

Preliminary

第二届北师大暗物质研讨会

# Indirect detection signal from dark matter bound state formation

Speaker: Daneng Yang ( 杨大能 )

Collaborators: Haipeng An ( 安海鹏 ), Chen Yang ( 杨晨 )

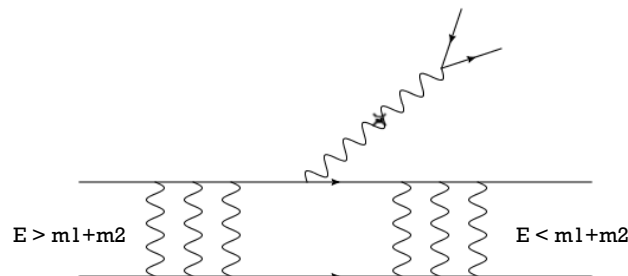
Department of Physics, Tsinghua University

2018年11月11日



# OUTLINE

- Introduction
- Model and parameter space
- Dark recombination and relic ionization fraction
- Photon spectrum from the Galactic center
- Solar signatures of the BS formation
- Summary



## 银河系中心光子在GeV存在超出

(arXiv:1704.03910 [astro-ph.HE] )

### DM annihilation benchmark

$$m_\chi = 50 \text{ GeV}$$

$$\langle\sigma v\rangle \sim 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

$$\chi\chi \rightarrow b\bar{b}$$

-----  
To obtain the same signal  
-----

能否用束缚态形成来解释? (简单估计)

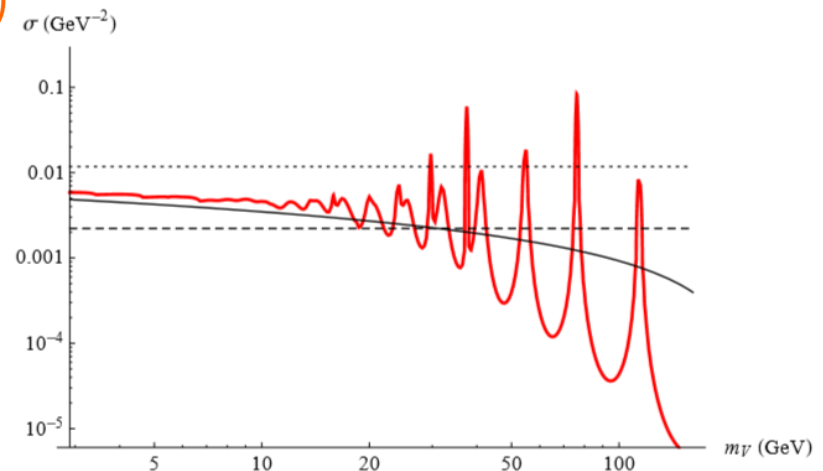
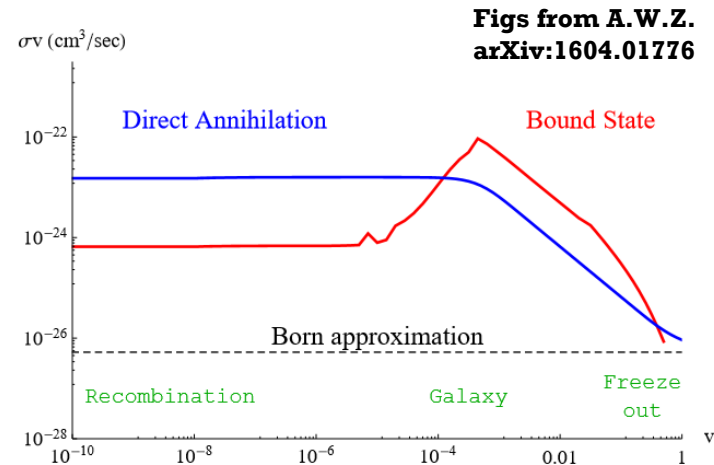
### Bound state formation benchmark

$$m_1 = m_2 = 1000 \text{ GeV}$$

$$\langle\sigma v\rangle \sim 400 \times 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

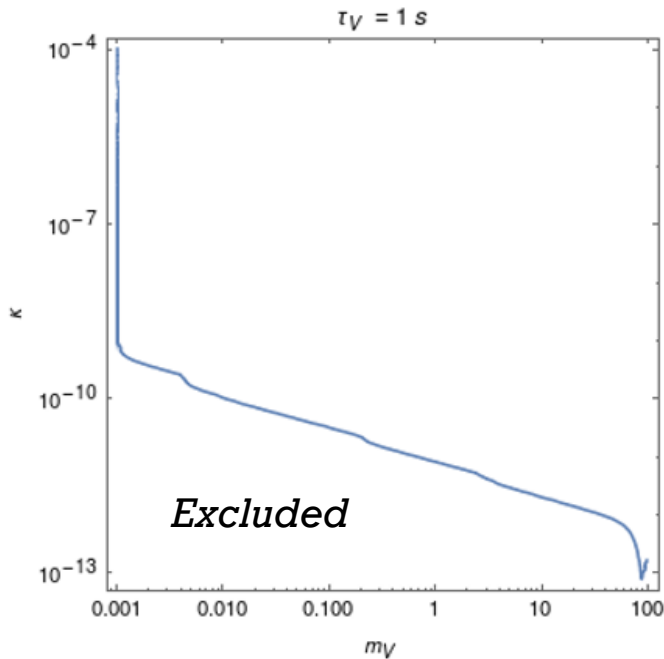
- Velocity dependence  $\sim 10$
- Mediator mass  $\sim 5$
- Other decay channels  $\sim 10$

$$\Phi_\gamma(E_\gamma, \psi) = \frac{dN_\gamma}{dE_\gamma} \frac{\langle\sigma v\rangle}{8\pi m_{\text{DM}}^2} \int_{\text{los}} \rho^2(r) dl,$$



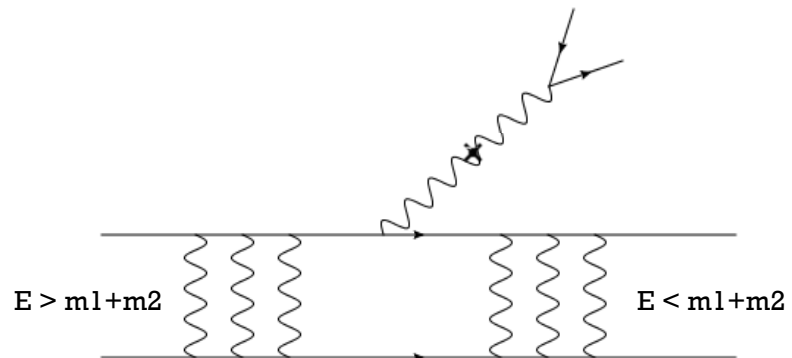
## 问题及解答

- 不考虑湮灭信号:  
Asymmetric dark matter
- 存在稳定束缚态:  
Long range force
- 束缚态在今天能形成:  
inefficient recombination
- 束缚态结合能  $\sim O(10 \text{ GeV})$ :  
 $m_1 = m_2$ , maximize
- 满足宇宙学限制  
 $O(\text{GeV})$  massive dark photon  
 $\tau_V < 1s$



## Asymmetric Dark Matter Model (ADM)

- Dark matter today is composed of particles without anti-particles
- SIDM with a primordial dark asymmetry  $\eta_D$   
(We consider two generations)
- With large enough  $\alpha_D$ , the symmetric part can be completely depleted
- Stable bound states
- **In depth studies exist for atomic dark matter, considering massless mediators**
  - **Astrophysicists:** Cyr-Racine Francis-Yan, Sigurdson Kris, Agrawal Prateek, Randall Lisa, Scholtz Jakub ...
  - **Particle physicists:** Petraki Kalliopi, von-Harling Benedict ...
  - ... ..
  - A few papers on cosmology indirect detection





# MODEL AND PARAMETER SPACE

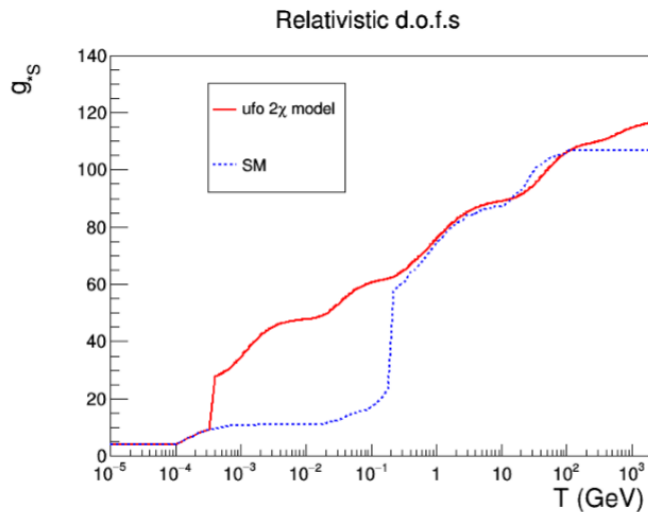
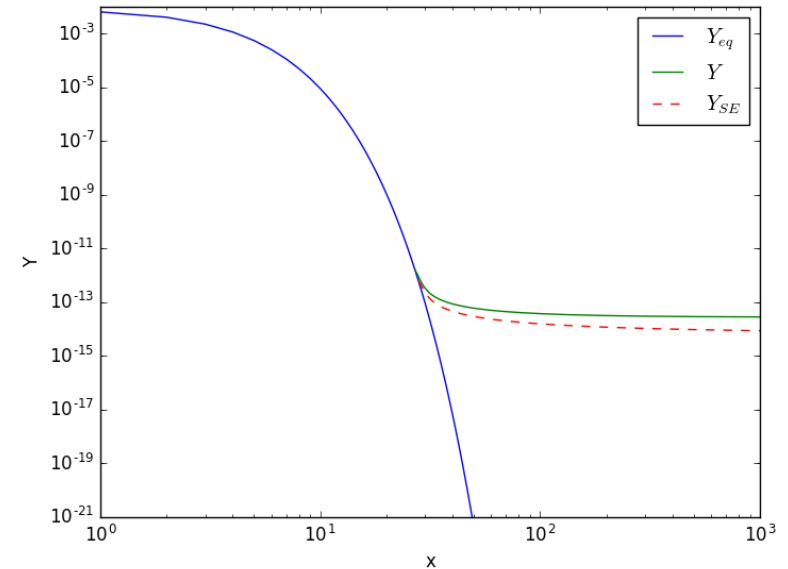
# Model and parameter space

## Dark Sector 演化

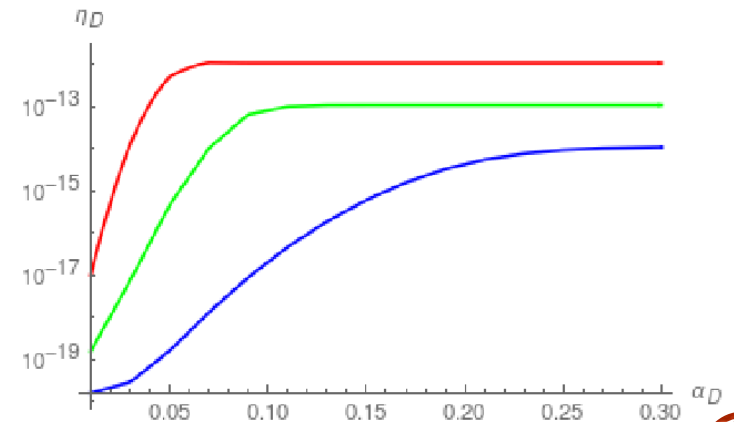
- Multi-TeV: UV model generates a primordial dark asymmetry  $\eta_D \equiv Y_\chi - Y_{\bar{\chi}}$
- Symmetric part annihilation
- Dark recombination
- Relic density composed of free  $\chi_1, \chi_2$  and  $(\chi_1 \chi_2)_{BS}$
- Bound state formation today
  - Galactic center
  - Sun
  - Dwarf galaxies

Symmetric part annihilation and freeze-out

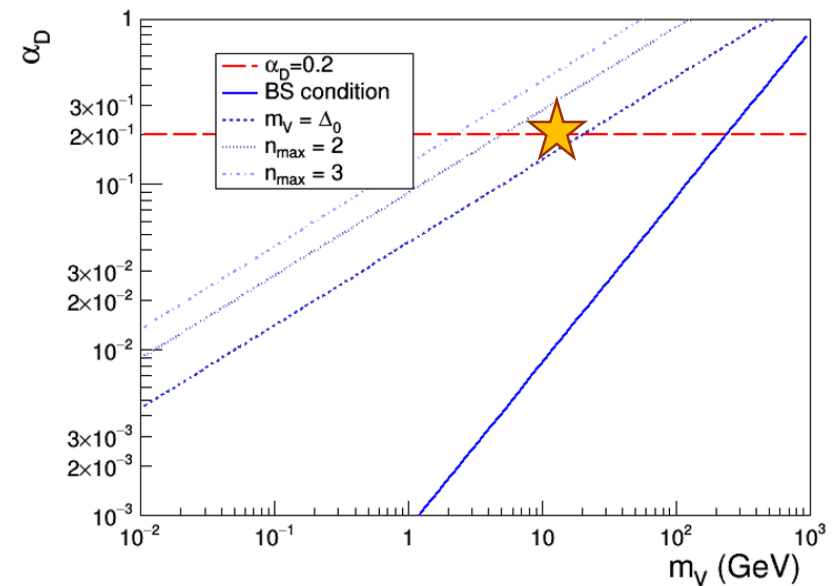
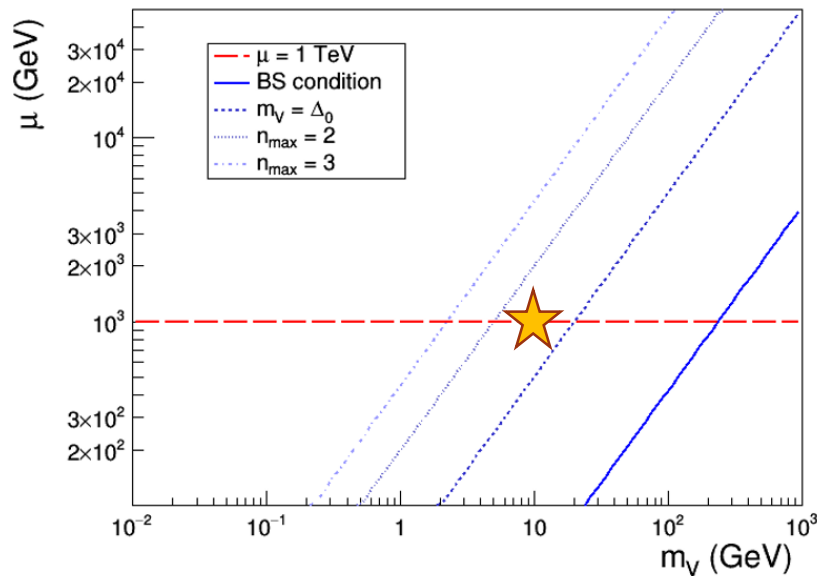
With/without Sommerfeld Enhancement



选取右下参数空间使  
得对称部分完全湮灭  
Red :  $\mu = 0.1$  TeV  
Blue :  $\mu = 1$  TeV  
Green :  $\mu = 10$  TeV



# Model and parameter space



The  $m_V$  is chosen s.t. (discuss later)

- the dark recombination is not efficient
- Bound state formation cross section is large

Size of the bound state:  $a_n = \frac{n}{\alpha_D \mu}$ .

Screening length of the Yukawa potential:  $\lambda = \frac{1}{m_V}$ .

Existence condition:  $a_n < \lambda$ .

Binding energy:  $\Delta_{nl} = E_{nl} \sim \frac{\alpha_D^2 \mu}{2n^2}$ .

Mass of bound state:  $M^{(n)} = 2m_D - \Delta_n$ .

Emitted dark photon energy:  $\omega^{(n)} = \Delta^{(n)} + \frac{\mu v_{rel}^2}{2}$ .



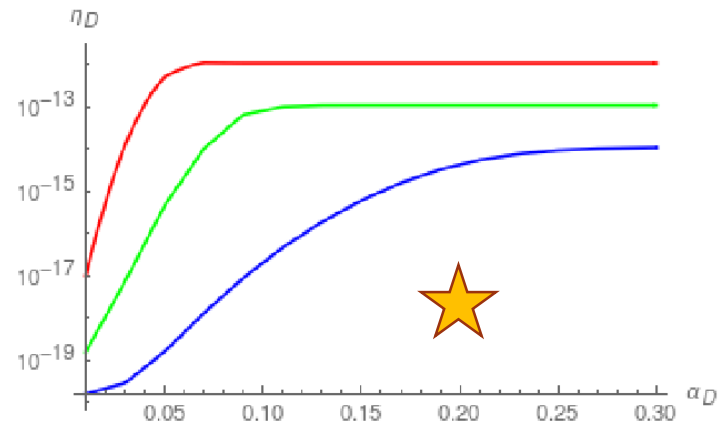
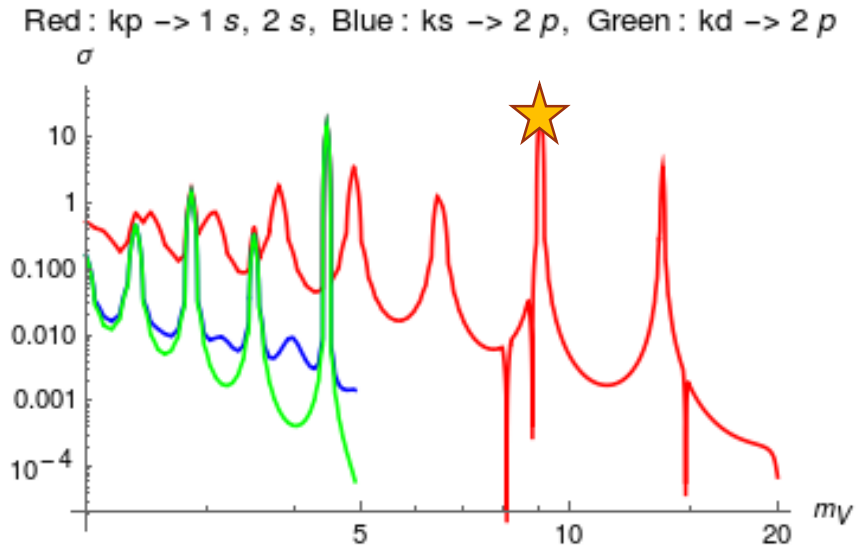
# Model and parameter space

The  $m_V$  is chosen s.t. (discuss later)

- the dark recombination is not efficient:
- Bound state formation cross section is large

$$\eta = \frac{m_p}{m_1 + m_2} \frac{\Omega_{DM}}{\Omega_b} \left( \frac{1 - r_\infty(\alpha_D, \eta)}{1 + r_\infty(\alpha_D, \eta)} \right) \eta_B$$

- $\alpha_D = 0.2$ ,
- $m_\chi = m_1 = m_2 = 2 \text{ TeV}$ ,  $\mu = 1 \text{ TeV}$ ,
- $m_V = 9.08 \text{ GeV}$ .
- $\eta_D$  chosen such that the DM relic density is completely generated by our model.
- $\kappa = 1 \times 10^{-8}$ .







# **DARK RECOMBINATION AND RELIC IONIZATION FRACTION**



# Bound state formation with a massive spin-1 mediator

In the non-relativistic limit:

$$\begin{aligned}
 H &= H_0 + H_{int}, \\
 H_0 &= -\frac{\nabla_{cm}^2}{2M} - \frac{\nabla^2}{2\mu} + V(r)_{static}, \\
 H_{int}^T &= \frac{g_D \mathbf{k}}{m_D} \left( \mathbf{A}^T\left(\frac{\mathbf{r}}{2}\right) + \mathbf{A}^T\left(-\frac{\mathbf{r}}{2}\right) \right), \\
 H_{int}^L &= -g_D \frac{m_V}{|\mathbf{q}|} \left( \phi_L\left(\frac{\mathbf{r}}{2}\right) - \phi_L\left(-\frac{\mathbf{r}}{2}\right) \right).
 \end{aligned}$$

The bound and scattering wavefunctions that solve the Schrodinger equation are written as

$$\begin{aligned}
 \psi_n(\mathbf{r}) &= \sum_{\ell m} R_{n\ell}(r) Y_{\ell m}(\hat{r}), \\
 \psi_k(\mathbf{r}) &= \sum_{\ell m} R_{k\ell}(r) Y_{\ell m}(\hat{r}) Y_{\ell m}^*(\hat{k})
 \end{aligned}$$

Bound state formation cross section as a function of  $m_V$ .

**Red line:**  $kp \rightarrow 1, 2s$ .

**Blue line:**  $ks \rightarrow 2p$ .

**Green line:**  $kd \rightarrow 2p$ .

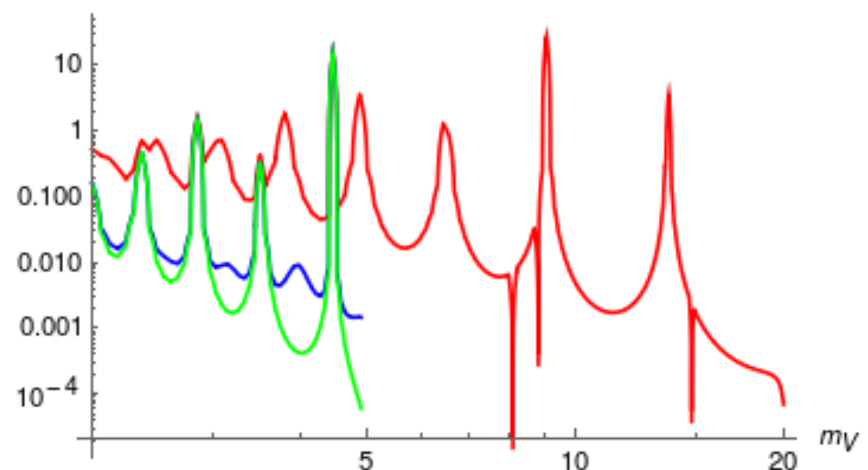
The amplitudes are calculated using time dependent perturbation theory, with dipole approximation

$$c_T = g_D \int d^3\mathbf{r} \Psi_f^*(\mathbf{r}) (E_i - E_f) \mathbf{r} \cdot \epsilon_T \Psi_i(\mathbf{r}),$$

$$c_L = -ig_D \int d^3\mathbf{r} \Psi_f^*(\mathbf{r}) m_V \frac{\mathbf{r} \cdot \mathbf{q}}{|\mathbf{q}|} \Psi_i(\mathbf{r}),$$

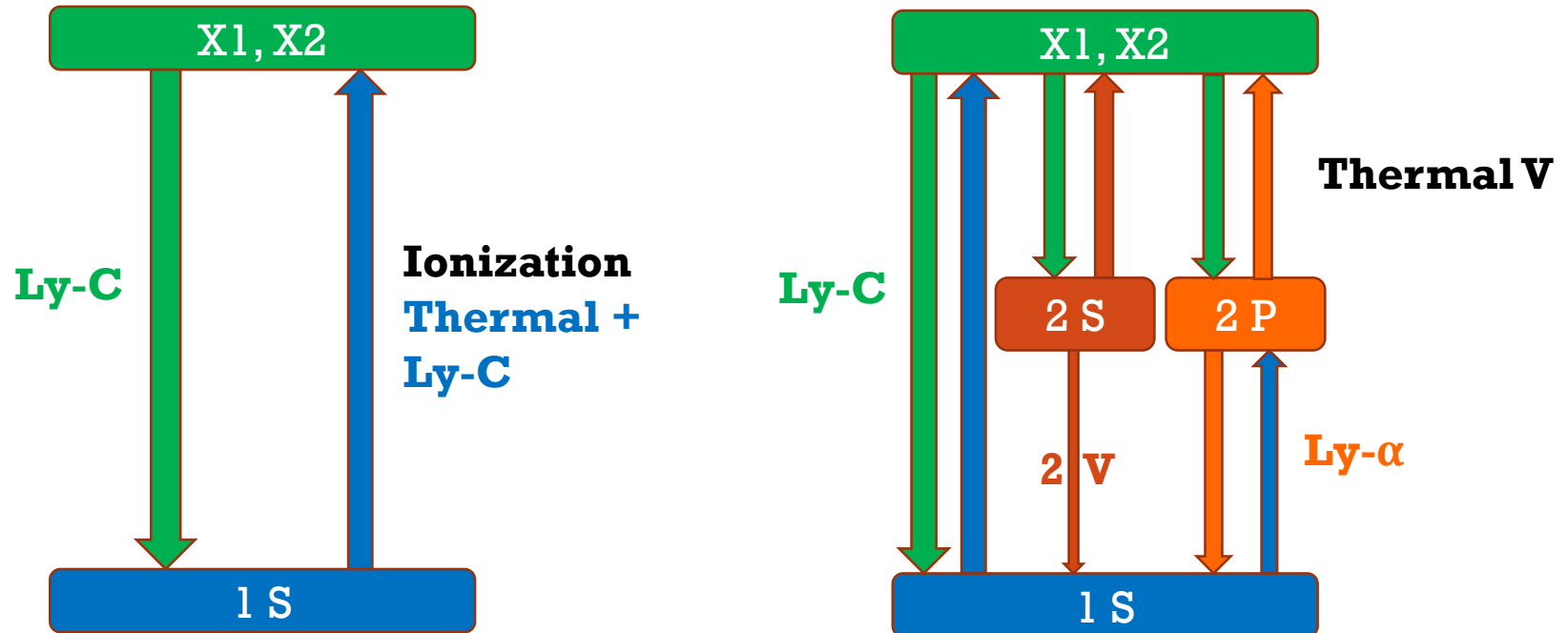
$$\begin{aligned}
 (\sigma v)_B &= \frac{\alpha_D}{3\pi} \sum_{n,l} \left( \omega_{nl}^2 + \frac{1}{2} m_V^2 \right) \sqrt{\omega_{nl}^2 - m_V^2} \\
 &\times \left[ l \left| \int dr r^3 R_{n,l}(r) R_{k,l-1} \right|^2 + (l+1) \left| \int dr r^3 R_{n,l}(r) R_{k,l+1} \right|^2 \right]
 \end{aligned}$$

Red :  $kp \rightarrow 1s, 2s$ , Blue :  $ks \rightarrow 2p$ , Green :  $kd \rightarrow 2p$



## Dark recombination and relic ionization fraction

- Temperature of thermal photons  $\ll$  Binding energy  $\Delta$
- $2S \rightarrow 1S$  dominant in the standard Hydrogen recombination
- Ly-Continuum photon are non-thermal, can ionize another bound state



Our present benchmark only allows for  $n=1$  bound state

## Dark recombination and relic ionization fraction

Assuming that our dark sector constitutes the dark matter energy density today

$$n_c = 2n_{BS} + n_{D1} + n_{D2}$$

$$s(T) = \frac{2\pi^2}{45} g_{\star S}(T) T^3,$$

$$n_c = \frac{n_c(T_0)}{n_B(T_0)} \eta_\gamma \frac{n_\gamma(T_0)}{s(T_0)} s(T).$$

$$X_D \equiv \frac{n_D}{n_c} \quad x \equiv \frac{\Delta}{T}$$

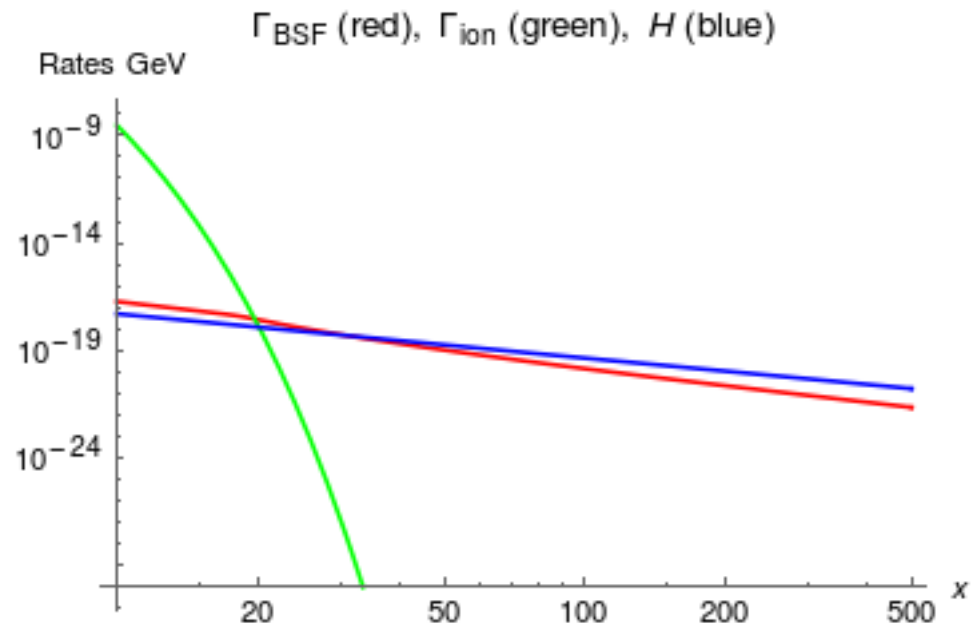
*We hope to have a significant relic  $XD$  s.t. the BS can still form efficiently today inside galaxies*

- Large formation and ionization rate

The ionization cross section can be obtained from the principle of detailed balance

$$(\sigma v)_B n_1^{eq} n_2^{eq} = g_V (\sigma v)_I n_{BS}^{eq} n_V^{eq}$$

$$\Gamma_I(T) = \langle \sigma v \rangle_I (n_V^{th} + n_V^{nt})$$



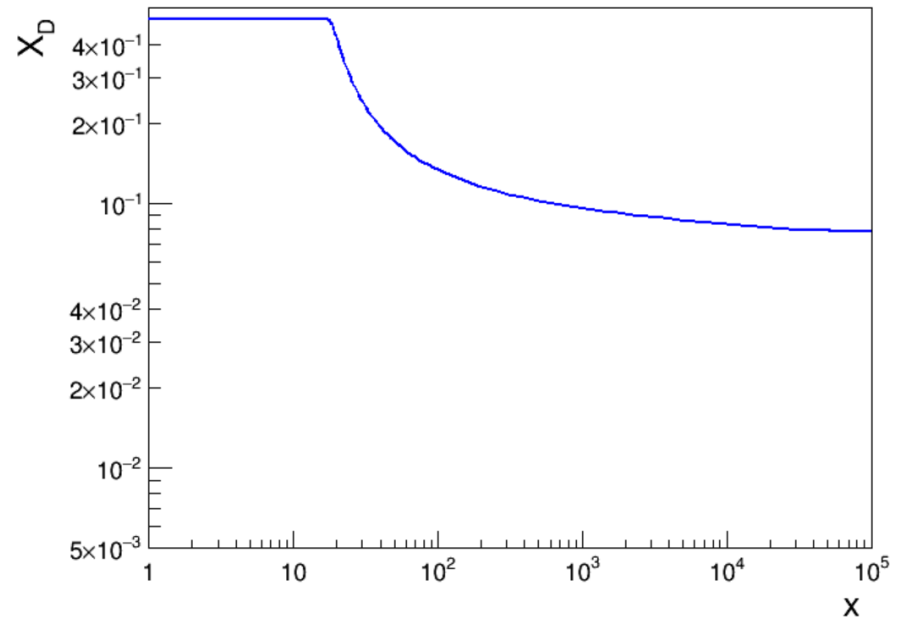
# Dark recombination and relic ionization fraction

Assuming that our dark sector constitutes the dark matter energy density today

$$n_c = 2n_{BS} + n_{D1} + n_{D2}$$

$$X_D \equiv \frac{n_D}{n_c} \quad x \equiv \frac{\Delta}{T}$$

- The dark recombination happened at GeV scale
- The recombination is slow
- Ionization fraction can be tuned to order 1: ~20% free states
- Non-thermal photons can prolong the recombination, but leads to negligible change in the relic ionization fraction



$$\frac{d(n_V^{nt}/n_c)}{dx} = \frac{x}{H(\Delta)} \left[ \langle \sigma v \rangle_B n_c X_D^2 - \langle \sigma v \rangle_I n_V^{nt} \left( \frac{1}{2} - X_D \right) \right]$$

$$\frac{dX_D}{dx} = -\frac{x n_c}{H(\Delta)} \langle \sigma v \rangle_B X_D^2 + \frac{x}{H(\Delta)} \Gamma_I \left( \frac{1}{2} - X_D \right)$$



# **GALACTIC CENTER GEV EXCESS**



## Photon flux from the BS formation

- We obtain the photon spectrum using the PPC program
- The predicted Gamma Ray signal from the DM is strongest in the GC owing to its proximity and the enhanced density of the DM.
- J-factor taken from Ref. [*Astrophys. J.* 840, 43 (2017), 1704.03910], (calculated assuming a local DM density of 0.4 GeV/cm<sup>3</sup>)

$$1.53 \times 10^{22} \text{ GeV}^2 \text{ cm}^{-5} \quad \gamma = 1$$

$$\begin{aligned} \frac{d\Phi}{dE_\gamma}(E_\gamma, \theta) &= \frac{1}{4\pi} \int_{\Delta\Omega} d\Omega \int_{l.o.s} ds \sum_{n,l} \left( \frac{\rho(r(s, \theta))}{m_1 + m_2} 2X_D \right)^2 \langle \sigma v \rangle_{n,l} \times \frac{dN_\gamma^{n,l}}{dE_\gamma} \\ &= \frac{X_D^2}{\pi} \sum_{n,l} \frac{\langle \sigma v \rangle_{n,l}}{(m_1 + m_2)^2} \frac{dN_\gamma^{n,l}}{dE_\gamma} \\ &\quad \times \int_{\Delta\Omega} \int_{l.o.s} ds d\Omega \rho^2(r(s, \theta)), \\ J\text{-factor} &= \int_{\Delta\Omega} \int_{l.o.s} ds d\Omega \rho^2(r(s, \theta)), \\ \frac{dN_\gamma^{n,l}}{dE_\gamma} &= \sum_f \text{Br}(V \rightarrow f\bar{f}) \frac{dN_\gamma^f}{dE_\gamma}. \end{aligned}$$

$$\begin{aligned} \rho_{DM}(r) &= \frac{\rho_0}{\left(\frac{r}{r_s}\right)^\gamma \left(1 + \frac{r}{r_s}\right)^{3-\gamma}} \\ f(v) &= \frac{1}{N_{esc}} \left(\frac{3}{2\pi\sigma_v^2}\right)^{\frac{3}{2}} e^{-\frac{3v^2}{2\sigma_v^2}}, \quad |v| < v_{esc} \\ f(v) &= 0: \text{ otherwise} \\ N_{esc} &= \text{erf}(z) - \frac{2}{\sqrt{\pi}} z e^{-z^2}, \\ z &\equiv \frac{v_{esc}}{\sigma_v} \sqrt{\frac{3}{2}}, \\ v_{esc} &= 553 \text{ km/s}, \\ \sigma_v &= 150 \text{ km/s}, \\ \sigma_r &= \sqrt{2}\sigma_v. \end{aligned}$$



# Model and parameter space

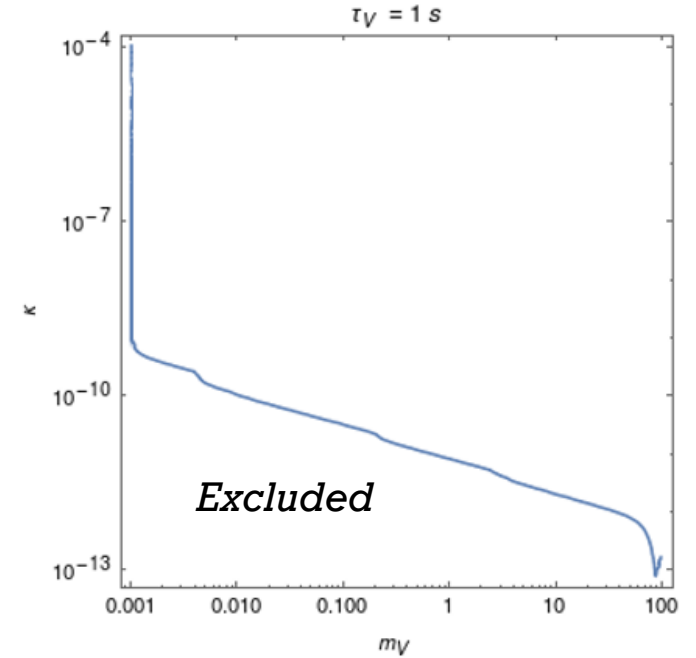
## V boson life time

We choose kappa and mV s.t. there is no impact on the BBN, CMB, and structure formation...

$$\frac{\kappa}{2 \cos \theta_W} B_{\mu\nu} V^{\mu\nu} = \frac{\kappa}{2} F_{\mu\nu} V^{\mu\nu} + \frac{\kappa \tan \theta_W}{2} Z_{\mu\nu} V^{\mu\nu}$$

$$\Gamma_{V \rightarrow f \bar{f}} = \frac{1}{2m_V} N_c^f \frac{1}{8\pi} \sqrt{1 - \frac{4m_f^2}{m_V^2}} \int \frac{d \cos \theta}{2} |\mathcal{M}_{L/R}(\cos \theta)|^2$$

$$= \frac{N_c}{24\pi m_V} \sqrt{1 - \frac{4m_f^2}{m_V^2}} (m_V^2 (g_L^2 + g_R^2) - m_f^2 (-6g_L^2 g_R^2 + g_L^2 + g_R^2))$$



Decay channel	Branching ratio (MadWidth)	Branching ratio (analytic)	Mass (MeV)	PDG ID1	PDG ID2
$u\bar{u}$	0.204	0.215901	2.2	2	-2
$c\bar{c}$	0.204	0.209764	1280	4	-4
$e^+e^-$	0.153	0.161767	0.511	11	-11
$\mu^+\mu^-$	0.153	0.161767	106	13	-13
$\tau^+\tau^-$	0.134	0.14253	1777	15	-15
$d\bar{d}$	0.0511	0.0541345	4.7	1	-1
$s\bar{s}$	0.0511	0.0541345	96	3	-3
$b\bar{b}$	0.0511	0	4180	5	-5
$\nu_e \bar{\nu}_e$	2.953107e-07	3.11121e-07	0	12	-12
$\nu_\mu \bar{\nu}_\mu$	2.953107e-07	3.11121e-07	0	14	-14
$\nu_\tau \bar{\nu}_\tau$	2.953107e-07	3.11121e-07	0	16	-16

TABLE II: Branching ratios for  $m_V = 5$  GeV. The MadWidth considers only  $\tau, t$  masses. They are listed here for validation.

## Photon flux from the BS formation

Validation, spectrum from Pythia (Left) and from PPPC (right)  
Electroweak correction has been included in the latter case.

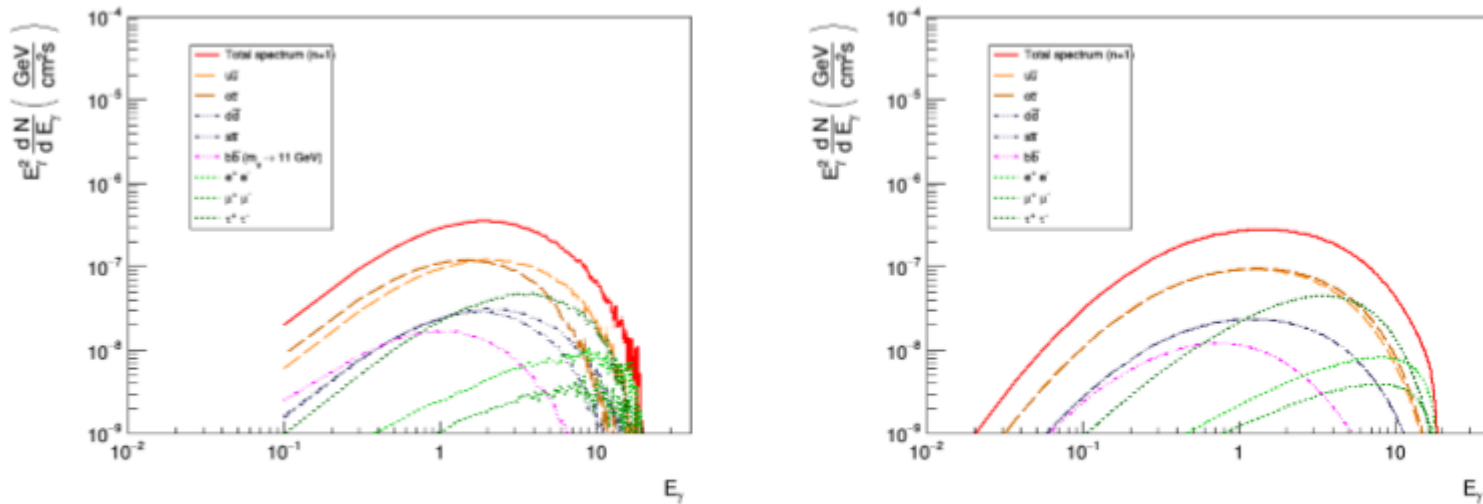
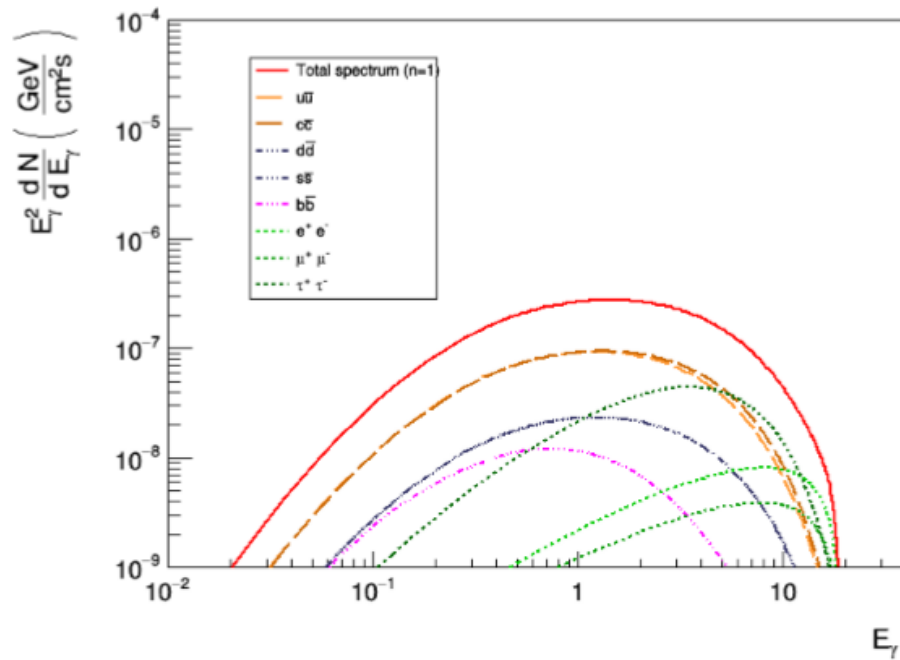


FIG. 16: Comparison of photon energy spectrum between pythia and pppc. The  $b\bar{b}$  mode does not appear in pythia until about 11 GeV.

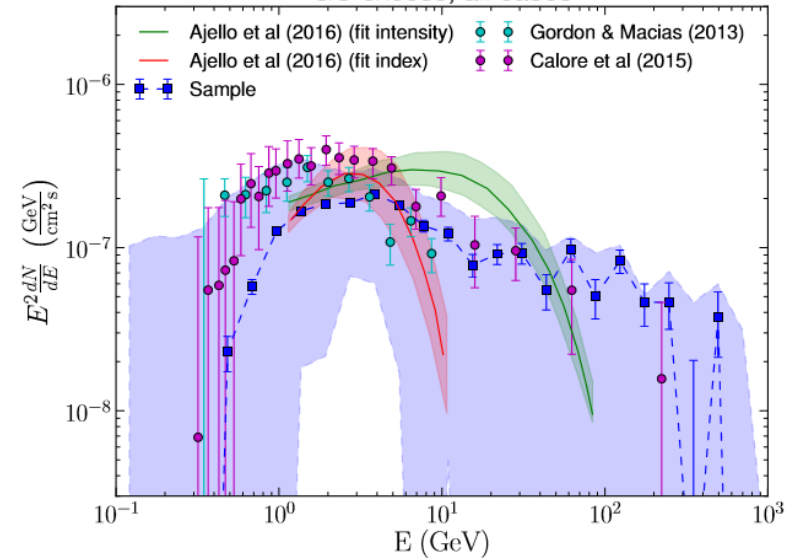
# Photon flux from the BS formation

BS formation can generate a photon flux as strong as the GC gamma-ray excess

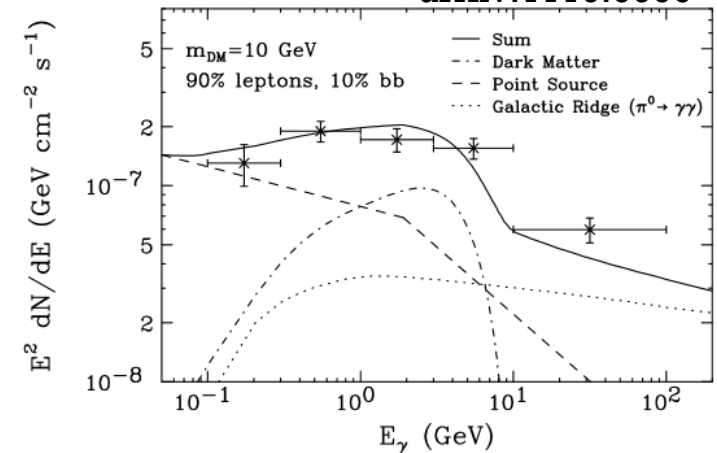


arXiv:1704.03910 [astro-ph.HE]

GC excess, all cases

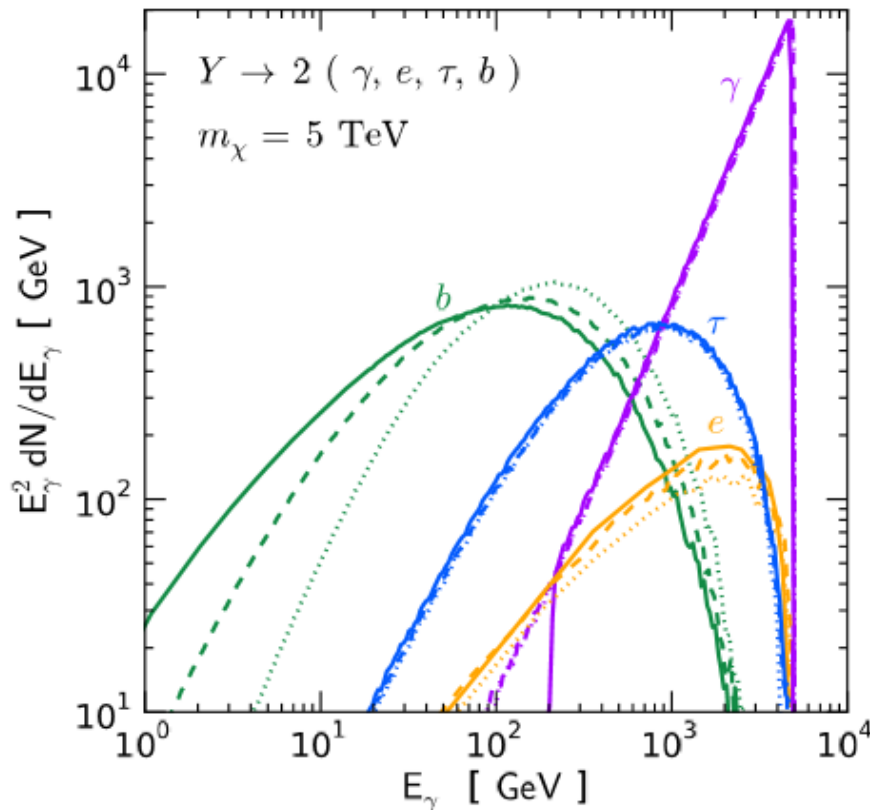


arXiv:1110.0006

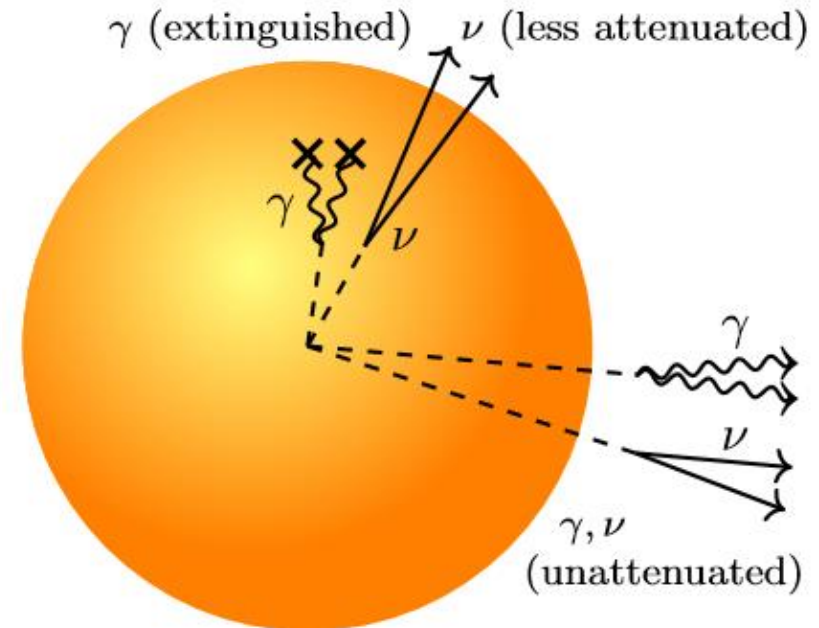


Figs from L.N.B arXiv: 1703.04629:

## Powerful Solar Signatures of Long-Lived Dark Mediators



Gamma-ray spectra for various final states, per DM annihilation, with mediator masses  $m_Y = 2 \text{ TeV}$  (solid),  $m_Y = 200 \text{ GeV}$  (dashed), and  $m_Y = 20 \text{ GeV}$  (dotted)



Long-lived mediators

- BBN generally require  $\tau < 1 \text{ s}$
- Solar Radius =  $2.32 \text{ s}$
- Consider  $\bar{p}, \bar{e}$  spectrum to reduce backgrounds

## SUMMARY

- Asymmetric Dark Matter model provides a possible way to generate the Galactic Center GeV photon excess through the BS formation process.
- The relic dark recombination can be made inefficient, making the bound state formation possible today.
- If the dark photon has a relatively long life time, there are interesting solar signatures.