

# Study of *KN* interaction from the hadron-hadron correlation in high-energy nuclear collisions



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- Introduction: Hadron correlation in high energy nuclear collisions
- Introduction:  $K^-p$  correlation function
- $K^-p$  correlation function with Koonin-Pratt formula
- Results and discussion

#### High energy nuclear collision and FSI



Hadron-hadron correlation

$$C_{12}(k_1, k_2) = \frac{N_{12}(k_1, k_2)}{N_1(k_1)N_2(k_2)}$$
  
= 
$$\begin{cases} 1 & (\text{w/o correlation}) \\ \text{Others (w/ correlation)} \end{cases}$$

#### High energy nuclear collision and FSI



#### Hadron-hadron correlation

$$C_{12}(\mathbf{k}_1, \mathbf{k}_2) = \frac{\int d^4 x_1 \int d^4 x_2 S_1(x_1, \mathbf{k}_1) S_2(x_2, \mathbf{k}_2) |\Psi^{(-)}(\mathbf{r}, \mathbf{q})|^2}{\int d^4 x_1 S_1(x_1, \mathbf{k}_1) \int d^4 x_2 S_2(x_2, \mathbf{k}_2)}$$

 $S_i(x, \mathbf{k})$ : Source function  $\Psi^{(-)}(\mathbf{r}, \mathbf{q})$ : Relative wave function

High energy nuclear collision and FSI



#### High energy nuclear collision and FSI



#### Hadron-hadron correlation

$$C_{12}(\mathbf{k}_1, \mathbf{k}_2) = \frac{\int d^4 x_1 \int d^4 x_2 S_1(x_1, \mathbf{k}_1) S_2(x_2, \mathbf{k}_2) |\Psi^{(-)}(\mathbf{r}, \mathbf{q})|^2}{\int d^4 x_1 S_1(x_1, \mathbf{k}_1) \int d^4 x_2 S_2(x_2, \mathbf{k}_2)}$$

 $S_i(x, \mathbf{k})$ : Source function

 $\Psi^{(-)}(\mathbf{r}, \mathbf{q})$ : Relative wave function

• Depends on ...

Interaction (strong and Coulomb)

quantum statistics (Fermion, boson)

#### How to study the hadron interaction $C_{12}(\mathbf{k}_1, \mathbf{k}_2) = \frac{\int d^4 x_1 \int d^4 x_2 S_1(x_1, \mathbf{k}_1) S_2(x_2, \mathbf{k}_2) |\varphi^{(-)}(\mathbf{r}, \mathbf{q})|^2}{\int d^4 x_1 S_1(x_1, \mathbf{k}_1) \int d^4 x_2 S_2(x_2, \mathbf{k}_2)}$ $S_i(x, \mathbf{k})$ : Source function $\varphi^{(-)}(\mathbf{r}, \mathbf{q})$ : Relative wave function Introduce relative source: $S(\mathbf{k}_i)$ Move to the center-of-mass frame Koonin-Pratt formula : $C(\mathbf{q}) \simeq \left| d^3 \mathbf{r} S(\mathbf{r}) \right| \varphi^{(-)}(\mathbf{r}, \mathbf{q}) \right|^2$ S.E. Koonin, PLB 70 (1977) S. Pratt et. al. PRC 42 (1990) Simple model with LL formula 3 $R/a_0 = -1 - -$ R. Lednicky, et al. Sov. J. Nucl. Phys. 35(1982). $R/a_0 = 0.1$ $R/a_0 = -0.1$ • Static Gaussian source 2 $R/a_0=1$ C(q) • Effective range expansion • C(q) is sensitive to $R/a_0$ r<sub>eff</sub>/R=0 • *R* : Gaussian source size 0 0.5 1.5 1 2 0 • *a*<sub>0</sub>: scattering length qR Morita, et al., arXiv:1908.05414

Powerful tool to study hadron interaction in low energy region

• Example:  $\Lambda\Lambda$  interaction from  $\Lambda\Lambda$  correlation function



- ΛΛ correlation ; measured by STAR collabo. (later by ALICE collabo.) L. Adamczyk, et al. PRL 114 (2015).
   S. Acharya et al., PRC 99 (2019)
- Detailed analysis of correlation function Morita et al. PRC 91(2015) Ohnishi et. al. Nuclear Physics A 954 (2016)

Constraint to the  $\Lambda\Lambda$  threshold parameters  $(a_0, r_e)$ 

### K pa correlation



- p+p collision at  $\sqrt{s} = 5$ , 7, 13 TeV, source radius R = 1.13 fm 1.18 fm
- Strong enhancement (*C* > 1) at small momenta ==> <u>Coulomb interaction</u>
- But larger than prediction with pure Coulomb ==> Strong interaction
- Characteristic cusp at the  $\bar{K}^0 n$  threshold (k = 58 MeV) ==> isospin sym. breaking

Can we derive constraint on  $\overline{KN}$  interaction from  $K^-p$  correlation?

# K<sup>-</sup>p correlation

0.8  $K^{-}p_{.7} \text{ correlation:} \\ correlation: constrained and constraine and constraine and constrained and co$ <sup>0.7</sup>ALICE, S. ALICE pp s = 13 TeV  $r_0 = 1.18 \pm 0.01 \pm 0.12$  fm  $C(k^*)$ ALICE\$β<sup>64</sup>/₅ ≌.86τeV AL<sup>2</sup>ICE<sup>0</sup>pp<sup>1</sup> <del>₹s<sup>0</sup>=07</del> TeV ALICE 61 ₹8 913 TeV 0.5 0.5  $r_0 = 1^{1} \cdot 18^{\oplus} \pm 0^{\circ} \cdot 01 \pm 0.12 \text{ fm}$  $r_0 = 1.13 \pm 0.02 \stackrel{+0.17}{_{-0.15}} \text{ fm}$  $r_0 = 1.13 \pm 0.02 \stackrel{+0.17}{_{-0.15}}$  fm  $\lambda = 0.64^{\text{omb}} \pm 0.06$ 0.2  $\lambda = 0.68 \pm 0.07$ 0.4  $\lambda = 0.76 \pm 0.08$ 0.4 Kp<sup>-</sup>Ceulomb+Strong (Julich Model) 0.3 0.3 0.3 Coulomb 0.7 ₹0\$5\_ < 1150 200 250<sup>-</sup> Coulomb+Strong (Kyoto Mo¢**k**A¢V/c) RT-149Sthreshold 50 **0**.7 **ℓ**0**S**<sub>T</sub> < 1150 - Model 200 Coulomb+Strong (Jülich Model) 250 250 250 k\* (MeV/c) k\* (MeV/c) k\* (MeV/c)- $0.7 < S_{T} < 1$  $0.7 < S_{T} < 1$  $0.7 < S_{T} < 1$ 0.8 50 100 150 50 100 2500 50 100 200 250 200 0 2500 150 200 150 *k*\* (MeV/*c*) *k*\* (MeV/*c*) k\* (MeV/c)

Kyoto Model

Ohnishi et al. NPA 954 (2016) Cho, et al., PPNP 95 (2017)

- Interaction: Based on Chiral SU(3) dynamics
  Ikeda, Hyodo, Weise, NPA881 (2012)
- Calculated with
  - Coulomb + Strong int.
  - $\overline{K}N (K^-p + \overline{K}^0n)$  w/ isospin ave. mass

#### Jülich Model

Haidenbauer NPA 981 (2018)

- Interaction: Jülich meson exchange model Refitted ver. of Müller-Groeling, et al., NPA 513 (1990)
- Calculated with
  - Coulomb (Gamow) + Strong int.
  - $\bar{K}N + \pi\Sigma + \pi\Lambda$  with particle mass

# K<sup>-</sup>p correlation

 $K^{-}p_{7}\text{correlation:} p_{0} \in \mathbb{C}$ <sup>0.7</sup>ALICE, S. Acharya et al., (2019), 1905.13470.  $r_0 = 1.18 \pm 0.01 \pm 0.12$  fm ALICE<sup>0</sup>pp<sup>1</sup> <del>₹</del>s<sup>0</sup>=<sup>0</sup>7 TeV ALICE964/₅ 9.96TeV ALICE 61 ₹8 913 TeV 0.5 0.5  $r_0 = 1^{1} \cdot 18^{\oplus} \pm 0^{\circ} \cdot 01 \pm 0.12 \text{ fm}$  $r_0 = 1.13 \pm 0.02 \stackrel{+0.17}{_{-0.15}}$  fm  $r_0 = 1.13 \pm 0.02 \stackrel{+0.17}{_{-0.15}}$  fm  $\lambda = 0.64^{\text{omb}} \pm 0.06$ 0.2  $\lambda=0.68\pm0.07$ 0.4  $\lambda = 0.76 \pm 0.08$ 0.4 K<sup>-</sup>pCeulomb+Strong (Julich Model)
 0.5 0.3 0.3 Coulomb 0.7 ±09\_ < 1150 200 250<sup>−</sup> Coulomb+Strong (Kyoto №o¢kplev/*c*) 50 R749Sthreshold. 50 **0**.7 **ℓ**0**S**<sub>T</sub> < 1150 Model 200 Coulomb+Strong (Jülich Model) k\* (MeV/c) k\* (MeV/c) k\* (MeV/c)- $0.7 < S_{T} < 1$  $0.7 < S_{T} < 1$  $0.7 < S_{T} < 1$ 0.8 50 100 150 50 100 2500 50 100 250 0 200 2500 150 200 150 200 *k*\* (MeV/*c*) *k*\* (MeV/*c*) k\* (MeV/c)

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- What we need to include
  - Coulomb interaction
  - Coupled-channel (decay channels)
  - Threshold energy difference  $(K^-p, \bar{K}^0n)$

and open channel contribution  $(\bar{K}^0 n)$ 

We update the Kyoto model calculation and discuss the coupled-channel effect and interaction dependence.

### *K<sup>-</sup>p* correlation with Koonin-Pratt Formula

#### Koonin-Pratt formula

$$C(\mathbf{k}_1, \mathbf{k}_2) = \frac{\int d^4 x_1 \int d^4 x_2 S_1(x_1, \mathbf{k}_1) S_2(x_2, \mathbf{k}_2) |\varphi^{(-)}(\mathbf{r}, \mathbf{q})|^2}{\int d^4 x_1 S_1(x_1, \mathbf{k}_1) \int d^4 x_2 S_2(x_2, \mathbf{k}_2)}$$

- Introduce relative source:  $S(\mathbf{k}_i)$
- Move to the center-of-mass frame

Koonin-Pratt formula : 
$$C(\mathbf{q}) \simeq \left[ d^3 \mathbf{r} S(\mathbf{r}) | \varphi^{(-)}(\mathbf{r}, \mathbf{q}) |^2 \right]$$

S.E. Koonin, PLB 70 (1977) S. Pratt et. al. PRC 42 (1990)

- Static Gaussian source function:  $S \propto \exp(-r^2/4R^2)$ 
  - *R* : Relative source radius
- Consider only *s*-wave
- Single-channel, no Coulomb int., non-identical particles

$$C(\mathbf{q}) = 1 + \int d^3 \mathbf{r} \ S(\mathbf{r}) \left[ -|j_0(qr)|^2 + |\chi^{(-)}(r,q)|^2 \right]$$

Free *s*-wave: (Spherical Bessel fcn.)

Scattering wave function:

obtained by solving the Schrödinger Eq. 12

### $K^-p$ correlation with Koonin-Pratt Formula

$$C(\mathbf{q}) = \underline{1} + \int d^3 \mathbf{r} \ S(\mathbf{r}) \left[ - \left| j_0(qr) \right|^2 + \left| \chi^{(-)}(r,q) \right|^2 \right]$$

• With Coulomb interaction:  $V_{\text{Coulomb}} = \alpha/r$ 



### $K^-p$ correlation with Koonin-Pratt Formula

$$C(\mathbf{q}) = \underline{1} + \int d^3 \mathbf{r} \ S(\mathbf{r}) \left[ - |j_0(qr)|^2 + |\chi^{(-)}(r,q)|^2 \right]$$

• With Coulomb interaction:  $V_{\text{Coulomb}} = \alpha/r$ 

$$C(\mathbf{q}) = \int d^3 \mathbf{r} \ S(\mathbf{r}) \left[ \frac{|\varphi^{C,\text{full}}(\mathbf{r},\mathbf{q})|^2}{\text{Full Coulomb}} - \frac{|j_0^C(qr)|^2}{\text{Coulomb}} + \frac{|\chi^{C,(-)}(r,q)|^2}{\text{Scattering wave}} \right]$$
  
wave s-wave: function with Coulomb int.

$$C_{i}(\mathbf{q}) = \int d^{3}\mathbf{r} \ S_{i}(\mathbf{r}) \left[ \left| \varphi^{C,\text{full}}(\mathbf{r},\mathbf{q}) \right|^{2} - \left| j_{0}^{C}(qr) \right|^{2} + \left| \chi_{i}^{C,(-)}(r,q) \right|^{2} \right] + \sum_{j \neq i} \omega_{j} \left[ d^{3}\mathbf{r} \ S_{j}(\mathbf{r}) \left| \chi_{j}^{C,(-)}(r,q) \right|^{2} \right]$$
  
Coupled channel effect

- $\omega_j$ : weight of channel j => We assume  $\omega_j = 1$ .
- $\chi_j^{(-)}(r,q)$ : channel *j* component of wave function with channel *i* outgoing boundary condition

R. Lednicky, et. al. Phys. At. Nucl. 61 (1998)

### *K*<sup>-</sup>*p* correlation with Koonin-Pratt Formula

#### • Chiral SU(3) based $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential

Miyahara, Hyodo, Weise, PRC 98 (2018)

• Constructed based on the amplitude with chiral SU(3) dynamics

Ikeda, Hyodo, Weise, NPA881 (2012)

• Coupled-channel, energy dependent as

$$V_{ij}(r,E) = e^{-(b_i/2 + b_j/2)r^2} \sum_{\alpha=0}^{\alpha_{\max}} K_{\alpha,ij} (E/100 \text{ MeV})^{\alpha}$$

• Reproduce the coupled-channel amplitude around the  $\bar{K}N$  sub-threshold region

© Coupled-channel boundary condition R. Lednicky, et. al. Phys. At. Nucl. 61 (1998)

• Asymptotic waves

open channels :  $\chi_j^{(C)}(r,q) \to a_j(\text{outgoing wave}) + b_j(\text{incoming wave})$ closed channels :  $\chi_j^{(C)}(r,q) \to a_j(\text{diverg. solution}) + b_j(\text{converg. solution})$ 

• Out-going wave boundary condition for  $K^-p$  channel;  $|a_{K^-p}| = 1$ ,  $|a_{others}| = 0$ 





#### • $K^-p$ correlation in particle basis w/ Coulomb



- C(q) calculated with  $K^-p$ ,  $K^-p + \bar{K}^0 n$ ,  $\bar{K}N + \pi\Sigma + \pi\Lambda$  components
- Inclusion of  $\overline{K}^0 n ==$  enhance correlation and the cusp structure
- Inclusion of decay channels ==> non-negligible enhancement

Inclusion of coupled channel effect is essential for  $K^-p$  correlation





## Results

Interaction dependence of  $\overline{K}N$  correlation

- $I = 0 \bar{K}N$  interaction <== strongly constrained by the SIDDHARTA constraint M. Bazzi, et al., NPA 881 (2012)
- $I = 1 \ \bar{K}N$  interaction is not well known ==> vary  $V_{\bar{K}N-\bar{K}N}^{I=1} \rightarrow \beta V_{\bar{K}N-\bar{K}N}^{I=1}$
- SIDDHARTA constraint on  $a_0^{K^-p} ==>$  Varied region of  $\beta$  as  $-0.24 < \beta < 1.09$



• For 
$$\beta = -0.24$$
,

- Remarkable suppression around  $\bar{K}^0 n$  threshold ( $q \simeq 58 \text{ MeV}$ )
- Moderate cusp structure

= 1  $\overline{K}N$  interaction can be determined with the detailed analysis!

Preliminary!



# Conclusion

- Hadron correlation in high energy nuclear collisions is a powerful tool to study the interaction of strangeness systems.
- $K^-p$  correlation function has been studied with the Koonin-Pratt formula and interaction based on chiral dynamics, including Coulomb interaction, coupled-channel effect, and threshold energy difference between  $K^-p$  and  $\bar{K}^0n$ . Characteristic cusp structure is found at  $\bar{K}^0n$ threshold, as found in ALICE data.
- C(q) is also sensitive to  $I = 1 \overline{KN}$  interaction, so that detailed analysis of  $K^-p$  correlation can give the constraint on I = 1 interaction.
- Detailed comparison with experimental data ==> Future study



# Thank you!



#### How to study the hadron interaction $C_{12}(\mathbf{k}_1, \mathbf{k}_2) = \frac{\int d^4 x_1 \int d^4 x_2 S_1(x_1, \mathbf{k}_1) S_2(x_2, \mathbf{k}_2) |\varphi^{(-)}(\mathbf{r}, \mathbf{q})|^2}{\int d^4 x_1 S_1(x_1, \mathbf{k}_1) \int d^4 x_2 S_2(x_2, \mathbf{k}_2)}$ $S_i(x, \mathbf{k})$ : Source function $\varphi^{(-)}(\mathbf{r}, \mathbf{q})$ : Relative wave function Introduce relative source: $S(\mathbf{k}_i)$ Move to the center-of-mass frame Koonin-Pratt formula : $C(\mathbf{q}) \simeq \left| d^3 \mathbf{r} S(\mathbf{r}) \right| \varphi^{(-)}(\mathbf{r}, \mathbf{q}) \right|^2$ S.E. Koonin, PLB 70 (1977) S. Pratt et. al. PRC 42 (1990) Simple model with LL formula 3 $R/a_0 = -1 - -$ R. Lednicky, et al. Sov. J. Nucl. Phys. 35(1982). $R/a_0 = 0.1$ $R/a_0 = -0.1$ • Static Gaussian source 2 $R/a_0=1$ C(q) • Effective range expansion • C(q) is sensitive to $R/a_0$ r<sub>eff</sub>/R=0 • *R* : Gaussian source size 0 0.5 1.5 1 2 0 • *a*<sub>0</sub>: scattering length qR Morita, et al., arXiv:1908.05414

Powerful tool to study hadron interaction in low energy region



#### •*K*<sup>-</sup>*p* correlation in isospin basis w/o Coulomb



Results

- C(q) calculated only with  $\overline{KN}(K^-p + \overline{K}^0n)$  component with R = 1.2 fm
- Re  $a_0^{I=0} > 0 \Longrightarrow C_{\bar{K}N}^{I=0} < 1$ , Re  $a_0^{I=1} < 0 \Longrightarrow C_{\bar{K}N}^{I=1} > 1$  at small q

• 
$$C_{K^-p}(q) = (C_{\bar{K}N}^{I=0} + C_{\bar{K}N}^{I=1})/2$$

### Results



- C(q) calculated with  $K^-p$ ,  $K^-p + \overline{K}^0 n$ , all of coupled-channel components
- Inclusion of  $\bar{K}^0 n ==>$  enhance correlation and the cusp structure
- Inclusion of decay channels ==> non-negligible enhancement

### Results

#### $\sim K^- p$ correlation in particle basis w/ Coulomb



- C(q) calculated with  $K^-p$ ,  $K^-p + \overline{K}^0 n$ , all of coupled-channel components
- Strong enhancement at the small  $q \leq 20 \text{ MeV}$
- Cusp structure is still clear













- Strong dependence around  $R \sim 1.0$  fm <== • Sensitive region:  $|R/a_0| \leq 1$ .
- Larger system ( > 3 fm) : cusp effect is moderate

Preliminary!