

Global analysis of the $\Delta(1232)$ contribution in the pion photo-production on nucleons

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August 17th, 2019
HADRON 2019, Guilin, China

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Astrid H. Blin,
Deliang Yao.*

Outline

1 Motivation

2 Theoretical approach

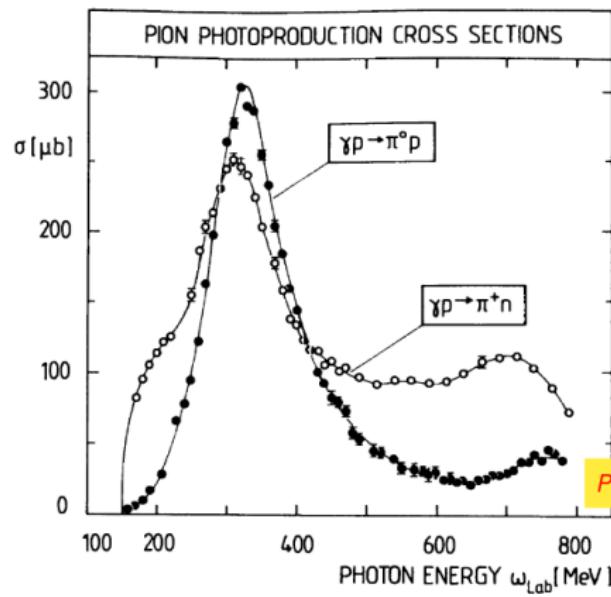
3 Results

4 Outlook

5 Summary

Motivation

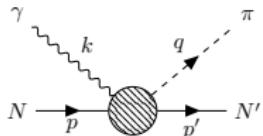
PION PHOTOPRODUCTION ON NUCLEONS



- At the peak the neutral channel has bigger cross section than the charged channel
- Near threshold there are huge cancellations for the neutral channel (not well described at low energy)
- For the charged channel there are not such cancellations
- How much contributes the resonances at low energies for neutral and charged channels ?
- We want to know how to deal properly with this energy region

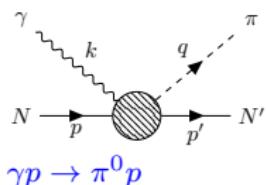
Pions and Nuclei, Ericson & Weise

Experimental data



$$E_\gamma \sim 145 \text{ MeV} - \sim 215 \text{ MeV}$$

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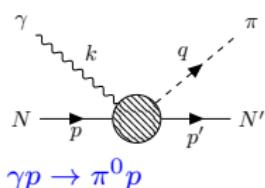


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Measured observables :

$$\sigma \quad , \quad \frac{d\sigma}{d\Omega} \quad , \quad \Sigma = \frac{d\sigma_\perp - d\sigma_\parallel}{d\sigma_\perp + d\sigma_\parallel}, \quad T = \frac{d\sigma_+ - d\sigma_-}{d\sigma_+ + d\sigma_-}$$

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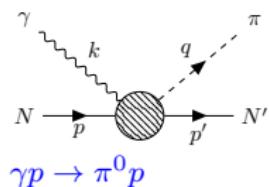
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779 data points

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$$\gamma p \rightarrow \pi^+ n$$

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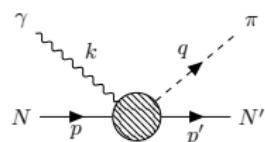
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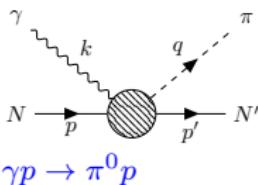
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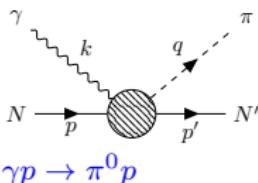
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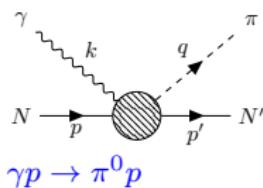
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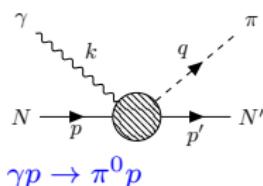
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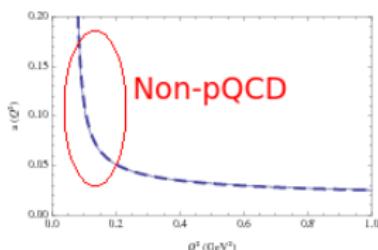
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Nothing yet
0 data points

Approaching the experimental data



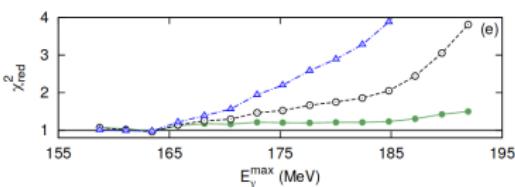
NON-PERTURBATIVE REGIME

- Low energy regime
- $E_\gamma \approx 145 \text{ MeV} - 215 \text{ MeV} \implies \alpha_S >> 1$ perturbative QCD breakdown
 - We need an Effective Theory approach \implies Chiral Perturbation Theory.
- At low energies : Relevant degrees of freedom



- We use as expansion parameters the relative $\frac{p}{\Lambda}$ and pion mass $\frac{m_\pi}{\Lambda}$

Previous works : Theoretical models for $\gamma p \rightarrow \pi^0 p$

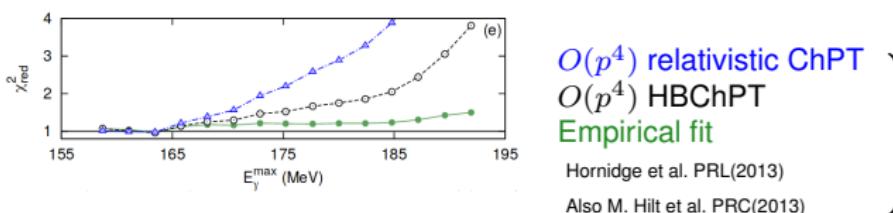


$O(p^4)$ relativistic ChPT
 $O(p^4)$ HBChPT
Empirical fit

Hornidge et al. PRL(2013)
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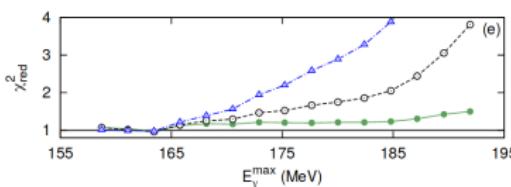
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But including Delta explicitly improved the situation

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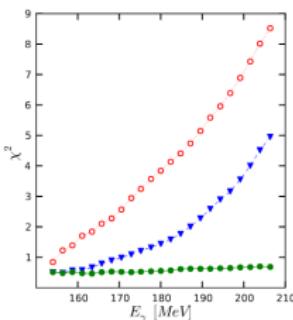
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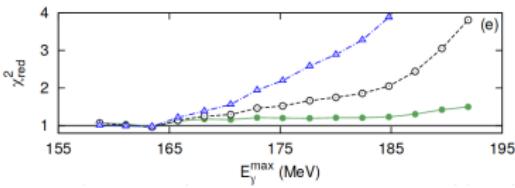
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But including Delta explicitly improved the situation



- $O(p^3)$ relativistic [tree level] ChPT
 - $O(p^3)$ relativistic [tree+loops] ChPT
 - relativistic [tree+loops+ Δ] H. Blin et al. PLB(2015)

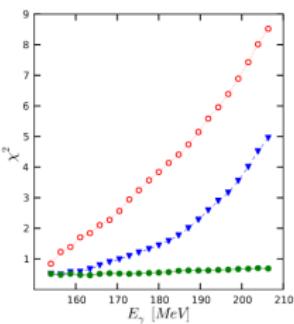
Some problems : Theoretical models for $\gamma p \rightarrow \pi^0 p$



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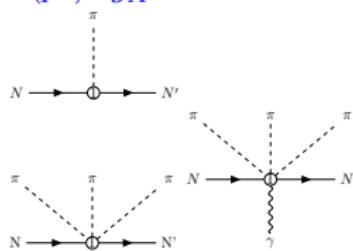
- Our aim is to extend this framework, relativistic ChPT with explicitly $\Delta(1232)$ inclusion, to the charged channels making a global analysis

The nucleon Lagrangian

Each chiral order brings new LECs with it. For this process $\gamma N \rightarrow \pi N'$

$$\mathcal{L}_N^{(1)} = \bar{\Psi} \left(i \not{D} - m + \frac{g_A}{2} \not{\psi} \gamma_5 \right) \Psi,$$

$O(p^1) : g_A$



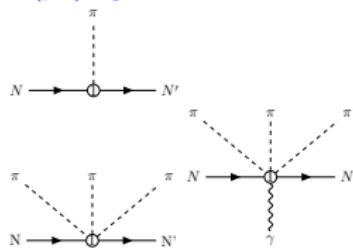
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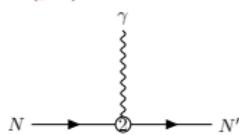
$$\mathcal{L}_N^{(1)} = \bar{\Psi} \left(i \not{D} - m + \frac{g_A}{2} \not{\psi} \gamma_5 \right) \Psi,$$

$$\mathcal{L}_N^{(2)} = \bar{\Psi} \frac{1}{8m} \left(c_6 F_{\mu\nu}^+ + c_7 \text{Tr} [F_{\mu\nu}^+] \right) \sigma^{\mu\nu} \Psi + \dots,$$

$$O(p^1) : g_A$$



$$O(p^2) : c_6, c_7$$

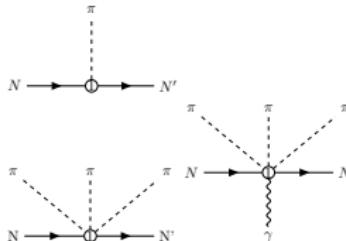


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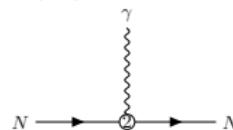
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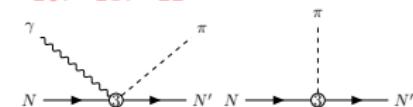
$O(p^1) : g_A$



$O(p^2) : c_6, c_7$



$O(p^3) : d_8, d_9, d_{16}, d_{18}, d_{20}, d_{21}, d_{22}$

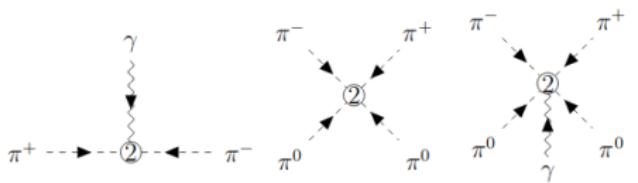


The pion Lagrangian

Up to $O(p^3)$ we do **not need extra fitting Low-Energy-Constants**

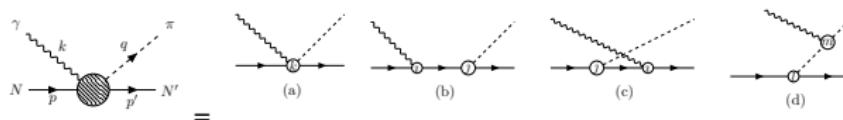
$$\mathcal{L}_{\pi\pi}^{(2)} = \frac{F_\pi^2}{4} \text{Tr} \left[D^\mu U (D_\mu U)^\dagger + \chi U^\dagger + U \chi^\dagger \right].$$

All can be expanded in terms of **no fitting parameters**



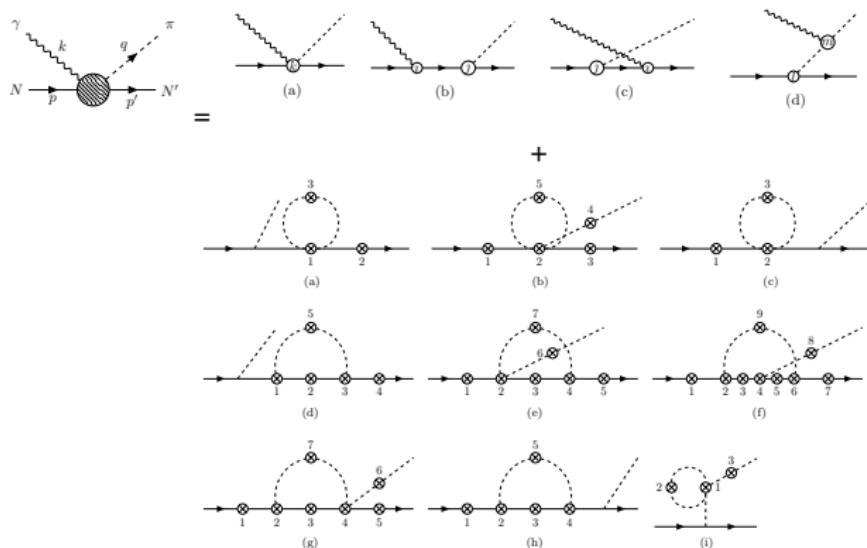
Putting the pieces together

- Using all the ingredients, we are able to calculate all possible diagrams up to $O(p^3)$ order



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- Using all the ingredients, we are able to calculate all possible diagrams up to $O(p^3)$ order



- This approach is the same as used in previous ChPT works

The Strategy

- We keep our expansion up to $O(p^3)$: Avoids inclusion of too many LECs at $O(p^4)$

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Loop diagrams regularization

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Loop diagrams regularization

- Loop diagrams : UV divergences and power counting breaking terms

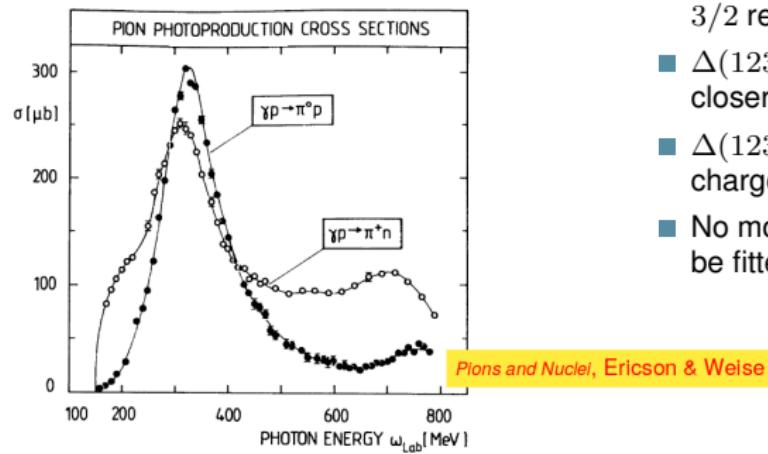
$$\frac{1}{\epsilon_{UV}} = \frac{1}{4-dim} \text{ and terms } \propto p^1, p^2, \text{ in amplitudes at order } O(p^3)$$

- Renormalization $\widetilde{\text{MS}}$ -EOMS :

- $\widetilde{\text{MS}}$: Subtracts multiples of $R = \gamma_E - 1/\epsilon_{UV} - \log(4\pi) - 1$,
- EOMS : Subtracts terms of lower order than the nominal one for loop diagrams

The Strategy to improve the approach

PION PHOTOPRODUCTION ON NUCLEONS

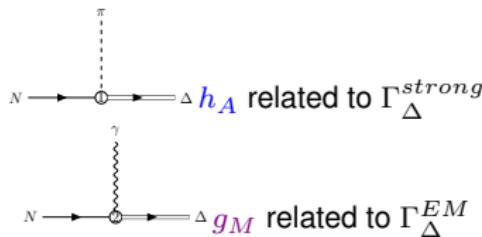


- Explicit inclusion of the $\Delta(1232)$ spin 3/2 resonance
- $\Delta(1232)$ becomes more relevant the closer we are to its mass
- $\Delta(1232)$ inclusion is important in charged and neutral channels
- No more Low-Energy-Constants to be fitted (keep the model simple)

Contributions generated by the $\Delta(1232)$

$$\mathcal{L}_{\Delta\pi N}^{(1)} = \frac{i h_A}{2Fm_\Delta} \bar{\Psi} T^a \gamma^{\mu\nu\lambda} (\partial_\mu \Delta_\nu) \partial_\lambda \pi^a + \text{h.c.},$$

$$\mathcal{L}_{\Delta\gamma N}^{(2)} = \frac{3ie g_M}{2m(m+m_\Delta)} \bar{\Psi} T^3 (\partial_\mu \Delta_\nu) \tilde{f}^{\mu\nu} + \text{h.c.},$$



Only one constrained parameter added, $g_M h_A$

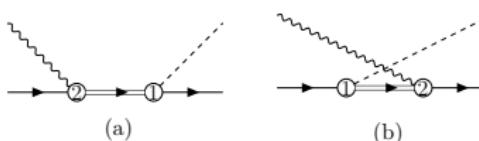
Power Counting Scheme

- For energies close to threshold, far from the $\Delta(1232)$ mass :

Lensky & Pascalutsa EPJ (2010)

$$D = 4L + \sum_{k=1}^{\infty} kV^{(k)} - 2N_{\pi} - N_N - \frac{1}{2}N_{\Delta}.$$

- No loop diagrams with $\Delta(1232)$ up to $O(p^3)$.



- $\Delta(1232)$ is only included at tree level $\rightarrow O(p^{5/2})$ order

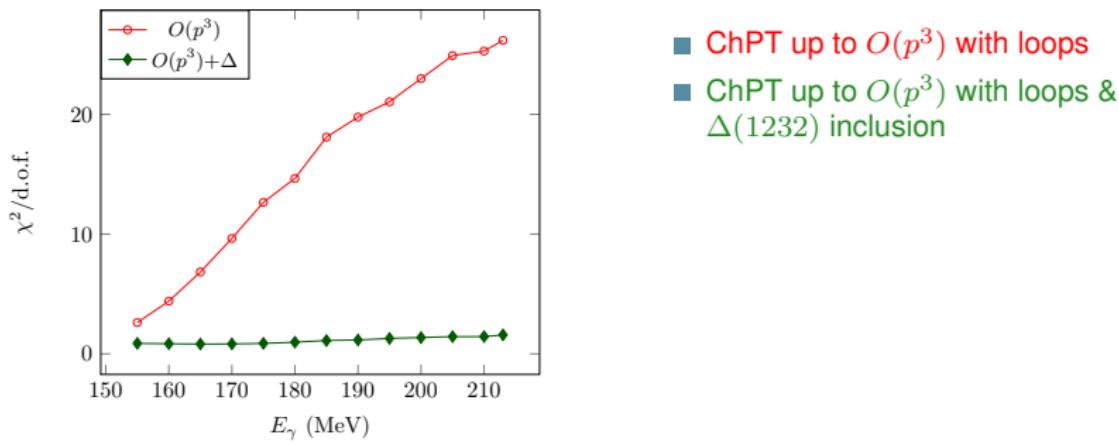
The first message

Fitting LECs using data for all channels

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Fitting LECs using data for all channels

Now we can see easily what can be obtained with the $\Delta(1232)$ inclusion, **without including extra fitting Low-Energy-Constants**



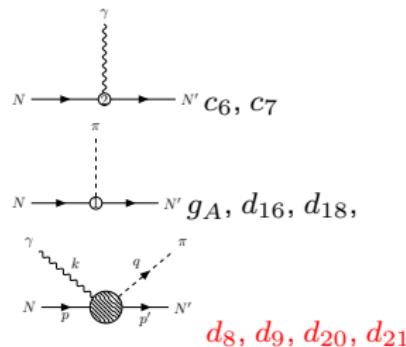
- ChPT up to $O(p^3)$ with loops
- ChPT up to $O(p^3)$ with loops & $\Delta(1232)$ inclusion

What about low-energy constants ?

Fitted LECs from other works using the same approach

LEC	Source
$\mathcal{L}_N^{(2)}$	\tilde{c}_6 \tilde{c}_7
$\mathcal{L}_N^{(3)}$	d_{16}^r
	d_{18}^r
	d_{22}^r
$\mathcal{L}_{\pi N \Delta}^{(1)}$	h_A
$\mathcal{L}_{\pi N \Delta}^{(2)}$	g_M

- c_6 and c_7 directly connected to nucleon magnetic moments
 - d_{16} , related to the πN axial-vector coupling, g_A
 - d_{18} , studied in πN scattering processes (Golderberg-Treiman relation in πN coupling)
 - We can fit few LECs appearing in $\gamma N \rightarrow \pi N'$



Preliminary results

TABLE – LECs and χ^2 for calculations at different chiral orders. Bold numbers are fixed and depend only on physical quantities such as g_A , the proton and neutron magnetic moment and the Δ decay width. Thus, they are not fitted to the pion photoproduction process.

LECs	$O(p^1)$	$O(p^2)$	$O(p^{5/2})$	$O(p^3)$, Fit I
g	1.27	1.27	1.27	1.11
c_6	-	3.706	3.706	5.07
c_7	-	-1.913	-1.913	-2.68
d_{18}	-	-	-	0.60
d_{22}	-	-	-	0.96
$d_8 + d_9$	-	-	-	1.16 ± 0.01
$d_8 - d_9$	-	-	-	1.02 ± 0.13
d_{20}	-	-	-	14.9 ± 2.5
d_{21}	-	-	-	-2.65 ± 0.18
h_A	-	-	2.87	2.87
g_M	-	-	3.16	2.90 ± 0.01
χ^2_{TOT}/dof	165.	310.	60.7	3.25
$\chi^2_{\pi^0}/dof$	208.	392.	76.6	3.58
$\chi^2_{\pi^+}/dof$	10.7	9.15	2.88	1.76
$\chi^2_{\pi^-}/dof$	5.73	6.29	2.51	2.49

Preliminary results

TABLE – The values of the LECs are dimensionless for g_M , in units of GeV^{-2} for d 's, and in units of GeV^{-3} for e_{48} . Fit I refers to the standard setting, Fit II removes Δ mechanisms, Fit III leaves d_{18} free and Fit IV includes an $O(p^4)$ piece (e_{48})

LECs	Fit I	Fit II - Δ	Fit III	Fit IV
$d_8 + d_9$	1.16 ± 0.01	3.53 ± 0.01	0.98 ± 0.02	0.90 ± 0.01
$d_8 - d_9$	1.02 ± 0.14	1.84 ± 0.24	1.72 ± 0.13	-0.09 ± 0.15
d_{18}	0.60	-1.00	5.40 ± 0.13!!	0.60
d_{20}	14.9 ± 2.5	-17.6 ± 2.4	29.7 ± 2.6	6.94 ± 2.5
d_{21}	-2.65 ± 0.18	0.01 ± 0.17	-1.52 ± 0.19	-2.46 ± 0.18
g_M	2.90 ± 0.01	-	3.13 ± 0.02	3.20 ± 0.01
e_{48}	-	-	-	1.97 ± 0.04
χ^2_{TOT}/dof	3.25	30.0	1.59	1.58
$\chi^2_{\pi^0}/dof$	3.58	37.2	1.33	1.48
$\chi^2_{\pi^+}/dof$	1.76	3.66	2.40	1.67
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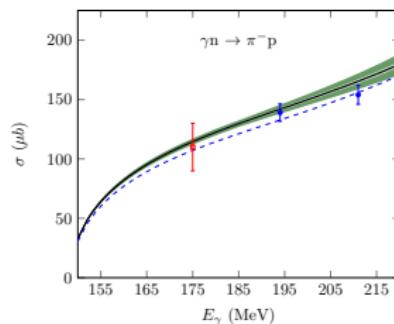
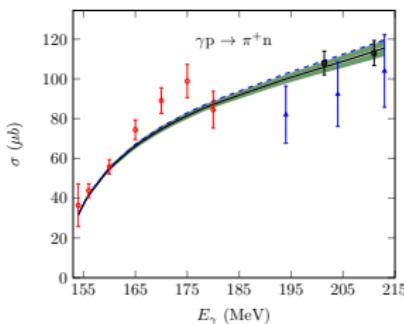
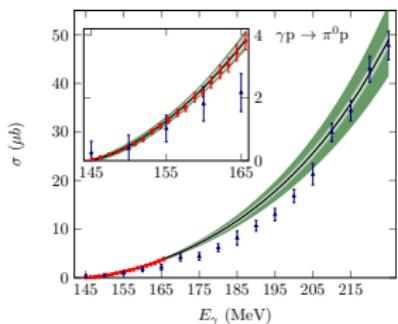
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Comparing theoretical results vs experimental data

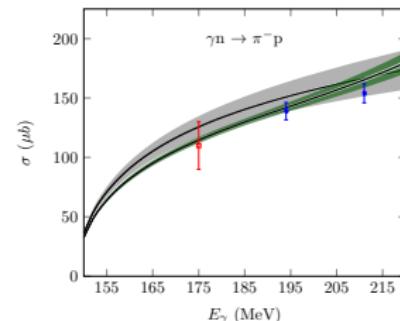
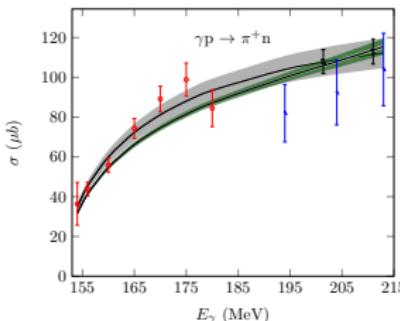
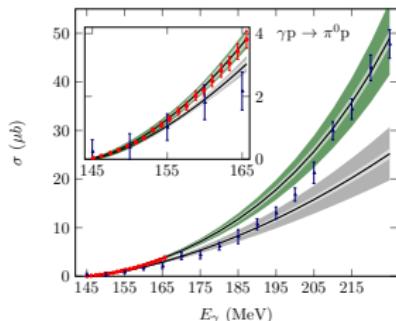
Cross sections



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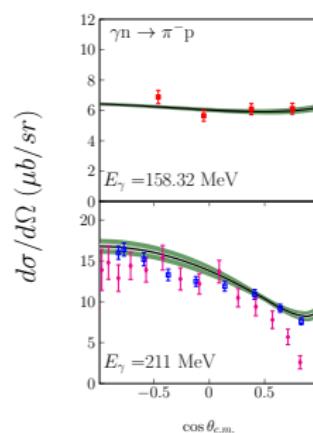
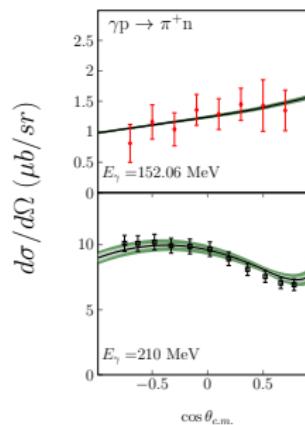
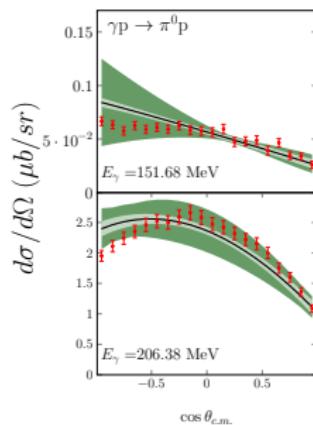
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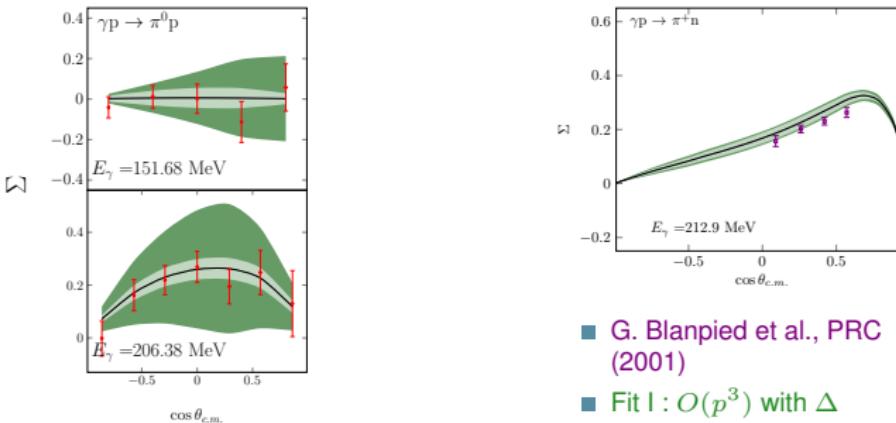
- Hornidge et al., (MAMI)
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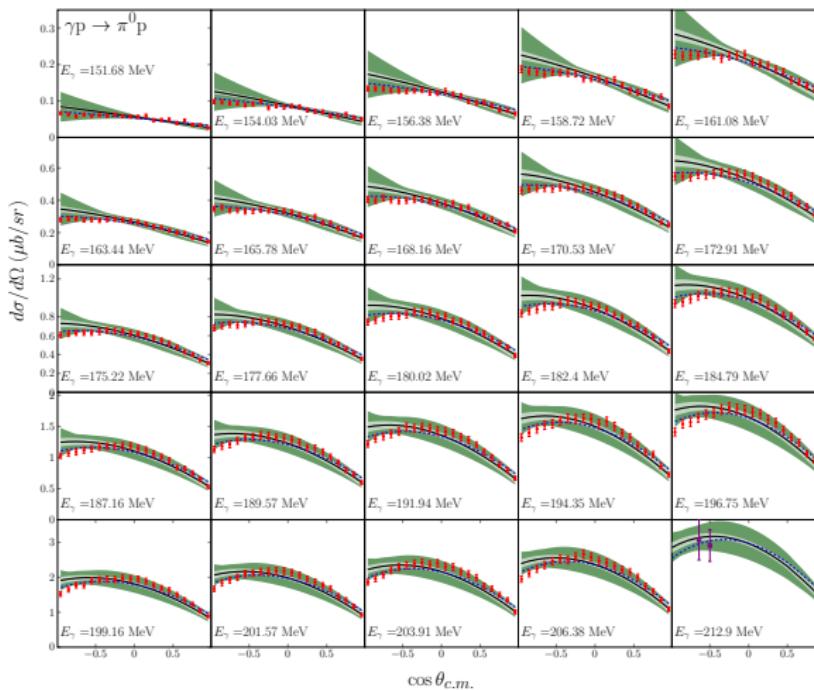
Beam asymmetries



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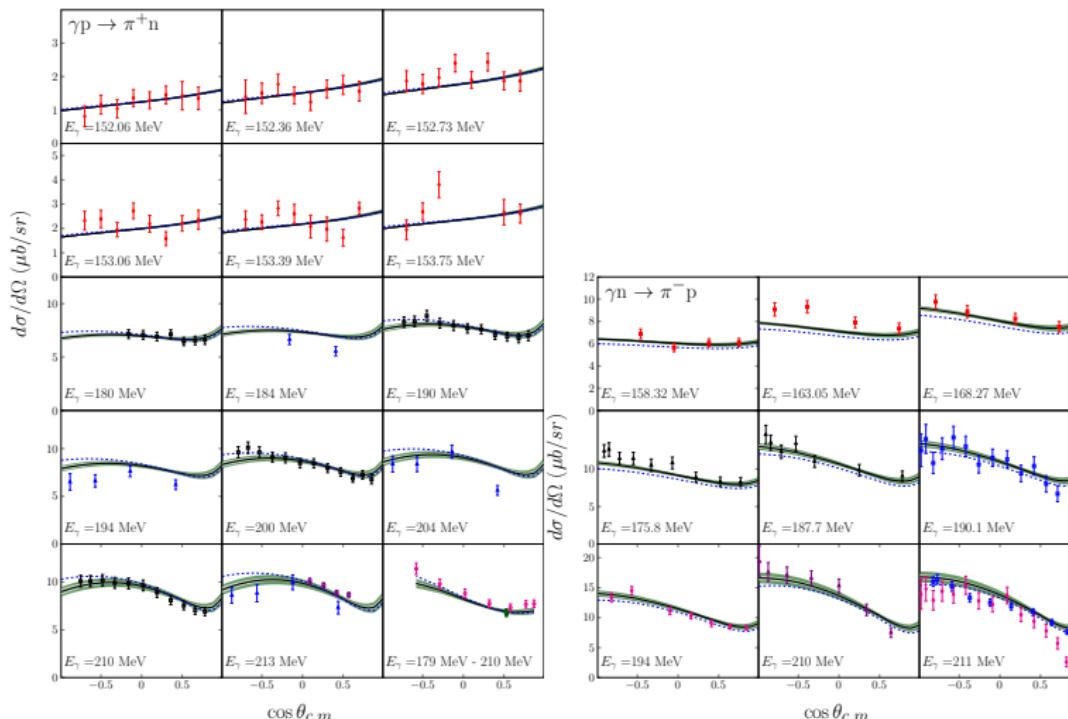
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Differential cross sections



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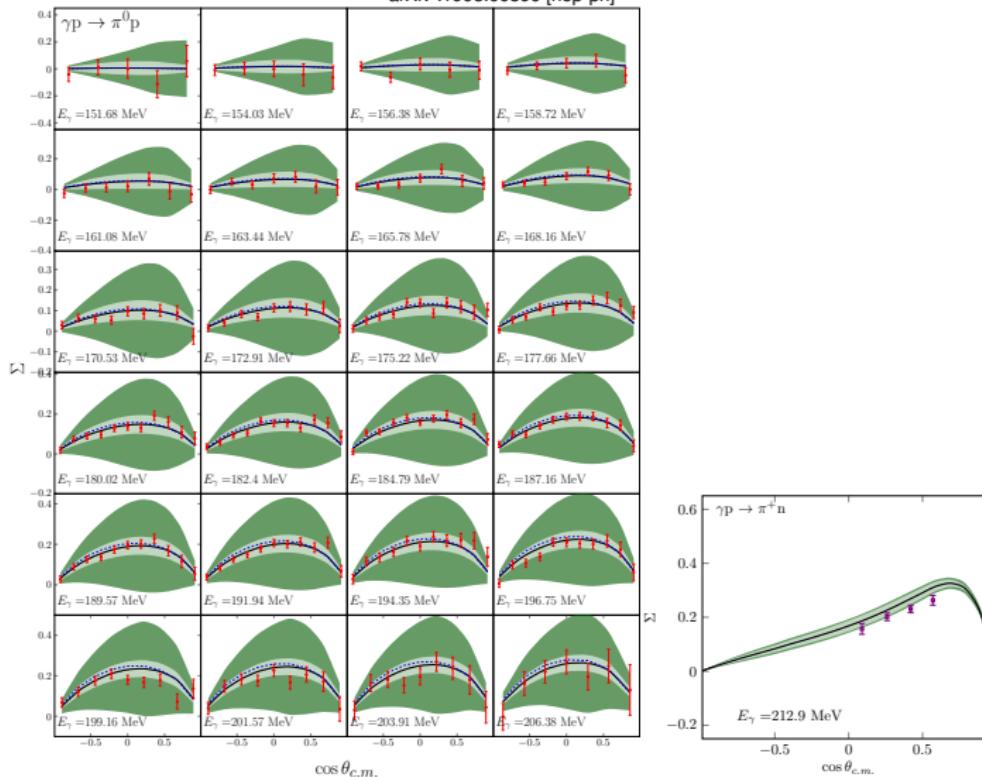
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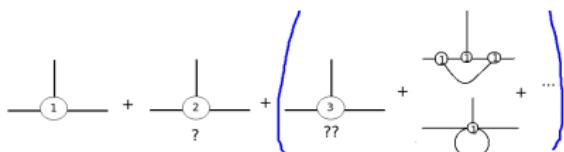
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What is the next?

- Other reactions :
 π -electro-production off nucleons $\gamma^* N \rightarrow \pi N'$
 - Trying to simplify the model by the $\Delta(1232)$ inclusion. No need to reach higher orders at this energy (This would add many unknown LECs)
 - Increase the prediction capability and useful accuracy of this approach for other processes



Fit LECs for first → then make predictions for other processes

Other processes in the same framework ?

Nucleon magnetic moments and Δ EM decays

$$N \rightarrow N' \gamma, \quad \Delta \rightarrow N \gamma$$

Pion Photoproduction (Prediction/LECs Fitting)

$$\gamma N \rightarrow \pi N'$$

$\chi^2/dof = 3.25$ with Δ , $\chi^2/dof = 30.0$ Δ

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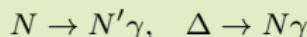
$$eN \rightarrow \pi N'$$

Weak Pion production → Neutrino high precision processes (Predictions)

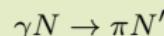
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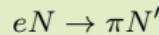


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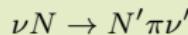


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 - Wider agreement with data near Δ mass using $O(p^3)$ calculation with Δ inclusion than previous higher order $O(p^4)$ calculation without Δ .
 - Even the inclusion of a single tree level order $O(p^4)$ piece can improve the results being good enough as a full $O(p^4)$ calculation can do with many parameters.

Thank you

Error estimations

$$\delta \mathcal{O}_{LEC} = \left(\sum_{i,j} [\text{Corr}(i,j)] \frac{\partial \mathcal{O}(\bar{x}_i)}{\partial x_i} \delta x_i \frac{\partial \mathcal{O}(\bar{x}_j)}{\partial x_j} \delta x_j \right)^{1/2},$$

$$\delta \mathcal{O}_{th}^{(n)} = \max \left(\left| \mathcal{O}^{(n_{LO})} \right| Q^{n-n_{LO}+1}, \left\{ \left| \mathcal{O}^{(k)} - \mathcal{O}^{(j)} \right| Q \right\} \right)$$

where $Q = m_\pi/\Lambda_b$, Λ_b is the breakdown scale of the chiral expansion. We have $\Lambda_b = 4\pi F_\pi \sim 1 \text{ GeV}$.

Other fitted LECs

TABLE – Values of the LECs determined from other processes.

LEC	Value	Source
$\mathcal{L}_N^{(2)}$	5.07 ± 0.15	μ_p and μ_n [Bauer :2012pv, Yao :2018pzc, PDG :2016]
	-2.68 ± 0.08	μ_p and μ_n [Bauer :2012pv, Yao :2019avf, PDG :2016]
d_{18}	$-0.20 \pm 0.80 \text{ GeV}^{-2}$	πN scattering [Alarcon :2012kn]
	$5.20 \pm 0.02 \text{ GeV}^{-2}$	$\langle r_A^2 \rangle_N$ [Yao :2017fym]
$\mathcal{L}_{\pi N \Delta}^{(1)}$	h_A	$\Gamma_\Delta^{\text{strong}}$ [Bernard :2012hb]
$\mathcal{L}_{\pi N \Delta}^{(2)}$	g_M	$\Gamma_\Delta^{\text{EM}}$ [Blin :2015era]

Chiral Perturbation Theory

χ PT

- At large distances (low energy) we are encouraged to use baryons and light mesons, instead of quarks and gluons, as degrees of freedom
 - We use pion mass in the same spirit of the Chiral Symmetry Breaking in QCD with quark masses



- We use as expansion parameters the relative $\frac{p}{\Lambda}$ and pion mass $\frac{m_\pi}{\Lambda}$
 - We can construct Lagrangians that preserves Chiral Symmetry for each order in expansion.
 - Appears Low Energy Constants (should be extracted from QCD, but QCD cannot be solved)

