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Implication of Chiral Symmetry on Neutral Weak Pion Production off a Nucleon

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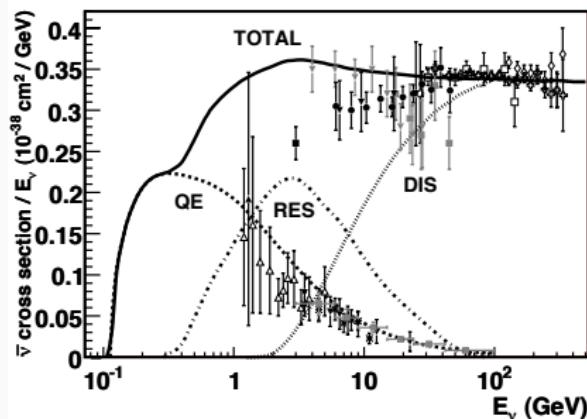
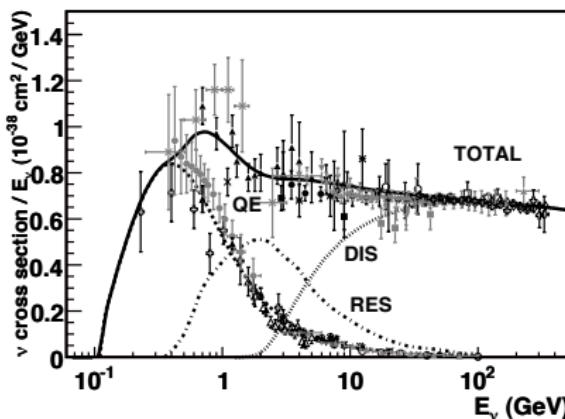
- Inputs
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4 Summary and Outlook

WHY SINGLE PION PRODUCTION?

- ❑ Two types of processes:
Charged-Current (CC) & Neutral-Current (NC) induced.
- ❑ Important contribution to the inclusive neutrino-nuclei (νA) cross section
 - ☒ RES: Predominantly $\Delta(1232)$ excitation $\implies \Delta \rightarrow \pi N$ (99.4%)
 - ☒ Prediction by NUANCE generator

[Formaggio, Zeller, Rev. Mod. Phys. (2012)]

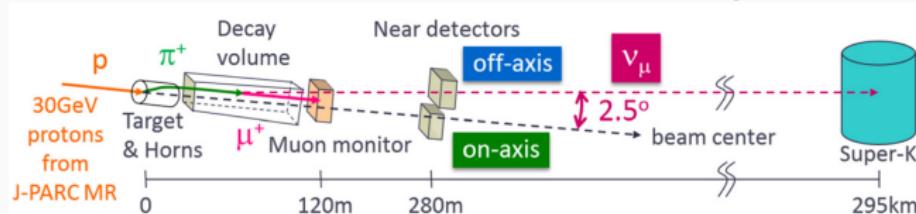


Total cross section per nucleon. Left: ν -CC; Right: $\bar{\nu}$ -CC.

WHY SINGLE PION PRODUCTION?

❑ Oscillation experiments (e.g. T2K)

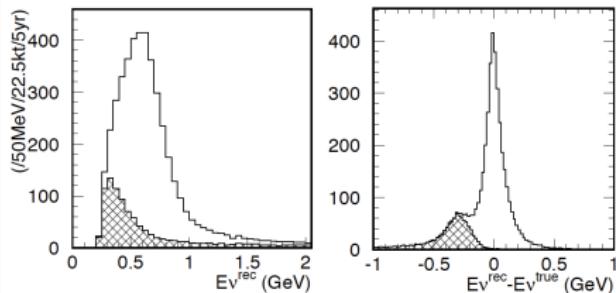
► survival probability of ν_μ : $P(\nu_\mu) = 1 - \sin^2 2\theta_{\mu\tau} \cdot \sin^2 \frac{\Delta m_{23} L}{E_\nu}$



❑ Source of experimental uncertainties

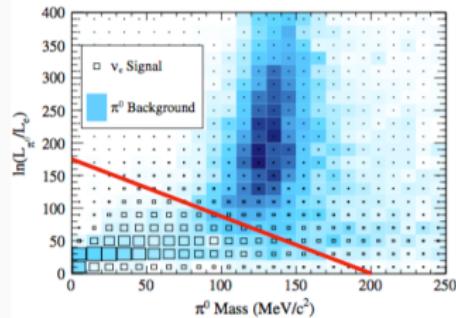
CC 1 π :

- CCQE-like events: misiden. of pion
- to be subtracted for a good E_ν



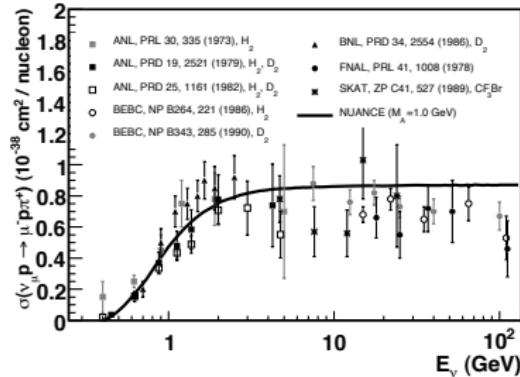
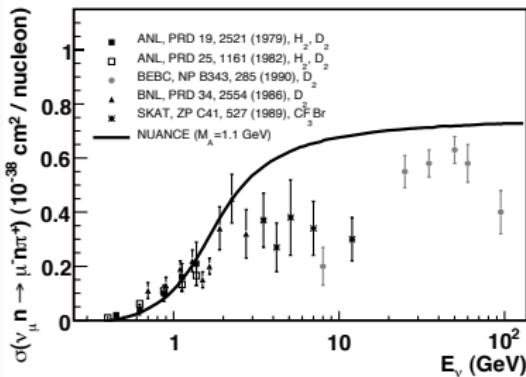
NC 1 π :

- e-like background to $\nu_\mu \rightarrow \nu_e$ searches
- improved at T2K with a π^0 rejection cut

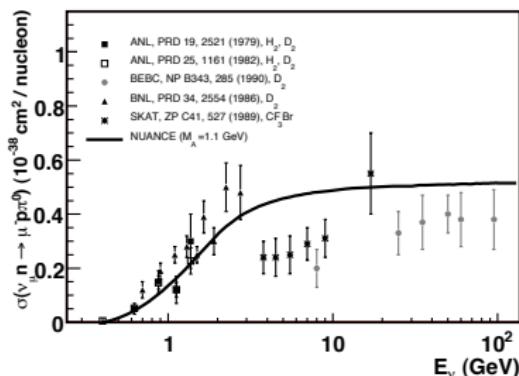


EXPERIMENTAL DATA

CC1 π



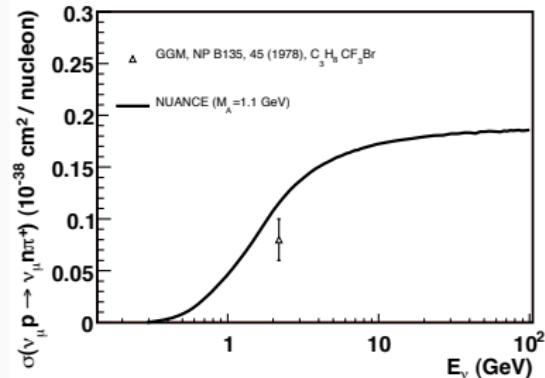
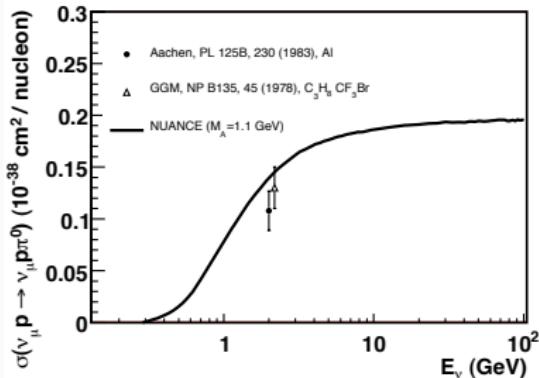
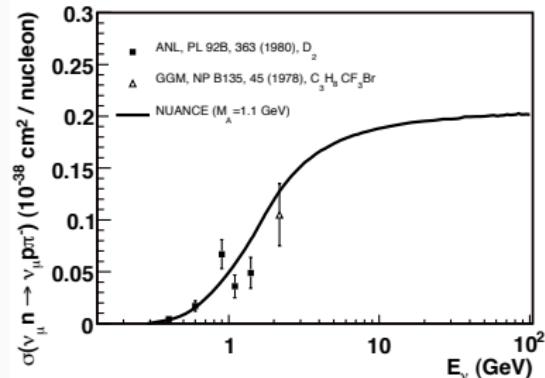
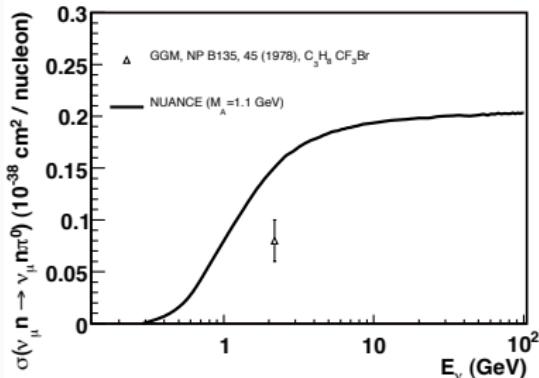
[Formaggio, Zeller, Rev. Mod. Phys. (2012)]



EXPERIMENTAL DATA

- NC1 π : Very rare data below 1 GeV

[Formaggio, Zeller, Rev. Mod. Phys. (2012)]



STATUS OF THEORETICAL STUDIES

❑ Isobar Models

☞ Δ and heavier resonances → nucleon-to-resonance form factors:

[e.g., Llewellyn Smith, Phys. Rep. 3 (1972)] [Fogli and Nardulli, Nucl. Phys. B160 (1979)] [Rein and Sehgal, Ann. Phys. (1981)]

- Real form factor from quark models
- Conserved vector current → related to electromagnetic ones extracted from electron scattering data
- PCAC → off-diagonal Goldberger-Treiman (GT) relation for the axial couplings

☞ Nonresonant mechanisms

[Fogli and Nardulli, Nucl. Phys. B160 (1979)]

[Bijtebier, Nucl. Phys. B21 (1970)]

[Alevizos et al., J. Phys. G 3(1977)]

❑ Hernandez-Nieves-Valverde (HNV) Model

☞ Δ resonances & non-resonant terms → constrained by chiral symmetry at threshold

[Hernandez, Nieves and Valverde, Phys. Rev. D (2007)]

☞ Final state interaction: imposing Watson's theorem [Alvarez-Ruso et al., Phys. Rev. D 93 (2016)]

☞ Unphysical spin-1/2 components: adding new contact terms

[Hernandez and Nieves, Phys. Rev. D (2017)]

STATUS OF THEORETICAL STUDIES

❑ Other Models:

☞ Dynamical model: coupled-channel Lippmann Schwinger equation

- Fulfilling Watson's theorem
- PCAC → partially constrain the axial current in terms of πN scattering amplitude fitted to data

[Nakamura, Kamano and Sato, Phys. Rev. D (2015)]

☞ Chiral effective model with π, N, Δ together with σ, ρ, ω

- Power counting only for tree diagrams

[Serot and Zhang, Phys. Rev. C (2012)]

☞ etc.

❑ Low energy regime:

Chiral symmetry + Power counting + Perturbative Unitarity

❑ Baryon Chiral Perturbation Theory (BChPT)

☞ Low-Energy theorems (axial only) at threshold using heavy baryon formalism

[Bernard, Kaiser and Meiñner, Phys. Lett. B (1994)]

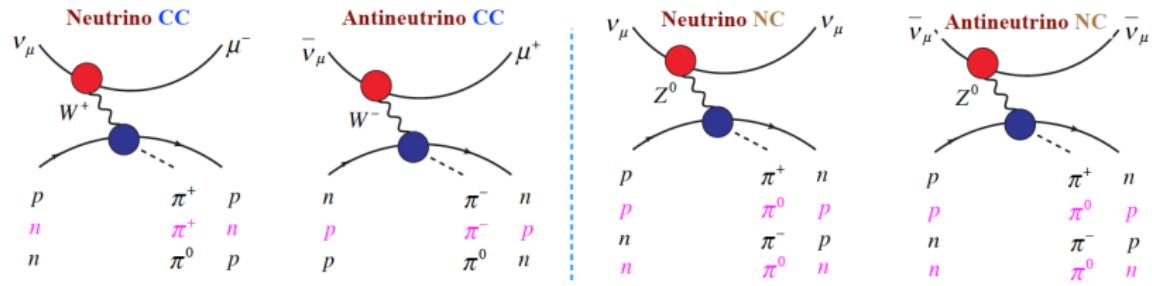
☞ Our work: One-loop analyses in relativistic BChPT with explicit Δs

[DLY, Alvarez-Ruso, Hiller-Blin and Vicent-Vacas, Phys. Rev. D (2018)]

[DLY, Alvarez-Ruso and Vicent-Vacas, Phys. Lett. B (2019)]

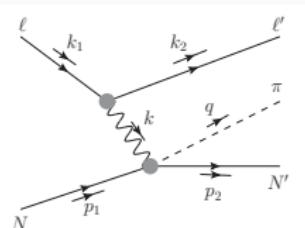
LEPTONIC AND HADRONIC PARTS

□ Physical channels (3 for CC & 4 for NC)



□ Amplitude structure:

- One-boson approximation and $k^2 \ll M_B^2$
- Leptonic part L_ν is well-known; Hadronic part H_μ needs to be investigated.



$$= i(2\pi)^4 \delta^{(4)}(k_1 + p_1 - k_2 - p_2 - q) \frac{iN^2}{M_B^2} \underbrace{\langle \ell' | J_\nu(0) | \ell \rangle}_{L^\mu} \underbrace{\langle \pi N' | J_\mu(0) | N \rangle}_{H_\mu}$$

CONVENIENT ISOSPIN DECOMPOSITION

- Isospin even (+), isospin odd (-), isoscalar (0)

$$\langle \pi^b N' | J_\mu^a(0) | N \rangle = \chi_f^\dagger [\delta^{ba} H_\mu^+ + i\epsilon^{bac} \tau^c H^- + \tau^b H_\mu^0] \chi_i$$

- The physical amplitudes constructed from the isospin amplitudes

$$H_\mu(\text{physical process}) = a_+ H_\mu^+ + a_- H_\mu^- + a_0 H_\mu^0$$

		Physical Process	a_+	a_-	a_0
NC	$Z^0 p \rightarrow p \pi^0$		1	0	1
	$Z^0 n \rightarrow n \pi^0$		1	0	-1
	$Z^0 n \rightarrow p \pi^-$		0	$-\sqrt{2}$	$\sqrt{2}$
	$Z^0 p \rightarrow n \pi^+$		0	$\sqrt{2}$	$\sqrt{2}$
CC	$W^+ p \rightarrow p \pi^+ / W^- n \rightarrow n \pi^-$		1	-1	0
	$W^+ n \rightarrow n \pi^+ / W^- p \rightarrow p \pi^-$		1	1	0
	$W^+ n \rightarrow p \pi^0 / W^- p \rightarrow n \pi^0$		0	$\sqrt{2}$	0

- The CC and NC amplitudes are related to each other

For CC, $H_\mu^\pm = \sqrt{2} \cos \theta_C (V_\mu^\pm - A_\mu^\pm)$, $H_\mu^0 = 0$.

For NC, $H_\mu^\pm = (1 - 2 \sin^2 \theta_W) V_\mu^\pm - A_\mu^\pm$, $H_\mu^0 = (-2 \sin^2 \theta_W) V_\mu^0$

THE LAGRANGIAN

- ❑ Covariant baryon chiral perturbation theory in SU(2) case.
- ❑ Nucleonic Lagrangian

[Fettes et al Ann. Phys. (2000)]

$$\mathcal{L}_N = \bar{N} [iD - m + \frac{g}{2}\psi\gamma_5] N + \bar{N} [c_j \mathcal{O}_j^{(2)} + d_k \mathcal{O}_k^{(3)}] N + \dots$$

- ❑ Purely mesonic Lagrangian [Gasser and Leutwyler, Ann. Phys. (1984)] [Gasser et al., Nucl. Phys. B307 (1988)]

$$\mathcal{L}_\pi = \frac{F^2}{4} \text{Tr}[\Delta_\mu U (\Delta^\mu U)^\dagger + \chi U^\dagger + U \chi^\dagger] + \sum_{j=3,4,6} \ell_j \mathcal{O}_j^{(4)}$$

- ❑ Electro-weak interactions enter through external fields [c.f. Scherer and Schindler, 2011, Springer]

☞ Charged weak bosons W^\pm :

$$r_\mu = 0, \quad l_\mu = -\frac{g}{\sqrt{2}}(V_{ud} W_\mu^+ \tau_+ + h.c.)$$

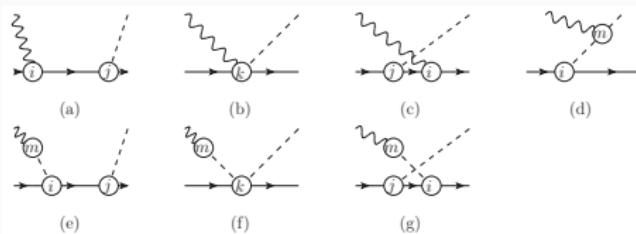
☞ Neutral weak boson Z^0 :

$$r_\mu = e \tan(\theta_W) Z_\mu^0 \frac{\tau_3}{2}, \quad l_\mu = -\frac{g}{\cos(\theta_W)} Z_\mu^0 \frac{\tau^3}{2} + e \tan(\theta_W) Z_\mu \frac{\tau_3}{2},$$

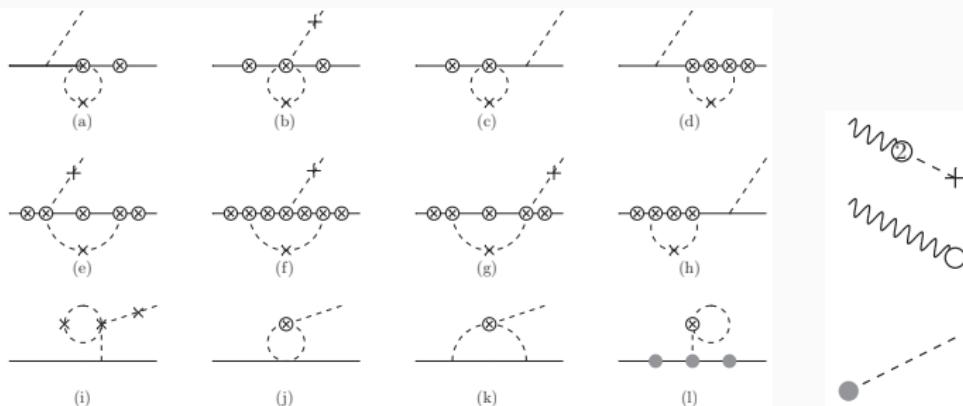
$$v_\mu^{(s)} = \frac{e \tan(\theta_W)}{2} Z_\mu^0$$

THE HADRONIC AMPLITUDE

- Tree diagrams up through $O(p^3)$:



- All possible loop diagrams at $O(p^3)$:



89 diagrams & wave function renormalization & EOMS

NECESSITY OF THE Δ RESONANCE

- ❑ Δ is strongly coupled to the final πN system
 - ☞ $\text{BR}(\Delta \rightarrow \pi N) \simeq 99.4\%$
 - ☞ Close to πN threshold: $\Delta = m_\Delta - m_N \sim 300 \text{ MeV}$

❑ Strategy: the δ -counting

[Pascalutsa and Phillips, Phys. Rev. C67 (2002)]

- ☞ hierarchy of scales: $M_\pi \sim p \ll \Delta \ll \Lambda \sim 4\pi F_\pi$
- ☞ expanding parameter: $\delta = \frac{\Delta}{\Lambda} \sim \frac{M_\pi}{\Delta} \sim \frac{p}{\Delta} \longrightarrow \frac{1}{p-m_\Delta} = \frac{1}{p-m_N-\Delta} \sim p^{-\frac{1}{2}}$

❑ Counting rule:

$$\text{chiral order } D = 4L + \sum_k kV^{(k)} - 2I_\pi - I_N - \frac{1}{2}I_\Delta$$

- ☞ only trees of $O(p^{3/2})$ and $O(p^{5/2})$
- ☞ No loop diagrams with explicit Δ up through $O(p^3)$

❑ The width effect

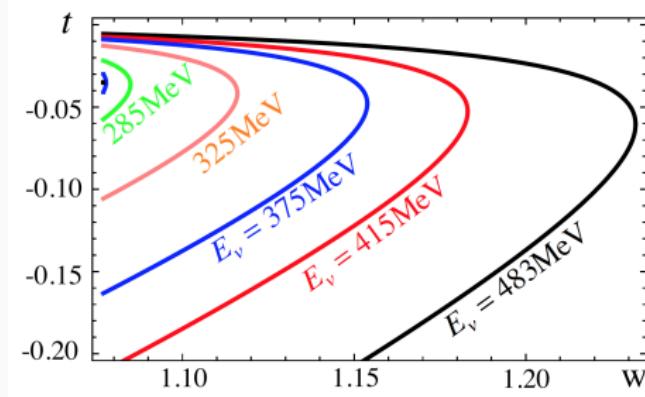
$$\frac{1}{m_\Delta^2 - s_\Delta} \rightarrow \frac{1}{m_\Delta^2 - im_\Delta \Gamma_\Delta(s_\Delta) - s_\Delta}$$

Energy dependent width $\Gamma_\Delta(s_\Delta)$ calculated in the same scheme

[Gegelia et al, Phys. Lett. B(2016)]

NUMERICAL SETTINGS

- ❑ Energies considered for $E_\nu \in [E_{\nu,th}, E_{\nu,max} \equiv E_{\nu,th} + M_\pi]$
 - ↳ E.g., $E_{\nu,max} = 415$ MeV for CC; $E_{\nu,max} = 289$ MeV for NC
 - ↳ Well below the Δ peak $\rightarrow \delta$ -counting is valid



W : CM energy of the final πN system (CC for example)

- ❑ Data for neutrino-induced single pion production off nucleons are very rare
- ❑ Values of the leading order constants

F_π	M_π	m_N	m_Δ		g_A	h_A
92.21	138.04	938.9	1232	MeV	1.27	1.43 ± 0.02

LOW ENERGY CONSTANTS BEYOND LO

- ❑ Most of the LECs (16 out of 23) are previously determined from other processes or observables

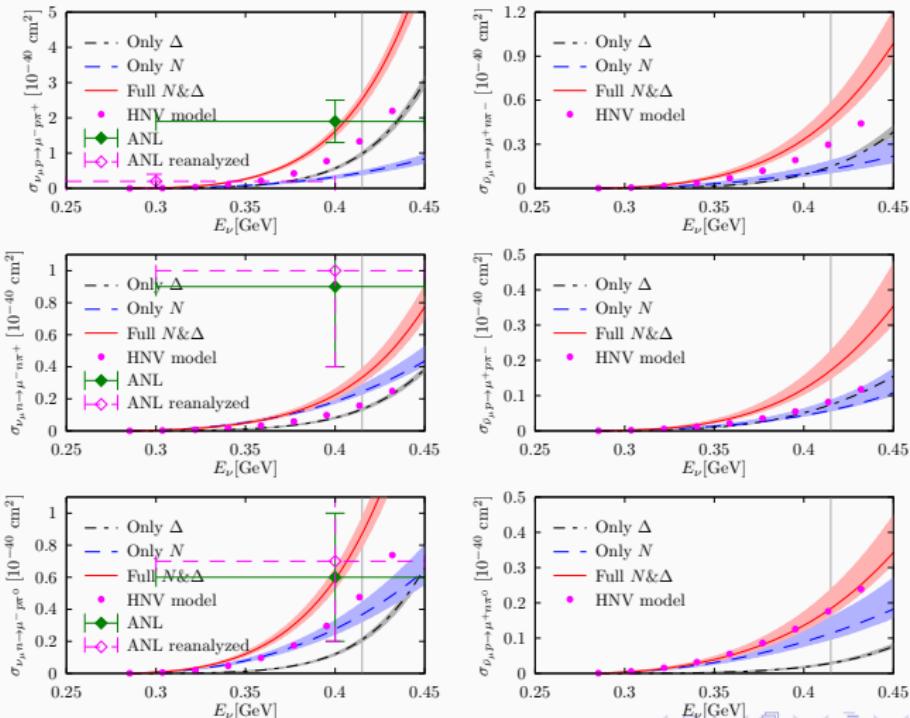
LEC	Value	Source
$\mathcal{L}_{\pi\pi}^{(4)}$		
$\bar{\ell}_6$	16.5 ± 1.1	$\langle r^2 \rangle_\pi$ [Gasser, Leutwyler 1984]
\tilde{c}_1	-1.00 ± 0.04	
\tilde{c}_2	1.01 ± 0.04	
$\mathcal{L}_{\pi N}^{(2)}$		πN scattering [Alarcon et al. 2013 & Chen et al. 2013]
\tilde{c}_3	-3.04 ± 0.02	
\tilde{c}_4	2.02 ± 0.01	
\tilde{c}_6	1.35 ± 0.04	
\tilde{c}_7	-2.68 ± 0.08	μ_p and μ_n [Bauer et al. 2012 & PDG2016]
$\mathcal{L}_{\pi N}^{(3)}$		
d_{1+2}^r	0.15 ± 0.20	
d_3^r	-0.23 ± 0.27	
d_5^r	0.47 ± 0.07	πN scattering [Alarcon et al. 2013 & Chen et al. 2013]
d_{14-15}^r	-0.50 ± 0.50	
d_{18}^r	-0.20 ± 0.80	
d_6^r	-0.70	
d_7^r	-0.47	$\langle r_E^2 \rangle_N$ [Fuchs et al. 2014]
d_{22}^r	0.96 ± 0.03	$\langle r_A^2 \rangle_N$ [Yao et al. 2017]
$\mathcal{L}_{\pi N \Delta}^{(2)}$	$(4.98 \pm 0.27)/m_N$	$\Gamma_\Delta^{\text{em}}$ [Bernard et al 2012]

- ❑ The remaining unknown LECs → set to natural size

$$d_j^r = 0.0 \pm 1.0 \text{ GeV}^{-2}, \quad j \in \{1, 8, 9, 14, 20, 21, 23\}$$

CROSS SECTIONS FOR CC1 π

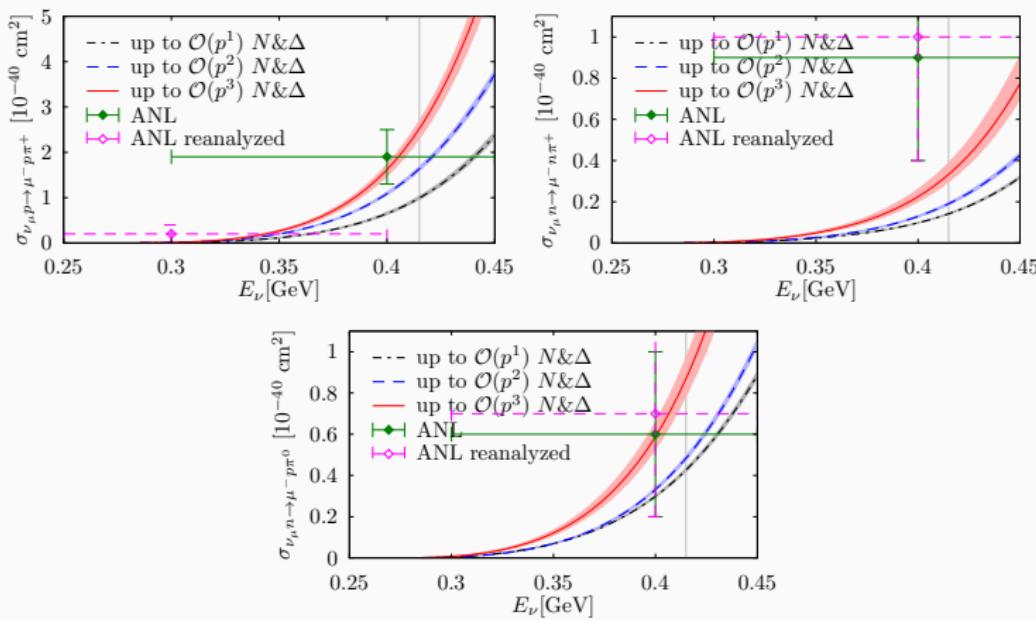
- Fairly good agreement with the ANL data for most of the channels except for $\nu_\mu n \rightarrow \mu^- n\pi^+$



CROSS SECTIONS FOR CC1 π

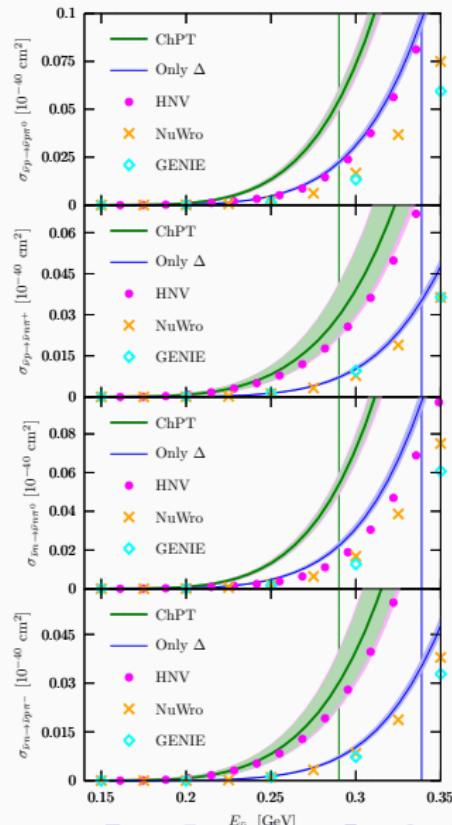
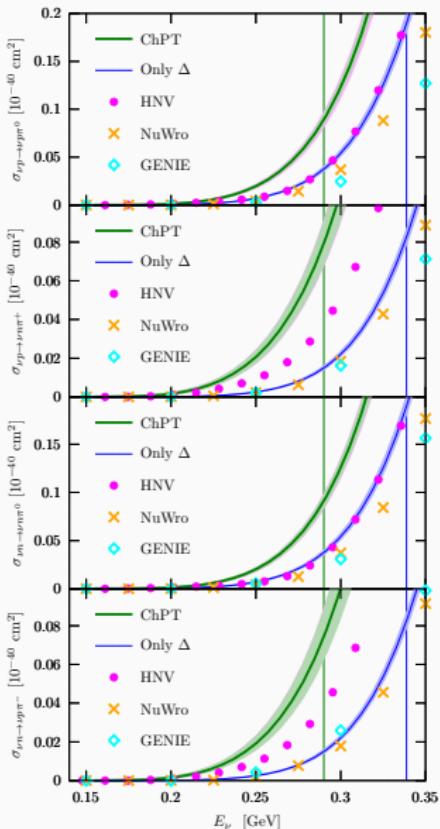
❑ Order by order

- ↳ Quite significant contribution when stepping from $O(p^2)$ and $O(p^3)$
- ↳ Next-order effects could still be relevant (especially loops that πN can be put on-shell)



CROSS SECTIONS FOR NC1 π

- ❑ The $O(p^3)$ ChPT calculation produces considerably larger cross sections with respect to the HNV model in all reaction channels
- ❑ Nuwro and GIENE results agree with the ChPT ones with Δ contribution
- ❑ Non-resonant contribution is sizeable, not accounted by Nuwro and GIENE



SUMMARY AND OUTLOOK

❑ Systematically study the weak single pion production off the nucleon for the first time within covariant BChPT up to $O(p^3)$ with explicit Δ s

- ☞ The Δ -mechanism contributes significantly to all production channels
- ☞ NC1 π : Non-resonant contribution is sizeable which is not implemented in events generators like NuWro and GIENE
- ☞ Provide a well-founded low energy benchmark for phenomenological models aimed at the description of weak pion production in the broad kinematic range of interest for current and future neutrino-oscillation experiments

❑ Future application and improvement

- ☞ Applied to study various low-energy theorems
- ☞ Applied to study electro-pion production for which there exists a wealth of experimental data and more LECs can be determined
- ☞ etc

