## Freed-Isobar Analysis of Light Mesons at COMPASS



Fabian Krinner for the COMPASS collaboration

Max Planck Institut für Physik



XVIII Conference on Hadron Spectroscopy and Structure

Guilin, China

August 17th 2019

#### The COMPASS experiment common Muon Proton Apparatus for Structure and Spectroscopy





#### The COMPASS experiment common Muon Proton Apparatus for Structure and Spectroscopy





Fabian Krinner (MPP)























 COMPASS: Large data set for the diffractive process

 $\pi_{\text{beam}}^- + p \rightarrow \pi^- \pi^+ \pi^- + p$ 



- COMPASS: Large data set for the diffractive process  $\pi^{-}_{\text{beam}} + \mathbf{p} \rightarrow \pi^{-}\pi^{+}\pi^{-} + \mathbf{p}$
- Squared four-momentum transferred t' by Pomeron  $\mathbb{P}$



- COMPASS: Large data set for the diffractive process  $\pi^{-}_{\text{beam}} + p \rightarrow \pi^{-}\pi^{+}\pi^{-} + p$
- Squared four-momentum transferred t' by Pomeron  $\mathbb{P}$
- Exclusive measurement





- COMPASS: Large data set for the diffractive process  $\pi^{-}_{\text{beam}} + p \rightarrow \pi^{-}\pi^{+}\pi^{-} + p$
- Squared four-momentum transferred t' by Pomeron  $\mathbb{P}$
- Exclusive measurement
- $46 \times 10^6$  exclusive events





- COMPASS: Large data set for the diffractive process  $\pi^{-}_{\text{beam}} + p \rightarrow \pi^{-}\pi^{+}\pi^{-} + p$
- Squared four-momentum transferred t' by Pomeron  $\mathbb{P}$
- Exclusive measurement
- $46 \times 10^6$  exclusive events
- Rich structure in π<sup>-</sup>π<sup>+</sup>π<sup>-</sup> mass spectrum: Intermediate states X<sup>-</sup>







- COMPASS: Large data set for the diffractive process  $\pi_{beam}^- + p \rightarrow \pi^- \pi^+ \pi_{bachelor}^- + p$
- Squared four-momentum transferred t' by Pomeron  $\mathbb{P}$
- Exclusive measurement
- $46 \times 10^6$  exclusive events
- Rich structure in π<sup>-</sup>π<sup>+</sup>π<sup>-</sup> mass spectrum: Intermediate states X<sup>-</sup>
- Also structure in π<sup>+</sup>π<sup>-</sup> subsystem: Intermediate states ξ (isobar)



COMPASS collaboration, PR **D95** (2017) 032004



• Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state



- Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

## Modelling resonances



#### Example: dynamic isobar amplitude for $\rho(770)$





- Intermediate states: Dynamic amplitudes △ (m): (line shape)
   Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

• Analysis performed independently in bins of  $m_{\chi^-} = m_{3\pi}$ .



- Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

- Analysis performed independently in bins of  $m_{\chi^-} = m_{3\pi}$ .
  - Learn  $m_{3\pi}$  dependence of partial-wave amplitudes
  - Dynamic amplitude of X<sup>-</sup> inferred form the data



- Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

- Analysis performed independently in bins of  $m_{\chi^-} = m_{3\pi}$ .
  - Learn  $m_{3\pi}$  dependence of partial-wave amplitudes
  - Dynamic amplitude of X<sup>-</sup> inferred form the data
- Dynamic amplitude of  $\xi \to \pi^+\pi^-$ : Model input in Partial-Wave Analysis



- Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

- Analysis performed independently in bins of  $m_{\chi^-} = m_{3\pi}$ .
  - Learn  $m_{3\pi}$  dependence of partial-wave amplitudes
  - Dynamic amplitude of X<sup>-</sup> inferred form the data
- Dynamic amplitude of  $\xi \to \pi^+\pi^-$ : Model input in Partial-Wave Analysis
  - ► Physical dynamic isobar amplitudes may differ from the model



- Intermediate states: Dynamic amplitudes  $\Delta(m)$ : (line shape) Complex-valued functions of the invariant mass of the state
- Simplest example: Breit-Wigner amplitude with nominal mass m<sub>0</sub> and width Γ<sub>0</sub> of a resonance:

$$\Delta_{\rm BW}\left(m\right) = \frac{m_0\Gamma_0}{m_0^2 - m^2 - im_0\Gamma_0}$$

- Analysis performed independently in bins of  $m_{\chi^-} = m_{3\pi}$ .
  - Learn  $m_{3\pi}$  dependence of partial-wave amplitudes
  - Dynamic amplitude of X<sup>-</sup> inferred form the data
- Dynamic amplitude of  $\xi \to \pi^+\pi^-$ : Model input in Partial-Wave Analysis
  - ► Physical dynamic isobar amplitudes may differ from the model
  - Parameterizations neglect e.g. final-state interactions

$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{ extsf{waves}} \mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$



$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{\mathsf{waves}} \mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$



$$\mathcal{I}\left(ec{ au}
ight)=\left|\sum_{i}^{ extsf{waves}}\mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$
- Decay amplitudes  $A_i(\vec{\tau})$ :
  - ► Describe \(\tilde{\tau}\) distributions of single partial waves
  - Known functions



$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{\mathsf{waves}} \mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$
- Decay amplitudes  $A_i(\vec{\tau})$ :
  - ► Describe \(\tilde{\tau}\) distributions of single partial waves
  - Known functions



$$\mathcal{A}_{i}(\vec{\tau}) = \psi_{i}(\vec{\tau}) \Delta_{i}(m_{\pi^{-}\pi^{+}})$$





$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{\mathsf{waves}}\mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$
- Decay amplitudes  $A_i(\vec{\tau})$ :
  - ► Describe \(\tilde{\tau}\) distributions of single partial waves
  - Known functions

Factorize decay amplitudes:

 $\mathcal{A}_{i}\left(\vec{\tau}\right) = \psi_{i}\left(\vec{\tau}\right) \Delta_{i}\left(m_{\pi^{-}\pi^{+}}\right)$ 

 Angular amplitudes ψ<sub>i</sub> (τ): Determined by angular momentum and spin quantum numbers of the waves

Ap. Ag > it

 Expand intensity distribution *I*(*τ*) over phase-space variables *τ* as sum over partial waves:

$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{\mathsf{waves}} \mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$
- Decay amplitudes  $A_i(\vec{\tau})$ :
  - ► Describe tributions of single partial waves
  - Known functions

Factorize decay amplitudes:

 $\mathcal{A}_{i}\left(\vec{\tau}\right) = \psi_{i}\left(\vec{\tau}\right) \Delta_{i}\left(m_{\pi^{-}\pi^{+}}\right)$ 

- Angular amplitudes ψ<sub>i</sub> (τ
   <sup>i</sup>): Determined by angular momentum and spin quantum numbers of the waves
- Dynamic isobar amplitudes  $\Delta_i(m_{\pi^-\pi^+})$ : Model input
  - Example: ρ(770) with fixed mass m<sub>0</sub>, width Γ<sub>0</sub> and quantum numbers J<sup>PC</sup><sub>ξ</sub> = 1<sup>-−</sup>

Ap. Ag > it

Expand intensity distribution *I*(*τ*) over phase-space variables *τ* as sum over partial waves:

$$\mathcal{I}\left(ec{ au}
ight) = \left|\sum_{i}^{\mathsf{waves}} \mathcal{T}_{i}\mathcal{A}_{i}\left(ec{ au}
ight)
ight|^{2}$$

- Production amplitudes  $T_i$ :
  - Encode strengths and relative phases of the single partial waves i
  - Free parameters in the analysis
  - Independent of  $\vec{\tau}$
- Decay amplitudes  $A_i(\vec{\tau})$ :
  - ► Describe \(\tilde{\tau}\) distributions of single partial waves
  - Known functions

Factorize decay amplitudes:

 $\mathcal{A}_{i}\left(\vec{\tau}\right) = \psi_{i}\left(\vec{\tau}\right)\Delta_{i}\left(m_{\pi^{-}\pi^{+}}\right)$ 

- Angular amplitudes ψ<sub>i</sub> (τ
   <sup>i</sup>): Determined by angular momentum and spin quantum numbers of the waves
- Dynamic isobar amplitudes  $\Delta_i(m_{\pi^-\pi^+})$ : Model input
  - Example: ρ(770) with fixed mass m<sub>0</sub>, width Γ<sub>0</sub> and quantum numbers J<sup>PC</sup><sub>ξ</sub> = 1<sup>-−</sup>
  - Not given by first principles
  - Have to be known beforehand









•  $J_{X^{-}}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^{-}$ 





J<sup>PC</sup><sub>X−</sub>: Spin and eigenvalues under parity and charge conjugation of X<sup>-</sup>
 M<sup>ε</sup>: Spin projection on the beam and naturality of the exchange particle Natural: P = (-1)<sup>J</sup>





- $J_{X}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^-$ •  $M^{\varepsilon}$ : Spin projection on the beam and naturality of the exchange particle
- $\xi$ : Appearing isobar, e.g.  $f_0$ ,  $\rho(770)$ ,  $f_2(1270)$  with  $J_{\xi}^{PC}$





•  $J_{X^{-}}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^{-}$ 

- $M^{\varepsilon}$ : Spin projection on the beam and naturality of the exchange particle
- $\xi$ : Appearing isobar, e.g.  $f_0$ ,  $\rho(770)$ ,  $f_2(1270)$  with  $J_{\xi}^{PC}$
- $\pi$ : Indicating the bachelor  $\pi^-$ . Always the same





- $J_{X^{-}}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^{-}$
- $M^{\varepsilon}$ : Spin projection on the beam and naturality of the exchange particle
- $\xi$ : Appearing isobar, e.g.  $f_0$ ,  $\rho(770)$ ,  $f_2(1270)$  with  $J_{\xi}^{PC}$
- $\pi$ : Indicating the bachelor  $\pi^-$ . Always the same
- L: Orbital angular momentum between isobar and bachelor pion





- $J_{X^-}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^-$
- $M^{\varepsilon}$ : Spin projection on the beam and naturality of the exchange particle
- $\xi$ : Appearing isobar, e.g.  $f_0$ ,  $\rho(770)$ ,  $f_2(1270)$  with  $J_{\xi}^{PC}$
- $\pi$ : Indicating the bachelor  $\pi^-$ . Always the same
- L: Orbital angular momentum between isobar and bachelor pion

Various coherent possibilities for quantum-number combinations
#### Wave naming scheme





- $J_{X^{-}}^{PC}$ : Spin and eigenvalues under parity and charge conjugation of  $X^{-}$
- $M^{\varepsilon}$ : Spin projection on the beam and naturality of the exchange particle
- $\xi$ : Appearing isobar, e.g.  $f_0$ ,  $\rho(770)$ ,  $f_2(1270)$  with  $J_{\xi}^{PC}$
- $\pi$ : Indicating the bachelor  $\pi^-$ . Always the same
- L: Orbital angular momentum between isobar and bachelor pion

Various coherent possibilities for quantum-number combinations



• Total intensity in one  $(m_{3\pi}, t')$  bin as function of phase-space variables  $\vec{\tau}$ :

$$\mathcal{I}(ec{ au}) = \left|\sum_{i}^{waves} \mathcal{T}_{i}[\psi_{i}(ec{ au}) \Delta_{i}(m_{\pi^{-}\pi^{+}}) + \text{Bose sym.}]
ight|^{2}$$

Fit parameters: Transition amplitudes  $T_i$ 

Fixed: Angular amplitudes  $\psi_i(\vec{\tau})$ , dynamic isobar amplitudes  $\Delta_i(m_{\pi^-\pi^+})$ 



• Total intensity in one  $(m_{3\pi}, t')$  bin as function of phase-space variables  $\vec{\tau}$ :

$$\mathcal{I}(ec{ au}) = \left|\sum_{i}^{waves} \mathcal{T}_{i}[\psi_{i}(ec{ au}) \Delta_{i}(m_{\pi^{-}\pi^{+}}) + \text{Bose sym.}]
ight|^{2}$$

Fit parameters: Transition amplitudes  $T_i$ Fixed: Angular amplitudes  $\psi_i(\vec{\tau})$ , dynamic isobar amplitudes  $\Delta_i(m_{\pi^-\pi^+})$ 

• Fixed isobar amplitudes  $\rightarrow$  Sets of  $m_{\pi^-\pi^+}$  bins:

$$\Delta_i (m_{\pi^-\pi^+}) 
ightarrow \sum_{\text{bins}} \mathscr{T}_i^{\text{bin}} \Delta_i^{\text{bin}} (m_{\pi^-\pi^+}) \equiv [\pi\pi]_{J^{PC}}$$
 $\Delta_i^{\text{bin}} (m_{\pi^-\pi^+}) = egin{cases} 1, & ext{if } m_{\pi^-\pi^+} & ext{in the bin.} \\ 0, & ext{otherwise.} \end{cases}$ 



• Total intensity in one  $(m_{3\pi}, t')$  bin as function of phase-space variables  $\vec{\tau}$ :

$$\mathcal{I}(ec{ au}) = \left|\sum_{i}^{waves} \mathcal{T}_{i}[\psi_{i}(ec{ au}) \Delta_{i}(m_{\pi^{-}\pi^{+}}) + \text{Bose sym.}]
ight|^{2}$$

Fit parameters: Transition amplitudes  $T_i$ Fixed: Angular amplitudes  $\psi_i(\vec{\tau})$ , dynamic isobar amplitudes  $\Delta_i(m_{\pi^-\pi^+})$ 

• Fixed isobar amplitudes  $\rightarrow$  Sets of  $m_{\pi^-\pi^+}$  bins:

$$\Delta_i (m_{\pi^-\pi^+}) 
ightarrow \sum_{\text{bins}} \mathscr{T}_i^{\text{bin}} \Delta_i^{\text{bin}} (m_{\pi^-\pi^+}) \equiv [\pi\pi]_{J^{PC}}$$
 $\Delta_i^{\text{bin}} (m_{\pi^-\pi^+}) = egin{cases} 1, & ext{if } m_{\pi^-\pi^+} & ext{in the bin.} \\ 0, & ext{otherwise.} \end{cases}$ 

• Each  $m_{\pi^-\pi^+}$  bin behaves like an independent partial wave with  $\mathcal{T}_i^{\text{bin}} = \mathcal{T}_i \mathscr{T}_i^{\text{bin}}$ :

$$\mathcal{I}(\vec{\tau}) = \left|\sum_{i}^{\text{waves}} \sum_{\text{bin}}^{\text{bin}} \mathcal{T}_{i}^{\text{bin}} \left[\psi_{i}(\vec{\tau}) \Delta_{i}^{\text{bin}}(m_{\pi^{-}\pi^{+}}) + \text{Bose sym.}\right]\right|^{2}$$









- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0 (GeV/c)<sup>2</sup>
  - 200 independent fits



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0 (GeV/c)<sup>2</sup>
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0 (GeV/c)<sup>2</sup>
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S$$
  
 $0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P$ 



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0 (GeV/c)<sup>2</sup>
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S$$
  

$$0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P$$
  

$$1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P$$
  

$$1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S$$



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0 (GeV/c)<sup>2</sup>
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$\begin{array}{lll} 0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S & 2^{-+}0^{+}[\pi\pi]_{0^{++}}\pi D \\ 0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P \\ 1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi F \\ 1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S & 2^{-+}0^{+}[\pi\pi]_{2^{++}}\pi S \end{array}$$



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0  $(\text{GeV}/c)^2$ 
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$\begin{array}{lll} 0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S & 2^{-+}0^{+}[\pi\pi]_{0^{++}}\pi D & 1^{++}1^{+}[\pi\pi]_{1^{--}}\pi S \\ 0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \\ 1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi F & 2^{++}1^{+}[\pi\pi]_{1^{--}}\pi D \\ 1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S & 2^{-+}0^{+}[\pi\pi]_{2^{++}}\pi S \end{array}$$



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0  $(\text{GeV}/c)^2$ 
  - ► 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$\begin{array}{lll} 0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S & 2^{-+}0^{+}[\pi\pi]_{0^{++}}\pi D & 1^{++}1^{+}[\pi\pi]_{1^{--}}\pi S \\ 0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \\ 1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi F & 2^{++}1^{+}[\pi\pi]_{1^{--}}\pi D \\ 1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S & 2^{-+}0^{+}[\pi\pi]_{2^{++}}\pi S & 1^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \end{array}$$



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0  $(\text{GeV}/c)^2$ 
  - 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- I2 waves freed (72 remaining waves still with fixed isobars):

$$\begin{array}{lll} 0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S & 2^{-+}0^{+}[\pi\pi]_{0^{++}}\pi D & 1^{++}1^{+}[\pi\pi]_{1^{--}}\pi S \\ 0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \\ 1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi F & 2^{++}1^{+}[\pi\pi]_{1^{--}}\pi D \\ 1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S & 2^{-+}0^{+}[\pi\pi]_{2^{++}}\pi S & 1^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \end{array}$$

• 40 MeV  $m_{\pi^-\pi^+}$  bins for freed waves, finer binnings in regions of known resonances:  $f_0(980)$ ,  $\rho(770)$ ,  $f_2(1270)$ 



- 50 bins in  $m_{3\pi}$  from 0.5 to 2.5 GeV, 4 bins in t' from 0.1 to 1.0  $(\text{GeV}/c)^2$ 
  - 200 independent fits
- Wave set: Based on 88 partial-waves model COMPASS collaboration, PRD 95, (2017) 032004
- 12 waves freed (72 remaining waves still with fixed isobars):

$$\begin{array}{lll} 0^{-+}0^{+}[\pi\pi]_{0^{++}}\pi S & 2^{-+}0^{+}[\pi\pi]_{0^{++}}\pi D & 1^{++}1^{+}[\pi\pi]_{1^{--}}\pi S \\ 0^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi P & 2^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \\ 1^{++}0^{+}[\pi\pi]_{0^{++}}\pi P & 2^{-+}0^{+}[\pi\pi]_{1^{--}}\pi F & 2^{++}1^{+}[\pi\pi]_{1^{--}}\pi D \\ 1^{++}0^{+}[\pi\pi]_{1^{--}}\pi S & 2^{-+}0^{+}[\pi\pi]_{2^{++}}\pi S & 1^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P \end{array}$$

- 40 MeV  $m_{\pi^-\pi^+}$  bins for freed waves, finer binnings in regions of known resonances:  $f_0(980)$ ,  $\rho(770)$ ,  $f_2(1270)$
- Depending on  $m_{3\pi}$  and wave, up to 62  $m_{\pi^-\pi^+}$  bins per freed wave





- Focus here: Spin-exotic wave with quantum numbers:  $J_{\chi^-}^{PC} = 1^{-+}$ 
  - Spin-exotic: Not a simple quark-model  $q\bar{q}$  state





- Focus here: Spin-exotic wave with quantum numbers:  $J_{X^-}^{PC} = 1^{-+}$ 
  - Spin-exotic: Not a simple quark-model  $q\bar{q}$  state
- Resonance content disputed for a long time
- Expected:  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$





- Focus here: Spin-exotic wave with quantum numbers:  $J_{X^{-}}^{PC} = 1^{-+}$ 
  - Spin-exotic: Not a simple quark-model  $q\bar{q}$  state
- Resonance content disputed for a long time
- Expected:  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$
- Dynamic isobar amplitude: Dominated by  $\rho(770)$





- Focus here: Spin-exotic wave with quantum numbers:  $J_{\chi^-}^{PC} = 1^{-+}$ 
  - Spin-exotic: Not a simple quark-model  $q\bar{q}$  state
- Resonance content disputed for a long time
- Expected:  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$
- Dynamic isobar amplitude: Dominated by  $\rho(770)$
- Possible additional effects on dynamic amplitude:

  - Rescattering with the third (bachelor)  $\pi^-$
  - Non-resonant contributions





- Focus here: Spin-exotic wave with quantum numbers:  $J_{\chi^-}^{PC} = 1^{-+}$ 
  - Spin-exotic: Not a simple quark-model  $q\bar{q}$  state
- Resonance content disputed for a long time
- Expected:  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$
- Dynamic isobar amplitude: Dominated by  $\rho(770)$
- Possible additional effects on dynamic amplitude:

  - Rescattering with the third (bachelor)  $\pi^-$
  - Non-resonant contributions

Ambiguity in this fit model!



• Freed-isobar analysis: Much more parameters than fixed-isobar analysis



- Freed-isobar analysis: Much more parameters than fixed-isobar analysis
  - May causes continuous mathematical ambiguities in the model

# Zero mode in the spin-exotic wave What is a "zero mode"?

- Ap. Dg > 1 t
- Freed-isobar analysis: Much more parameters than fixed-isobar analysis
  - ► May causes continuous mathematical ambiguities in the model
- Example: Spin-exotic wave:  $\mathcal{A}(\vec{\tau}) = \psi(\vec{\tau}) \Delta(m_{\pi^-\pi^+})$ with  $\psi(\vec{\tau}) \propto \vec{\rho}_1 \times \vec{\rho}_3 \quad (\propto \sin \theta_{\mathsf{HF}})$

# Zero mode in the spin-exotic wave What is a "zero mode"?

- Ar Dy>it
- Freed-isobar analysis: Much more parameters than fixed-isobar analysis
  - May causes continuous mathematical ambiguities in the model
- Example: Spin-exotic wave:  $\mathcal{A}(\vec{\tau}) = \psi(\vec{\tau}) \Delta(m_{\pi^-\pi^+})$ with  $\psi(\vec{\tau}) \propto \vec{p}_1 \times \vec{p}_3 \quad (\propto \sin \theta_{\text{HF}})$
- Bose symmetrized amplitude for  $\pi_1^- \pi_2^+ \pi_3^-$  final state:

$$\mathcal{A}_{1^{-+}}^{\text{symm}}\left(\vec{\tau}\right) = \left(\vec{p}_{1} \times \vec{p}_{3}\right) \Delta^{0}(m_{\pi_{1}^{-}\pi_{2}^{+}}) + \left(\vec{p}_{3} \times \vec{p}_{1}\right) \Delta^{0}(m_{\pi_{2}^{+}\pi_{3}^{-}})$$

# Zero mode in the spin-exotic wave What is a "zero mode"?

- Ar Ag > it
- Freed-isobar analysis: Much more parameters than fixed-isobar analysis
  - May causes continuous mathematical ambiguities in the model
- Example: Spin-exotic wave:  $\mathcal{A}(\vec{\tau}) = \psi(\vec{\tau}) \Delta(m_{\pi^-\pi^+})$ with  $\psi(\vec{\tau}) \propto \vec{p}_1 \times \vec{p}_3 \quad (\propto \sin \theta_{\text{HF}})$
- Bose symmetrized amplitude for  $\pi_1^- \pi_2^+ \pi_3^-$  final state:

$$\begin{split} \mathcal{A}_{1-+}^{\text{symm}}\left(\vec{\tau}\right) &= \left(\vec{\rho}_{1} \times \vec{\rho}_{3}\right) \Delta^{0}(m_{\pi_{1}^{-}\pi_{2}^{+}}) + \left(\vec{\rho}_{3} \times \vec{\rho}_{1}\right) \Delta^{0}(m_{\pi_{2}^{+}\pi_{3}^{-}}) \\ &= \left(\vec{\rho}_{1} \times \vec{\rho}_{3}\right) \left[\Delta^{0}(m_{\pi_{1}^{-}\pi_{2}^{+}}) - \Delta^{0}(m_{\pi_{2}^{+}\pi_{3}^{-}})\right] \end{split}$$

vanishes at every point  $\vec{\tau}$  in phase space, if  $\Delta^0$  (*m*) is constant



- Freed-isobar analysis: Much more parameters than fixed-isobar analysis
  - ► May causes continuous mathematical ambiguities in the model
- Example: Spin-exotic wave:  $\mathcal{A}(\vec{\tau}) = \psi(\vec{\tau}) \Delta(m_{\pi^-\pi^+})$ with  $\psi(\vec{\tau}) \propto \vec{\rho}_1 \times \vec{\rho}_3 \quad (\propto \sin \theta_{\text{HF}})$
- Bose symmetrized amplitude for  $\pi_1^- \pi_2^+ \pi_3^-$  final state:

$$\begin{split} \mathcal{A}_{1-+}^{\text{symm}}\left(\vec{\tau}\right) &= \left(\vec{\rho}_{1} \times \vec{\rho}_{3}\right) \Delta^{0}(m_{\pi_{1}^{-}\pi_{2}^{+}}) + \left(\vec{\rho}_{3} \times \vec{\rho}_{1}\right) \Delta^{0}(m_{\pi_{2}^{+}\pi_{3}^{-}}) \\ &= \left(\vec{\rho}_{1} \times \vec{\rho}_{3}\right) \left[\Delta^{0}(m_{\pi_{1}^{-}\pi_{2}^{+}}) - \Delta^{0}(m_{\pi_{2}^{+}\pi_{3}^{-}})\right] \end{split}$$

vanishes at every point  $\vec{\tau}$  in phase space, if  $\Delta^0$  (*m*) is constant

• "Zero mode": Shift of dynamic isobar amplitude by arbitrary  ${\mathcal C}$ 

$$\Delta^{\mathrm{meas}}\left(m_{\xi}
ight)=\Delta^{\mathrm{phys}}\left(m_{\xi}
ight)+\mathcal{C}\Delta^{0}\left(m_{\xi}
ight)$$

leaves  $\mathcal{A}^{symm}_{1-+}$  and therefore intensity and likelihood invariant  $\Rightarrow$  Ambiguous solutions



#### • Superfluous degree of freedom C: Indistinguishable by fit



- Superfluous degree of freedom C: Indistinguishable by fit
- Physical solution: Conditions on dynamic isobar amplitudes Δ<sup>meas</sup>



- Superfluous degree of freedom C: Indistinguishable by fit
- Physical solution: Conditions on dynamic isobar amplitudes Δ<sup>meas</sup>
- In the case of the  $1^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P$  wave:
  - ► Use the Breit-Wigner for the  $\rho(770)$  resonance with fixed parameters as in the fixed-isobar analysis
  - ► limit fit range to  $m_{\pi^-\pi^+} < 1.12 \,\text{GeV}$  to minimize effects from possible excited  $\rho'$  states



- Superfluous degree of freedom C: Indistinguishable by fit
- Physical solution: Conditions on dynamic isobar amplitudes Δ<sup>meas</sup>
- In the case of the  $1^{-+}1^{+}[\pi\pi]_{1^{--}}\pi P$  wave:
  - ► Use the Breit-Wigner for the  $\rho(770)$  resonance with fixed parameters as in the fixed-isobar analysis
  - ► limit fit range to  $m_{\pi^-\pi^+}$  < 1.12 GeV to minimize effects from possible excited  $\rho'$  states
- Note: Resolving the ambiguity fixes only a single complex-valued degree of freedom, C. n<sub>bins</sub> 1 complex-valued degrees of freedom remain free.
   FK, D. Greenwald, D. Ryabchikov, B. Grube, S. Paul, PRD 97 (2018) 114008





Resolving the ambiguity





Resolving the ambiguity


























 $0.3260 < t' < 1.000 \left( {\rm GeV} / c \right)^2$ 



Ambiguity resolved in every  $m_{3\pi}$  bin separately



## $0.3260 < t' < 1.000 \left( \text{GeV}/c \right)^2$



Ambiguity resolved in every  $m_{3\pi}$  bin separately



## $0.3260 < t' < 1.000 \left( \text{GeV}/c \right)^2$



Ambiguity resolved in every  $m_{3\pi}$  bin separately



## $0.3260 < t' < 1.000 \left( \text{GeV}/c \right)^2$



Ambiguity resolved in every  $m_{3\pi}$  bin separately Correlation of  $\pi_1$  (1600) with  $\rho$ (770) confirmed Freed-isobar result Comparison to fixed-isobar PWA



 Coherently sum up all m<sub>π<sup>-</sup>π<sup>+</sup></sub> bins to obtain m<sub>3π</sub> spectrum





- Coherently sum up all m<sub>π<sup>-</sup>π<sup>+</sup></sub> bins to obtain m<sub>3π</sub> spectrum
- Zero mode exactly cancels





- Coherently sum up all m<sub>π<sup>-</sup>π<sup>+</sup></sub> bins to obtain m<sub>3π</sub> spectrum
- Zero mode exactly cancels
- Similar to fixed-isobar result for π<sub>1</sub> (1600)





- Coherently sum up all m<sub>π<sup>-</sup>π<sup>+</sup></sub> bins to obtain m<sub>3π</sub> spectrum
- Zero mode exactly cancels
- Similar to fixed-isobar result for π<sub>1</sub> (1600)
- Isobar model: Valid assumption





- Coherently sum up all  $m_{\pi^-\pi^+}$  bins to obtain  $m_{3\pi}$  spectrum
- Zero mode exactly cancels
- Similar to fixed-isobar result for  $\pi_1$  (1600)
- Isobar model: Valid assumption
- Observed deviations hint to:
  - Excited isobar resonances
  - Final state interactions
  - Non-resonant contributions
































































































































































Conclusion: Extended freed-isobar analysis with 12 out of 88 freed waves

- In total 200 independent fits in  $m_{3\pi}$  and t' bins
- Independent dynamic isobar amplitudes obtained in every fit
- Zero mode ambiguities resolved



Conclusion: Extended freed-isobar analysis with 12 out of 88 freed waves

- In total 200 independent fits in  $m_{3\pi}$  and t' bins
- Independent dynamic isobar amplitudes obtained in every fit
- Zero mode ambiguities resolved
- Decay  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$  reconfirmed
  - Independent of  $\rho(770)$  parameterization



Conclusion: Extended freed-isobar analysis with 12 out of 88 freed waves

- In total 200 independent fits in  $m_{3\pi}$  and t' bins
- Independent dynamic isobar amplitudes obtained in every fit
- Zero mode ambiguities resolved
- Decay  $\pi_1$  (1600)  $\rightarrow \rho$ (770)  $\pi^-$  reconfirmed
  - Independent of  $\rho(770)$  parameterization

Next step: Analyze extracted dynamic isobar amplitudes

- Pin down resonance parameters of  $\rho'$
- Study 2-body resonances in 3-body environment
  - Study 3-body re-scattering effects
  - ► e.g. Khuri-Treiman amplitudes (Bonn group)