

Igor Danilkin

August 20, 2019, Gulin, China



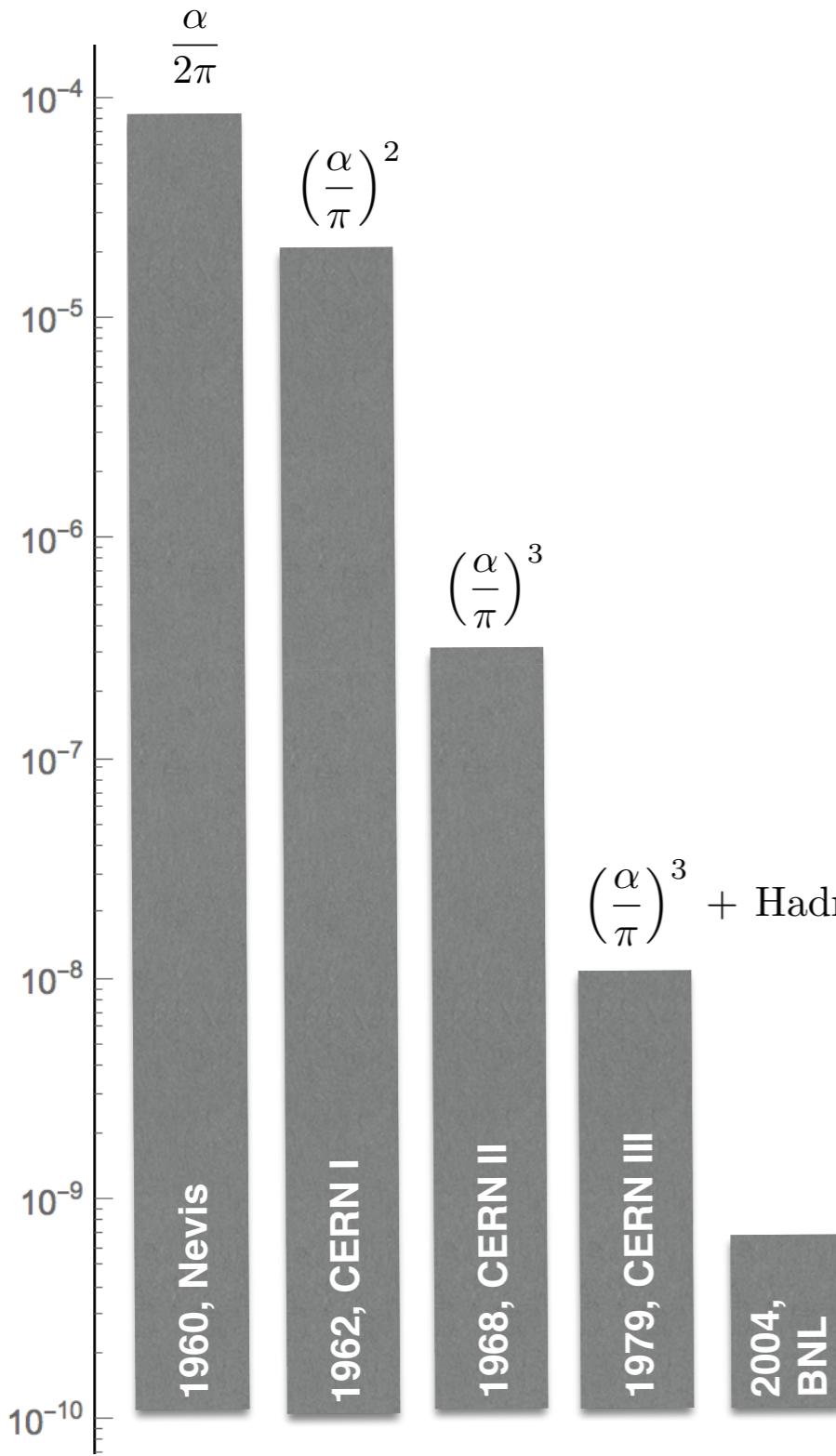
THE LOW-ENERGY FRONTIER
OF THE STANDARD MODEL

JOHANNES GUTENBERG
UNIVERSITÄT MAINZ



Motivation

- Accuracy



- Magnetic moment of the muon

$$\vec{\mu} = \frac{Q}{2m} g \vec{S}$$

- Today

Experiment (BNL 2004)

Theory (Standard model)

Difference

$$a_\mu^{\text{exp}} = 0.0011659209(6)$$

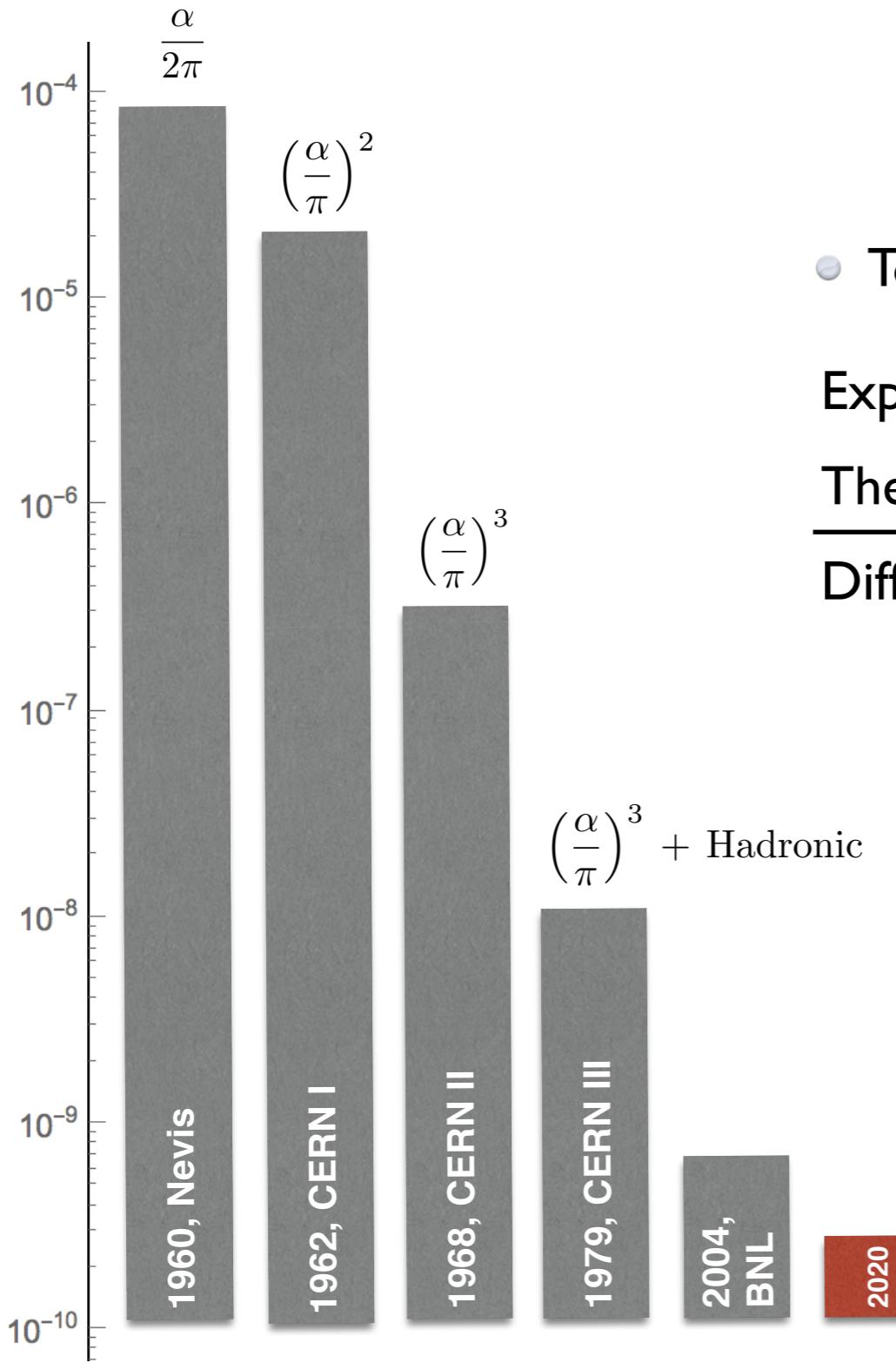
$$a_\mu^{\text{th}} = 0.0011659183(5)$$

$\sim (3 - 4) \sigma$

Hagiwara et al. (2011)

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$$a_\mu^{\text{exp}} = 0.0011659209(6)$$

$$a_\mu^{\text{th}} = 0.0011659183(5)$$

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Hagiwara et al. (2011)

$$\Delta a_\mu^{\text{exp}} = 6.3 \times 10^{-10}$$

$$\rightarrow 1.6 \times 10^{-10}$$

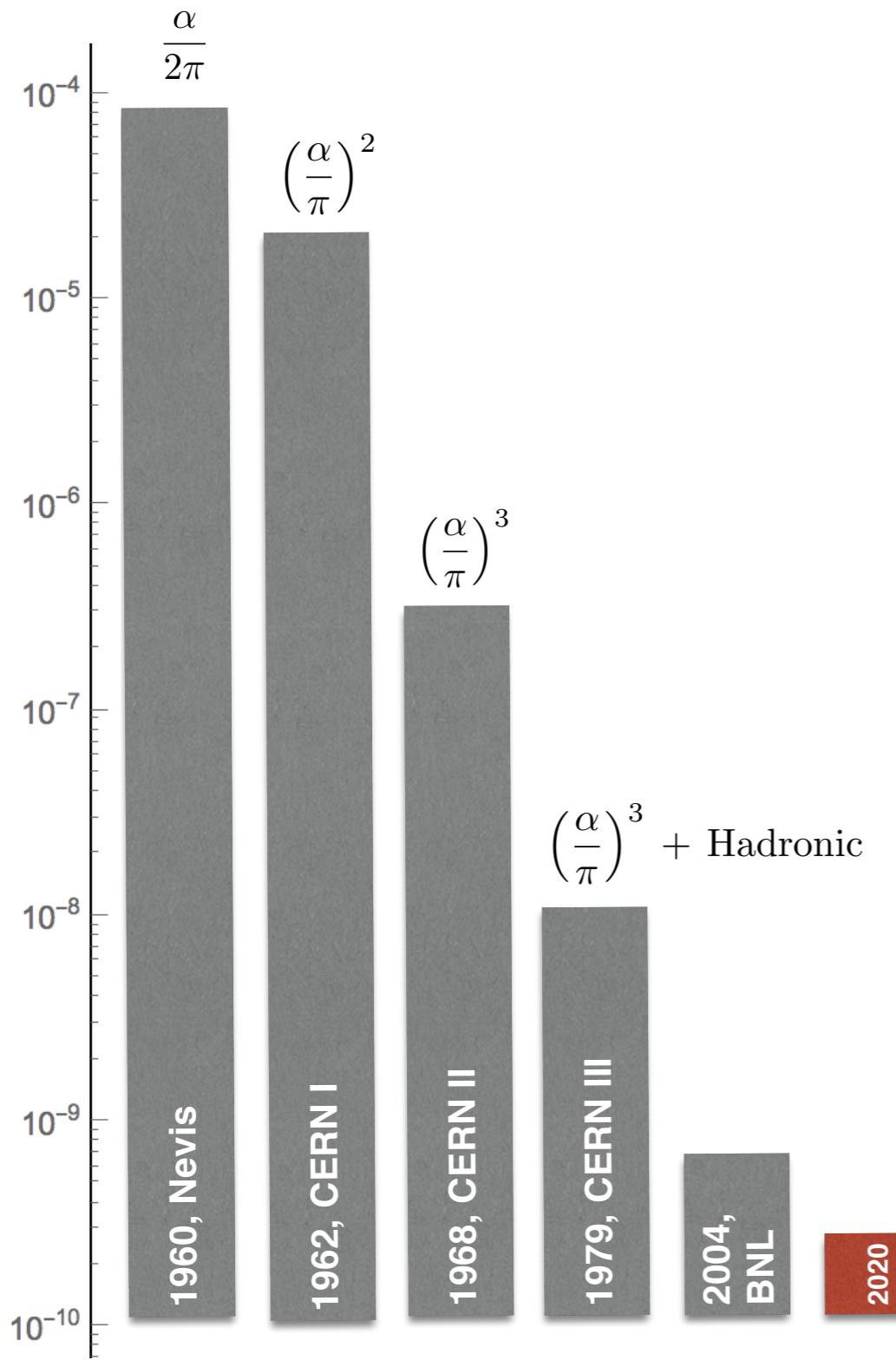


- Anomalous part

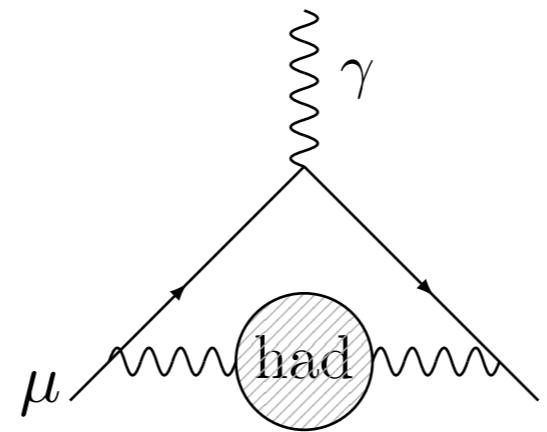
$$a_\mu = \frac{(g - 2)_\mu}{2}$$

History of achieved accuracy

- Accuracy



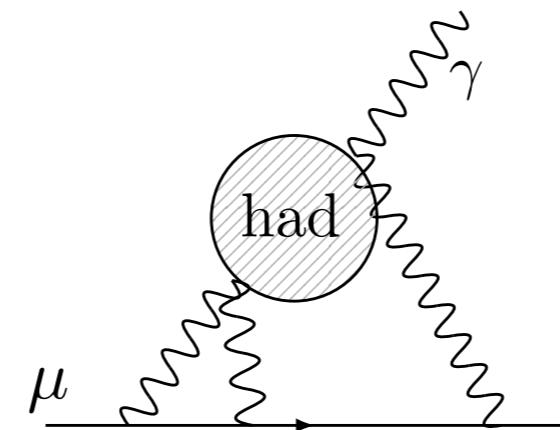
- Hadronic corrections



- hadronic vacuum polarization (LO + NLO)

$$a_\mu^{\text{had, VP}} = 685.1(4.3) \times 10^{-10}$$

Hagiwara et al. (2011)



- hadronic light-by-light scattering

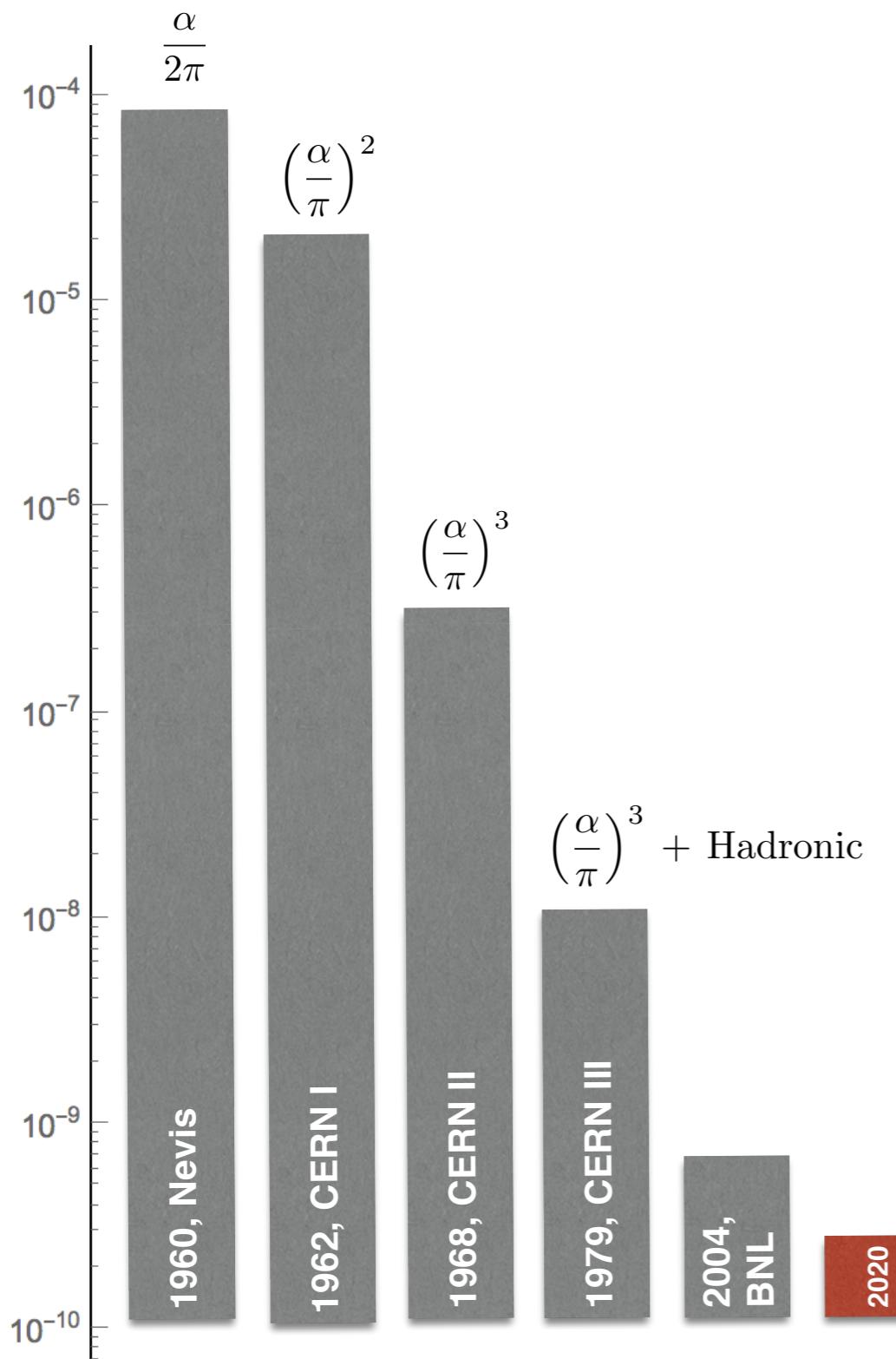
$$a_\mu^{\text{had, LbL}} = 10.2(3.9) \times 10^{-10}$$

Prades et al. (2009)
Jegerlehner, Nyffeler (2009, 2015)

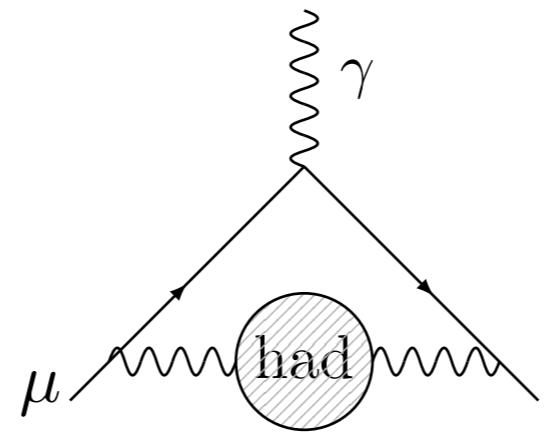
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History of achieved accuracy

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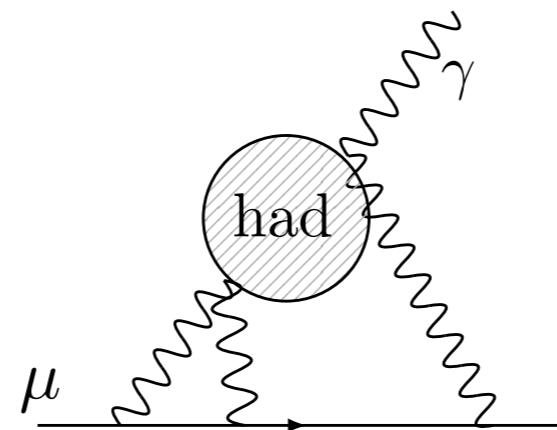
- Hadronic corrections



- hadronic vacuum polarization (LO + NLO)

$$a_\mu^{\text{had, VP}} = 685.1(4.3) \times 10^{-10}$$
$$683.4(2.5) \times 10^{-10}$$

Hagiwara et al. (2011)
Keshavarzi et al. (2018)



- hadronic light-by-light scattering

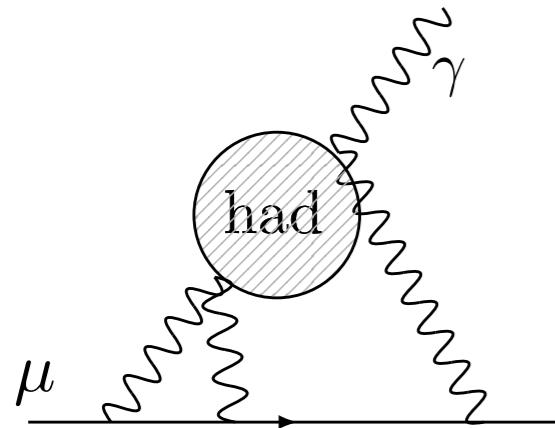
$$a_\mu^{\text{had, LbL}} = 10.2(3.9) \times 10^{-10}$$
$$8.7(1.3) \times 10^{-10}$$

Prades et al. (2009)
Jegerlehner, Nyffeler (2009, 2015)
I.D, Redmer, Vanderhaeghen (2019)

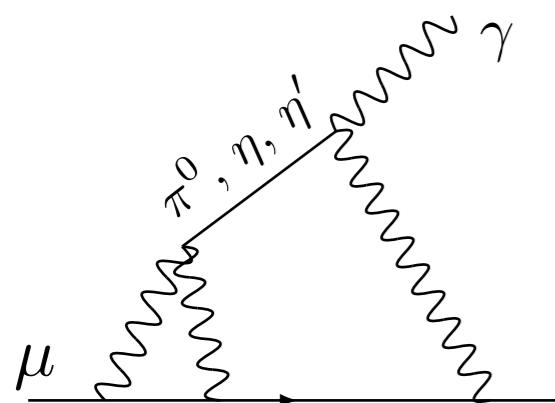
$$\Delta a_\mu^{\text{exp}} = 6.3 \times 10^{-10}$$
$$\rightarrow 1.6 \times 10^{-10}$$

White paper (ongoing)

QCD contribution to (g-2)

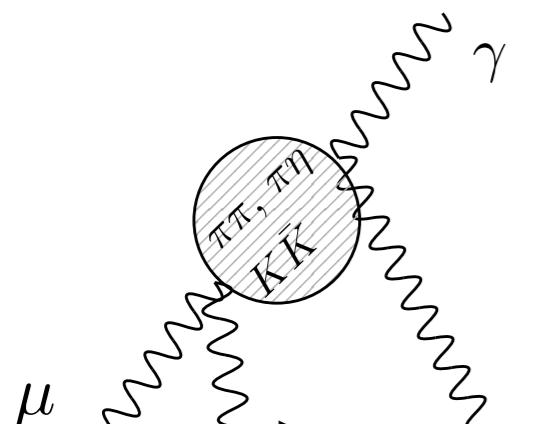


=



+

Relies on measurements of **TFF** $\pi^0\gamma^*\gamma^{(*)}, \eta\gamma^*\gamma^{(*)}, \dots$
to reduce the model dependence

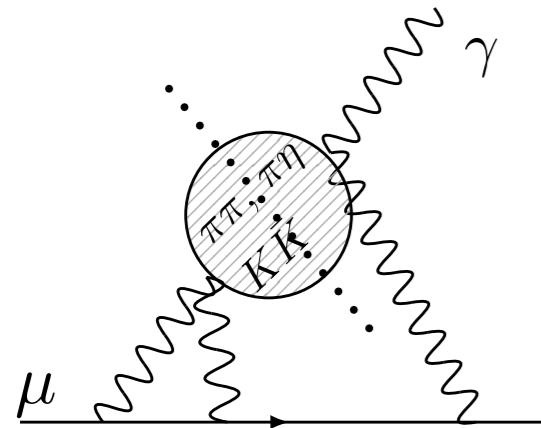


+

Dispersive analysis for $\pi\pi, \pi\eta, \dots$ loops is needed

...

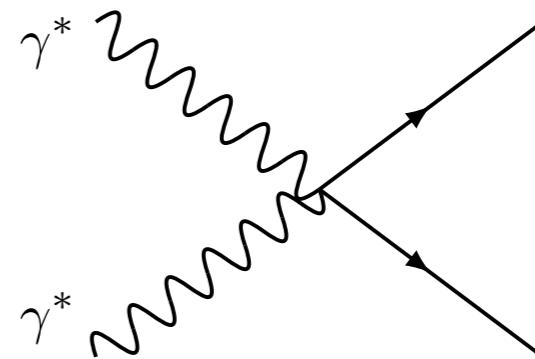
Multi-meson production



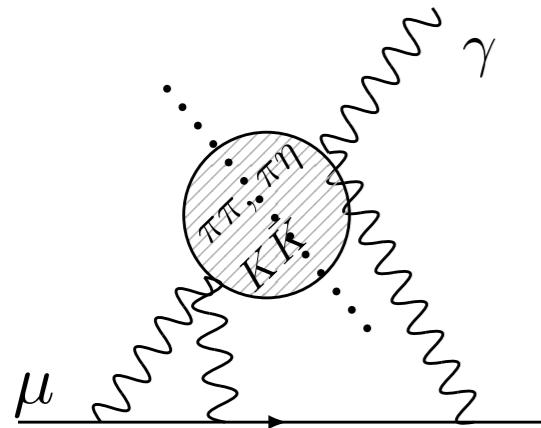
Important ingredient:

$$\gamma^* \gamma^* \rightarrow \pi\pi, \pi\eta, \dots$$

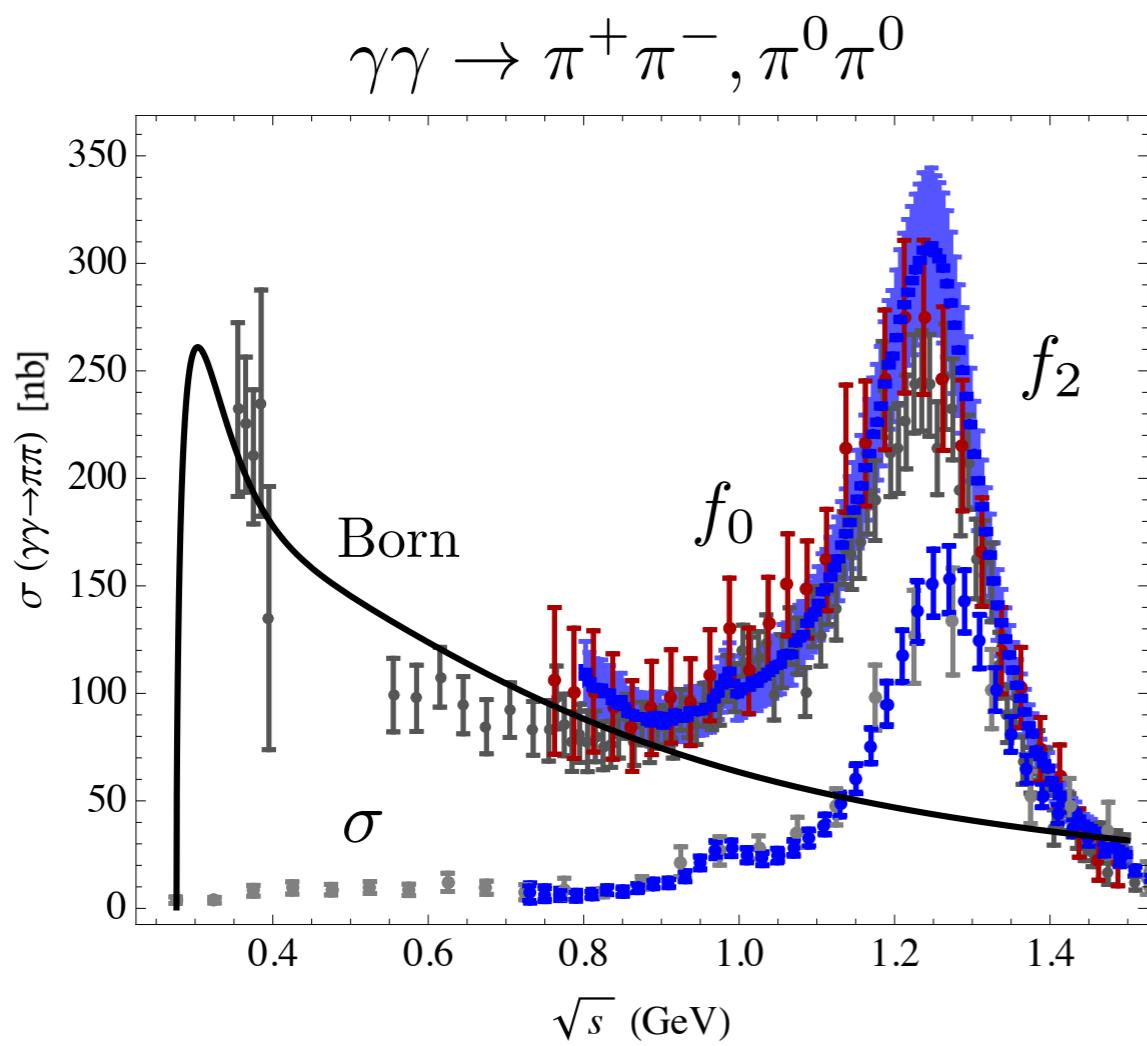
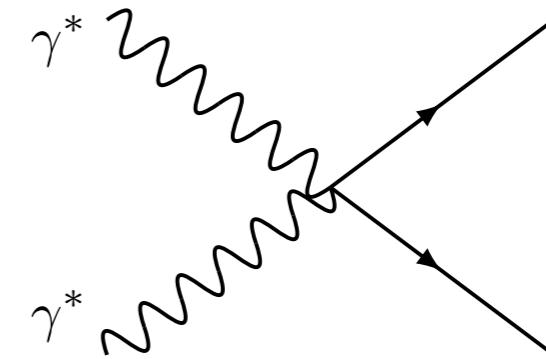
$$q^2 = -Q^2 < 0 \quad \text{space-like } \gamma^*$$



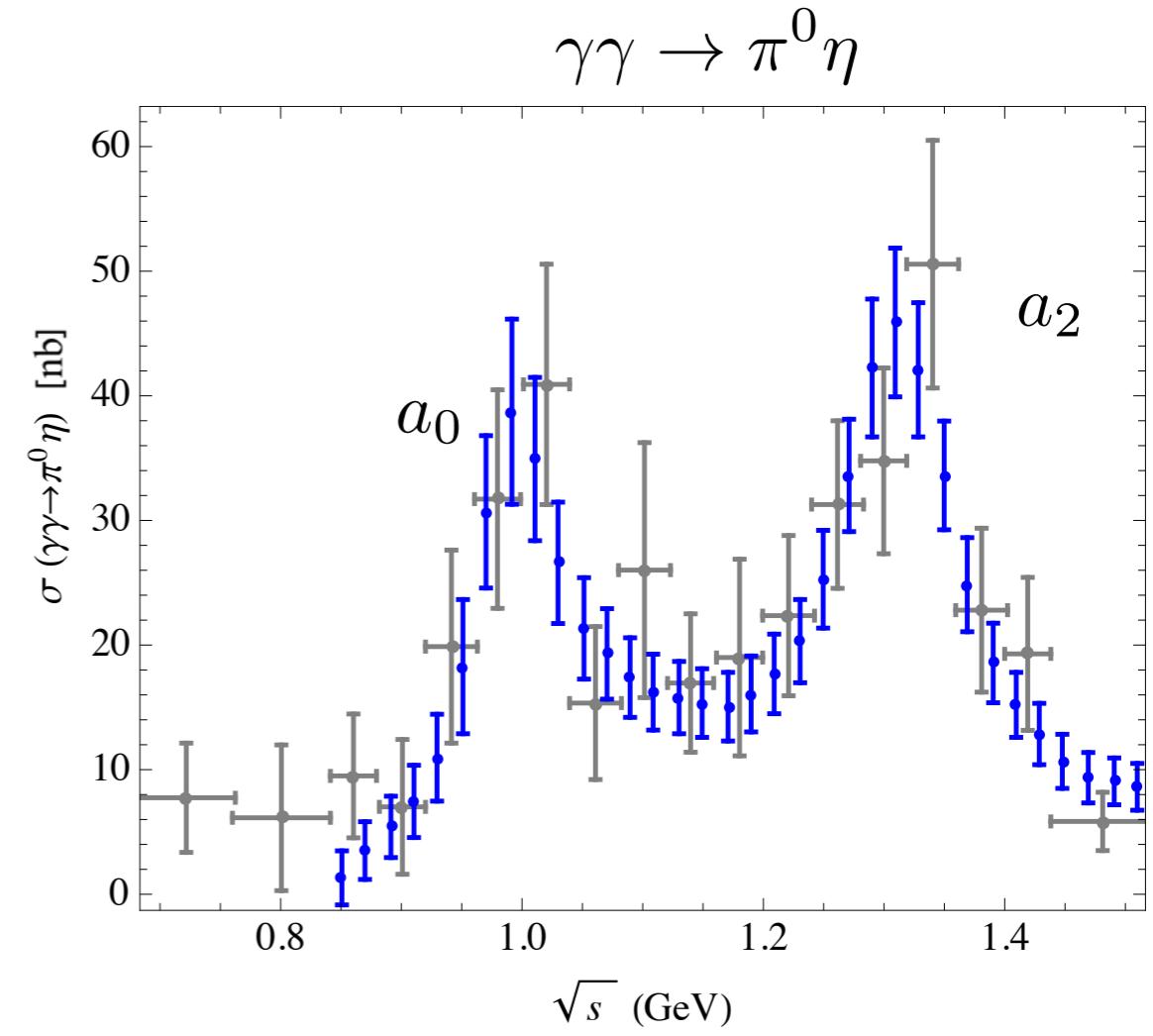
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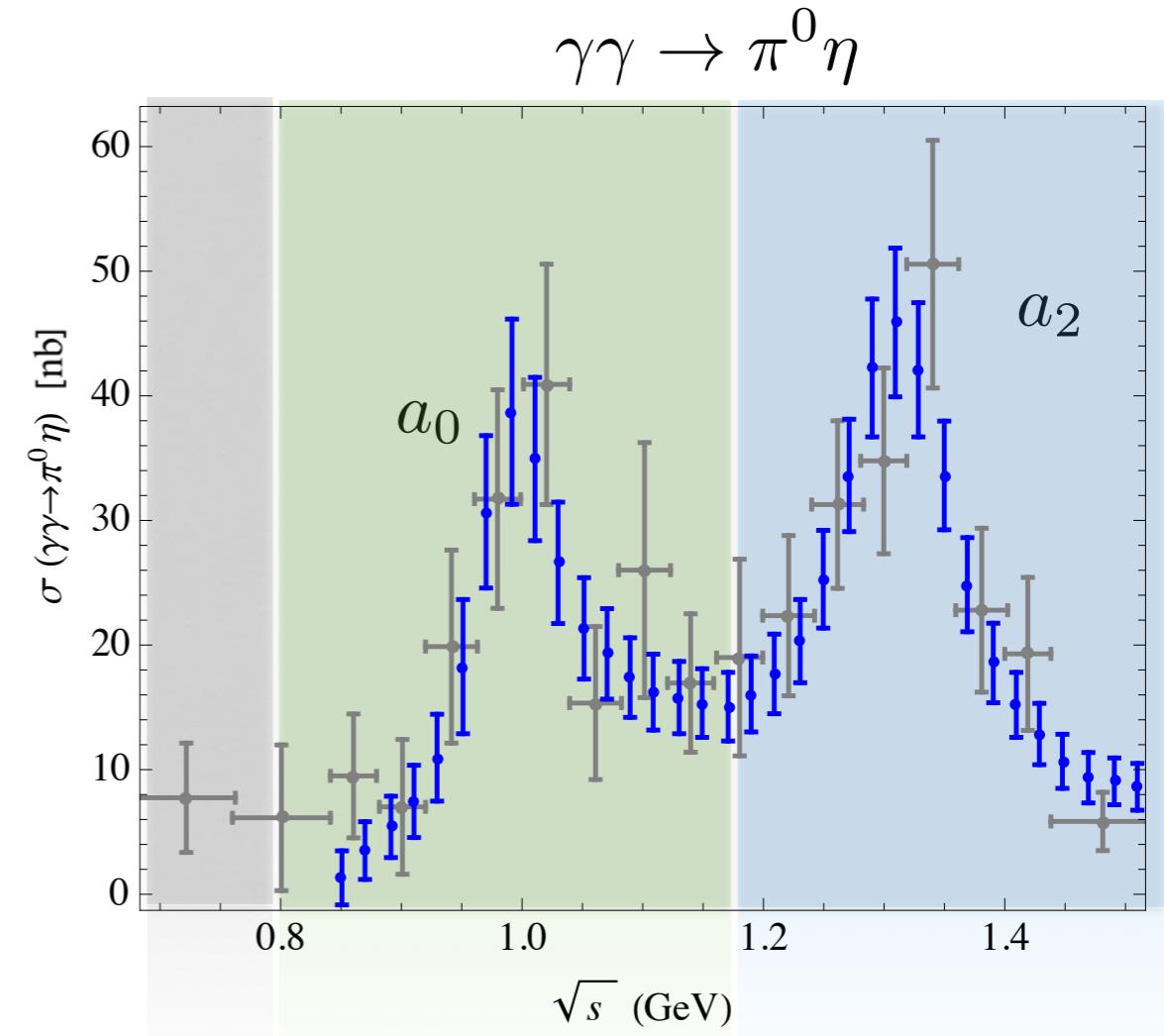
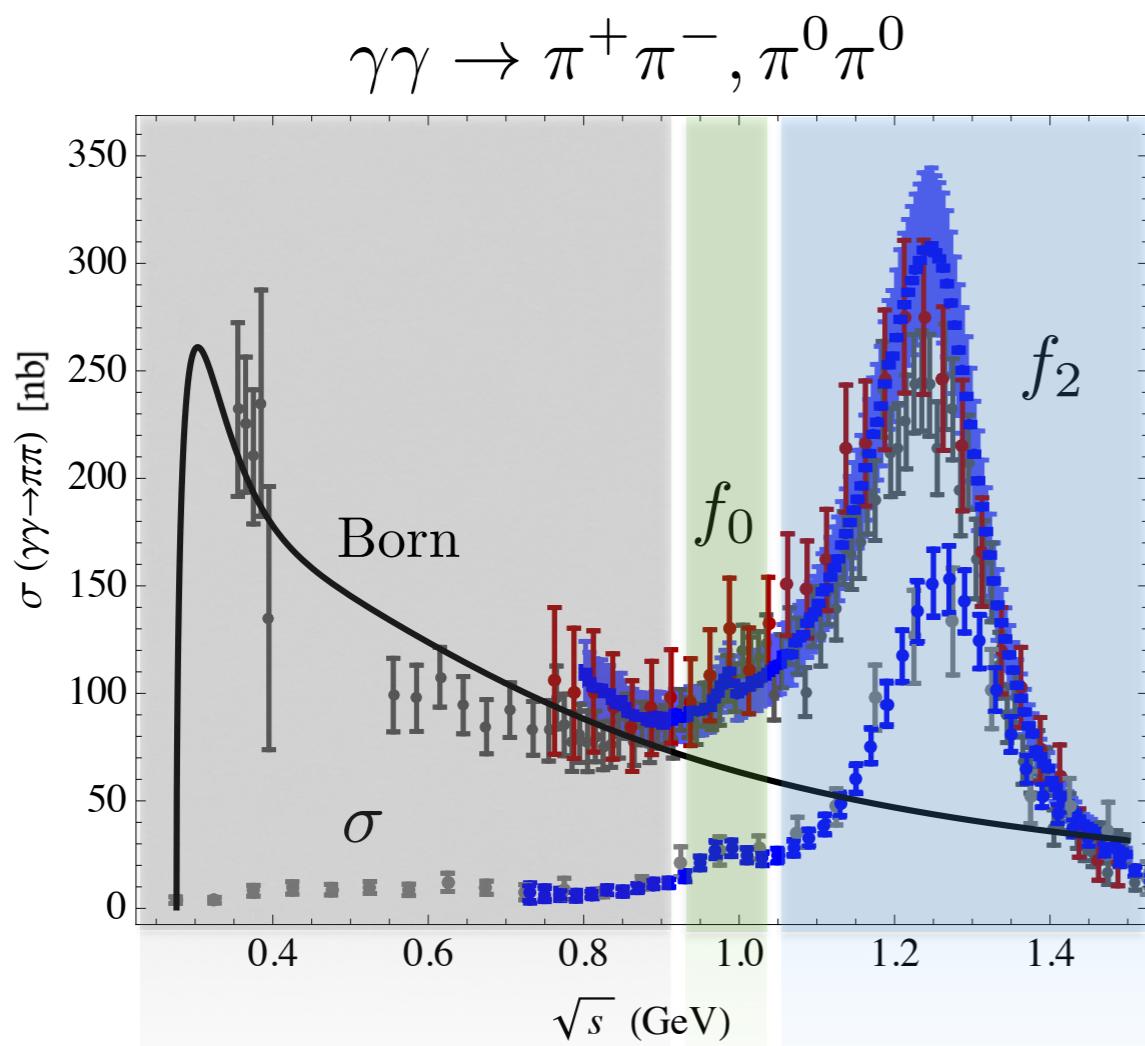
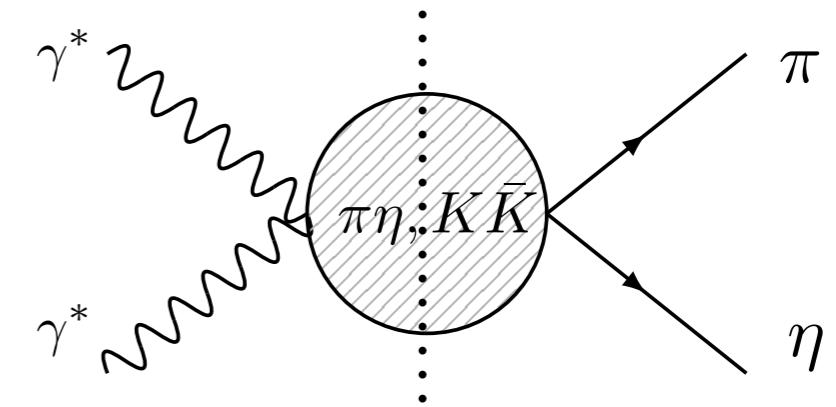
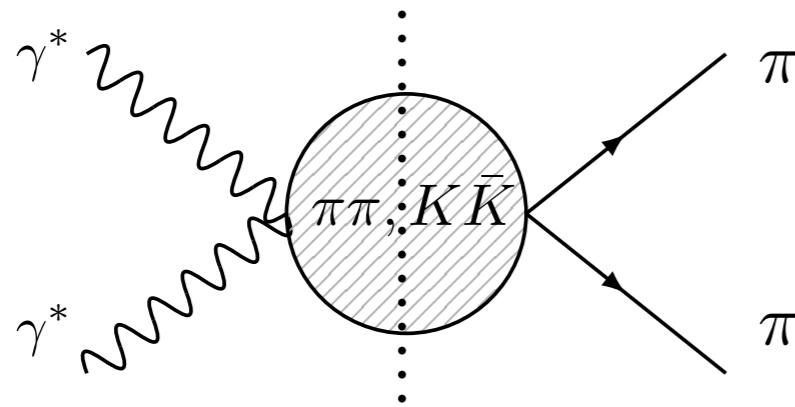


MarkII ('90), CELLO ('92), Crystal Ball ('90), Belle ('07, '09)
 Crystal Ball ('90), Belle ('09)

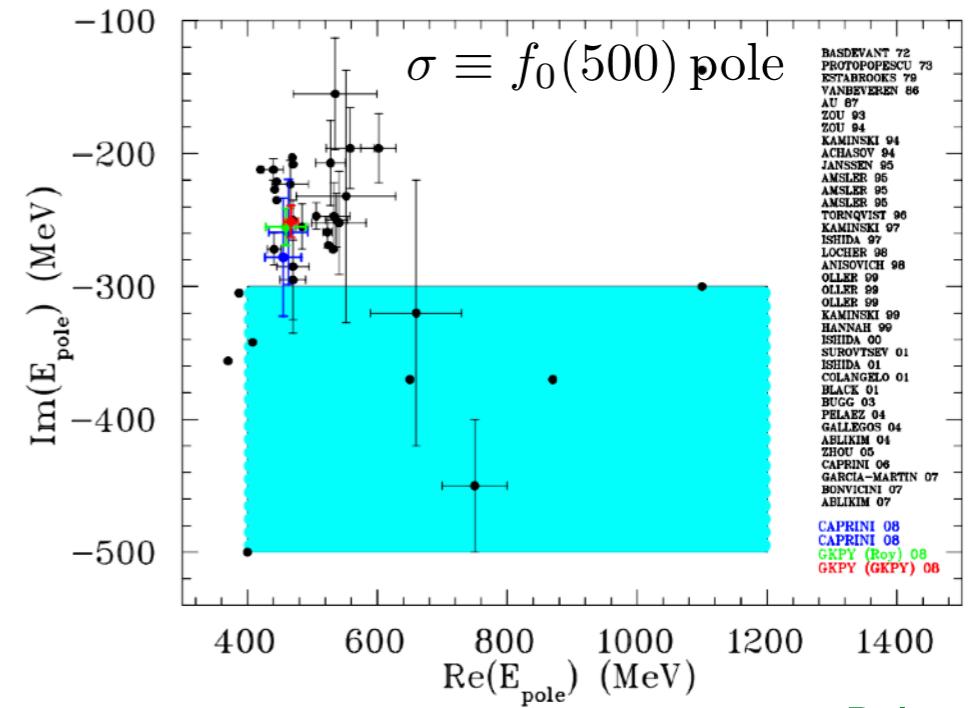
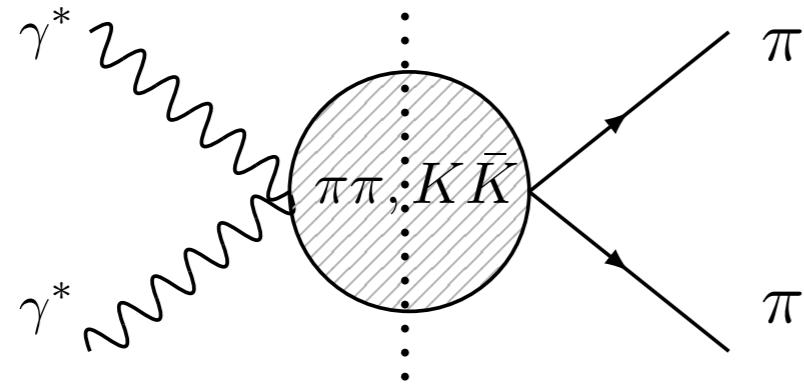


+ Belle (2016) $\gamma\gamma^* \rightarrow \pi^0\pi^0, Q^2 > 3.5$ GeV 2

Inelastic contributions



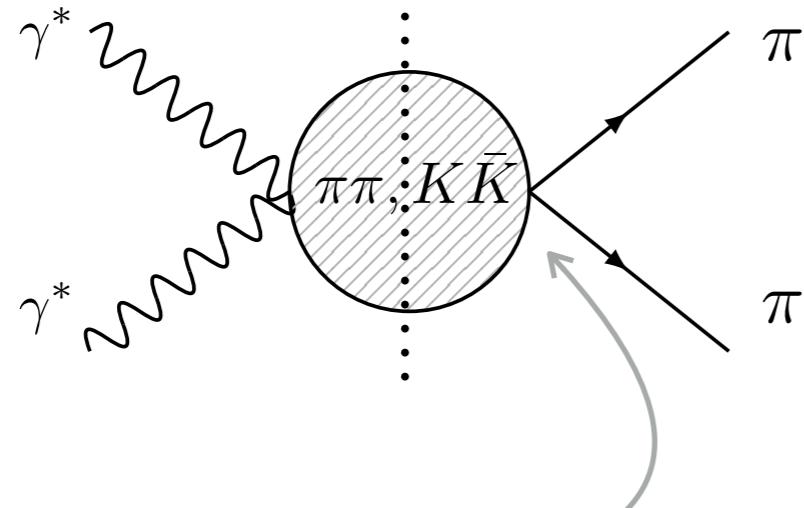
Right-hand cuts (hadronic input)



Pelaez (2016)

$$\sqrt{s_{\text{Roy}}^{\text{II}}} = (449_{-16}^{+22}) \pm i(275 \pm 12) \text{ MeV}$$

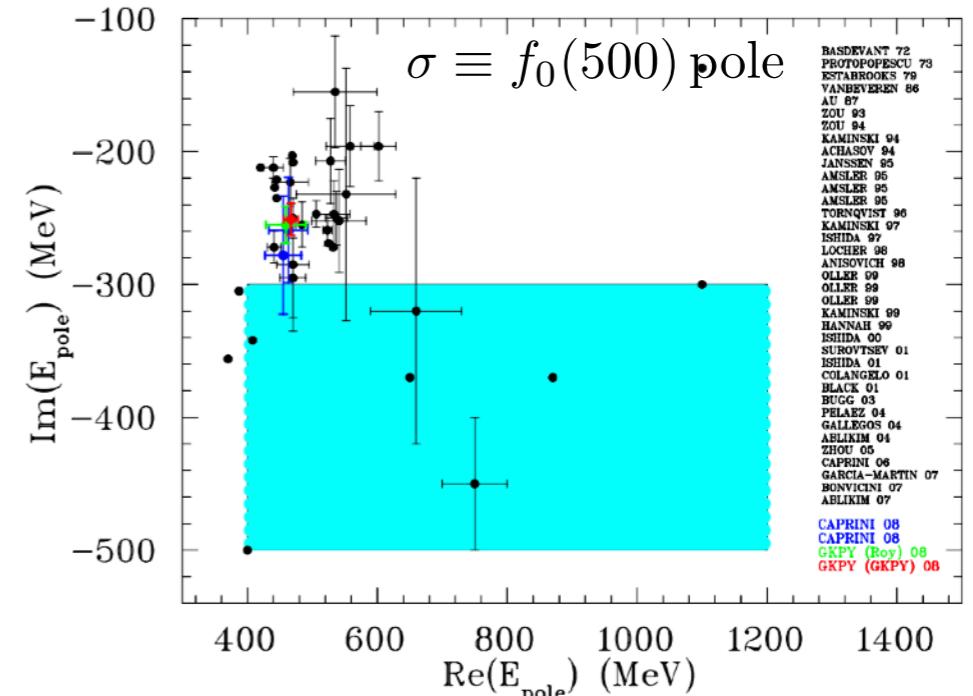
Right-hand cuts (hadronic input)



- Coupled-channel Omnès formalism

$$t(s) = U(s) + \int_R \frac{ds'}{\pi} \frac{\rho(s') |t(s')|^2}{s' - s}$$

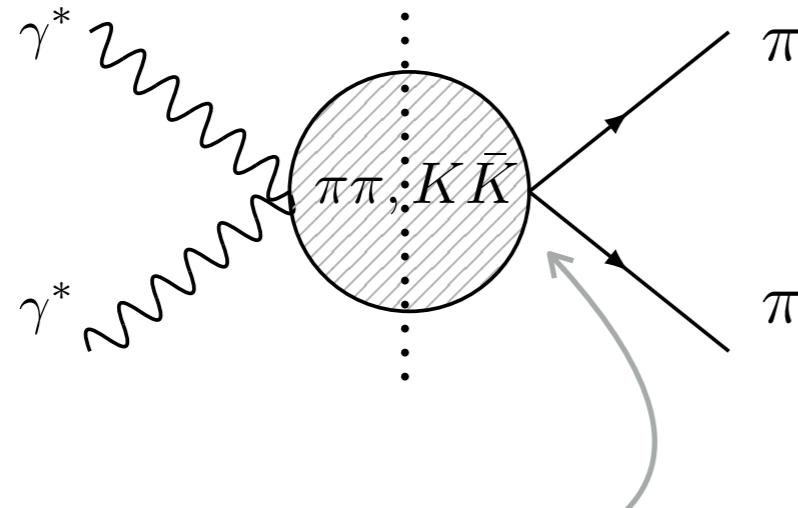
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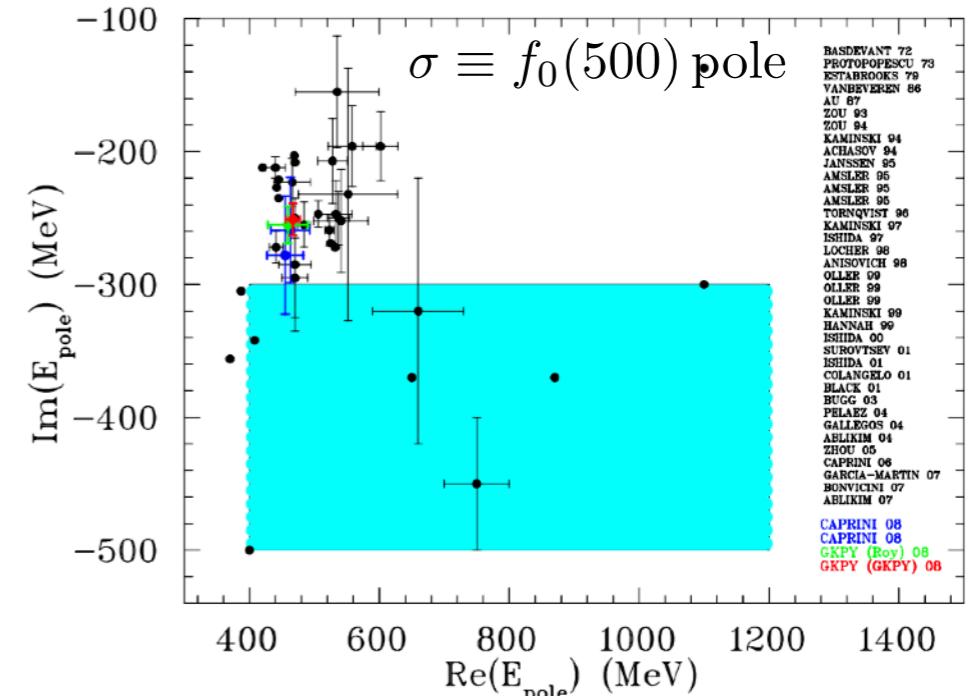
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- Model independent form of the left-hand cuts: conformal mapping expansion

$$U(s) = \sum_k C_k \xi(s)^k$$

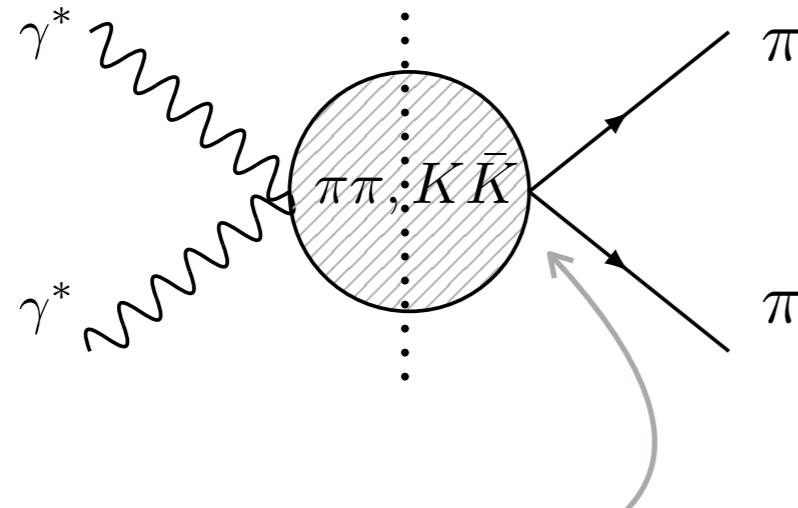
I.D., Lutz, Gasparyan (2011)



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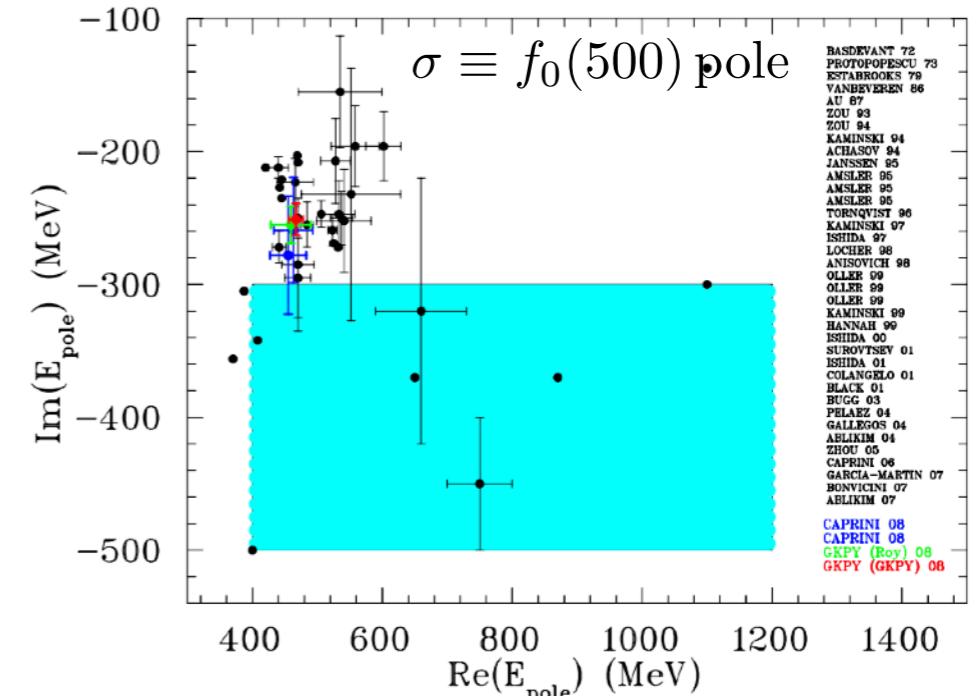
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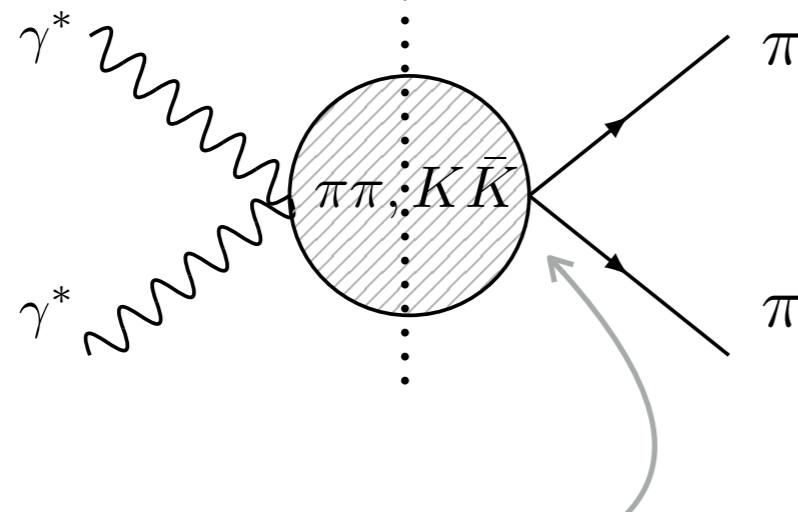
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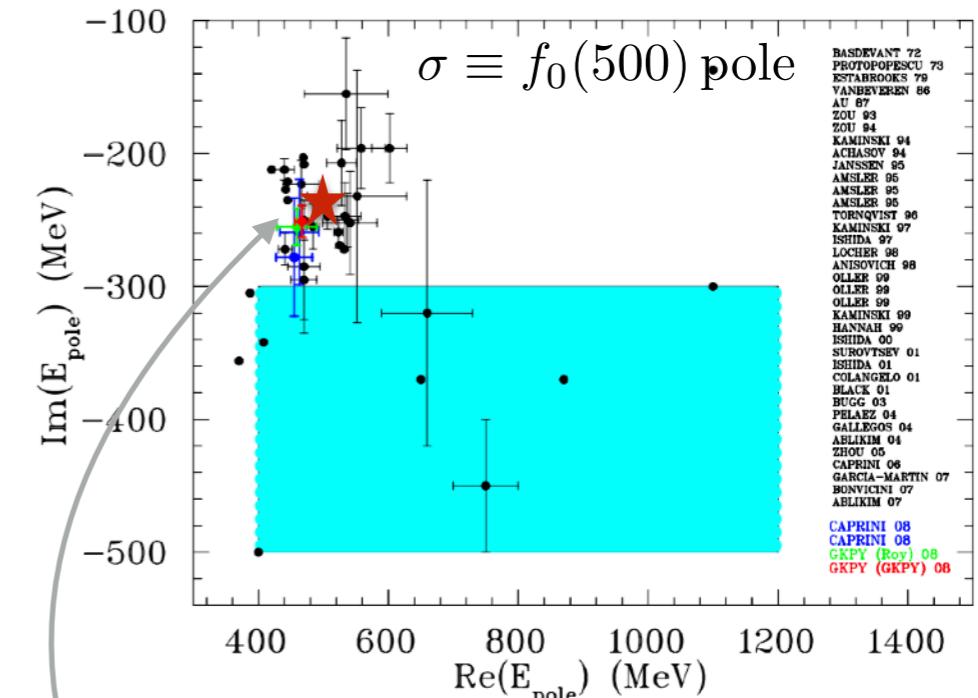
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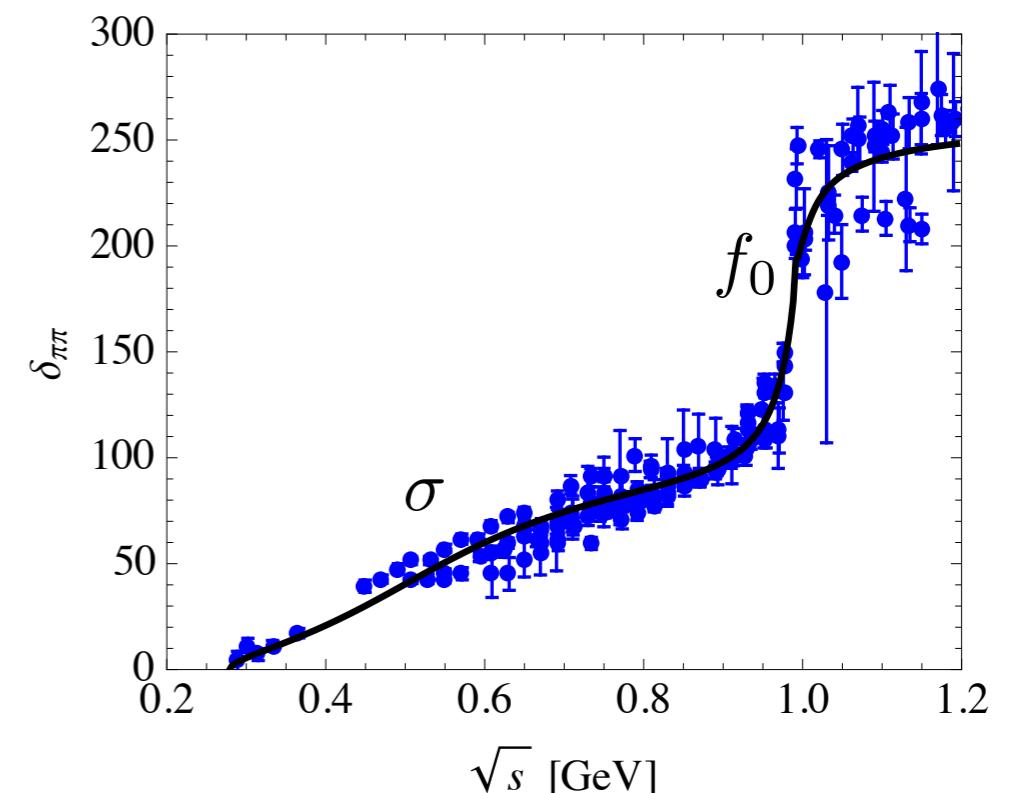
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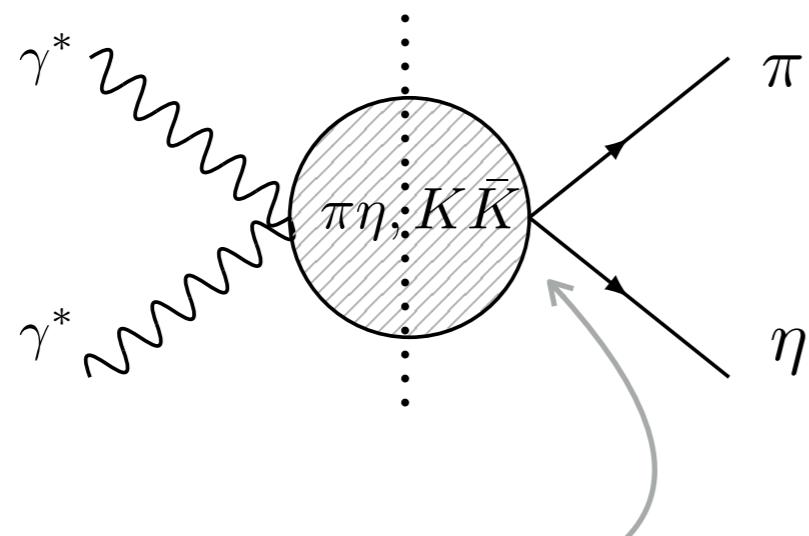
Pelaez (2016)

$$\sqrt{s_{\text{Roy}}^{\text{II}}} = (449_{-16}^{+22}) \pm i(275 \pm 12) \text{ MeV}$$

$$\sqrt{s_{N/D}^{\text{II}}} = (496 \pm 48) \pm i(226 \pm 30) \text{ MeV}$$



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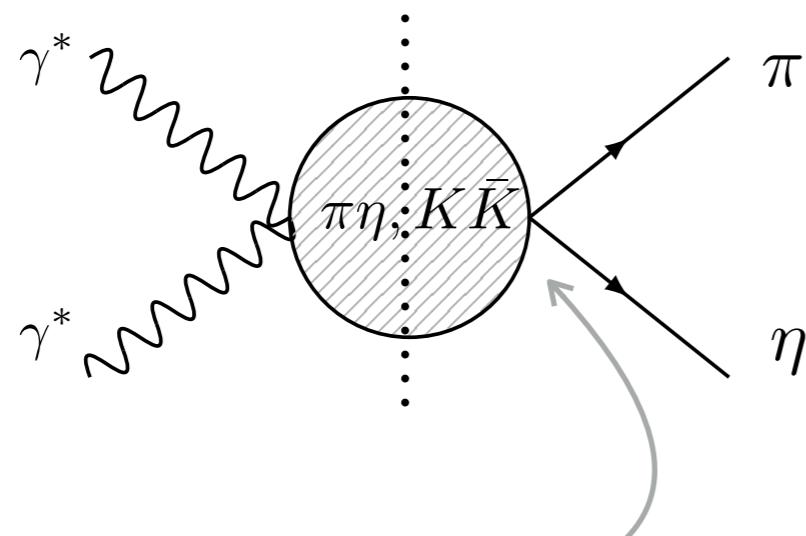
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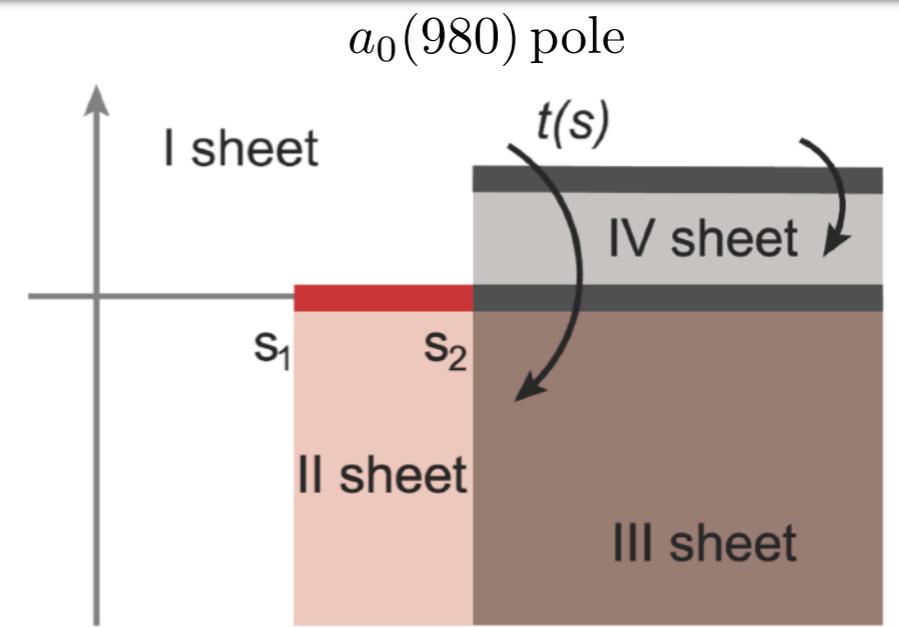
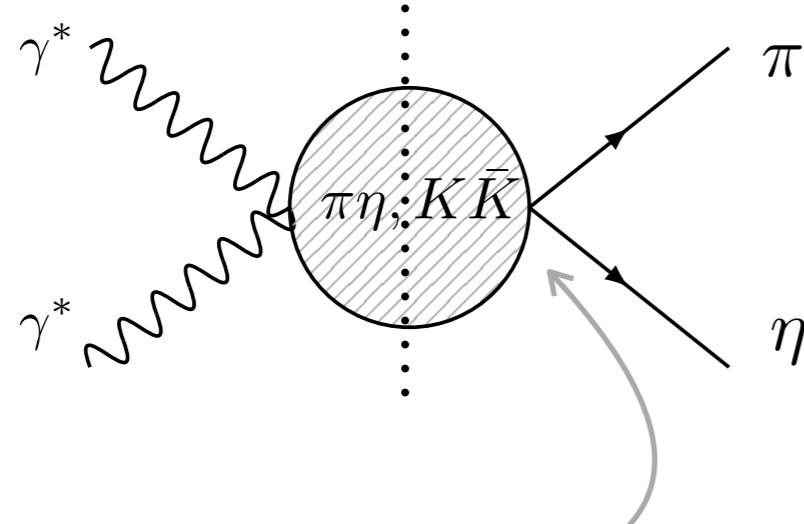
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I.D., Lutz, Gasparyan (2011)

- Coefficients C_k matched to SU(3)
Chiral Perturbation Theory at threshold

I.D., Gil, Lutz (2011, 2013)

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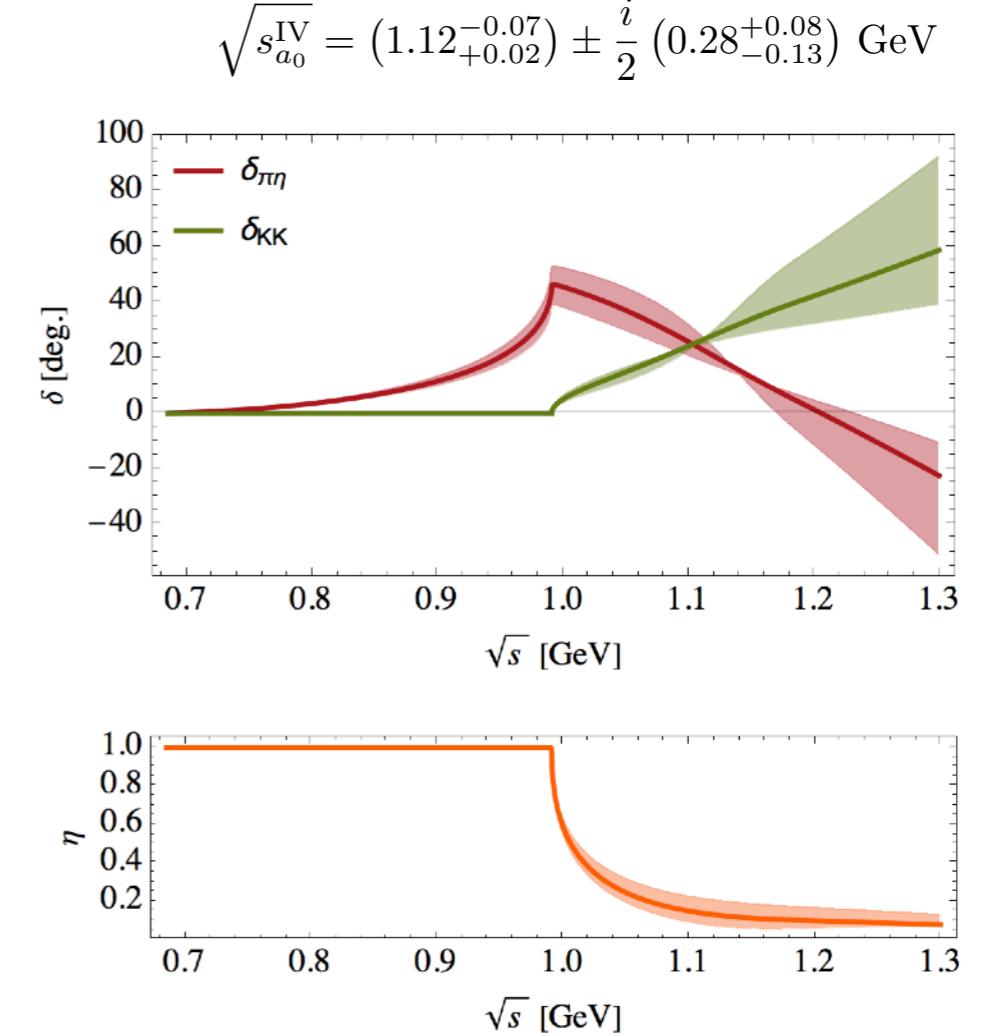
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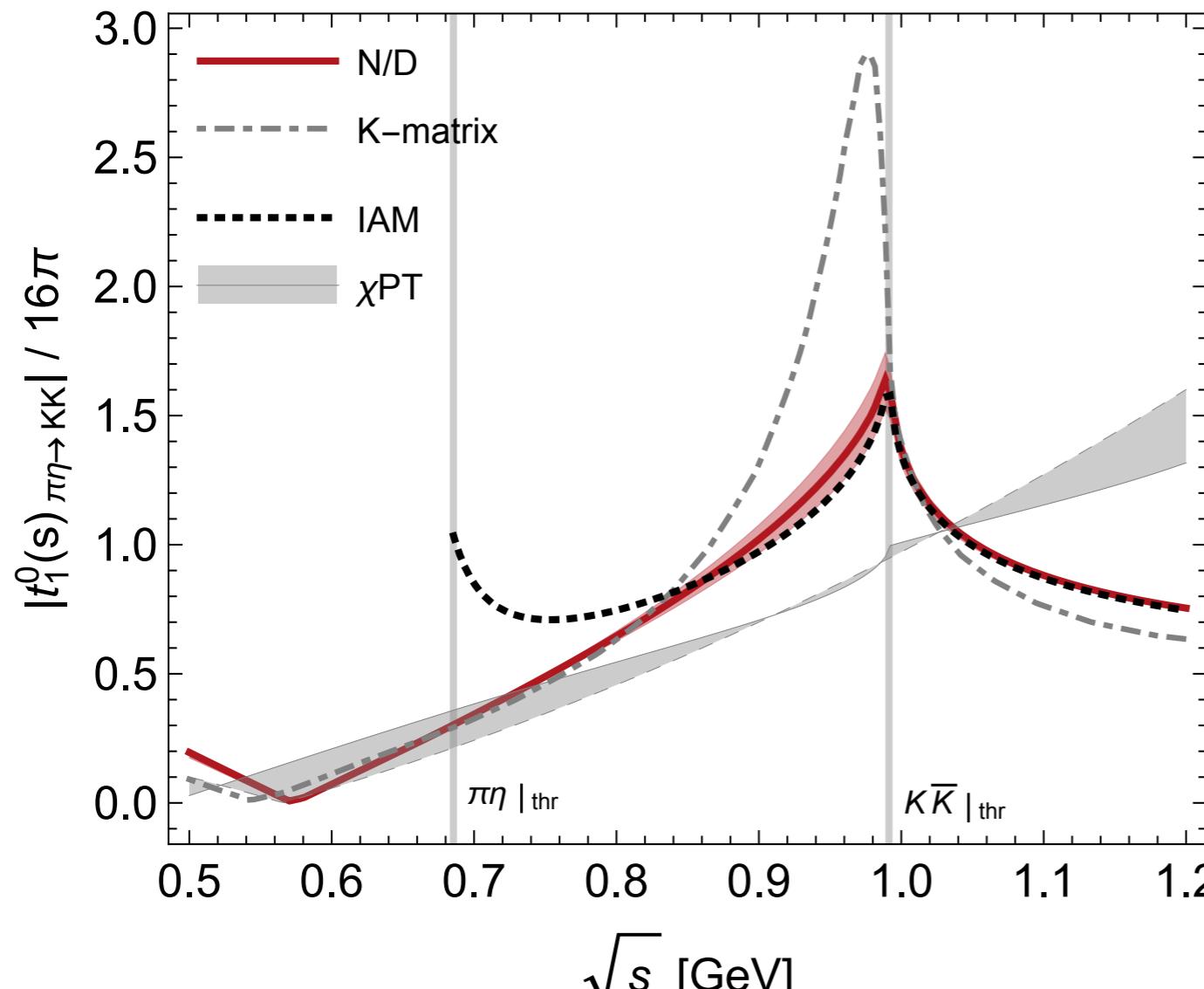
- Coefficients C_k matched to SU(3) Chiral Perturbation Theory at threshold

I.D., Gil, Lutz (2011, 2013)



cf. also HadSpec Coll. (2016)

Scattering amplitude $\pi\eta \rightarrow K\bar{K}$



N/D

I.D., Gil, Lutz (2013), I.D., Deineka, Vanderhaeghen (2017)

K-matrix

Albaladejo et. al. (2017)

Inverse Amplitude Method (IAM)

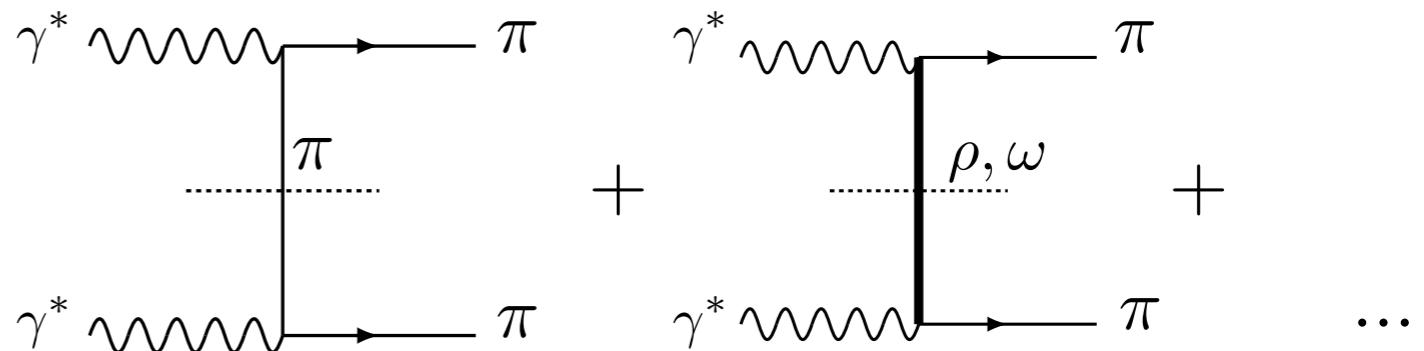
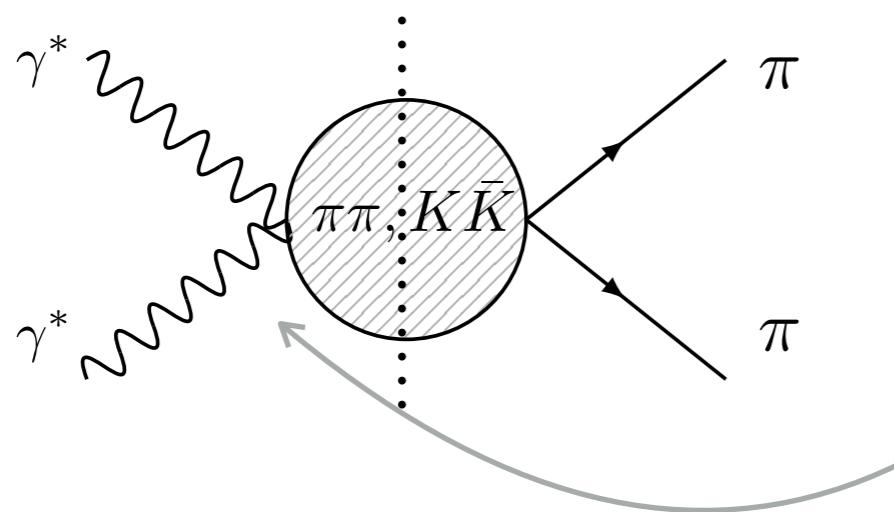
Gomez Nicola et.al. (2002)

Chiral Perturbation Theory

Gasser et. al. (1985)

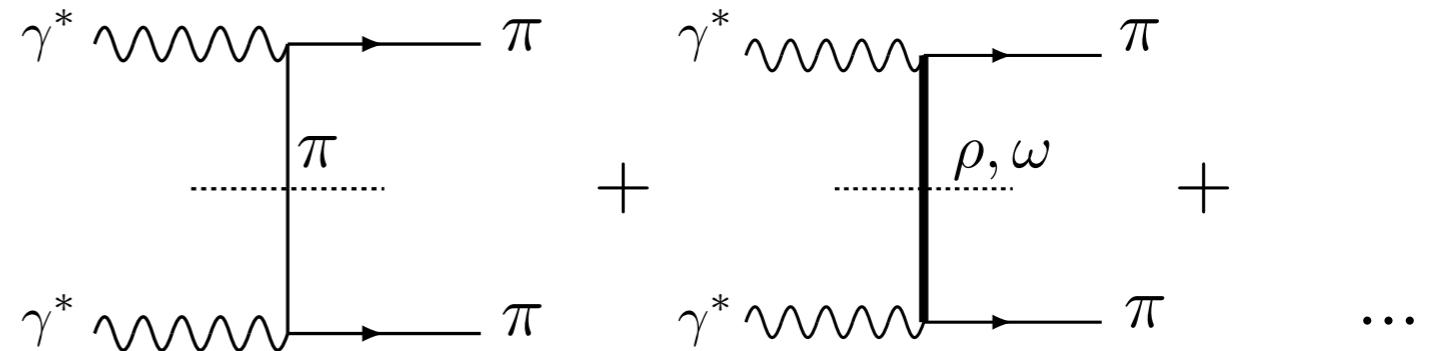
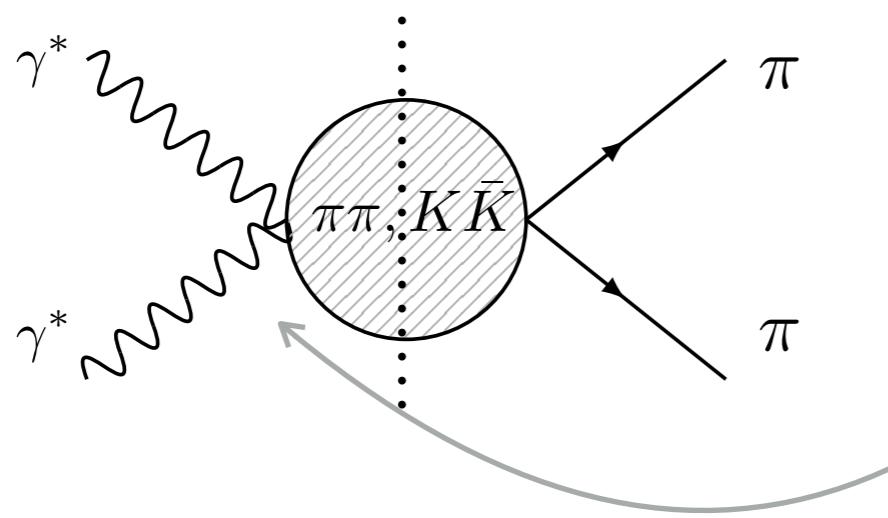
- First lattice analysis for $m_\pi=391$ MeV HadSpec Coll. (2016)
- Chiral extrapolation of the lattice results Zhi-Hui Guo et. al. (2017)

Left-hand cuts (pion pole)

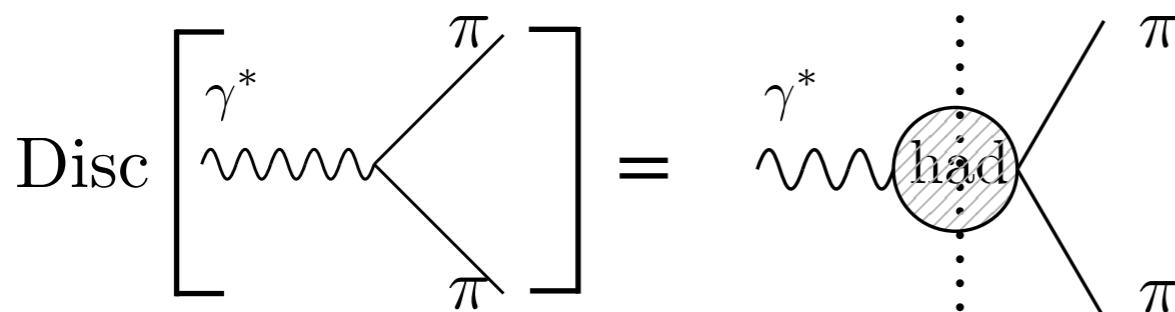


- Left-hand cuts requires knowledge from $\gamma^*\pi\pi$, $\gamma^*\pi\omega$, $\gamma^*\pi\rho$ transition form factors

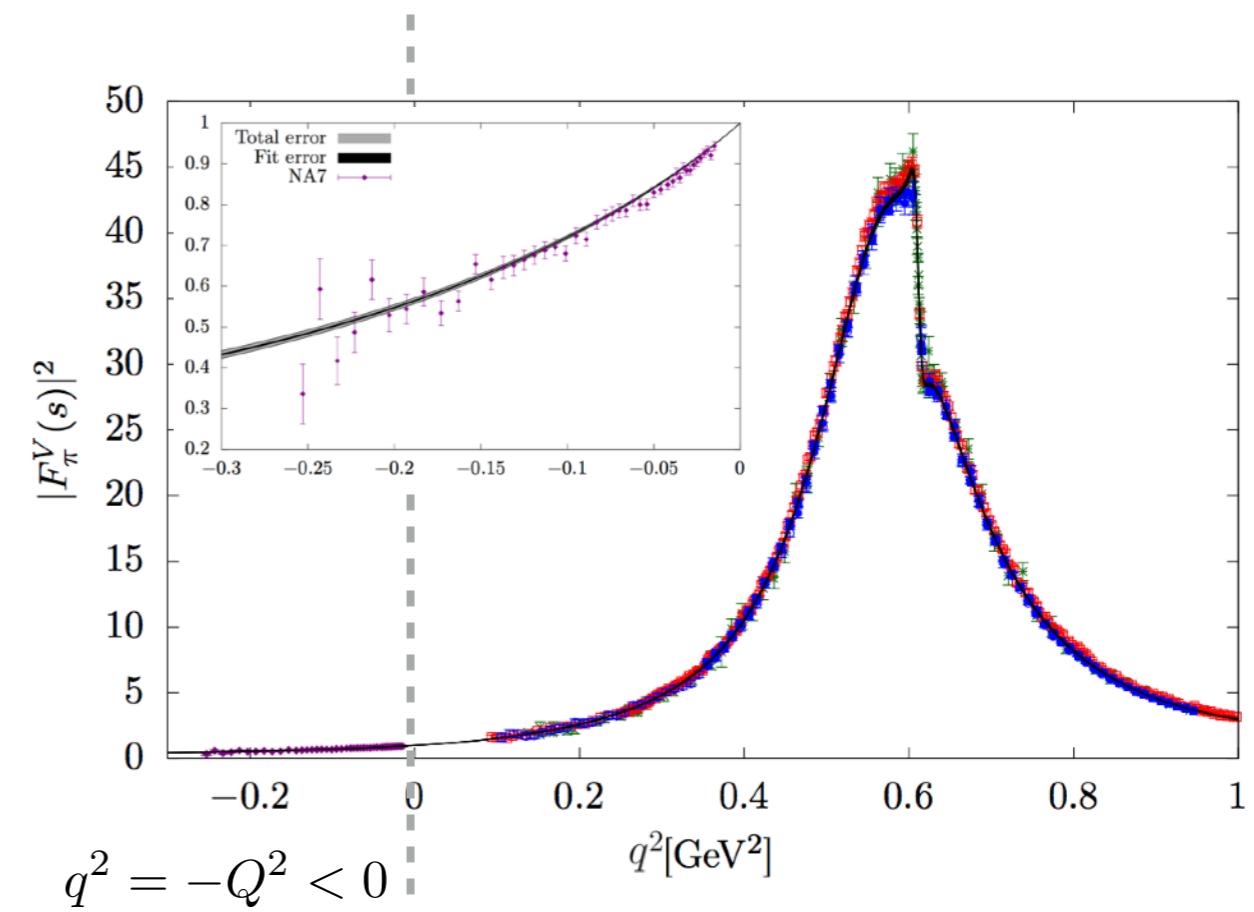
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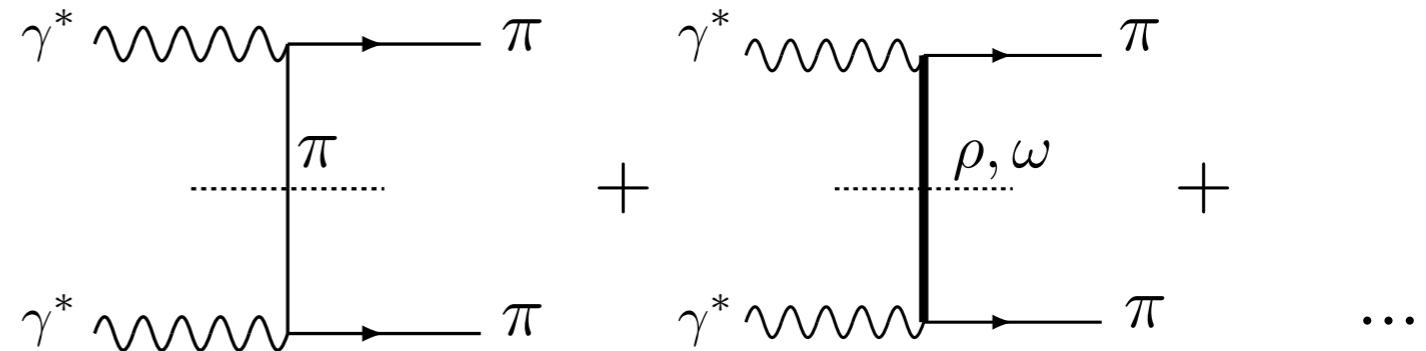
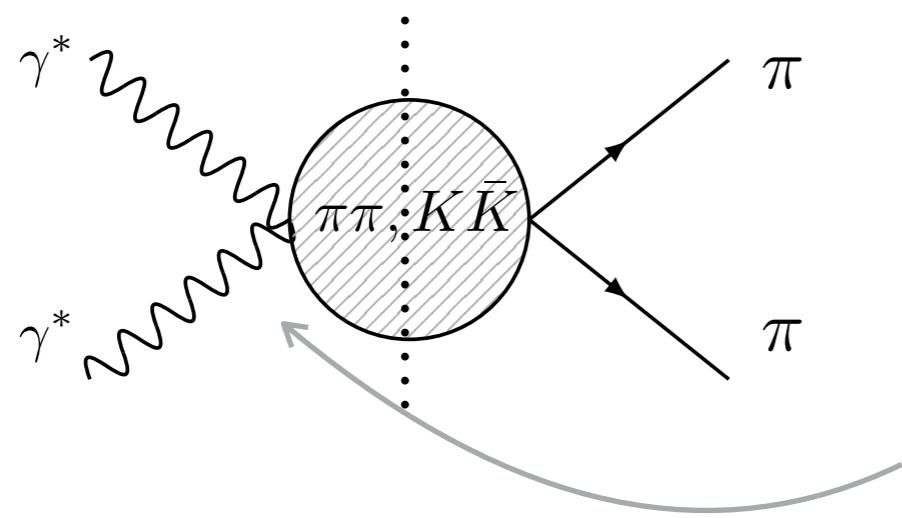


- Input from data: monopole TFF works well for space like region



Colangelo et al. (2019)

Left-hand cuts (vector poles)

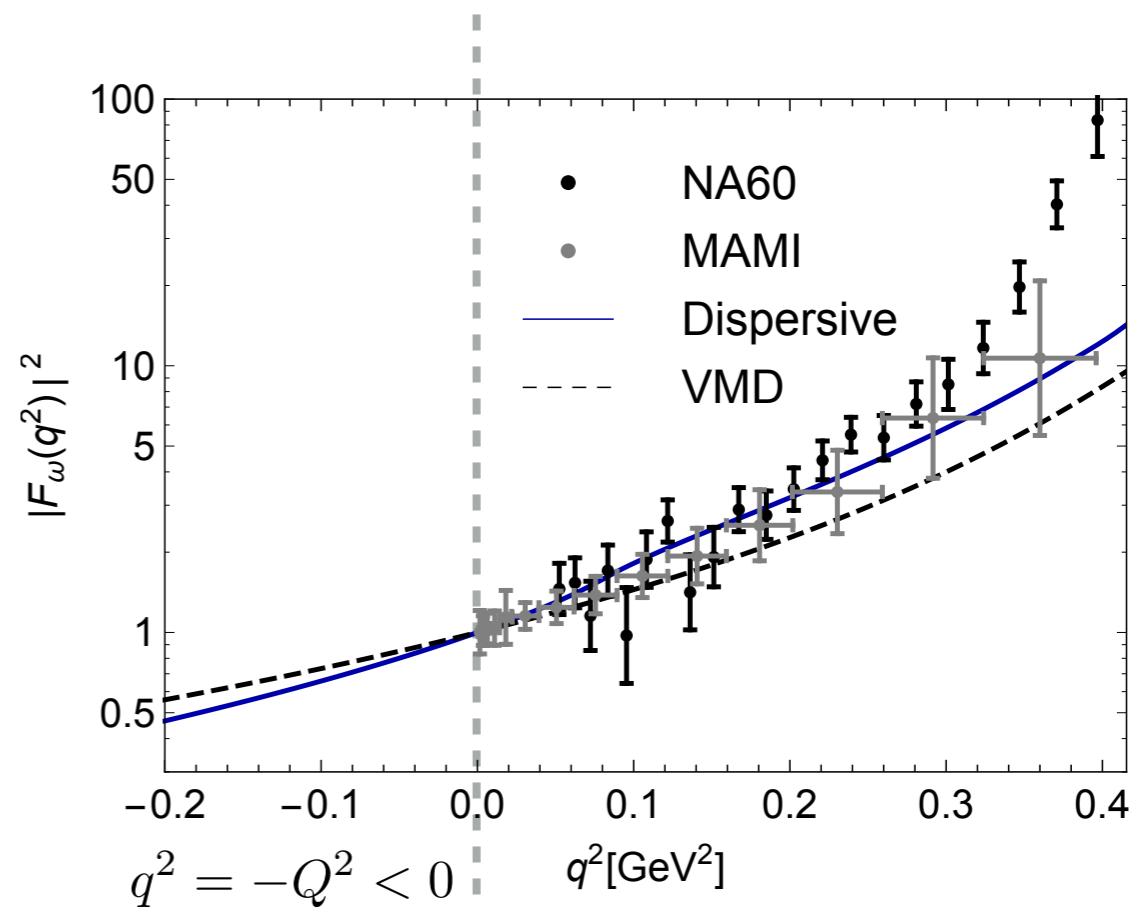


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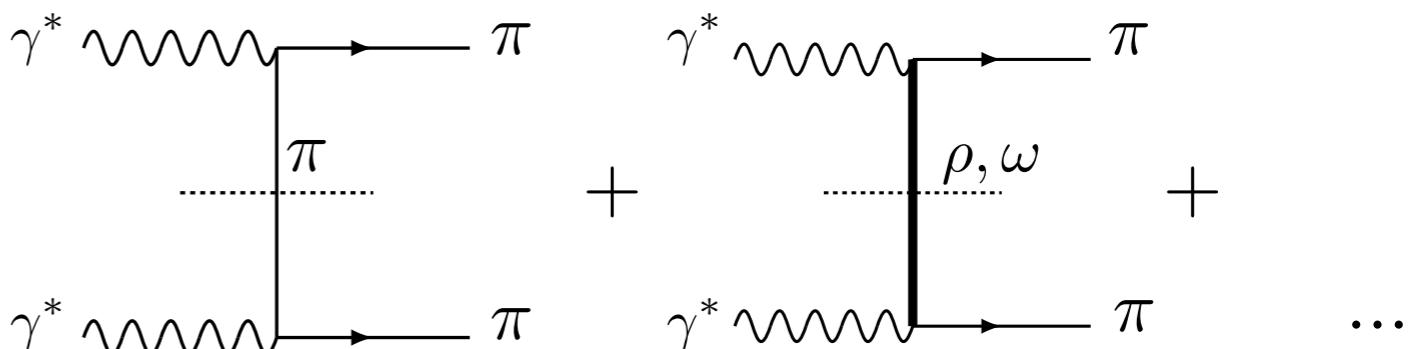
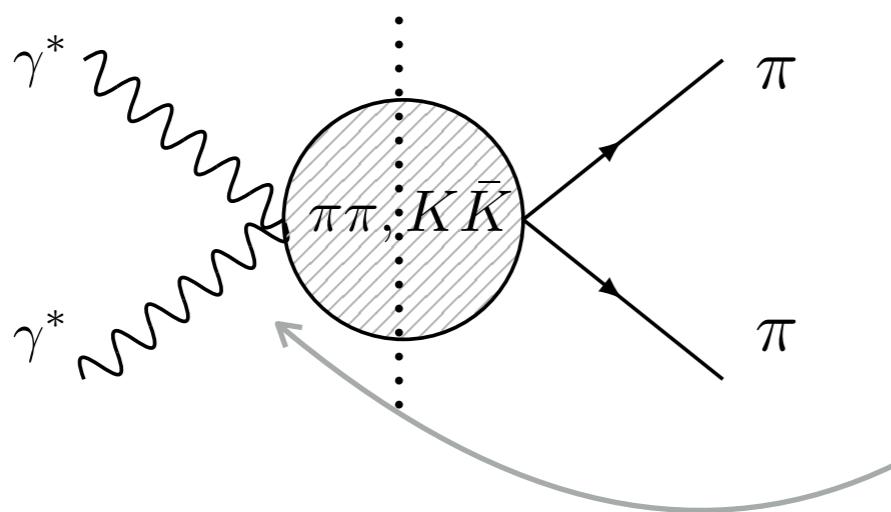
$$\text{Disc} \left[\begin{array}{c} \gamma^* \\ \hbox{\wavy line} \\ \gamma^* \end{array} \right] = \gamma^* \rightarrow \hbox{had} \rightarrow \omega \rightarrow \pi$$

- Input from: $\omega \rightarrow 3\pi$

Khuri-Treiman (1960)
Schneider et al. (2012)
I.D. & JPAC (2015)



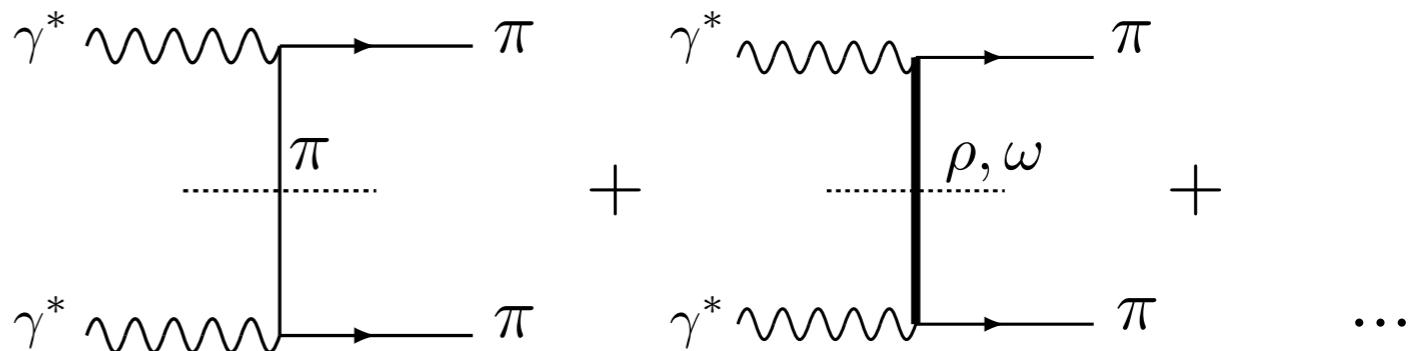
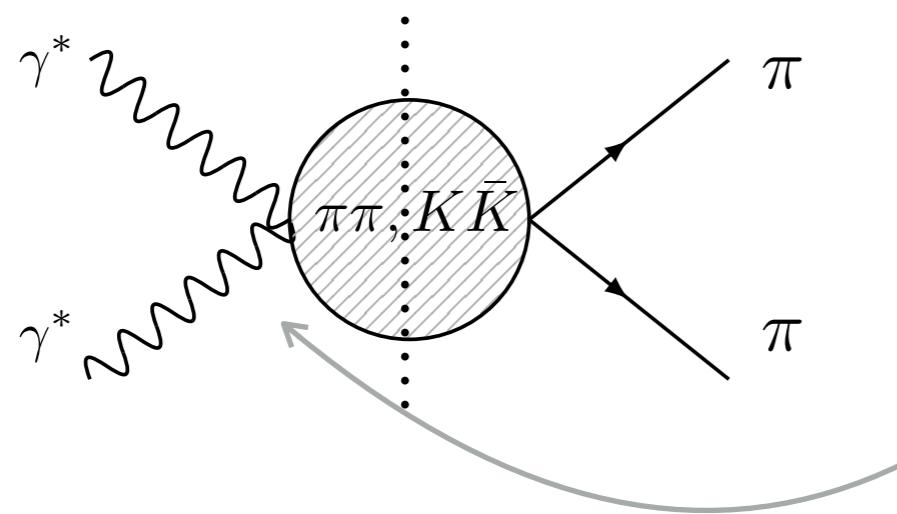
Left-hand cuts (vector poles)



- Left-hand cuts requires knowledge from $\gamma^*\pi\pi, \gamma^*\pi\omega, \gamma^*\pi\rho$ transition form factors

- Fitted parameter is the coupling: $g_{V \rightarrow \pi\gamma} \simeq C_{\rho^{\pm,0} \rightarrow \pi^{\pm,0}\gamma} \simeq \frac{1}{3} C_{\omega \rightarrow \pi^0\gamma} \stackrel{\text{PDG}}{\equiv} 0.37(2) \text{ GeV}^{-1}$
 $g_{V \rightarrow \pi\gamma} = 0.33 \text{ GeV}^{-1}$ PDG (2018)

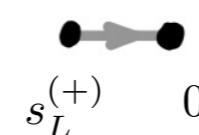
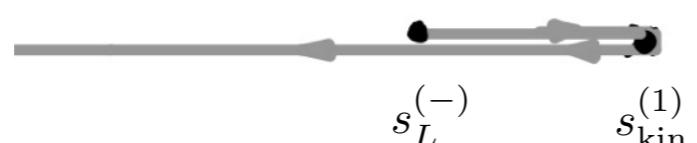
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 $g_{V \rightarrow \pi\gamma} = 0.33 \text{ GeV}^{-1}$ PDG (2018)

- Left-hand cuts: “anomalous thresholds” for large virtualities $Q_1^2 Q_2^2 > (M_V^2 - m_\pi^2)^2$



Hoferichter, Stoffer (2019)
Mandelstam (1960)

Kinematic constraints

- Helicity amplitudes

$$H_{\lambda_1, \lambda_2} = \epsilon_\mu(\lambda_1) \epsilon_\nu(\lambda_2) \sum_{i=1}^5 F_i(s, t) L_i^{\mu\nu}$$

Tarrach (1975)
Drechsel, Metz et al (1998)

where $\lambda_{1,2} = \pm 1, 0$ are photon helicities
(minimal basis for Born subtracted amplitudes)

Low et al. (1954)

- p.w. helicity amplitudes suffer from kinematic constraints

$$h_{\lambda_1 \lambda_2}^{(J)} = \int \frac{d \cos \theta}{2} d_{\lambda_1 - \lambda_2, 0}^J(\theta) H_{\lambda_1 \lambda_2}$$

$$A_n^{(J)} = \frac{1}{(p q)^J} \int \frac{d \cos \theta}{2} P_J(\theta) F_n(s, t) \quad \leftarrow \text{object free of kinematic constraints}$$

Lutz et al. (2010, 2014)

- Unconstrained basis for Born subtracted p.w. amplitudes

$$\bar{h}_i^{(J)} \equiv h_i^{(J)} - h_i^{(J), \text{Born}}$$

$$\bar{h}_i^{(J)} = K_{ij} \bar{h}_j^{(J)} \quad j \equiv \lambda_1 \lambda_2 = \{++, +-, +0, 0+, 00\}$$

K_{ij} is 5×5 matrix

Dispersion relation

- Unsubtracted dispersion relation for kinematically unconstrained p.w. amplitudes

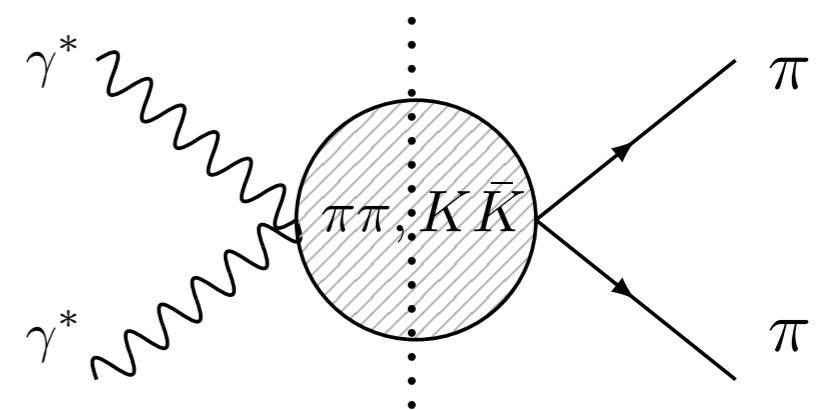
$$\bar{h}_i^{(J)} = \int_{-\infty}^0 \frac{ds'}{\pi} \frac{\text{Disc } \bar{h}_i^{(J)}(s')}{s' - s} + \int_{4m_\pi^2}^\infty \frac{ds'}{\pi} \frac{\text{Disc } \bar{h}_i^{(J)}(s')}{s' - s}$$

Garcia-Martin et. al (2010)
Hoferichter et. al. (2011,19)
Dai et al. (2014)
Moussallam (2013)

- Omnès solution of the unitarity relation

$$\text{Disc } h_i^{(J)} = h_i^{(J)} \rho t_{\pi\pi}^{(J)*}$$

$$\text{Disc } \Omega^{(J)} = \Omega^{(J)} \rho t_{\pi\pi}^{(J)*} \quad |_{s > 4m_\pi^2}$$

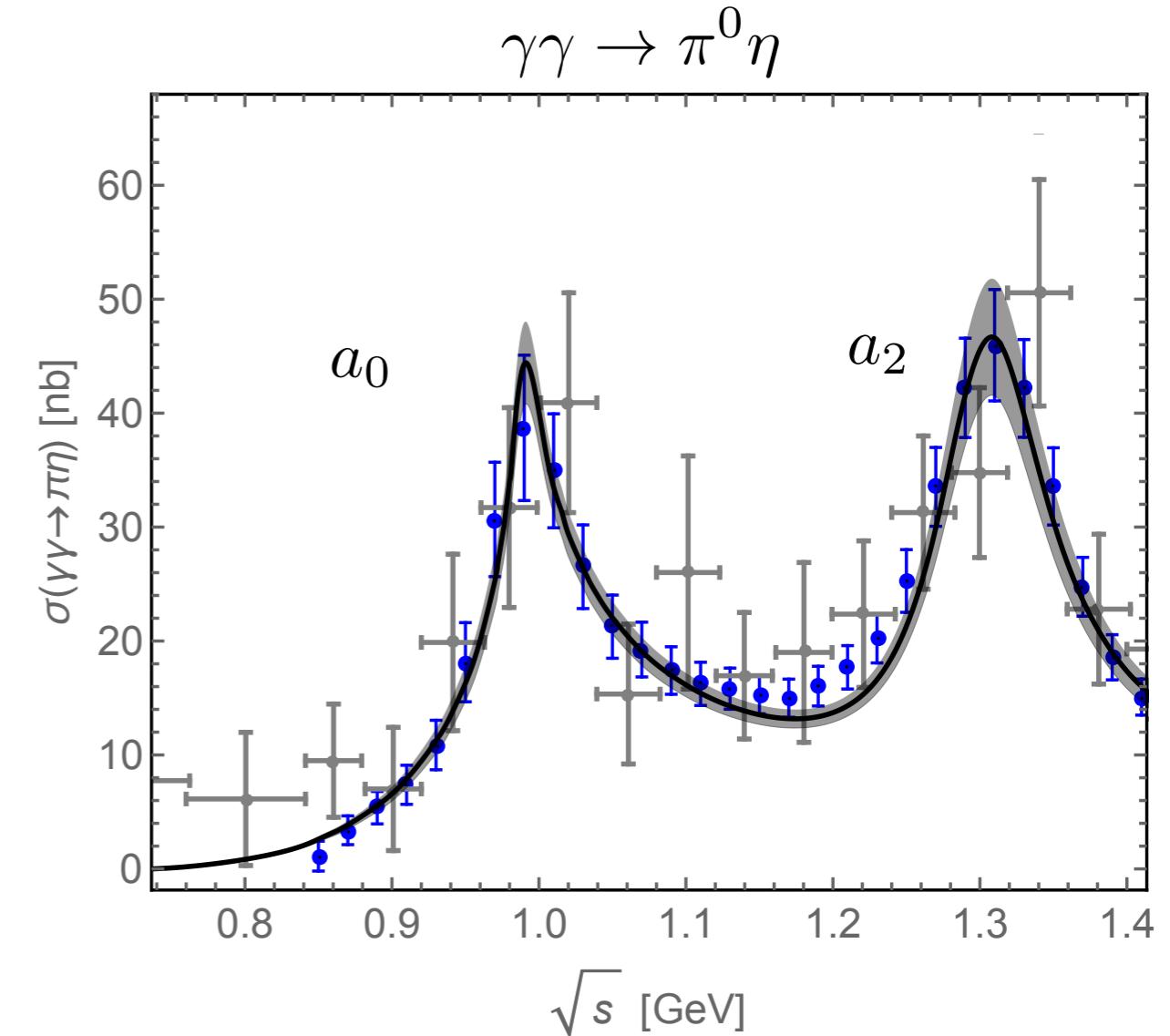
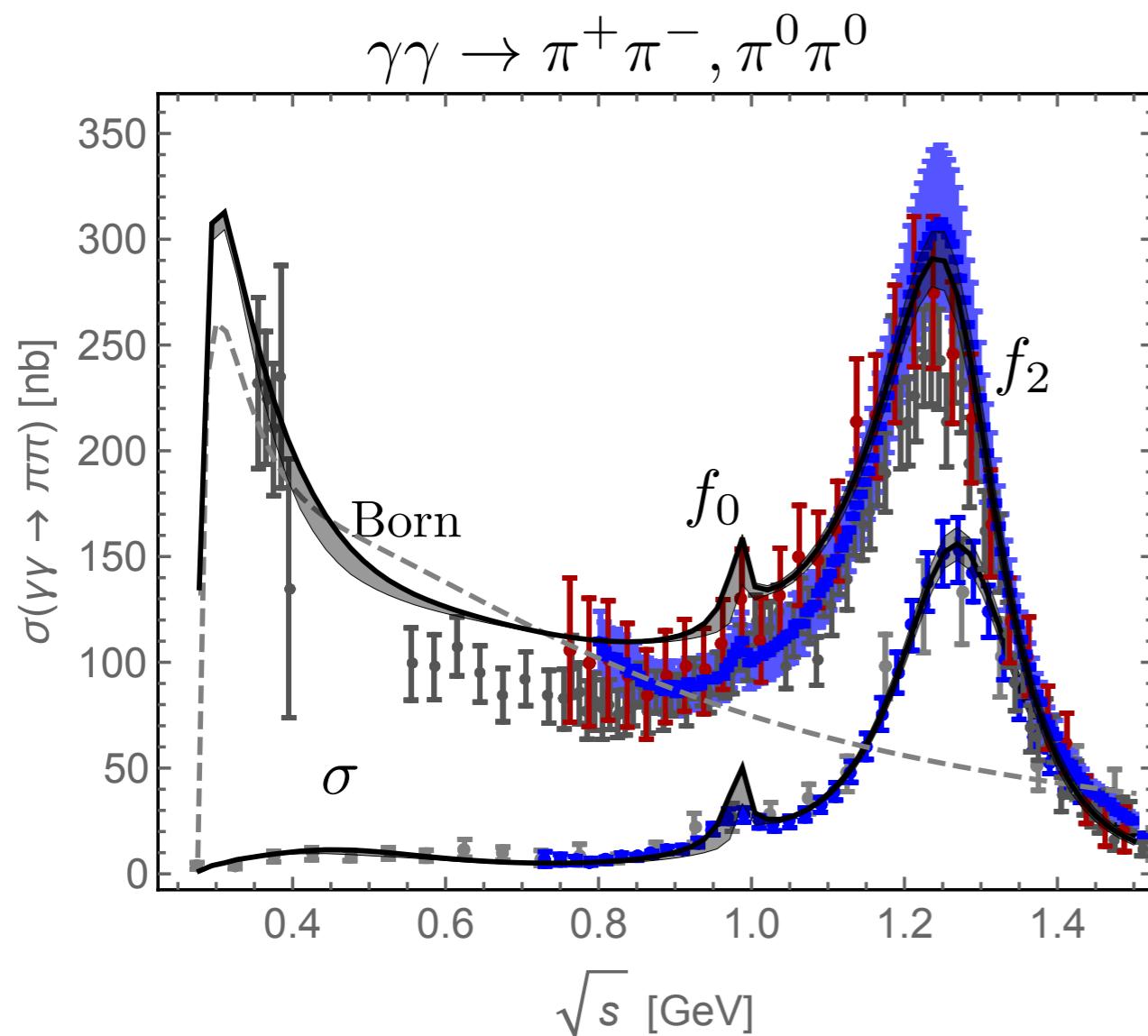


leads to

$$h_i^{(J)} = h_i^{(J),\text{Born}} + \Omega^{(J)} \left(\int_{-\infty}^0 \frac{ds'}{\pi} \frac{\text{Disc}(\bar{h}_i^{(J)}(s')) \Omega^{(J)}(s')^{-1}}{s' - s} - \int_{4m_\pi^2}^\infty \frac{ds'}{\pi} \frac{h^{(J),\text{Born}}(s') \text{Im} \Omega^{(J)}(s')^{-1}}{s' - s} \right)$$

Omnès (1958)
Muskhelishvili (1953)

Results for real photons

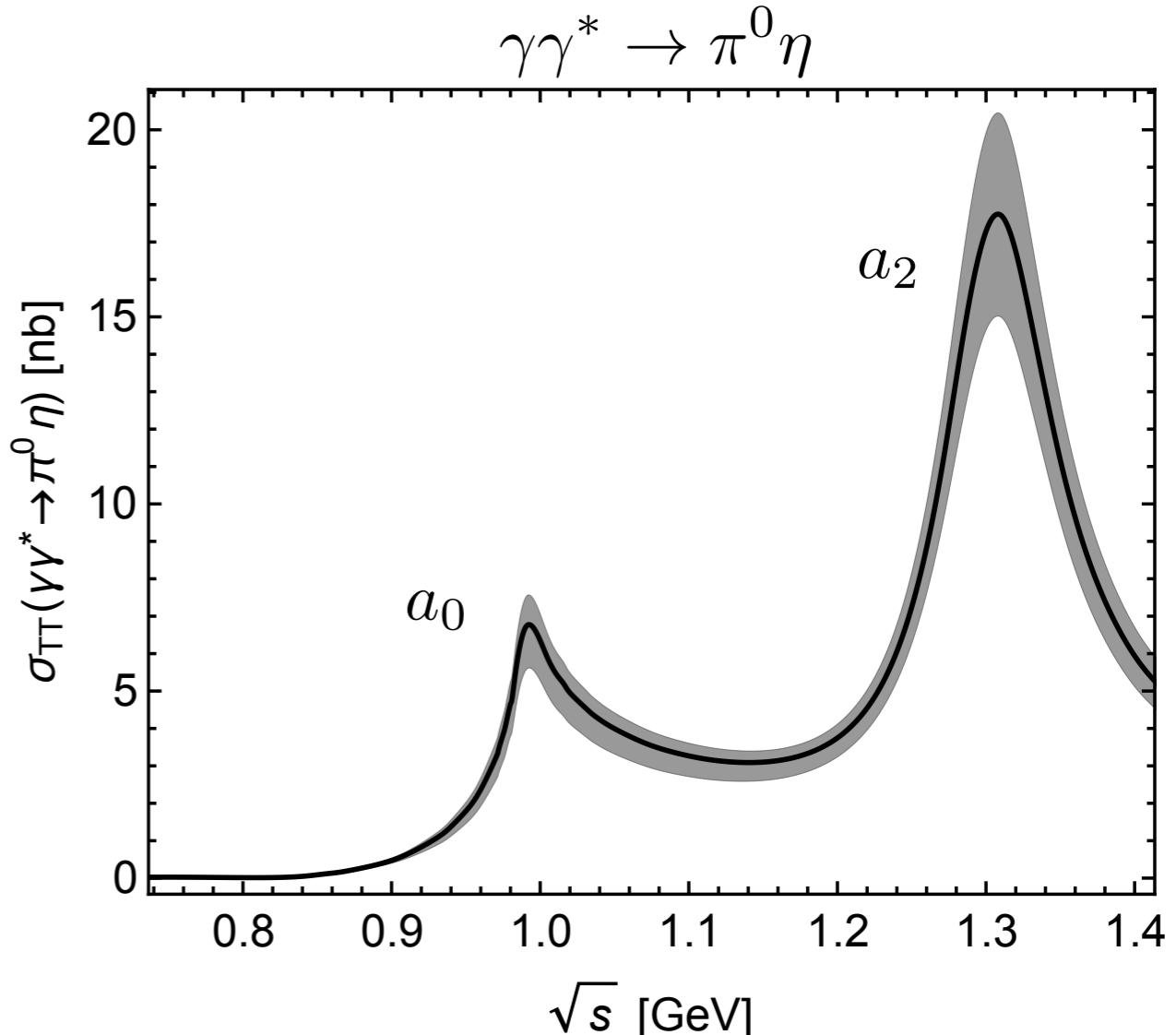
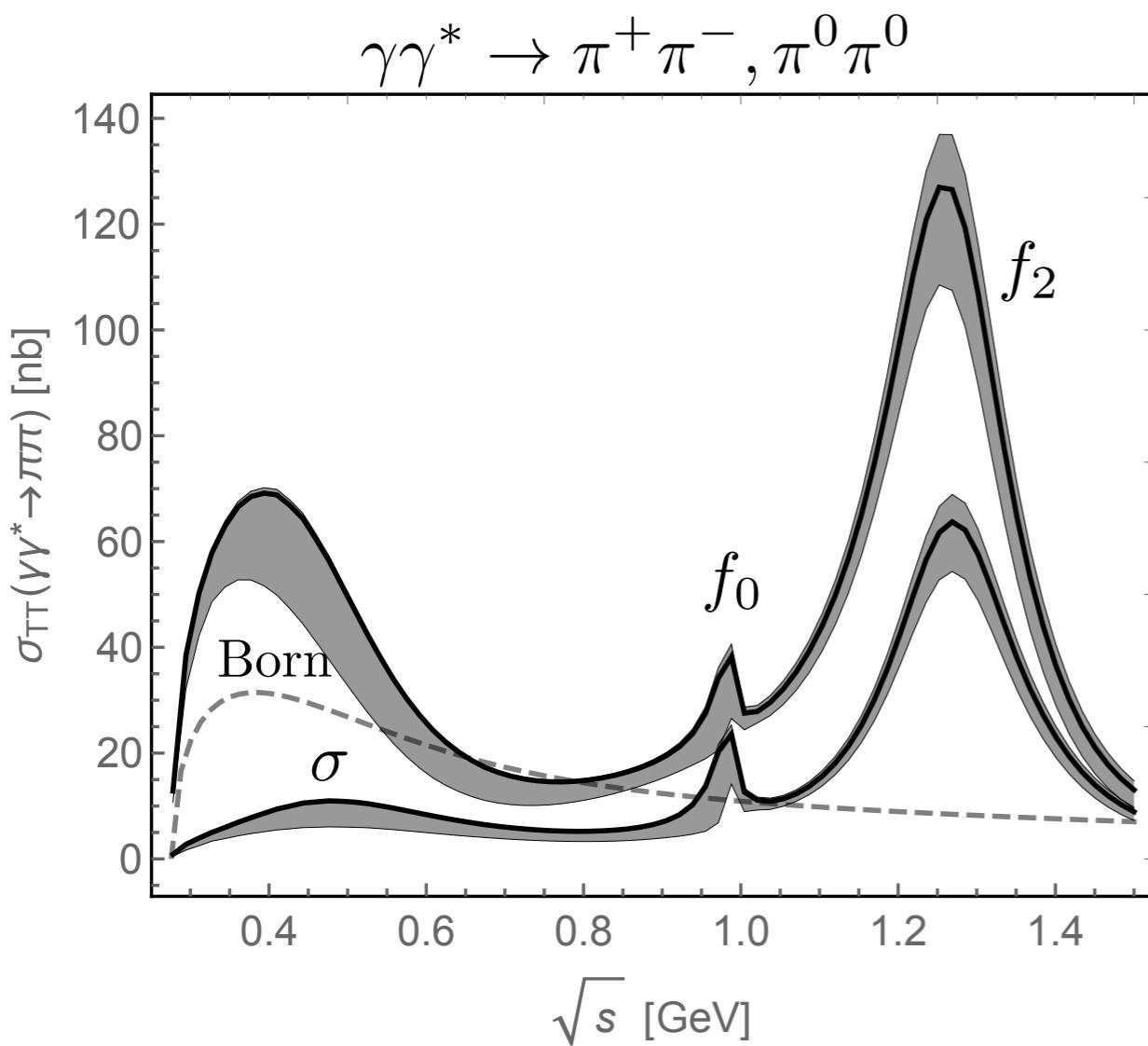


- Coupled-channel dispersive treatment of $f_0(980)$ and $a_0(980)$ is crucial
- $f_2(1270)$ described dispersively through Omnès function
- $a_2(1320)$ described as a Breit Wigner resonance

I.D., Deineka, Vanderhaeghen
(2017, 2018)

cf. also Dai et al. (2014)
Hoferichter et. al. (2011, 19)
Garcia-Martin et. al (2010)

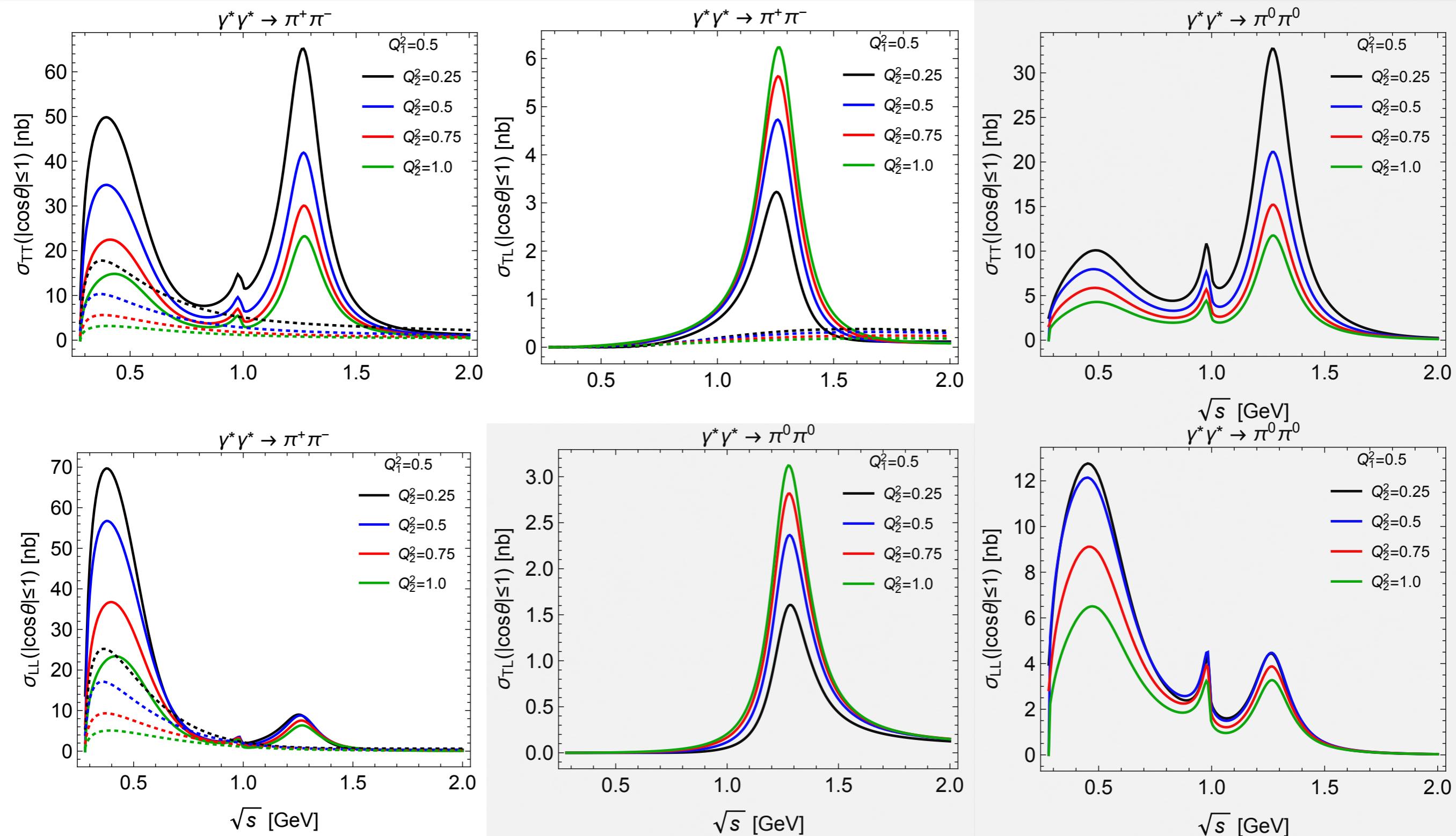
Results for single virtual photon ($Q^2=0.5$)



- Single tagged BESIII data for $\pi^+\pi^-$, $\pi^0\pi^0$ in range $0.1 \text{ GeV}^2 < Q^2 < 2 \text{ GeV}^2$ under analysis. It will validate left-hand cuts approximations.

I.D., Deineka, Vanderhaeghen (2018)
cf. also Moussallam (2013)
Hoferichter Stoffer (2019)

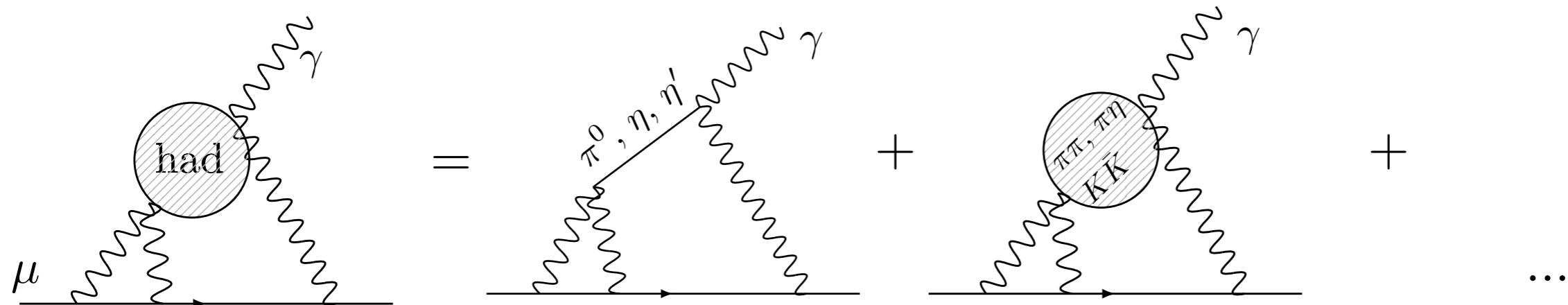
Results for double virtual photons



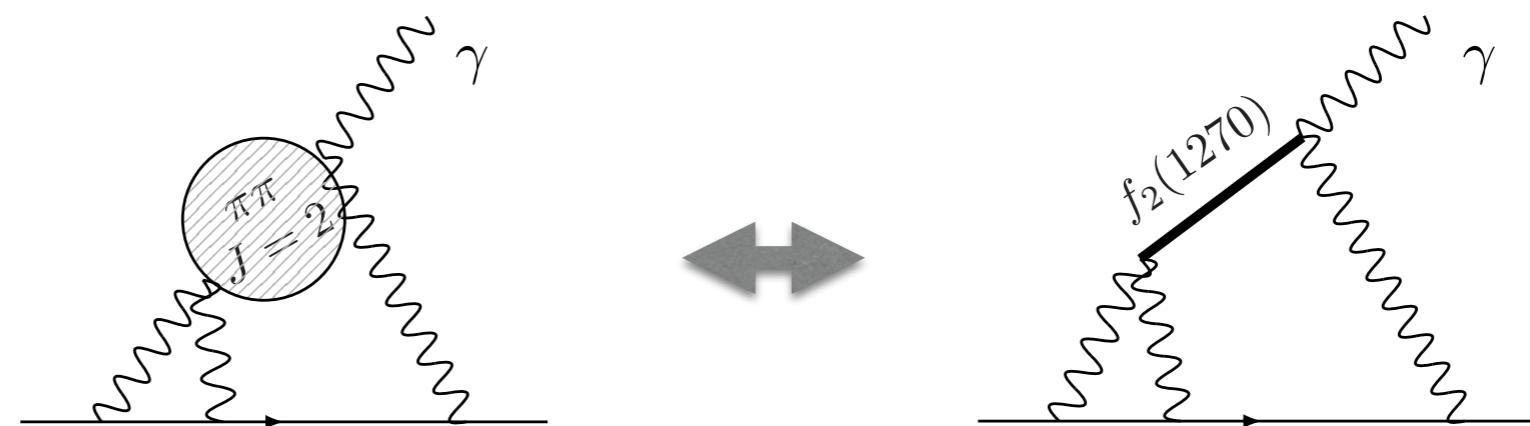
$$\frac{d\sigma_{TT}}{d\cos\theta} \sim |H_{++}|^2 + |H_{+-}|^2, \quad \frac{d\sigma_{TL}}{d\cos\theta} \sim |H_{+0}|^2, \quad \frac{d\sigma_{LL}}{d\cos\theta} \sim |H_{00}|^2$$

I.D., Deineka, Vanderhaeghen (2019)
cf. also Hoferichter Stoffer (2019)

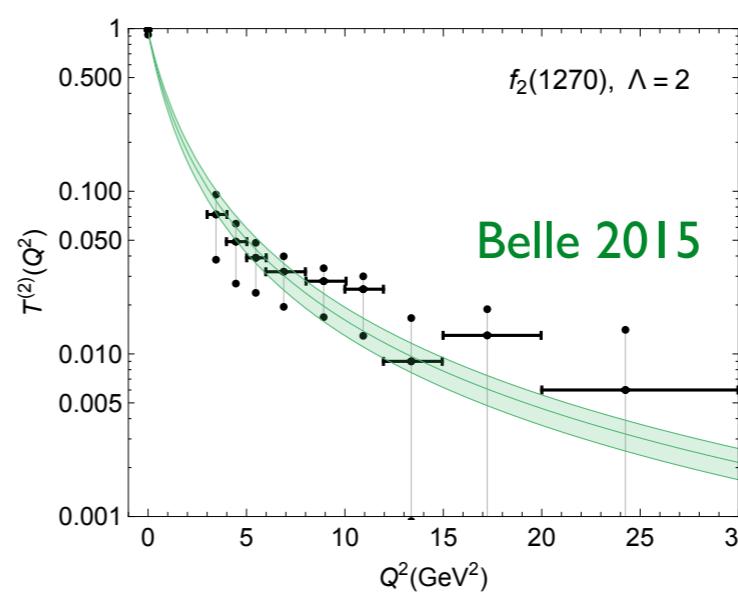
Effective resonance description



- One needs to compare



ongoing work



Pascalutsa, Pauk
Vanderhaeghen (2012)
I.D., Vanderhaeghen (2016)

$$a_\mu = (0.05 \pm 0.01) \times 10^{-10}$$

Summary and Outlook

- new a_μ Fermilab and J-Parc experiments ongoing:
aim: factor 4 improvement in experimental value
- Need to take into account $f_0(500)$, $f_0(980)$, $a_0(980)$, $f_2(1270)$... and non resonant contributions in a dispersive approach to $(g-2)$

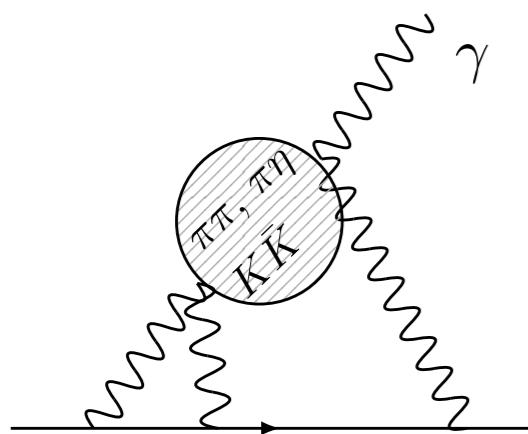
$$a_\mu^{\pi\text{-box}} + a_{\mu, J=0}^{\pi\pi, \pi\text{-pole LHC}} = -2.4(1) \times 10^{-10}$$

Colangelo et al. (2014-2017)

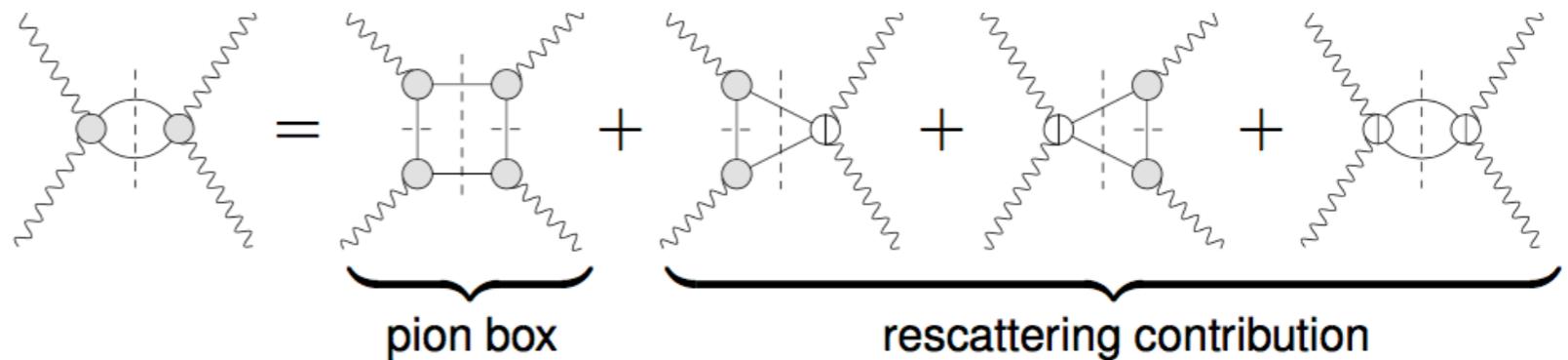
- Main ingredients: $\gamma^*\gamma^*\rightarrow\pi\pi$, $\pi\eta$ and (KK work in progress). Can be used in different $(g-2)$ dispersive approaches.
- Need to implement high energy asymptotic on TFF ($\gamma^*\pi\omega$, $\gamma^*\pi\rho$) used as input in dispersive formalism
- It is important to validate dispersive treatment of $\gamma\gamma^*\rightarrow\pi\pi$, $\pi\eta$, KK... with upcoming BESIII data

Thank you!

Multi-meson contribution to $(g-2)$



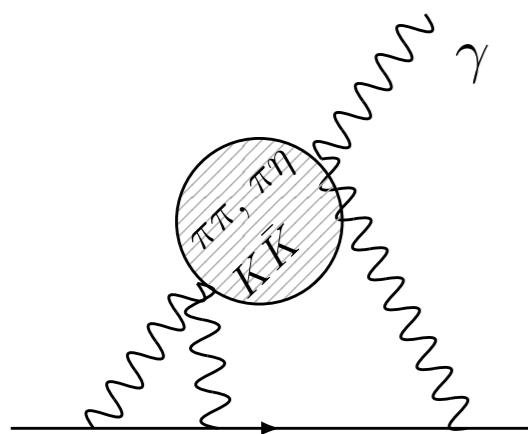
- Pioneering dispersive analyses for $\pi\pi\pi$ loop contribution to a_μ



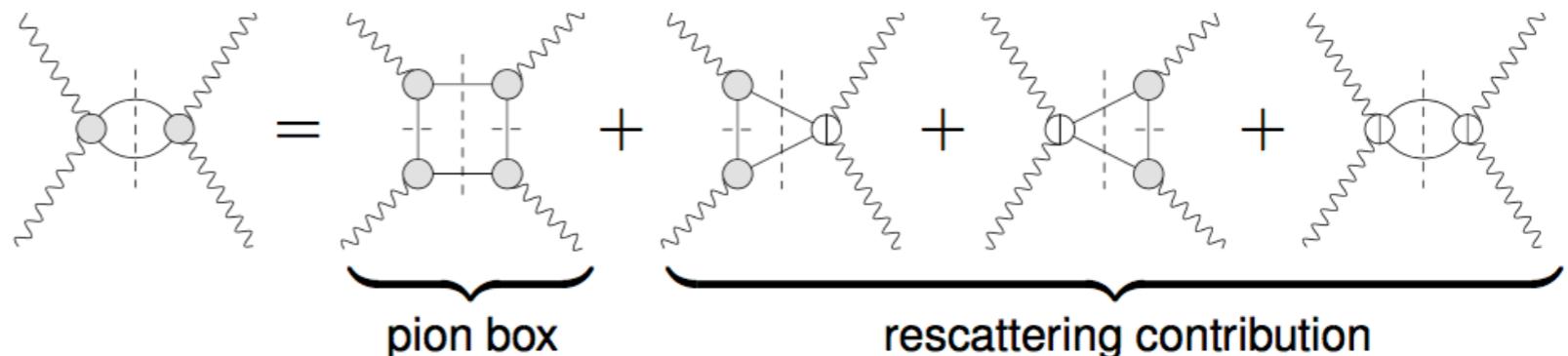
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Colangelo et al. (2014-2017)

Multi-meson contribution to $(g-2)$



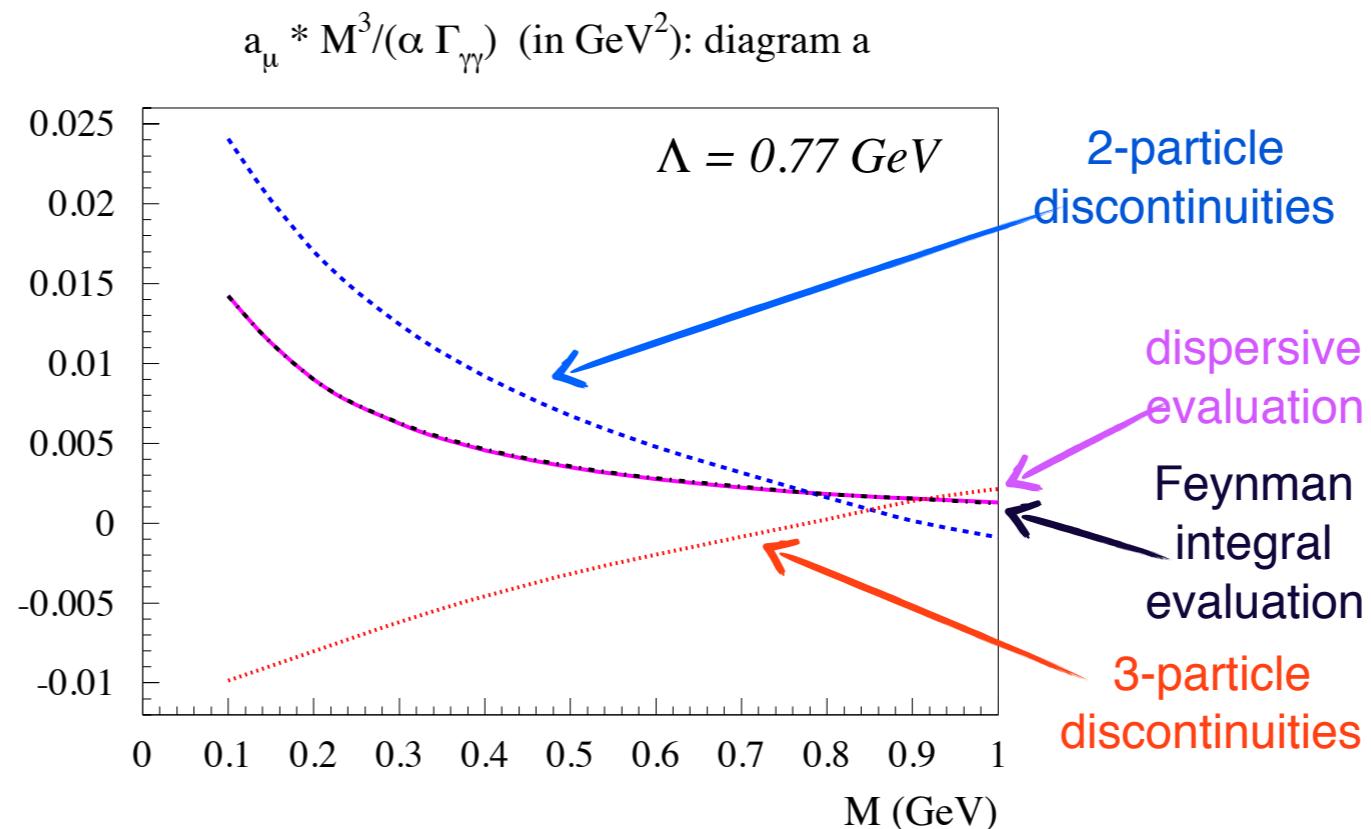
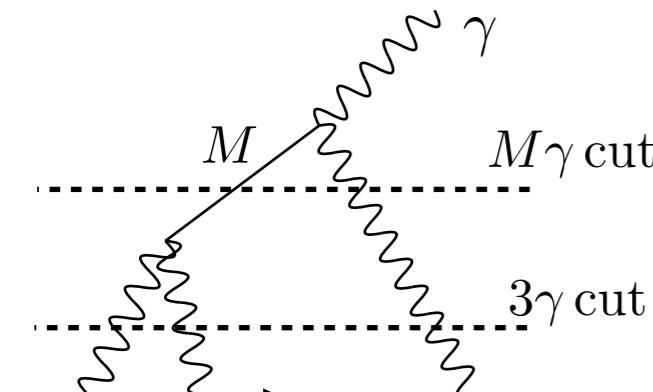
- Pioneering dispersive analyses for $\pi\pi\pi$ loop contribution to a_μ



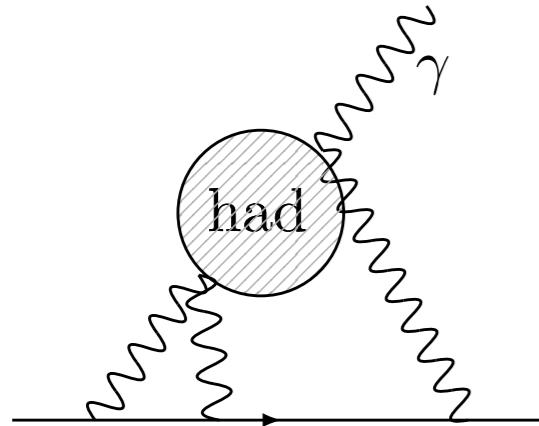
$$a_\mu^{\pi\text{-box}} + a_{\mu, J=0}^{\pi\pi, \pi\text{-pole LHC}} = -2.4(1) \times 10^{-10}$$

Colangelo et al. (2014-2017)

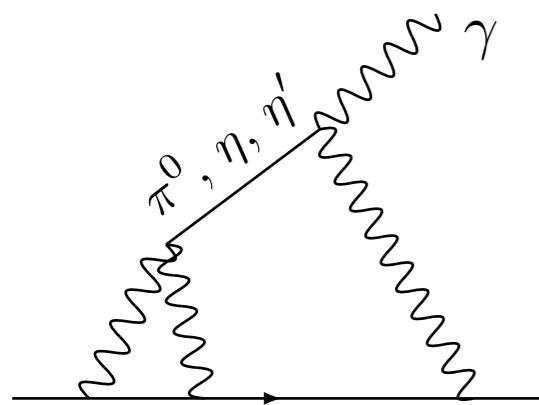
- Dispersive analysis for muon Pauli Form Factor



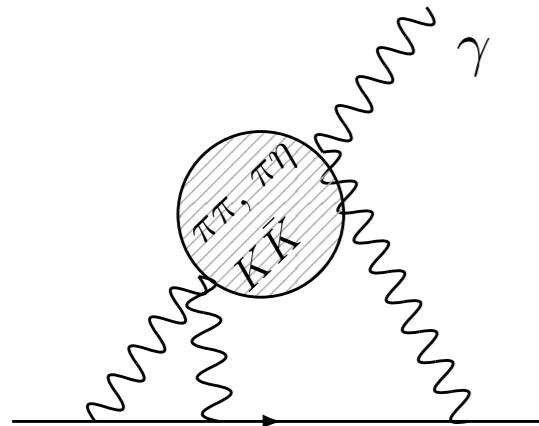
Observables in experiment



=



+



+

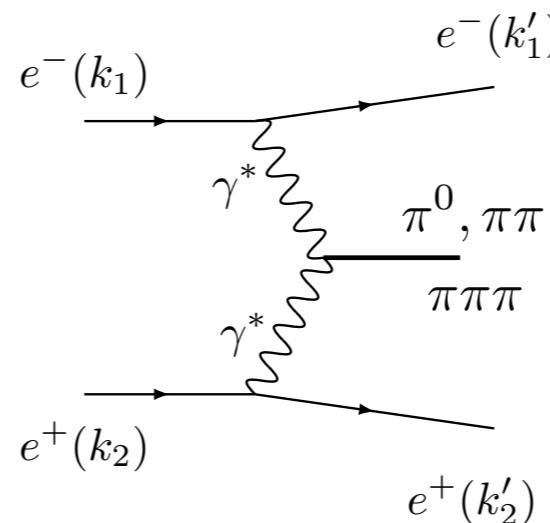
...

$$q^2 > 0$$

time-like γ^*

$$q^2 = -Q^2 < 0$$

space-like γ^*



$$\gamma\gamma^*\pi^0, \gamma\gamma^*\eta, \gamma\gamma^*\eta'$$

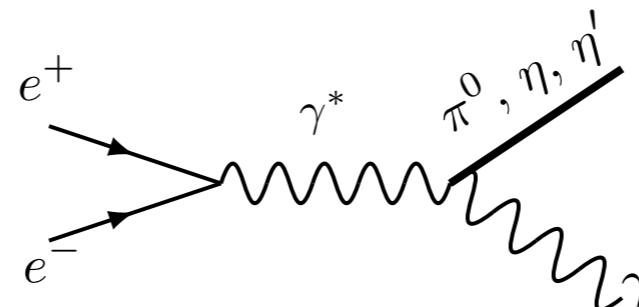
CLEO, BaBar, Belle, BESIII, ...

$$\gamma\gamma \rightarrow \pi\pi, \pi\eta, K\bar{K}$$

Mark II, CELLO, Crystal Ball, Belle

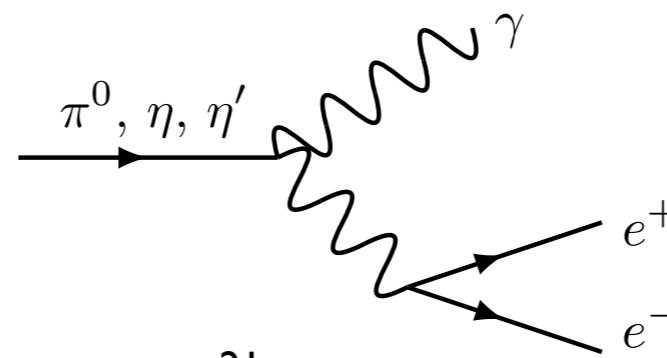
$$\gamma\gamma^* \rightarrow \pi\pi$$

BESIII



$$\gamma\gamma^*\pi^0, \gamma\gamma^*\eta$$

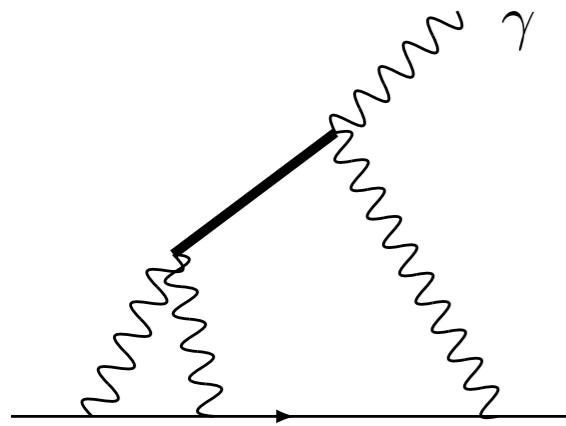
SND, CMD-2, BESIII, ...



$$\gamma\gamma^*\pi^0, \gamma\gamma^*\eta, \gamma\gamma^*\eta'$$

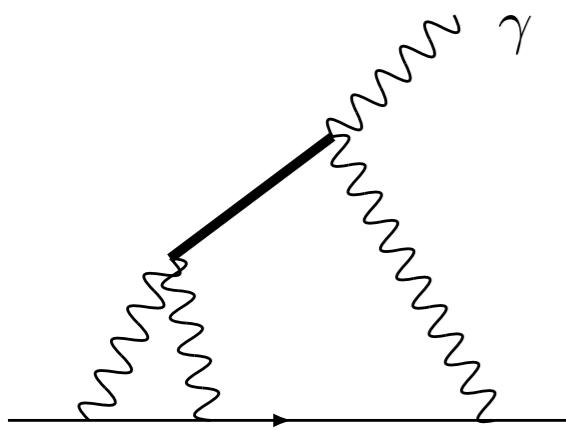
KLOE, MAMI/A2

Meson pole contributions to (g-2)



$$a_{\mu}^{LbL} = \lim_{k \rightarrow 0} ie^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2) \Pi_{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2)$$
$$\frac{1}{q_1^2} \frac{1}{q_2^2} \frac{1}{(k - q_1 - q_2)^2} \frac{1}{(p + q_1)^2 - m^2} \frac{1}{(p' - q_2)^2 - m^2}$$

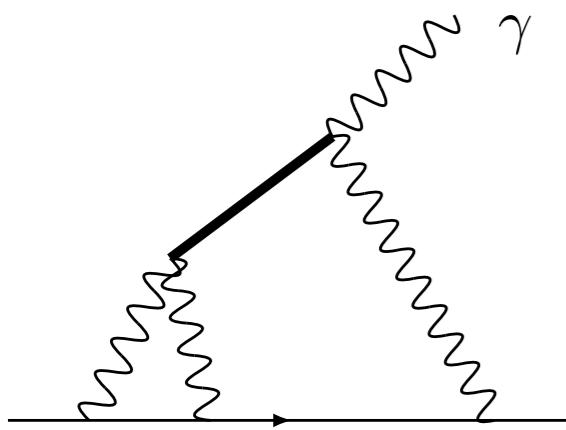
Meson pole contributions to (g-2)



Lepton tensor: well known

$$a_{\mu}^{LbL} = \lim_{k \rightarrow 0} ie^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2) \Pi_{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2)$$
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Meson pole contributions to (g-2)

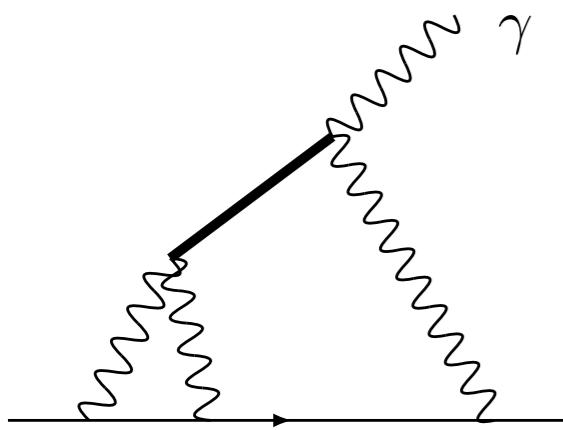


Lepton tensor: well known

Hadron tensor: requires input from **TFFs**

$$a_{\mu}^{LbL} = \lim_{k \rightarrow 0} ie^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2) \Pi_{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2)$$
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Meson pole contributions to (g-2)



Lepton tensor: well known

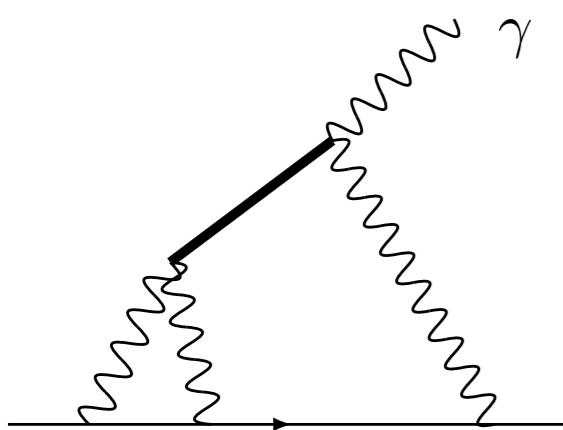
Hadron tensor: requires input from **TFFs**

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Results (excluding low energy region):

$$a_\mu[f_2(1270), f_2(1565)] = (0.1 \pm 0.01) \times 10^{-10}$$

Meson pole contributions to (g-2)



Lepton tensor: well known

Hadron tensor: requires input from **TFFs**

$$a_\mu^{LbL} = \lim_{k \rightarrow 0} ie^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2) \Pi_{\mu\nu\lambda\sigma}(q_1, k - q_1 - q_2, q_2)$$

$$\frac{1}{q_1^2} \frac{1}{q_2^2} \frac{1}{(k - q_1 - q_2)^2} \frac{1}{(p + q_1)^2 - m^2} \frac{1}{(p' - q_2)^2 - m^2}$$

Results (excluding low energy region):

$$a_\mu[f_2(1270), f_2(1565)] = (0.1 \pm 0.01) \times 10^{-10}$$

New evaluation of axial vector contributions (satisfying Landau-Yang theorem)

$$a_\mu[f_1(1285), f_1(1420)] = (0.64 \pm 0.20) \times 10^{-10}$$

$$= (0.75 \pm 0.27) \times 10^{-10}$$

Pauk, Vdh (2013)
Jegerlehner (2015)