GPD Measurements at COMPASS

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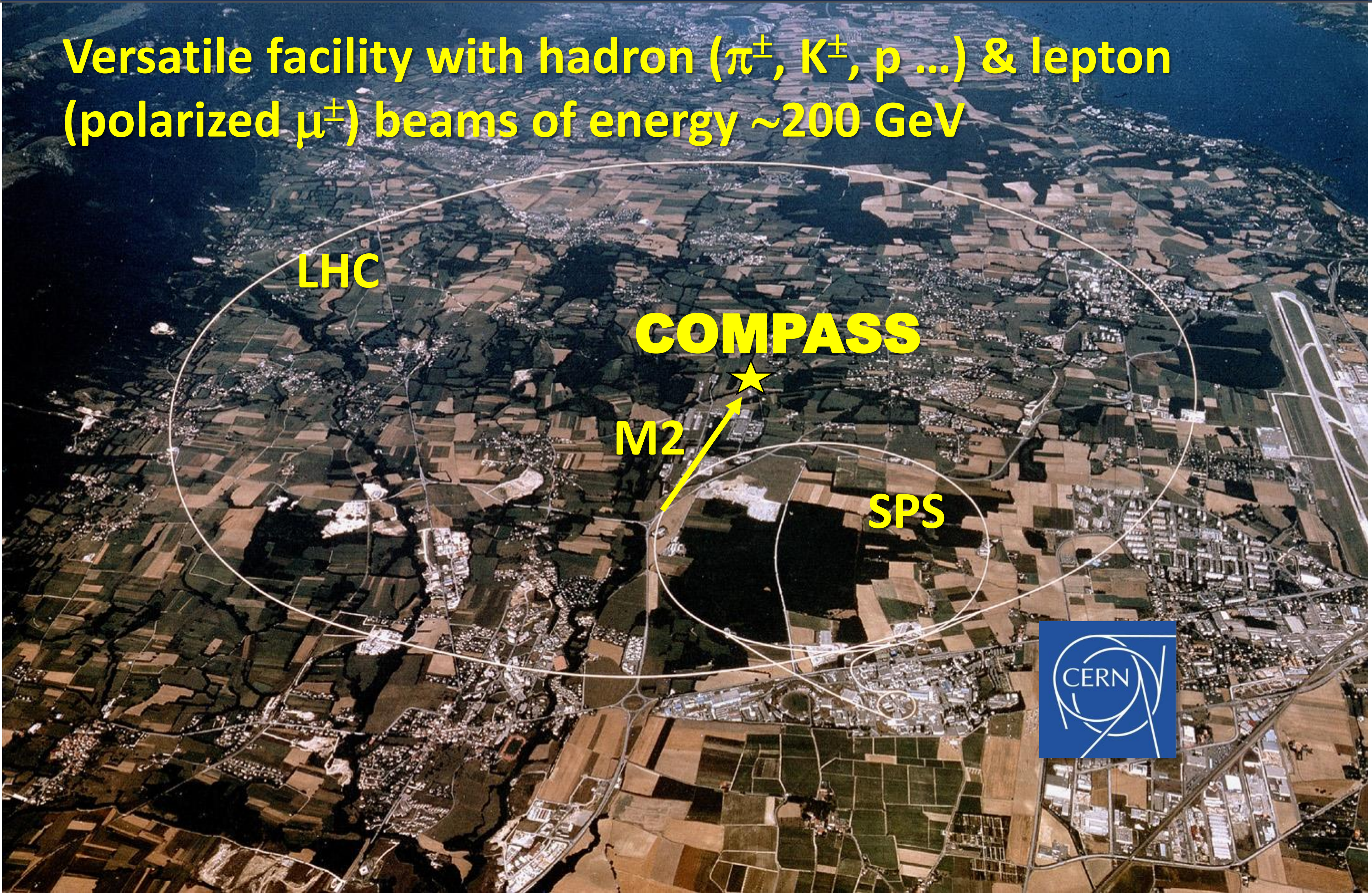
Outline

- The COMPASS Experiment
- Deeply Virtual Compton Scattering
- Hard Exclusive Meson Production
- Summary

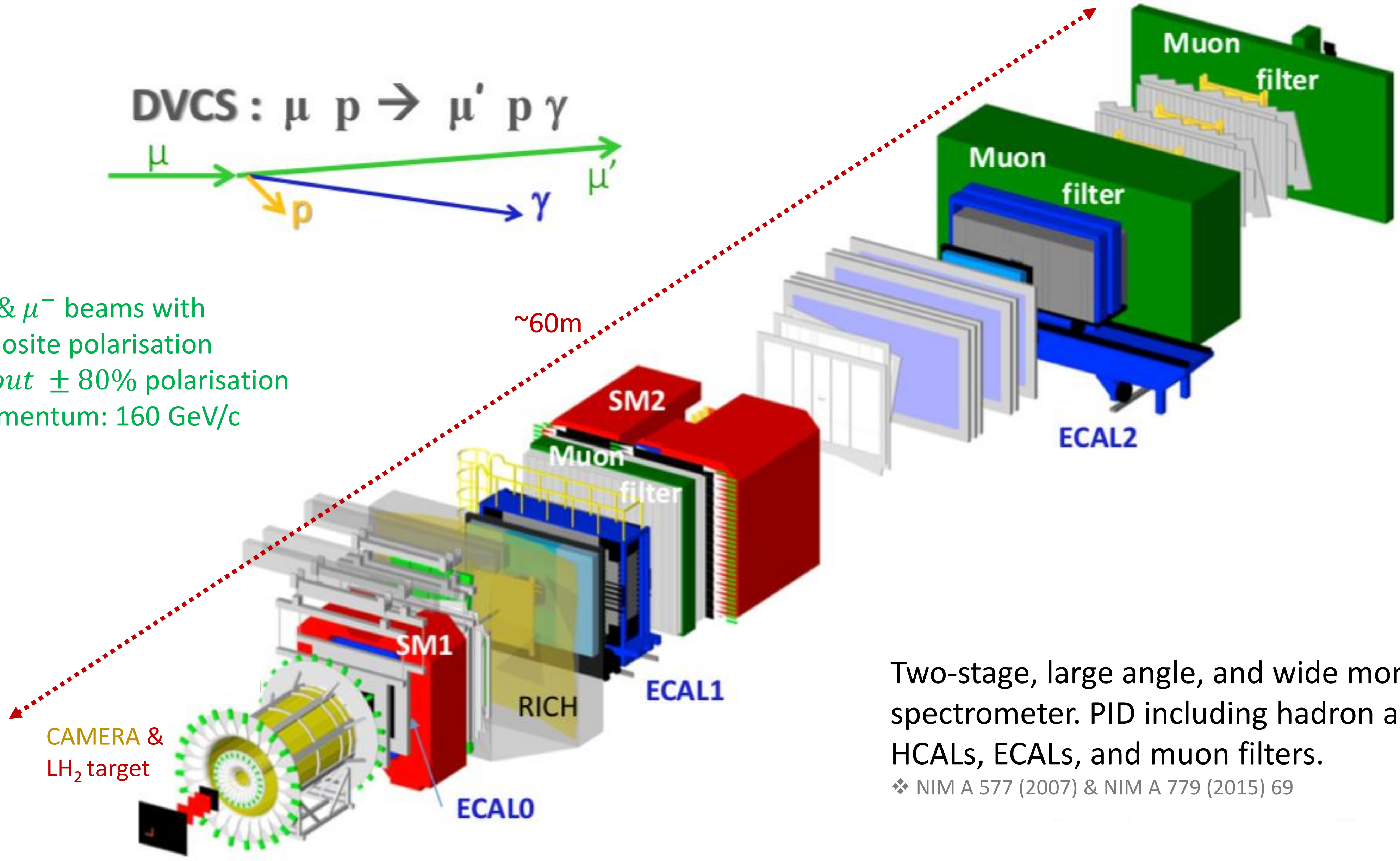
COMPASS Experiment

Versatile facility with hadron ($\pi^{\pm}$, $K^{\pm}$, p …) & lepton (polarized $\mu^{\pm}$) beams of energy ~200 GeV



COmmon
Muon and
Proton
Apparatus for
Structure and
Spectroscopy

COMPASS Setup for GPD Measurements



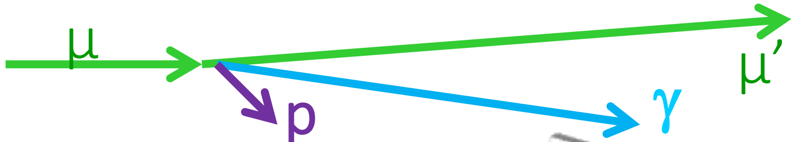
- μ^+ & μ^- beams with opposite polarisation
- $About\ \pm 80\%$ polarisation
- Momentum: 160 GeV/c

Two-stage, large angle, and wide momentum range spectrometer. PID including hadron absorbers, RICH, HCALs, ECALs, and muon filters.

❖ NIM A 577 (2007) & NIM A 779 (2015) 69

COMPASS Setup for GPD Measurement

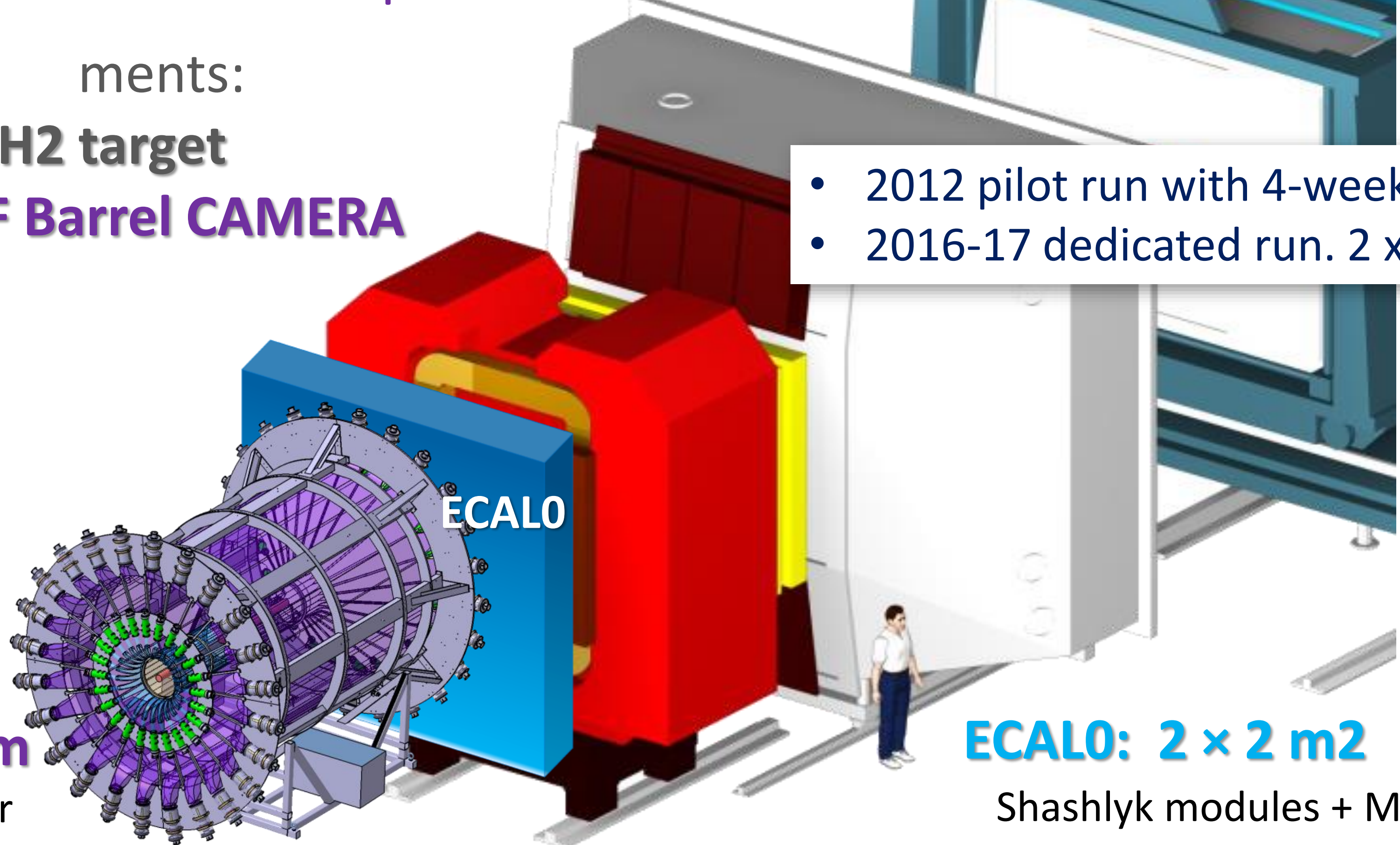
DVCS : $\mu \; p \rightarrow \mu' \; p \; \gamma$



ments:

- **2.5m LH2 target**
- **4m ToF Barrel CAMERA**
- **ECAL0**

- 2012 pilot run with 4-week data taking
- 2016-17 dedicated run. 2 x 6 months.



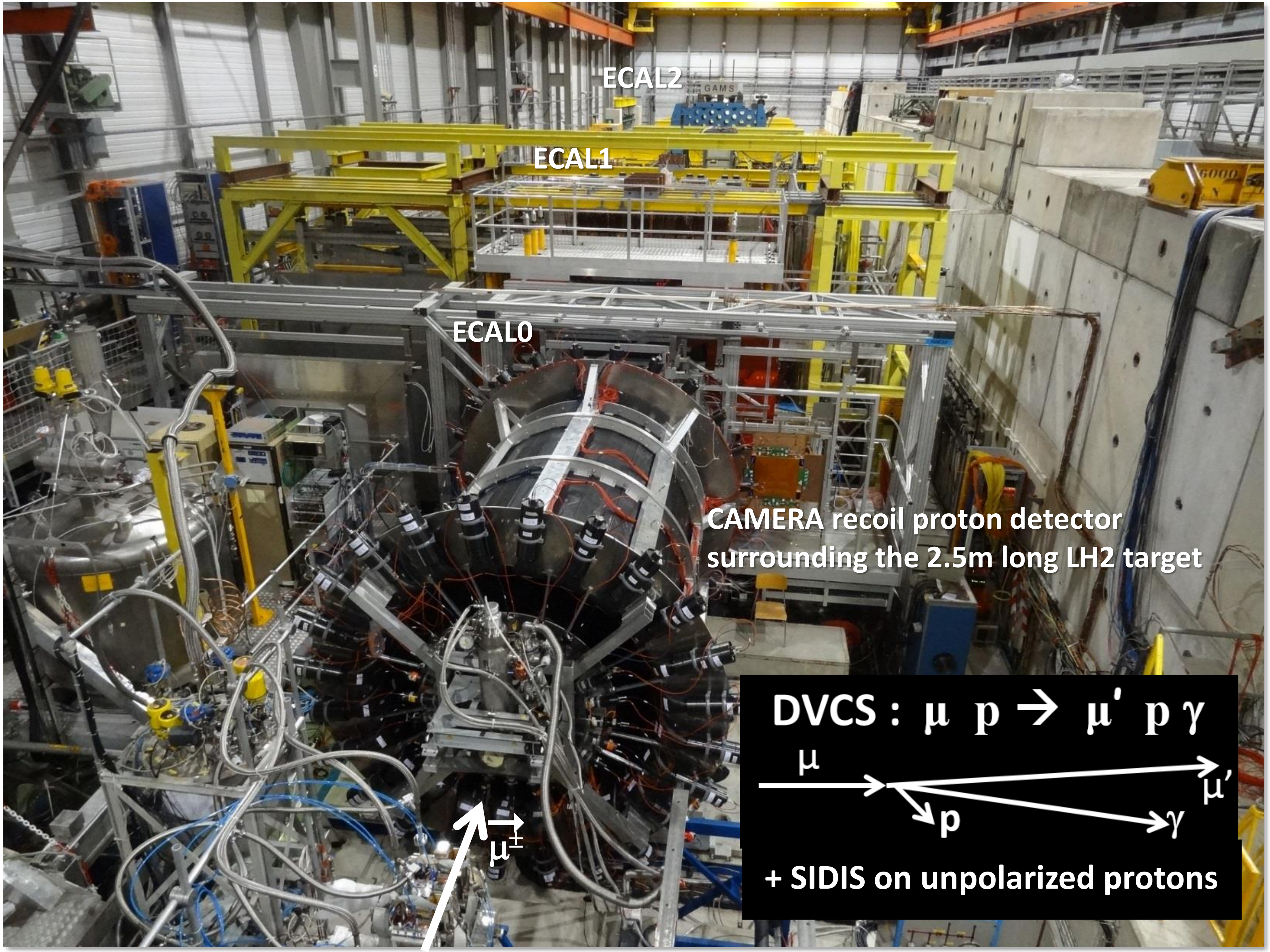
CAMERA

L=4m ∅=2m

24 inner & outer
scintillators separated by about 1m
1 GHz SADC readout, 330ps ToF resolution

ECAL0: 2 × 2 m2

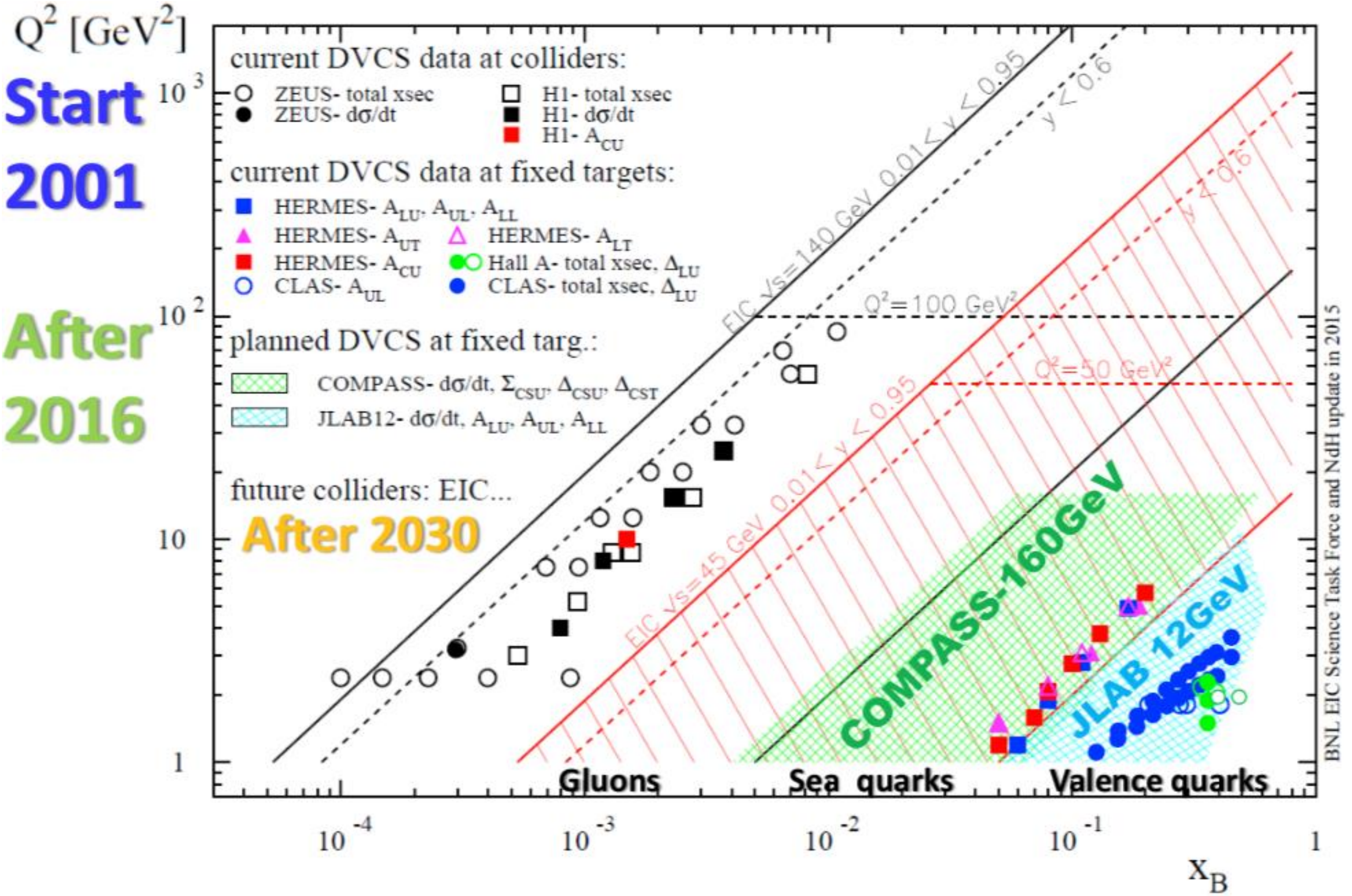
Shashlyk modules + MAPD readout
one module is made of 9 cells (4×4 cm^2)
= 194 modules or 1746 cells


ECAL2
GAMS
ECAL1
ECAL0
CAMERA recoil proton detector
surrounding the 2.5m long LH2 target
DVCS : μ p → μ' p γ
μ
μ'
p
γ
+ SIDIS on unpolarized protons
μ±

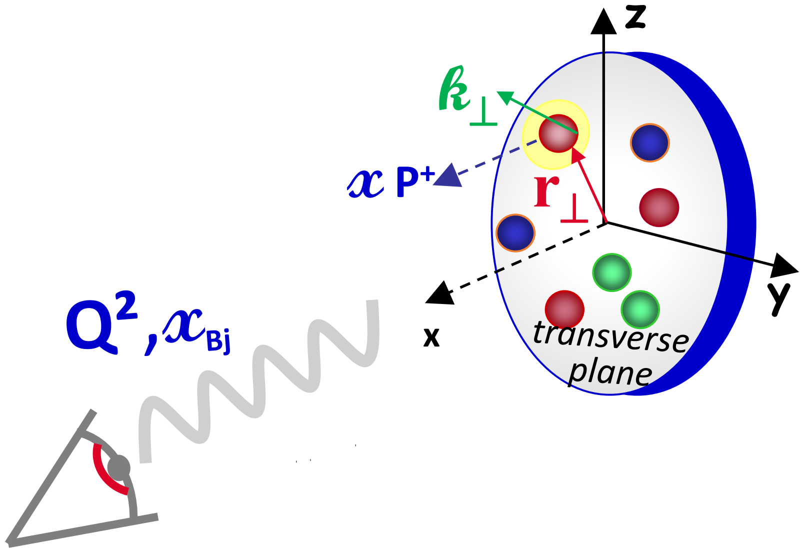
COMPASS Coverage

Start 2001

After 2016



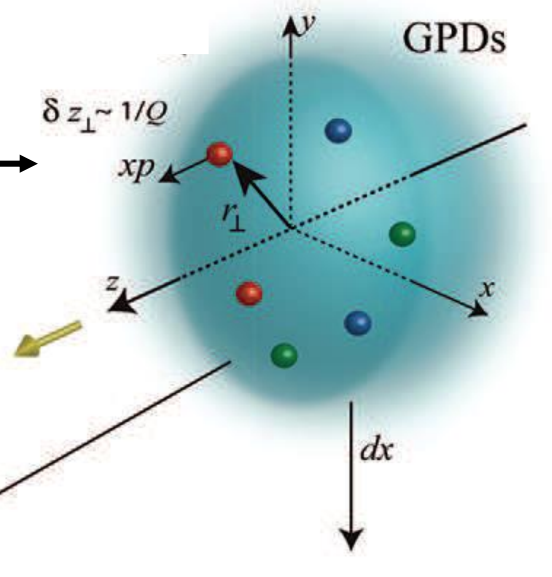
Proton Picture

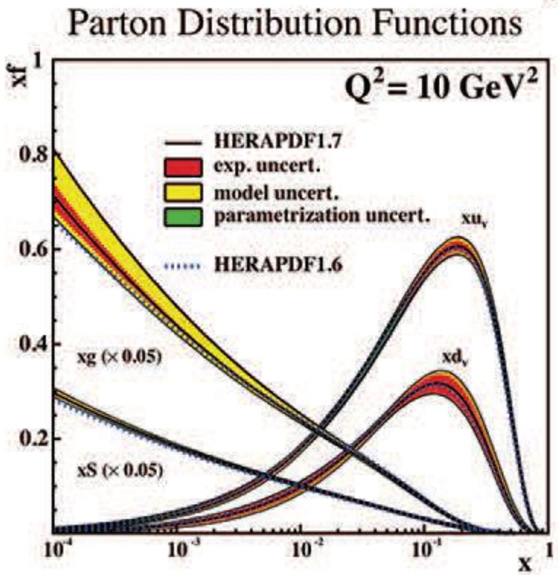


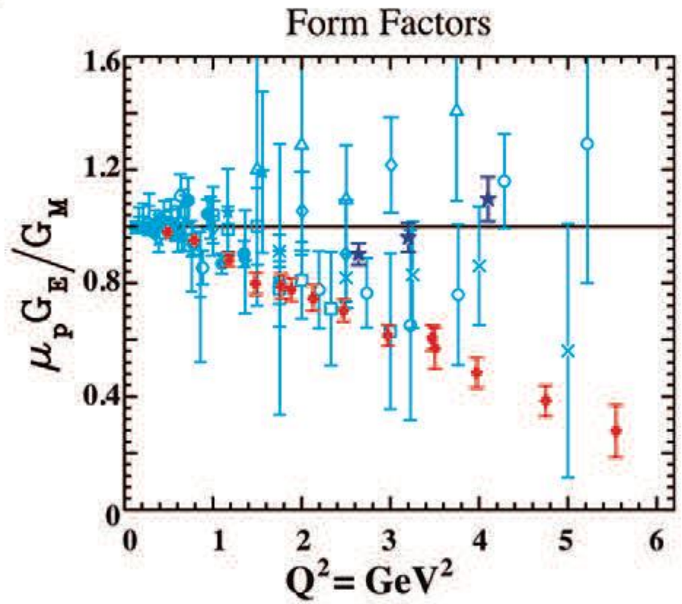
Wigner distributions

$$\mathcal{W}\,(x, k_\perp, r_\perp) \xrightarrow{\int dk_\perp}$$

- The complete partonic description of the nucleon



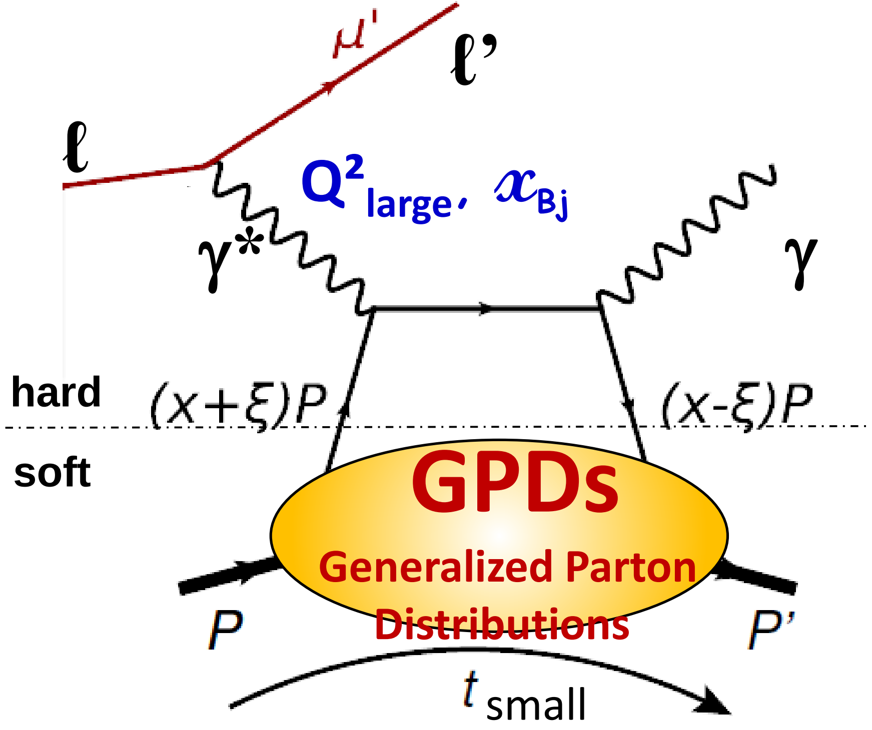




GPDs

- 8 in total
- Info. on transverse position
- Embody both PDFs and Form factors
- Accessible in exclusive processes

DVCS



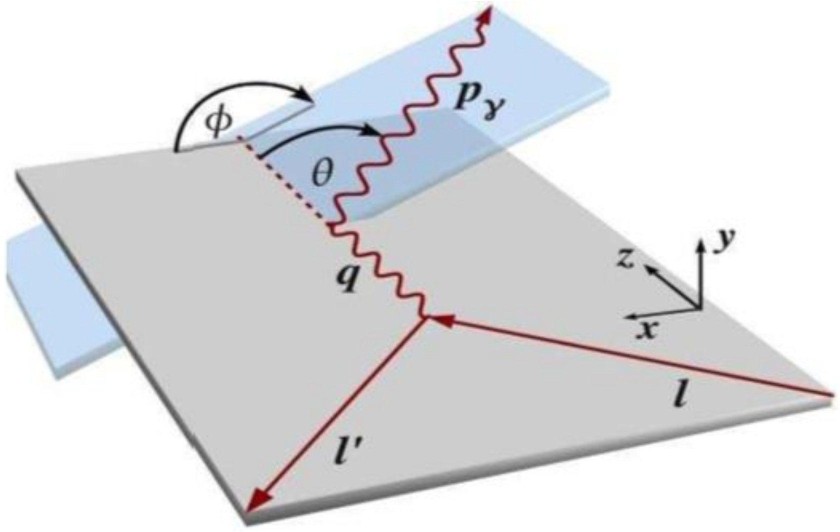
DVCS: $l + p \rightarrow l' + p' + \gamma$

DVCS: the golden channel to investigate the GPDs; Interference with the well-known BH process gives access to more info.

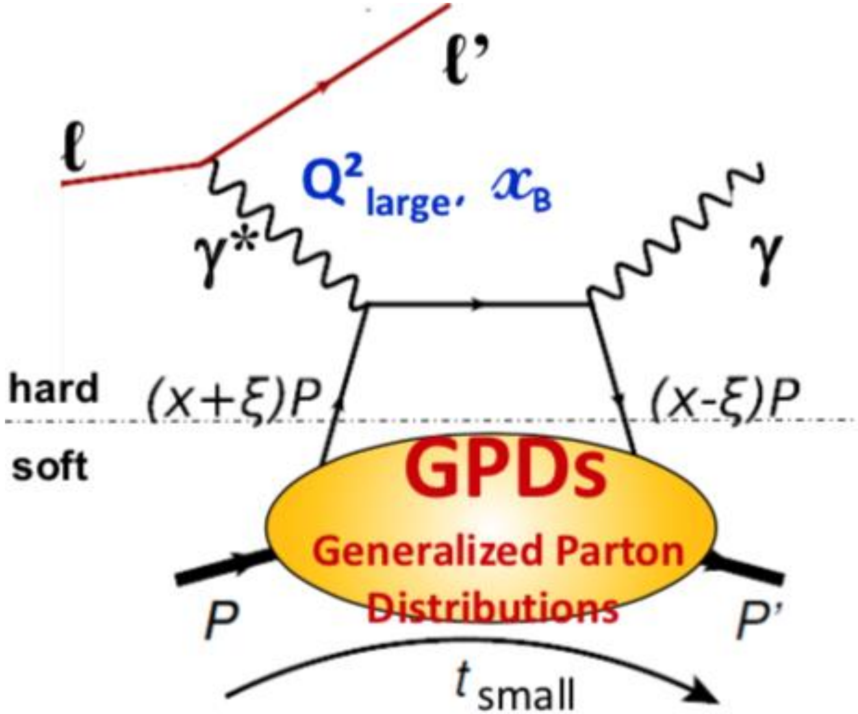
The variables measured in the experiment:

E_ℓ, Q^2, $x_{Bj} \sim 2\xi/(1+\xi)$,

t (or $\theta_{\gamma^*\gamma}$) and ϕ ($\ell\ell'$ plane/$\gamma\gamma^*$ plane)



DVCS



- The GPDs depends on the following variables:
  - x: average longitudinal momentum frac.
  - ξ: longitudinal momentum diff.
  - t: four momentum transfer
    (correlated to $b_\perp$ via Fourier transform)
  - Q^2: virtuality of γ^*

Sensible to 4 GPDs, at COMPASS with small x_B
→ focuses on **H**

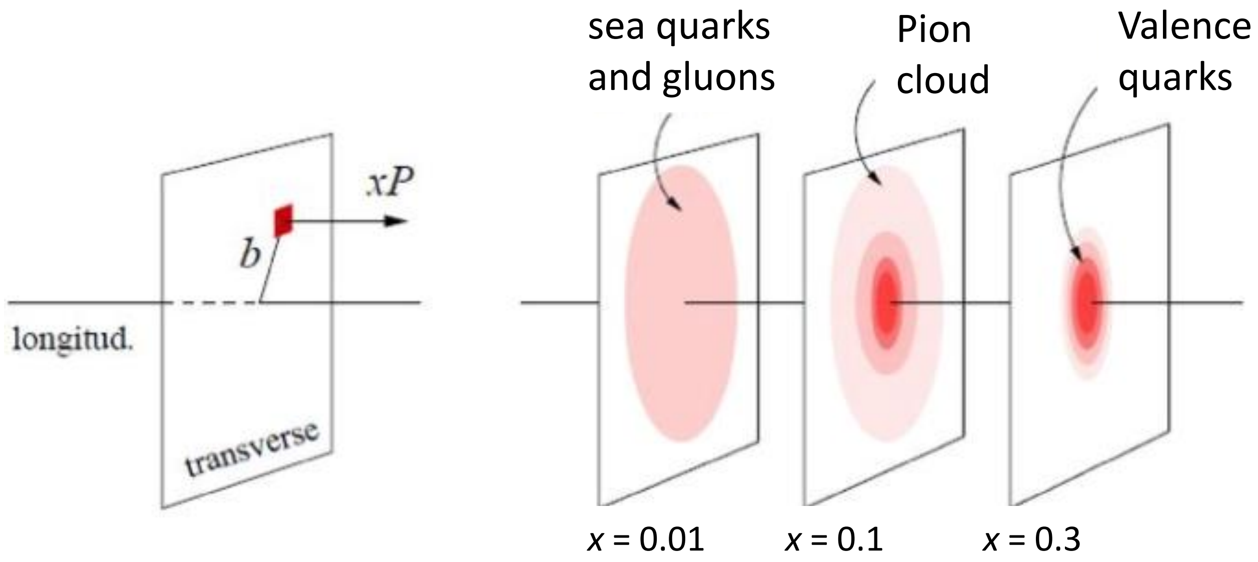
CFF GPD REAL part Imaginary part

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x,\xi,t)}{x-\xi+i\varepsilon} = \mathcal{P}\int_{-1}^{+1} dx \frac{\mathbf{H}(x,\xi,t)}{x-\xi} - i\pi\,\mathbf{H}(x=\pm\xi,\xi,t)$$

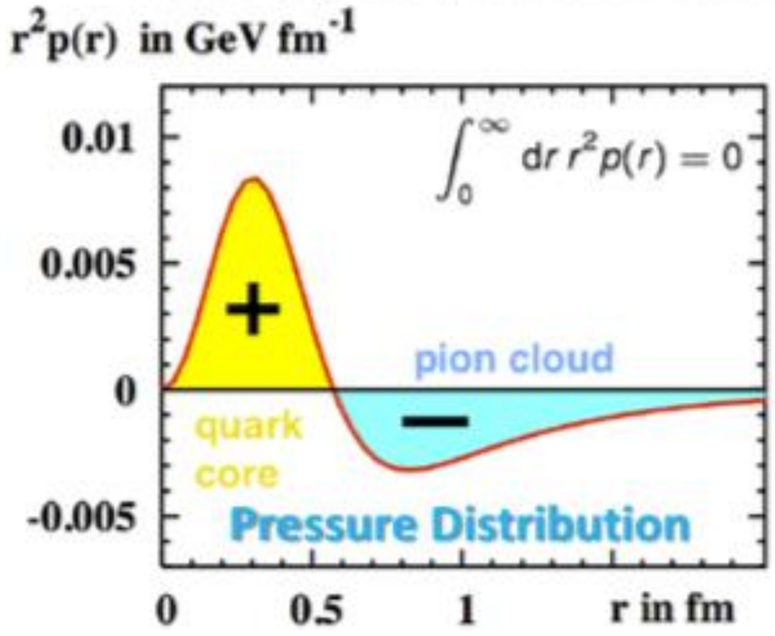
t, ξ fixed

$$\mathcal{R}e\,\mathcal{H}(\xi,t) = \mathcal{P}\int dx \frac{\mathcal{I}m\,\mathcal{H}(x,t)}{x-\xi} + d(t)$$

Transverse Imaging and Pressure Dist.



M. Polyakov, P. Schweitzer



CFF *GPD*

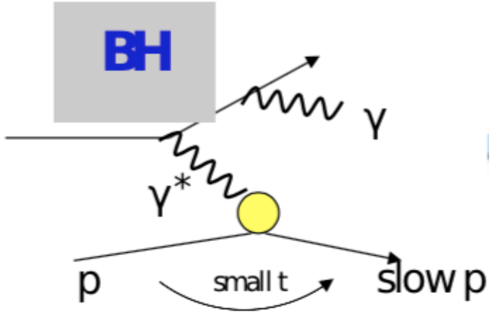
REAL part **Imaginary part**

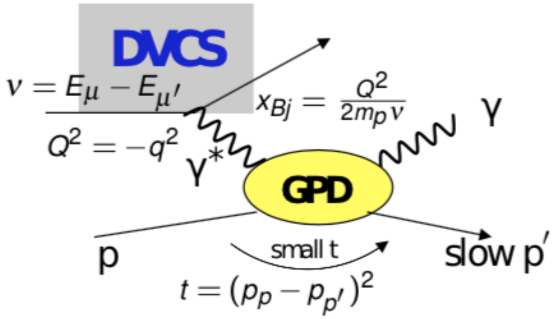
$$\mathcal{H} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x,\xi,t)}{x-\xi+i\varepsilon} = \mathcal{P}\int_{-1}^{+1} dx \frac{\mathbf{H}(x,\xi,t)}{x-\xi} - i\pi\, \mathbf{H}(x=\pm\xi,\xi,t)$$

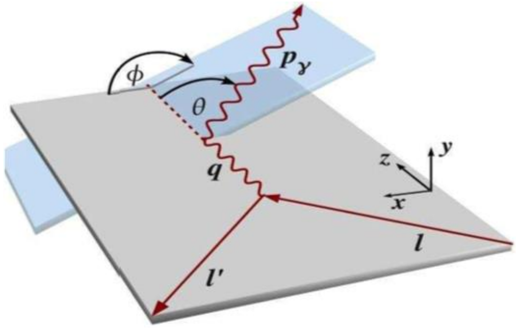
t, ξ fixed

$$\mathcal{Re}\,\mathcal{H}(\xi,t) = \mathcal{P}\int dx \frac{\mathcal{Im}\,\mathcal{H}(x,t)}{x-\xi} + \mathrm{d(t)}$$

Azimuthal Dependence







$$\frac{\mathrm{d}^4\sigma(\ell p \to \ell p \gamma)}{\mathrm{d}x_B \mathrm{d}Q^2 \mathrm{d}|t| \mathrm{d}\phi} = \underbrace{\mathrm{d}\sigma^{BH}}_{\text{Well known}} + \left(\mathrm{d}\sigma^{DVCS}_{unpol} + P_\ell\, \mathrm{d}\sigma^{DVCS}_{pol}\right) + (e_\ell \mathrm{Re}\, I + e_\ell P_\ell\, \mathrm{Im}\, I)$$

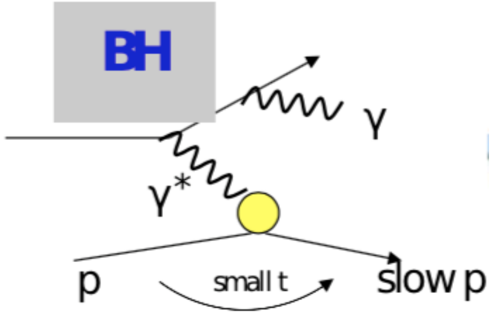
Beam Charge-spin difference & sum

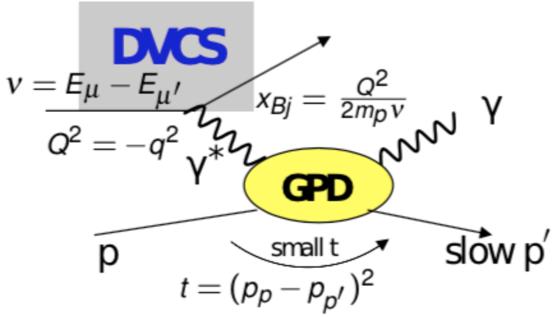
$$\boxed{\begin{aligned} \mathcal{D}_{CS,U}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \\ \mathcal{S}_{CS,U}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) \end{aligned}}$$

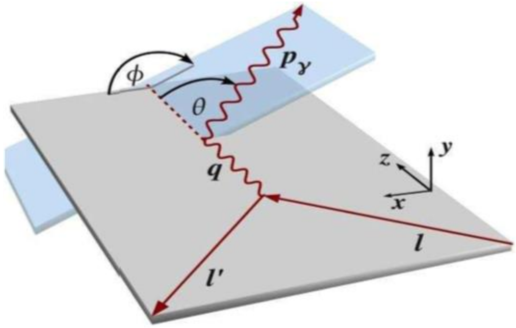
$$\begin{aligned}
\mathrm{d}\sigma^{BH} &\propto c_0^{BH} + c_1^{BH}\cos\phi + c_2^{BH}\cos 2\phi \\
\mathrm{d}\sigma^{DVCS}_{unpol} &\propto c_0^{DVCS} + c_1^{DVCS}\cos\phi + c_2^{DVCS}\cos 2\phi \\
\mathrm{d}\sigma^{DVCS}_{pol} &\propto s_1^{DVCS}\sin\phi \\
\mathrm{Re}\, I &\propto c_0^I + c_1^I\cos\phi + c_2^I\cos 2\phi + c_3^I\cos 3\phi \\
\mathrm{Im}\, I &\propto s_1^I\sin\phi + s_2^I\sin 2\phi
\end{aligned}$$

Twist-2 >> ■ Twist-3, ■ Twist-2 - double helicity flip for gluons (NLO)

Azimuthal Dependence







$$\frac{d^4\sigma(\ell p \to \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma^{DVCS}_{unpol} + P_\ell\, d\sigma^{DVCS}_{pol}\right) + (e_\ell \mathrm{Re}\, I + e_\ell P_\ell \mathrm{Im}\, I)$$

Beam Charge-spin difference & sum

$$\boxed{\begin{aligned} \mathcal{D}_{\text{CS, U}}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \\ \mathcal{S}_{\text{CS, U}}(\phi) &\equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) \end{aligned}}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos\phi + c_2^{BH} \cos 2\phi$$

$$d\sigma^{DVCS}_{unpol} \propto c_0^{DVCS} + c_1^{DVCS} \cos\phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma^{DVCS}_{pol} \propto s_1^{DVCS} \sin\phi$$

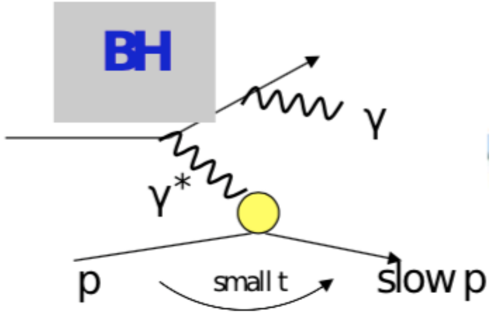
$$\mathrm{Re}\, I \propto c_0^I + c_1^I \cos\phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

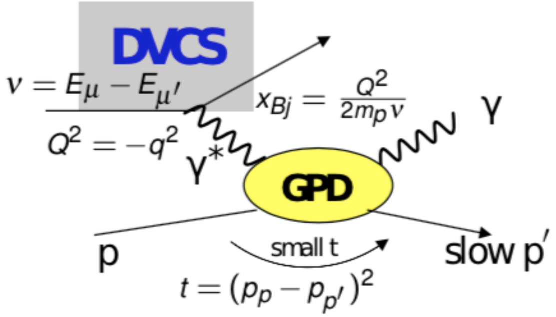
$\mathcal{D}_{\text{CS, U}}(\phi)$

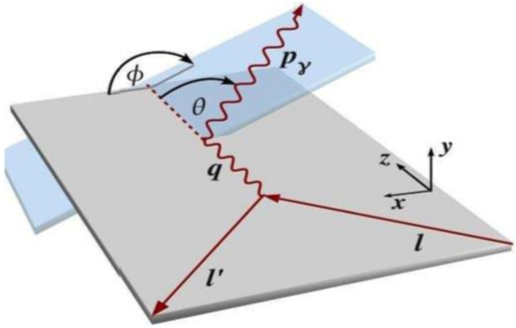
$$\mathrm{Im}\, I \propto s_1^I \sin\phi + s_2^I \sin 2\phi$$

Twist-2 >> ■ Twist-3, ■ Twist-2 - double helicity flip for gluons (NLO)

Azimuthal Dependence







$$\frac{d^4\sigma(\ell p \to \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma^{DVCS}_{unpol} + P_\ell\, d\sigma^{DVCS}_{pol}\right) + (e_\ell \mathrm{Re}\, I + e_\ell P_\ell \,\mathrm{Im}\, I)$$

Beam Charge-spin difference & sum

$$\mathcal{D}_{\text{CS, U}}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{\text{CS, U}}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH}\cos\phi + c_2^{BH}\cos 2\phi$$

$$d\sigma^{DVCS}_{unpol} \propto c_0^{DVCS} + c_1^{DVCS}\cos\phi + c_2^{DVCS}\cos 2\phi$$

$\mathcal{S}_{\text{CS, U}}(\phi)$

$$d\sigma^{DVCS}_{pol} \propto s_1^{DVCS}\sin\phi$$

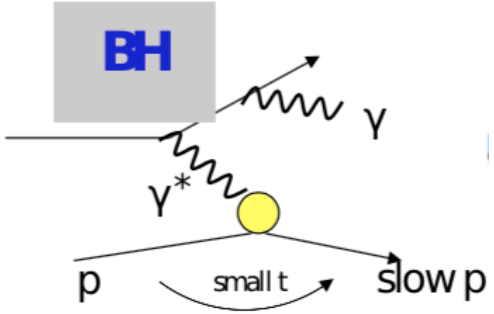
$$\mathrm{Re}\, I \propto c_0^I + c_1^I\cos\phi + c_2^I\cos 2\phi + c_3^I\cos 3\phi$$

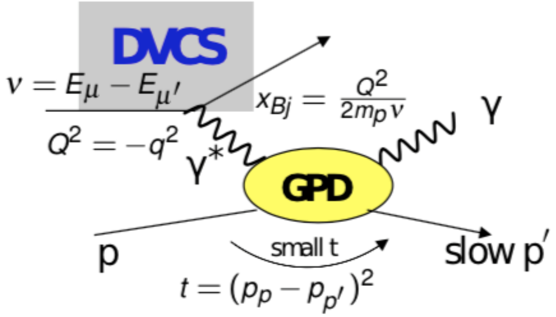
$$\mathrm{Im}\, I \propto s_1^I\sin\phi + s_2^I\sin 2\phi$$

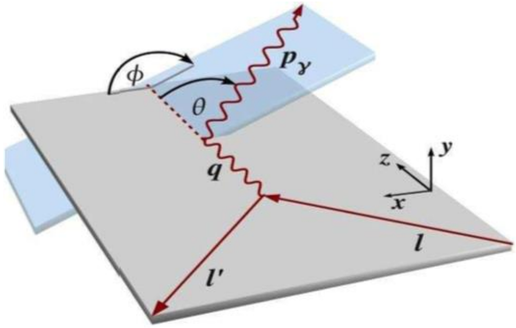
$\mathcal{S}_{\text{CS, U}}(\phi)$

Twist-2 >> Twist-3, Twist-2 - double helicity flip for gluons (NLO)

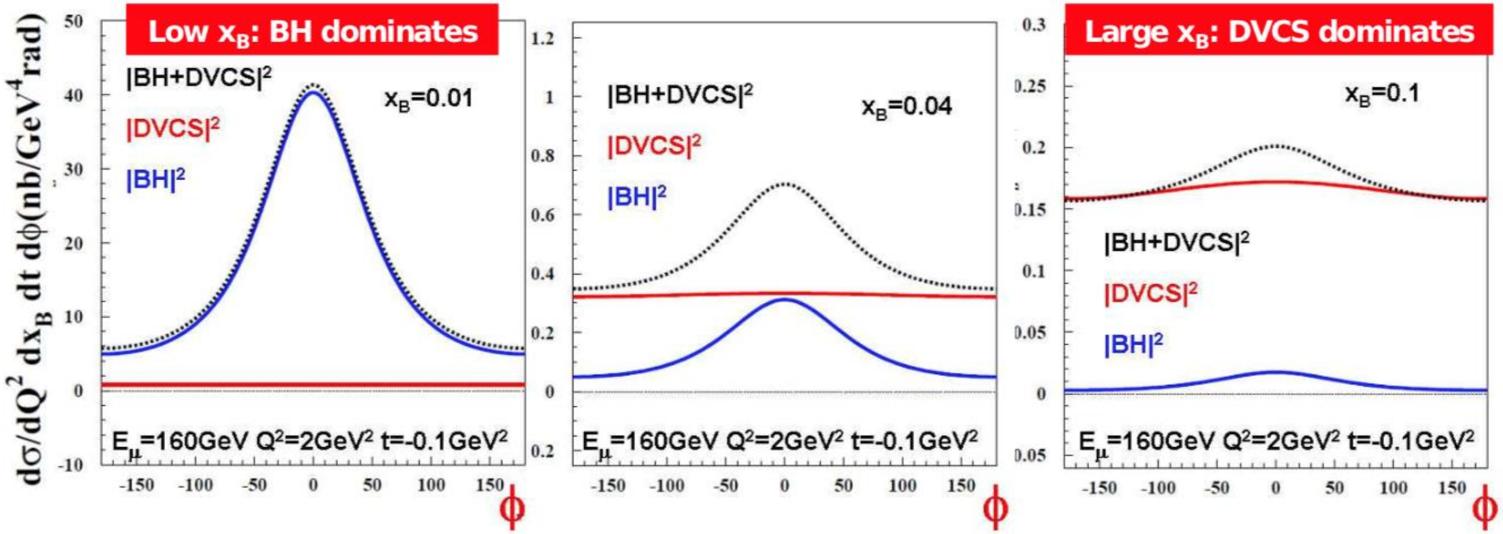
Azimuthal & x_{Bj} Dependence







Increasing x_{Bj}

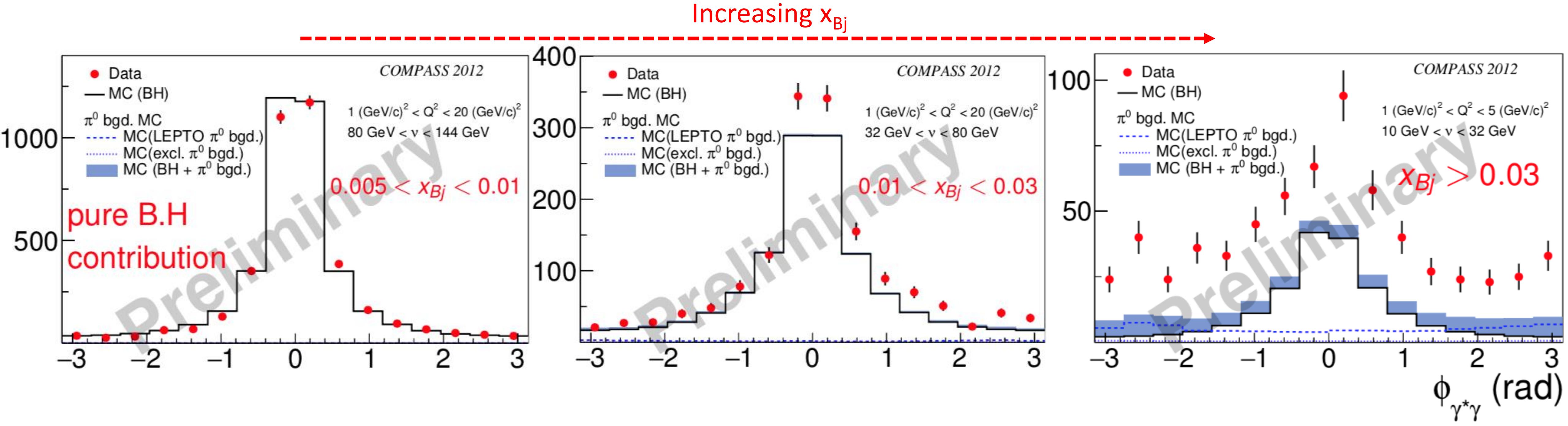


reference yield of almost pure Bethe-Heitler

Study DVCS with:
$\mathrm{Re}(T^{DVCS})$ & $\mathrm{Im}(T^{DVCS})$
via $(d\sigma^{+\leftarrow} \pm d\sigma^{-\rightarrow})$

Transverse Imaging:
$d\sigma^{DVCS}/dt$
via $(d\sigma^{+\leftarrow} + d\sigma^{-\rightarrow})$

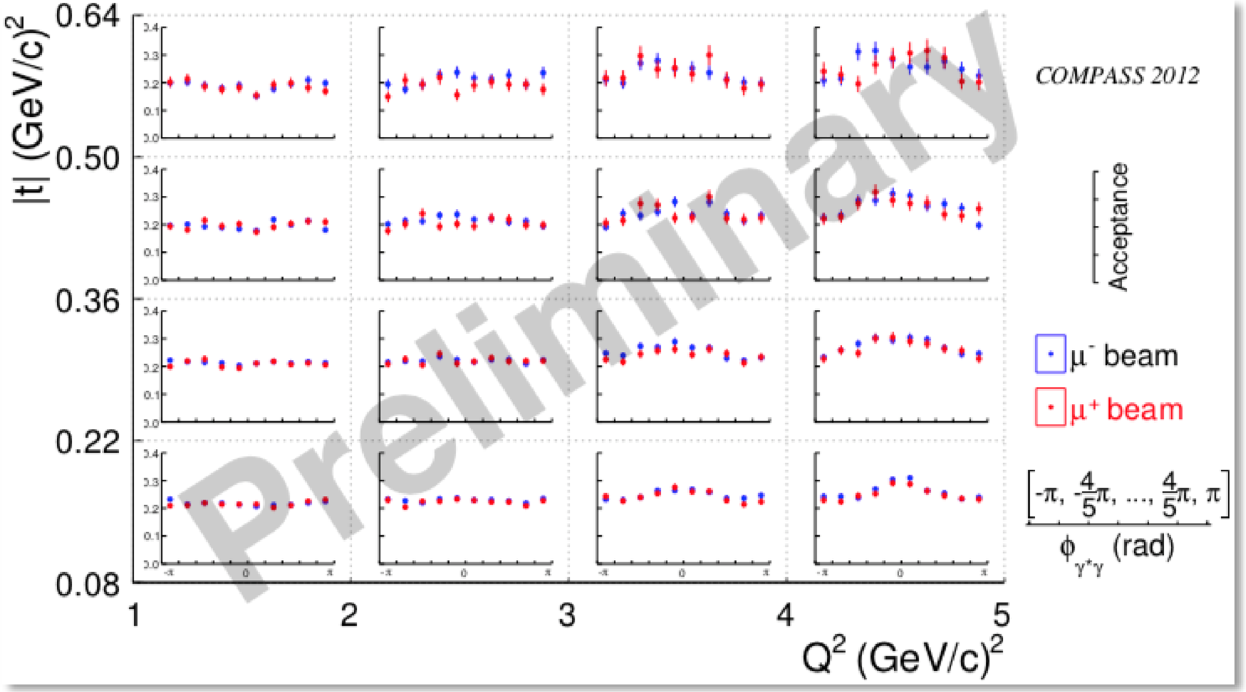
Extraction of DVCS Events - 2012 data

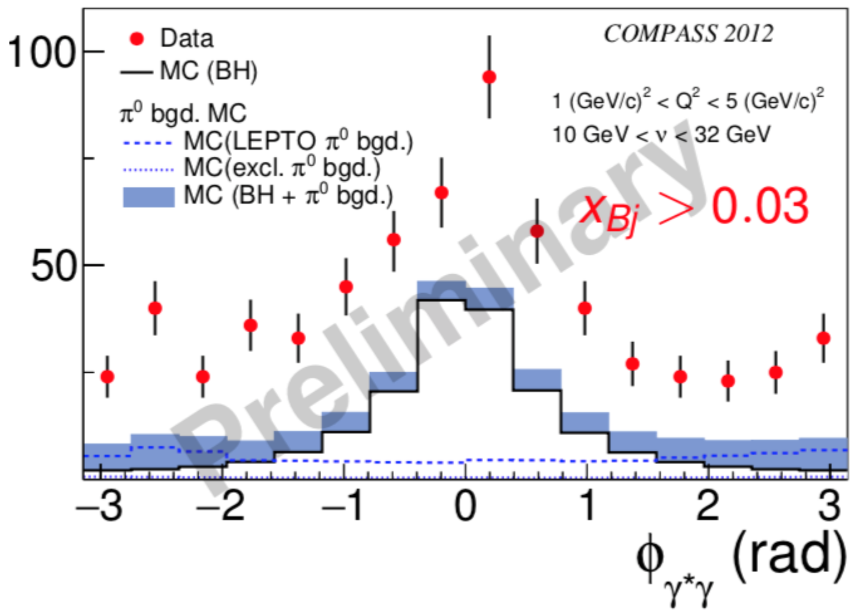


$$S_{CS,\,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) = 2[\, d\sigma^{BH} + d\sigma^{DVCS}_{unpol} + \mathrm{Im}\, I \,]$$

$$= 2[\, d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos\phi + c_2^{DVCS} \cos\phi + s_1^{I} \sin\phi + s_2^{I} \sin 2\phi \,]$$

Extraction of DVCS Events - 2012 data





$$S_{CS,\,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) = 2[\, d\sigma^{BH} + d\sigma^{DVCS}_{unpol} + \mathrm{Im}\, I \,]$$

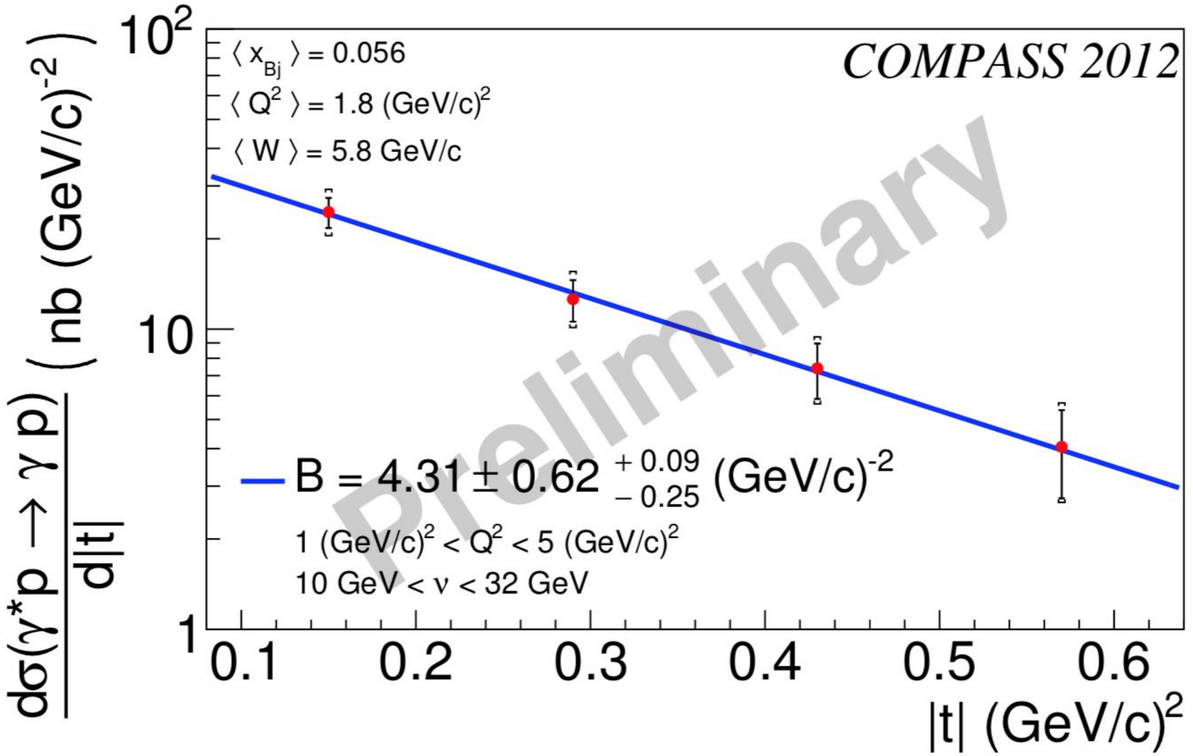
$$= 2[\, \cancel{d\sigma^{BH}} + c_0^{DVCS} + \cancel{c_1^{DVCS} \cos\phi + c_2^{DVCS} \cos\phi + s_1^{I} \sin\phi + s_2^{I} \sin 2\phi} \,]$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^* \rightarrow 4\ (Im\ \mathcal{H})^2$$

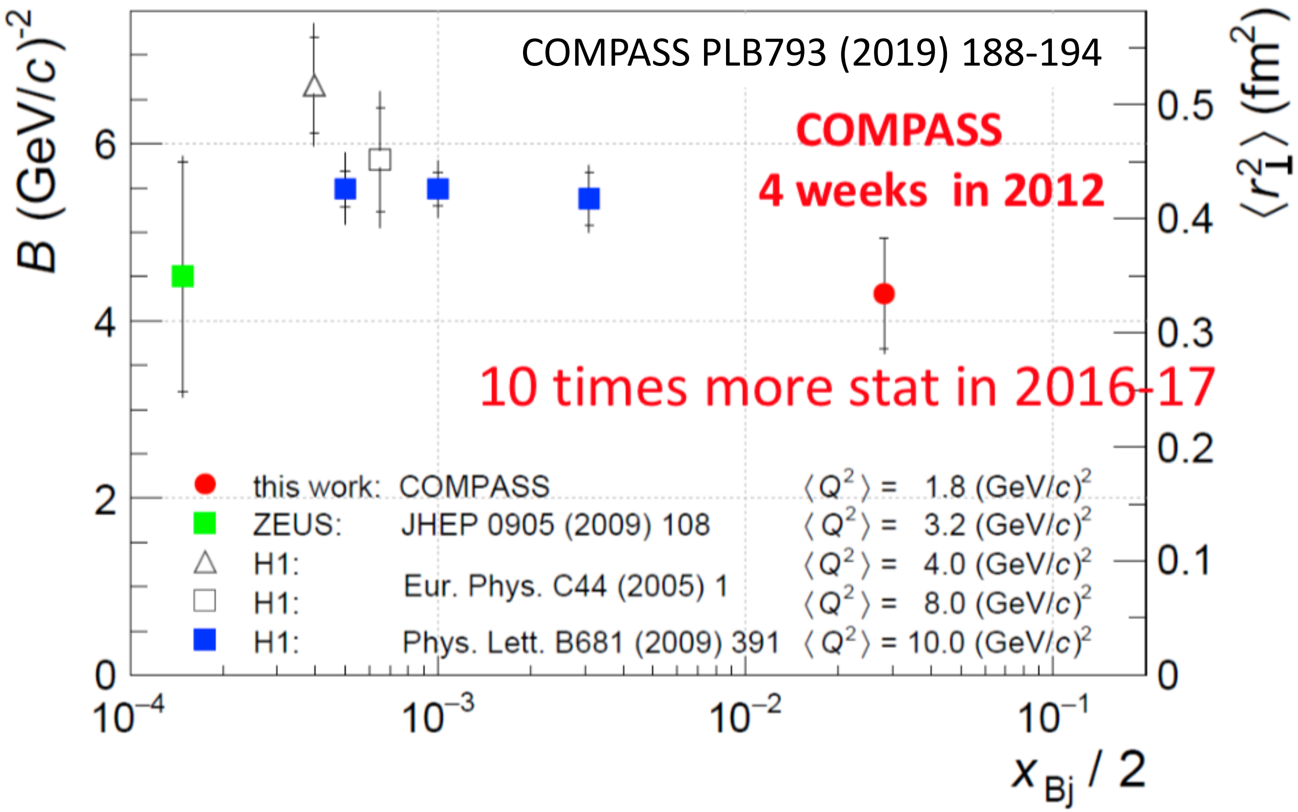
small x_{Bj} model dependent

Tranverse extension of partons – 2012 data

$$d\sigma^{DVCS}/d|t| \propto e^{-B|t|}$$



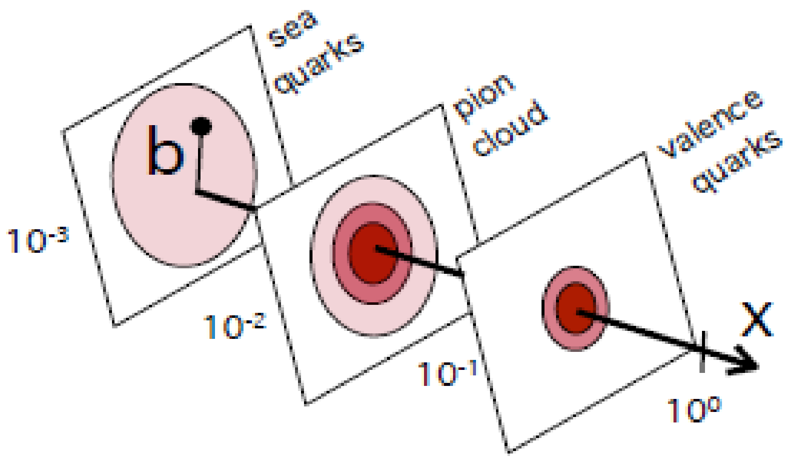
$$\langle r_\perp^2(x_{Bj})\rangle \approx 2B(x_{Bj})$$ At small x_{Bj}



$$\sqrt{\langle r_\perp^2\rangle} = \left(0.58 \pm 0.04_{\text{stat}} \,{}^{+\,0.01}_{-\,0.02}\Big|_{\text{sys}} \pm 0.04_{\text{model}}\right)\,\text{fm}$$

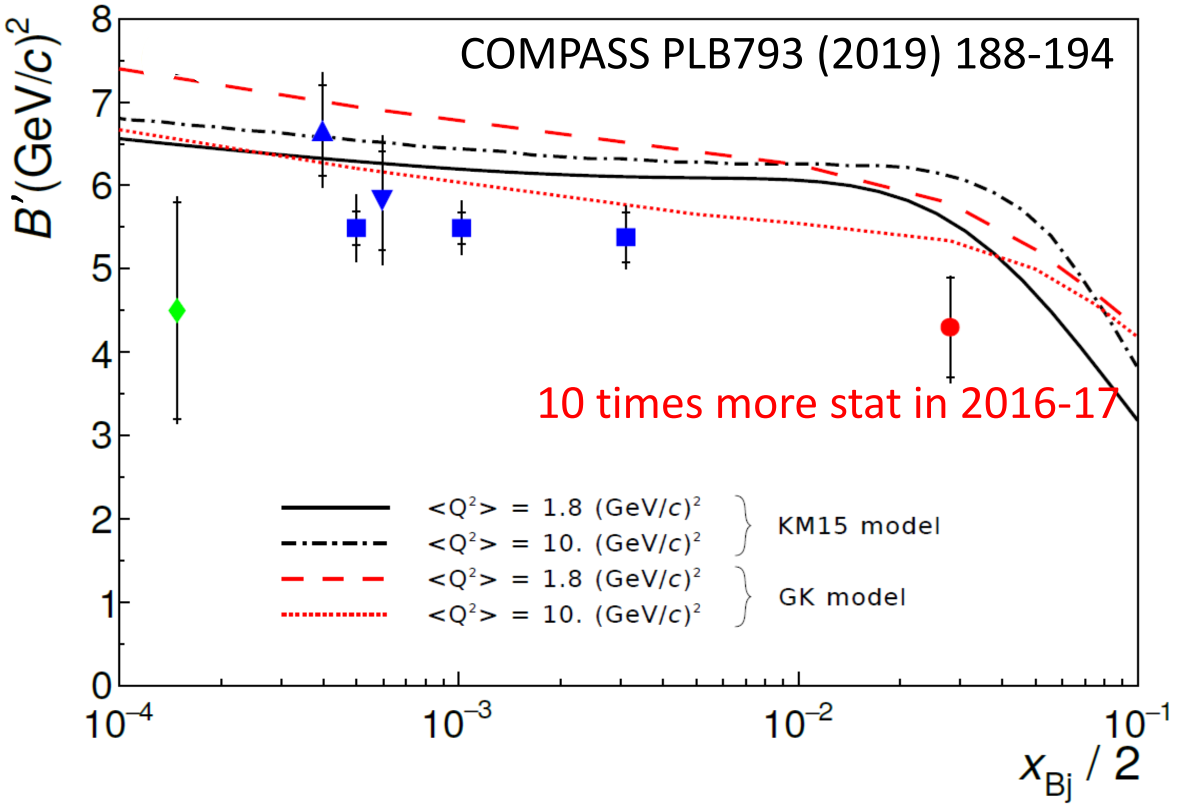
With $\langle x_{Bj}\rangle$ = 0.056

Tranverse extension of partons – 2012 data



- KM15 and GK parameterization shows the complete evolution of HERA and COMPASS data
- The transverse size evolution as a function of x_{Bj}
  → Expect at least 3 x_{Bj} bins from 2016-17 data

$$\langle r_\perp^2(x_{Bj})\rangle \approx 2B(x_{Bj})$$ At small x_{Bj}



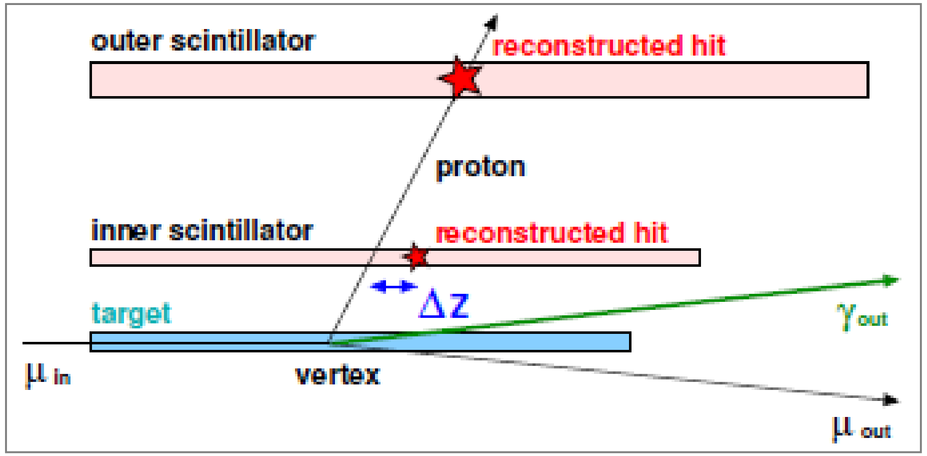
$$\sqrt{\langle r_\perp^2\rangle} = (0.58 \pm 0.04_{\mathrm{stat}}\ {}^{+\,0.01}_{-\,0.02}\big|_{\mathrm{sys}}\ \pm 0.04_{\mathrm{model}})\ \mathrm{fm}$$

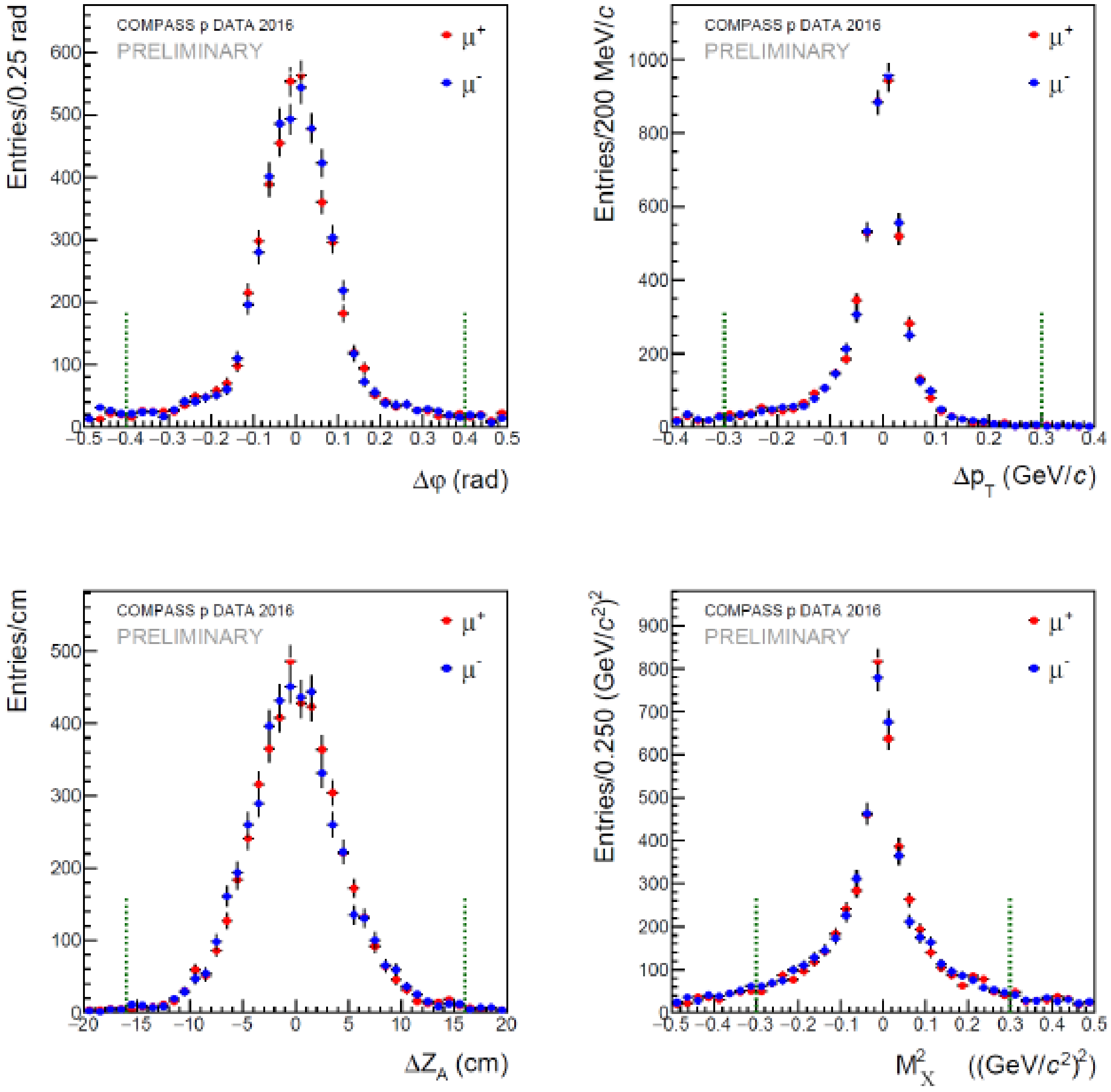
With $\langle x_{Bj}\rangle$ = 0.056

2016 – 2017 Data First Insight

DVCS : μ p → μ' p γ

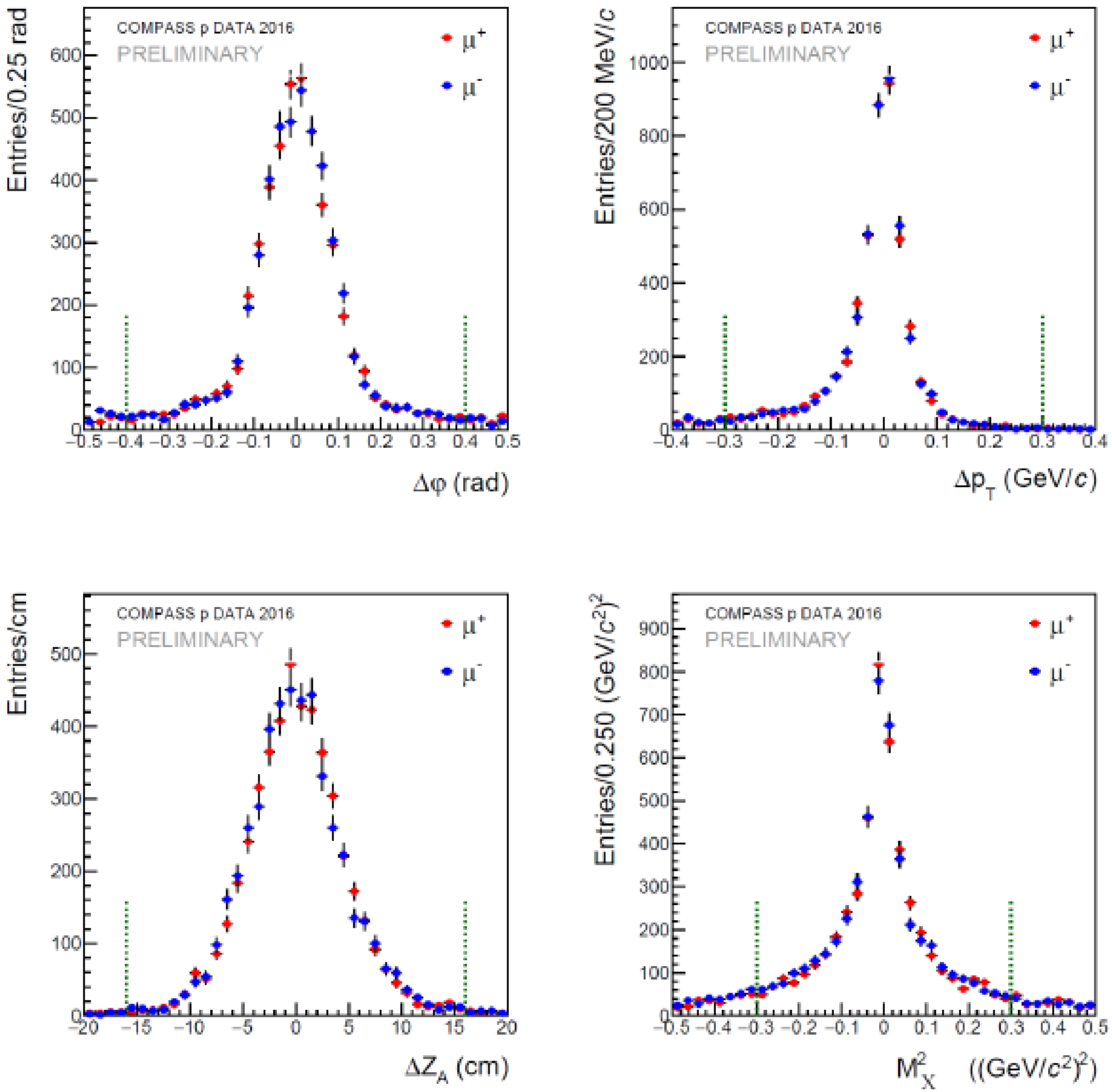
1) $\Delta p_T = p_T^{cam} - p_T^{spec}$
2) $\Delta\varphi = \varphi^{cam} - \varphi^{spec}$
3) $\Delta z_A = z_A^{cam} - z_A^{Z_B \text{ and vertex}}$
4) $M^2_{X=0} = (p_{\mu_{in}} + p_{p_{in}} - p_{\mu_{out}} - p_{p_{out}} - p_\gamma)^2$





2016 – 2017 Data First Insight

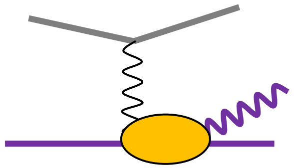
- All distributions obtained with μ^- are normalized to the same luminosity of the ones with μ^+
- Very good agreement between μ^+ and μ^- data observed.



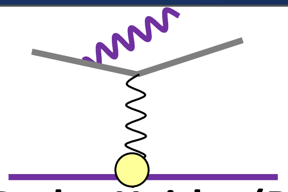
✓ $\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$

$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$

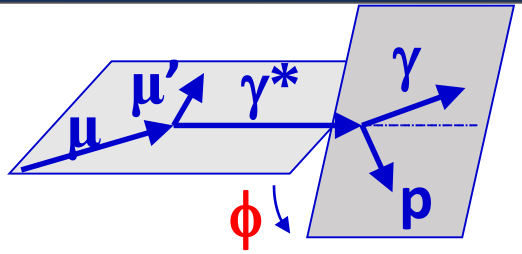
2016 – 2017 Data First Insight



DVCS

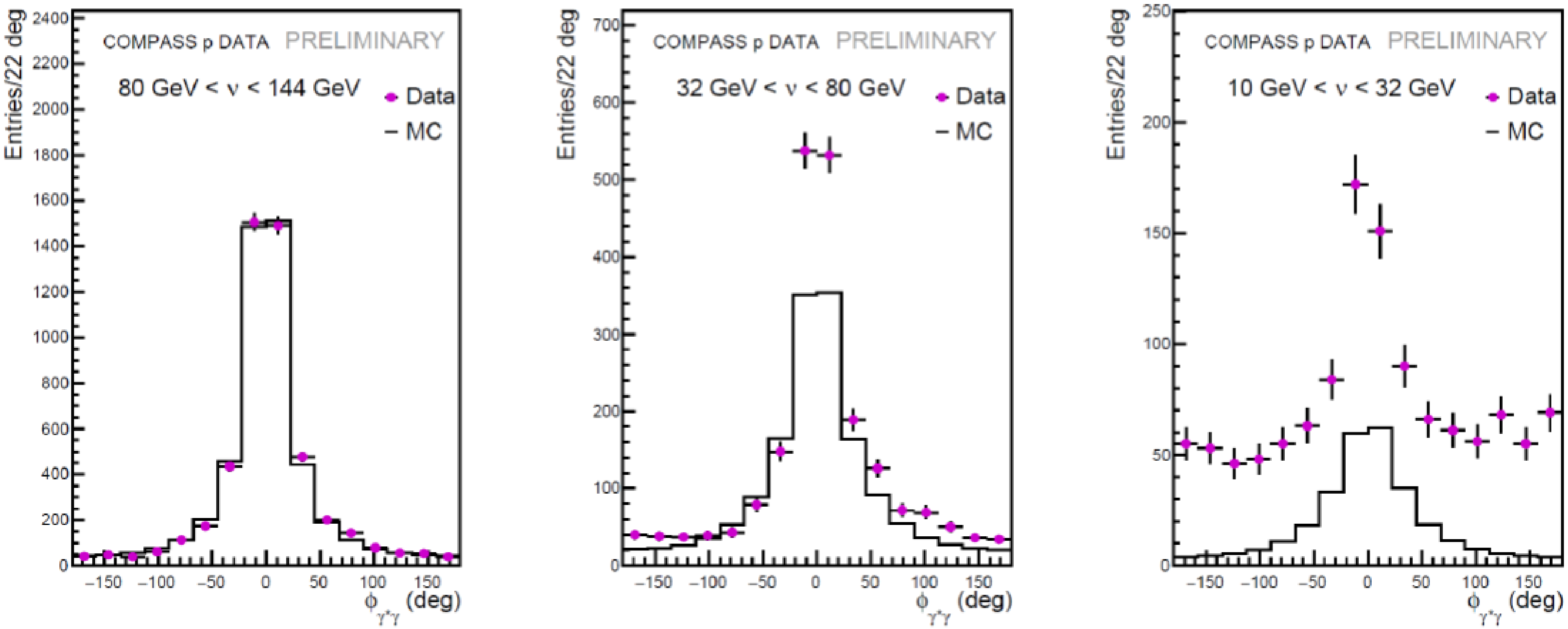


Bethe-Heitler (BH)



Only 13% of 2016-17 data

$0.005 < x_{Bj} < 0.01$ $\quad$ $0.01 < x_{Bj} < 0.03$ $\quad$ $x_{Bj} > 0.03$



No π^0 subtraction.

DVCS contribution at high x_{Bj} will allow to perform re-analysis

$d\sigma^{DVCS}/dt = e^{-B'|t|}$

$= c_0^{DVCS}$

BLUE WATERS

This research is part of the Blue Waters sustained-petascale computing project, which is supported by the National Science Foundation (awards OCI-0725070 and ACI-1238993) and the state of Illinois. Blue Waters is a joint effort of the University of Illinois at Urbana-Champaign and its National Center for Supercomputing Applications. This work is also part of the "Mapping Proton Quark Structure using Petabytes of COMPASS Data" PRAC allocation supported by the National Science Foundation (award number OCI 1713684).

Beam Charge-spin Difference

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \rightarrow c_0^I + c_1^I \cos\phi$$

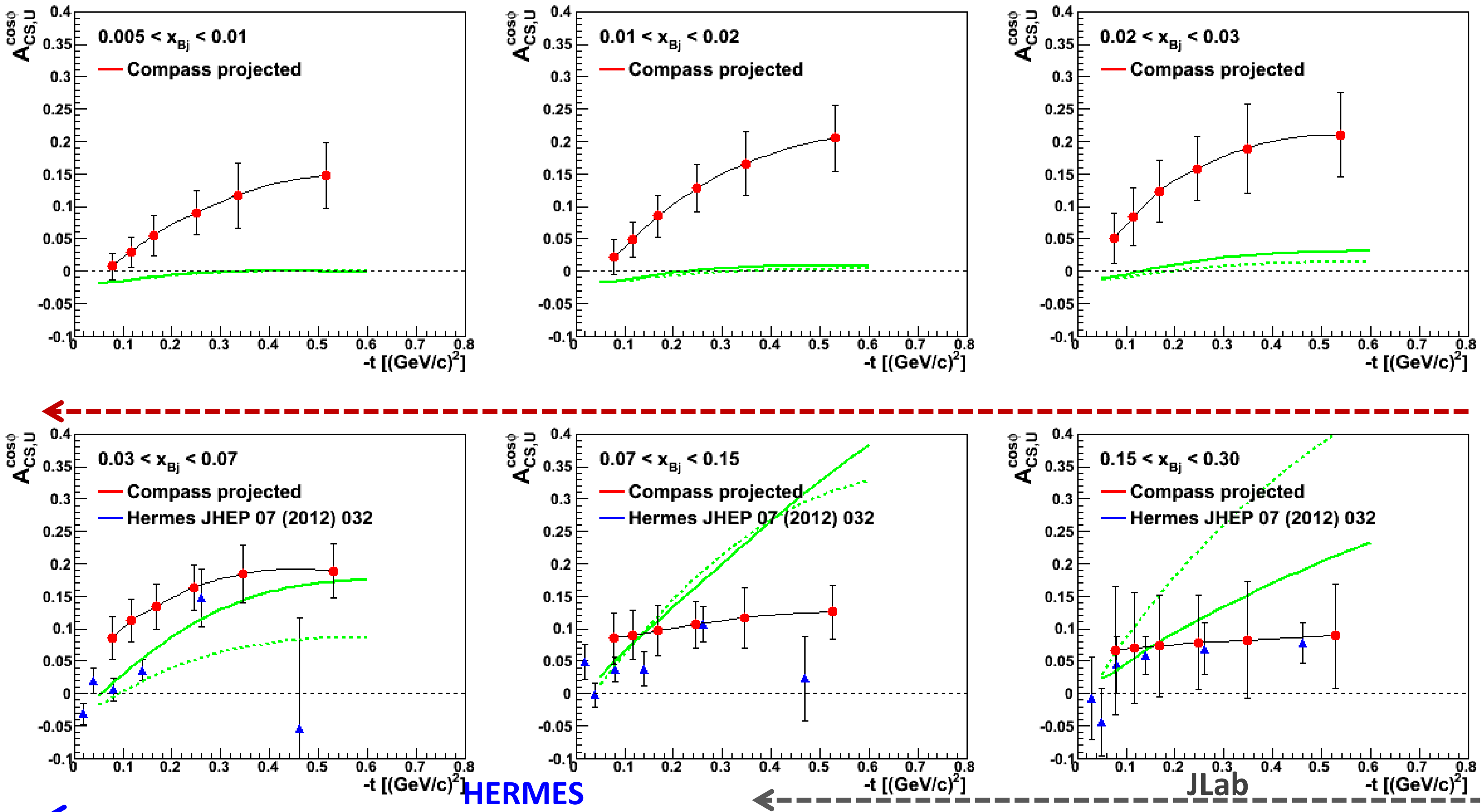
$$c_1^I = \mathcal{R}e\, F_1 \mathcal{H}$$

KM10

VGG

- With and $\mathcal{R}e\, F_1 \mathcal{H}$ and $\mathcal{I}m\, F_1 \mathcal{H}$
  → Extraction of D-term

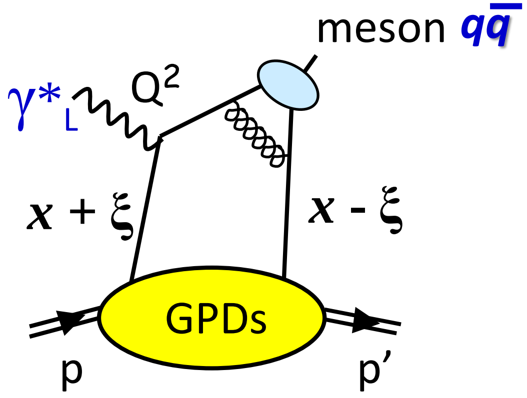
$\mathcal{R}e\, \mathcal{H}$ > 0 at H1
< 0 at HERMES
Value of x_{Bj} for the node?



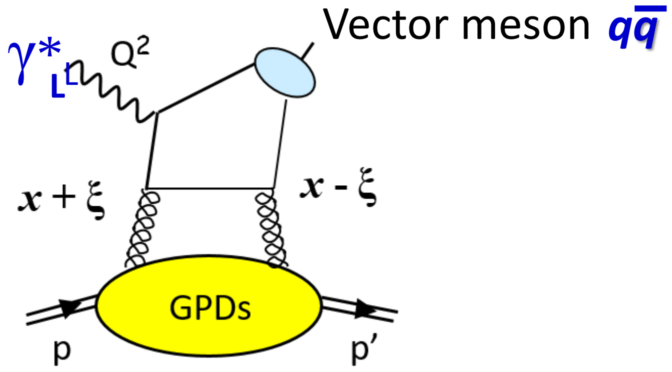
COMPASS 2 years of data Eμ= 160 GeV 1 < Q^2 < 8 GeV^2

GPDs in Hard Exclusive Meson Production

Quark contribution



Gluon contribution at the same order in α_s



4 chiral-even GPDs: helicity of parton unchanged

$H^q(x, \xi, t)$	$E^q(x, \xi, t)$	**For Vector Meson**
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$	**For Pseudo-Scalar Meson**

+ 4 chiral-odd or transversity GPDs: helicity of parton changed
(not possible in DVCS)

$H^q_T(x, \xi, t)$	$E^q_T(x, \xi, t)$	$\bar{E}^q_T = 2\tilde{H}^q_T + E^q_T$
$\tilde{H}^q_T(x, \xi, t)$	$\tilde{E}^q_T(x, \xi, t)$	

- Universality of GPDs, quark flavor filter
- Ability to probe the chiral-odd GPDs.
- Additional non-perturbative term from meson wave function
- In addition to nuclear structure, provide insights into reaction mechanism

Exclusive π^0 Production on Unpolarized Proton

e p → e π^0 p

$$\frac{d^2\sigma}{dt\,d\phi_\pi} = \frac{1}{2\pi}\left[\boxed{\left(\frac{d\sigma_T}{dt} + \epsilon\frac{d\sigma_L}{dt}\right)} + \epsilon\cos 2\phi_\pi \boxed{\frac{d\sigma_{TT}}{dt}} + \sqrt{2\epsilon(1+\epsilon)}\cos\phi_\pi \boxed{\frac{d\sigma_{LT}}{dt}}\right]$$

$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'}\frac{1}{Q^6}\left\{(1-\xi^2)\left|\langle\tilde{H}\rangle\right|^2 - 2\xi^2\mathrm{Re}\left[\langle\tilde{H}\rangle^*\langle\tilde{E}\rangle\right] - \frac{t'}{4m^2}\xi^2\left|\langle\tilde{E}\rangle\right|^2\right\}$$

Leading twist expected be dominant

But measured as ≈ only a few % of $\frac{d\sigma_T}{dt}$

The other contributions arise from coupling between chiral-odd (quark helicity flip) GPDs to the twist-3 pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'}\frac{\mu_\pi^2}{Q^8}\left[(1-\xi^2)\left|\langle H_T\rangle\right|^2 - \frac{t'}{8m^2}\left|\langle\bar{E}_T\rangle\right|^2\right]$$

$$\frac{\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'}\frac{\mu_\pi}{Q^7}\xi\sqrt{1-\xi^2}\frac{\sqrt{-t'}}{2m}\mathrm{Re}\left[\langle H_T\rangle^*\langle\tilde{E}\rangle\right]$$

$$\frac{\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'}\frac{\mu_\pi^2}{Q^8}\frac{t'}{16m^2}\left|\langle\bar{E}_T\rangle\right|^2$$

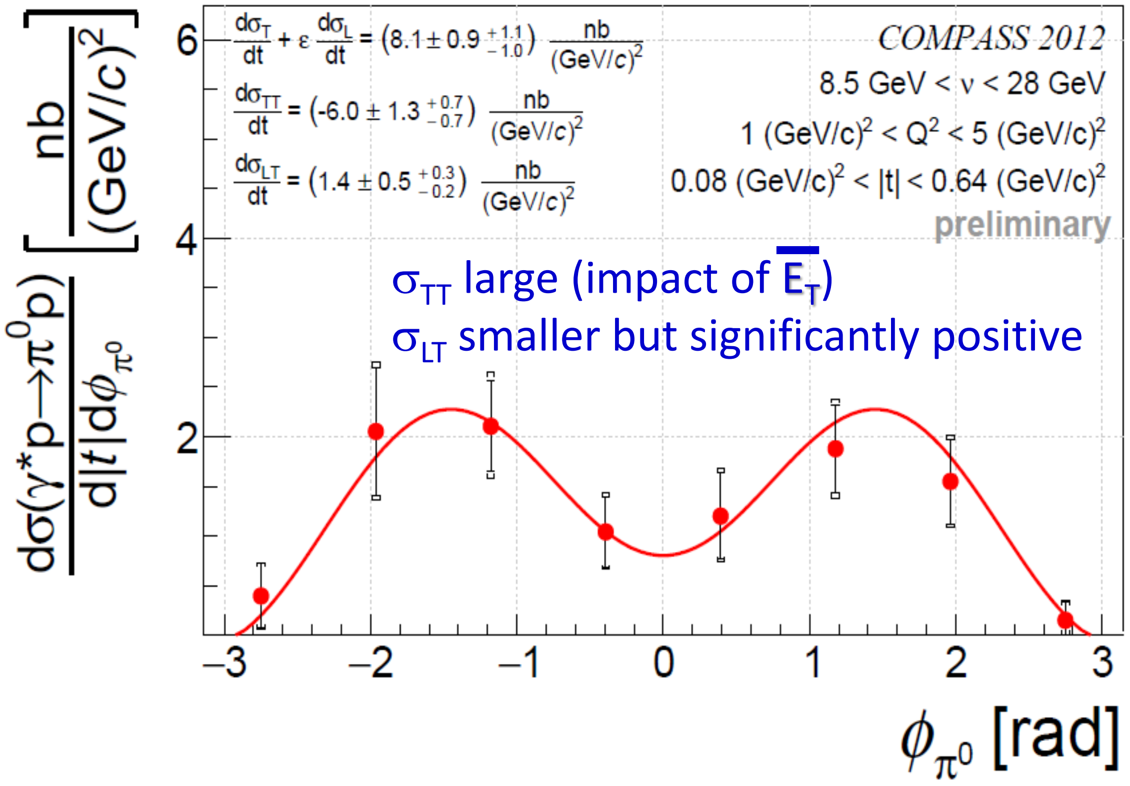
A large impact of $\bar{E}_T$ can be identified:

- σ_{TT} contribution
- The dip at small |t| of σ_T

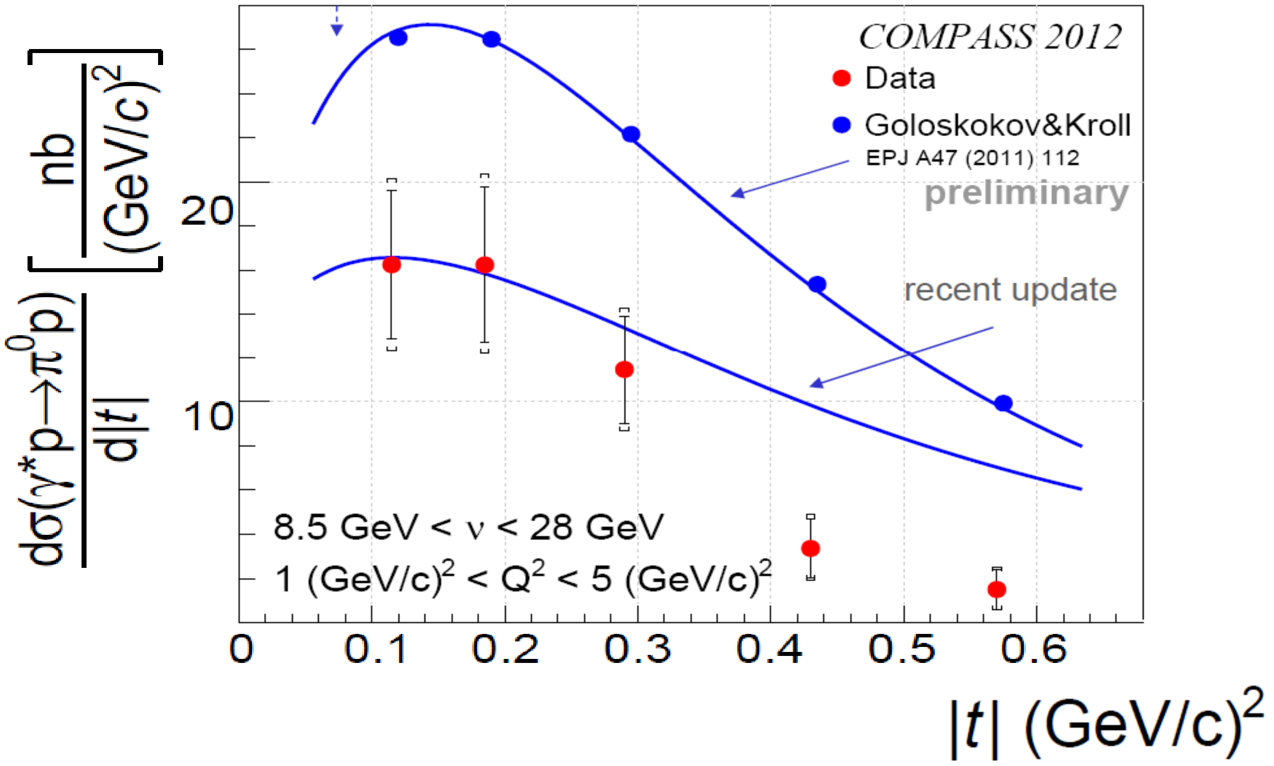
2012 Exclusive π^0 Prod. on Unpolarized Proton

e p → e π^0 p

$$\frac{d^2\sigma}{dt\,d\phi_\pi} = \frac{1}{2\pi}\left[\boxed{\left(\frac{d\sigma_T}{dt} + \epsilon\frac{d\sigma_L}{dt}\right)} + \epsilon\cos 2\phi_\pi \boxed{\frac{d\sigma_{TT}}{dt}} + \sqrt{2\epsilon(1+\epsilon)}\cos\phi_\pi \boxed{\frac{d\sigma_{LT}}{dt}}\right]$$



A dip at small t would indicate a large impact of $\bar{E}_T$



hep-ex/1903.12030, subm. to PLB

2012 Exclusive ω Prod. on Unpolarized Proton

SCHC ($\lambda_\gamma = \lambda_V$)

(S-Channel Helicity Conservation)

SCHC implies:

- $r^1_{1-1} + \mathrm{Im}\, r^2_{1-1} = 0$

= - 0.010 ± 0.032 ± 0.047 OK

- $\mathrm{Re}\, r^5_{10} + \mathrm{Im}^6_{10} = 0$

= 0.014 ± 0.011 ± 0.013 OK

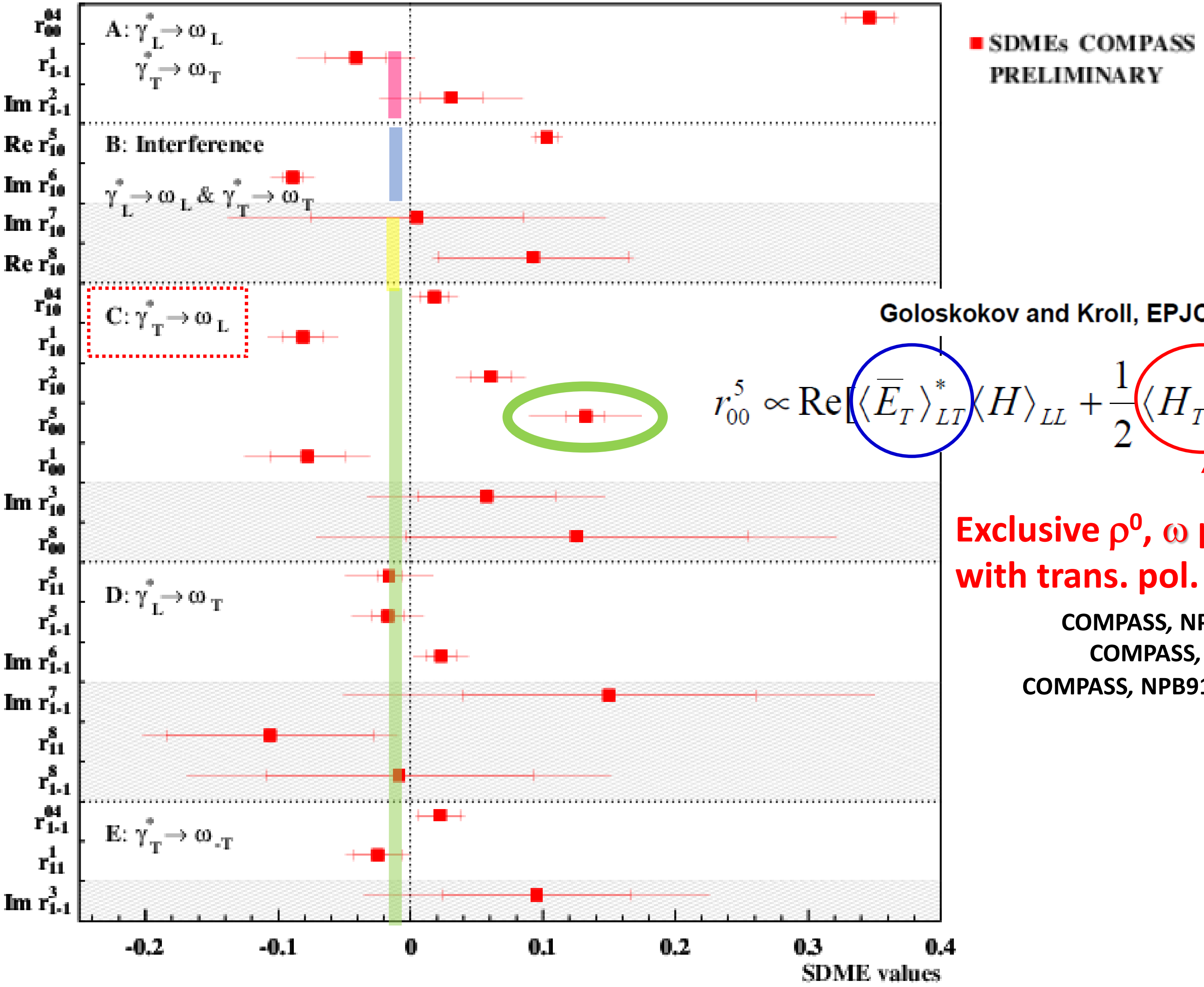
- $\mathrm{Im}\, r^7_{10} - \mathrm{Re}\, r^8_{10} = 0$

= - 0.088 ± 0.110 ± 0.196 OK

- all elements of classes C, D, E should be 0

for $\gamma^*_L \to \omega_T$ and $\gamma^*_T \to \omega_{-T}$ OK within errors

not obeyed for transitions $\gamma^*_T \to \omega_L$



Goloskokov and Kroll, EPJC 74 (2014) 2725

$$r^5_{00} \propto \mathrm{Re}[\langle \overline{E}_T \rangle^*_{LT} \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle^*_{LT} \langle E \rangle_{LL}]$$

Exclusive ρ^0, ω production with trans. pol. target

COMPASS, NPB865 (2012) 1-20
COMPASS, PLB731 (2014) 19
COMPASS, NPB915 (2017) 454-475

Summary

➢ GPD by DVCS and HEMP in COMPASS

Sum and difference of DVCS x-sections with polarized μ+ and μ-

→ Transverse extension of partons as a function of x_{Bj}

→ Im$\mathcal{H}(\xi,t)$ and Re$\mathcal{H}(\xi,t)$ for D-term and pressure distribution

HEMP of π^0, ρ, ω, ϕ, J/ψ

→ Universality of GPDs - Transverse GPDs - Flavor Decomposition

Results on 2016-17 data are coming!

Backup Slides

COMPASS++/AMBER

A new QCD facility at the M2 beam line of the CERN SPS

Letter of Intent - Draft 1.0: https://arXiv.org/abs/1808.0084

Expected to start at 2022

- Unique beam line with polarised $\mu^{\pm}$ and high-intensity **Pion** beam
- Possible high-intensity **antiproton** and **Kaon** beams, provided by RF-separation technique
- With upgraded apparatus

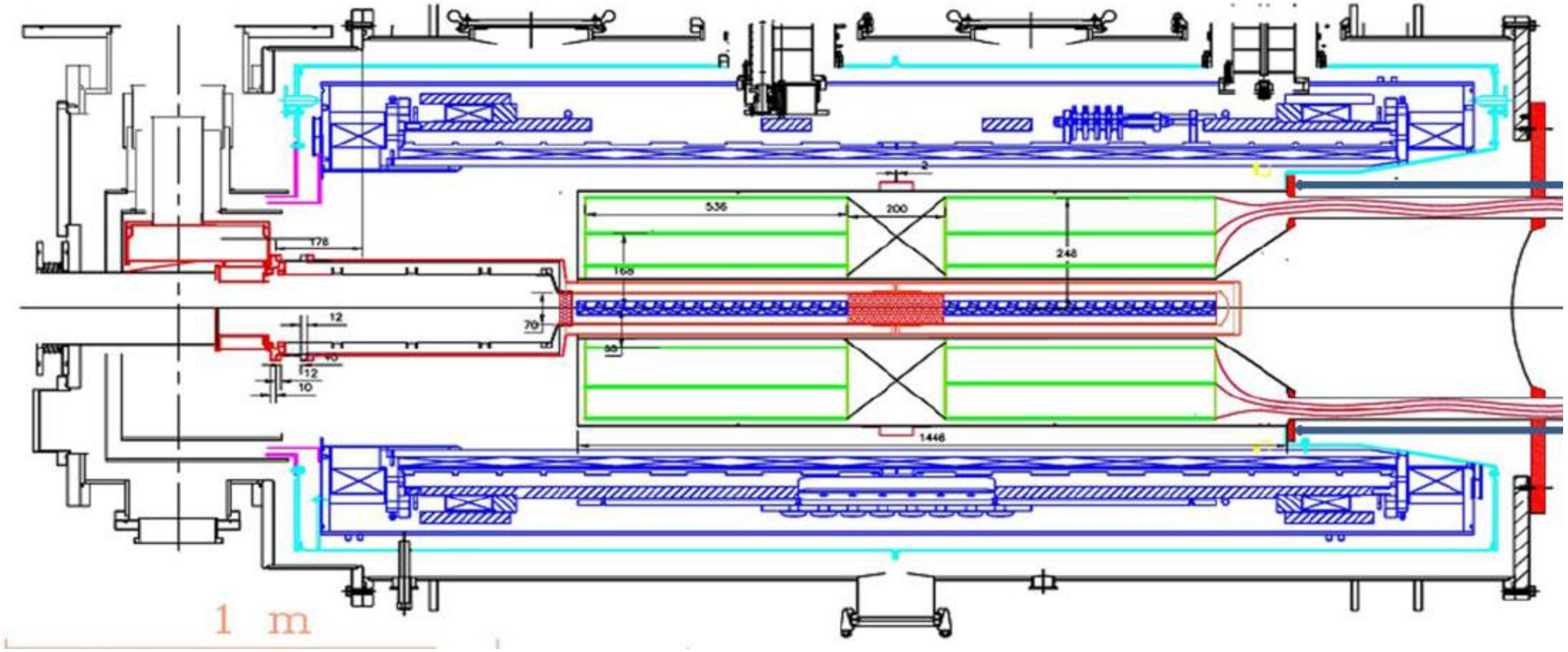
Proposed physics goals

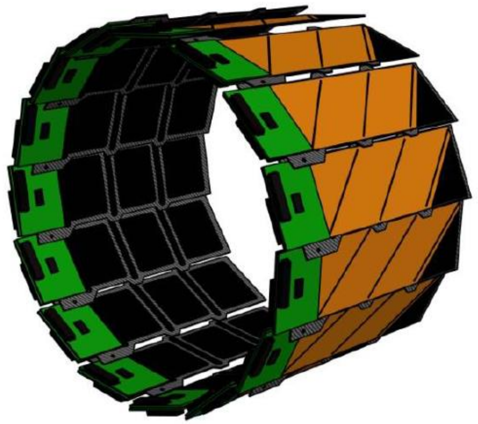
Proton Radius
Meson PDF – gluon PDF
Proton spin structure
3D imaging (TMDs and GPDs)
Hadron spectroscopy
Anti-matter cross section

Program	Physics Goals	Beam Energy [GeV]	Beam Intensity [s^{-1}]	Trigger Rate [kHz]	Beam Type	Target	Earliest start time, duration	Hardware Additions
μp elastic scattering	Precision proton-radius measurement	100	$4 \cdot 10^6$	100	$\mu^{\pm}$	high-pressure H2	2022 1 year	active TPC, SciFi trigger, silicon veto,
Hard exclusive reactions	GPD E	160	$2 \cdot 10^7$	10	$\mu^{\pm}$	$NH_3^{\uparrow}$	2022 2 years	recoil silicon, modified PT magnet
Input for Dark Matter Search	$\bar{p}$ production cross section	20-280	$5 \cdot 10^5$	25	p	LH2, LHe	2022 1 month	LHe target
$\bar{p}$-induced Spectroscopy	Heavy quark exotics	12, 20	$5 \cdot 10^7$	25	$\bar{p}$	LH2	2022 2 years	target spectr.: tracking, calorimetry
Drell-Yan	Pion PDFs	190	$7 \cdot 10^7$	25	$\pi^{\pm}$	C/W	2022 1-2 years	
Drell-Yan (RF)	Kaon PDFs & Nucleon TMDs	$\sim$100	10^8	25-50	$K^{\pm}, \bar{p}$	$NH_3^{\uparrow}$, C/W	2026 2-3 years	"active absorber", vertex det.
Primakoff (RF)	Kaon polarisa-bility & pion life time	$\sim$100	$5 \cdot 10^6$	> 10	K^-	Ni	non-exclusive 2026 1 year	
Prompt Photons (RF)	Meson gluon PDFs	≥ 100	$5 \cdot 10^6$	10-100	$K^{\pm}$ $\pi^{\pm}$	LH2, Ni	non-exclusive 2026 1-2 years	hodoscope
K-induced Spectroscopy (RF)	High-precision strange-meson spectrum	50-100	$5 \cdot 10^6$	25	K^-	LH2	2026 1 year	recoil TOF, forward PID
Vector mesons (RF)	Spin Density Matrix Elements	50-100	$5 \cdot 10^6$	10-100	$K^{\pm}, \pi^{\pm}$	from H to Pb	2026 1 year	

Possible RPD for COMPASS^{++}/AMBER

A recoil proton detector (RPD) is mandatory to ensure the exclusivity. A Silicon detector is included *between* the target surrounded by the modified MW cavity *and* the polarizing magnet





A technology developed at JINR for NICA for the BM@N experiment

No possibility for ToF ➔ PID of p/π with dE/dx
Momentum and trajectory measurments
$|t|_{min} \sim 0.1$ GeV

π^0 Background

π^0 are one of the main background sources for excl. photon events.

Two possible case:

- **Visible** (both γ detected ➔ subtracted)

  the DVCS photon after all exclusivity cuts is combined with all detected photons below the DVCS threshold: 4,5,10 GeV in ECAL0, 1, 2

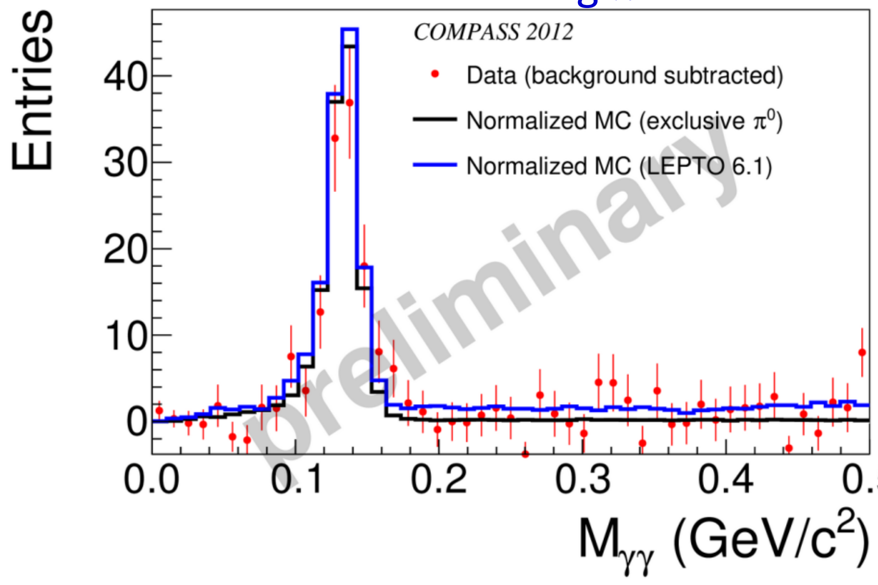
- **Invisible** (one γ lost ➔ estimated by MC)

  - **Semi-inclusive LEPTO 6.1**

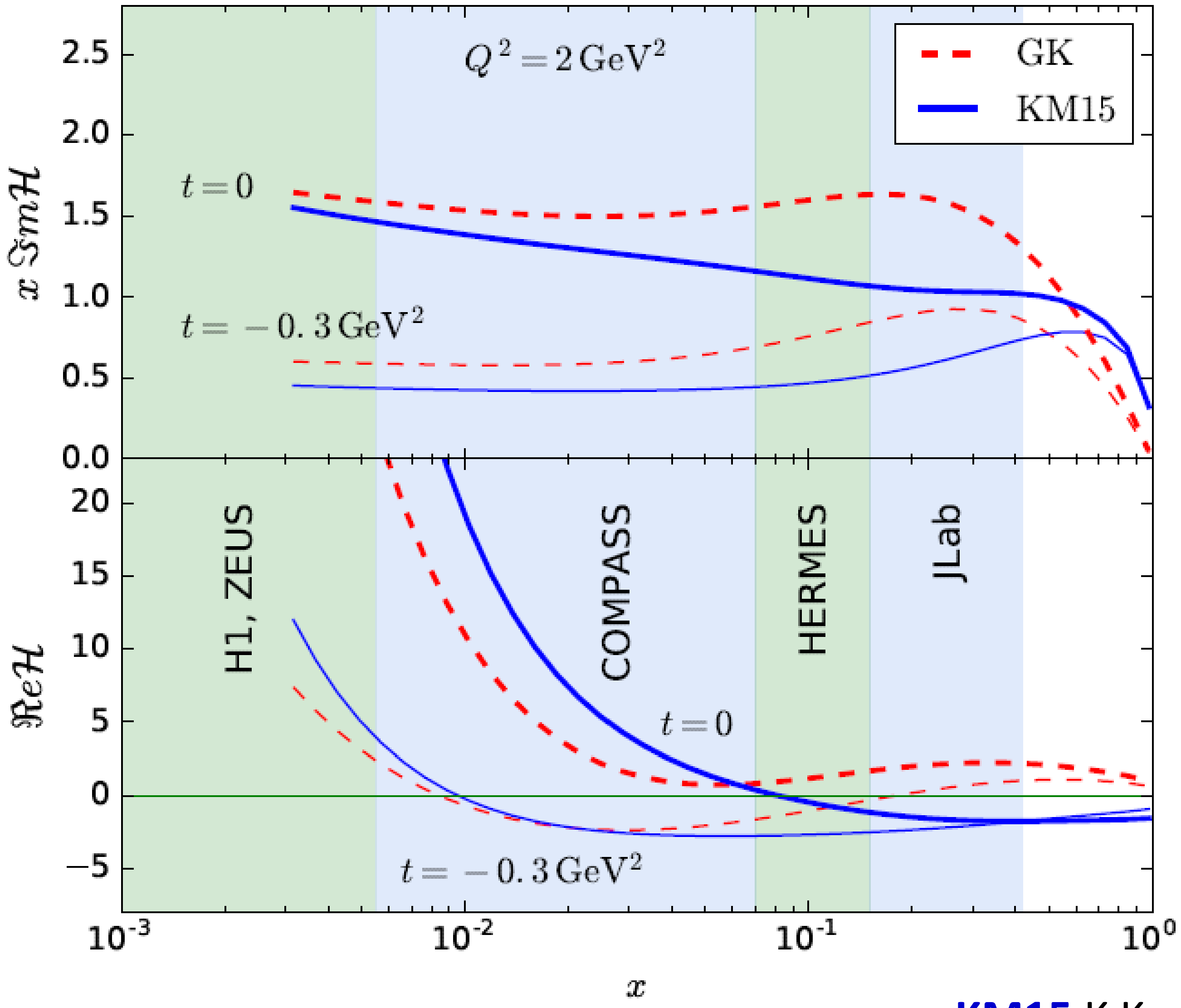
  - **Exclusive HEPGEN π^0** (Goloskokov-Kroll model)

Comparing the two components to the data allows the determination of their relative normalisation.
The sum of the 2 components is normalized to the visible π^0 contamination in the $M_{\gamma\gamma}$ peak

Visible leaking π^0 in the data



GPD H Global Analysis



Im $\mathcal{H}$
is better understood

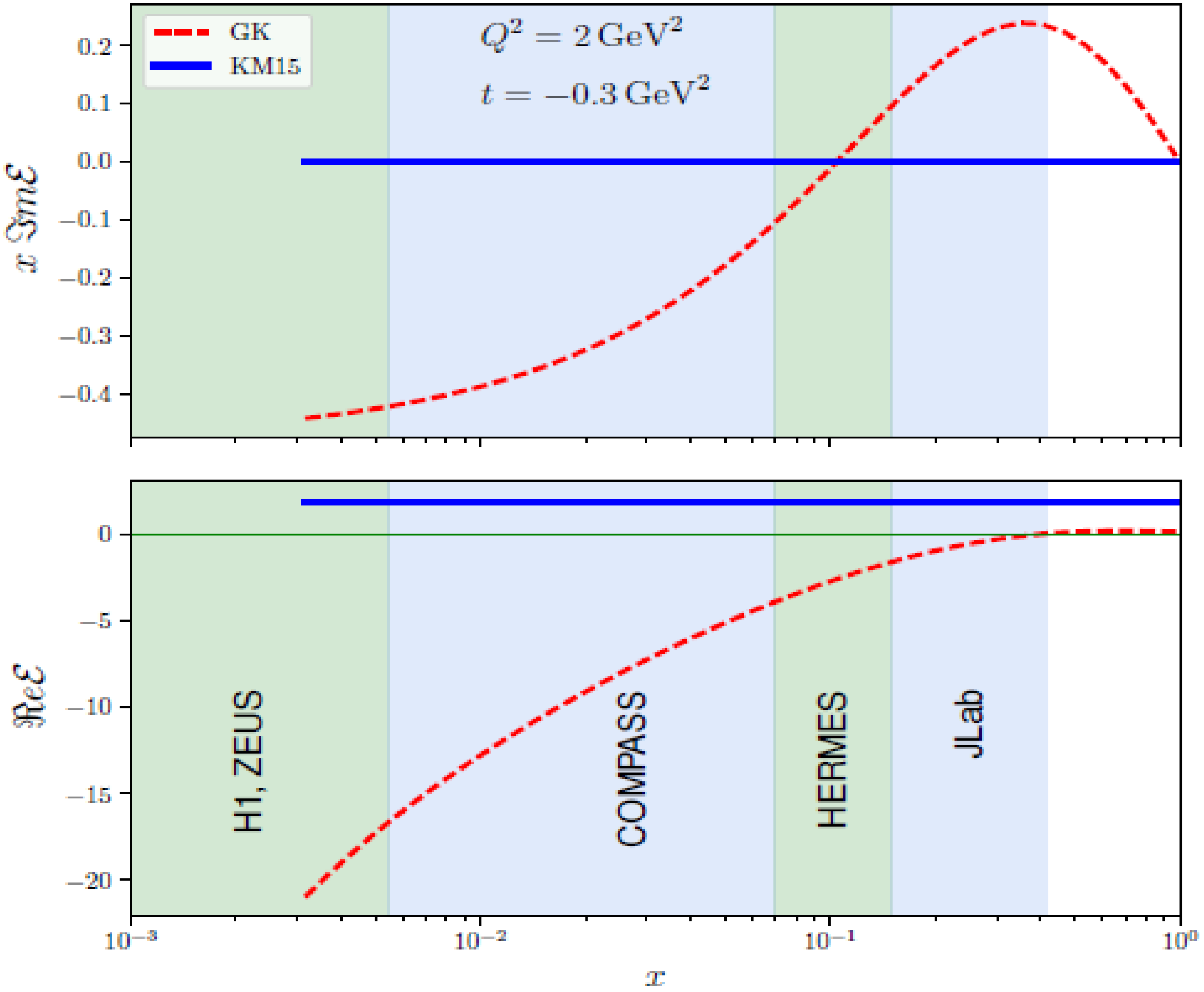
Re $\mathcal{H}$ linked to the
*d term*is still poorly
constrained

KM15 K Kumericki and D Mueller [arXiv:1512.09014v1]

GK S.V. Goloskokov, P. Kroll, EPJC53 (2008), EPJA47 (2011)

GPD E Global Analysis

Figure made by D. Mueller and K. Kumericki



Im E
is rather unknown

Re E
is rather unknown

KM15 K Kumericki and D Mueller [arXiv:1512.09014v1]

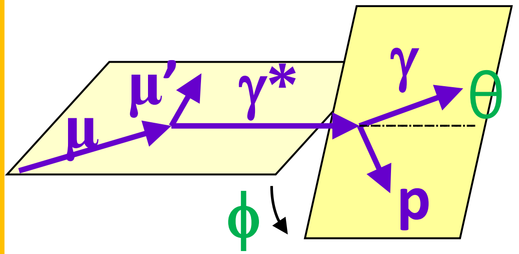
GK S.V. Goloskokov, P. Kroll, EPJC53 (2008), EPJA47 (2011)

ϕ Dep. of BH+DVCS with Unpol Target

$$\frac{d^4\sigma(\ell p \to \ell p \gamma)}{dx_B\, dQ^2\, d|t|\, d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma^{DVCS}_{unpol} + P_\ell\, d\sigma^{DVCS}_{pol}\right) + (e_\ell \mathrm{Re}\, I + e_\ell P_\ell\, \mathrm{Im}\, I)$$

e-	μ±
Σ	Σ
Σ	Σ
Δ	Δ
Σ	Δ
Δ	Σ
⇆	⇆

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH}\cos\phi + c_2^{BH}\cos 2\phi$$
$$d\sigma^{DVCS}_{unpol} \propto c_0^{DVCS} + c_1^{DVCS}\cos\phi + c_2^{DVCS}\cos 2\phi$$
$$d\sigma^{DVCS}_{pol} \propto s_1^{DVCS}\sin\phi$$
$$\mathrm{Re}\, I \propto c_0^I + c_1^I\cos\phi + c_2^I\cos 2\phi + c_3^I\cos 3\phi$$
$$\mathrm{Im}\, I \propto s_1^I\sin\phi + s_2^I\sin 2\phi$$



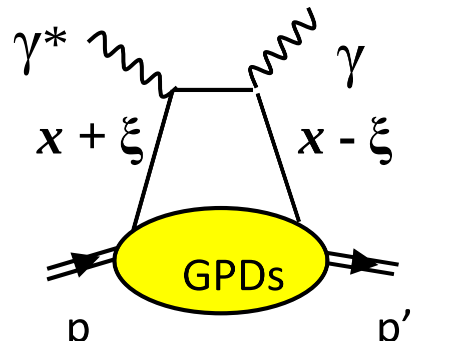
Twist-2 >>

■ Twist-3,

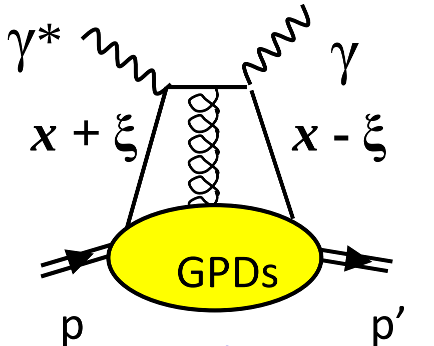
■ Twist-2

double helicity flip

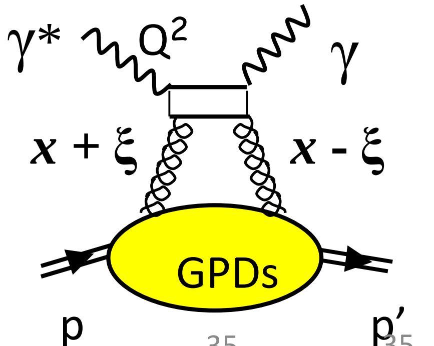
for gluons (NLO)



Twist-2



Twist-3



Twist-2, NLO

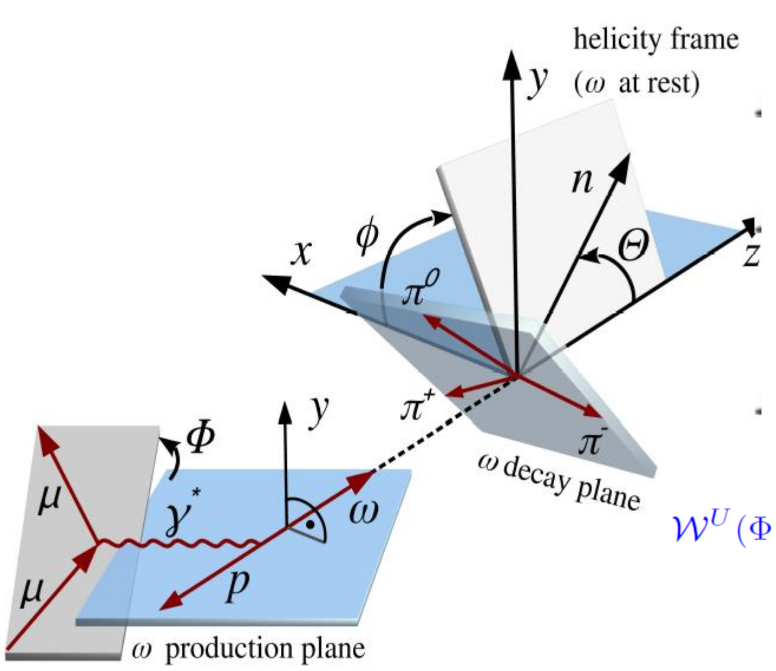
s_1^I = *Im* $\mathcal{F}$

c_1^I = *Re* $\mathcal{F}$

$$\mathcal{F} = F_1\mathcal{H} + \xi(F_1 + F_2)\tilde{\mathcal{H}} - t/4m^2 F_2\mathcal{E} \xrightarrow[\text{for proton}]{\text{at small } x_B} F_1\mathcal{H}$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^* \xrightarrow{\text{at small } x_B} 4\ (\mathit{Im}\ \mathcal{H})^2$$

Exclusive ω Production on Unpolarized Proton



experimental angular distributions

$$\mathcal{W}^{U+L}(\Phi,\phi,\cos\Theta) = \mathcal{W}^{U}(\Phi,\phi,\cos\Theta) + P_b \mathcal{W}^{L}(\Phi,\phi,\cos\Theta)$$

Self analysis... with phi and cosTheta

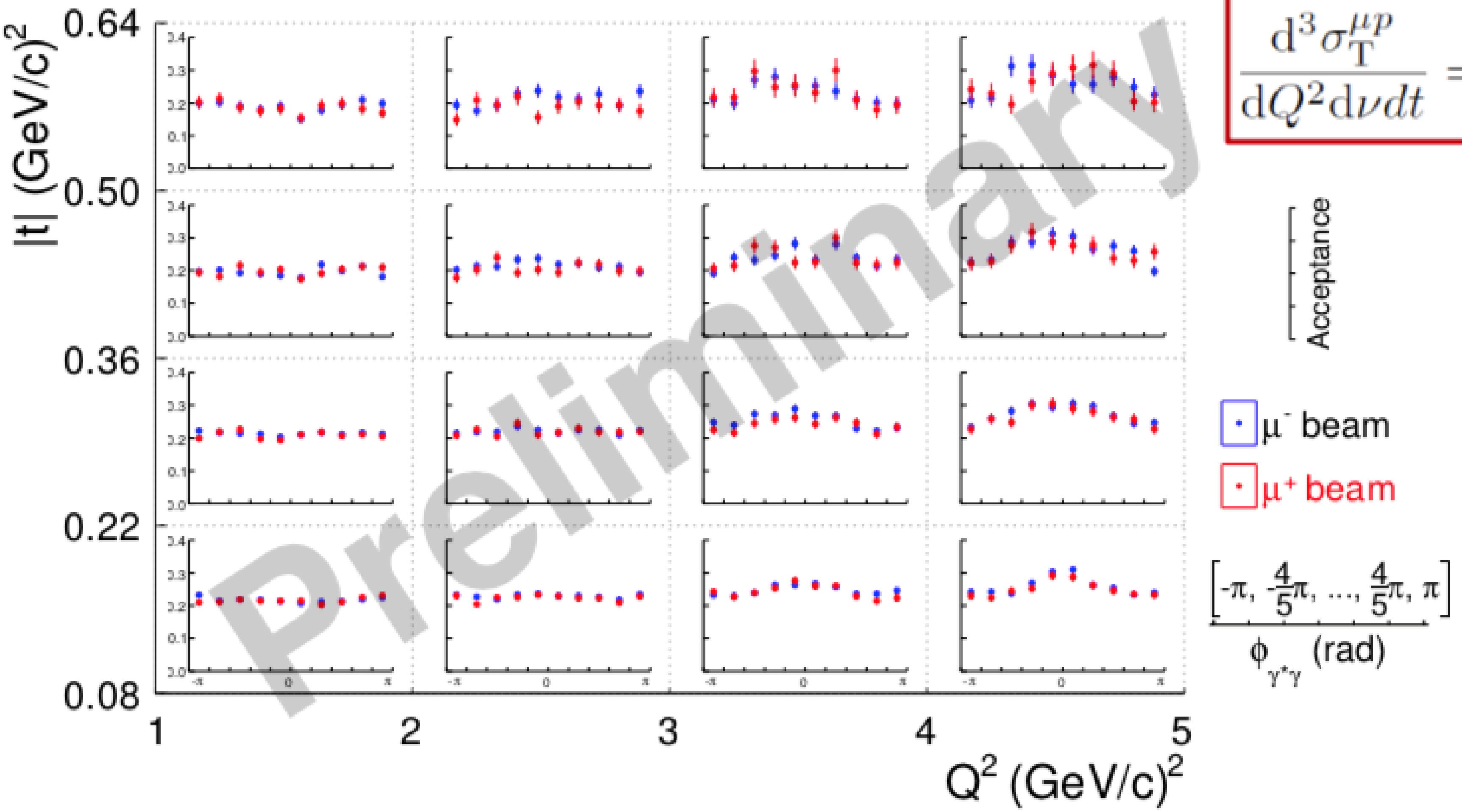
15 'unpolarized' and 8 'polarized' SDMEs

$$\begin{aligned}
\mathcal{W}^{U}(\Phi,\phi,\cos\Theta) = \frac{3}{8\pi^2}\Big[& \frac{1}{2}(1 - r^{04}_{00}) + \frac{1}{2}(3r^{04}_{00} - 1)\cos^2\Theta - \sqrt{2}\mathrm{Re}\{r^{04}_{10}\}\sin 2\Theta\cos\phi - r^{04}_{1-1}\sin^2\Theta \\
& -\epsilon\cos 2\Phi\Big(r^{1}_{11}\sin^2\Theta + r^{1}_{00}\cos^2\Theta - \sqrt{2}\mathrm{Re}\{r^{1}_{10}\}\sin 2\Theta\cos\phi - r^{1}_{1-1}\sin^2\Theta\cos 2\phi\Big) \\
& -\epsilon\sin 2\Phi\Big(\sqrt{2}\mathrm{Im}\{r^{2}_{10}\}\sin 2\Theta\sin\phi + \mathrm{Im}\{r^{2}_{1-1}\}\sin^2\Theta\sin 2\phi\Big) \\
& +\sqrt{2\epsilon(1+\epsilon)}\cos\Phi\Big(r^{5}_{11}\sin^2\Theta + r^{5}_{00}\cos^2\Theta - \sqrt{2}\mathrm{Re}\{r^{5}_{10}\}\sin 2\Theta\cos\phi - r^{5}_{1-1}\sin^2\Theta\cos 2\phi\Big) \\
& +\sqrt{2\epsilon(1+\epsilon)}\sin\Phi\Big(\sqrt{2}\mathrm{Im}\{r^{6}_{10}\}\sin 2\Theta\sin\phi + \mathrm{Im}\{r^{6}_{1-1}\}\sin^2\Theta\sin 2\phi\Big)\Big],
\end{aligned}$$

cos2ϕ

$$\begin{aligned}
\mathcal{W}^{L}(\Phi,\phi,\cos\Theta) = \frac{3}{8\pi^2}\Big[& \sqrt{1-\epsilon^2}\Big(\sqrt{2}\mathrm{Im}\{r^{3}_{10}\}\sin 2\Theta\sin\phi + \mathrm{Im}\{r^{3}_{1-1}\}\sin^2\Theta\sin 2\phi\Big) \\
& +\sqrt{2\epsilon(1-\epsilon)}\cos\Phi\Big(\sqrt{2}\mathrm{Im}\{r^{7}_{10}\}\sin 2\Theta\sin\phi + \mathrm{Im}\{r^{7}_{1-1}\}\sin^2\Theta\sin 2\phi\Big) \\
& +\sqrt{2\epsilon(1-\epsilon)}\sin\Phi\Big(r^{8}_{11}\sin^2\Theta + r^{8}_{00}\cos^2\Theta - \sqrt{2}\mathrm{Re}\{r^{8}_{10}\}\sin 2\Theta\cos\phi - r^{8}_{1-1}\sin^2\Theta\cos 2\phi\Big)\Big]
\end{aligned}$$

COMPASS Acceptance of ϕ for DVCS

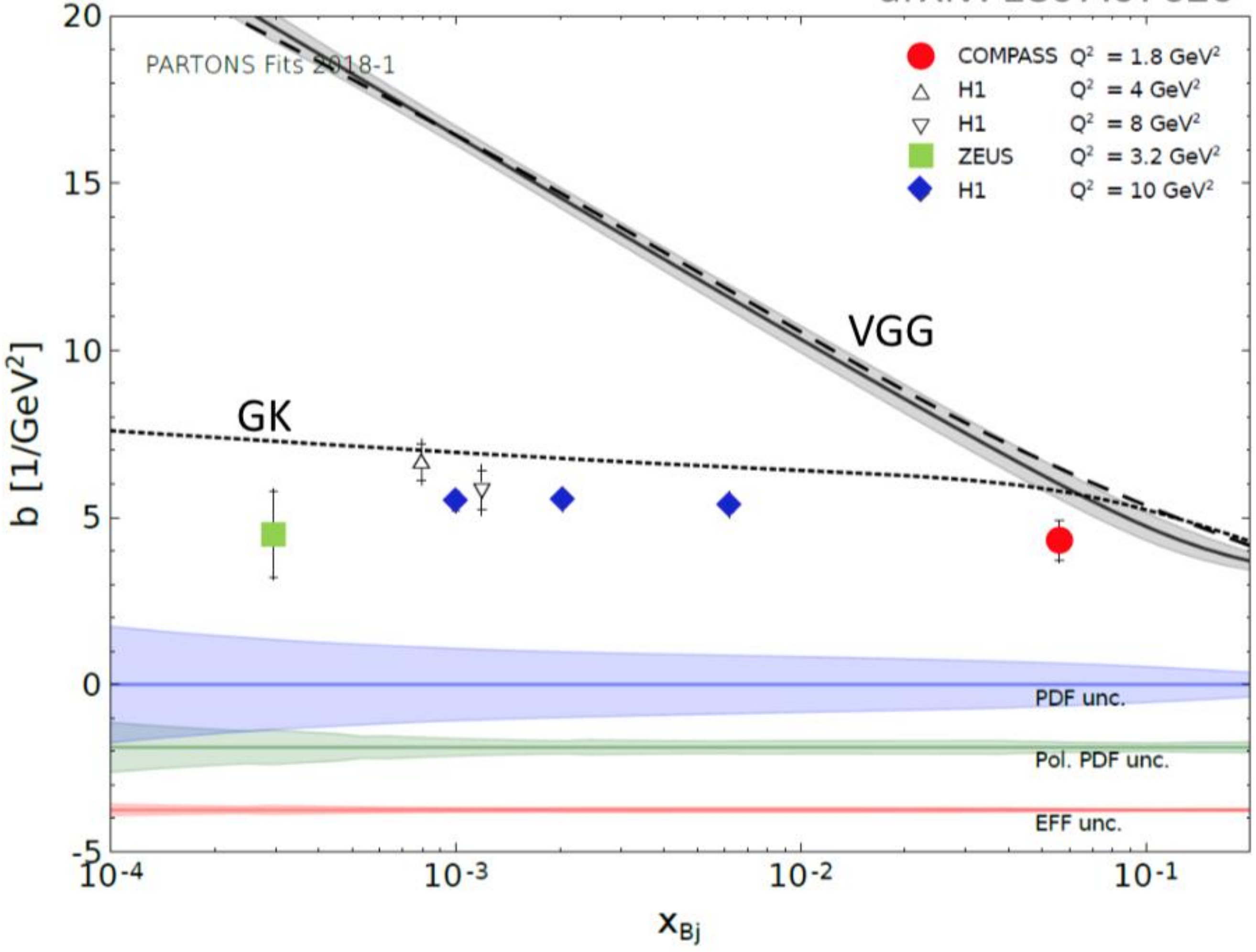


$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi \, (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \underbrace{\frac{1}{\Gamma(Q^2, \nu, E_\mu)}}_{\text{Flux of transverse virtual photons}} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

Results of 2012 Data

Figure from Moutarde, Sznajder, Wagner
arXiv: 1807.07620



- **Global Fit – MSW (grey band):** fit of CFF in the PARTON framework at LO and LT, using a GPD parametrization that only involves valence and sea quarks.
- **VGG Model :** with valence and sea quarks only.
- **GK Model** : includes also gluons.

Nice manifestation of gluons or NLO effect

Dominance of $Im\mathcal{H}$ (with respect of $Re\mathcal{H}$ and other $\mathcal{CFF}$) at small x_B

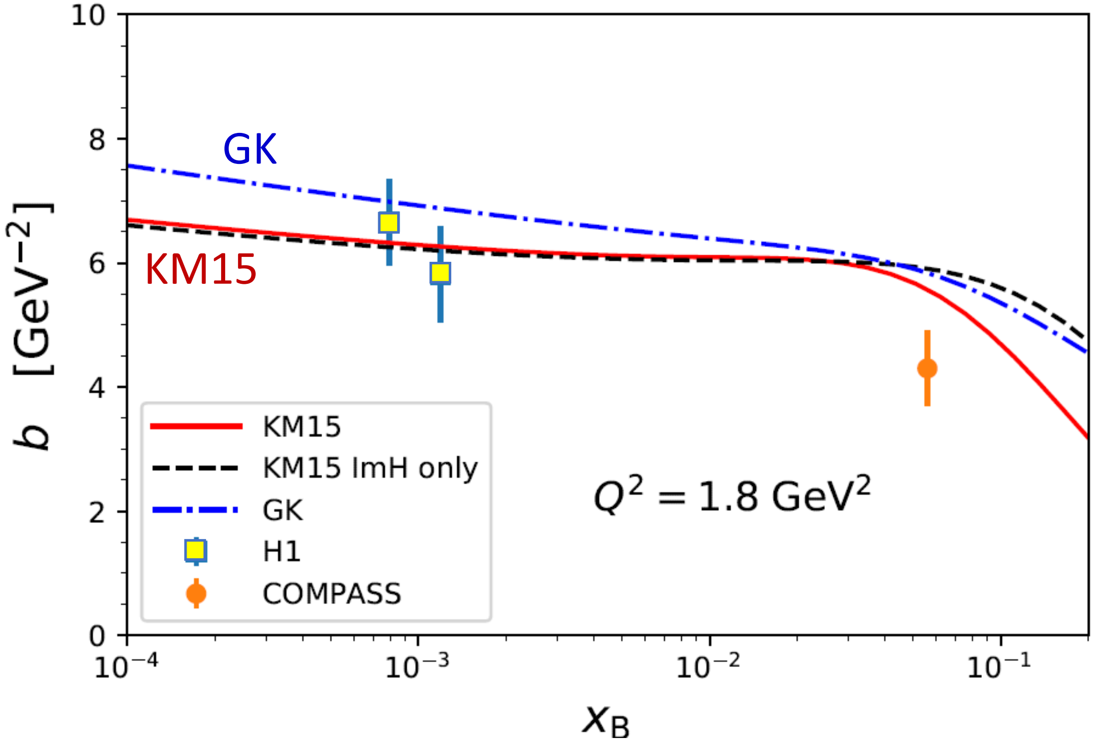


Figure from Moutarde, Sznajder, Wagner
arXiv: 1807.07620