Role of the tensor force in the heavy hadronic molecules

Yasuhiro Yamaguchi (RIKEN, Japan)

in collaboration with

Hugo García-Tecocoatzi (UNAM), Alessandro Giachino (INFN Genoa, Genoa Univ.), Atsushi Hosaka (RCNP, Osaka Univ.), Elena Santopinto (INFN Genoa), Sachiko Takeuchi (Japan Coll. Social Work), Makoto Takizawa (Showa Pharmaceutical Univ.).

The 18th International Conference on Hadron Spectroscopy and Structure (HADRON2019)

Guilin, China 16-21 August 2019

Outline

Introduction

- Exotic hadrons Hadronic molecules
- Hidden-charm pentaquark Pc
- Odel setup
 - One Pion Exchange Potential
 - Compact 5-quark potential
- Numerical results
 - Hidden-charm molecules
- Summary



Hadronic molecule



Pentaquark (Compact)

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Hadron structure: Constituent quark model

• Ordinary Hadrons: Baryon (qqq) and Meson $(q\bar{q})$



• Exotic Hadrons ($\neq qqq, q\bar{q}$): Multiquark? Multihadron?



Constituent quark picture and beyond Introduction



S. Godfrey and N. Isgur, PRD32(1985)189

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Constituent quark picture and beyond Introduction



N. Brambilla, et al. Eur. Phys. J.C 71(2011)1534,

S. Godfrey and N. Isgur, PRD32(1985)189

Exotics ≠ cc̄ have been observed in the Experiments (BaBar, Belle, BESIII, LHCb,...) ⇒ Q. Structure? Physics?

Hadronic molecules? Introduction

• Exotics as Hadronic molecule = Hadron composite system



Hadronic molecule

$$X(3872) = D\bar{D}^* ext{ molecule}?$$

 $D^+D_-^{*-} ext{ 3879.84 MeV}$
 $D^0\bar{D}_-^{*0} ext{ - 3871.69 MeV} ext{ 3871.68 MeV}$

- Expected near the thresholds
 - \rightarrow Hadron-hadron (quasi) bound state
- ⇒ Analogous to Atomic Nuclei

Deuteron $\sim pn$ bound state (B = 2.2 MeV)



Hadronic molecules and π exchange potential $_{\rm Introduction}$

- Driving force of Nuclei \Rightarrow long range force: π exchange
 - \longrightarrow generating the loosely bound state



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• Strong attraction from Tensor term (S – D mixing)

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Hadronic molecules and π exchange potential Introduction

- Driving force of Nuclei \Rightarrow long range force: π exchange
 - \rightarrow generating the loosely bound state



• Strong attraction from **Tensor term** (S - D mixing)Important role in the heavy hadronic moleucles? \Rightarrow

18 Aug 2019



Hidden-charm pentaquarks



R.Aaij, et al. (LHCb collaboration) PRL115(2015)072001, PRL122(2019)222001

Two hidden-charm pentaquarks !! (2015)

Introduction: pentaquark

• Observation of the Hidden-charm Pentaquark ($c\bar{c}uud$) in $\Lambda_b^0 \rightarrow J/\psi K^- p$ Decay? R.Aaij, et al. (LHCb collaboration) PRL115(2015)072001



- $P_c(4380)$: M = 4380 MeV $P_c(4450)$: M = 4449.8 MeV $\Gamma = 205 \text{ MeV}$ $\Gamma = 39 \text{ MeV}$
 - J^{P} ? (3/2⁻, 5/2⁺), (3/2⁺, 5/2⁻), or (5/2⁺, 3/2⁻)
 - There have been a lot of articles investigating the *P_c* states... Hadronic molecule? Compact state? Kinematical effect?

H.X.Chen, et al., Phys.Rept.639(2016)1, A.Esposito, et al., Phys.Rept.668(2016)1, A.Ali, et al., PPNP97(2017)123

New LHCb analysis in 2019!

Introduction: pentaquark

• R. Aaij, et al. Phys.Rev.Lett. 122 (2019) no.22, 222001



- $\begin{array}{c|c} \triangleright & P_c(4450) \text{ in } 2015 \longrightarrow P_c(4440) \text{ and } P_c(4457) \\ P_c(4440) & M = 4440.3 \text{ MeV} \quad P_c(4457) \quad M = 4457.3 \text{ MeV} \\ \Gamma = 20.6 & \Gamma = 6.4 \text{ MeV} \end{array}$
- ▷ Observation of New state! $P_c(4312)$ M = 4311.9 MeV $\Gamma = 9.8 \text{ MeV}$

18 Aug 2019

Yasuhiro Yamaguchi(RIKEN)

Hadron2019(Guilin)

Hidden-charm meson-baryon molecule...?

Introduction: pentaquark

 \triangleright $P_{\rm c}$ states reported close to the $\overline{D}^{(*)}\Sigma_c^{(*)}$ thresholds

 $\triangleright \pi$ exchange in the heavy hadron system

Hidden-charm meson-baryon molecule...?

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π exchange in the heavy hadron system
 enhanced by the heavy quark spin symmetry!

N.Isgur, M.B.Wise, PLB232(1989)113

 $\Rightarrow \bar{D}(0^{-}) - \bar{D}^{*}(1^{-}), \Sigma_{c}(1/2^{+}) - \Sigma_{c}^{*}(3/2^{-})$ mixing

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Compact state: 5-quark configuration

Introduction: pentaquark

- S. Takeuchi and M. Takizawa, PLB764 (2017) 254-259.
 - P_c states by the quark cluster model
- 5-quark configurations



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• $[q^3 8_c 3/2]$: Color magnetic int. is attractive!

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 $S_{q^3}=1/2, {3/2}, \; S_{car{c}}=0,1 \quad S_{q^3}=1/2, \; S_{car{c}}=0,1$

- $[q^3 8_c 3/2]$: Color magnetic int. is attractive!
 - ⇒ Couplings to (qqc) baryon- $(q\bar{c})$ meson, e.g. $\bar{D}\Sigma_c$, are allowed!

Mixing of Compact state and Hadronic Molecule!

• Hadronic molecule (MB) + Compact state (5q)



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Hadronic molecule (MB) + Compact state (5q)
 ⇒ MB coupled to 5q (Feshbach Projection)



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Hadronic molecule (MB) + Compact state (5q)
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Long range interaction: One pion exchange potential (OPEP)
 Short range interaction: 5q potential

Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, PRD96(2017)114031

Y.Y, H.G-Tecocoatzi, A.Giachino, A.Hosaka, E.Santopinto, S.Takeuchi, M.Takizawa, 1907.04684[hep_ph] 👘

Hadronic molecule (MB) + Compact state (5q)
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Long range interaction: One pion exchange potential (OPEP)
 Short range interaction: 5q potential (→Local Gaussian)

Spin dependence \rightarrow Spin structure of 5q

Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, PRD96(2017)114031

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$\bar{D}^{(*)}Y_c$ Interaction: Long range force HQS and OPEP

• One pion exchange potential

$$\begin{array}{c|c}
\bar{D}^{(*)} & Y_{c} \\
\mathcal{L}_{\pi\bar{D}^{(*)}\bar{D}^{(*)}} & \bar{D}^{(*)} : \bar{D} \text{ or } \bar{D}^{*} \\
\mathcal{L}_{\pi\bar{D}^{(*)}\bar{D}^{(*)}} & Y_{c} \\
\bar{D}^{(*)} & Y_{c} \\
V_{\bar{D}^{(*)}Y_{c}-\bar{D}^{(*)}Y_{c}} = \frac{g_{D^{*}D\pi}g_{Y_{c}Y_{c}\pi}}{f_{\pi}^{2}} \left[\vec{S}_{1} \cdot \vec{S}_{2}C(r) + S_{S_{1}S_{2}}T(r) \right] \\
\end{array}$$
(Contact term is removed)

• Form factor with Cutoff Λ (determined by the hadron size)

$$F(q^2)=rac{\Lambda^2-m_\pi^2}{\Lambda^2-q^2}, \hspace{1em} \Lambda_{ar{D}}\sim 1130 \hspace{1em} ext{MeV}, \Lambda_{Y_{ ext{c}}}\sim 840 \hspace{1em} ext{MeV}$$

Y.Y. A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, Phys. Rev. D96 (2017), 114031

2. Short range force: 5-quark potential



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- 5-quark potential \Rightarrow **Local Gaussian potential** is employed.
- \triangleright Massive M_{5q} (few hundred MeV above $\bar{D}^*\Sigma^*_{\mathrm{c}}) \rightarrow$ Attractive



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Free Parameters

Strength f and Gaussian para. α (\rightarrow may be fixed in the future) (f vs E will be shown latter. $\alpha = 1$ fm⁻² is fixed.)

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Relative strength S_i

Spectroscopic factors \Rightarrow determined by the spin structure of 5q

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Spectroscopic factor S_i (Spin structure) 5q potential

• S-factor:
$$S_i = \langle (\bar{D}Y_c)_i | 5q \rangle$$

Table: Spectroscopic factors S_i for each meson-baryon channel.

| J | | $S_{c\bar{c}}$ | S _{3q} | $\bar{D}\Lambda_{ m c}$ | $ar{D}^* \Lambda_{ m c}$ | $ar{D}\Sigma_{ m c}$ | $ar{D}\Sigma_{ m c}^{*}$ | $ar{D}^*\Sigma_{ m c}$ | $ar{D}^*\Sigma^*_{ m c}$ |
|-----|-------|----------------|-----------------|-------------------------|--------------------------|----------------------|--------------------------|------------------------|--------------------------|
| 1/2 | (i) | 0 | 1/2 | 0.4 | 0.6 | -0.4 | _ | 0.2 | -0.6 |
| | (ii) | 1 | 1/2 | 0.6 | -0.4 | 0.2 | — | -0.6 | -0.3 |
| | (iii) | 1 | 3/2 | 0.0 | 0.0 | -0.8 | — | -0.5 | 0.3 |
| 3/2 | (i) | 0 | 3/2 | _ | 0.0 | | -0.5 | 0.6 | -0.7 |
| | (ii) | 1 | 1/2 | | 0.7 | — | 0.4 | -0.2 | -0.5 |
| | (iii) | 1 | 3/2 | _ | 0.0 | | -0.7 | -0.8 | -0.2 |
| 5/2 | (i) | 1 | 3/2 | — | — | | | — | -1.0 |

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• $\overline{D}Y_{c}$ with Large S_{i} will play an important role,

Numerical Results for Hidden-charm sector



Bound state and Resonance

- Coupled-channel Schrödinger equation for $\bar{D}\Lambda_c$, $\bar{D}^*\Lambda_c$, $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$, $\bar{D}^*\Sigma_c^*$ (6 *MB* components).
- For $J^P = 1/2^-$, $3/2^-$, $5/2^-$ (Negative parity)

Results (f^{5q} vs E) of charm $\overline{D}Y_c$ for $J^P = 1/2^-$

• Energy with $V_{\pi} + V^{5q}(f^{5q})$. (Y.Yamaguchi *et al*, PRD**96** (2017), 114031)



Dashed line: Thresholds, Red line: Energy obtained

- For small f^{5q} , **No bound state**
 - \Rightarrow The OPEP attraction is not enough to generate a state
- 5q potential helps to generate the states near the thresholds
 ⇔ Large S-factor (Spin structure)

Results $(f^{5q} \text{ vs } E)$ for $J^P = 3/2^-, 5/2^-$

• Energy with $V_{\pi} + V^{5q}(f^{5q})$. (Y.Yamaguchi *et al*, PRD**96** (2017), 114031)



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In 2019, New P_c states by LHCb

• $f^{5q} = 45$ is fixed to reproduce new P_c 's

Y.Y., H.Garcia-Tecocoatzi, A.Giachino, A.Hosaka, E.Santopinto, S.Takeuchi, M.Takizawa, 1907.04684[hep-ph]

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- Missing three $(1/2^-, 3/2^-, 5/2^-)$ states below $ar{D}^*\Sigma_c^*$
- Broad $P_c(4380) \leftrightarrow \text{Our prediction with } 4376 \text{ MeV}$
 - \rightarrow Further theoretical and experimental studies are necessary...

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J^P assignment for $P_c(4440)$ and $P_c(4457)$



* C: Central force, T: Tensor force

J^P assignment for $P_c(4440)$ and $P_c(4457)$



* C: Central force, T: Tensor force

• S - D, D - D couplings producing the attraction from the tensor force

$$\Rightarrow 1/2^{-}: {}^{2}S - {}^{4}D$$

3/2⁻: {}^{4}S - {}^{2}D, {}^{4}S - {}^{4}D, {}^{2}D - {}^{4}D \leftarrow \text{more attractive!}

Summary

- Hidden-charm pentaquarks as Hadronic molecule + Compact multiquark
- $\bar{D}^{(*)}Y_c$ Interaction

Long range force: OPEP with the tensor force, enhanced by the heavy quark symmetry Short range force: Coupling to Compact 5q states

- The OPEP is not enough to generate the bound state. \rightarrow OPEP + 5q potential generates the states
- Applying this model to New hidden-charm pentaquarks by LHCb in 2019

 \Rightarrow our prediction is consistent with EXP

• The J^P assignment $P_c(4440)$: $3/2^-$ and $P_c(4457)$: $1/2^-$ understood by the tensor force of the OPEP

Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, Phys.Rev. D96 (2017), 114031

Y.Y., H.Garcia-Tecocoatzi, A.Giachino, A.Hosaka, E.Santopinto, S.Takeuchi, M.Takizawa, 1907,04684[hep-ph] 🦏

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| State | Mass | Width | Our pred. (M,Γ,J^P) |
|-------------------|--------------------------------|-----------------------------|--------------------------------|
| $P_{c}(4312)^{+}$ | $4311.9\pm0.7^{+6.8}_{-0.6}$ | $9.8\pm2.7^{+3.7}_{-4.5}$ | $(4312, 5, \frac{1}{2}^{-})$ |
| $P_{c}(4380)^{+}$ | $4380\pm8\pm29$ | $205\pm18\pm86$ | $(4376,8,\frac{3}{2}^{-})$ |
| $P_{c}(4440)^{+}$ | $4440.3 \pm 1.3^{+4.1}_{-4.7}$ | $20.6\pm4.9^{+8.7}_{-10.1}$ | $(4442, 26, \frac{3}{2})$ |
| $P_{c}(4457)^{+}$ | $4457.3\pm0.6^{+4.1}_{-1.7}$ | $6.4\pm2.0^{+5.7}_{-1.9}$ | $(4462, 6.6, \frac{1}{2}^{-})$ |
| | | | $(4524, 1.5, \frac{1}{2}^{-})$ |
| | | | $(4521, 23, \frac{3}{2})$ |
| | | | $(4511,55, \frac{5}{2}^{-})$ |

(in units of MeV)

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Coupled-channels

| Channels | $\bar{D}Y_{c}(^{2S+1}L)$ | |
|-----------|--|---------|
| $1/2^{-}$ | $\bar{D}\Lambda_{\rm c}(^2S),\ \bar{D}^*\Lambda_{\rm c}(^2S),$ | |
| | $\bar{D}\Sigma_{\rm c}(^2S), \ \bar{D}\Sigma_{\rm c}^*(^4D),$ | |
| | $ar{D}^* \Sigma_{ m c}({}^2S, {}^4D), \ ar{D}^* \Sigma_{ m c}^* ({}^2S, {}^4D, {}^6D)$ | (10 ch) |
| 3/2- | $\bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^4S,^2D,^4D),$ | |
| | $ar{D}\Sigma_{ m c}(^2D),\ ar{D}\Sigma_{ m c}^*(^4S,^4D),$ | |
| | $ar{D}^* \Sigma_{ m c}({}^4S, {}^2D, {}^4D), \ ar{D}^* \Sigma_{ m c}^*({}^4S, {}^2D, {}^4D, {}^6D, {}^6G)$ | (15 ch) |
| $5/2^{-}$ | $\overline{D}\Lambda_{\mathrm{c}}(^{2}D),\ \overline{D}^{*}\Lambda_{\mathrm{c}}(^{2}D,^{4}D,^{4}G),$ | |
| | $\bar{D}\Sigma_{\rm c}(^2D),\ \bar{D}\Sigma_{\rm c}^*(^4D, ^4G),$ | |
| | $\bar{D}^*\Sigma_{\rm c}(^2D, ^4D, ^4G), \ \bar{D}^*\Sigma_{\rm c}^*(^6S, ^2D, ^4D, ^6D, ^4G, ^6G)$ | (16 ch) |
| _ | | |

• 6 $\bar{D}Y_{c}$ channels: $\bar{D}\Lambda_{c}$, $\bar{D}^{*}\Lambda_{c}$, $\bar{D}\Sigma_{c}$, $\bar{D}\Sigma_{c}^{*}$, $\bar{D}^{*}\Sigma_{c}$, $\bar{D}^{*}\Sigma_{c}^{*}$.

• S - D mixing induced by the Tensor force (S_{12})

Heavy hadron- π coupling HQS and OPEP

• Effective Lagrangians: Heavy hadron and π

R. Casalbuoni *et al.*, Phys.Rept.**281** (1997)145, T. M. Yan, *et al.*, PRD**46**(1992)1148

Y.-R.Liu and M.Oka, PRD85(2012)014015



▷ Heavy meson: $\overline{D}^{(*)}\overline{D}^{(*)}\pi$ (*DD* π : Parity violation)

$$\mathcal{L}_{\pi HH} = -\frac{g_{\pi}}{2f_{\pi}} \text{Tr} \left[H \gamma_{\mu} \gamma_5 \partial^{\mu} \hat{\pi} \bar{H} \right], \quad H = \frac{1 + \not}{2} \left[D_{\mu}^* \gamma^{\mu} - D \gamma_5 \right]$$

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Heavy hadron- π coupling HQS and OPEP

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 $\mathcal{L}_{\pi HH} = -\frac{\underline{\mathbf{g}}_{\pi}}{2f_{\pi}} \mathrm{Tr} \left[H \gamma_{\mu} \gamma_{5} \partial^{\mu} \hat{\pi} \bar{H} \right], \quad \mathbf{H} = \frac{\mathbf{1} + \mathbf{\dot{\nu}}}{2} \left[\mathbf{D}_{\mu}^{*} \gamma^{\mu} - \mathbf{D} \gamma_{5} \right]$

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Heavy hadron- π coupling HQS and OPEP

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Spectroscopic factors *S_i* (**Spin structure**) ^{5*q* potential}

• Spin of 5q states $\rightarrow S_{c\bar{c}}$ and S_{3q} configuration e.g. for $J^P = 1/2^-$, (i), (ii), (iii)



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Spectroscopic factors S_i (Spin structure) 5*q* potential

• Spin of 5*q* states $\rightarrow S_{c\bar{c}}$ and S_{3q} configuration e.g. for $J^P = 1/2^-$, (i), (ii), (iii)



• Overlap of the spin wavefunctions of 5-quark state and $ar{D}Y_{
m c}$

$$S_i = \left\langle (\bar{D}Y_{
m c})_i \, \middle| \, 5q \right
angle$$

 \Rightarrow Relative strength of couplings to $\bar{D}Y_{c}$ channel

Volume integrals of the potentials

Bound and Resonant states appears for *f*^{5q} ≥ 25
 ⇔ Large? Small?

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Volume integrals of the potentials

- Bound and Resonant states appears for *f*^{5q} ≥ 25
 ⇔ Large? Small?
- ▷ Volume integral $V(q = 0) = \int V(r)dr^3$ Comparison with the *NN* interaction (Bonn potential) R. Machleidt, K. Holinde and C. Elster, Phys. Rept. **149**, 1 (1987).

$$ig| V_{f=25}^{5q}(0) ig| = 1.1 imes 10^{-4} \text{ MeV} \sim 0.03 |C_{NN}^{\sigma}(0)|$$

 $(C_{NN}^{\sigma}: \text{Central force of } \sigma \text{ exchange})$

• $\left|V_{f=25}^{5q}(0)\right|$ is much smaller than $|C_{NN}^{\sigma}(0)|$. However, the bound and resonant states are obtained!

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