academy of Sciences, Nanjing 210008, China



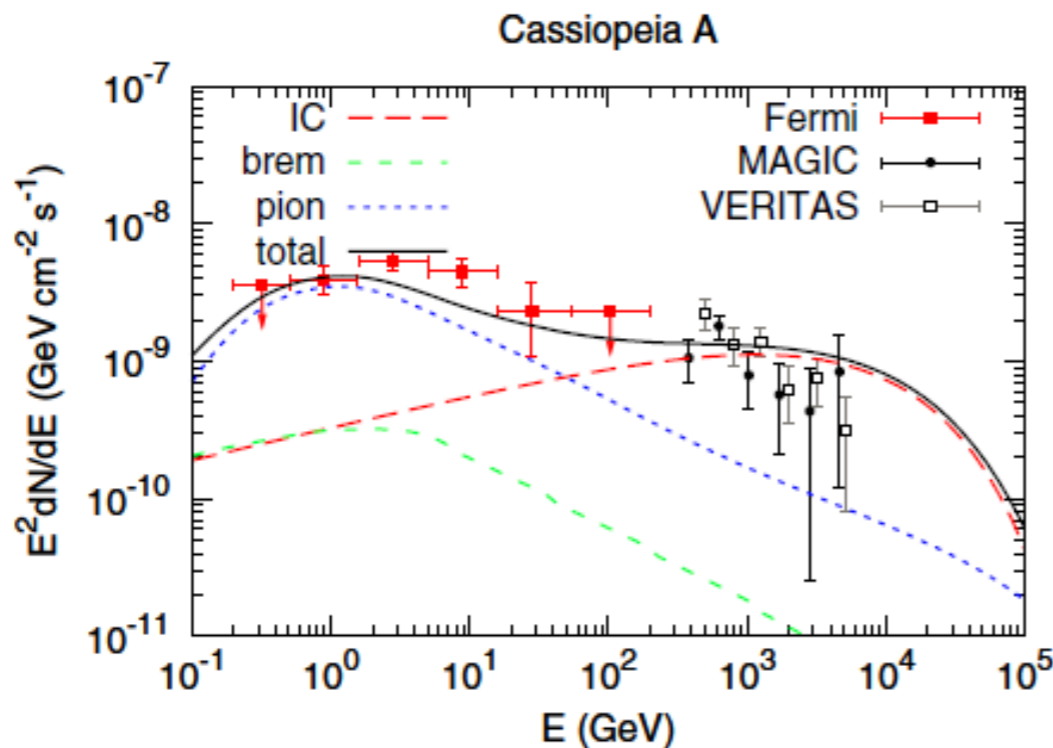
# Background

$n=0.01$

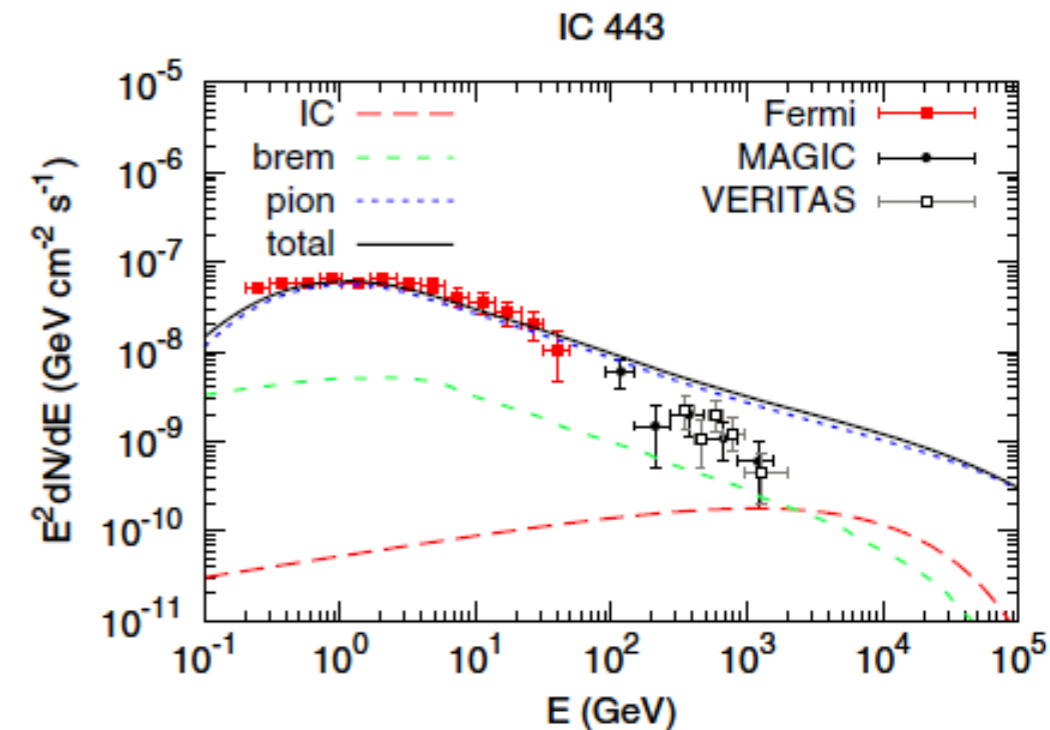


**Figure 2.** Expected  $\gamma$ -ray spectra for SNRs RX J1713.7-3946 (left) and RX J0852.0-4622 (right). The gas density is adopted to be  $n = 0.01 \text{ cm}^{-3}$ . References of the observational data—RX J1713.7-3946: *Fermi* (Abdo et al. 2011), HESS (Aharonian et al. 2007b); RX J0852.0-4622: *Fermi* (Tanaka et al. 2011), HESS (Aharonian et al. 2007a).

$n=1.0$



$n=100$



# Gamma-ray spectral evolution



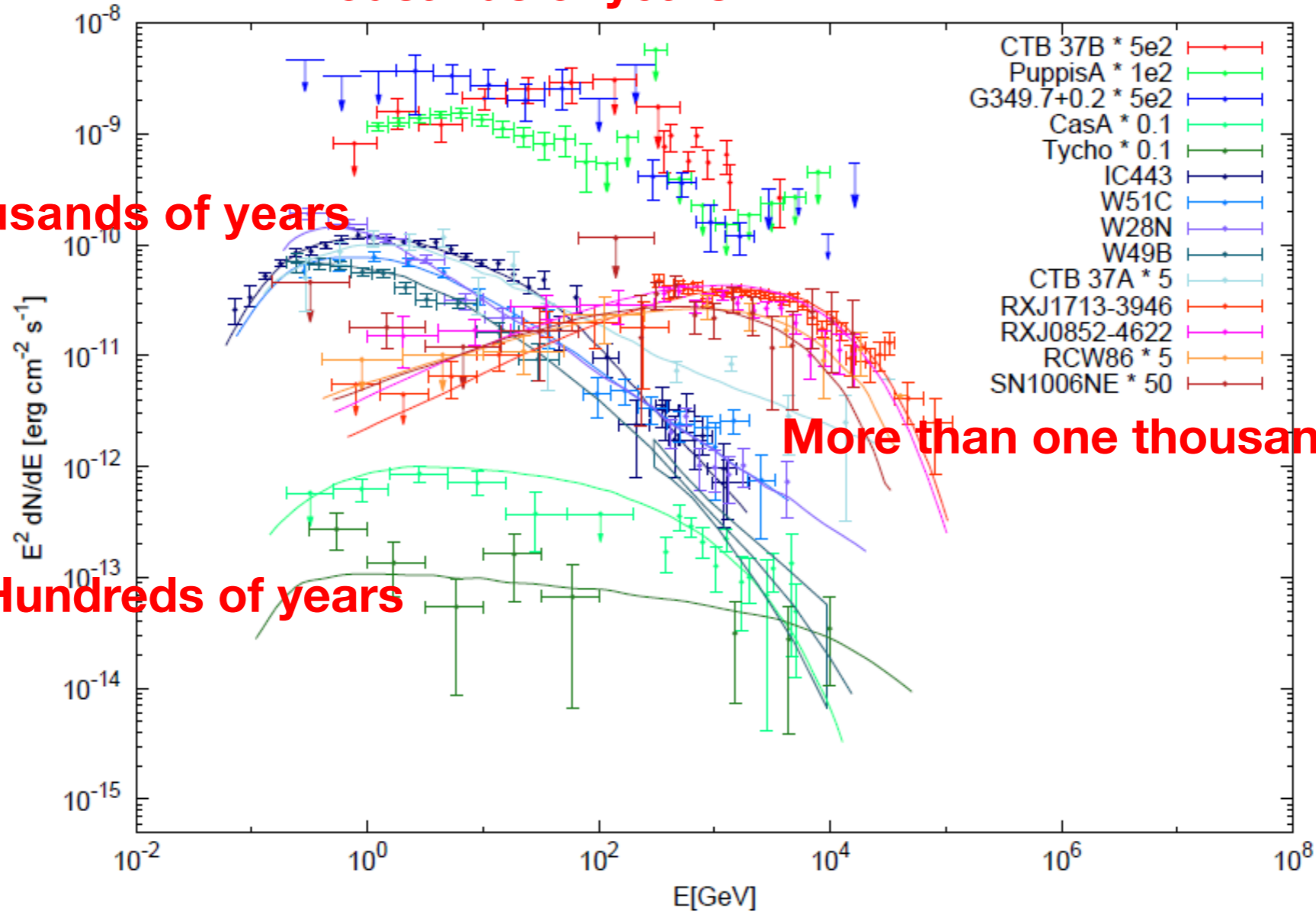
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Thousands of years

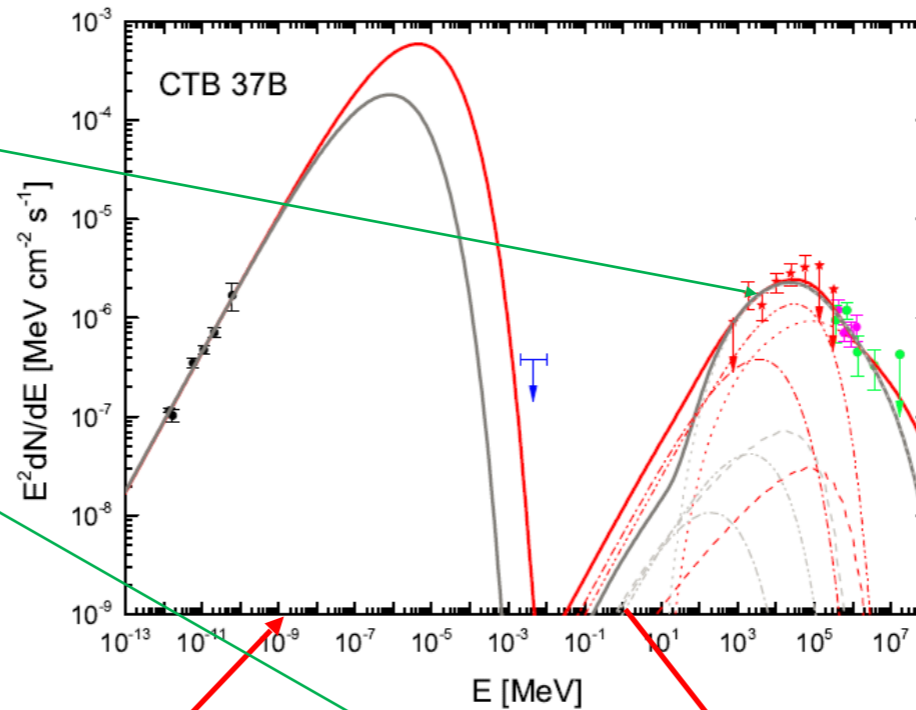
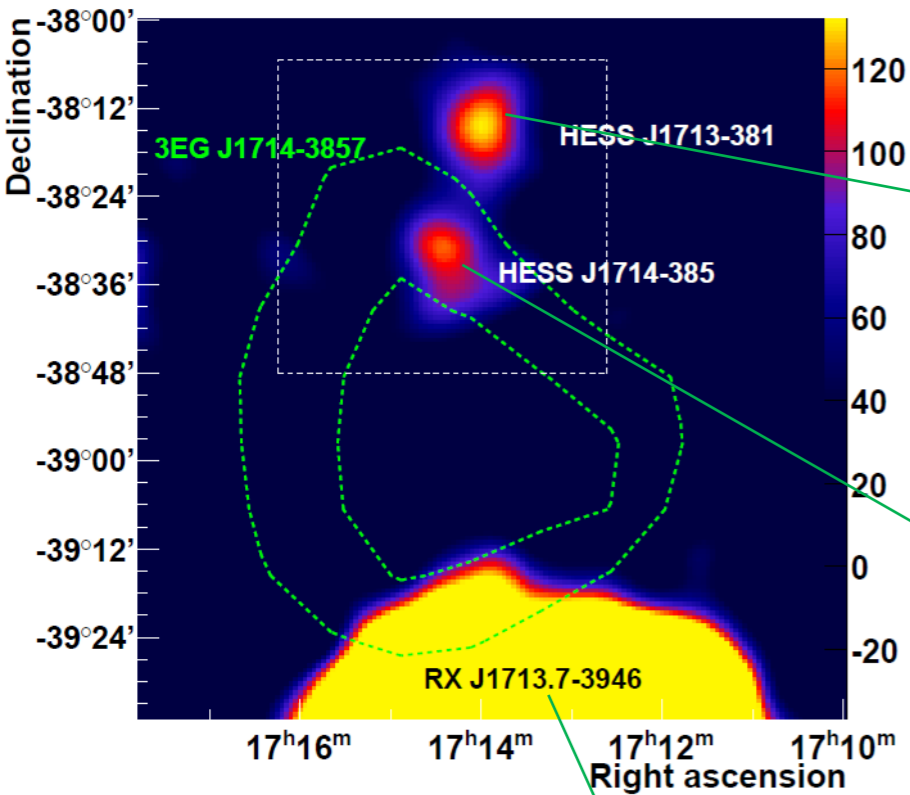
Tens of thousands of years

Hundreds of years

More than one thousand years





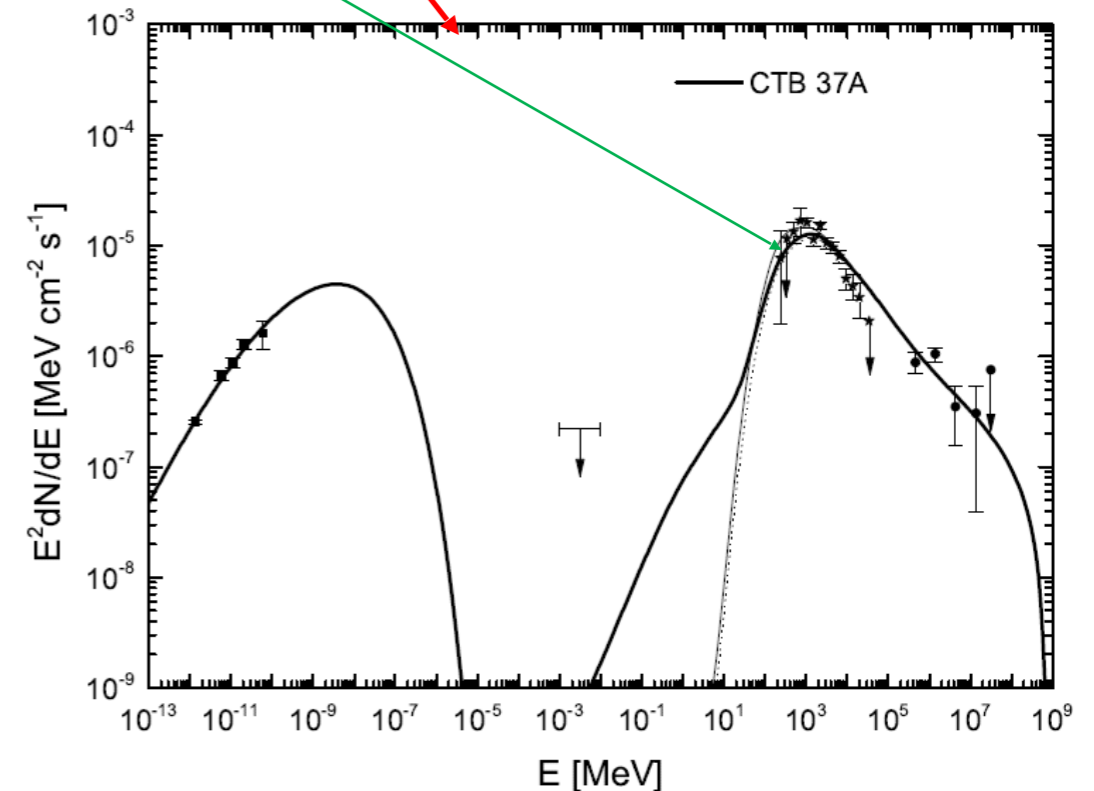
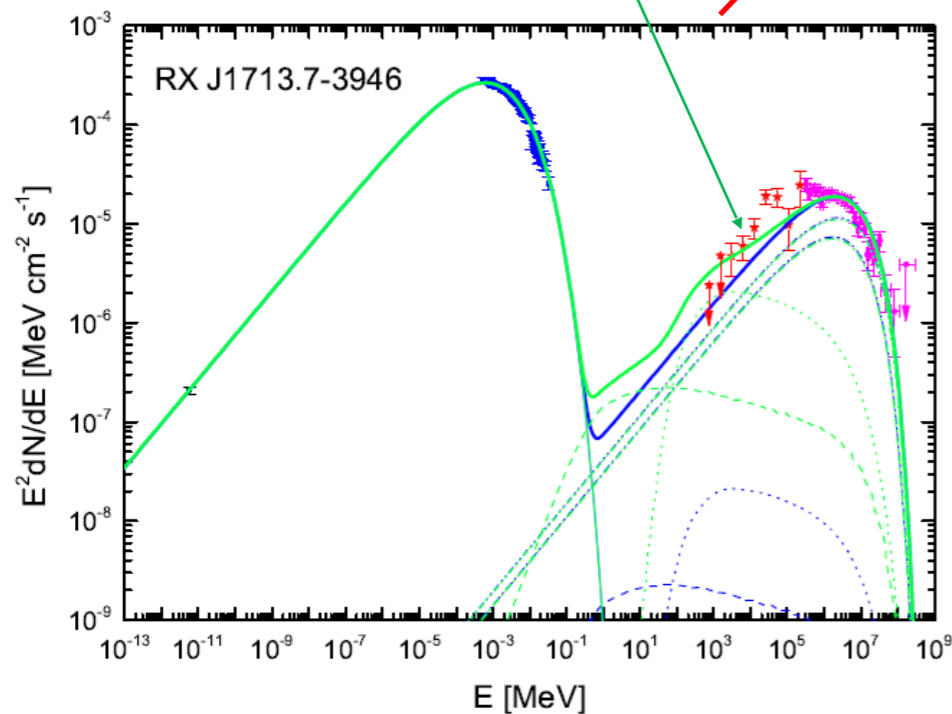


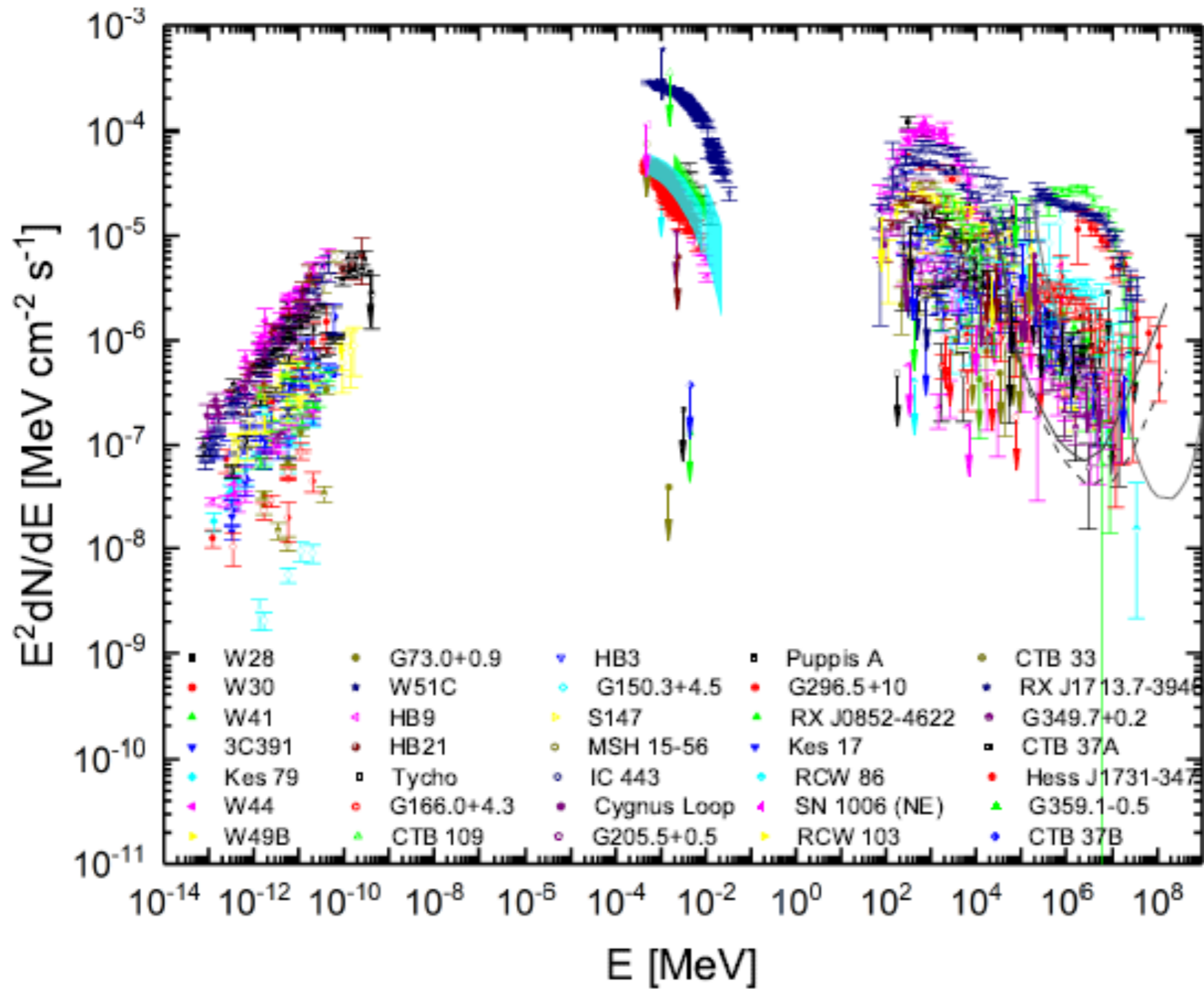
**CTB 37B:**  
 ~5000 Years  
**E<sub>br</sub>~400GeV**  
 Both IC and pp  
 contribute to  $\gamma$ -rays

**CTB 37A:**  
 >10000 Years  
**E<sub>br</sub>~2GeV**  
 pp dominates  $\gamma$ -rays

*The evolution  
 of SNRs' SEDs*

**RX J1713.7-3946**  
 1600 Years **E<sub>br</sub>~1TeV**  
 IC dominates  $\gamma$ -rays





**34 SNRs have been selected, 6 X-ray, 16 TeV data**

We normalize the flux at 100 GeV to  $1 \times 10^{-5} \text{ MeV cm}^{-2} \text{ s}^{-1}$  by a power-law fit to the spectrum from 1 GeV to 300 GeV to better demonstrate the spectral evolution.

### Break energy

#### 1-20GeV:

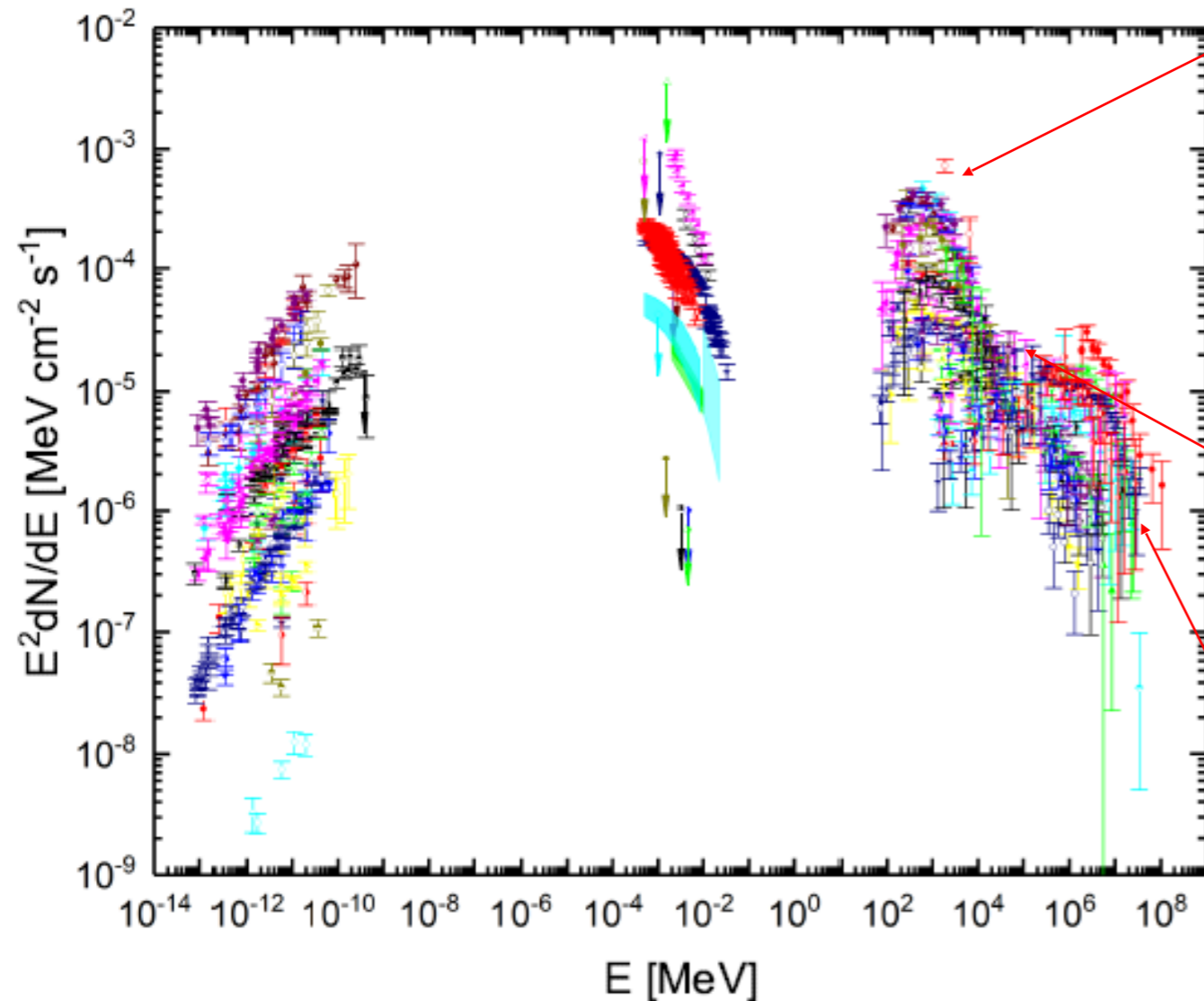
W28 (Abdo et al. 2010a),  
W51C (Abdo et al. 2009; Aleksic et al. 2012), W49B (Abdo et al. 2010c; Brun et al. 2011), IC 443 (Acciari et al. 2009; Ackermann et al. 2013)  
W41 (Aharonian et al. 2006b; Castro et al. 2013) and so.

#### 20-200GeV:

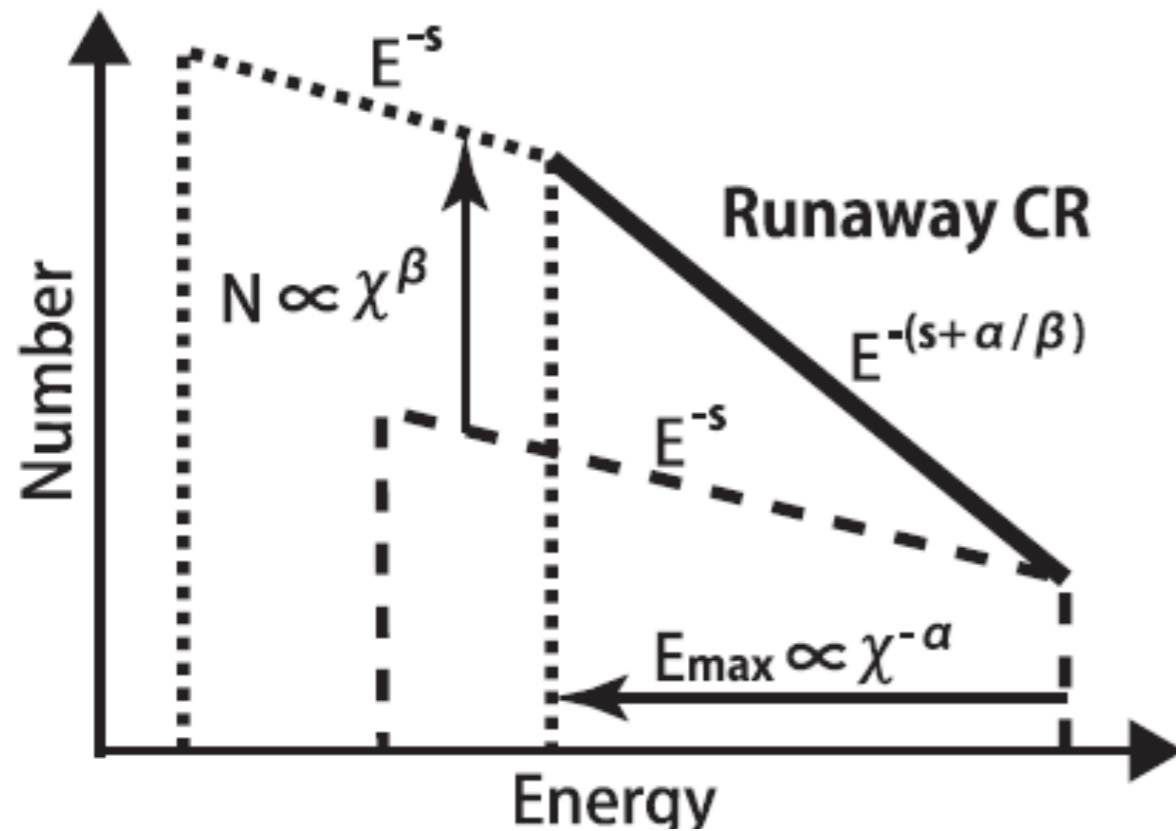
G349.7+0.2 (Hess Collaboration et al. 2015), CTB 37B (Xin et al. 2016),  
Puppis A (Xin et al. 2017)

#### 1TeV:

RX J1713.7-3946 (Hess Collaboration et al. 2017)



# A model for broken power-laws in SNRs



THE ASTROPHYSICAL JOURNAL LETTERS, 729:L13 (5pp), 2011 March 1  
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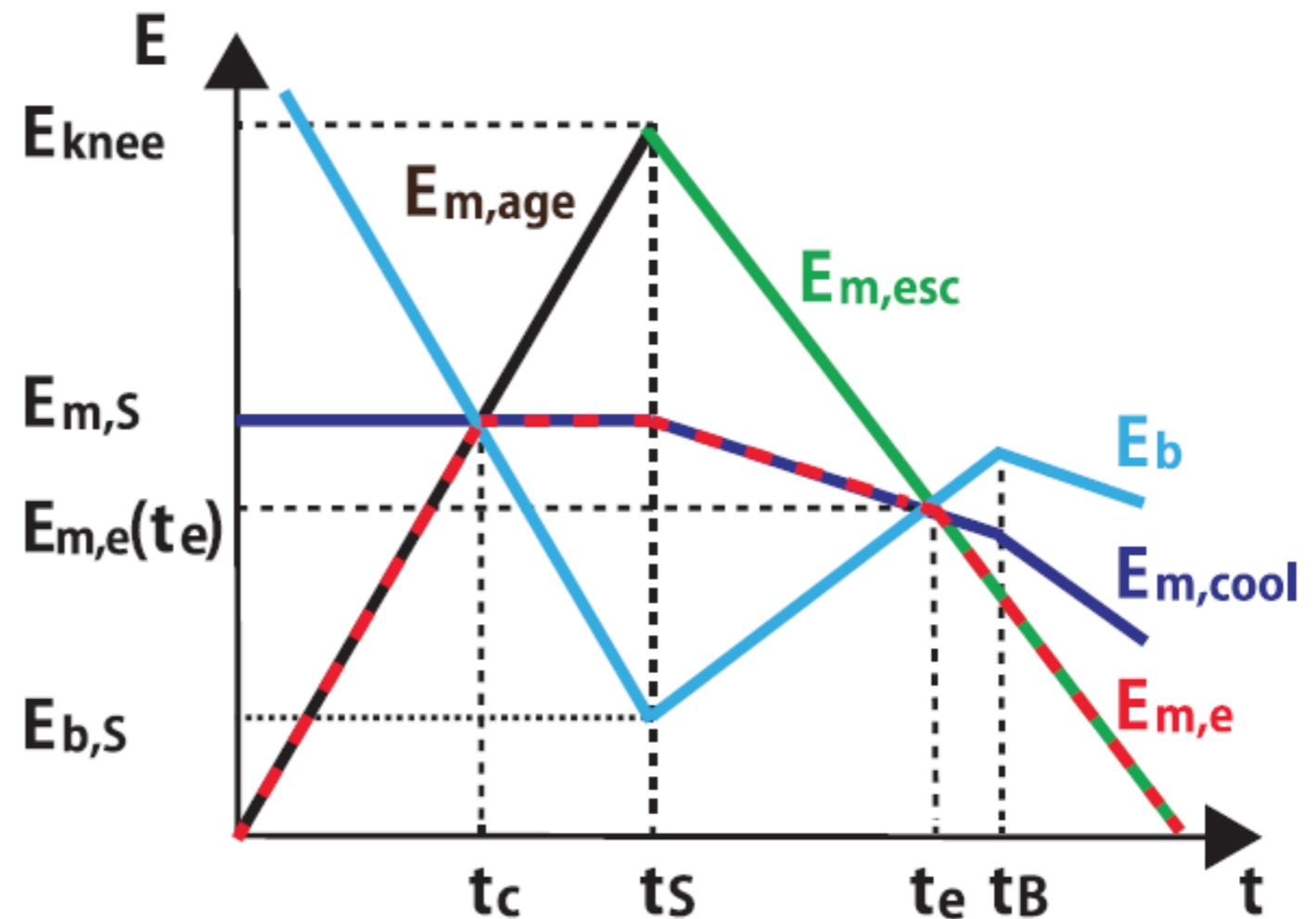
COSMIC-RAY HELIUM HARDENING  
 YUTAKA OHIRA AND KUNIHITO IOKA

Mon. Not. R. Astron. Soc. 427, 91–102 (2012)

doi:10.1111/j.

## Escape of cosmic-ray electrons from supernova remnants

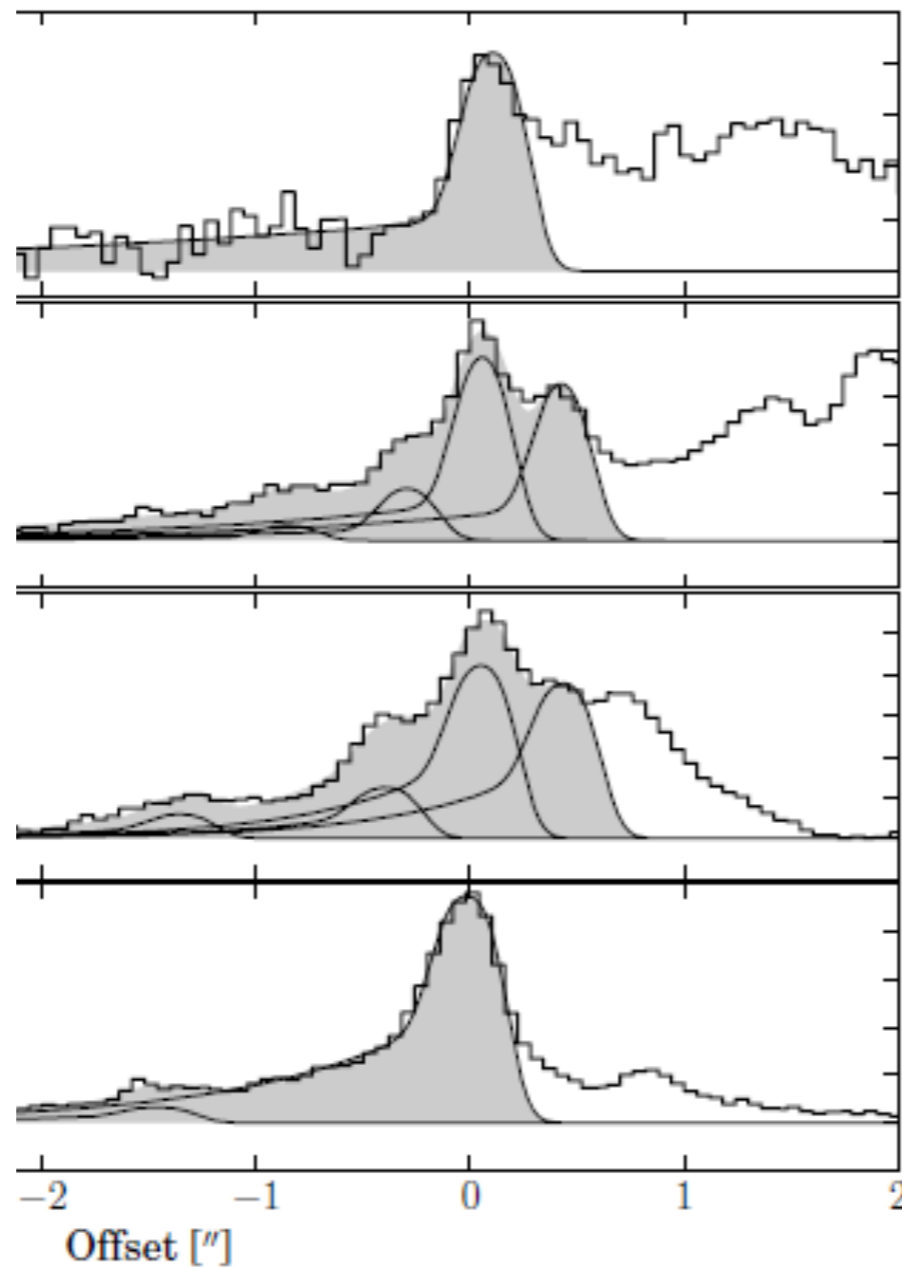
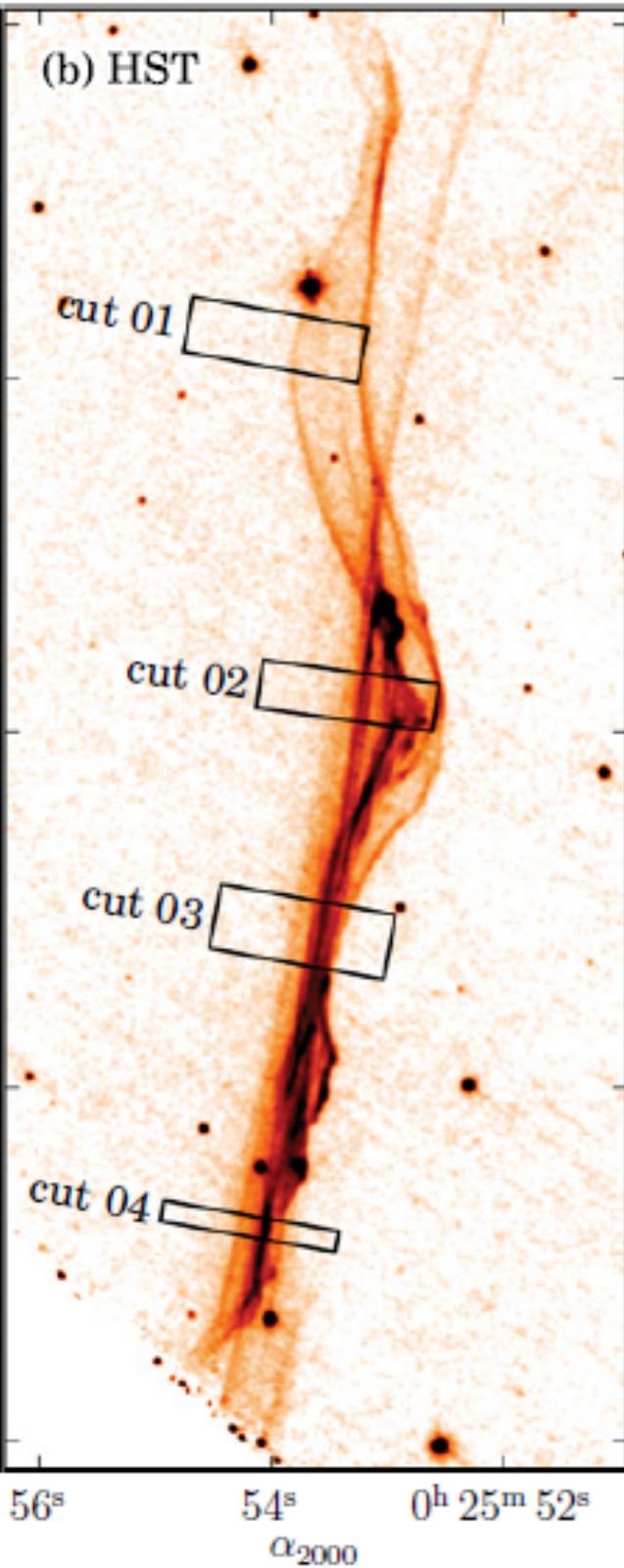
Yutaka Ohira,<sup>1\*</sup> Ryo Yamazaki,<sup>1</sup> Norita Kawanaka<sup>2</sup> and Kunihito Ioka<sup>3,4</sup>



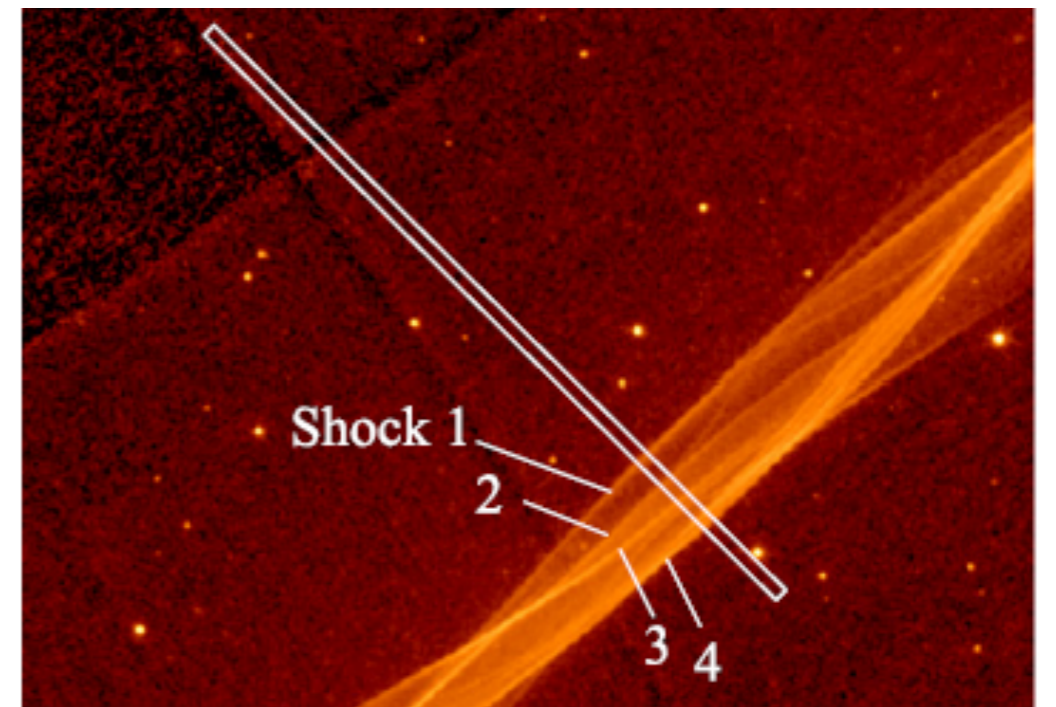
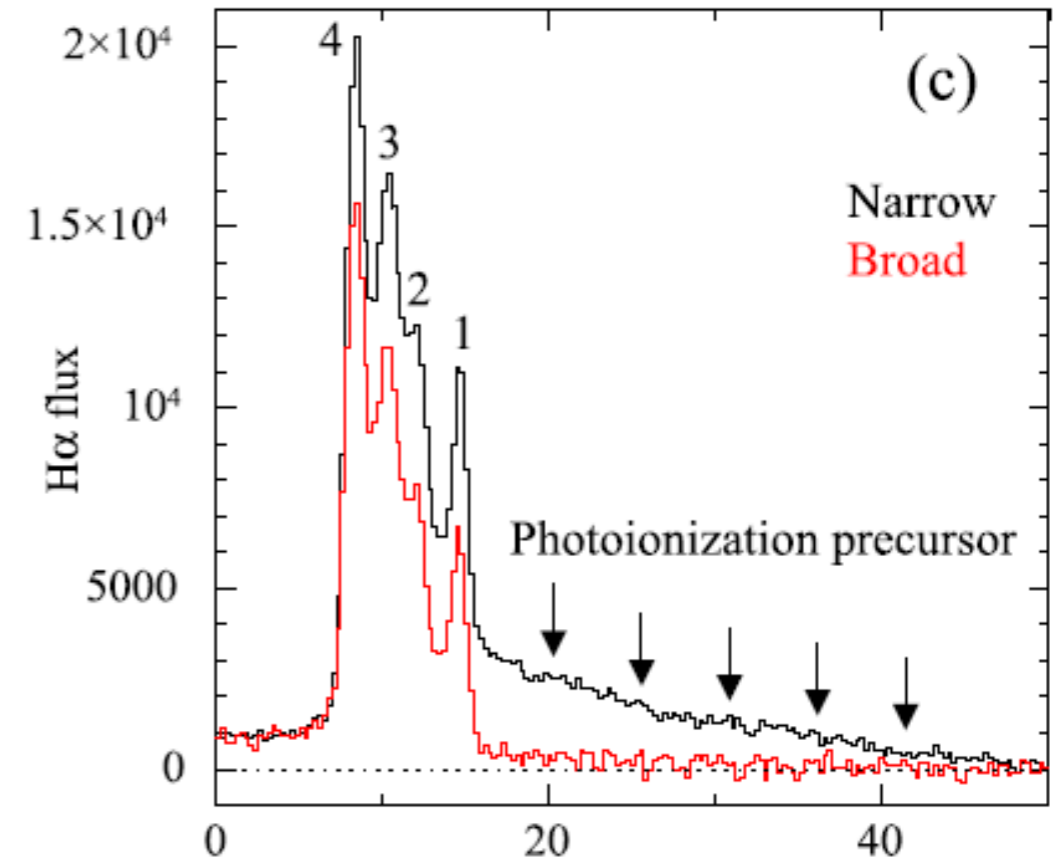
# H $\alpha$ observations

Tycho

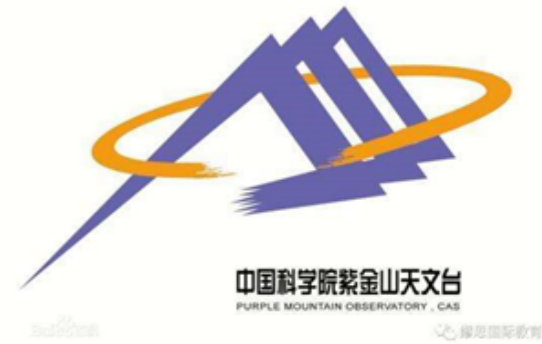
Young with wider precursor for higher energy acceleration



Cygnus loop old



# One-zone model



1. The distribution of particles: "i" represents e or p

$$N(P_i) = N_{0,i} \exp\left(-\frac{P_i}{P_{i,cut}}\right) \begin{cases} P_i^{-\alpha} & \text{if } P_i < P_{br} \\ P_{i,br} P_i^{-(\alpha+1)} & \text{if } P_i \geq P_{br} , \end{cases}$$

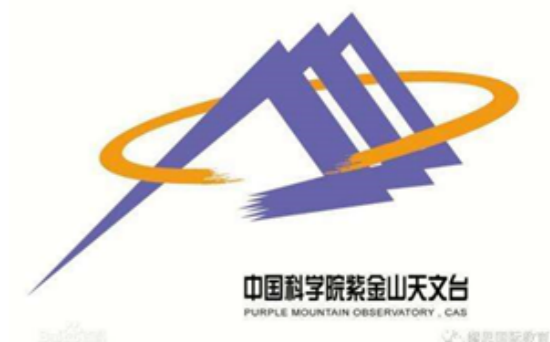
2.  $K_{ep}=0.01$ , and both leptonic (Synch, Brem and IC) and hadronic (pp) emissions are considered in our fitting.
3. For comparison, the same background photon field is assumed for all SNRs (CMB and IR) .
4. Distance, Age, Shock velocity, The gas density from literature

$$\frac{P_{e,cut}}{\text{GeV}/c} = 1.25 \times 10^6 \left(\frac{T_{age}}{\text{Year}}\right)^{-1} \left(\frac{B}{100\mu\text{G}}\right)^{-2}$$

5-6 free parameters

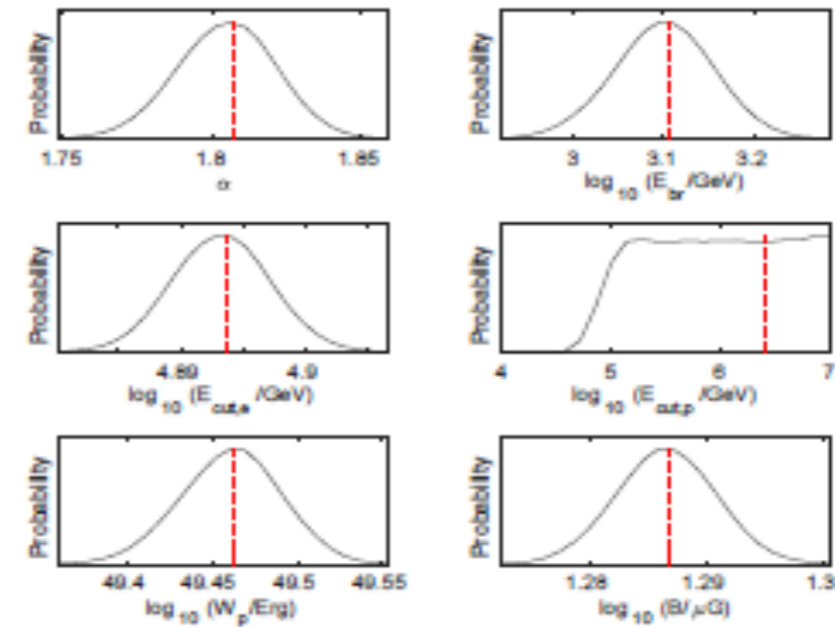
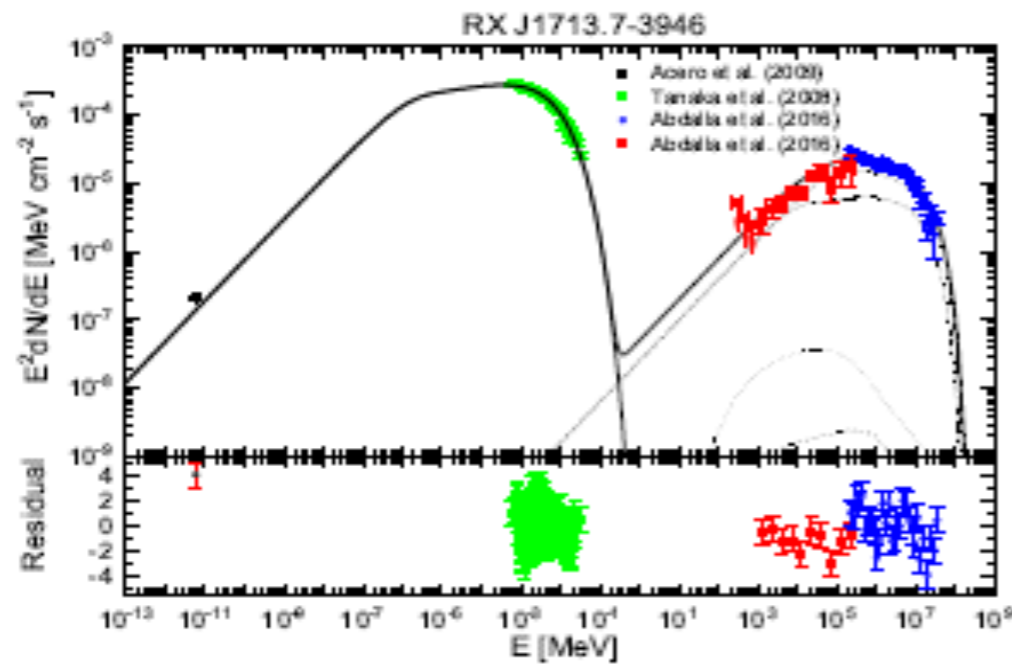
**MCMC method** is applied to constrain these model parameters.

# The Sample

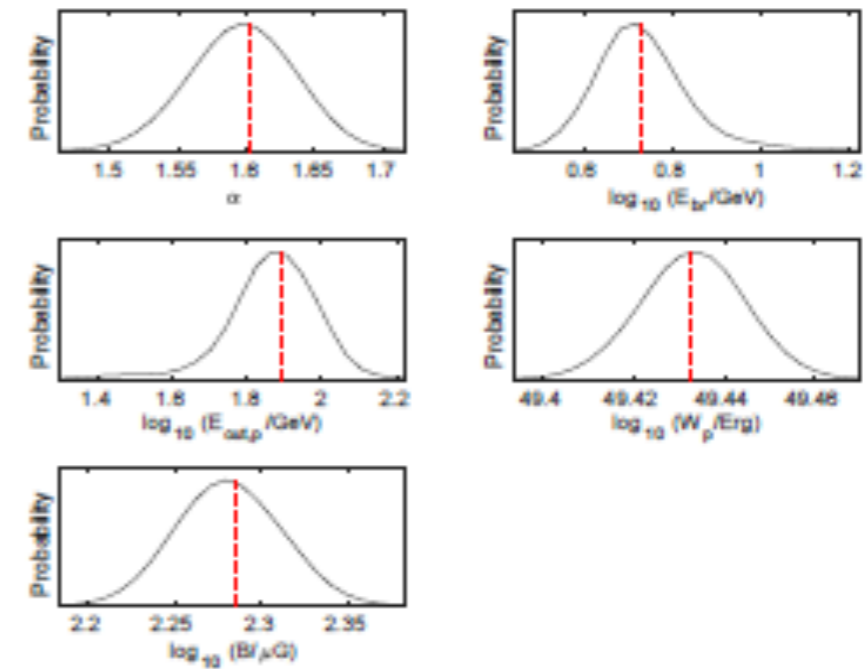
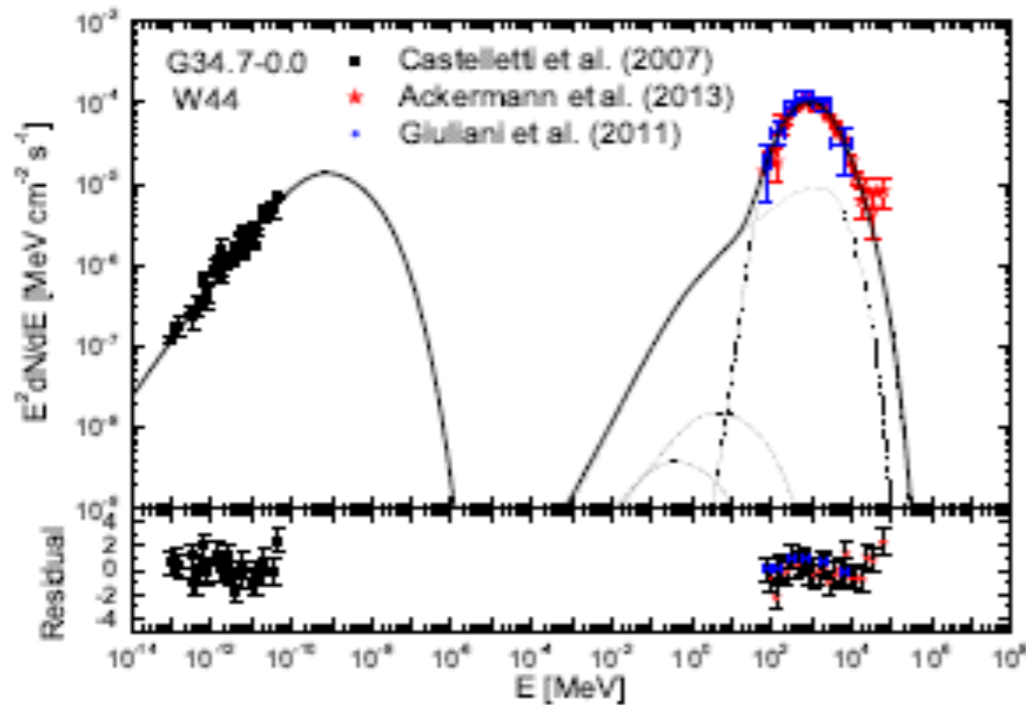


SNR Name	Other Name	Radius pc	Distance Kpc	Age kyr	n $\text{cm}^{-3}$	$V_{shock}$ Km/s	The data of observations				References
							Radio	X-ray	GeV	TeV	
C006.4-00.1	W28	~ 13	~ 2.0	40(33-150)	~ 100	60-80	✓		✓	✓	[1-4]
C008.7-00.1	W30	~ 26	~ 4.0	25(15-28)	~ 100	530-750	✓		✓	✓	[5][6]
C023.3-00.3	W41	~ 19	~ 4.2	~ 100	~ 10	110	✓	T	✓	✓	[7][8]
C031.9-00.0	3C 391	~ 7	~ 7.2	~ 4	~ 300	620-730	✓		✓	✓	[9-12]
C033.6+00.1	Kes 79	~ 9.6	~ 7.0	~ 4.4-6.7	~ 3(1-5)	400 ± 5	✓		✓	✓	[13-15]
C034.7-00.0	W44	~ 12.5	~ 3.0	~ 20	~ 200	100-150	✓		✓	✓	[16-18]
C043.3-00.2	W49B	~ 5	~ 10	~ 5.7(5-6)	~ 700	~ 400	✓		✓	✓	[19][20]
C049.2-00.7	W51C	~ 18	~ 4.3	~ 30	~ 10	~ 100	✓	T	✓	✓	[21-24]
C073.9+00.9		~ 16/5.2	~ 4.0/1.3	~ 11-12	~ 10	~ 200-300	✓	T	✓	✓	[25][26]
C074.0-08.5	Cygnus loop	~ 16	~ 0.54	~ 14	~ 5.0	240-330	✓		✓	✓	[27-30]
C089.0-04.7	HB21	~ 26	~ 1.7	~ 40(36 or 45)	~ 15	~ 125	✓		✓	✓	[31-35]
C109.1-0.1	CTB109	~ 16	~ 3.1	~ 9.0(9.0-9.2)	~ 1.1	~ 230 ± 5	✓	T	✓	✓	[36][37]
C120.1-01.4	Tycho	~ 3.3	~ 3.0	~ 0.44	~ 10/0.3	4600-4800	✓	✓	✓	✓	[38][39]
C132.7-00.3	HB3	~ 26.4	~ 2.2	~ 30.0	~ 2.0	303-377	✓		✓	✓	[40-42]
C150.3+04.5		~ 9.4	~ 0.40	~ 1.5(0.5-5)	~ 1.0	?	✓		✓	✓	[43]
C160.9-02.6	HB9	~ 15	~ 0.8	5.3(4-7)	~ 0.1	~ 740	✓	T	✓	✓	[44][45]
C166.0+04.3		~ 26	~ 4.5	24.0	~ 0.01	~ 680	✓		✓	✓	[46][47]
C180.0+01.7	S147	~ 38	~ 1.3	30(20-10)	~ 250(100-500)	~ 500	✓		✓	✓	[48][49]
C189.1-03.0	IC 443	~ 11	~ 1.5	~ 30	~ 140	60-100	✓		✓	✓	[50-52]
C205.5+0.5	Monoceros	~ 63.36	~ 1.98	~ 30	~ 3.6	~ 50	✓		✓	✓	[53-55]
C260.4-03.4	Puppis A	~ 15	~ 2.2	4.45(3.75-5.20)	~ 4.0	700-2500	✓		✓	T	[56-59]
C266.2-01.2	RX J0852-4622	~ 13	~ 0.75	2.7(1.7-4.3)	~ 3.8	~ 3000	✓	✓	✓	✓	[60][61]
C296.5+10.0		~ 26	~ 2.1	~ 10.0	~ 13.0	~ >35	✓		✓	✓	[62][63]
C304.6-00.1	Kes 17	~ 10	~ 10	4.2(2-5.2)	~ 10	150-200	✓		✓	✓	[64][65]
C315.4-02.3	RCW 86	~ 15	~ 2.5	~ 1.8	~ 0.1-2.0	700-2000	✓	✓	✓	✓	[66-68]
C326.3-01.8	MSH 15-56	~ 22.2	~ 4.1	~ 10.0(10-16.5)	~ 0.1/1.0	500-860	✓		✓	✓	[67][69][70]
C327.6+14.6	SN 1006 (NE)	~ 9.0	~ 2.2	~ 1.0	~ 0.085	3200-5800	✓	✓	✓	✓	[71][72]
C332.4-00.4	RCW 103	~ 5	~ 3.3	~ 2.0	~ 10	~ 1100	✓		✓	✓	[73-75]
C337.0-00.1	CTB 33	~ 2.55	~ 11.0	~ 5.0	~ 60	?	✓		✓	✓	[76-78]
C347.3-00.5	RX 1713.7-3946	~ 10	~ 1.0	~ 1.6	~ 0.01	~ 5000	✓	✓	✓	✓	[82-84]
C348.5+00.1	CTB 37A	~ 10	~ 7.9	~ 30	~ 100	75-100	✓	T	✓	✓	[85-89]
C348.7+00.3	CTB 37B	~ 20	~ 13.2	~ 5	~ 10/0.5	~ 800	✓	T	✓	✓	[88-91]
C349.7+00.2		~ 3.3	~ 11.5	~ 2.8	~ 35.0	700-900	✓	T	✓	✓	[92-95]
C353.6-00.7	Hess J1731-347	~ 14.0	~ 3.2	~ 2-6	~ 0.01	~ 2100	✓	✓	✓	✓	[96][97]
C359.1-00.5	Hess J1745-303	~ 16.0	~ 4.6	~ 70	~ 100	~ 300	✓	T	✓	✓	[98-101]

# Examples of spectral fits

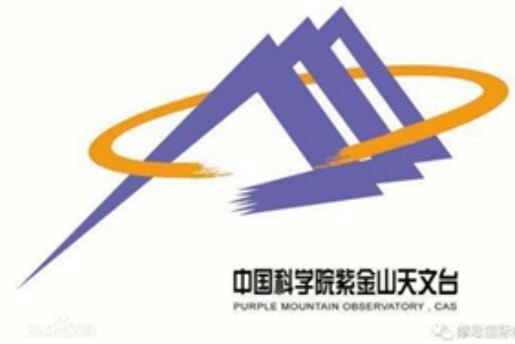


RX J1713.7-3946



W44





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Source Name	$\alpha$	$\log_{10} \frac{E_{br}}{\text{GeV}}$	$\log_{10} \frac{E_{e,cut}}{\text{GeV}}$	$\log_{10} \frac{E_{p,cut}}{\text{GeV}}$	$\log_{10} \frac{R}{\mu\text{G}}$	$\log_{10} \frac{W_p}{\text{erg}}$	$\frac{W_B}{W_e}$	$\frac{n}{\text{cm}^{-3}}$	$\frac{\chi^2}{\text{NDF}}$
W28	$1.76^{+0.03}_{-0.03}$	$0.18^{+0.11}_{-0.11}$	1.63	> 5.72	$1.94^{+0.04}_{-0.04}$	$49.36^{+0.02}_{-0.02}$	$3.0 \times 10^3$	100	$\frac{24.9}{10} = 2.49$
W30	$1.63^{+0.10}_{-0.11}$	$0.24^{+0.33}_{-0.38}$	2.06	> 4.29	$1.86^{+0.12}_{-0.13}$	$49.69^{+0.07}_{-0.07}$	736	100	$\frac{6.0}{7} = 0.86$
W41	$1.22^{+0.04}_{-0.04}$	< -0.30	1.97	$4.52^{+0.20}_{-0.20}$	$1.87^{+0.11}_{-0.11}$	$50.21^{+0.05}_{-0.05}$	464	10	$\frac{12.6}{12} = 1.05$
3C391	$1.99^{+0.05}_{-0.05}$	$1.15^{+0.14}_{-0.14}$	1.86	> 3.81	$2.31^{+0.04}_{-0.04}$	$49.03^{+0.03}_{-0.03}$	619	300	$\frac{37.1}{19} = 1.95$
Kes79	$2.11^{+0.05}_{-0.05}$	$0.71^{+0.15}_{-0.15}$	2.89	> 4.32	$1.74^{+0.04}_{-0.04}$	$49.47^{+0.04}_{-0.04}$	28.9	100.0	$\frac{66}{23} = 2.54$
W44	$1.60^{+0.04}_{-0.04}$	$0.73^{+0.09}_{-0.09}$	1.23	$1.87^{+0.11}_{-0.09}$	$2.28^{+0.03}_{-0.03}$	$49.43^{+0.01}_{-0.01}$	$1.48 \times 10^3$	200	$\frac{44.4}{48} = 0.93$
W49B	$1.47^{+0.04}_{-0.04}$	$-0.21^{+0.23}_{-0.24}$	1.55	$3.70^{+0.13}_{-0.13}$	$2.40^{+0.06}_{-0.06}$	$49.43^{+0.02}_{-0.02}$	235	700	$\frac{18.9}{20} = 1.00$
W51C	$1.56^{+0.02}_{-0.02}$	$0.31^{+0.08}_{-0.08}$	1.64	$4.39^{+0.30}_{-0.29}$	$2.08^{+0.03}_{-0.03}$	$49.83^{+0.01}_{-0.01}$	708	100	$\frac{59.7}{29} = 2.06$
W51C <sup>b</sup>	$1.64^{+0.02}_{-0.02}$	$0.32^{+0.05}_{-0.05}$	1.57	> 5.78	$2.02^{+0.03}_{-0.03}$	$49.79^{+0.01}_{-0.01}$	201	100	$\frac{34.9}{29} = 1.20$
G73.9+0.9 <sup>c</sup>	$0.78^{+0.19}_{-0.19}$	$0.96^{+0.09}_{-0.09}$	$0.96^{+0.09}_{-0.09}$	$0.96^{+0.09}_{-0.09}$	$1.57^{+0.05}_{-0.05}$	$49.34^{+0.04}_{-0.04}$	393	10	$\frac{22.4}{13} = 1.72$
Cygnus Loop	$2.01^{+0.05}_{-0.05}$	$0.69^{+0.09}_{-0.09}$	2.85	> 4.09	$1.47^{+0.02}_{-0.02}$	$48.72^{+0.02}_{-0.02}$	232	5.0	$\frac{21.9}{19} = 1.15$
HB21	$1.21^{+0.10}_{-0.10}$	$0.61^{+0.08}_{-0.08}$	$0.73^{+0.07}_{-0.07}$	$0.77^{+0.07}_{-0.06}$	$1.74^{+0.02}_{-0.02}$	$49.42^{+0.01}_{-0.01}$	562	15	$\frac{36.4}{21} = 1.73$
CTB109	$1.94^{+0.09}_{-0.09}$	$2.66^{+0.38}_{-0.81}$	3.28	> 4.82	$1.47^{+0.17}_{-0.20}$	$49.84^{+0.12}_{-0.12}$	19.6	1.1	$\frac{20.9}{8} = 2.61$
Tycho	$2.15^{+0.02}_{-0.02}$	$3.37^{+0.12}_{-0.12}$	$4.14^{+0.08}_{-0.09}$	> 5.04	$2.15^{+0.04}_{-0.05}$	$49.01^{+0.08}_{-0.08}$	23.5	0.3	$\frac{55}{35} = 1.57$
Tycho	$2.16^{+0.02}_{-0.02}$	$3.36^{+0.11}_{-0.11}$	$4.06^{+0.07}_{-0.07}$	> 4.93	$2.29^{+0.04}_{-0.05}$	$48.78^{+0.07}_{-0.07}$	92.2	10.0	$\frac{44}{35} = 1.26$
HB3	$2.07^{+0.10}_{-0.10}$	$0.76^{+0.14}_{-0.14}$	3.50	> 4.57	$1.07^{+0.03}_{-0.03}$	$50.03^{+0.03}_{-0.03}$	7.49	2.0	$\frac{16.8}{18} = 0.93$
G150.3+4.5	$1.73^{+0.22}_{-0.22}$	$2.65^{+0.36}_{-0.42}$	6.04	> 6.37	$0.45^{+0.13}_{-0.13}$	$48.33^{+0.05}_{-0.05}$	1.42	1.0	$\frac{11.5}{12} = 0.96$
HB9	$2.23^{+0.06}_{-0.06}$	$0.89^{+0.12}_{-0.12}$	5.07	> 5.68	$0.67^{+0.04}_{-0.04}$	$50.10^{+0.05}_{-0.05}$	0.16	0.1	$\frac{15.2}{14} = 1.09$
G166.0+4.3 <sup>c</sup>	$1.32^{+0.17}_{-0.18}$	1.87	1.87	$1.87^{+0.14}_{-0.15}$	$0.57^{+0.24}_{-0.24}$	$50.92^{+0.26}_{-0.25}$	0.12	0.01	$\frac{6.92}{5} = 1.38$
G166.0+4.3 <sup>c</sup>	$1.26^{+0.17}_{-0.18}$	1.18	1.18	$1.18^{+0.16}_{-0.16}$	$1.62^{+0.10}_{-0.10}$	$49.18^{+0.07}_{-0.07}$	717	10.0	$\frac{7.70}{5} = 1.54$
S147	$1.36^{+0.06}_{-0.06}$	$-0.14^{+0.12}_{-0.13}$	0.09	> 3.86	$2.77^{+0.09}_{-0.09}$	$47.71^{+0.05}_{-0.05}$	$2.7 \times 10^8$	250	$\frac{17.3}{17} = 1.02$
S147	$1.53^{+0.11}_{-0.11}$	$0.51^{+0.12}_{-0.12}$	3.57	> 4.65	$1.03^{+0.05}_{-0.05}$	$49.94^{+0.04}_{-0.04}$	31.6	1.0	$\frac{19.8}{17} = 1.16$
IC 443	$1.38^{+0.03}_{-0.03}$	$0.12^{+0.07}_{-0.07}$	1.35	$3.22^{+0.10}_{-0.10}$	$2.14^{+0.02}_{-0.02}$	$48.96^{+0.01}_{-0.01}$	$2.28 \times 10^3$	140	$\frac{92.0}{64} = 1.44$
Monoceros Loop	$1.63^{+0.02}_{-0.02}$	$0.74^{+0.11}_{-0.11}$	2.97	> 5.77	$1.31^{+0.03}_{-0.03}$	$50.29^{+0.03}_{-0.03}$	224	3.6	$\frac{42.5}{16} = 1.63$
Puppis A	$2.08^{+0.02}_{-0.02}$	$3.23^{+0.48}_{-0.56}$	2.50	> 4.57	$1.97^{+0.02}_{-0.02}$	$49.53^{+0.04}_{-0.04}$	500	4.0	$\frac{48.8}{30} = 1.46$
RX J0852-4622 <sup>d</sup>	$2.21^{+0.04}_{-0.04}$		$4.30^{+0.06}_{-0.06}$	> 5.15	$1.03^{+0.04}_{-0.04}$	$49.61^{+0.05}_{-0.05}$	2.79	0.01	$\frac{27.8}{16} = 1.74$
RX J0852-4622	$1.33^{+0.05}_{-0.05}$	$1.13^{+0.18}_{-0.16}$	$4.38^{+0.06}_{-0.06}$	> 5.15	$1.04^{+0.04}_{-0.04}$	$49.70^{+0.04}_{-0.04}$	2.6	0.01	$\frac{18.6}{15} = 1.24$
G296.5+10.0	$1.86^{+0.08}_{-0.08}$	> 3.75	0.59	> 3.99	$2.73^{+0.13}_{-0.13}$	$48.55^{+0.14}_{-0.14}$	$1.15 \times 10^7$	13.0	$\frac{4.3}{5} = 0.86$
Kes 17	$2.01^{+0.18}_{-0.17}$	> 3.52	3.03	> 4.20	$1.79^{+0.17}_{-0.17}$	$50.33^{+0.11}_{-0.11}$	7.0	10.0	$\frac{1.06}{2} = 0.53$
RCW 86	$2.26^{+0.02}_{-0.02}$	$3.92^{+0.08}_{-0.09}$	$4.42^{+0.04}_{-0.03}$	> 5.23	$1.44^{+0.02}_{-0.02}$	$49.82^{+0.03}_{-0.03}$	15.3	0.01	$\frac{31.5}{22} = 1.43$
MSH 15-56	$1.43^{+0.14}_{-0.14}$	$2.13^{+0.17}_{-0.20}$	2.40	> 3.06	$1.81^{+0.09}_{-0.08}$	$51.05^{+0.13}_{-0.13}$	34.3	0.1	$\frac{10.1}{7} = 1.44$
MSH 15-56	$1.61^{+0.12}_{-0.12}$	$1.17^{+0.08}_{-0.11}$	3.02	> 3.83	$1.60^{+0.10}_{-0.10}$	$50.75^{+0.06}_{-0.06}$	9.7	1.0	$\frac{8.0}{7} = 1.14$
SN 1006	$2.09^{+0.04}_{-0.04}$	> 4.84	$3.86^{+0.10}_{-0.10}$	> 4.92	$1.77^{+0.05}_{-0.06}$	$48.72^{+0.07}_{-0.07}$	240	0.085	$\frac{49.0}{35} = 1.40$
RCW 103	$2.11^{+0.08}_{-0.08}$	> 3.86	3.90	> 4.87	$1.45^{+0.08}_{-0.08}$	$50.00^{+0.06}_{-0.06}$	0.44	10	$\frac{1.0}{6} = 0.17$
CTB 33	$2.00^{+0.38}_{-0.33}$	$1.41^{+0.56}_{-0.55}$	4.67	> 5.43	$0.87^{+0.08}_{-0.08}$	$50.47^{+0.07}_{-0.07}$	0.001	60	$\frac{9.06}{5} = 1.81$
RX J1713.7-3946 <sup>a</sup>	$1.81^{+0.02}_{-0.02}$	$3.10^{+0.05}_{-0.05}$	$4.89^{+0.004}_{-0.004}$	> 5.57	$1.29^{+0.004}_{-0.004}$	$49.46^{+0.03}_{-0.03}$	6.0	0.01	$\frac{445}{240} = 1.85$
CTB 37A	$1.47^{+0.02}_{-0.02}$	$0.36^{+0.19}_{-0.17}$	1.0	> 5.96	$2.40^{+0.12}_{-0.10}$	$49.82^{+0.02}_{-0.02}$	607	100	$\frac{23.4}{16} = 1.46$
CTB 37B	$1.49^{+0.11}_{-0.11}$	$2.40^{+0.33}_{-0.34}$	0.81	> 5.34	$2.84^{+0.15}_{-0.15}$	$50.51^{+0.04}_{-0.04}$	$1.04 \times 10^5$	10	$\frac{15.6}{14} = 1.11$
CTB 37B	$1.58^{+0.07}_{-0.07}$	$3.06^{+0.19}_{-0.20}$	2.47	> 5.32	$1.97^{+0.06}_{-0.06}$	$51.60^{+0.04}_{-0.04}$	28.3	0.5	$\frac{14.1}{14} = 1.00$
G349.7+0.2	$2.06^{+0.13}_{-0.12}$	$2.82^{+0.30}_{-0.38}$	2.70	> 5.00	$2.00^{+0.12}_{-0.12}$	$50.09^{+0.04}_{-0.04}$	1.30	35	$\frac{5.2}{10} = 0.52$
Hess J1731-347	$1.86^{+0.04}_{-0.04}$	$3.65^{+0.10}_{-0.10}$	$4.27^{+0.02}_{-0.02}$	> 5.19	$1.46^{+0.02}_{-0.02}$	$49.42^{+0.04}_{-0.04}$	45.1	0.01	$\frac{283.9}{322} = 0.88$
Hess J1745-303	$1.64^{+0.04}_{-0.04}$	$0.52^{+0.20}_{-0.18}$	2.03	> 5.37	$1.66^{+0.08}_{-0.08}$	$49.53^{+0.08}_{-0.08}$	167	100	$\frac{3.62}{8} = 0.45$

# Parameters

# Results

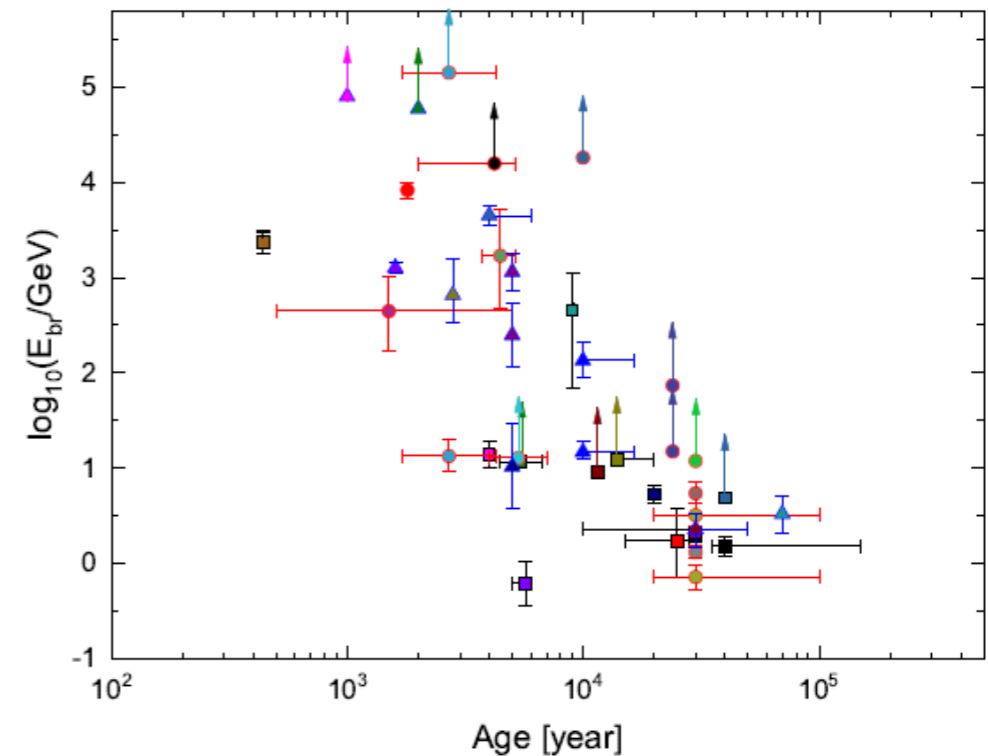
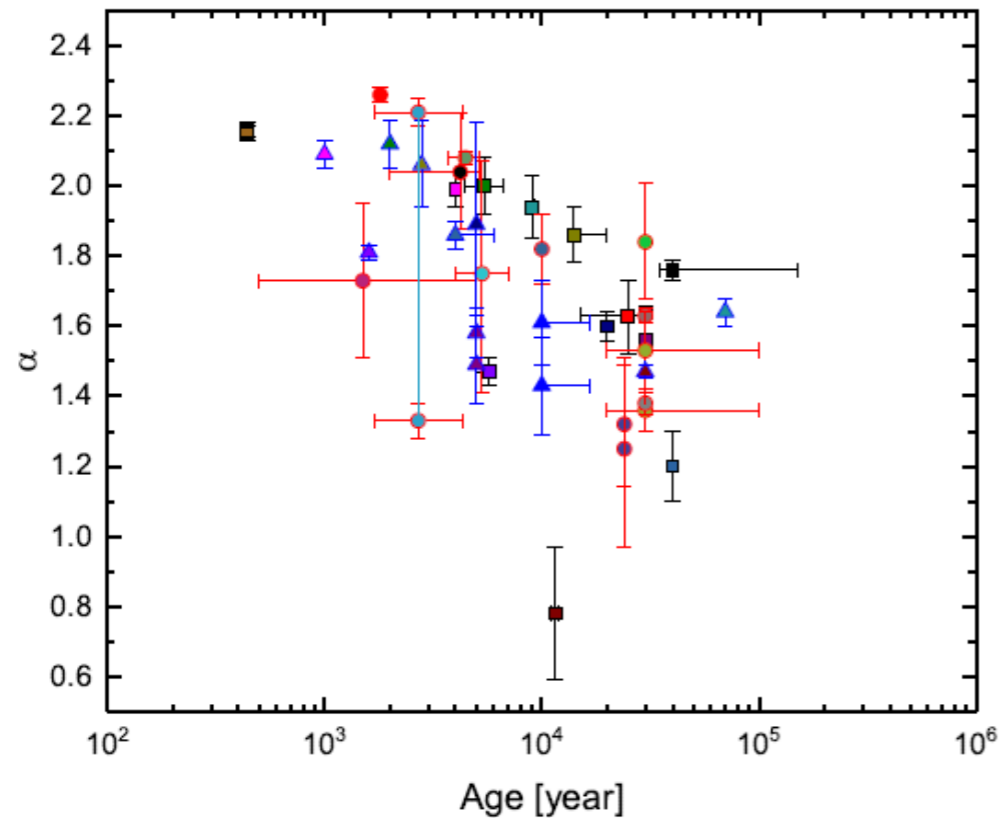
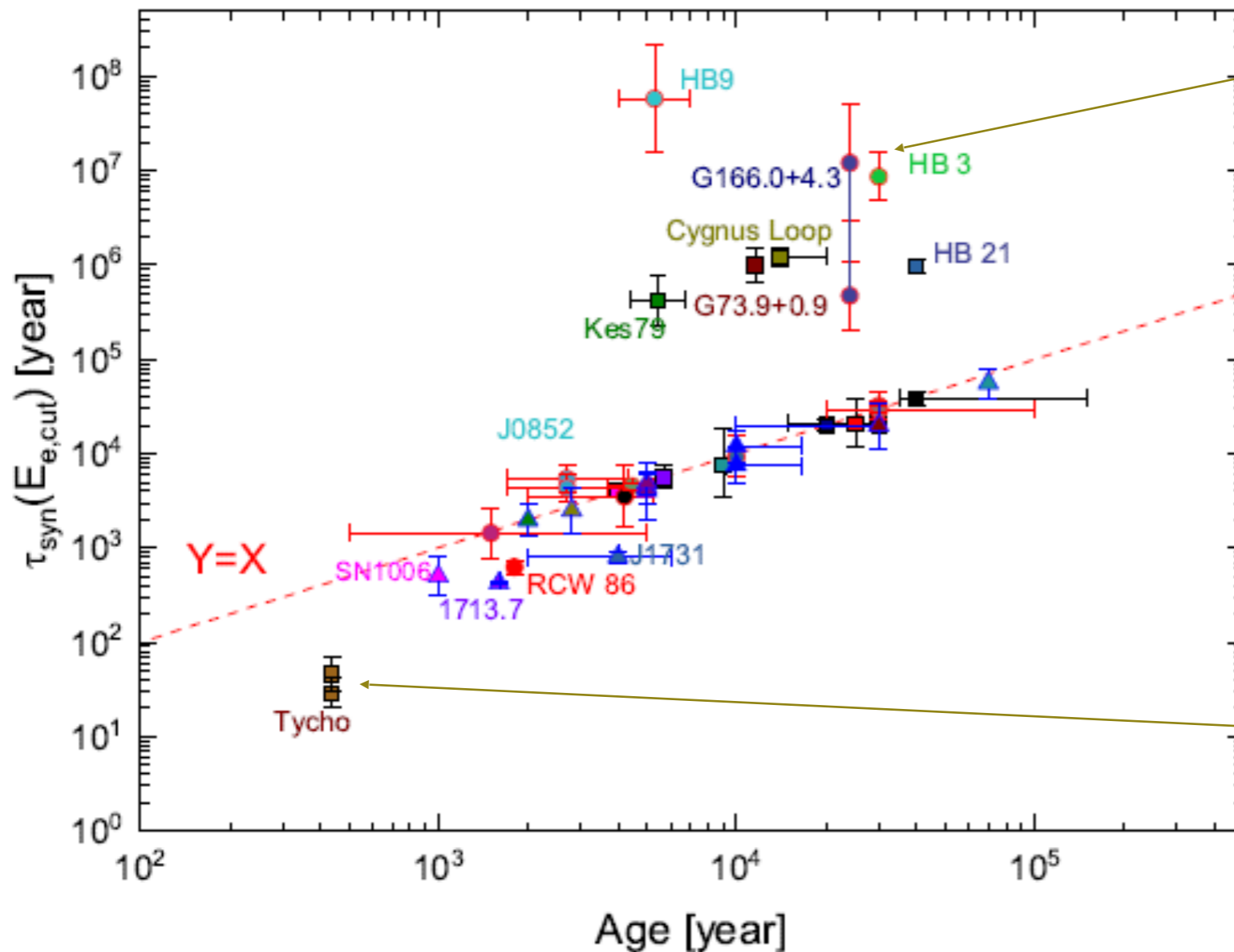


Figure (a) show there is an inverse correlation between the low-energy spectra index and age, and the spectra become harder with aging of SNRs. This result agrees with the observational fact that the radio spectrum hardens with aging of SNRs (e.g. Dubner & Giacani (2015)).

Figure (b): The break energy of particle distribution decreases with the age of SNRs, which may be related to the gradual weakening of shock waves with aging in SNRs and agrees with the particle acceleration model proposed by Ohira et al. (see also Zhang et al. 2017)

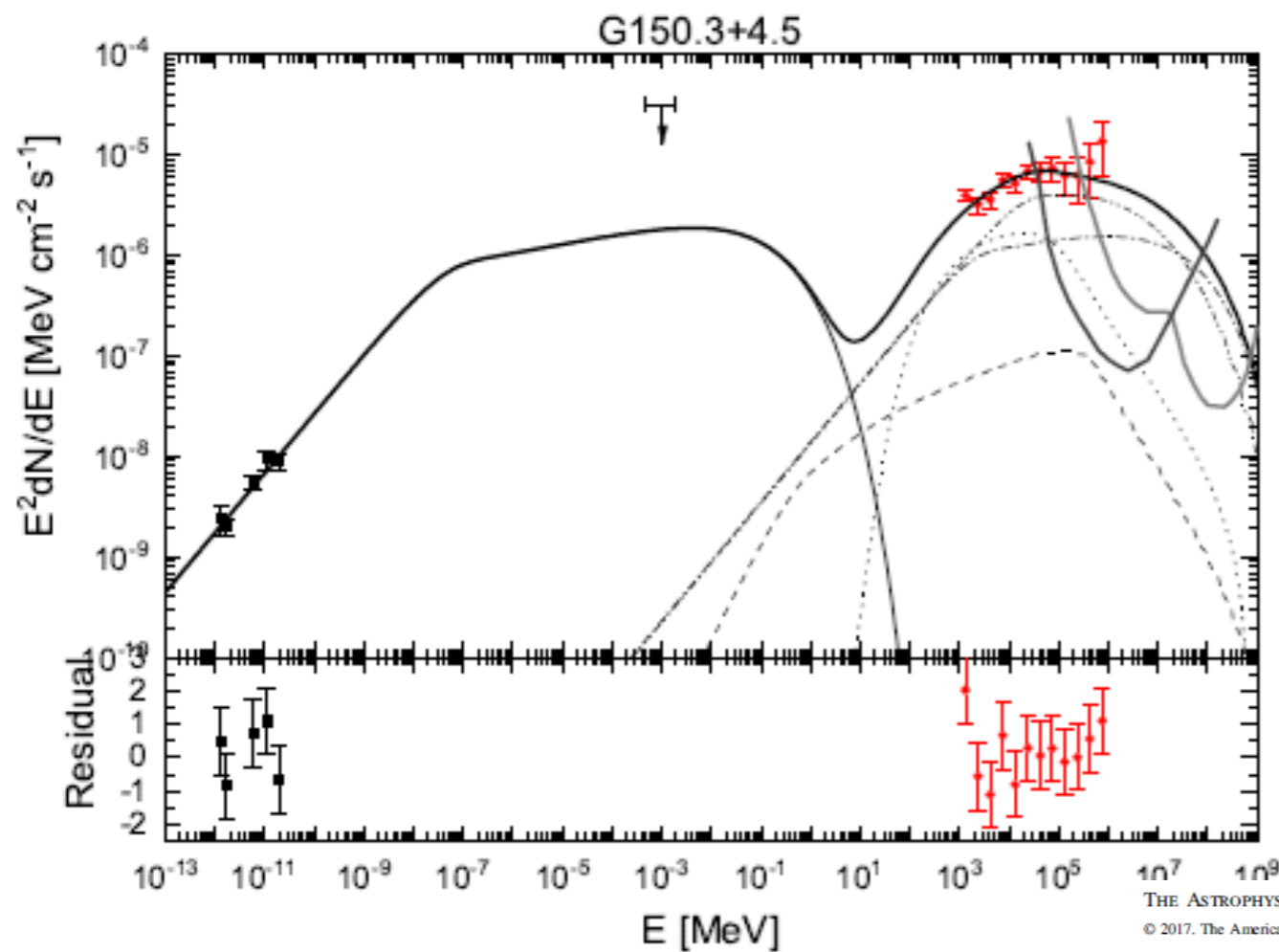
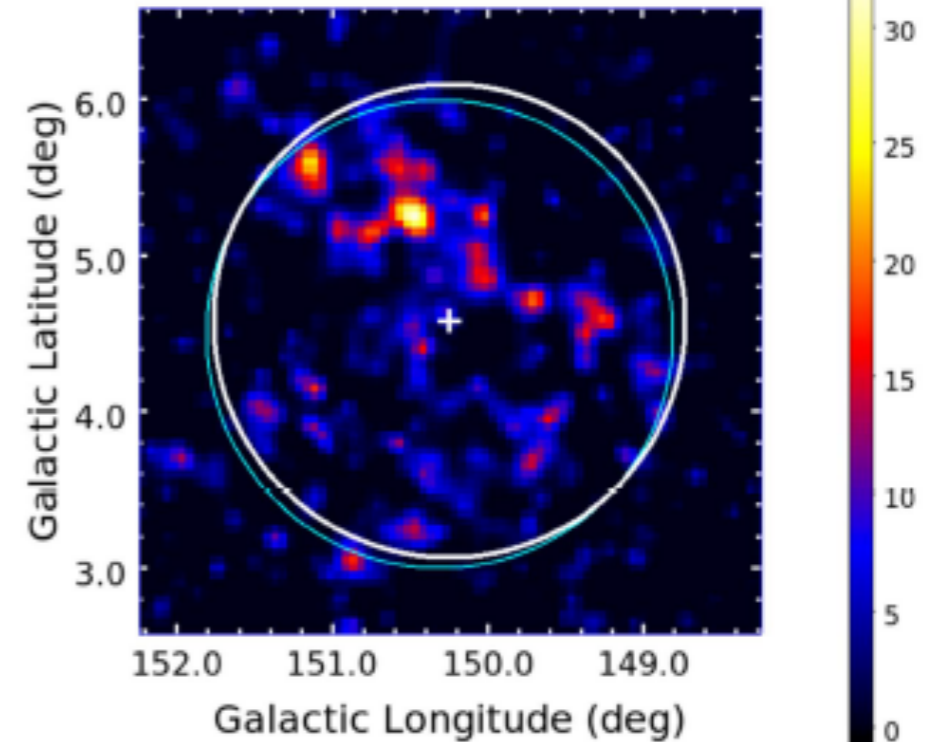
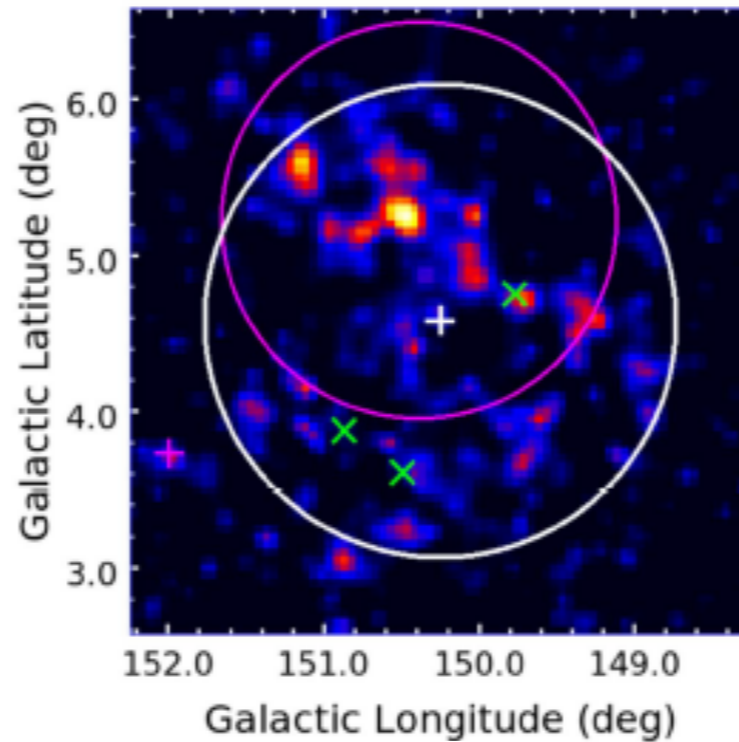
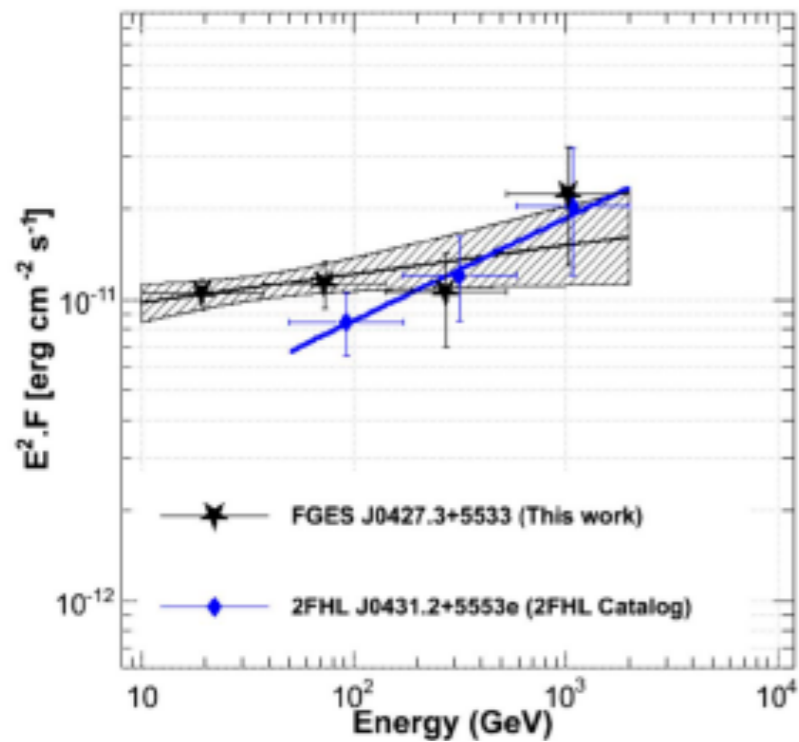
# Results



**Energy loss timescale  $>$  Age  
may suggest escape of highest  
energy particles from SNRs.**

**Energy loss timescale  $<$   
Age implying continuous  
particle acceleration at  
the cutoff energy of SNRs**

**For most middle-age SNRs, the energy loss timescale of electrons at the high-energy cutoff is approximately equal to the age of SNRs, implying quenching of high-energy electron acceleration (Ohira et al).**



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<https://doi.org/10.3847/1538-4357/aa775a>



CrossMark

**G150.3+4.5**

**A good LHAASO target?**

Search for Extended Sources in the Galactic Plane Using Six Years of *Fermi*-Large Area Telescope Pass 8 Data above 10 GeV

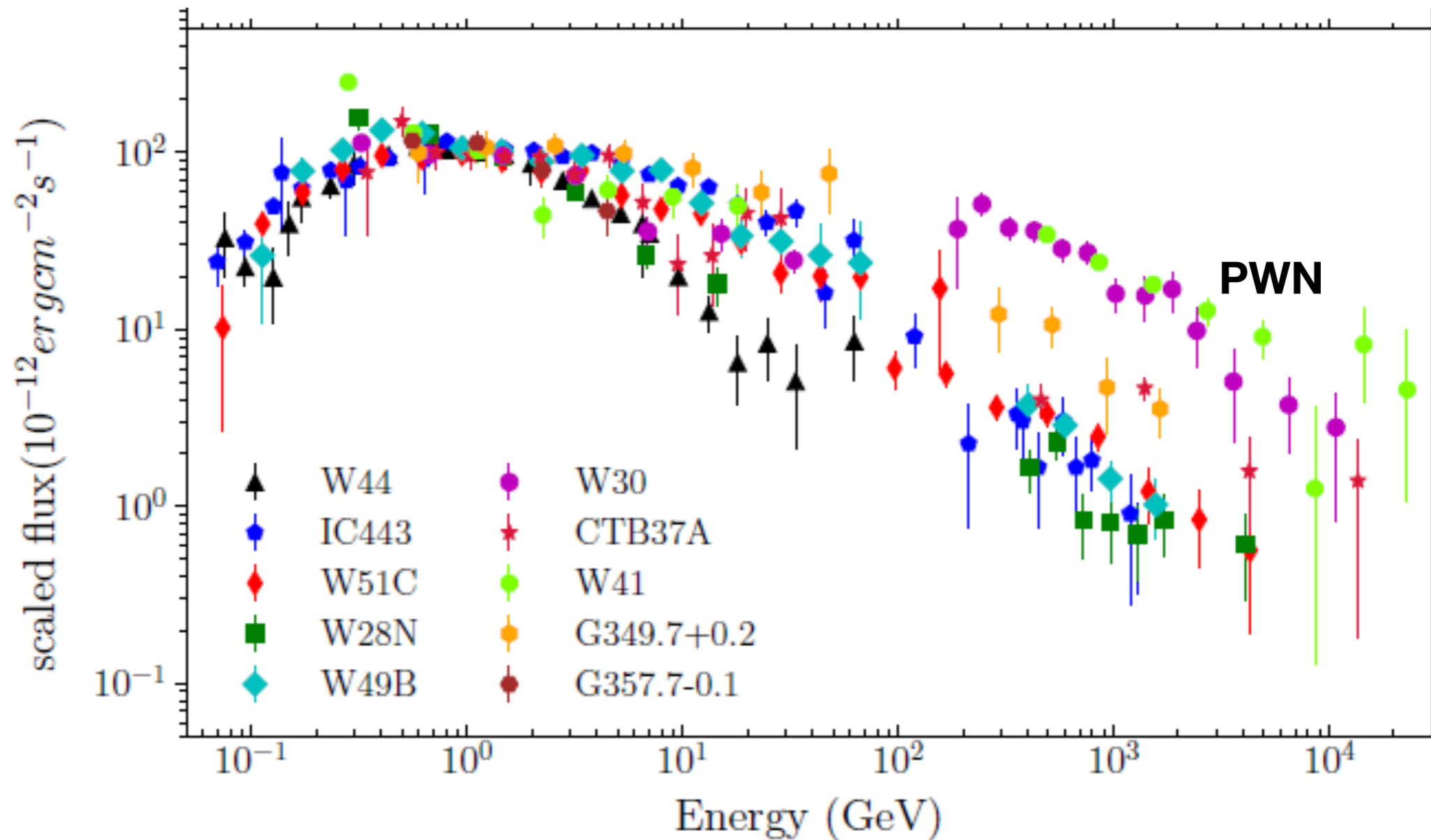
M. Ackermann<sup>1</sup>, M. Ajello<sup>2</sup>, L. Baldini<sup>3</sup>, J. Ballet<sup>4</sup>, G. Barbiellini<sup>5,6</sup>, D. Bastieri<sup>7,8</sup>, R. Bellazzini<sup>9</sup>, E. Bissaldi<sup>10</sup>, E. D. Bloom<sup>11</sup>,

# Gamma-ray emission from middle-aged supernova remnants interacting with molecular clouds: the challenge for current models

Xiaping Tang,<sup>1,2</sup>★

<sup>1</sup>Max Planck Institute for Astrophysics, Karl-Schwarzschild-Str. 1, D-85741 Garching, Germany

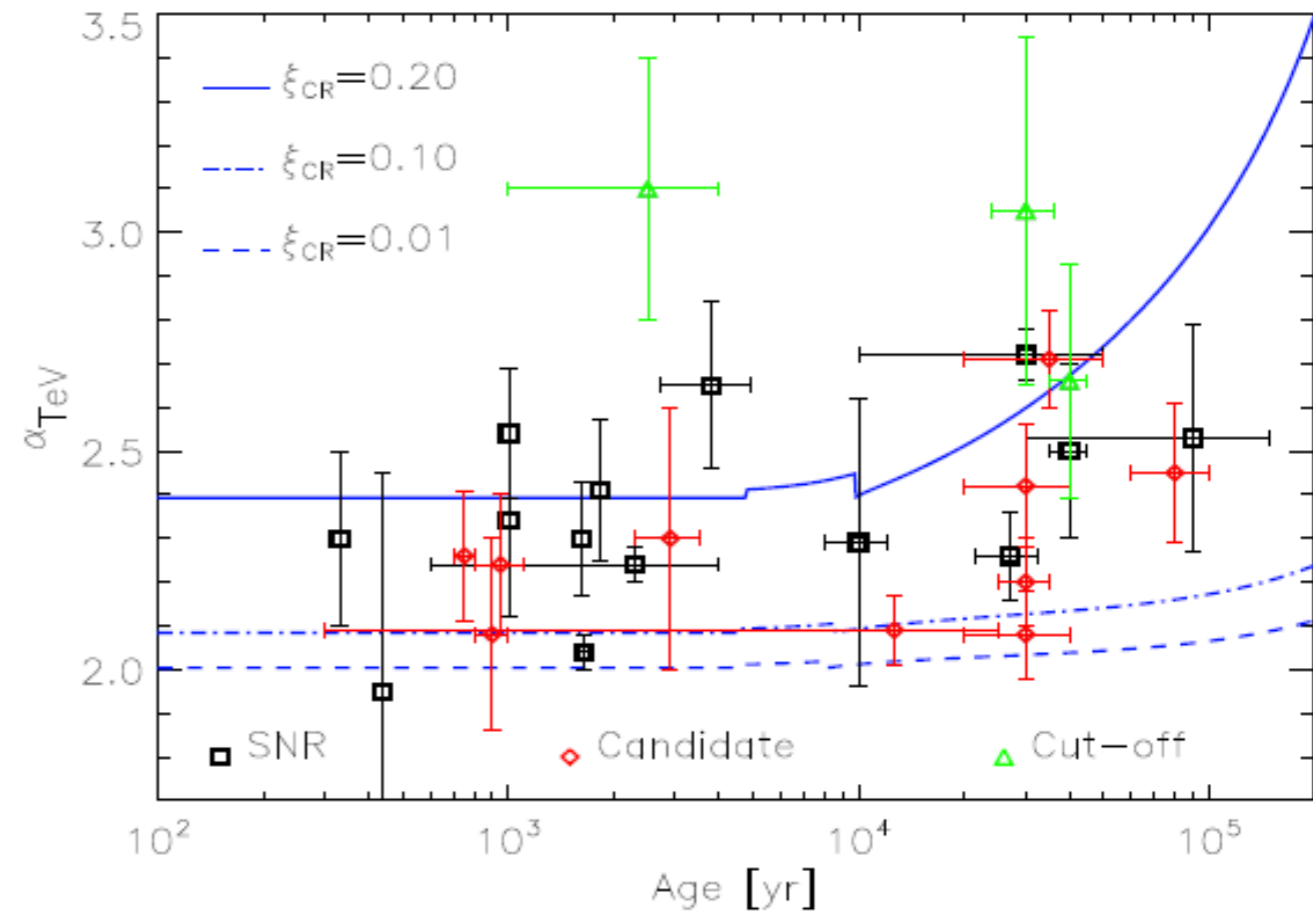
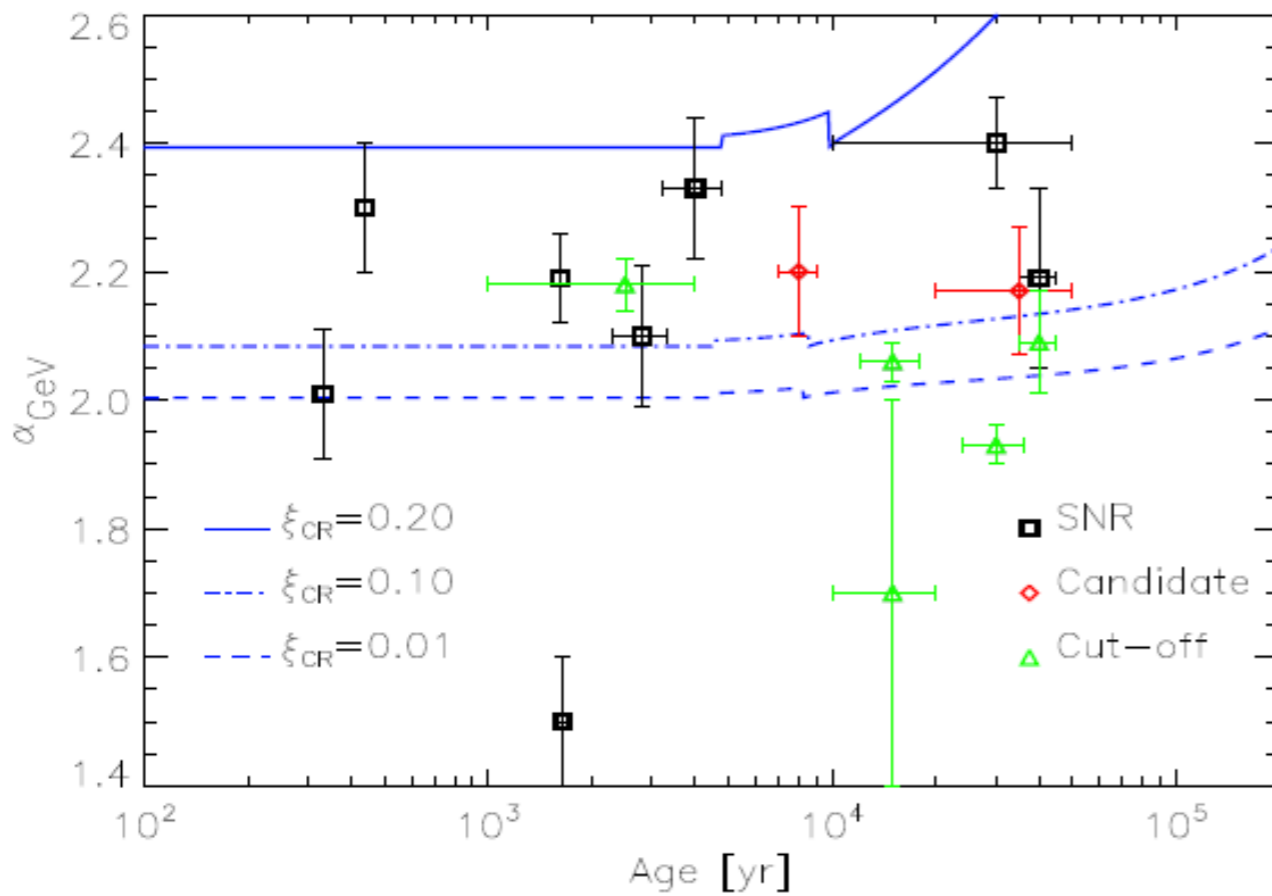
<sup>2</sup>The Racah Institute of physics, The Hebrew University of Jerusalem, Jerusalem 91904, Israel



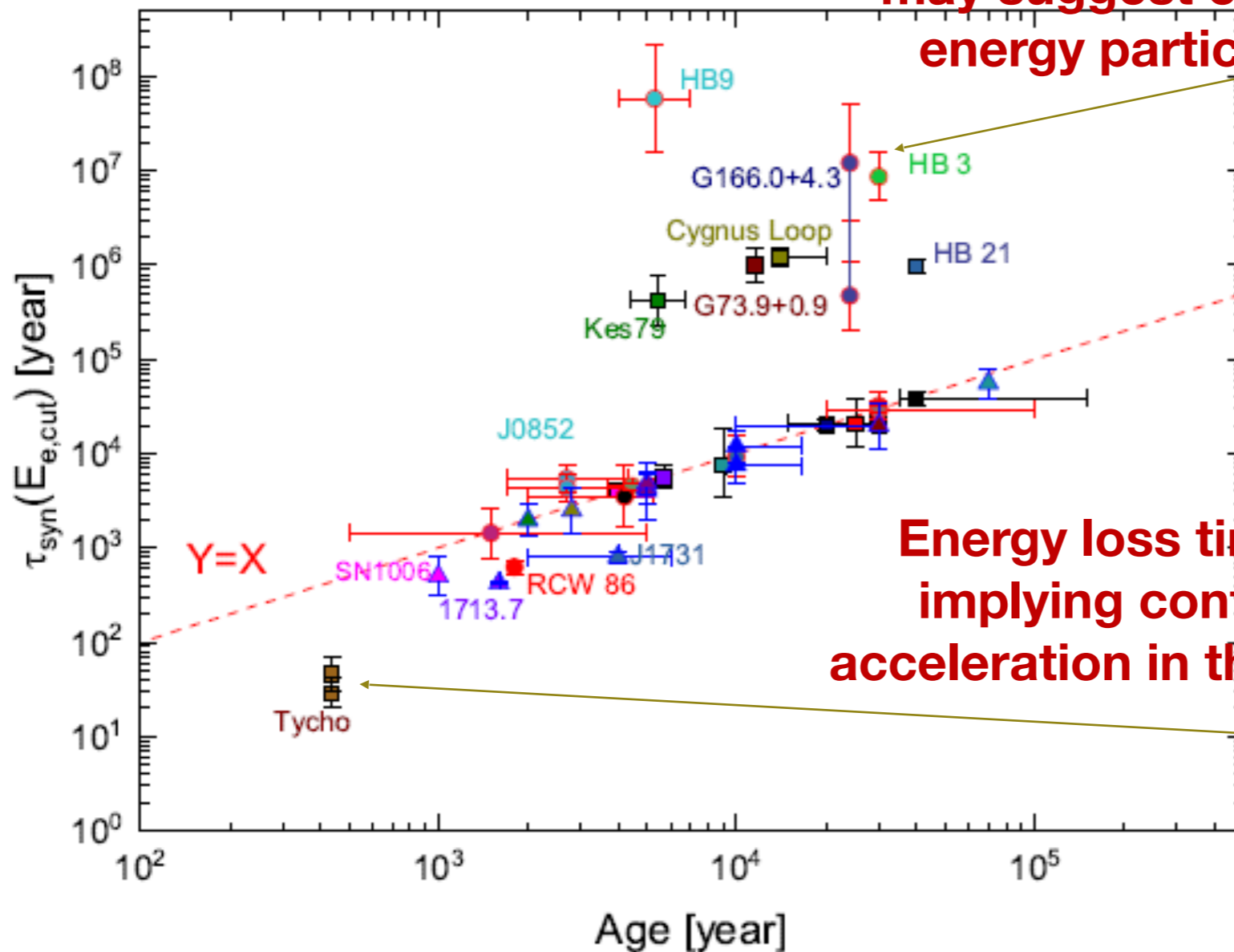
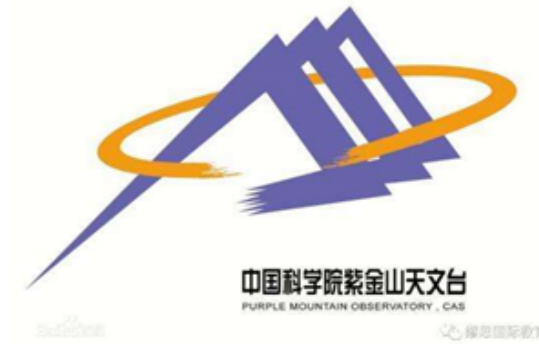
**Gamma ray index does not change significantly: Implying energy-independent escape**

# Understanding hadronic gamma-ray emission from supernova remnants

Caprioli Damiano



# Cosmic ray injection

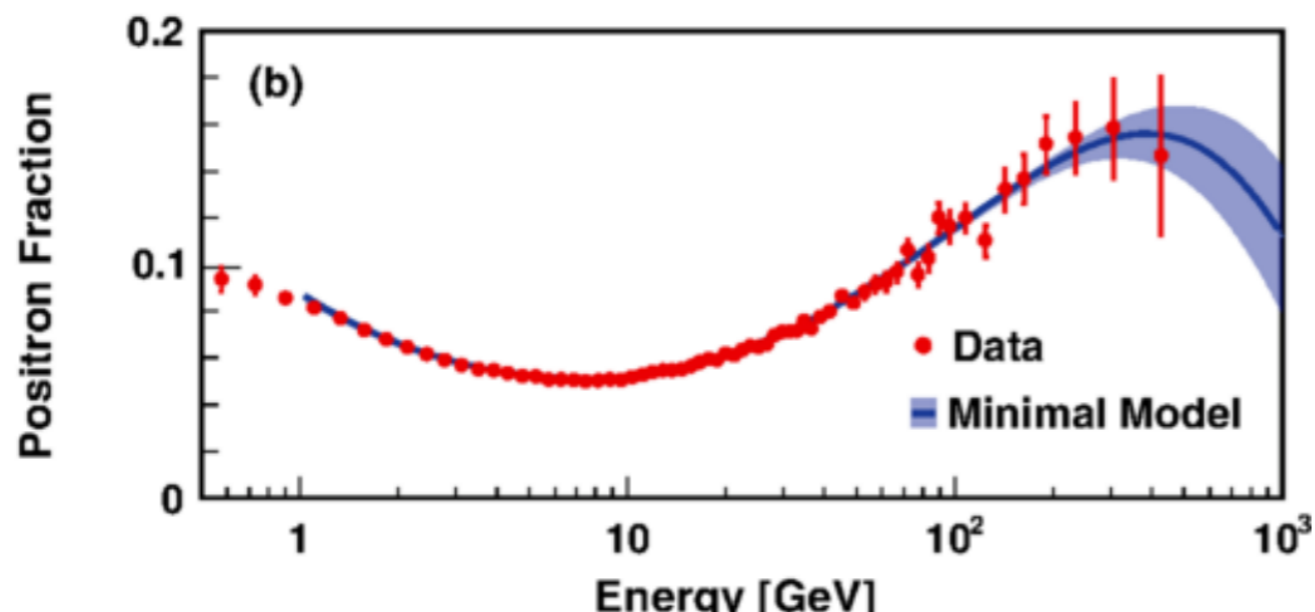
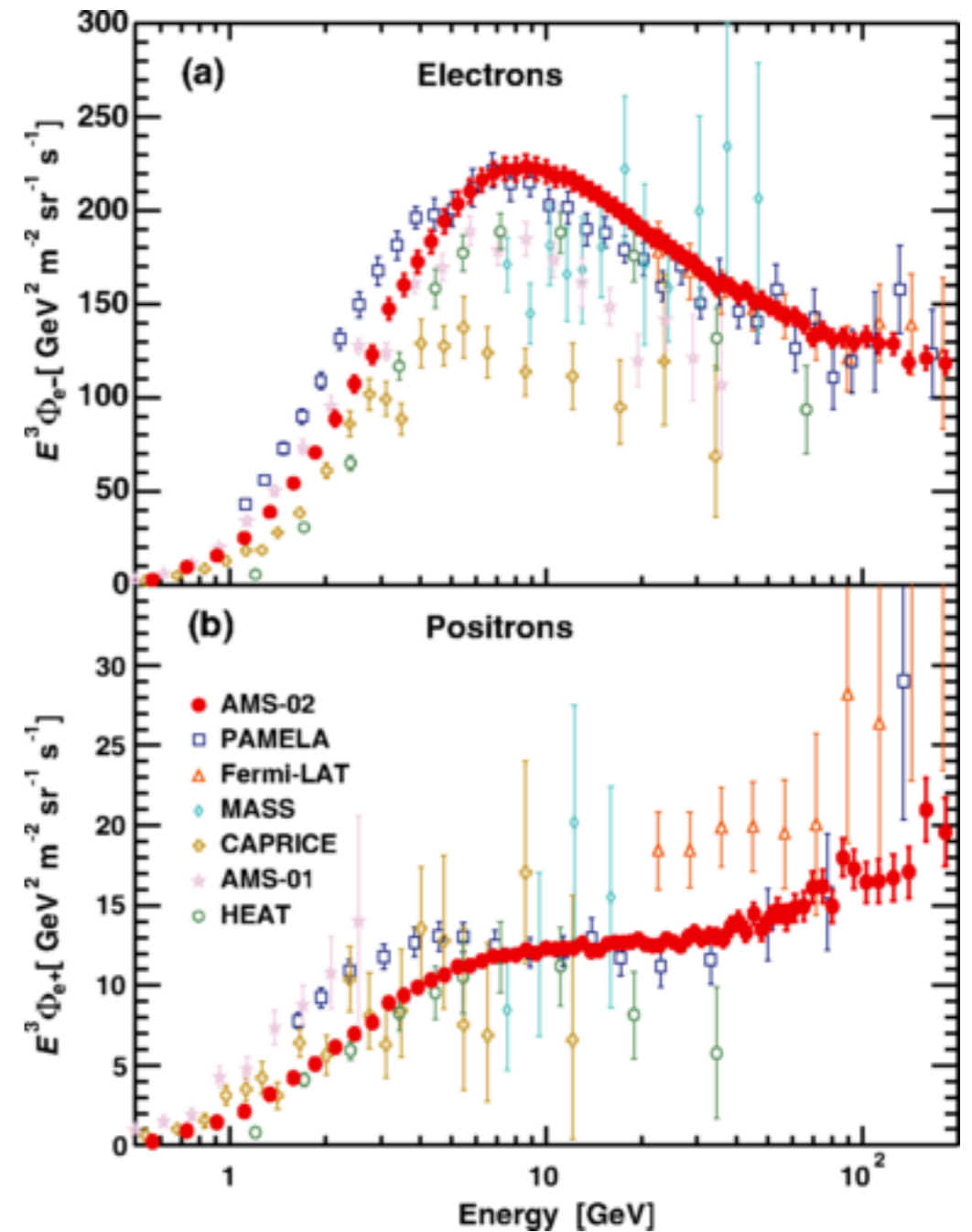
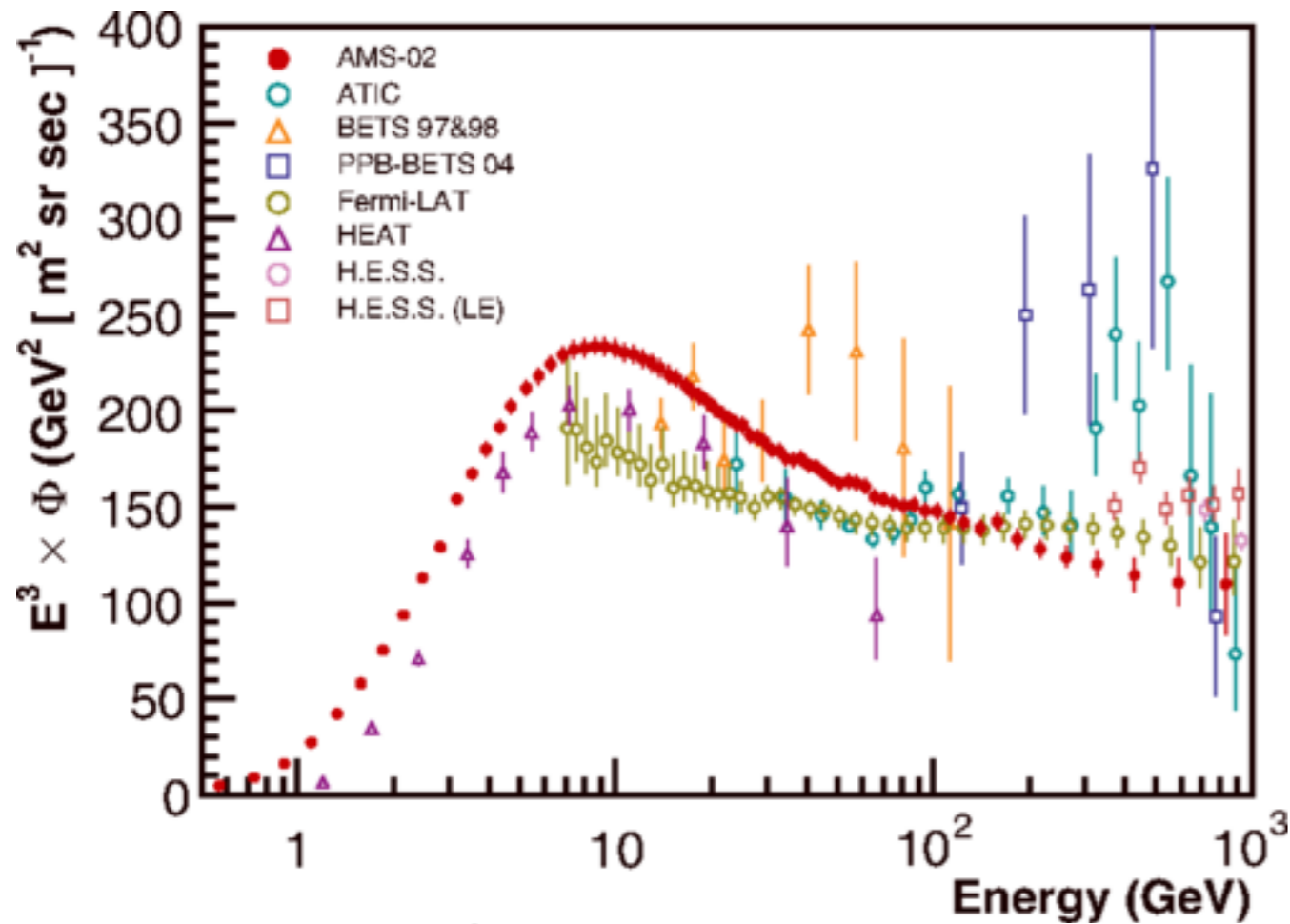


Energy loss timescale > Age  
may suggest escape of highest  
energy particles from SNRs.

Energy loss timescale < Age  
implying continuous particle  
acceleration in the interior of SNRs

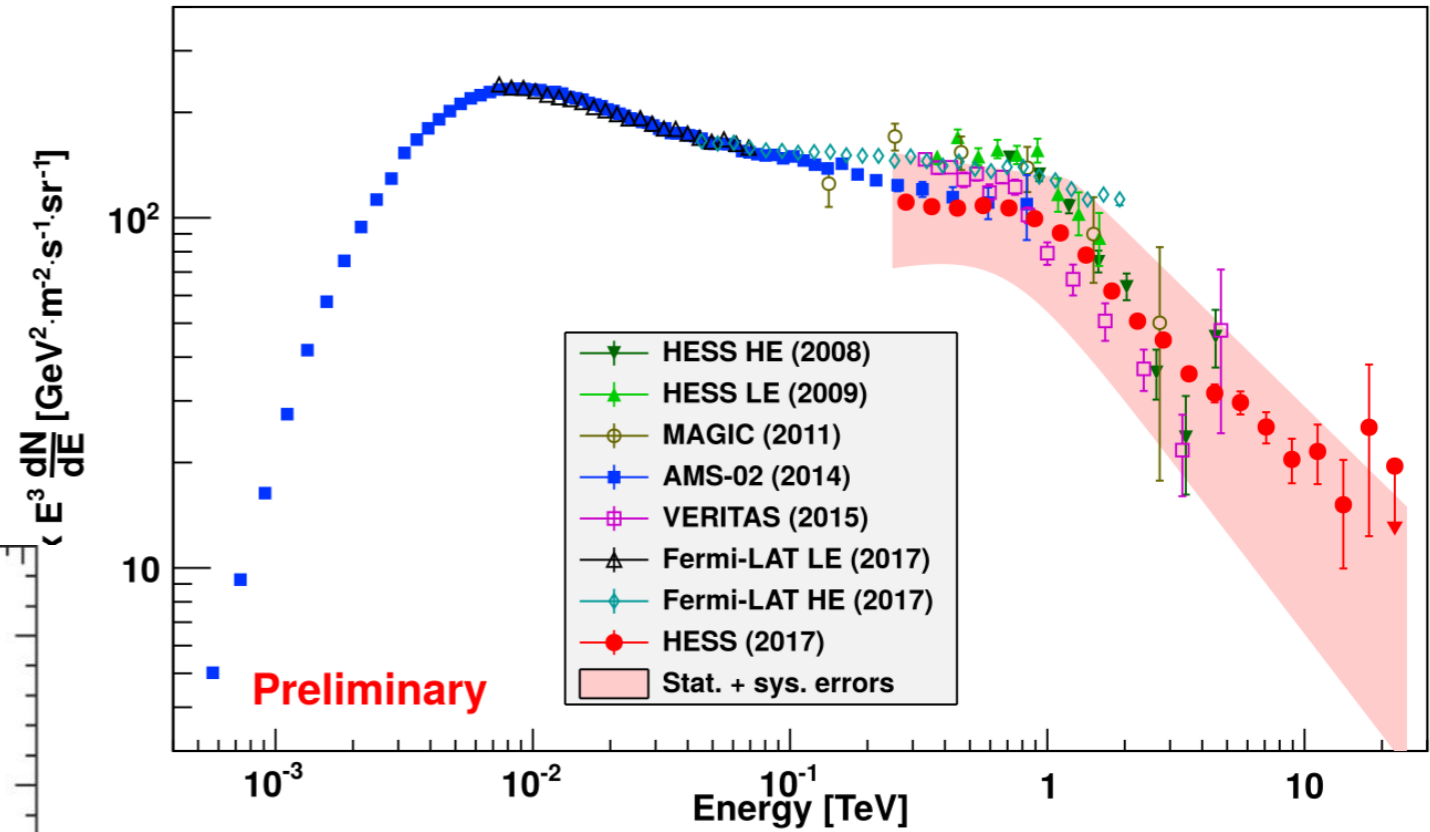
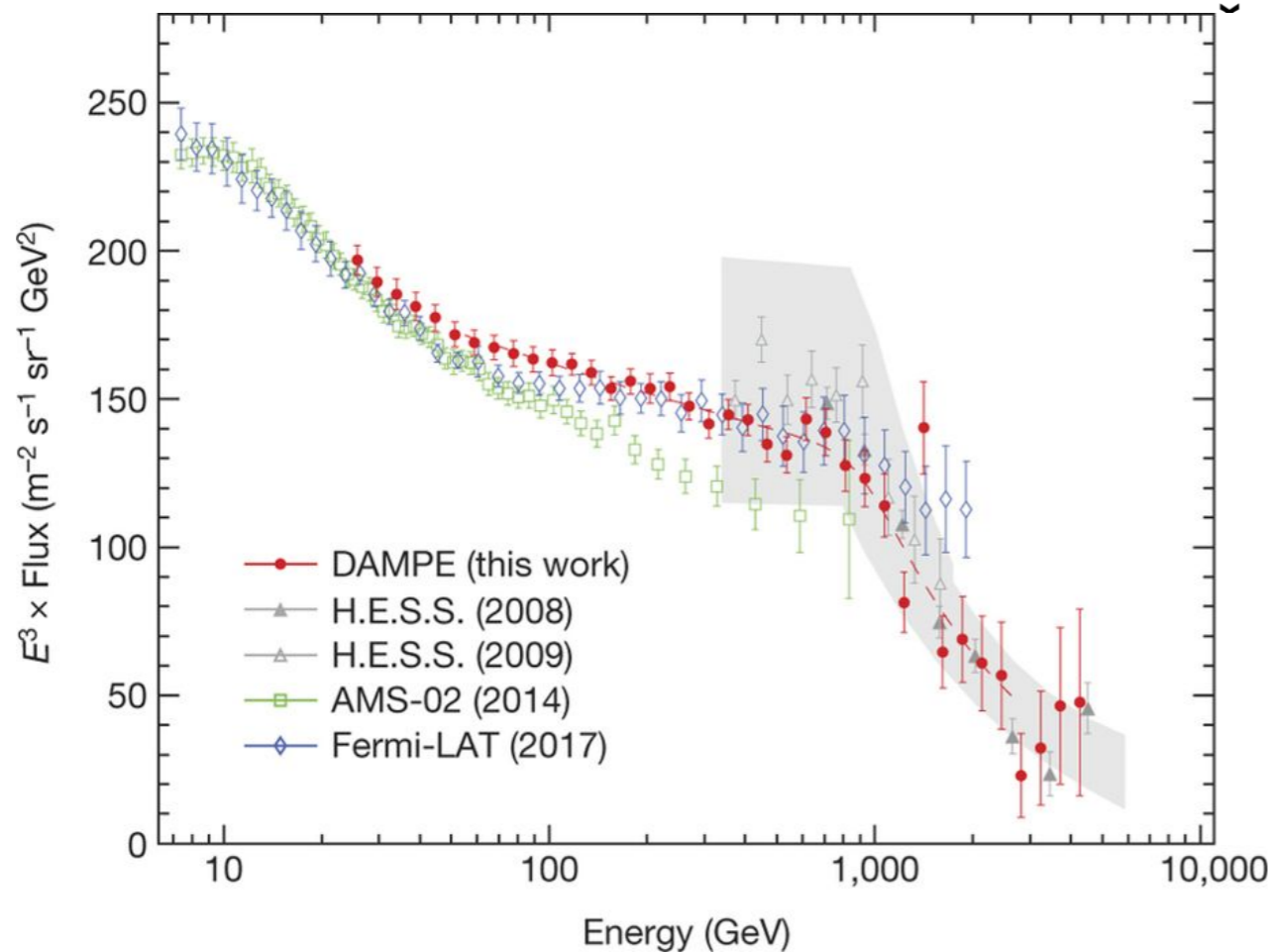
For most middle-age SNRs, the energy loss timescale of electrons at the high-energy cutoff is approximately equal to the age of SNRs, implying quenching of high-energy electron acceleration (Ohira et al).

# Cosmic ray electron & positron





# Cosmic ray electron & positron



# Modeling electron & positron spectra

- Propagated Fluxes of Electrons & Positrons

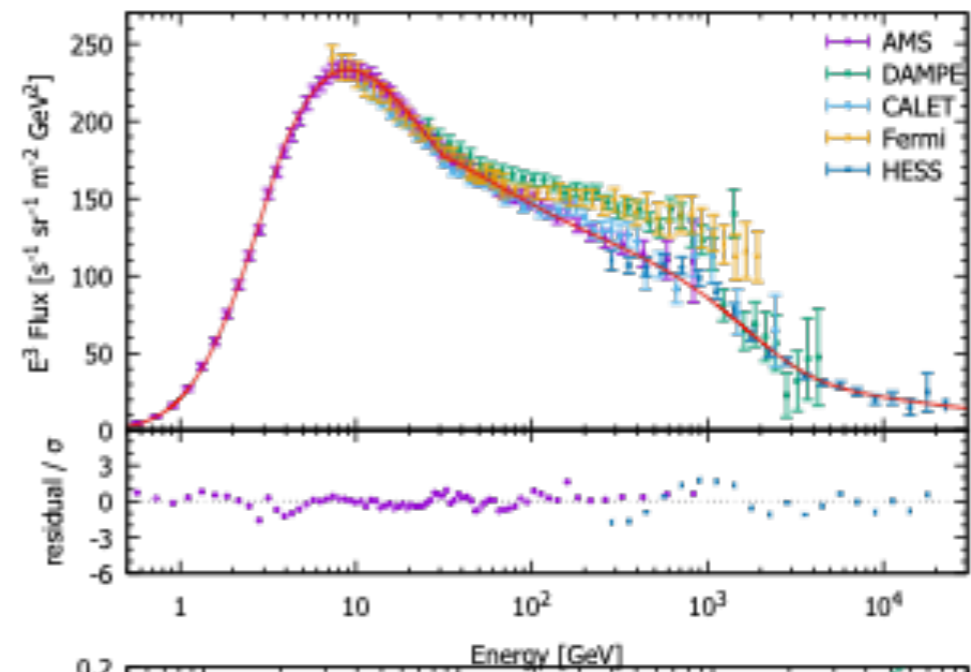
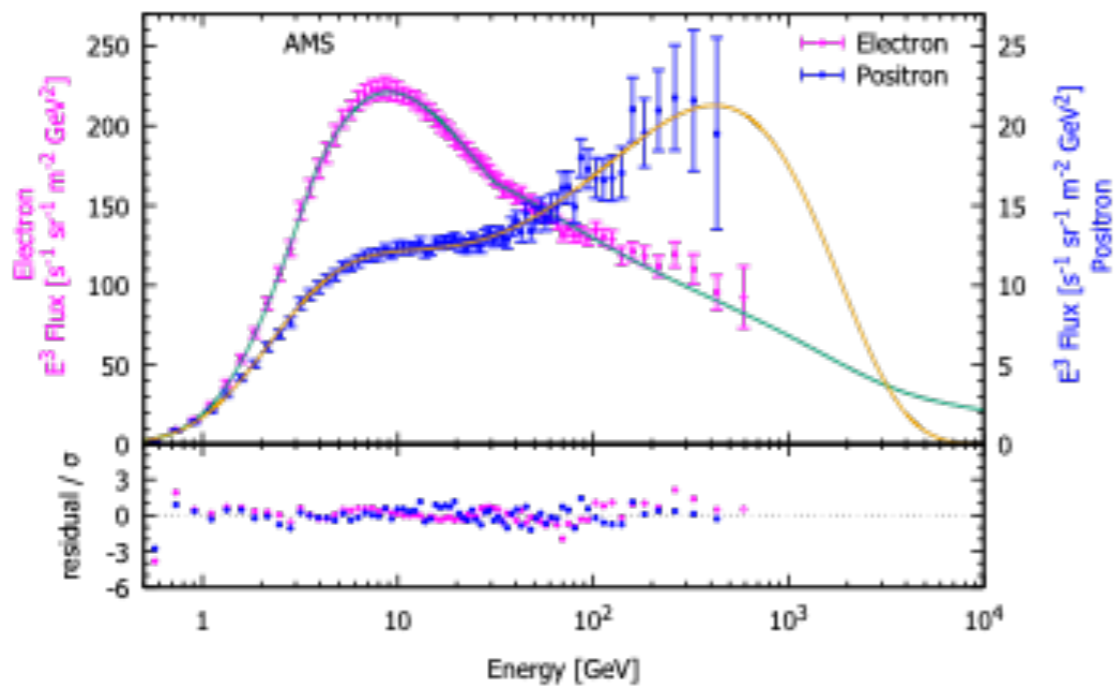
$$J_{e^-} = CE^{-\gamma} e^{-\frac{E}{E_{cut}}} + 0.6C_{e^+} E^{-\gamma_{e^+}} + C_{e^-} \begin{cases} E^{-\gamma_{e^-}^1} & E \leq E_{br1} \\ E_{br1}^{\gamma_{e^-}^2 - \gamma_{e^-}^1} E^{-\gamma_{e^-}^2} & E_{br1} < E \leq E_{br2} \\ E_{br1}^{\gamma_{e^-}^2 - \gamma_{e^-}^1} E_{br2}^{\gamma_{e^-}^3 - \gamma_{e^-}^2} E^{-\gamma_{e^-}^3} & E > E_{br2} \end{cases}$$

$$J_{e^+} = CE^{-\gamma} e^{-\frac{E}{E_{cut}}} + C_{e^+} E^{-\gamma_{e^+}}$$

- Solar Modulation

$$J(E) = \left( \frac{E}{E + e\phi} \right)^2 J_0(E + e\phi)$$

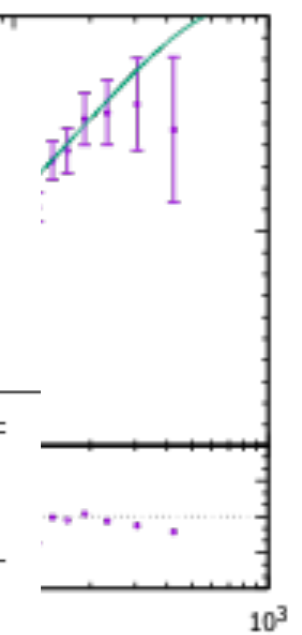
# Modeling CR Electron & Positron Spectra



**solar modulation  $\phi = 1.140$  GV**

$C_{e^-}^\dagger$	$\gamma_{e^-}^1$	$E_{br1}$	$\gamma_{e^-}^2$	$E_{br2}$	$\gamma_{e^-}^3$	$C_{e^+}^\dagger$	$\gamma_{e^+}$	$C^\dagger$	$\gamma$	$E_{cut}$
1063	3.347	4.960	3.643	32.355	3.370	163.25	4.048	3.139	2.620	1100

$C_{e^-}^{inj\dagger}$	$\gamma_{e^-}^{1,inj}$	$E_{br}^{inj}$	$\gamma_{e^-}^{2,inj}$	$C_{e^+}^{inj\dagger}$	$\gamma_{e^+}^{inj}$	$C^{inj\dagger}$	$\gamma^{inj}$	$E_{cut}^{inj}$
$6.612 \times 10^{41}$	3.059	40.917	2.644	$1.053 \times 10^{41}$	3.683	$1.050 \times 10^{39}$	2.040	1555



$^\dagger$  with units of  $s^{-1} sr^{-1} m^{-2} GeV^{-1}$ .

$^\ddagger$  with units of  $kyr^{-1} pc^{-3} GeV^{-1}$ .

# 1D Propagation Model

- 1D Steady-state Transport Equation

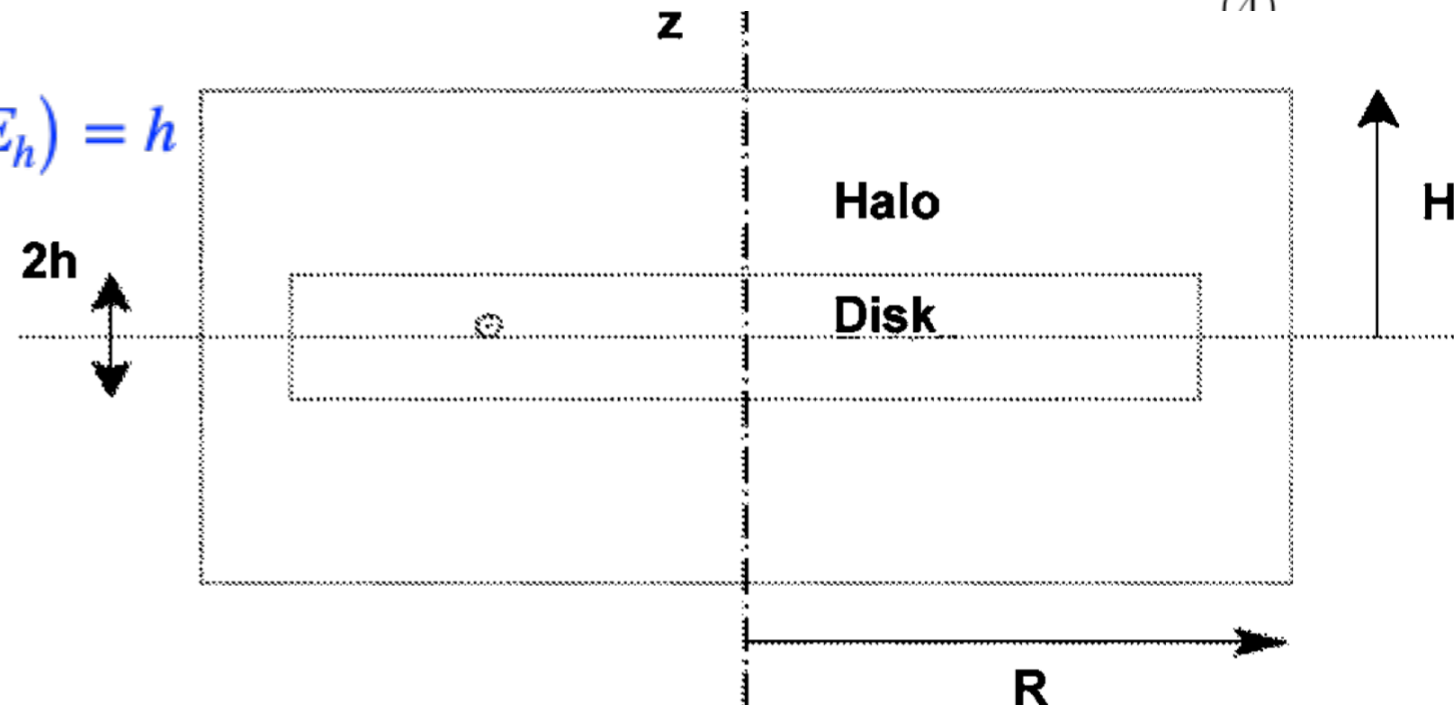
$$D(E) \frac{\partial^2 N(z, E)}{\partial z^2} + \frac{\partial b(E) N(z, E)}{\partial E} + Q(z, E) = 0$$

$$D(E) = D_0 (E/1 \text{ GeV})^\delta, \quad b(E) = -dE/dt = b_0 E^2$$

$$N(\pm H, E) = 0, \quad Q(z, E) = q(E) \theta(h - |z|)$$

$$N(z, E) = \frac{2}{b_0 E^2} \sum_{m=0}^{\infty} \frac{\sin \left[ \left( m + \frac{1}{2} \right) \pi \frac{h}{H} \right]}{\left( m + \frac{1}{2} \right) \pi} \cos \left[ \left( m + \frac{1}{2} \right) \pi \frac{z}{H} \right] \int_E^{\infty} dE' q(E') \exp \left\{ \left[ \left( m + \frac{1}{2} \right) \pi \frac{\ell(E)}{H} \right]^2 \left[ \left( \frac{E}{E'} \right)^{1-\delta} - 1 \right] \right\}.$$

$$\ell(E) \equiv \sqrt{\frac{D_0 E^{\delta-1}}{(1-\delta)b_0}}, \quad \ell(E_H) = H, \quad \ell(E_h) = h$$



Credit: Ptuskin (2001)

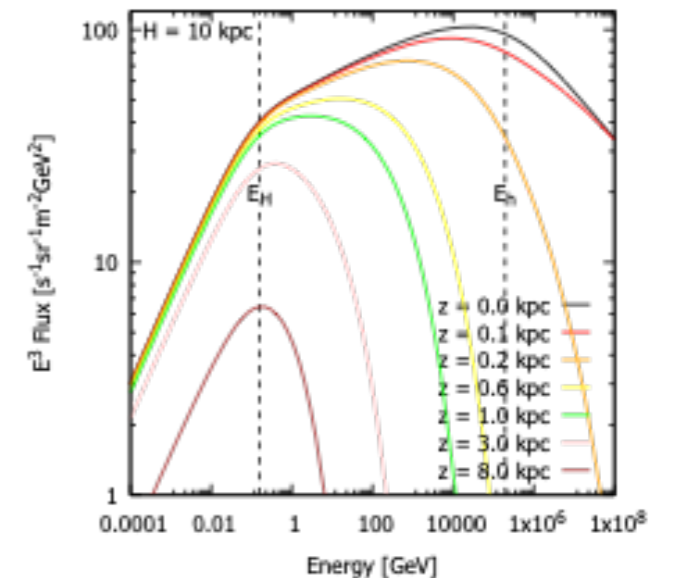
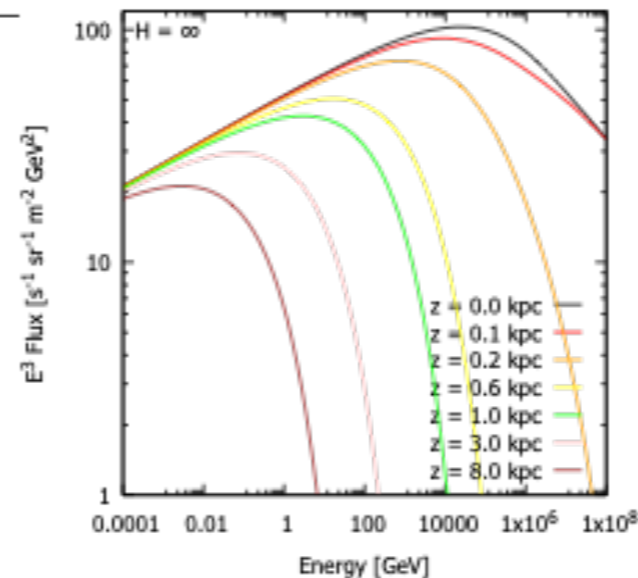
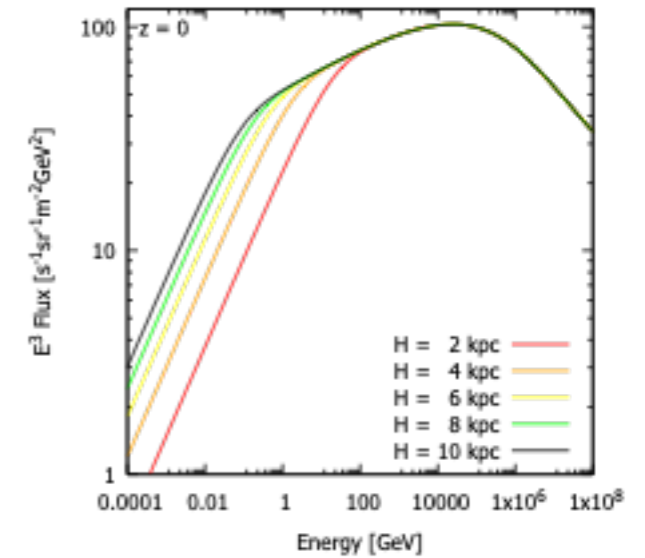
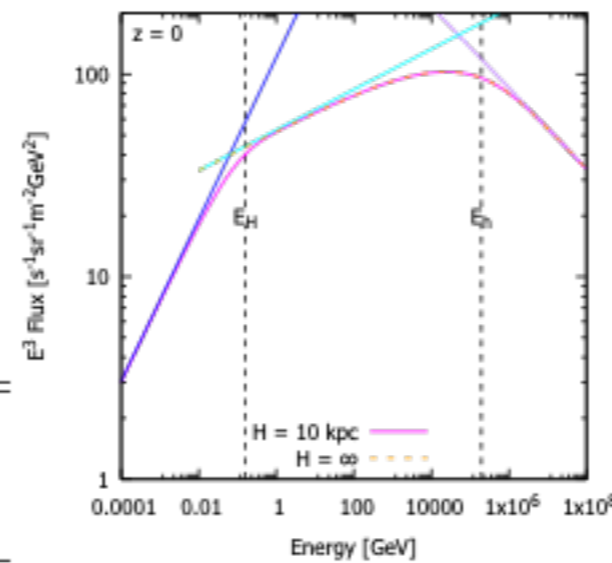
# 1D Propagation Model

al index ( $z = 0$ ) for a power-law injection

$$\gamma_a = \begin{cases} \gamma + \delta, & E < E_H \\ \gamma + \frac{\delta+1}{2}, & E_H < E < E_h \\ \gamma + 1, & E > E_h \end{cases}$$

$$\ell(E) \equiv \sqrt{\frac{D_0 E^{\delta-1}}{(1-\delta)b_0}}, \quad \ell(E_H) = H, \quad \ell(E_h) = h$$

$D_0$	$\delta$	$b_0$	$\gamma$	$K$
$[\text{pc}^2 \text{ kyr}^{-1}]$		$[\text{GeV}^{-1} \text{ kyr}^{-1}]$		$[\text{kyr}^{-1} \text{ pc}^{-3} \text{ GeV}^{-1}]$
100	0.4	$5 \times 10^{-6}$	2.2	$1.0 \times 10^{40}$



# Injection Spectra of Electrons & Positrons

- Kolmogorov's Spectrum  $\delta = 1/3$

- Best-fit Propagation Parameters

$D_0$	$b_0$	$h$	$H$
[pc <sup>2</sup> kyr <sup>-1</sup> ]	[GeV <sup>-1</sup> kyr <sup>-1</sup> ]	[pc]	[kpc]
88.386	$6.417 \times 10^{-6}$	145.164	2.704

- Electron Injection Spectrum is double power-law

$$E_H = 4.75 \text{ GeV} \sim E_{br1} = 4.96 \text{ GeV} \quad E_h = 30.7 \text{ TeV}$$

- Electron-Positron Excess

**solar modulation  $\phi = 1.273 \text{ GV}$**

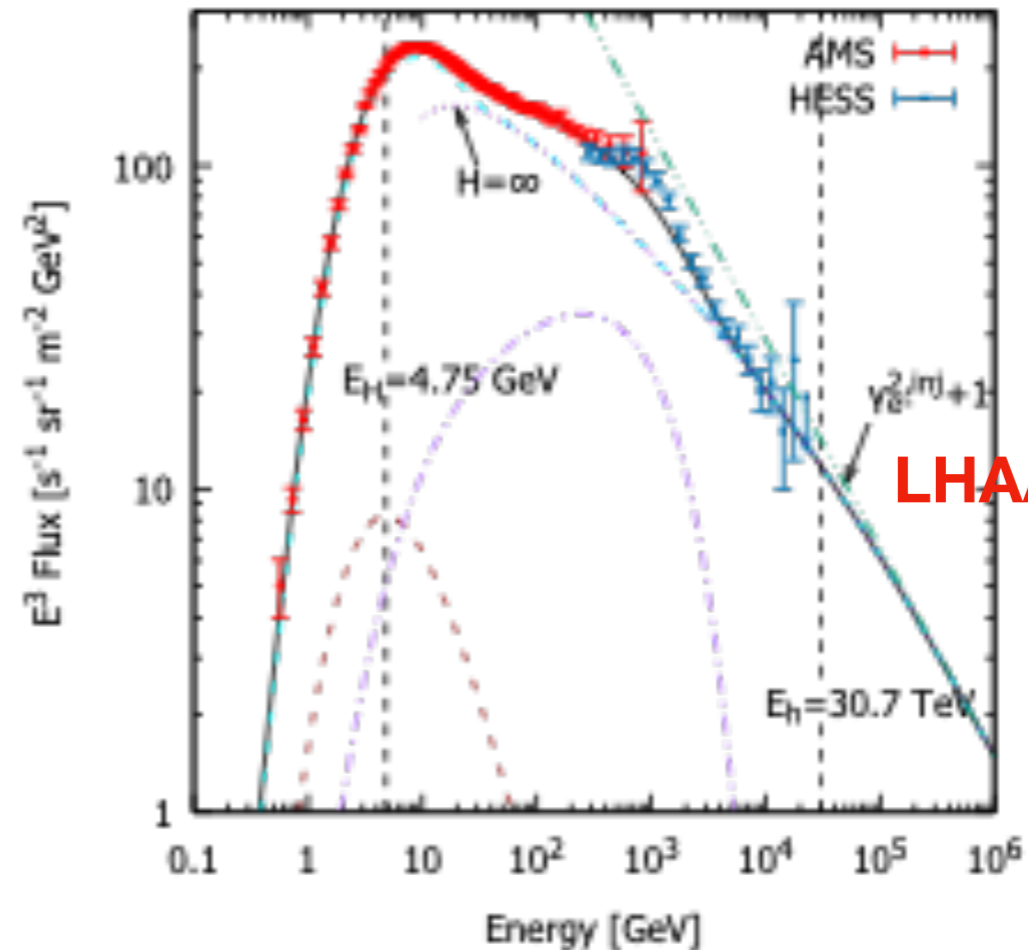
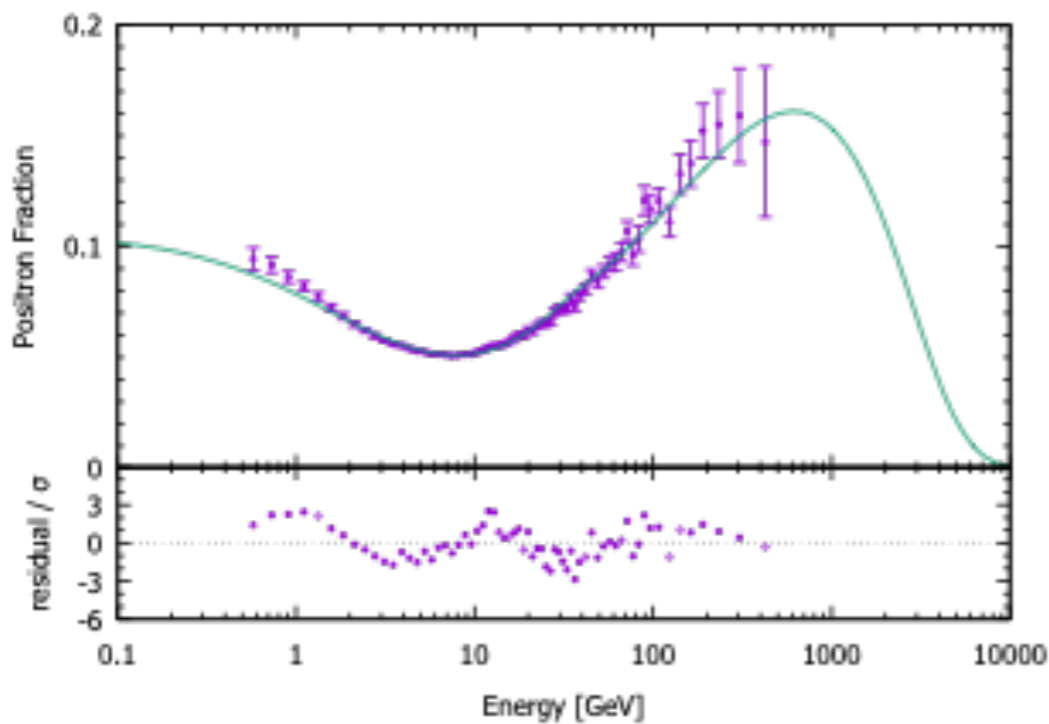
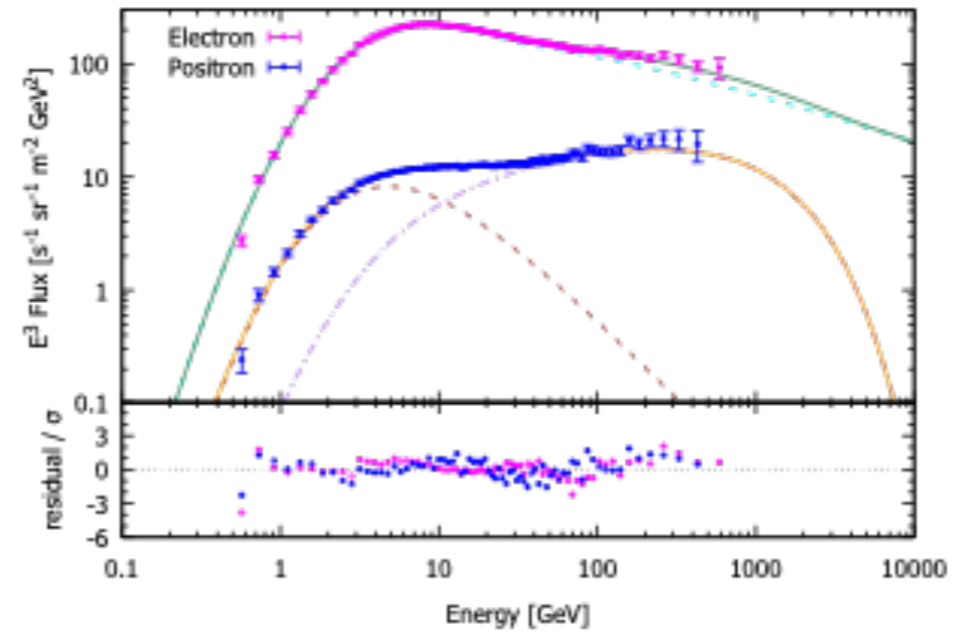
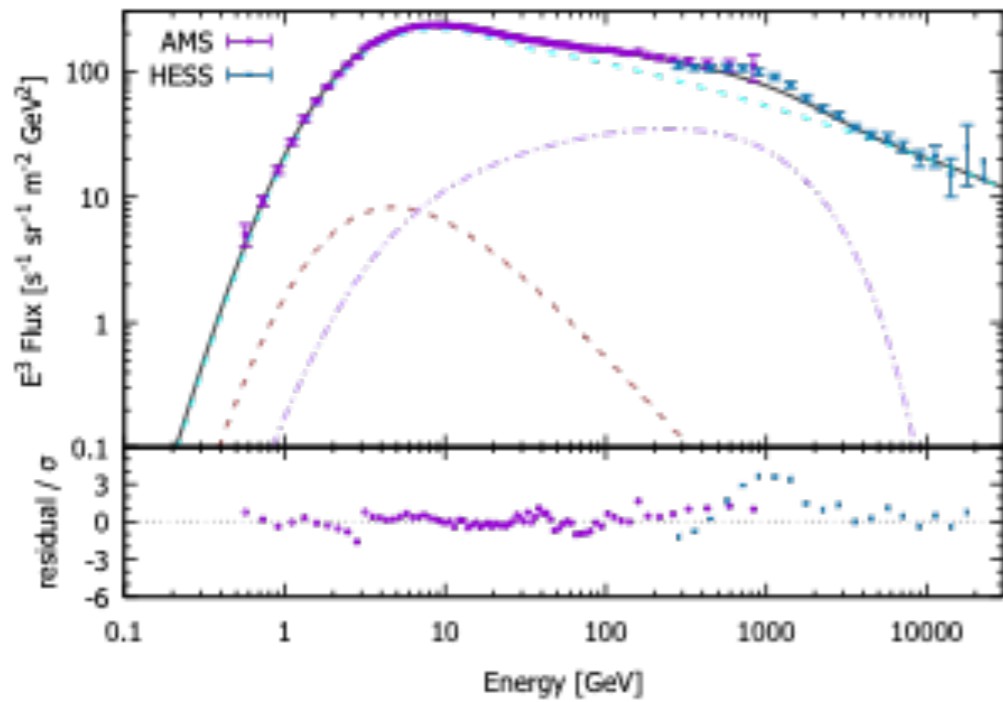
$C_{e^-}^\dagger$	$\gamma_{e^-}^1$	$E_{br1}$	$\gamma_{e^-}^2$	$E_{br2}$	$\gamma_{e^-}^3$	$C_{e^+}^\dagger$	$\gamma_{e^+}$	$C^\dagger$	$\gamma$	$E_{cut}$
1063	3.347	4.960	3.643	32.355	3.370	163.25	4.048	3.139	2.620	1100

$C_{e^-}^{inj\dagger}$	$\gamma_{e^-}^{1,inj}$	$E_{br}^{inj}$	$\gamma_{e^-}^{2,inj}$	$C_{e^+}^{inj\dagger}$	$\gamma_{e^+}^{inj}$	$C^{inj\dagger}$	$\gamma^{inj}$	$E_{cut}^{inj}$
$6.612 \times 10^{41}$	3.059	40.917	2.644	$1.053 \times 10^{41}$	3.683	$1.050 \times 10^{39}$	2.040	1555

<sup>†</sup> with units of s<sup>-1</sup> sr<sup>-1</sup> m<sup>-2</sup> GeV<sup>-1</sup>.

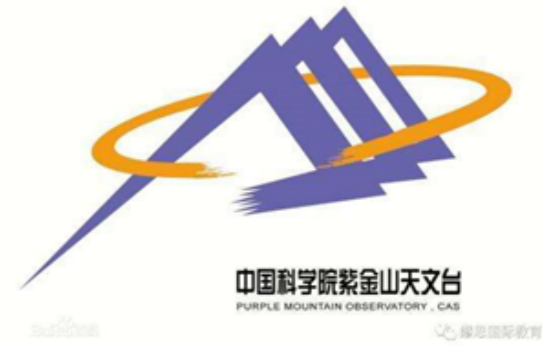
<sup>‡</sup> with units of kyr<sup>-1</sup> pc<sup>-3</sup> GeV<sup>-1</sup>.

# Spectra of Electron & Positron



LHAASO Target?

# 结论



银河系宇宙线起源问题的核心是把宇宙线的观测特征和银河系高能源的观测特征联系起来

- 1：如果银河系宇宙线主要来自于超新星遗迹，有迹象表明能量较高的宇宙线主要在遗迹激波演化的早期被加速，而大部分的且能量相对较低的宇宙线在超新星遗迹演化的晚期被加速。
- 2：年轻超新星遗迹中高能粒子的逃逸概率似乎不依赖于粒子的能量，而在年老遗迹中能谱的高能部分可能从遗迹中完全逃逸。
- 3：宇宙线正电子很可能主要来源于和脉冲星有关的过程。
- 4：超新星遗迹中的粒子加速和逃逸模型需要进一步完善以解决银河系宇宙线的起源问题。
- 5：10TeV以上电子谱测量以及超新星遗迹G150.3+4.5是重要的LHAASO观测对象。



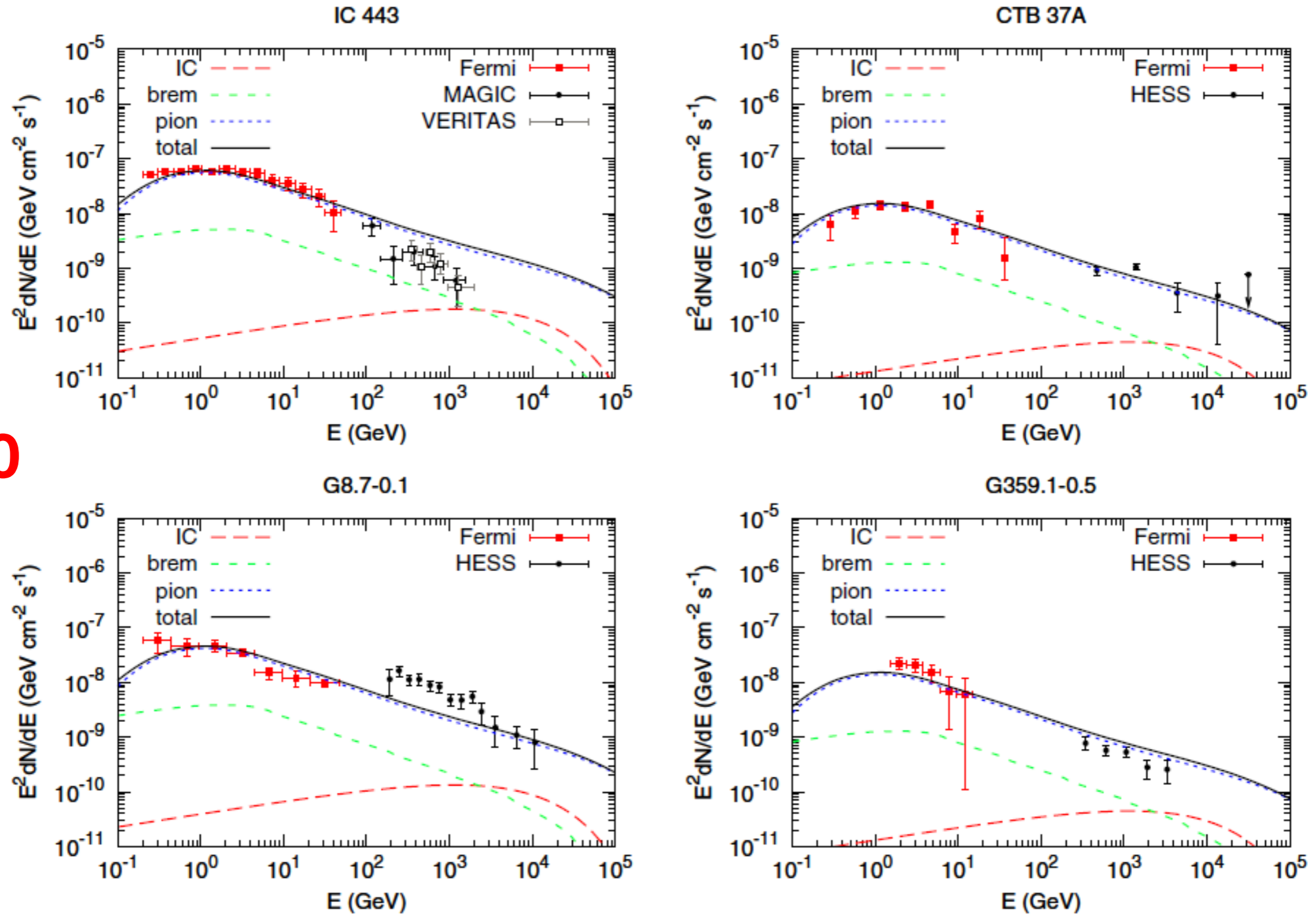


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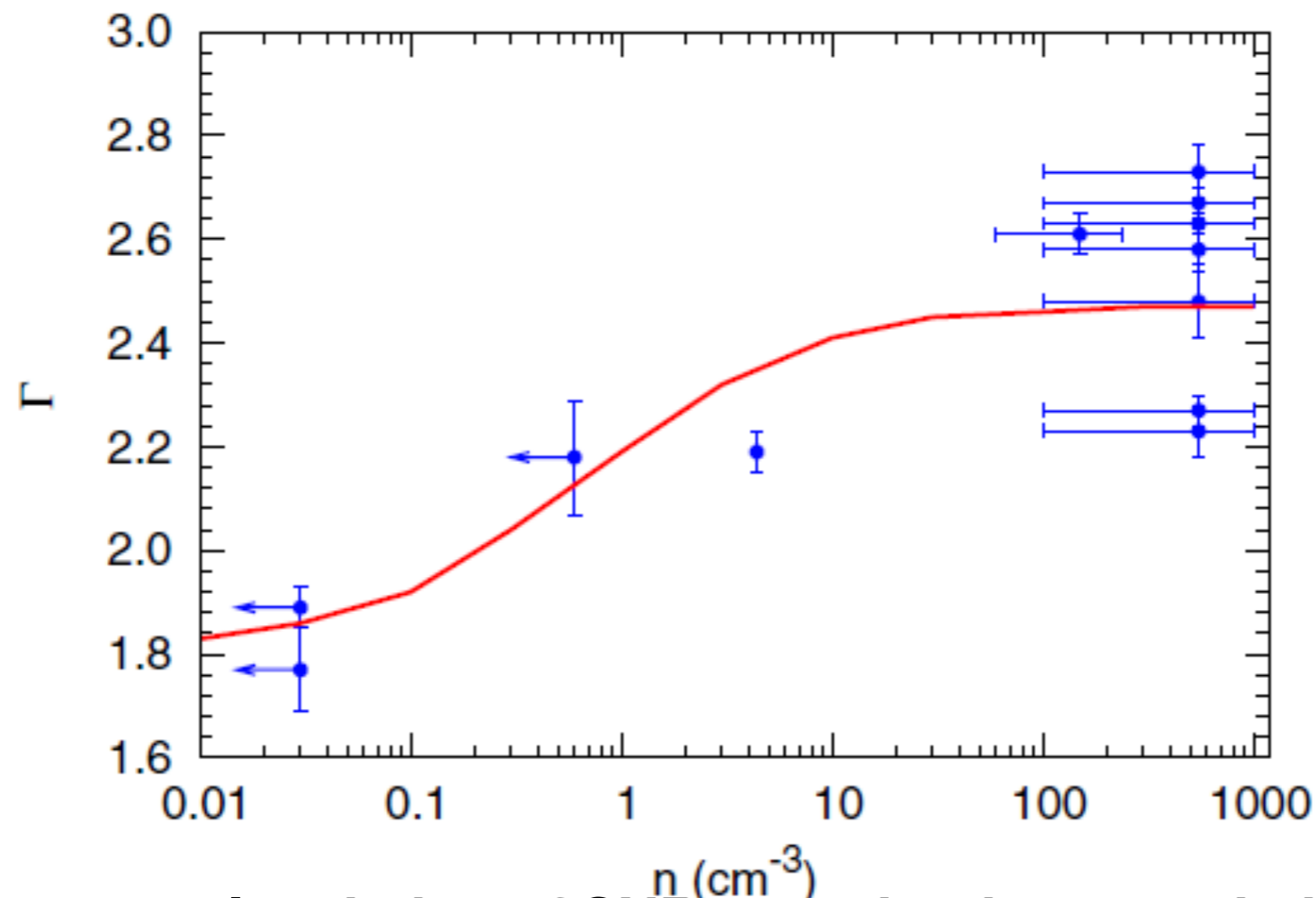
Purple Mountain Observatory, Chinese Academy of Sciences

***Thanks For Your Attention***

**n=100**



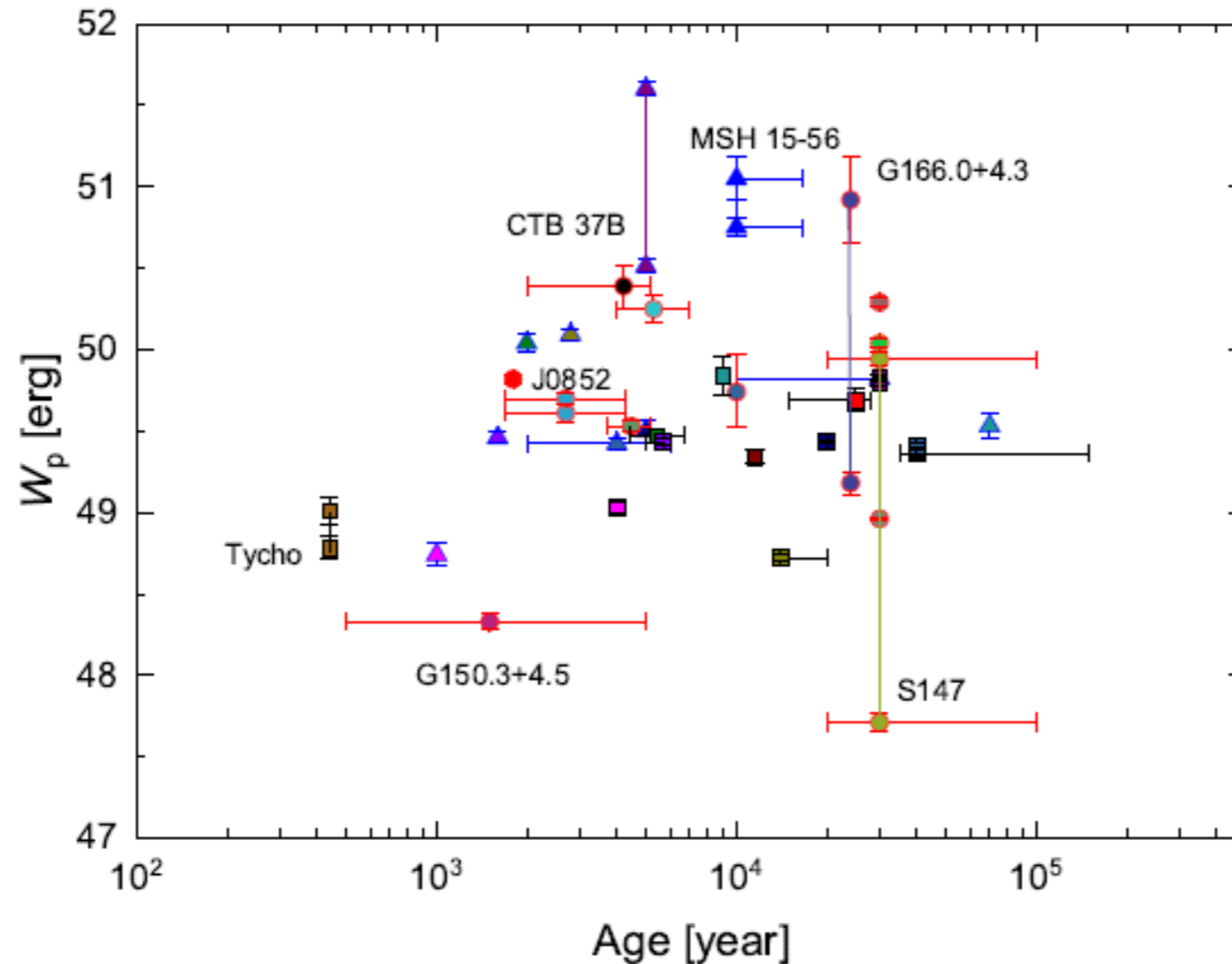
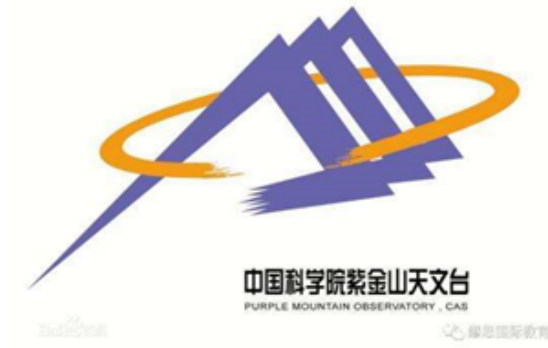
**Figure 4.** Same as Figure 2, but for SNR–MC interacting systems. The gas density is adopted to be  $n = 100 \text{ cm}^{-3}$ . References of the observational data—W28: *Fermi* (Abdo et al. 2010a), HESS (Aharonian et al. 2008c); W41: *Fermi* (Mehault et al. 2011), HESS (Mehault et al. 2011); W49B: *Fermi* (Abdo et al. 2010c), HESS (Brun et al. 2011); W51C: *Fermi* (Abdo et al. 2009), HESS (Fiasson et al. 2009), MAGIC (Carmona et al. 2011); IC 443: *Fermi* (Abdo et al. 2010e), MAGIC (Albert et al. 2007a), VERITAS (Acciari et al. 2009); CTB 37A: *Fermi* (Castro & Slane 2010), HESS (Aharonian et al. 2008a); G8.7-0.1: *Fermi* (Ajello et al. 2012), HESS



**Gamma-ray spectral variation of SNRs may be due to variation in density of the surrounding environment.**  
 Figure 5. Photon index  $\Gamma$  (between 1 GeV and 1 TeV) vs. the gas density  $n$  of the 12 SNRs studied in this work. The solid line is the model expected result.

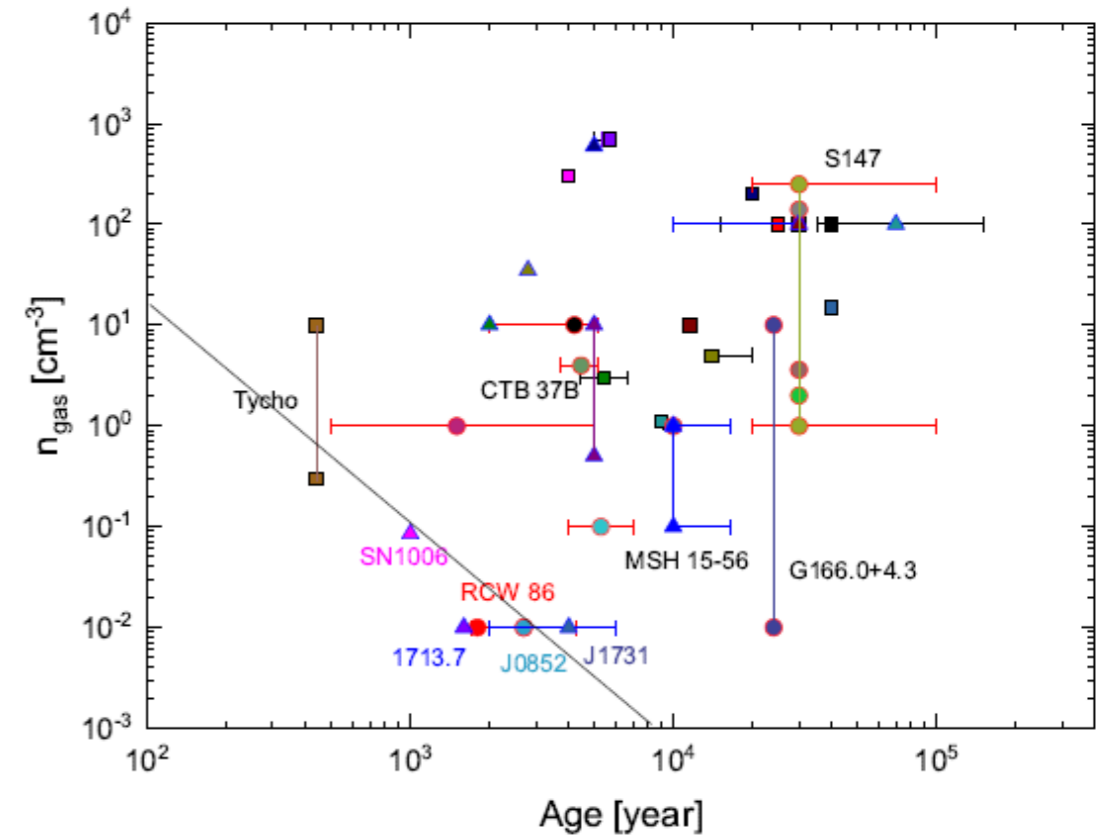
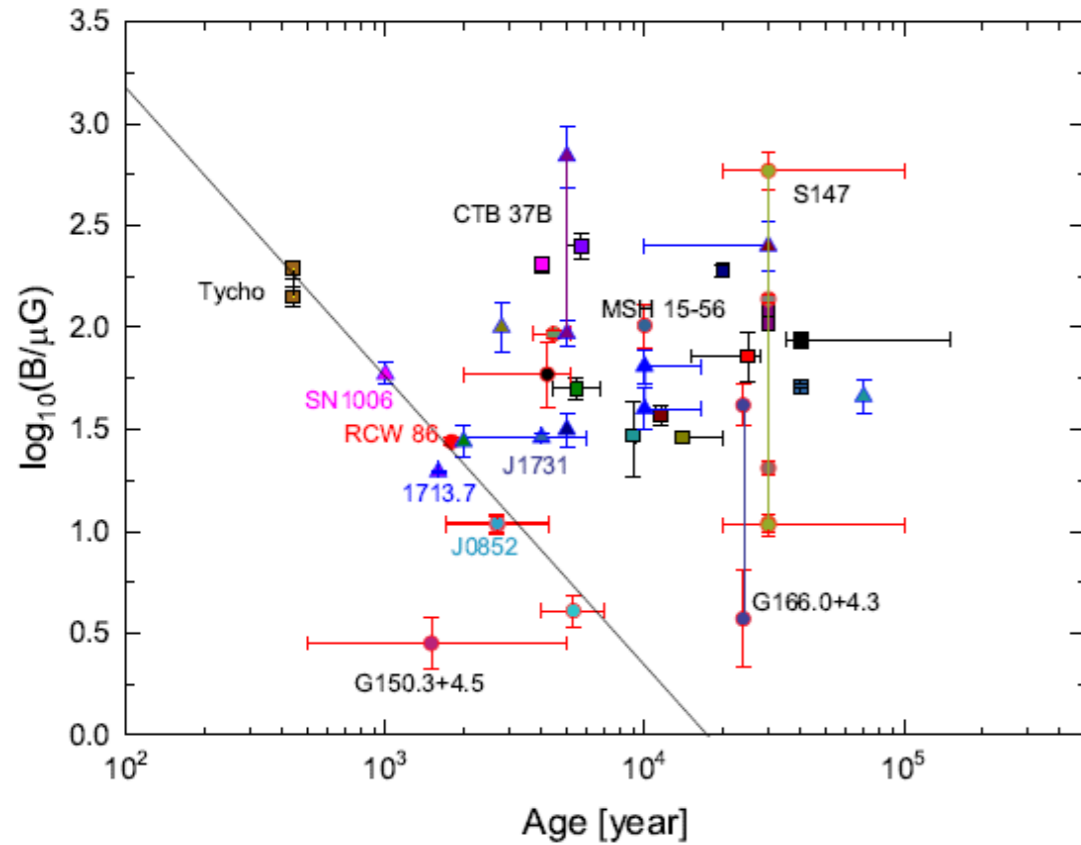
**These results can be considered as evidence for the SNR origin of Galactic cosmic rays.**

# Results



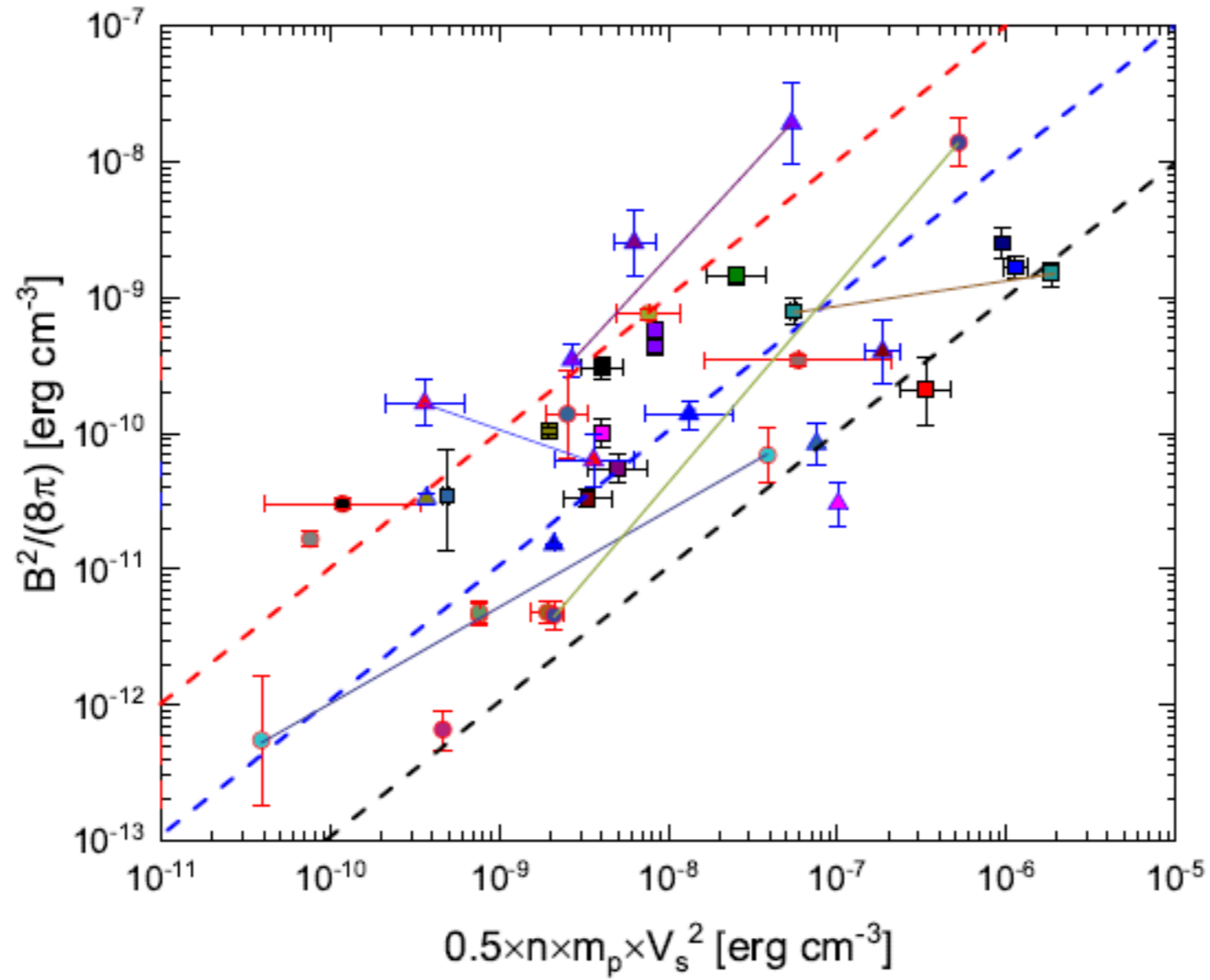
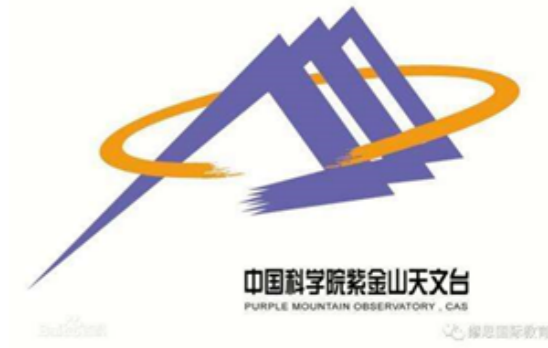
**The total energy of particles for most of SNRs are greater than erg which can be regarded as the lower limit of the cosmic rays produced by SNRs, and also supports that SNRs are the sources of Galactic cosmic rays.**

# Results

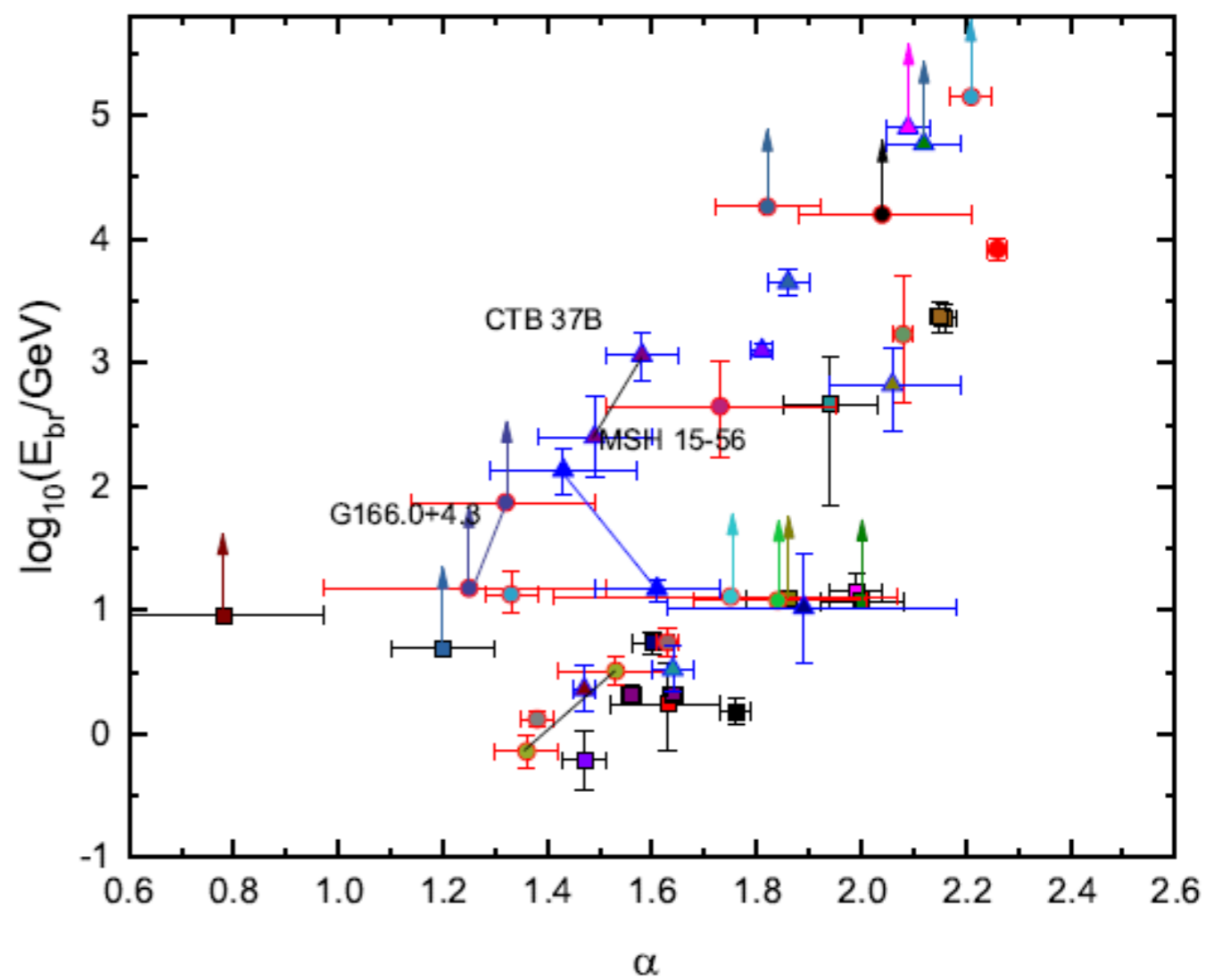
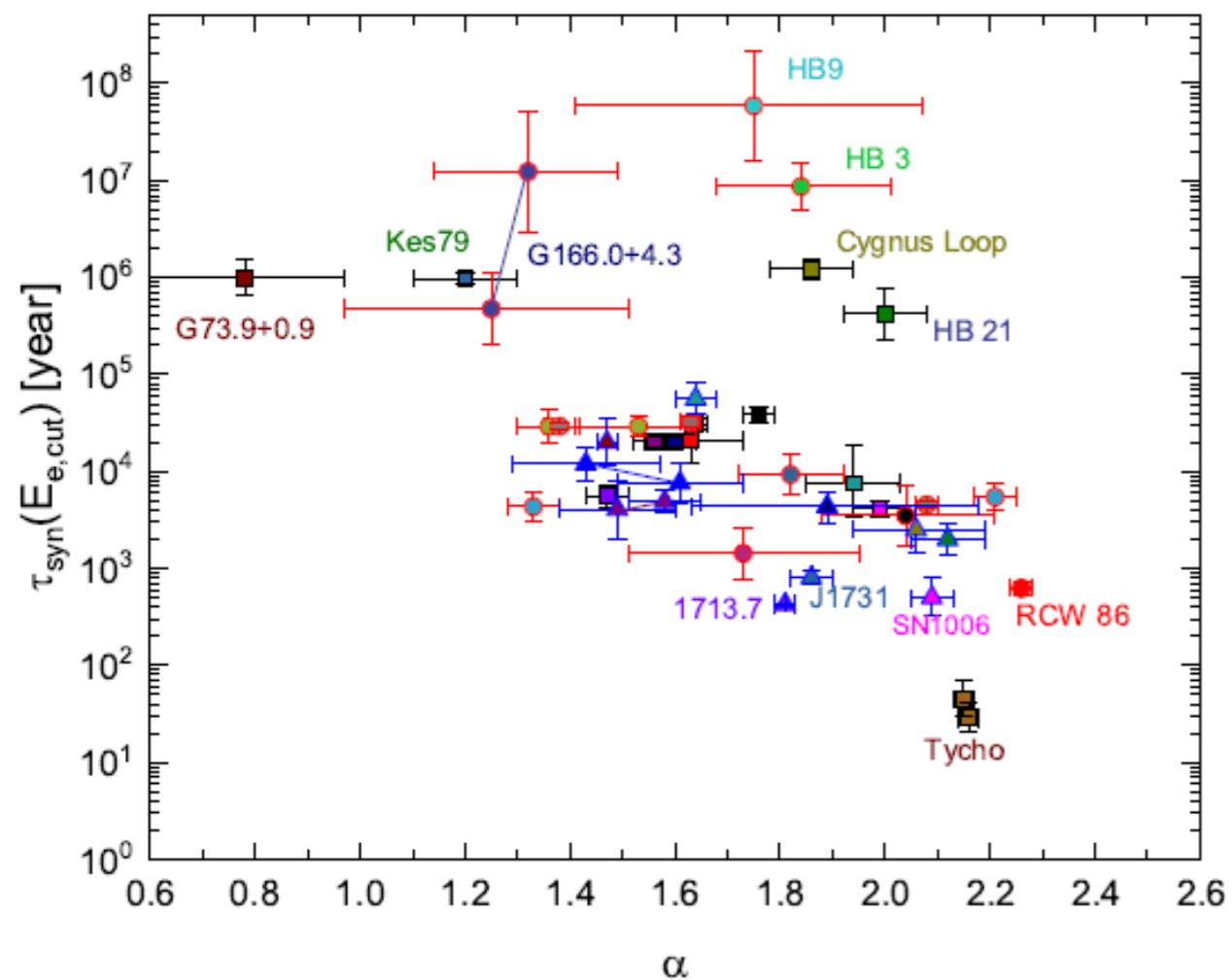


These two figures respectively show that the relationship between the magnetic field, the gas density and the age of SNRs. The results show there are no correlation between them, but an anti-correlation may be existence for several young shell-type SNRs, consistent with evolution in wind bubbles.

# Results and discussion

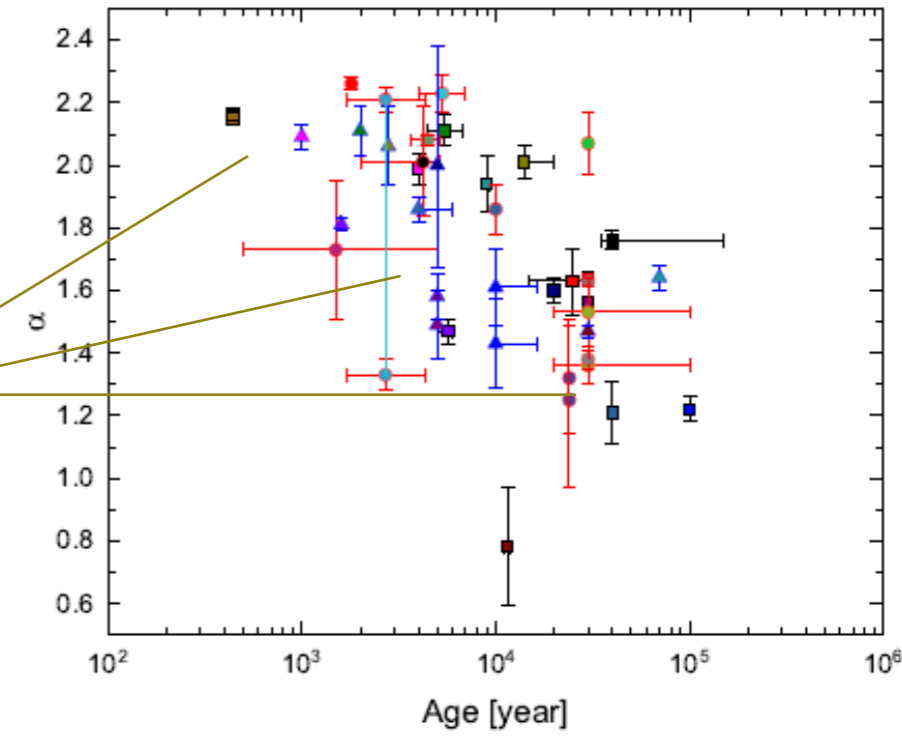
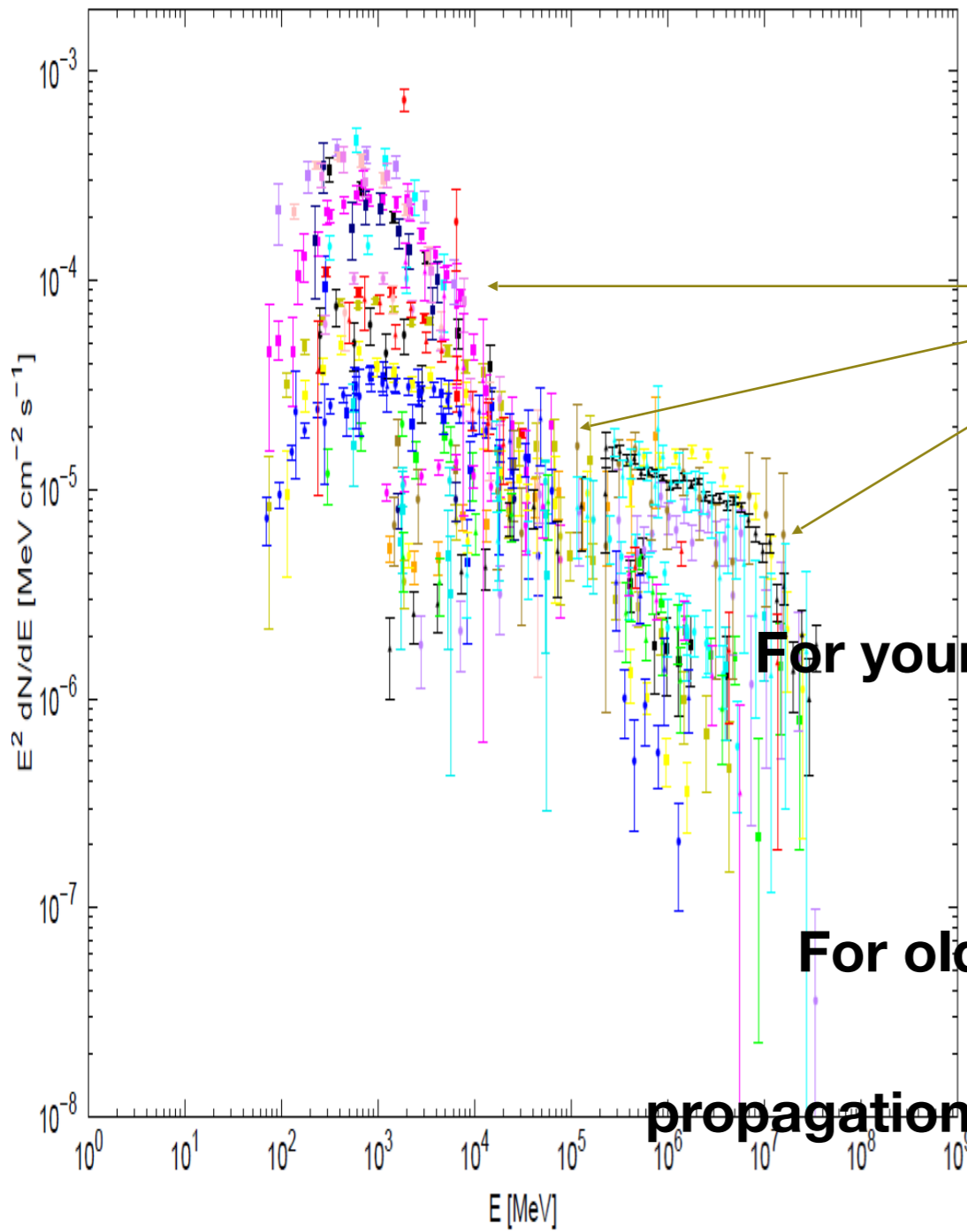


# Results



# Discussion

(Possible relation between Particles and CRs spectrum)



For young SNRs:

For old SNRs:

propagation effect 0.3-0.6

