

Partial Wave Mixing in Hamiltonian Effective Field Theory (HEFT)

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Outline

- 1 Background and Motivation
- 2 Partial Wave Mixing in HEFT
- 3 Example of Isospin-2 $\pi\pi$ Scattering
- 4 Summary

Observables are different in experiments and lattice

- **Experiments:** Scattering in the infinite volume
 - Quantum states: Scattering states
 - Observables: Scattering observables like phase shifts $\delta_l(E)$

- **Lattice:** Scattering in a finite periodic box
 - Quantum states: Bound states
 - Observables: Finite volume spectra $E_n(L)$

Symmetry is different in experiments and lattice

- **Experiments:** Spherical symmetry → Rotation group: $O(3)$
 - Phase shifts $\delta_l(E)$ are classified by the **irreps** l^\pm of $O(3)$
 - Irreps of $O(3)$: $0^\pm, 1^\pm, 2^\pm, \dots$
- **Lattice:** Cubic symmetry → Cubic group: O_h
 - Spectra $E_n(\Gamma, L)$ are classified by the irreps Γ of O_h
 - Irreps of O_h : $A_1^\pm, A_2^\pm, E^\pm, T_1^\pm, T_2^\pm$.

Irreducible representations

Symmetry is different in experiments and lattice

- O_h is a subgroup of $O(3)$
- The irreps of $O(3)$ can be reduced into the irreps of O_h
- For each l , we have $m \rightarrow (\Gamma, f, \alpha)$
 - Γ denotes the irrep of O_h
 - f counts the number of Γ in the l (It's always 1 for $l \leq 4$)
 - α counts the dimension of Γ

$$\mathbf{0}^+ = \mathbf{A}_1^+$$

$$\mathbf{1}^- = \mathbf{T}_1^-$$

$$\mathbf{2}^+ = \mathbf{E}^+ \oplus \mathbf{T}_2^+$$

$$\mathbf{3}^- = \mathbf{A}_2^- \oplus \mathbf{T}_1^- \oplus \mathbf{T}_2^-$$

$$\mathbf{4}^+ = \mathbf{A}_1^+ \oplus \mathbf{E}^+ \oplus \mathbf{T}_1^+ \oplus \mathbf{T}_2^+$$

What is partial wave mixing

- Use l_{cut} to ignore higher partial wave ($l > l_{\text{cut}}$) contributions
- $l_{\text{cut}} = 0$: Only s-wave phase shifts are related to \mathbf{A}_1^+ spectra
- $l_{\text{cut}} = 4$: Both s- and g-wave phase shifts are related to \mathbf{A}_1^+ spectra
- **Partial wave mixing** between s- and g-waves.

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Model independent relation: Lüscher's formula

- $\delta_l(E)$ and $E_n(\Gamma, L)$ are related model independently (Up to exponentially small corrections)
- It's a **one-to-one** relation only in the simplest cases
 - $l_{\text{cut}} = 0$: $E_n^{\text{lat}}(\mathbf{A}_1^+, L) = 100 \text{ MeV}$
 - \rightarrow equation of $\delta_0(100 \text{ MeV})$
 - $\rightarrow \delta_0(100 \text{ MeV})$
- It's **not one-to-one** in general cases
 - $l_{\text{cut}} = 4$: $E_n^{\text{lat}}(\mathbf{A}_1^+, L) = 100 \text{ MeV}$
 - \rightarrow equation of $\delta_0(100 \text{ MeV})$ and $\delta_4(100 \text{ MeV})$
 - Need more than one energy levels with the **same** energy value
- Most data fail to find their partners to apply Lüscher's formula
- Fitting process is necessary (to relate quantities at **different** energy values)
- Normal approach: Parametrize the phase shifts

M. Lüscher:

Commun. Math. Phys., 105:153-188, 1986.

Commun. Math. Phys., 104:177, 1986.

Nucl. Phys., B354:531-578, 1991.

What is Hamiltonian effective field theory (HEFT)

- Build the infinite and finite volume Hamiltonians for a system

$$\text{Infinite: } \hat{H} = \hat{H}_0 + \hat{V}$$

$$\text{Finite: } \hat{H}_L = \hat{H}_{0L} + \hat{V}_L$$

- Parametrize the potentials of the Hamiltonians

$$\hat{V} \leftarrow \text{parameters} \rightarrow \hat{V}_L$$

- Fit the eigenvalues of \hat{H}_L to $E_n^{\text{lat}}(\Gamma, L)$ to set the parameters
- Use \hat{H} to calculate $\delta_l(E)$, which are potential model independent
- Use \hat{H} and \hat{H}_L to study the system

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The infinite volume potential

- We start with the plane wave states: $|\mathbf{k}\rangle$
- We should deal with $|k; l, m\rangle := \int d\Omega_{\hat{\mathbf{k}}} Y_{lm}(\hat{\mathbf{k}}) |\mathbf{k}\rangle$
- Rotational invariance only allows interaction like $|k'; l, m\rangle \langle k; l, m|$
- Wigner-Eckart theorem \rightarrow different m correspond to the same v_l

$$\hat{V} = \int \frac{k'^2 dk'}{(2\pi)^3} \int \frac{k^2 dk}{(2\pi)^3} \sum_{l,m} v_l(k', k) |k'; l, m\rangle \langle k; l, m|$$

How does potential change from infinite to finite volume

$$\hat{V} = \int \frac{k'^2 dk'}{(2\pi)^3} \int \frac{k^2 dk}{(2\pi)^3} \sum_{l,m} v_l(k', k) |k'; l, m\rangle \langle k; l, m|$$

- $\mathbf{k} \rightarrow \frac{2\pi}{L}\mathbf{n}$ $\int_0^\infty \frac{k^2 dk}{(2\pi)^3} \rightarrow \sum_{N=0}^\infty$ N denotes $|\mathbf{n}|^2$
- $v_l(k', k) \rightarrow \tilde{v}_l(k_{N'}, k_N) := \frac{v_l(k_{N'}, k_N)}{4\pi L^3}$ k_N denotes $\frac{2\pi}{L}\sqrt{N}$
- $|k; l, m\rangle := \int d\Omega_{\hat{\mathbf{k}}} Y_{lm}(\hat{\mathbf{k}}) |\mathbf{k}\rangle$
 $\Rightarrow |N; l, m\rangle := \sqrt{4\pi} \sum_{|\mathbf{n}|^2=N} Y_{lm}(\hat{\mathbf{n}}) |\mathbf{n}\rangle$

$$\hat{V}_L = \sum_{N'} \sum_N \sum_{l,m} \tilde{v}_l(k_{N'}, k_N) |N'; l, m\rangle \langle N; l, m|$$

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How does potential change from infinite to finite volume

$$\hat{V} = \int \frac{k'^2 dk'}{(2\pi)^3} \int \frac{k^2 dk}{(2\pi)^3} \sum_{l,m} v_l(k', k) |k'; l, m\rangle \langle k; l, m|$$

- $\mathbf{k} \rightarrow \frac{2\pi}{L}\mathbf{n}$ $\int_0^\infty \frac{k^2 dk}{(2\pi)^3} \rightarrow \sum_{N=0}^\infty$ N denotes $|\mathbf{n}|^2$
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- $|k; l, m\rangle := \int d\Omega_{\hat{\mathbf{k}}} Y_{lm}(\hat{\mathbf{k}}) |\mathbf{k}\rangle$
 $\implies |N; l, m\rangle := \sqrt{4\pi} \sum_{|\mathbf{n}|^2=N} Y_{lm}(\hat{\mathbf{n}}) |\mathbf{n}\rangle$

$$\hat{V}_L = \sum_{N'} \sum_N \sum_{l,m} \tilde{v}_l(k_{N'}, k_N) |N'; l, m\rangle \langle N; l, m|$$

How are partial waves mixed

$$\hat{V}_L = \sum_{N'} \sum_N \sum_{l,m} \tilde{v}_l(k_{N'}, k_N) |N'; l, m\rangle \langle N; l, m|$$

- (l, m) is no longer a good quantum number in the finite volume
- $|N; l, m\rangle$ belong to the reducible representations of O_h
- $|N; l, m\rangle$ have components of states belonging to the irreps of O_h

$$\mathbf{0}^+ = \mathbf{A}_1^+$$

$$\mathbf{1}^- = \mathbf{T}_1^-$$

$$\mathbf{2}^+ = \mathbf{E}^+ \oplus \mathbf{T}_2^+$$

$$\mathbf{3}^- = \mathbf{A}_2^- \oplus \mathbf{T}_1^- \oplus \mathbf{T}_2^-$$

$$\mathbf{4}^+ = \mathbf{A}_1^+ \oplus \mathbf{E}^+ \oplus \mathbf{T}_1^+ \oplus \mathbf{T}_2^+$$

- v_l are coupling to those states, hence coupling to the spectra of those irreps
- $\mathbf{0}^+$ and $\mathbf{4}^+$ can be mixed in \mathbf{A}_1^+

How are partial waves mixed

- The infinite volume states $|k; l, m\rangle$ are orthogonal to each other

$$\langle k'; l', m' | k; l, m \rangle \propto \int d\Omega_{\hat{\mathbf{k}}} Y_{l'm'}^*(\hat{\mathbf{k}}) Y_{lm}(\hat{\mathbf{k}}) = \delta_{l',l} \delta_{m',m}$$

- The finite volume states $|N; l, m\rangle$ are not

$$\langle N'; l', m' | N; l, m \rangle = 4\pi \sum_{|\mathbf{n}|^2=N} Y_{l'm'}^*(\hat{\mathbf{n}}) Y_{lm}(\hat{\mathbf{n}})$$

- $|N; l, m\rangle$ with different (l, m) can contain an overlap of their components \implies they are mixed

P-Matrix: Measure the degree of partial wave mixing

- $[P_N]_{l',m';l,m} := \langle N; l', m' | N; l, m \rangle = 4\pi \sum_{|\mathbf{n}|^2=N} Y_{l'm'}^*(\hat{\mathbf{n}}) Y_{lm}(\hat{\mathbf{n}})$
- $[P_N]_{0,0;0,0} = \langle N; 0, 0 | N; 0, 0 \rangle = \sum_{|\mathbf{n}|^2=N} \mathbf{1} = C_3(N)$
- Recovery of spherical symmetry:

$$L \rightarrow \infty \text{ with } \frac{2\pi}{L} \sqrt{N} \text{ fixed} \quad \Rightarrow \quad N \rightarrow \infty$$

$$\sum_{|\mathbf{n}|^2=N} \rightarrow \int d\Omega \quad \Rightarrow \quad [P_N]_{l',m';l,m} \rightarrow C_3(N) \delta_{l',l} \delta_{m',m}$$

P-Matrix: Measure the degree of partial wave mixing

s-wave		p-wave		d-wave		f-wave		g-wave	
1.00	0	0	0	0	0	0	0	0	1.05
0	1.00	0	0	0	0	0	-0.94	0	0
0	0	1.00	0	0	0	0	1.53	0	0
0	0	0	1.00	0	0	-1.21	0	0	-1.08
0	0	0	0	1.25	0	0	0	0	0
0	0	0	0	0	2.50	0	0	0	0
0	0	0	0	0	0	0	0	1.40	-1.17
0	0	0	0	0	1.25	0	0	0	-1.08
0	0	0	-1.21	0	0	0	1.46	0	0
0	0	0	0	0	0	0	1.13	0	0
0	-0.94	0	0	0	0	0	0.88	0	0
0	0	1.53	0	0	0	0	2.33	0	0
0	0	0	-0.94	0	0	0	1.13	0	0
0	0	0	0	0	0	0	0.88	0	0
0	0	0	-1.21	0	0	0	0	0	0
1.05	0	0	0	0	-1.17	0	0	0	1.18
0	0	0	0	0	0	0	0	0	0
0	0	0	-1.08	0	0	-1.08	0	0	0.94
0	0	0	0	0	0	0	0	0	0.94
1.75	0	0	0	0	1.40	0	0	0	3.84
0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	-1.08	0	0	0	0.94
0	0	0	0	0	0	-1.08	0	0	0.94
1.05	0	0	0	0	-1.17	0	0	0	0

$$[P_{N=1}] \Big/ C_3(1)$$

$$C_3(1) = 6$$

$$[P_N]_{l',m';l,m} := \langle N; l', m' | N; l, m \rangle = 4\pi \sum_{|\mathbf{n}|^2=N} Y_{l'm'}^*(\hat{\mathbf{n}}) Y_{lm}(\hat{\mathbf{n}})$$

25 × 25 matrix ordered as $(l, m) = (0, 0), (1, -1), (1, 0), (1, 1), \dots, (4, 4)$

P-Matrix: Measure the degree of partial wave mixing

s-wave		p-wave		d-wave		f-wave		g-wave		
1.00	0	0	0	0	0	0	0	0	0	0.07
0	1.00	0	0	0	0	0	-0.06	0	0	-0.08
0	0	1.00	0	0	0	0	0.10	0	0	0
0	0	0	1.00	0	0	0	-0.06	0	0	0
0	0	0	0	1.02	0	0	0.08	0	0	-0.07
0	0	0	0	0	0.94	0	0	0	0	-0.01
0	0	0	0	0	1.10	0	0	0	0	0.10
0	0	0	0	0	0.94	0	0	0	0	-0.01
0	0	0	0	0.08	0	0	1.02	0	0	-0.07
0	0	0	-0.08	0	0	0	1.03	0	0	0
0	0	0	0	0	0	0.94	0	0	0	-0.03
0	-0.06	0	0	0	0	0	0.98	0	0	0.09
0	0	0.10	0	0	0	0	0	1.10	0	0
0	0	0	-0.06	0	0	0	0.09	0	0	0.98
0	0	0	0	0	0	-0.03	0	0	0	0
0	-0.08	0	0	0	0	0	0.09	0	0	1.03
0.07	0	0	0	0	-0.09	0	0	0	0.04	0.11
0	0	0	0	0	0.03	0	0	0	0.93	0
0	0	0	-0.07	0	0	-0.09	0	0	0	1.03
0	0	0	0	-0.01	0	0	0	0	0.85	0
0.11	0	0	0	0	0.10	0	0	0.11	0	1.30
0	0	0	0	0	-0.01	0	0	0	-0.04	0.85
0	0	0	0	-0.09	0	0	-0.07	0	0	0
0	0	0	0	0.03	0	0	0	0	-0.04	1.03
0.07	0	0	0	0	-0.09	0	0	0.20	0	0.93
									0.11	1.04

$$[P_{N=581}] \Big/ C_3(581)$$

$$C_3(581) = 336$$

$$[P_N]_{l',m';l,m} := \langle N; l', m' | N; l, m \rangle = 4\pi \sum_{|\mathbf{n}|^2=N} Y_{l'm'}^*(\hat{\mathbf{n}}) Y_{lm}(\hat{\mathbf{n}})$$

25 × 25 matrix ordered as $(l, m) = (0, 0), (1, -1), (1, 0), (1, 1), \dots, (4, 4)$

P-Matrix: Measure the degree of partial wave mixing

s-wave		p-wave		d-wave		f-wave		g-wave		
1.00	0	0	0	0	0	0	0	-0.01	0	0
0	1.00	0	0	0	0	0	0.01	0	0	0
0	0	1.00	0	0	0	0	-0.02	0	0	0
0	0	0	1.00	0	0	0.02	0	0	0	0
0	0	0	0	1.00	0	0	0.01	0	0	0
0	0	0	0	0	1.01	0	0	0.01	0	0
0	0	0	0	0	0.98	0	0	0.00	0	0
0	0	0	0	0	1.01	0	0	-0.02	0	0.02
0	0	0	0	0	-0.02	0	0	0.00	0	-0.00
0	0	0	0	0.02	0	0	1.00	0	0	0
0	0	0	0	0.99	0	0	-0.02	0	0	0
0	0	0	0	0	1.01	0	0	0.00	0	0
0	0.01	0	0	0	0	0	1.00	0	-0.02	0
0	0	-0.02	0	0	0	0	0.98	0	0	0
0	0	0.01	0	0	0	-0.02	0	1.00	0	0
0	0	0	0	0	0.00	0	0	1.01	0	0
0	0.02	0	0	0	0	-0.02	0	0	0.99	0
-0.01	0	0	0	0	0.02	0	0	0	0.99	-0.02
0	0	0	0	0	-0.00	0	0	0	1.01	0
0	0	0	0	0.01	0	0	0	0	0	0.00
0	0	0	0	0.00	0	0	0	0	1.00	0
-0.02	0	0	0	-0.02	0	0	0	-0.02	0	0.00
0	0	0	0	0.00	0	0	0	0.00	0	1.02
0	0	0	0	0.02	0	0	0	0	0.95	0
0	0	0	0	-0.00	0	0	0	0	-0.01	1.02
-0.01	0	0	0	0.02	0	0	0	0	0	1.00

$$[P_{N=941}] \Big/ C_3(941)$$

$$C_3(941) = 552$$

$$[P_N]_{l',m';l,m} := \langle N; l', m' | N; l, m \rangle = 4\pi \sum_{|\mathbf{n}|^2=N} Y_{l'm'}^*(\hat{\mathbf{n}}) Y_{lm}(\hat{\mathbf{n}})$$

25 × 25 matrix ordered as $(l, m) = (0, 0), (1, -1), (1, 0), (1, 1), \dots, (4, 4)$

P-Matrix: Measure the degree of partial wave mixing

- P-Matrix also respects the cubic symmetry
- $m \rightarrow (\Gamma, f, \alpha) \implies |N; l, m\rangle \rightarrow |N, l; \Gamma, f, \alpha\rangle$

$$|N, l; \Gamma, f, \alpha\rangle = \sum_m [C_l]_{\Gamma, f, \alpha; m} |N; l, m\rangle$$

- P-Matrix in new basis is obtained through a unitary transformation
 - The transformation matrix is given by $[C] = \text{diag}([C_0], [C_1], \dots)$
 - P-Matrix will be block diagonal according to the irreps of O_h

$$[C][P_N][C]^\dagger \rightarrow \delta_{\Gamma', \Gamma} \delta_{\alpha', \alpha} \langle N, l'; \Gamma, f', \alpha | N, l; \Gamma, f, \alpha \rangle$$

How is the dimension of the Hamiltonian matrix reduced

- Original basis $|\mathbf{n}\rangle$: $\sum_{N=0}^{N_{\text{cut}}=600} C_3(N) \sim 60,000$
- Basis $|N; l, m\rangle$ with $l_{\text{cut}} = 4$: $\sum_{N=0}^{600} 25 \sim 600 \times 25$
- Basis $|N, l; \Gamma = \mathbf{A}_1^+, f, \alpha\rangle$: $\sum_{N=0}^{600} 2 \sim 600 \times 2$
- Orthonormalization needs the inner products——P-Matrix

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The process

- $l_{\text{cut}} = 4$, only s-, d- and g-waves are present

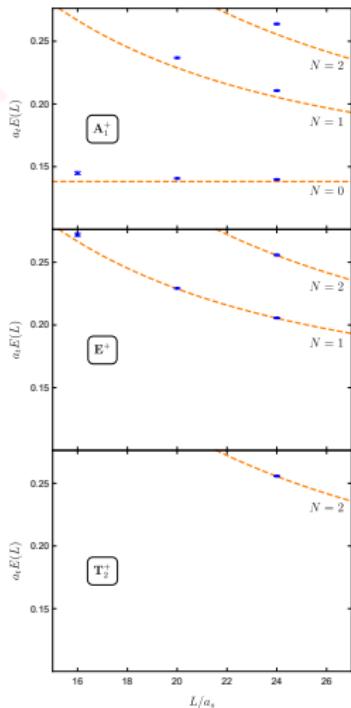
- Separable potential model:

$$v_l(p, k) = f_l(p) G_l f_l(k)$$

$$f_l(k) \sim \frac{(d_l \times k)^l}{(1 + (d_l \times k)^2)^{l/2+2}}$$

- 6 parameters: $G_0, G_2, G_4, d_0, d_2, d_4$
 - Dimensions of Hamiltonians ($N_{\text{cut}} = 600$):
- $\mathbf{A}_1^+ : 923 \quad \mathbf{E}^+ : 965 \quad \mathbf{T}_2^+ : 963$
- The fitted data: 11 energy levels

PRD 86, 034031 (2012)
Jozef J. Dudek et al.

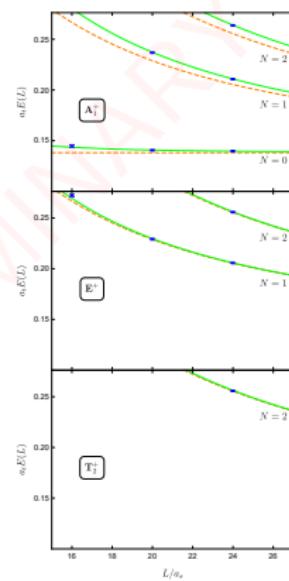


The results

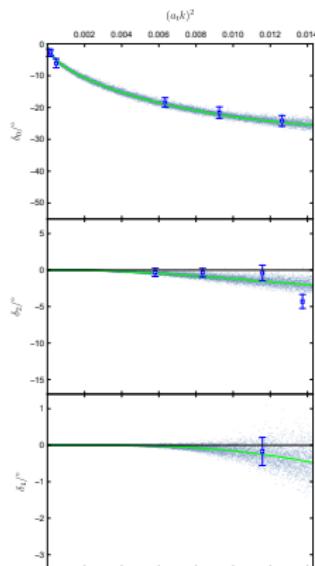
Components of eigenstates

\mathbf{A}_1^+	$N = 0$	$N = 1$	$N = 2$	$N = 3$	$N = 4$	\dots
1st	99.7	0.2	0.0	0.0	0.0	\dots
2nd	0.1	97.4	1.9	0.2	0.0	\dots
3rd	0.0	1.5	94.5	2.8	0.3	\dots

Volume dependent spectra



Phase shifts with errors



- Fitting \rightarrow Parameters $\rightarrow \hat{H}$ and \hat{H}_L
- $\hat{H} \rightarrow \delta_l(E)$
- $\hat{H}_L \rightarrow E_n(\Gamma, L)$
- $\hat{H}_L \rightarrow$ Eigenstates

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Summary

- Partial wave mixing is discussed in HEFT
 - How it happens in the viewpoint of HEFT
 - P-Matrix is defined to measure the degree of partial wave mixing
 - The dimension of the Hamiltonian can be highly reduced
- Example of isospin-2 $\pi\pi$ scattering is discussed

ACKNOWLEDGMENTS

Finite-volume energy levels taken from [PRD 86, 034031 (2012), Jozef J. Dudek et al.] were provided by the Hadron Spectrum Collaboration – no endorsement on their part of the analysis presented in the current paper should be assumed.

Thank you!

Backup

$ \mathbf{n}\rangle$	$N = \mathbf{n} ^2$	0	1	2	3	\dots
$\sum_{\hat{\mathbf{n}}} Y_{lm}(\hat{\mathbf{n}}) \mathbf{n}\rangle$	$ N; l, m\rangle$	1	6	12	8	\dots
		1	∞	∞	∞	\dots

$$N = 0 : (0, 0, 0)$$

$$N = 1 : (+1, 0, 0) \quad (0, +1, 0) \quad (0, 0, +1)$$

$$l = 0, 1, 2, \dots$$

$$m = -l, \dots, +l$$

$$(-1, 0, 0) \quad (0, -1, 0) \quad (0, 0, -1)$$

.....

$ \mathbf{n}\rangle$	$N = \mathbf{n} ^2$	0	1	2	3	\dots
$\sum_{\hat{\mathbf{n}}} Y_{lm}(\hat{\mathbf{n}}) \mathbf{n}\rangle$	$ \mathbf{n}\rangle$ with l_{cut}	1	6	12	8	\dots
		1	∞	∞	∞	\dots
		1	$(l_{\text{cut}} + 1)^2$	$(l_{\text{cut}} + 1)^2$	$(l_{\text{cut}} + 1)^2$	\dots

$|N; l, m\rangle$

$$N = 0 : (0, 0, 0)$$

$$N = 1 : (+1, 0, 0) \quad (0, +1, 0) \quad (0, 0, +1)$$

$$m = -l, \dots, +l$$

$$(-1, 0, 0) \quad (0, -1, 0) \quad (0, 0, -1)$$

.....

$$\sum_{l=0}^{l_{\text{cut}}} (2l+1) = (l_{\text{cut}}+1)^2 = 25 \text{ for } l_{\text{cut}} = 4$$

Dimension: big number $\rightarrow (l_{\text{cut}}+1)^2 \times N_{\text{cut}} + 1$

$\sim 25 \times 600$ for $l_{\text{cut}} = 4$, $N_{\text{cut}} = 600$

