

# Non-minimal interactions of Higgs boson

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# Do we live at an exciting time as particle physicist?

- We found Higgs boson  $m_h = 125$  GeV.
- We only found Higgs boson.

Is there something special about the Higgs boson from a gravitational point of view?

gravity might help us control hierarchy problem!

Three Generations of Matter (Fermions) spin 1/2

	I		II		III		
mass	2.4 MeV		1.27 GeV		171.2 GeV		0
charge	2/3		2/3		2/3		0
name	u		c		t		g
	Left	Right	Left	Right	Left	Right	gluon
	up		charm		top		
	d		s		b		γ
	Left	Right	Left	Right	Left	Right	photon
	down		strange		bottom		
	ν <sub>e</sub>		ν <sub>μ</sub>		ν <sub>τ</sub>		Z
	Left	Right	Left	Right	Left	Right	91.2 GeV
	electron neutrino		muon neutrino		tau neutrino		weak force
	e		μ		τ		H
	Left	Right	Left	Right	Left	Right	>114 GeV
	electron		muon		tau		Higgs boson
	0.511 MeV		105.7 MeV		1.777 GeV		spin 0
	-1		-1		-1		
	e		μ		τ		W
	Left	Right	Left	Right	Left	Right	80.4 GeV
	electron		muon		tau		weak force

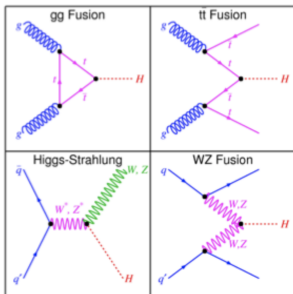
Bosons (Forces) spin 1

Einstein gravity

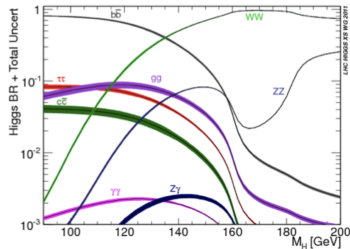


# It is the Higgs boson that we all know

- The Standard Model predicts precisely how the Higgs boson should be produced.



- The Standard Model predicts precisely how the Higgs boson should decay.



# Something special about the Higgs boson

Einstein-Hilbert action

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} (R - 2\Lambda) + \xi R H^\dagger H \right]$$

Non-minimal coupling: required by renormalisation

Is there any bound on its value?

This is a dimension-4 operator: it is a fundamental constant of nature.

# The non-decoupling effect

- Consider Standard Model with a non-minimal coupling to R in Jordan frame

$$S_J = \int d^4x \sqrt{-g^J} \left( \frac{1}{2} M^2 + \xi H^\dagger H \right) R^J - \frac{1}{4} \sum_j F_{\mu\nu j} F_j^{\mu\nu} + (D_\mu H)^\dagger (D^\mu H) - V(H)$$

Gravity is not canonical

Equation of motion after differentiating with action is

$$\left\{ R_{\mu\nu}^J - \frac{1}{2} g_{\mu\nu}^J R^J + \left( g_{\mu\nu}^J \Delta_\lambda \Delta^\lambda - \Delta_\mu \Delta_\nu \right) \right\} \Omega^2 = -\frac{1}{M_P^2} T_{\mu\nu}$$

# Obtain a canonical Einstein equation

## Weyl transformation

We go from the Jordan frame to the Einstein frame

$$g_{\mu\nu}^J = \Omega^{-2} g_{\mu\nu}$$

$$g^{\mu\nu J} = \Omega^2 g^{\mu\nu}$$

$$\sqrt{-g^J} = \Omega^{-4} \sqrt{-g}$$

$$R^J = \Omega^2 \left( R - 6\Delta^\lambda \Delta_\lambda \log \Omega + 6g^{\mu\nu} \Delta_\mu \log \Omega \Delta_\nu \log \Omega \right)$$

## Scale factor

$$\Omega^2 = \frac{M^2 + 2\xi H^\dagger H}{M_P^2}$$

# Non-decoupling effect in Einstein frame

graviton is now canonical with higgs not being canonical.

Here is the action after Weyl transformation

$$S_E = - \int d^4x \sqrt{-\tilde{g}} \left[ \frac{1}{2} M_P^2 \tilde{R} - \frac{3\xi^2}{M_P^2 \Omega^4} \partial^\mu (H^\dagger H) \partial_\mu (H^\dagger H) \right. \\ \left. - \frac{1}{\Omega^2} (D^\mu H)^\dagger (D_\mu H) + \frac{\mathcal{L}_{SM}}{\Omega^4} \right]$$

One notices that the Higgs boson kinetic term is not canonically normalized. We need to diagonalize this term.

- use the unitary gauge

$$H = \frac{1}{\sqrt{2}} (0, h + v)^\top$$

- The Planck mass is defined by

$$(M^2 + \xi v^2) = M_P^2$$

Additional transformation?

Higgs doublet must be transformed too

# Ways Forward: canonical normalization

- To diagonalize the Higgs boson kinetic term

$$\frac{d\chi}{dh} = \sqrt{\frac{1}{\Omega^2} + \frac{6\xi^2 v^2}{M_P^2 \Omega^4}}$$

- To leading order in  $\Omega^{-1}$

$$h = \frac{1}{\sqrt{1 + \beta}} \chi \quad \beta = 6\xi^2 v^2 / M_P^2$$

- The couplings of the Higgs boson to particles of the SM are rescaled!

$$yh\bar{\psi}\psi \rightarrow \frac{y}{\sqrt{1 + \beta}} \chi \bar{\psi}\psi$$



# Jordan Frame Detection

Schematic form of the Jordan frame equation

The decoupling can also be seen in the Jordan frame

$$\mathcal{L}^{(2)} = -\frac{M^2 + \xi v^2}{8} (h^{\mu\nu} \square h_{\mu\nu} + 2\partial_\nu h^{\mu\nu} \partial^\rho h_{\mu\rho} - 2\partial_\nu h^{\mu\nu} \partial_\mu h^\rho_\rho - h^\mu_\mu \square h^\nu_\nu) + \frac{1}{2} (\partial_\mu h)^2 + \xi v (\square h^\mu_\mu - \partial_\mu \partial_\nu h^{\mu\nu}) h$$

same renormalization factor!

$$h = \frac{1}{\sqrt{1 + \beta}} \chi$$
$$h_{\mu\nu} = \frac{1}{M_P} \tilde{h}_{\mu\nu} - \frac{2\xi v}{M_P^2 \sqrt{1 + \beta}} \bar{g}_{\mu\nu} \chi$$

# Bound on the non-minimal coupling from the LHC

- The LHC experiments produce fits to the data assuming that all Higgs boson couplings are modified by a single parameter

$$\kappa = 1/\sqrt{1 + \beta}$$

- In the narrow width approximation, one finds

$$\begin{aligned}\sigma(ii \rightarrow h \rightarrow ff) &= \sigma(ii \rightarrow h) \cdot \text{BR}(h \rightarrow ff) \\ &= \kappa^2 \sigma_{\text{SM}}(ii \rightarrow h) \cdot \text{BR}_{\text{SM}}(h \rightarrow ff)\end{aligned}$$

# Bound on the non-minimal coupling from the LHC

- The signal strength of a single Higgs boson in a specific analysis

$$\mu = \sum_i c_i \omega_i$$

- For  $i$ -th channel, the signal strength can be calculated

$$c_i = \frac{[\sigma \times BR]_i}{[\sigma_{SM} \times BR_{SM}]_i}, \quad \omega_i = \frac{\epsilon_i [\sigma_{SM} \times BR_{SM}]_i}{\sum_j \epsilon_j [\sigma_{SM} \times BR_{SM}]_j}$$

# Actual bound from numerical analysis

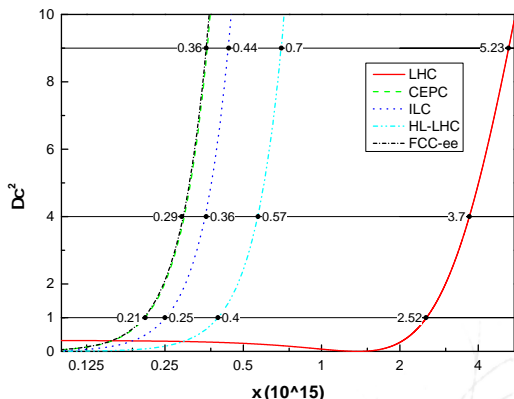


Figure: Bounds on non-minimal coupling  $\xi$  ( $\times 10^{15}$ ) from the LHC (run1+run2), HL-LHC(14 TeV,  $3 \text{ fb}^{-1}$ ), ILC(250 GeV,  $2 \text{ fb}^{-1}$ ), FCC-ee (240 GeV,  $5 \text{ fb}^{-1}$ ) and CEPC(240 GeV,  $5 \text{ fb}^{-1}$ ).

# Unitary constraint on cut-off

Does the theory become unnatural at the scale  $\Lambda = M_P/\xi$

In terms of power counting and unitary constraint, one gets the cut-off  $\Lambda = M/\xi$

- For larger  $\xi$  satisfying the bound, we will have a much lower cut-off i.e.  $\Lambda = 10^3$  GeV

Does our bound make sense?

Do we need new physics above the bound like supersymmetry or strong coupling?

# Self-healing of unitarity

There is no need to worry about cut-off problem!

- Perturbative unitarity of the S-matrix implies at one-loop.

$$|T^{\text{tree}}|^2 = \text{Im}(T^{1\text{-loop}})$$

- This is optical theorem and is a test of unitarity
- Look at the gravitational scattering between the Higgs boson and pseudo-Goldstone particles.

If so

New physics is not needed above  $M/\xi$ . The actual cut-off can be arbitrary as Plank scale.

- at tree-levels

$$A_{\text{tree}} = \frac{8\pi G_N(\bar{\phi})}{s} [s^2 (6\xi_1\xi_2 + \xi_1 + \xi_2) + ut]$$

- At one-loop

$$A_{1\text{-loop}} = -\frac{G_N^2(\phi)}{15} [s^2 F(\xi_A, \xi_B, \xi) - ut] \log(-s)$$

- we thus verify the cutting relation implied by unitarity

$$|a_{0,\text{tree},\text{max}}|^2 = \text{Im}(a_{0,1\text{-loop},\text{max}})$$

# Implication on naturalness

- Since cut-off can be very high, Higgs naturalness problems re-emerges!
- non-minimal coupling  $\xi$  provides a possible method to solve hierarchy problem.
- at spontaneously broken phase.

$$v^2 = \frac{-m_h^2 M_P^2 - 4V_0 \xi}{M_P^2 \lambda_h + m_h^2 \xi}$$

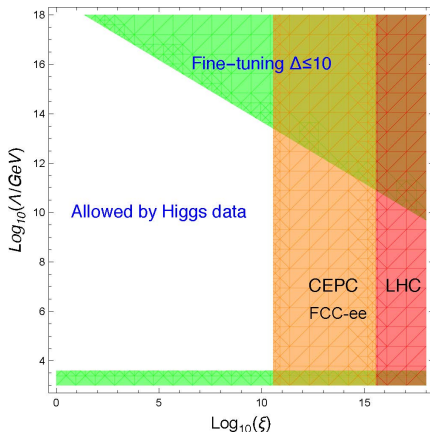
we have the one-loop correction to vev

$$\delta v^2 \sim \left(1 - \frac{\Lambda^2}{M_P^2} \xi\right) \Lambda^2$$

# Fine-tuning measure

Here we define a dimensionless fine-tuning measure

$$\Delta = \delta v^2 / v^2$$





# Conclusions

- Large non-minimal coupling of higgs boson to gravity can be bounded by LHC data.
- Large non-minimal coupling of higgs boson to gravity does not suffer from unitarity constraints.
- Large non-minimal coupling of higgs boson to gravity can soften the fine-tuning problem of itself.