

Partonic structure of neutral pseudoscalars via two photon transition form factors

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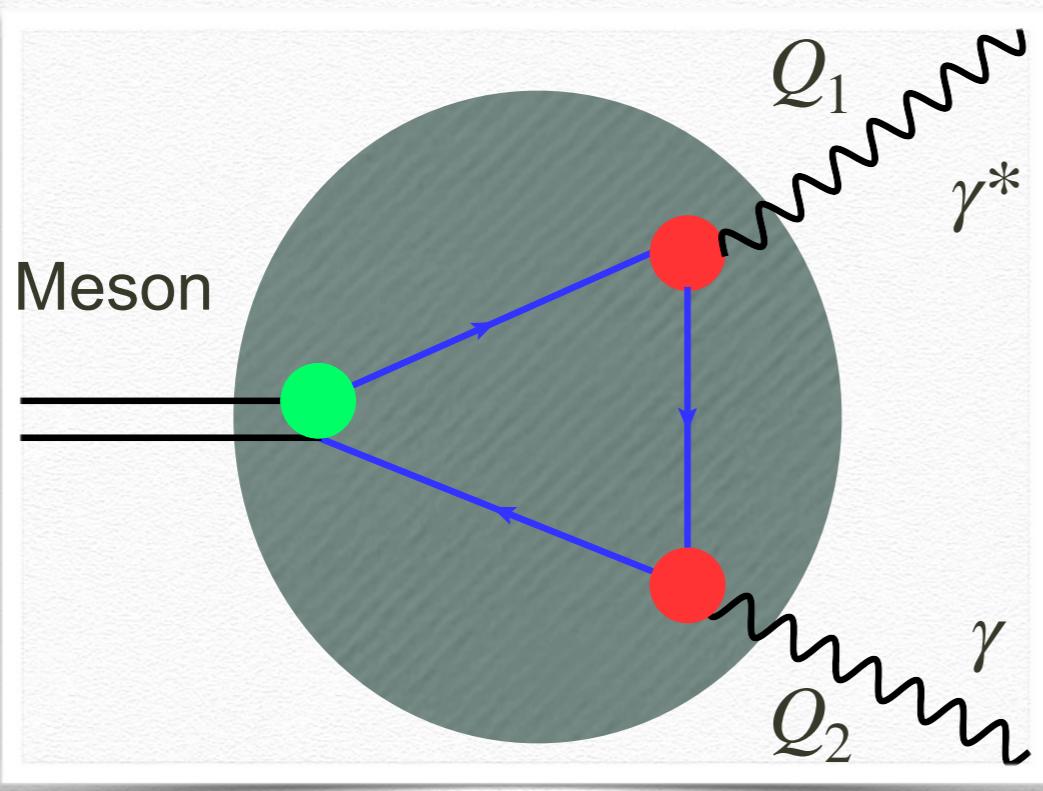
Nankai University, Tianjin, China & ECT*, Trento, Italy

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Two photon transition form factors

- ❖ Understanding the **structure of hadrons** in terms of QCD's **quarks and gluons** is one of the central goals of modern hadron physics.
- ❖ Form factors: the closest thing to a snapshot, the size, shape and makeup of hadrons.

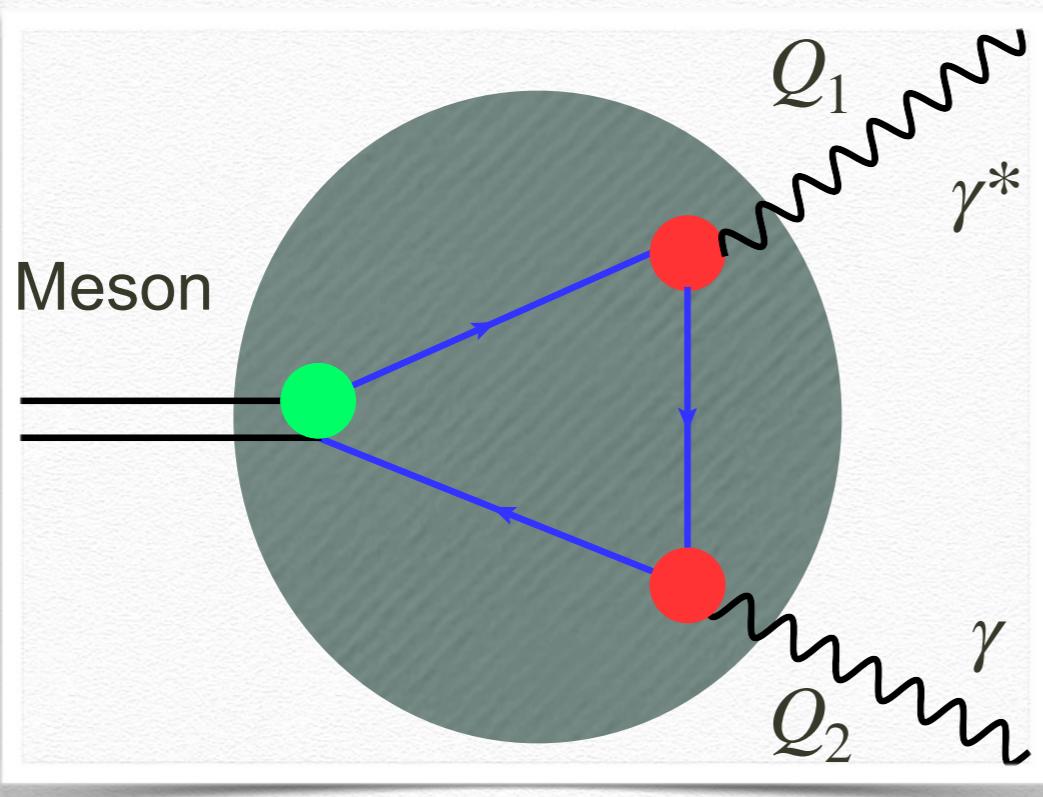


- ❖ Two photon transition form factors
- ❖ Transfer photon momentum Q_1, Q_2
 - $Q_1^2 \approx Q_2^2 \approx 0$, two photon decay width
 - $Q_1^2 \neq 0, Q_2^2 \approx 0$, single-virtual transition form factor (TFF)
 - $Q_1^2 \neq 0, Q_2^2 \neq 0$, double-virtual TFF

Single-virtual transition form factors: $Q_1^2 = Q^2 \neq 0, Q_2^2 \approx 0$

- ❖ Strong-QCD prediction: $Q^2 F_M^q(Q^2) \stackrel{Q^2 > Q_0^2}{\approx} 4\pi^2 f_M^q e_q^2 \tilde{w}_M^q(Q^2)$
- ❖ TFF: $F_M^q(Q^2)$, $q\bar{q}$ component of M, power-law scaling.

PhysRevD.22.2157



- $Q_0 > \Lambda_{QCD}, e_q$, electric charge
- $f_M^q, q\bar{q}$ —component leptonic decay constant
- $\tilde{w}_M^q(Q^2) = \int_0^1 dx \frac{1}{x} \varphi_M^q(x, Q^2)$
- $\varphi_M^q(x, Q^2)$ dressed-valence q-parton contribution to the meson's **distribution amplitude (DA)**

Transition form factors (TFF) & Distribution amplitude (DA)

❖ Strong-QCD prediction: $Q^2 F_M^q(Q^2) \stackrel{Q^2 > Q_0^2}{\approx} 4\pi^2 f_M^q e_q^2 \tilde{w}_M^q(Q^2)$

In the limit $\tau^2 := \Lambda_{QCD}^2/Q^2 \simeq 0$

■ Asymptotic TFF

$$Q^2 F_M^q(Q^2) \stackrel{\tau \simeq 0}{\approx} 12\pi^2 f_M^q e_q^2$$

■ Asymptotic DA:

$$\varphi_M(x, Q^2) \stackrel{\tau \simeq 0}{\approx} \varphi_\infty = 6x(1-x)$$

Models

Large- Q^2 :

$$F_M^q(Q^2) = \frac{A}{1 + Q^2/(12\pi^2 f_M^q e_q^2)}$$

Low- Q^2 :

$$F_M^q(Q^2) = \frac{B}{1 + Q^2/\Lambda_M^2}$$

Phys. Lett. B 87, 359 (1979)

Phys. Lett. B 94, 245 (1980)

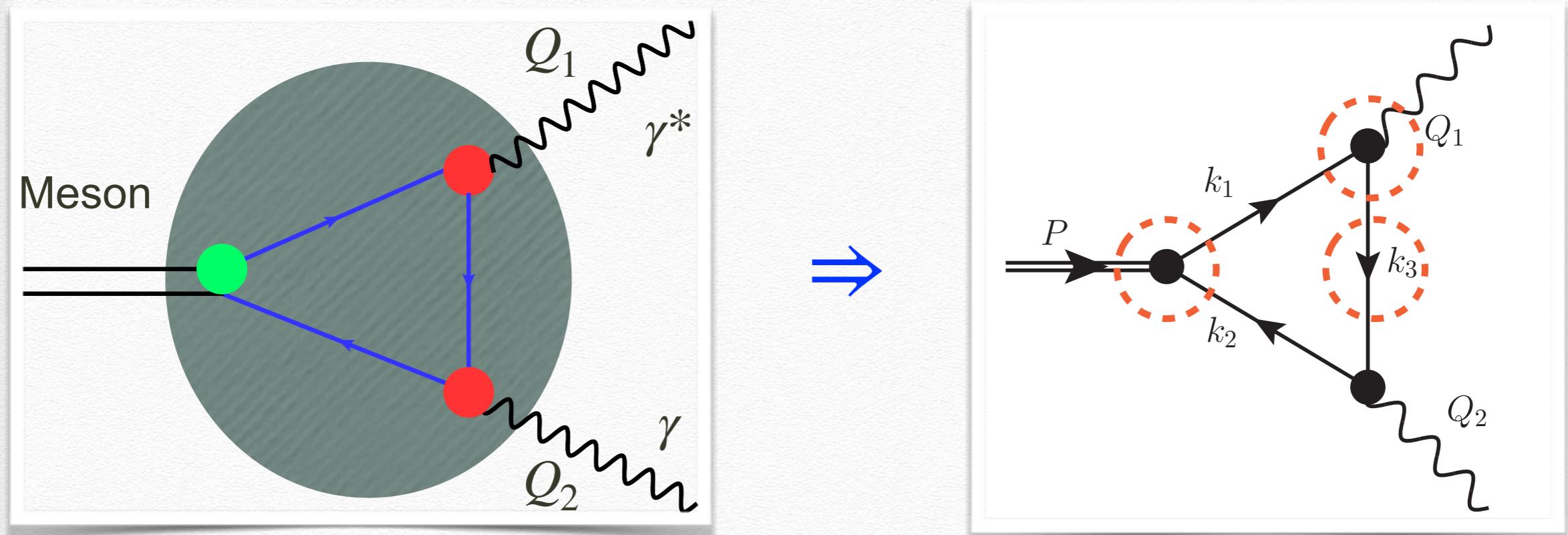
Single-virtual transition form factors with DSEs: formulae I

❖ Impulse approximation

$$\Lambda_{\alpha\beta}^q(Q_1^2, Q_2^2) = e^2 N_c \int_{dk}^\Lambda \text{tr} \left[S(k_1) \Gamma_M^{q\bar{q}}(k; P) S(k_2) i\Gamma_\beta(k + Q_1/2; Q_2) S(k_3) i\Gamma_\alpha(k - Q_2/2; Q_1) \right]$$

arXiv:nucl-th/0201017

$$❖ \text{TFF and vertex } \Lambda_{\mu\nu}^{\gamma^*\gamma^*\pi^0} = \frac{2i\alpha_{em}g_{\gamma\gamma\pi}}{\pi\tilde{f}_\pi} \epsilon_{\mu\nu\rho\sigma} Q_{1\rho} Q_{2\sigma} F_{\gamma^*\gamma^*\pi^0}(Q_1^2, Q_2^2)$$



Single-virtual transition form factors with DSEs: formulae II

❖ Impulse approximation

$$\Lambda_{\alpha\beta}^q(Q_1^2, Q_2^2) = e^2 N_c \int_{dk}^\Lambda \text{tr} \left[S(k_1) \boxed{\Gamma_M^{q\bar{q}}(k; P)} S(k_2) i \Gamma_\beta(k + Q_1/2; Q_2) S(k_3) i \Gamma_\alpha(k - Q_2/2; Q_1) \right]$$

- $S(k_i)$: dressed quark propagator

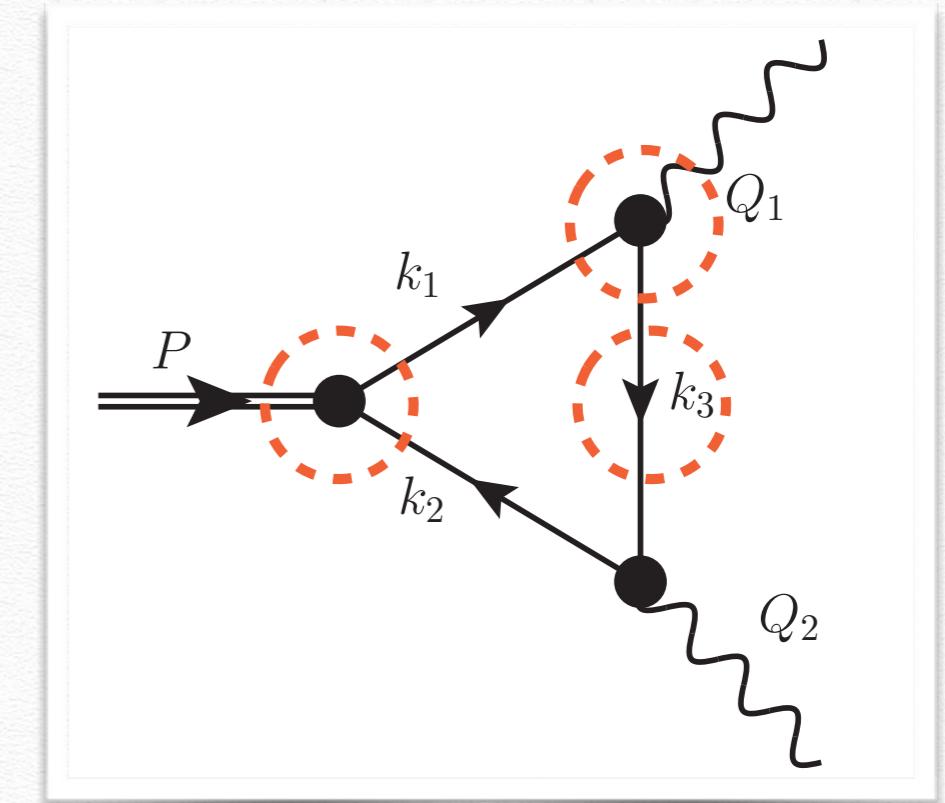
$$k_1 = k + P/2, k_2 = k - P/2$$

$$k_3 = k + (Q_1 - Q_2)/2$$

- $\Gamma_M^{q\bar{q}}$: meson Bethe-Salpeter amplitude

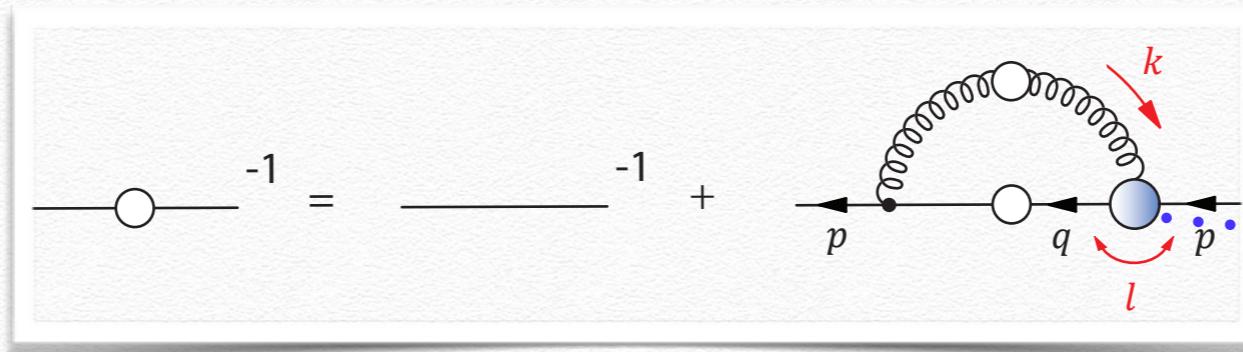
- $\Gamma_{\alpha/\beta}$: quark-photon vertex

$$\Lambda_{\alpha\beta}^{\pi^0 \rightarrow \gamma\gamma^*}(Q_1; Q_2) = \frac{2}{\sqrt{2}} \left((\hat{Q}^u)^2 - (\hat{Q}^d)^2 \right) \Lambda_{\alpha\beta}^u(Q_1; Q_2)$$



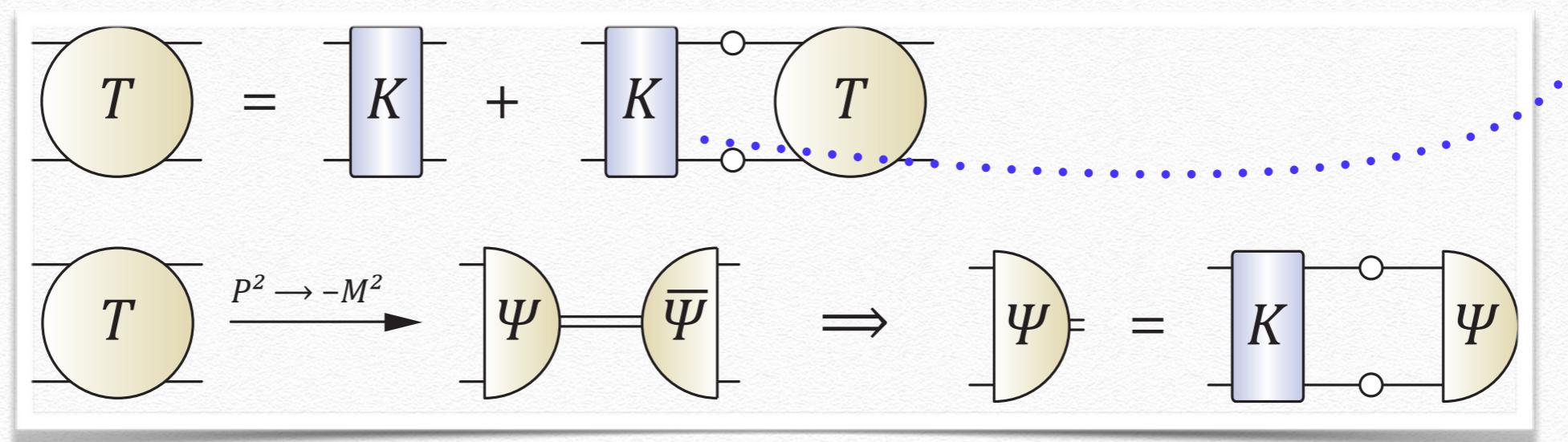
DSEs: Equations for mesons

- $S(k_i)$: dressed quark propagator



quark gluon vertex

- $\Gamma_M^{q\bar{q}}$: meson Bethe-Salpeter amplitude



scattering kernel

- $\Gamma_{\alpha/\beta}$: quark-photon vertex, inhomogeneous vector Bethe-Salpeter equation

picture from G. Eichmann, arXiv:0909.0703

DSEs: Axial-vector Ward-Takahashi identity

$$P_\mu \boxed{\Gamma_{5\mu}^a(k; P)} = \boxed{S^{-1}(k_+) i\gamma_5 \mathcal{F}^a + i\gamma_5 \mathcal{F}^a S^{-1}(k_-)} \\ - 2i\mathcal{M}^{ab} \boxed{\Gamma_5^b(k; P)} - \boxed{A^a(k; P)}$$

- ❖ An identity between correlation functions, symmetry conserved.
- ❖ Non-Abelian anomaly: $A^a(k; P)$.
- ❖ Neutral Pseudoscalars: quark gluon vertex \Leftrightarrow scattering kernel
- π^0, η_c, η_b , anomaly $A^a(k; P) = 0$. Rainbow-Ladder approximation.
- η, η' , anomaly $A^a(k; P) \neq 0$. Phenomenally introduce $A^a(k; P)$.

DSEs: Axial-vector Ward-Takahashi identity. π^0, η_c, η_b

❖ Anomaly $A^a(k; P) = 0$

$$P_\mu \Gamma_{5\mu}^a(k; P) = S^{-1}(k_+) i\gamma_5 \mathcal{F}^a + i\gamma_5 \mathcal{F}^a S^{-1}(k_-) - 2i\mathcal{M}^{ab} \Gamma_5^b(k; P)$$

❖ DSEs

$$\Gamma_{5\mu}(p, P) = \gamma_5 \gamma_\mu + \int_k K(p, k, P) S^a(k_+) \Gamma_{5\mu}(k, P) S^b(k_-)$$

$$S^{-1}(p) = i\gamma \cdot p + m + \Sigma(p) \quad \Gamma_5(p, P) = \gamma_5 + \int_k K(p, k, P) S^a(k_+) \Gamma_5(k, P) S^b(k_-)$$

❖ Self energy & scattering kernel

$$\Sigma(p_+) i\gamma_5 + i\gamma_5 \Sigma(p_-) = - \int_k K(p, k, P) (S^a(k_+) i\gamma_5 + i\gamma_5 S^a(k_-))$$

■ Rainbow

$$\Sigma(p) = \int_{dk} g^2 D_{\mu\nu}(p-k) \gamma_\mu S(k) \gamma_\nu$$

■ Ladder

$$K_{RL}(p, k, P) = -g^2 D_{\mu\nu}(p-k) \frac{\lambda^a}{2} \frac{\lambda^a}{2} \gamma_\mu \gamma_\nu$$

quark gluon vertex

↔ scattering kernel

DSEs: Axial-vector Ward-Takahashi identity. η, η'

- ❖ Anomaly $A^a(k; P) \neq 0$

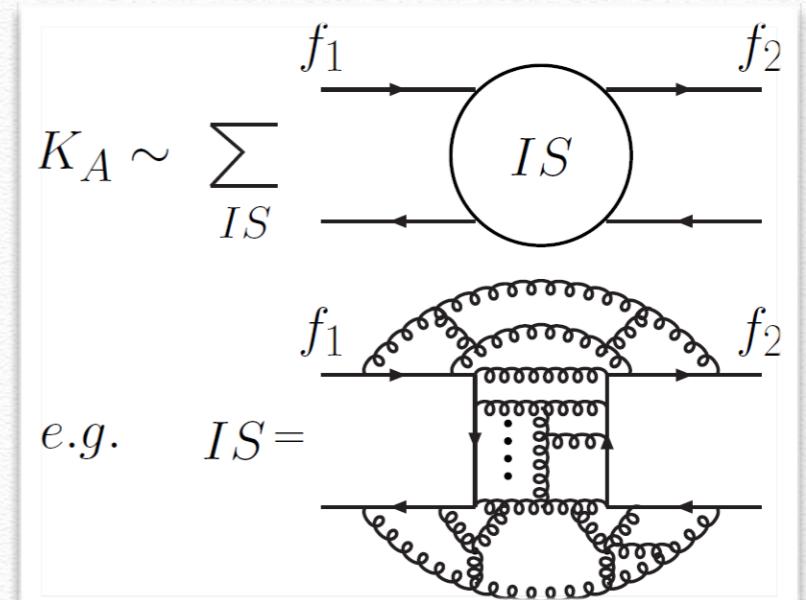
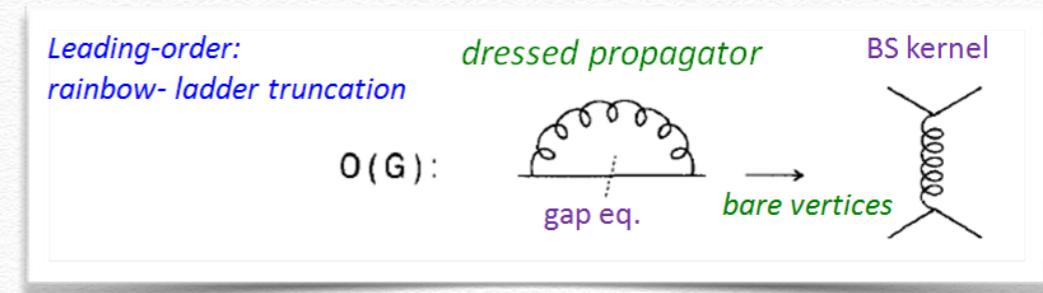
$$P_\mu \Gamma_{5\mu}^a(k; P) = S^{-1}(k_+) i\gamma_5 \mathcal{F}^a + i\gamma_5 \mathcal{F}^a S^{-1}(k_-) - 2i\mathcal{M}^{ab} \Gamma_5^b(k; P) - A^a(k; P)$$

- ❖ No simple relation between self energy and scattering kernel

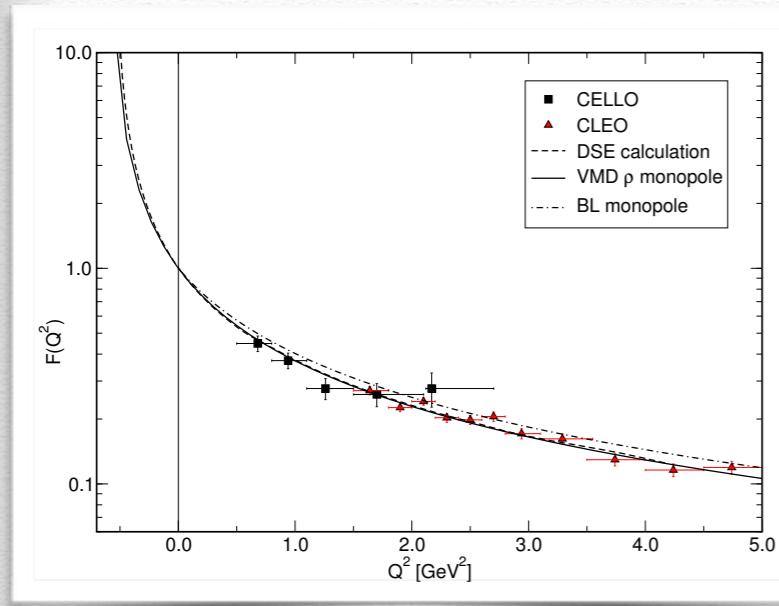
- ❖ "hairpin" structure: intermediate state (IS) involve infinitely many lines

$$\begin{aligned} & [K_A(k, q, P)]_{\rho\sigma}^{\alpha\beta} \\ &= \xi(s) \cos^2(\theta) \xi [zi\gamma_5]_{\alpha\beta} [zi\gamma_5]_{\rho\sigma} \\ &+ \frac{1}{\chi^2} \xi(s) \sin^2(\theta) \xi [zi\gamma_5 \gamma \cdot P]_{\alpha\beta} [zi\gamma_5 \gamma \cdot P]_{\rho\sigma} \end{aligned}$$

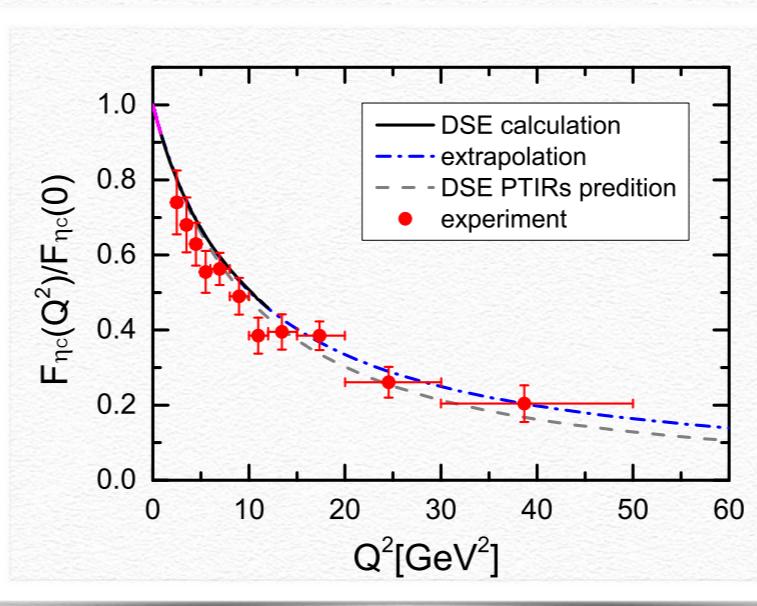
■ $\xi(s) = 0 \rightarrow$ anomaly kernel: $K_A = 0$



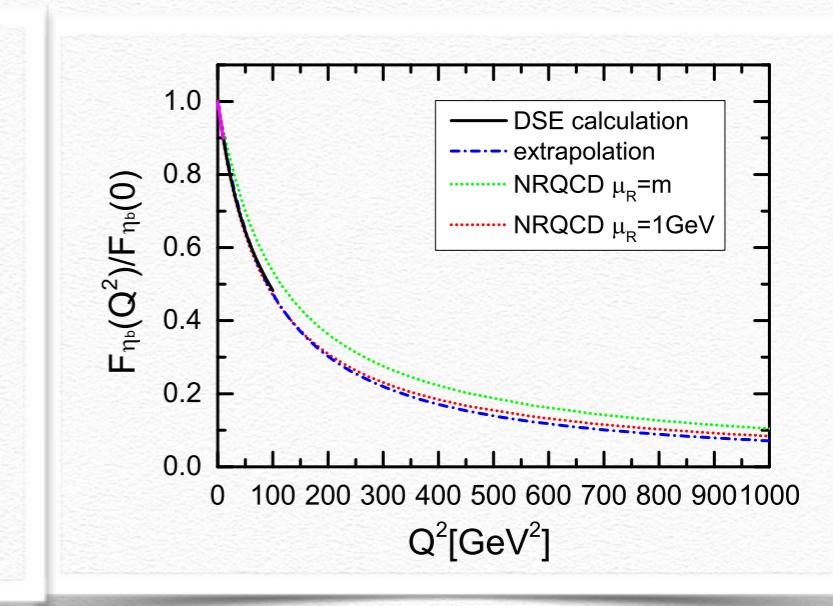
Single-virtual TFFs with DSEs: direct calculation



arXiv:nucl-th/0201017



arXiv:1611.05960



❖ Singularity of quark propagator

- F_{π^0} upper-limit $5 GeV^2$ ↵ BELLE
 $Q^2 \in [4,40] GeV^2$
- F_{η_c} upper-limit $12 GeV^2$ ↵ BABAR
 $Q^2 \in [2,50] GeV^2$
- F_{η_b} upper-limit $100 GeV^2$
 ↵ No experiment data

❖ Need extrapolation

Single-virtual TFFs with DSEs: extrapolation methods

- ❖ Method based on Distribution amplitude (DA)
- ❖ Off-shell method
 - Introducing a virtuality eigenvalue $\lambda(v)$ into the Bethe-Salpeter equations
 - Considering the v -dependence of the pointwise behaviour of the Bethe-Salpeter amplitude

arXiv:1702.06100

Single-virtual TFFs with DSEs: extrapolation methods

- ❖ Method based on Distribution amplitude (DA)

$$\begin{aligned}\langle 0 | \psi(-z) \gamma_5 \gamma \cdot n \psi(z) | \pi(P) \rangle &= f_\pi n \cdot P \int_0^1 dx e^{-i(2x-1)z \cdot P} \phi(x), \\ &= tr_{CD} Z_2 \int_{dq}^\Lambda e^{-iz \cdot q - iz \cdot (q-P)} \gamma_5 \gamma \cdot n \chi(q; P).\end{aligned}$$

- ❖ Projecting Bethe-Salpeter wave function onto the light front

$$f_\pi \phi(x) = tr_{CD} Z_2 \int_{dq}^\Lambda \delta(n \cdot q_+ - xn \cdot P) \gamma_5 \gamma \cdot n \chi(q; P)$$

- n , light-like four-vector, $n^2 = 0$. f_π , decay constant.
- $\chi(q; P)$ Bethe-Salpeter wave function

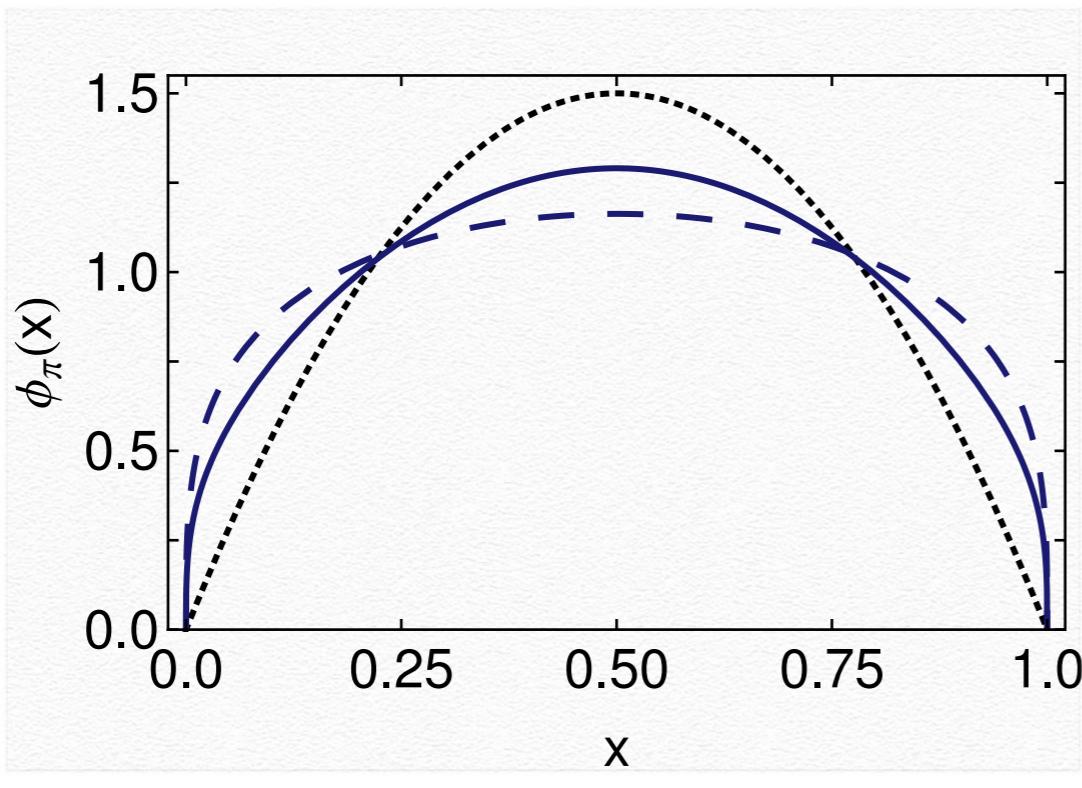
$$\chi(q, P) = S(q_+) \Gamma(q, P) S(q_-)$$

Single-virtual TFFs with DSEs: extrapolation methods

- ❖ Method based on Distribution amplitude (DA)
 - Perturbation theory integral representations (PTIRs):
 - Represent the Quark propagator with conjugate poles
 - Represent Bethe-Sapeter wave function with Nakanishi representation
 - Combine denominators of the integral by Feynman parameters

Prog. Theor. Phys. Suppl. 43, 1 (1969)

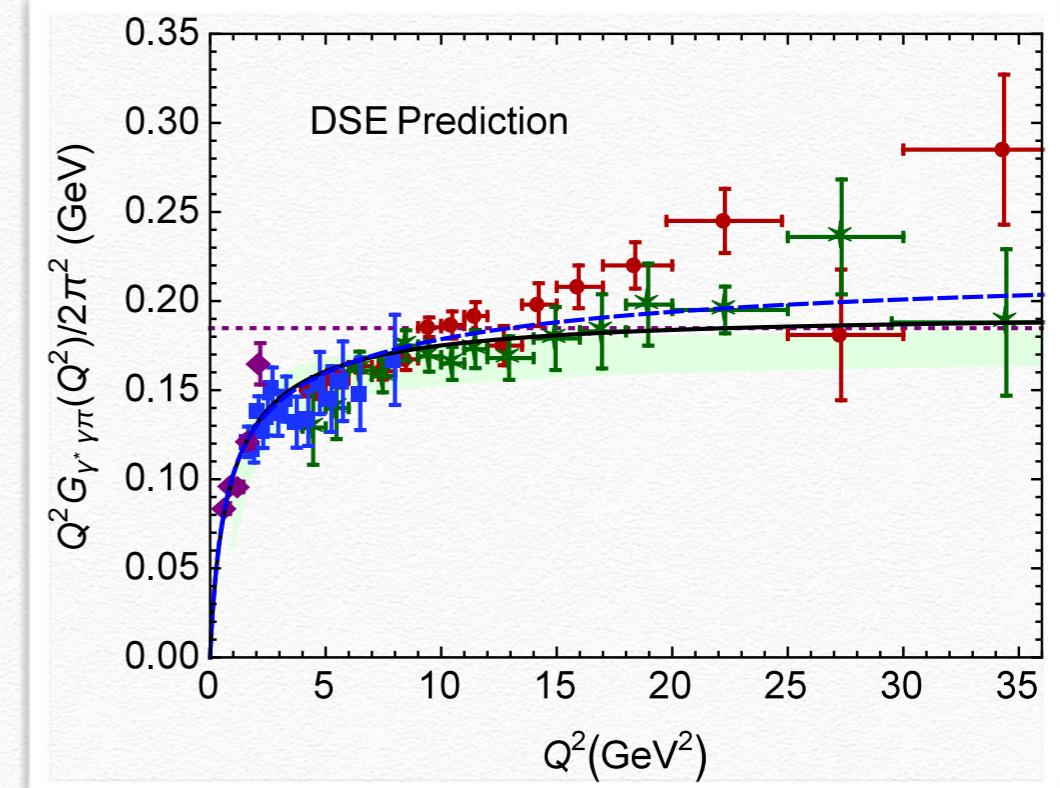
Single-virtual TFFs with DSEs: π^0



DA

arXiv:1301.0324

- Solid line: broad, concave function at hadronic scale;
- Dotted line: asymptotic prediction:
 $\phi_{asy}(x) = 6x(1 - x)$



TFF

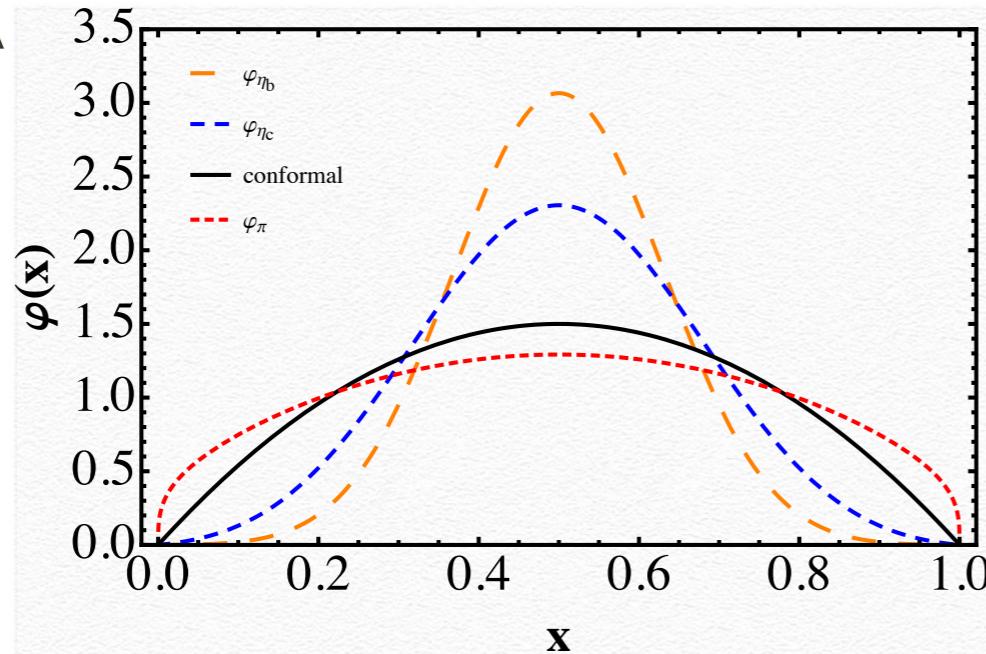
arXiv:1510.02799

- Consistent with all non-BaBar data
- Approach pQCD limit $2f_\pi$ from above

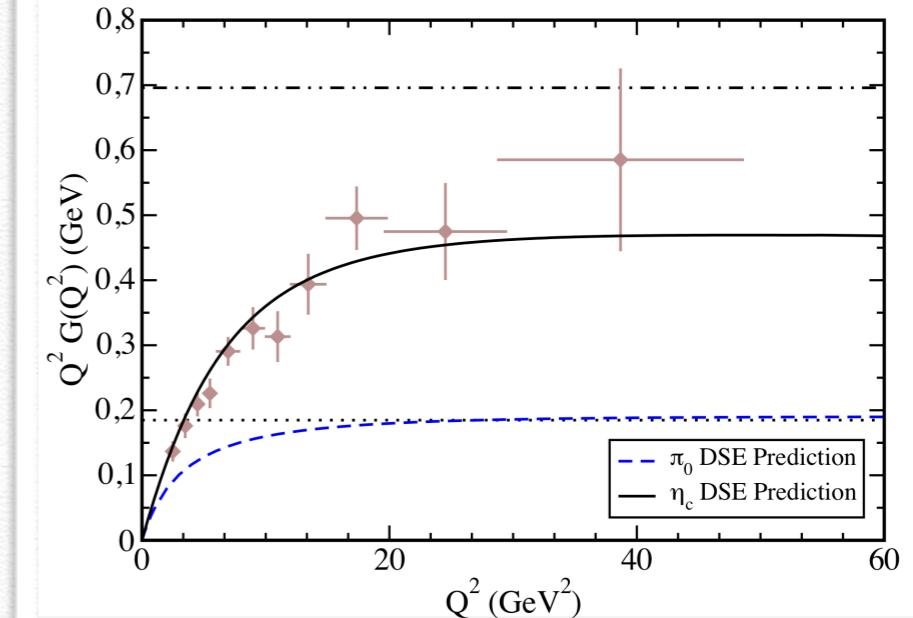
$$Q^2 F_\pi(Q^2) \stackrel{Q^2 > Q_0^2}{=} 4\pi^2 f_\pi \frac{1}{3} \int_0^1 dx \frac{\phi_\pi(x; Q^2)}{x}$$

Single-virtual TFFs with DSEs: η_c & η_b

DA



TFF



arXiv:1511.04943

- DA: piecewise convex-concave-convex function

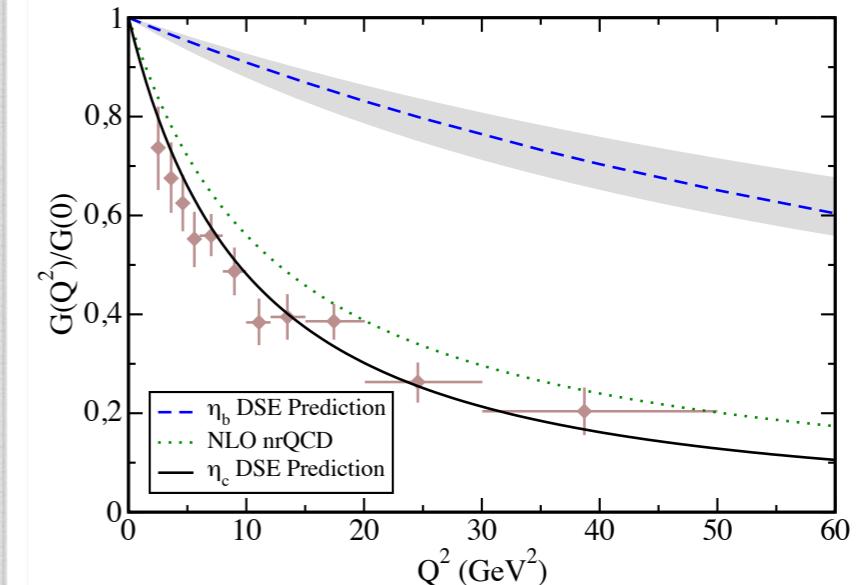
$$\Lambda_{QCD}/m_q(\zeta) \rightarrow 0, \phi(x) \rightarrow \delta(x - 1/2)$$

- TFF: consistent with BaBar data

- Approach pQCD limit $\frac{8}{3} f_{\eta_c}$ from below

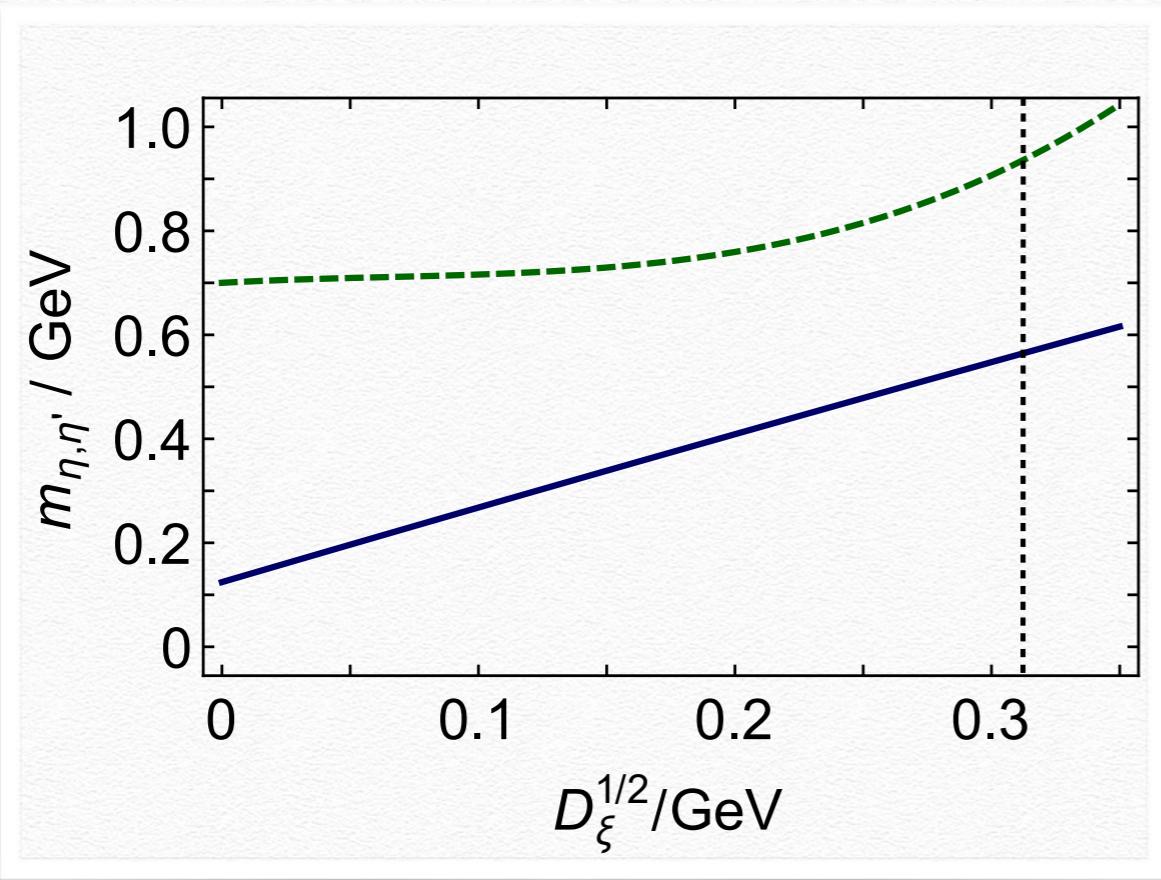
$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta_c}(Q^2) = 4\pi^2 \int_0^1 dx \frac{\frac{4}{9} f_{\eta_c} \phi_{\eta_c}(x)}{1-x}$$

arXiv:1610.06575



Single-virtual TFFs with DSEs: η & η'

◆ Meson Mass



- Ideal mixing: no anomaly
 $m_\eta = m_\pi, m_{\eta'} = m_{s\bar{s}} = 0.7 \text{GeV}$
- Meson mass grow with the anomaly strength, reach the empirical values
 $m_\eta = 0.56 \text{GeV}, m_{\eta'} = 0.96 \text{GeV}$

Single-virtual TFFs with DSEs: η & η'

- ❖ Decay constants: pseudovector projection of the Bethe-Salpeter wave function

$$f_{\eta,\eta'}^{l,s} P_\mu = Z_2 \text{tr} \int_{dk}^{\Lambda} \gamma_5 \gamma_\mu \chi_{\eta,\eta'}^{l,s}(k; P)$$

	f_η^l	f_η^s	$f_{\eta'}^l$	$f_{\eta'}^s$
herein-direct	0.072	-0.092	0.070	0.104
herein-fit	0.074	-0.094	0.068	0.101
phen. ^{1,2}	0.090(13)	-0.093(28)	0.073(14)	0.094(8)

- ❖ Flavor mixing angle

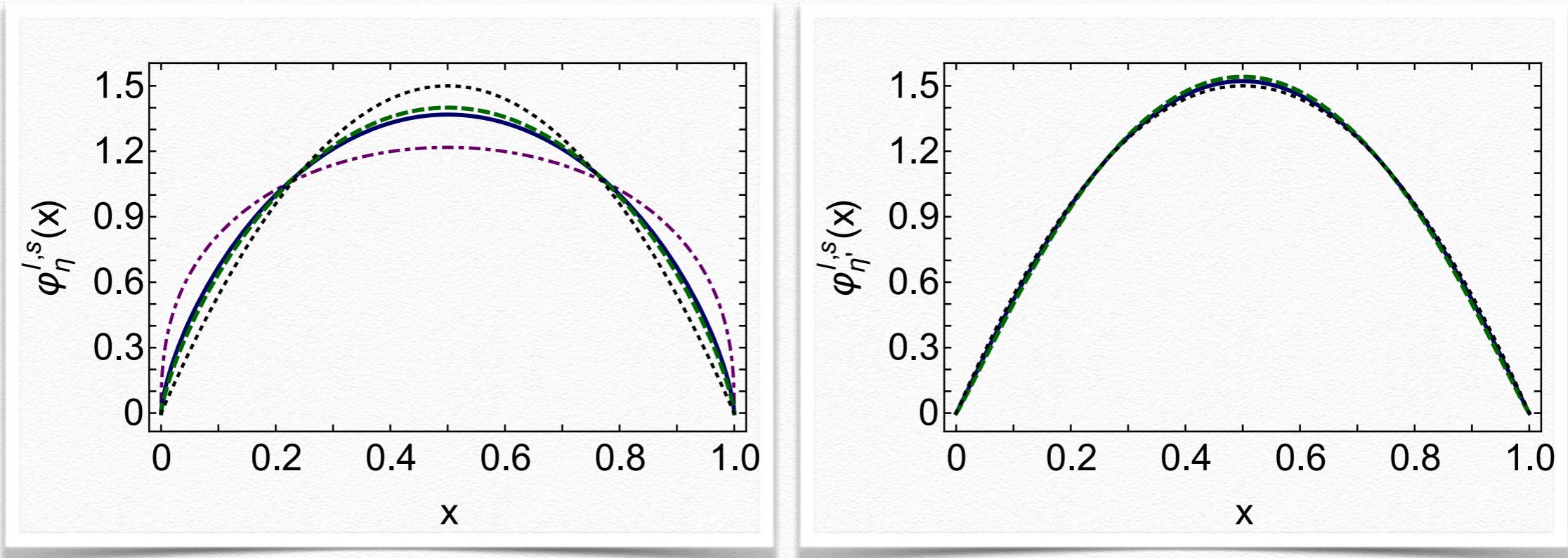
$$\begin{pmatrix} f_\eta^l & f_\eta^s \\ f_{\eta'}^l & f_{\eta'}^s \end{pmatrix} = \begin{pmatrix} f^l \cos \theta & -f^s \sin \theta \\ f^l \sin \theta & f^s \cos \theta \end{pmatrix}$$

- $\theta = 42.8^\circ$
- $f^l = 0.101 \text{GeV} = 1.08 f_\pi$
- $f^s = 0.138 \text{GeV} = 1.49 f_\pi$

$$\begin{pmatrix} f_\eta^8 & f_\eta^0 \\ f_{\eta'}^8 & f_{\eta'}^0 \end{pmatrix} = \begin{pmatrix} f^8 \cos \theta_8 & -f^0 \sin \theta_0 \\ f^8 \sin \theta_8 & f^0 \cos \theta_0 \end{pmatrix}$$

- $\theta_8 = -21^\circ, \theta_0 = -2.8^\circ$
- $f^8 = 1.34 f_\pi$
- $f^0 = 1.26 f_\pi$

Single-virtual TFFs with DSEs: η & η'



❖ DAs: ϕ_η
 $\phi^{asy} <_N \phi_\eta <_N \phi_\pi$

❖ DAs: $\phi_{\eta'}$
 $\phi_{\eta'} \sim \phi^{asy}$

❖ Light- and s-quark component DAs have similar profiles

Single-virtual TFFs with DSEs: η & η' Low Q^2

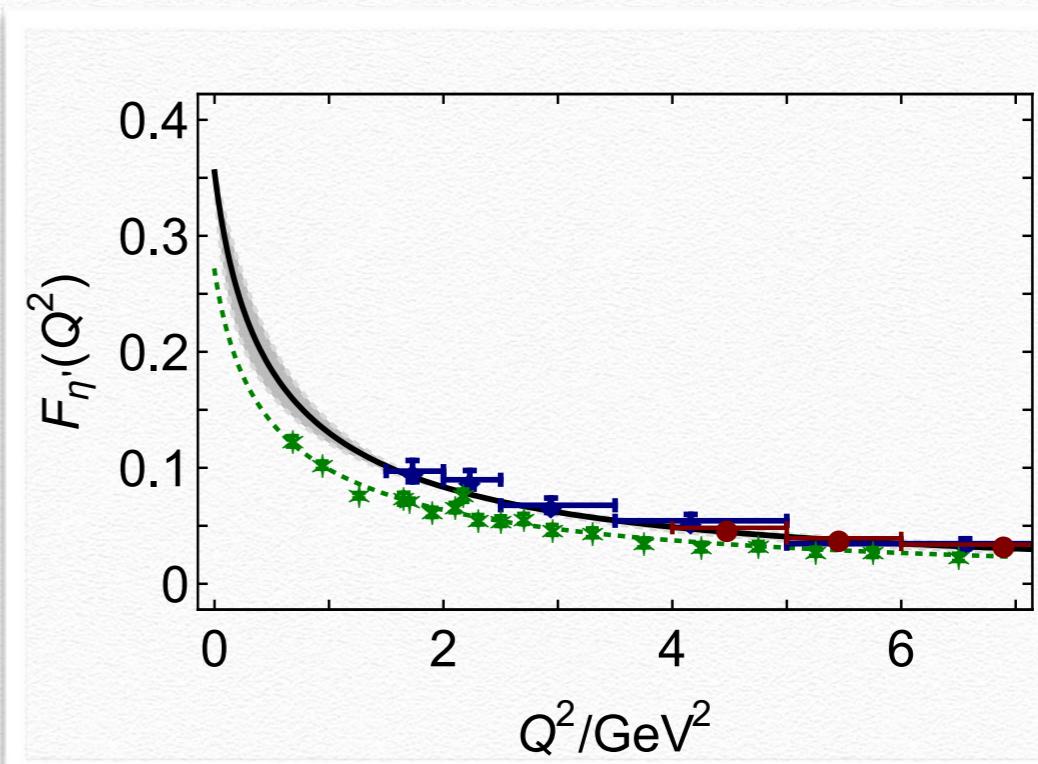
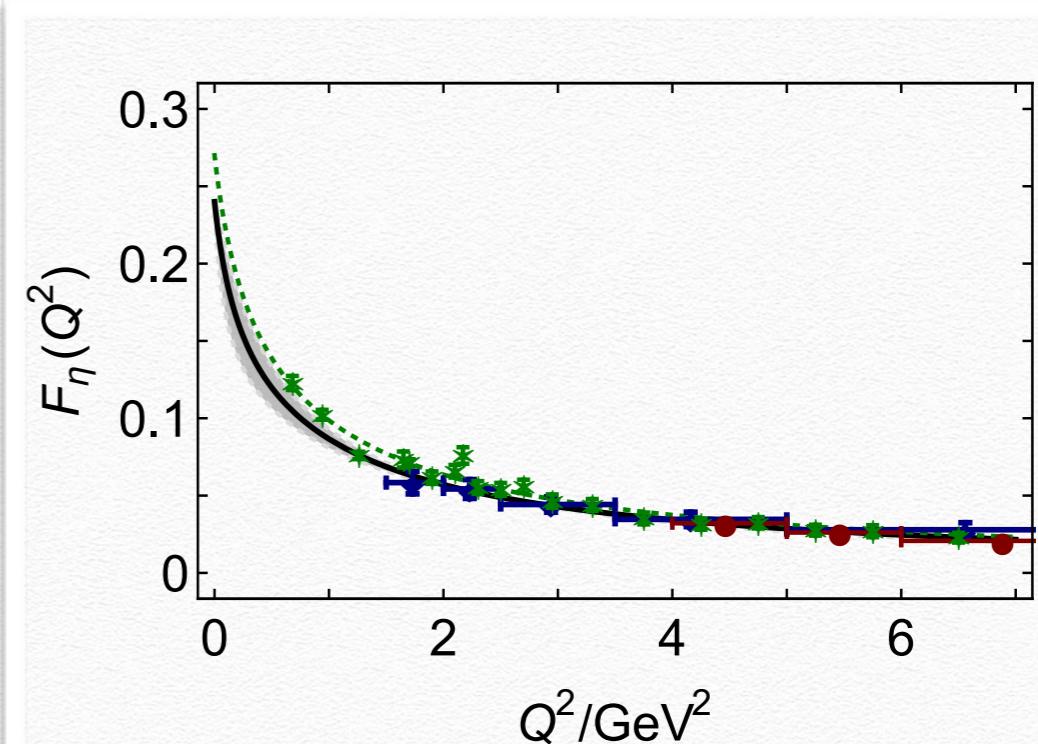
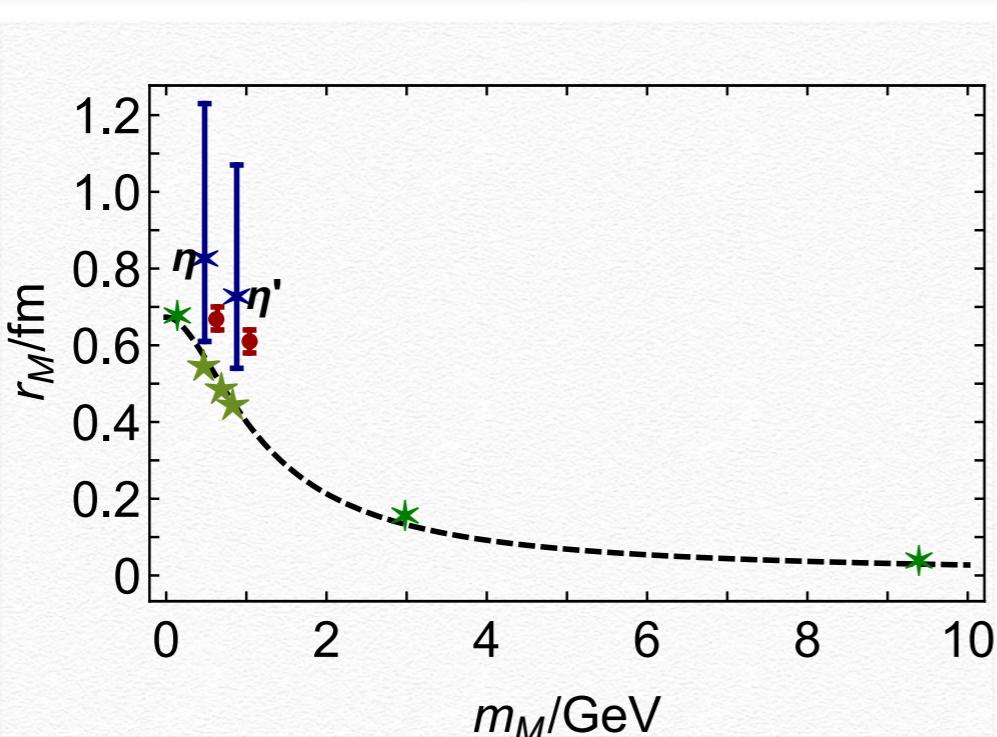
- ❖ TFFs consistent with Babar data

- ❖ Radii:

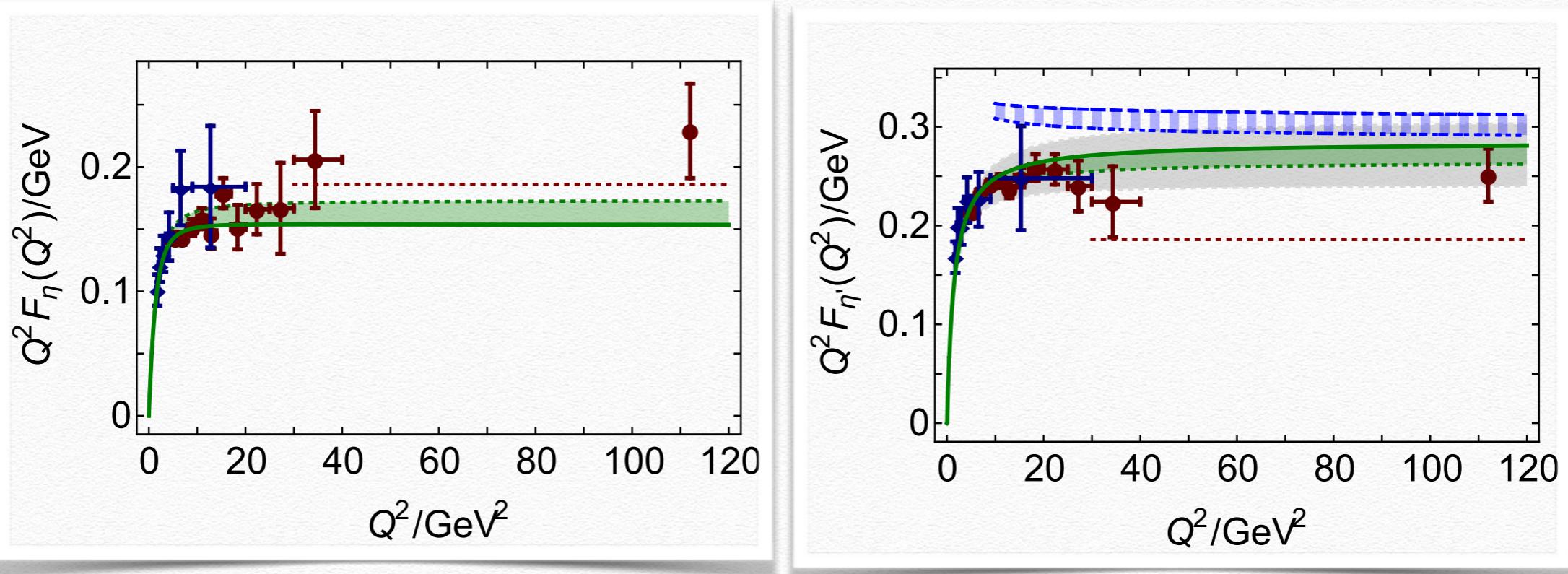
$$r_\eta = 0.83^{+0.40}_{-0.22} \text{ fm}, r_{\eta'} = 0.73^{+0.34}_{-0.19} \text{ fm}$$

$$r_{\eta,\eta'}^2 := \frac{-6}{F_{\eta,\eta'}(0)} \frac{d}{dQ^2} F_{\eta,\eta'}(Q^2) \Big|_{Q^2=0}$$

- ❖ $\pi^0, \eta, \eta', \eta_c, \eta_b$ & $m_0^-/\text{GeV} = 0.47, 0.69, 0.83$
- ❖ Non-Abelian anomaly: η, η' larger than a fit by 24% and 48%



Single-virtual TFFs with DSEs: η & η' Large Q^2



■ $Q^2 F_\eta(Q^2 \rightarrow \infty) = 0.15 \text{GeV}$

■ $Q^2 F_{\eta'}(Q^2 \rightarrow \infty) = 0.30 \text{GeV}$

■ Omission of $\varphi_{\eta,\eta'}^0 - \varphi_{\eta,\eta'}^g$ mixing

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta,\eta'}(Q^2) = 6 \left[\frac{5}{9} f_{\eta,\eta'}^l(Q^2) + \frac{\sqrt{2}}{9} f_{\eta,\eta'}^s(Q^2) \right]$$

η & η' topological charge

$$P_\mu \Gamma_{5\mu}^a(k; P) = S^{-1}(k_+) i\gamma_5 \mathcal{F}^a + i\gamma_5 \mathcal{F}^a S^{-1}(k_-)$$

$$-2i\mathcal{M}^{ab} \Gamma_5^b(k; P) - A^a(k; P)$$

↓

$$m_{\eta,\eta'}^2 f_{\eta,\eta'}^a = \delta^{a0} n_{\eta,\eta'} + 2\mathcal{M}^{ab} \rho_{\eta,\eta'}^b$$

❖ Axial-vector Ward-Takahashi identity

❖ General mass formula

■ Topological charge: $n_{\eta,\eta'}$

$$n_{\eta,\eta'} = \sqrt{\frac{3}{2}} \nu_{\eta,\eta'}, \nu_{\eta,\eta'} = \langle 0 | \mathcal{Q}(x) = i \frac{\alpha_s}{4\pi} \text{tr}_C [\epsilon_{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma(x)}] | \eta, \eta' \rangle$$

■ Pseudoscalar projection of the Bethe-Salpeter wave function: $\rho_{\eta,\eta'}^b$

$$\begin{array}{cccc} \rho_\eta^8 & \rho_\eta^0 & \rho_{\eta'}^8 & \rho_{\eta'}^0 \\ 0.50^2 & 0.033^2 & -0.37^2 & 0.57^2 \end{array} \quad \Rightarrow \quad \begin{aligned} \nu_\eta &= (0.29 GeV)^3 \\ \nu_{\eta'} &= (0.37 GeV)^3. \end{aligned}$$

■ Topological charge content of the η' is 2.1 times that of η

Summary

- ❖ Two photon **single-virtual transition form factors** of neutral pseudoscalars using DSEs, constrain Bethe-Salpeter kernel by **axial-vector Ward-Takahashi identity**.
- ❖ $\gamma\gamma^* \rightarrow \pi^0$: consistent with all non-BaBar data; approach pQCD limit $2f_\pi$ from above; **DA: broad concave function**.
- ❖ $\gamma\gamma^* \rightarrow \eta_c, \eta_b$: consistent with BaBar data; approach pQCD limit $\frac{8}{3}f_{\eta_c}$ from below; **DA: narrow piecewise convex-concave-convex function**; $\phi_{\eta_b} <_N \phi_{\eta_c} <_N \phi^{asy}$
- ❖ $\gamma\gamma^* \rightarrow \eta, \eta'$: consistent with BaBar data; **DA: $\phi^{asy} <_N \phi_\eta <_N \phi_\pi, \phi_{\eta'} \sim \phi^{asy}$** ; **non-abelian anomaly**: η, η' larger than a non-mixing pseudoscalar radii fit by 24% and 48%; $Q^2 F_\eta(Q^2 \rightarrow \infty) = 0.15 GeV, Q^2 F_{\eta'}(Q^2 \rightarrow \infty) = 0.30 GeV$; Topological charge content of the η' is 2.1 times that of η .

Outlook

- ❖ **Double virtual** transition form factors, $Q_1^2 \neq 0, Q_2^2 \neq 0$, BaBar, (2018)
- ❖ **Time like** form factor: $Q^2 < 0$.

Thank you!