valence quark distribution inside pion from lattice QCD

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the enigma of pion mass

- contribution of gluon trace-anomaly to pion mass: Ş
 - $m_q \to 0$: $\langle \pi(p) | \Theta^{\mu}_{\mu} | \pi(p) \rangle$
- but must be large contributions of gluons to proton mass even in the chiral limit

 - simultaneously explaining the *absence* of pions mass

$$\rangle\rangle = -p^{\mu}p_{\mu} = m_{\pi}^2 \to 0$$

a puzzling dichotomy !!

understanding of proton mass will remain incomplete without



large-x behavior of pion valance PDF







whitepaper: arXiv:1907.08218

Pion and Kaon Structure at the Electron-Ion Collider

abstract

1 GeV mass-scale that characterizes atomic nuclei appear; why does it have the observed value; and, enigmatically, why are the composite Nambu-Goldstone (NG) bosons in quantum chromodynamics (QCD) abnormally light in comparison? In this perspective, we provide an analysis of the mass budget of the pion and proton in QCD; discuss the special role of the kaon, which lies near the boundary between dominance of strong and Higgs mass-generation mechanisms; and explain the need for a coherent effort in QCD phenomenology and continuum calculations, in exa-scale computing as provided by lattice QCD, and in experiments to make progress in understanding the origins of hadron masses and the distribution of that mass within them. We compare the unique capabilities foreseen at the electron-ion collider (EIC) with those at the hadron-electron ring accelerator (HERA), the



EIC: Sullivan process with off-shell pion

... and Parton structure of pion from QCD

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journey from Euclidean time to light-cone ... LaMET





rest frame of operator

rest frame of hadron

bare matrix element, extract ground-state aptitude

 $Q_b(z, p_z) = \langle H(p_z, 0) | \overline{\psi}(0, \tau) \gamma_0 W(0, z) \psi(z, \tau) | H(p_z, T_{\text{sink}}) \rangle$

equal-time, non-local operator within boosted hadron

W(0,z): Wilson line from 0 to z

X. Ji: Phys. Rev. Lett. 110, 262002 (2013)



perturbative matching target-mass & higher-twist corrections









rest frame of operator

rest frame of hadron

renormalized operator:

renormalization condition: $Z(z, p^R, a)Q_b(z, p = p^R, a) \equiv e^{izp_z^R}$

by computing the same matrix element for Landau-gauge fixed off-shell quarks with $p^2 > 0$



perturbative matching target-mass & higher-twist corrections

 $Q(z, p_z, p^R) = Z(z, p^R)Q_b(z, p_z)$

non-perturbative RI-MOM







rest frame of operator

rest frame of hadron

RI-MOM renormalized quasi-PDF (qPE)

Fourier transform w.r.t z

fixed p_z , small x — Iarge |z| — Contaminations of $\Lambda_{\rm QCD}, m_H$



perturbative matching target-mass & higher-twist corrections

DF):
$$q(x, p_z, p^R) = \frac{1}{4\pi} \int_0^\infty dz e^{-izp_z x} Q(z, p_z, p^R)$$







rest frame of operator

rest frame of hadron

 \Re RI-MOM qPDF to \overline{MS} -PDF:

$$q(x, P_z, P^R) = \int \frac{dy}{|y|} f(x, \mu) C\left(\frac{x}{y}, \frac{\mu}{yP_z}, \frac{P_\perp^R}{P_z^R}, \frac{yP_z}{P_z^R}\right) + \mathcal{O}\left(\frac{m_h^2}{P_z^2}, \frac{\Lambda_{QCD}^2}{P_z^2}\right)$$

first continuum limit, then $p_z \rightarrow 0$

universal for all hadrons, presently only 1-loop



perturbative matching target-mass & higher-twist corrections

Stewart, Zhao: Phys. Rev. D 97, 054512 (2018)



from QCD to $f_v^{\pi}(x)$...

energy levels of boosted pion



- Wilson-Clover valance quarks: $m_{\pi}^{val} = 300 \text{ MeV}$ $p_z = 1.29, 1.72 \text{ GeV}$ Final alpha spacing: a = 0.06 fm Figure 4.4 Solution and the size: $3.84 \times 2.88^3 \text{ fm}^4$ (64×48^3) $m_{\pi}^{val}L = 4.4$
- 3 2+1 flavor HISQ HotQCD gauge configurations (1-HYP smeared): $m_{\pi}^{sea} \simeq 160$ MeV, $m_{K}^{sea} \simeq 500$ MeV



pion almost on the light-cone





renormalized 3-pt operator in position space



red: real part; blue: imaginary part

dashed black lines: reconstructions from JAM global fits by reversing the procedure Barry et. al. (JAM collaboration): Phys. Rev. Lett. 121, 152001 (2018)



BNL-SBU LQCD: $2\langle x \rangle = 0.43(6)$

BNL-SBU: Phys. Rev. D (in press); arXiv:1905.06349

JAM:

 $2\langle x \rangle = 0.437$

4 3 $f_v^{\pi}(x,\mu) \\ c$ 1

pion valance PDF



$$f_v^{\pi}(x) = f_u^{\pi}(x) - f_d^{\pi}(x) \dots 0 < x < 1$$



more is better ... alternative representation

- pseudo loffe-time distribution: Lorentz invariant
- approach to light-cone: $z^2 \rightarrow 0$, $\nu = \text{fixed}$

reduced loffe-time distribution:

$$rITD(z^{2},\nu) = \frac{M_{b}(z^{2},\nu)}{M_{b}(z^{2},0)} = \sum_{n=0}^{\infty} \frac{(-i\nu)^{n}}{n!}$$

renormalized

perturbative, 1-loop

Izubuchi et. al.: Phys. Rev. D98, 056004 (2018)

$M_b(z^2,\nu) = \langle H(p_z,0) | \overline{\psi}(0,\tau)\gamma_0 W(0,z)\psi(z,\tau) | H(p_z,T_{\rm sink}) \rangle$ $\nu \equiv z p_z$

$$\lim_{z^2 \to 0} M(z^2, \nu) \to M(x, \mu) = \int_{-1}^{1} C(\mu, z) f(x, \mu) e^{-\frac{1}{2}}$$



target-mass, higher-twist







reduced loffe-time distribution



Xiang Gao et. al.: BNL-SBU-Tsinghua, on-going

a = 0.04 fm, $p_z = 0 - 1.48$ GeV, z = a = 0.06 fm, $p_z = 0 - 1.29$ GeV, z = 0 - 1.29 GeV, z = 0

$$= 0 - 0.8 \text{ fm}, \ m_{\pi}^{val} = 300 \text{ MeV}, \ 64 \times 64^3$$
$$= 0 - 0.8 \text{ fm}, \ m_{\pi}^{val} = 300 \text{ MeV}, \ 64 \times 48^3$$





joint fit data for all z, p_z for $\nu = zP_z = 0 - 4$ sensitive only to $\langle x^2 \rangle$, $\langle x^4 \rangle$

$$rITD(zp_{z}, z^{2}) = \sum_{n=0}^{\infty} \frac{(-izP^{z})^{n}}{n!} \frac{C_{n}(\mu^{2}z^{2})}{C_{0}(\mu^{2}z^{2})} \langle x^{n} \rangle (\mu)$$

rITD reconstructed form JAM results using n=2-10 moments



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2nd moment of $f_v^{\pi}(x)$



BNL-SBU-Tsinghua, on-going

BNL-SBU-Tsinghua LQCD: $\langle x^2 \rangle \simeq 0.10(1)$

$$\mu = 3.2 \text{ GeV}$$

JAM: $\langle x^2 \rangle = 0.095$



4th moment of $f_v^{\pi}(x)$



BNL-SBU-Tsinghua, on-going

BNL-SBU-Tsinghua LQCD: $\langle x^4 \rangle \simeq 0.030(5)$

 $\mu = 3.2 \text{ GeV}$

JAM: $\langle x^4 \rangle = 0.032$



instead of a summary ... shape of $f_v^{\pi}(x)$



BNL-SBU-Tsinghua, on-going

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