

Large v_2 with small cross section

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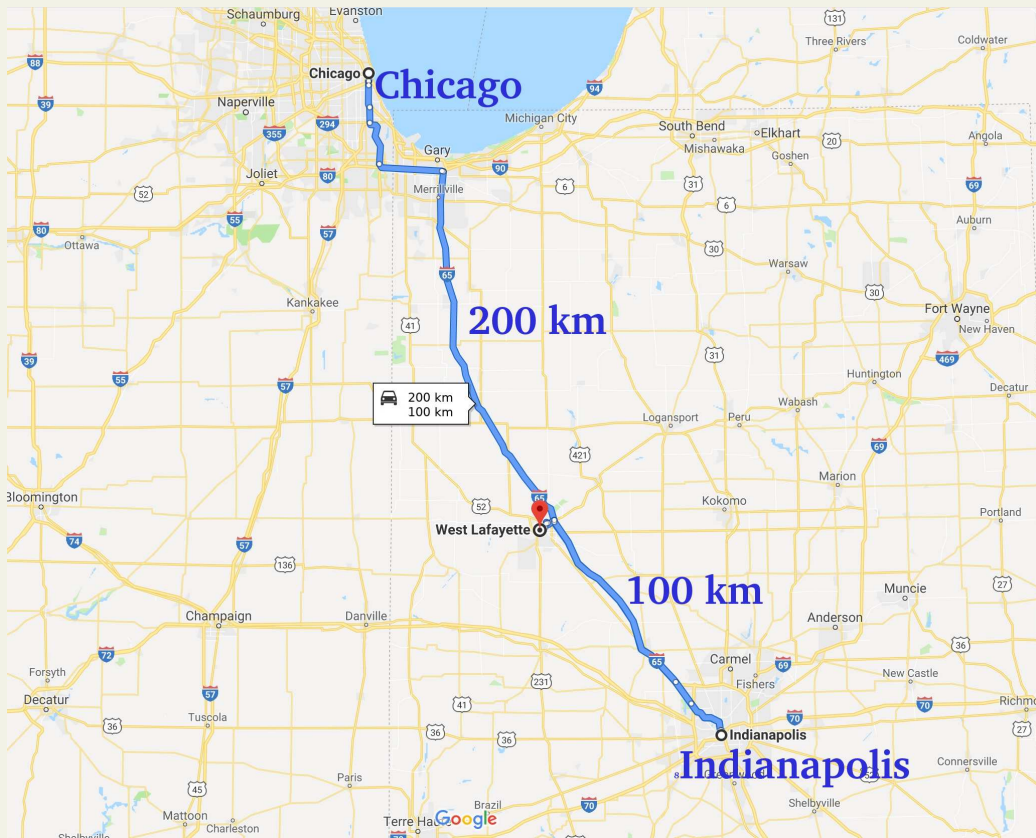
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Purdue University

→ www.purdue.edu

~ 2 hrs to Chicago (QM2017)

~ 1 hr to Indianapolis (F1 racing)

~ 50 000 students

~ 150 physics grad students

~ 60 physics professors



High-E Nuclear Physics Group

→ Fuqiang Wang, Wei Xie, DM

Outline

- I. AMPT puzzle (hydro vs small cross sections)**
- II. Covariant transport and MPC**
- III. Understanding AMPT through MPC comparisons**
- IV. Summary - how small cross sections can work**

A Multi-Phase Transport

Lin, Ko et al, PRC72 ('05)

full-fledged multicomponent event generator

**AMPT \approx Lund string model (HIJING)
+ $2 \rightarrow 2$ parton cascade (ZPC)
+ hadron transport (ART)**

version with “string melting”:

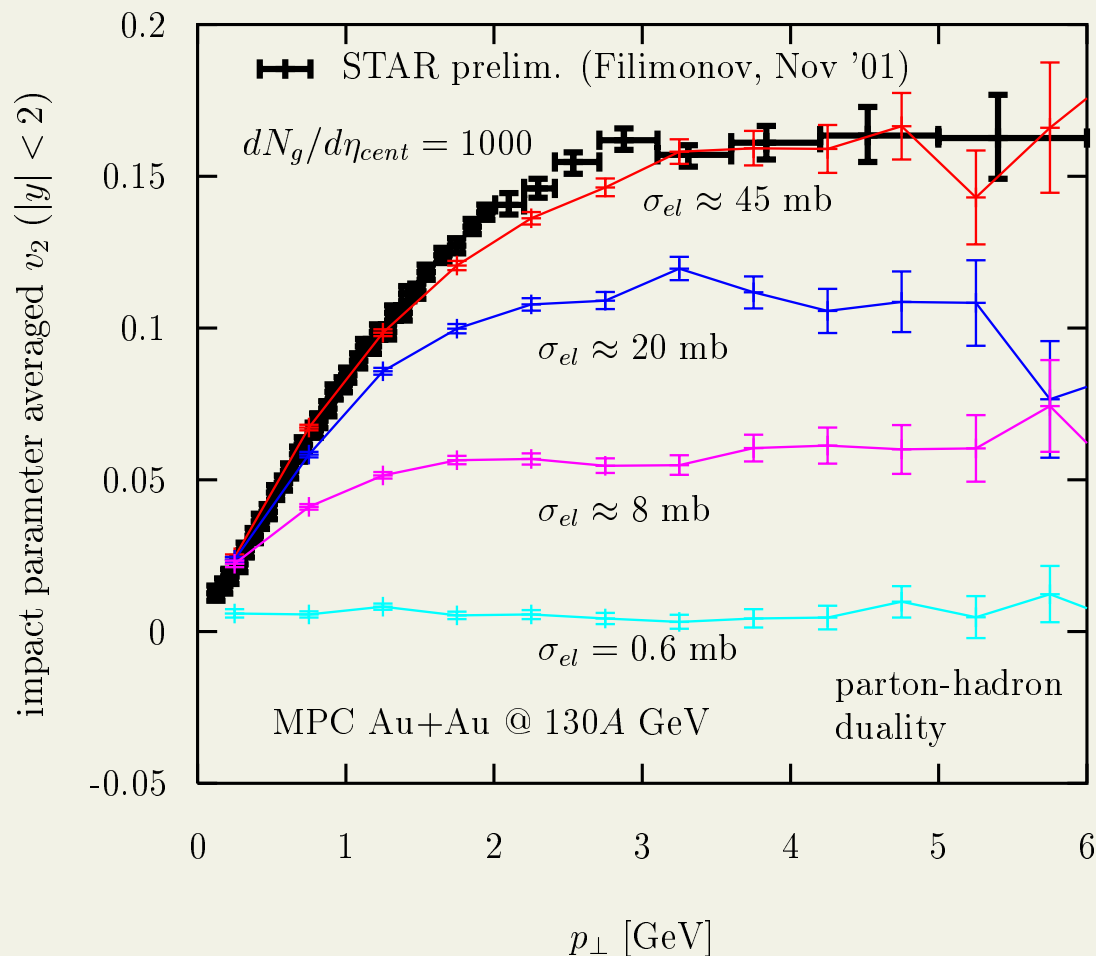
- **energy density in strings converted to quanta (quarks/antiquarks)**
- **\rightarrow fluctuating initial geometry (random nucleon positions)**
- **hadronization via coalescence**

explains quite well a wide range of A+A observables

- **using small ~ 3 mb partonic cross sections (!?)**

Parton opacity puzzle

DM & Gyulassy, NPA 697 ('02): $v_2(p_T, \chi)$ in Au+Au at RHIC



parton transport model MPC

$2 \rightarrow 2$ only, forward-peaked

$\sigma_{tr} \approx 0.3\sigma_{tot}$

Au+Au @ 130 GeV, $b = 8 \text{ fm}$

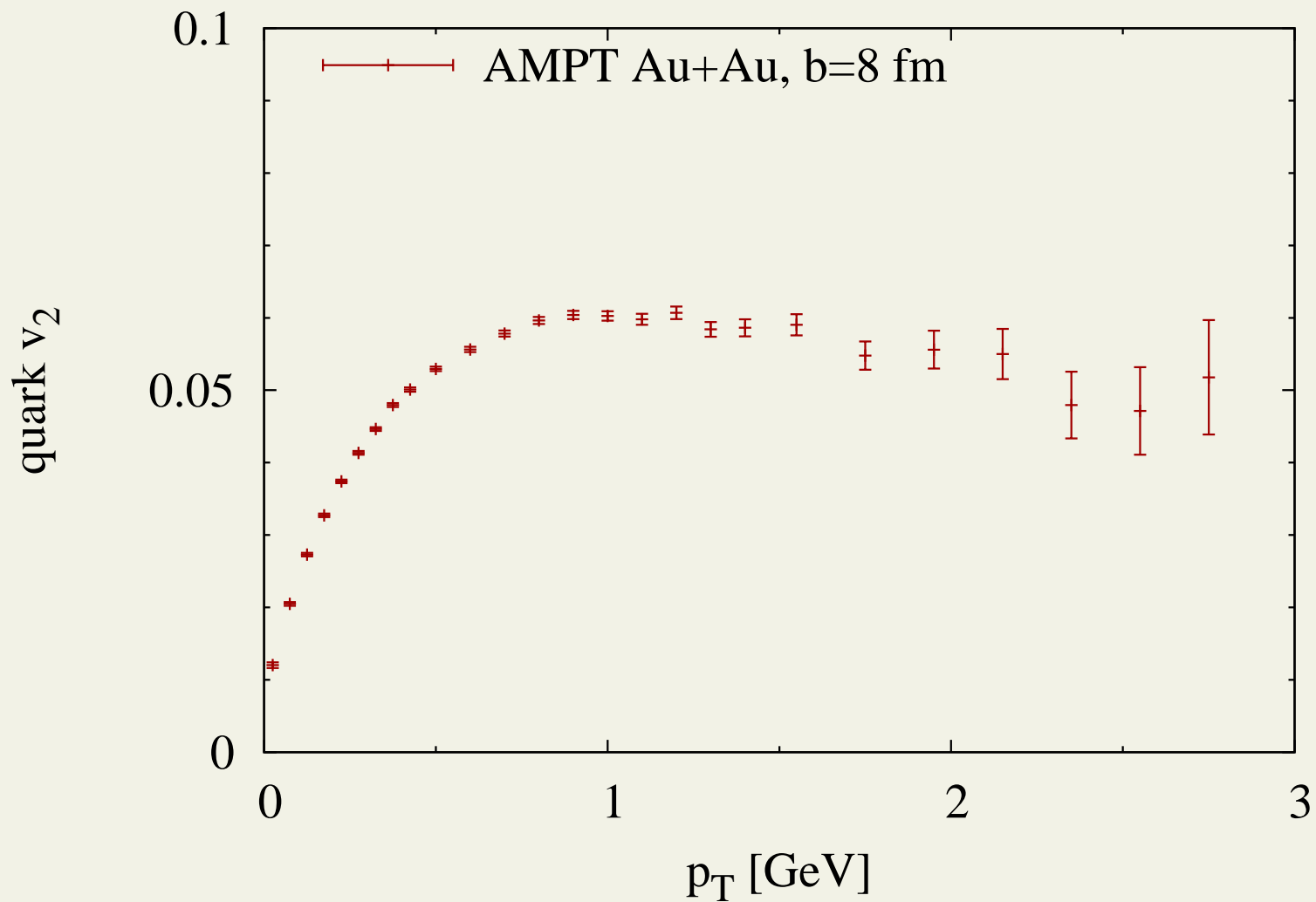
- minijet initconds

- 1 parton \rightarrow 1 π hadronization

perturbative $\sigma_{gg \rightarrow gg} \approx 3 \text{ mb}$ gives $v_2 \approx 2\%$ \rightarrow **need 15 \times higher opacity**

radiative $gg \leftrightarrow ggg$ helps (e.g., BAMPS)... but AMPT has pure elastic $2 \rightarrow 2$

quark v_2 from AMPT for 200-GeV Au+Au, $b=8$ fm



$v_2 \approx 6\%$ with only $\sigma_{qq} = 3$ mb... yet $\sim 15\%$ hadron v_2 (30% centrality)

v_2 amplified by coalescence

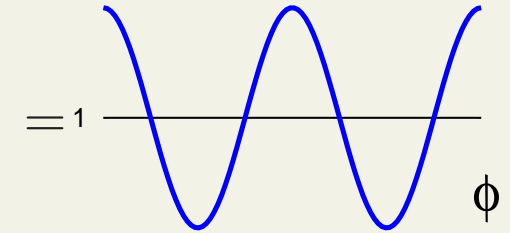
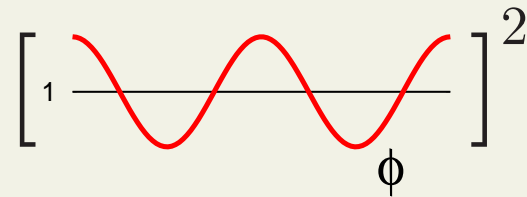
Ko, Lin, Voloshin, DM, Greco, Levai, Mueller, Fries, Bass, Nonaka, Asakawa ...

coalescence of comoving quarks: $q\bar{q} \Rightarrow \longrightarrow M$ $3q \Rightarrow \longrightarrow B$

DM & Voloshin, PRL91 ('03)

$$\frac{dN_M(p_T)}{d\phi} \propto \left[\frac{dN_q(p_T/2)}{d\phi} \right]^2$$

$$\frac{dN_B(p_T)}{d\phi} \propto \left[\frac{dN_q(p_T/3)}{d\phi} \right]^3$$

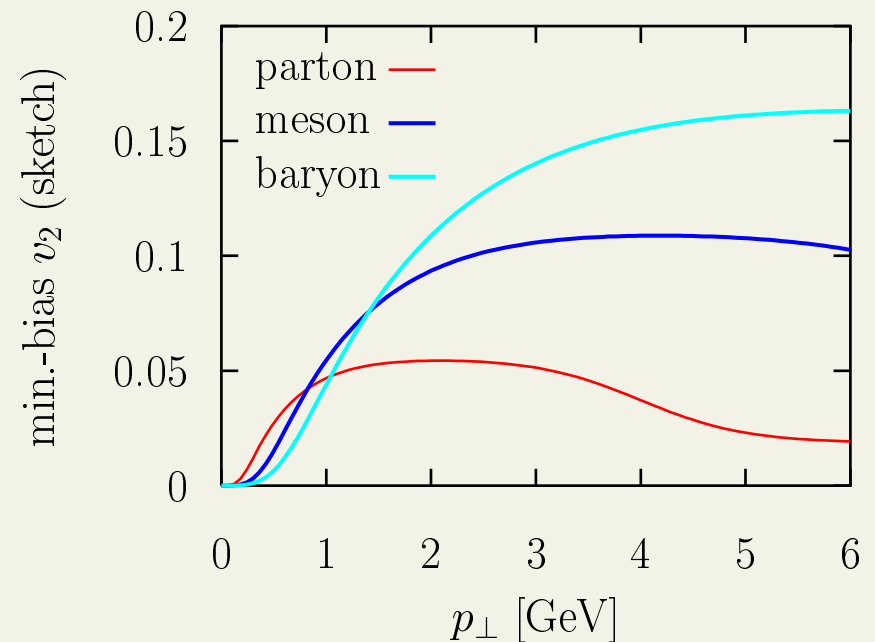


squared/cubed probability \rightarrow amplified v_2

$$v_2^{\text{hadron}}(p_\perp) \approx n \times v_2^{\text{quark}}(p_\perp/n)$$

$3 \times$ for baryons } 50% larger v_2
 $2 \times$ for mesons } for baryons

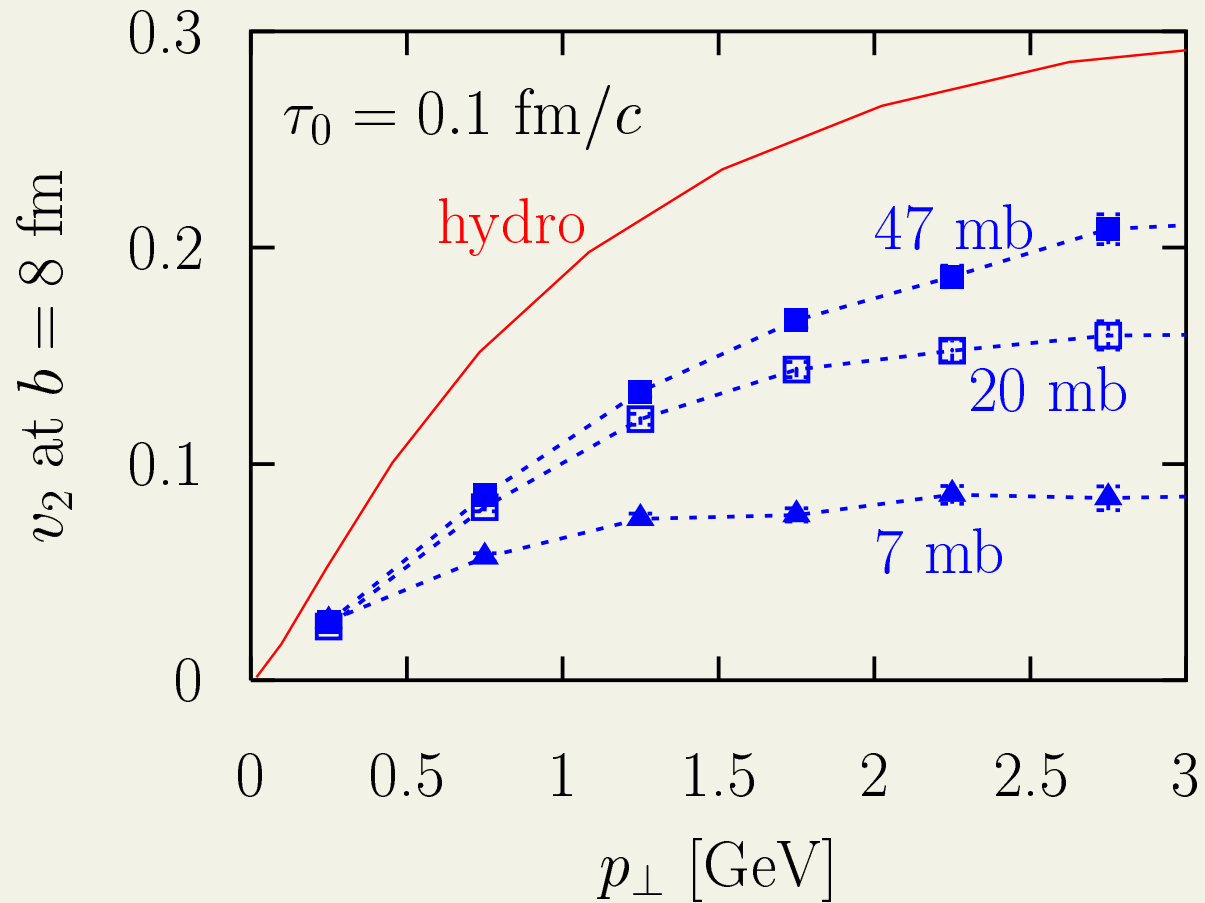
$\rightarrow 5 \times$ for pentaquark, $6 \times$ for deuteron



\Rightarrow AMPT can get hadron v_2 with lower opacity... in principle

could we be in or near the hydro limit? → **No.**

DM & Huovinen, PRL94 ('05)

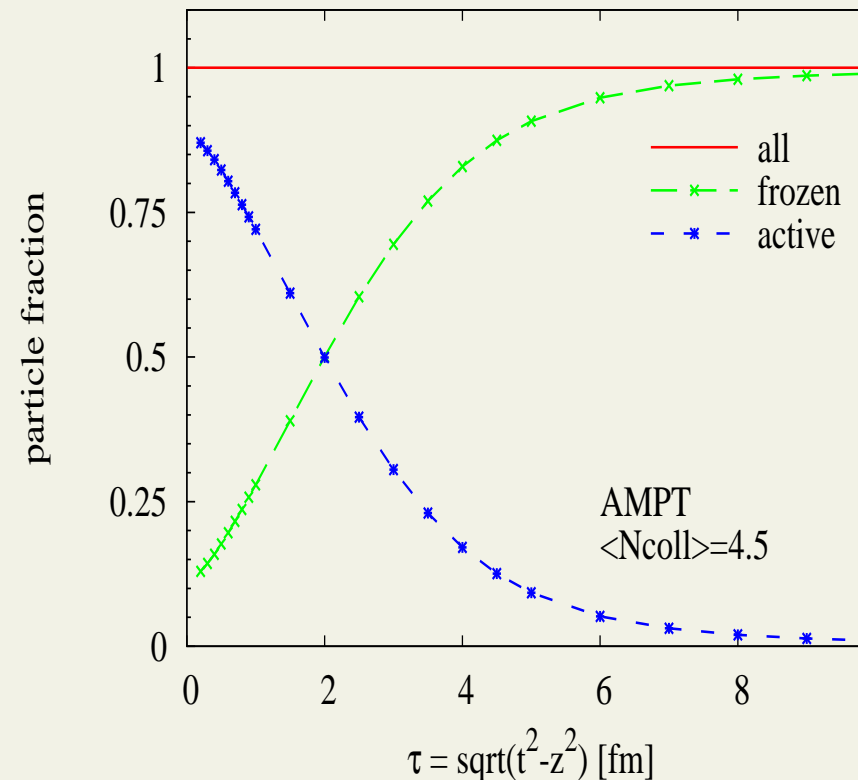
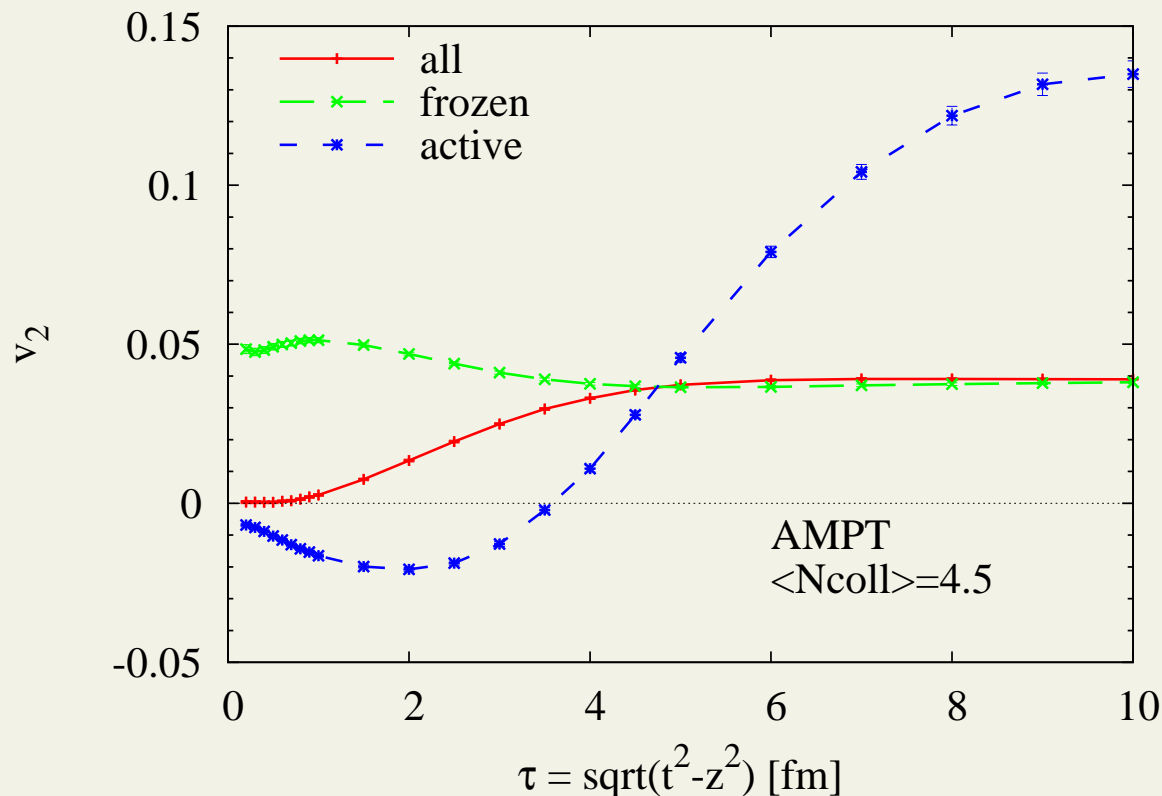


- $2 \rightarrow 2$ parton kinetic theory is definitely not ideal hydro
- for large $\sigma \sim 40 - 50 \text{ mb}$, viscous hydro is reachable Huovinen & DM, JPG35 ('08)

AMPT \neq Hydro

AMPT's v_2 comes from "anisotropic escape" He et al, PLB 753 ('15); Li et al, ...

DM et al ('15): AMPT v_2 vs Bjorken τ



- for long time, **still interacting** partons carry nearly zero or even negative v_2
- almost all v_2 is carried by **frozen-out** partons
- the **combined** parton v_2 rises then saturates as expected Ko, Zhang, Gyulassy ('99)

AMPT still gets $v_2 \sim 6\%$ for quarks with only 3 mb... but how?

Test: use parton transport MPC to check AMPT's partonic stage

1) dynamics - issue of parton subdivision

2) initconds - AMPT vs just minijets

focus on **200-GeV Au+Au** at RHIC, fixed $b = 8$ fm impact parameter

Covariant transport

(on-shell) phase-space density $f(x, \vec{p}) \equiv \frac{dN(\vec{x}, \vec{p}, t)}{d^3x d^3p}$

transport equation (BTE):

$$p^\mu \partial_\mu f_i(x, p) = C_{2 \rightarrow 2}^i[\{f_j\}](x, p) + C_{2 \leftrightarrow 3}^i[\{f_j\}](x, p) + \dots$$

with, e.g.,

$$C_{2 \rightarrow 2}^i = \frac{1}{2} \sum_{jkl} \int_{234} (f_3^k f_4^l - f_1^i f_2^j) W_{12 \rightarrow 34}^{ij \rightarrow kl} \left(\int_j \equiv \int \frac{d^3 p_j}{2E_j}, \quad f_a^k \equiv f^k(x, p_a) \right)$$

thermalizes (in box), fully causal and stable \rightarrow can derive hydro eqns
e.g., Denicol, Rischke et al

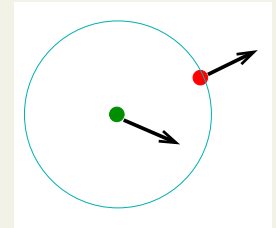
handles both high or low opacities \rightarrow usable for fluid-to-particle conversion
e.g., Teaney, Moore & Dusling; DM & Wolff, ...

hydro limit: transport coeffs & rel. times ($\eta \approx 1.2T/\sigma$, $\tau_\pi \approx 1.2\lambda_{tr}$...)

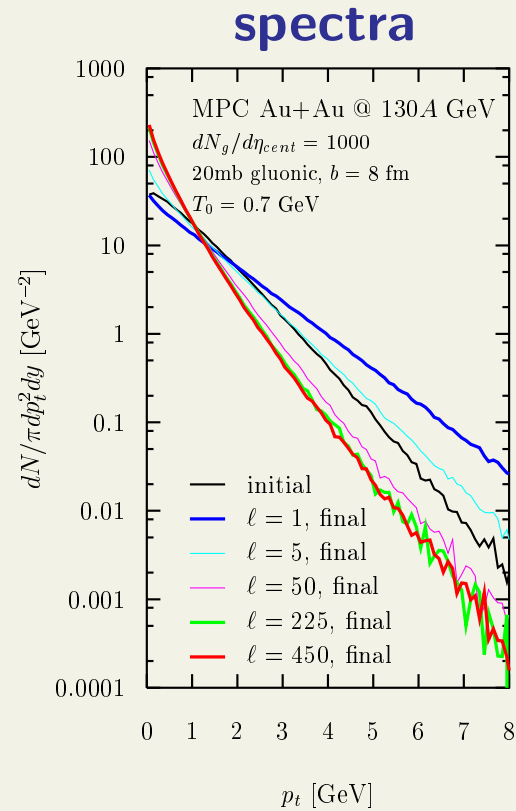
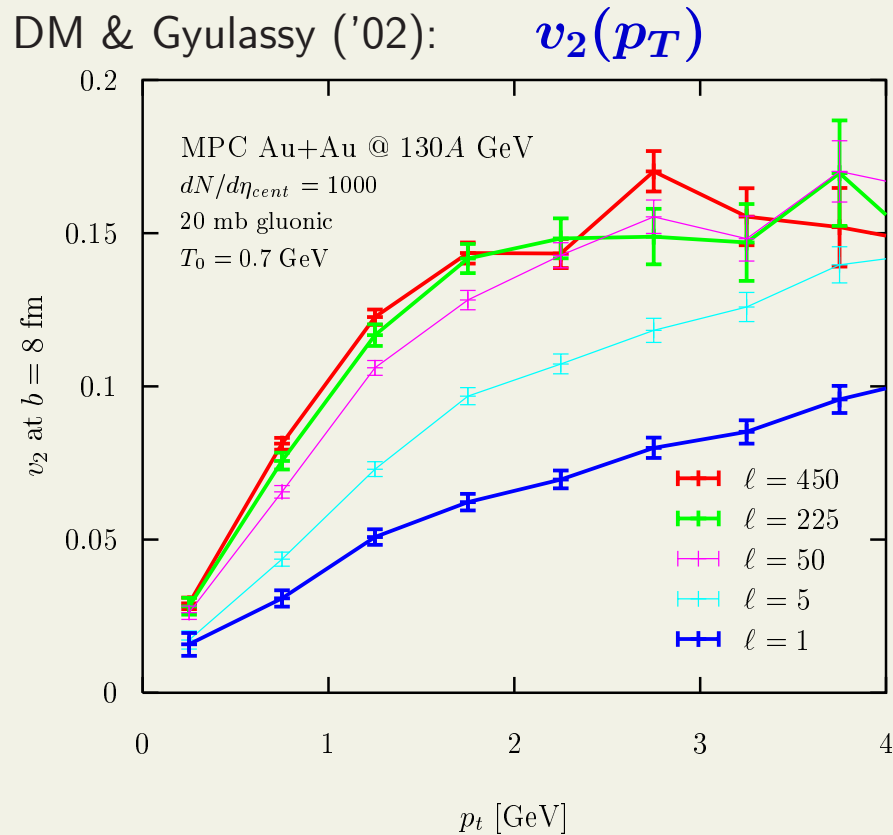
\exists covariant transport codes: **ZPC** (Zhang), **MPC** (Molnar), **BAMPS** (Xu), ...

Parton subdivision

Nonlocal artifacts: due to action at distance $d < \sqrt{\frac{\sigma}{\pi}}$



subdivision: rescale $f \rightarrow f \cdot \ell$, $\sigma \rightarrow \sigma/\ell \Rightarrow d \propto \ell^{-1/2}$ local as $\ell \rightarrow \infty$



- ZPC could do subdivision, but **AMPT** runs it with $\ell = 1$
- high RHIC opacities: need subdivision $\ell \sim \mathcal{O}(100)$ to remove artifacts in v_2

Initial conditions (Au+Au at RHIC)

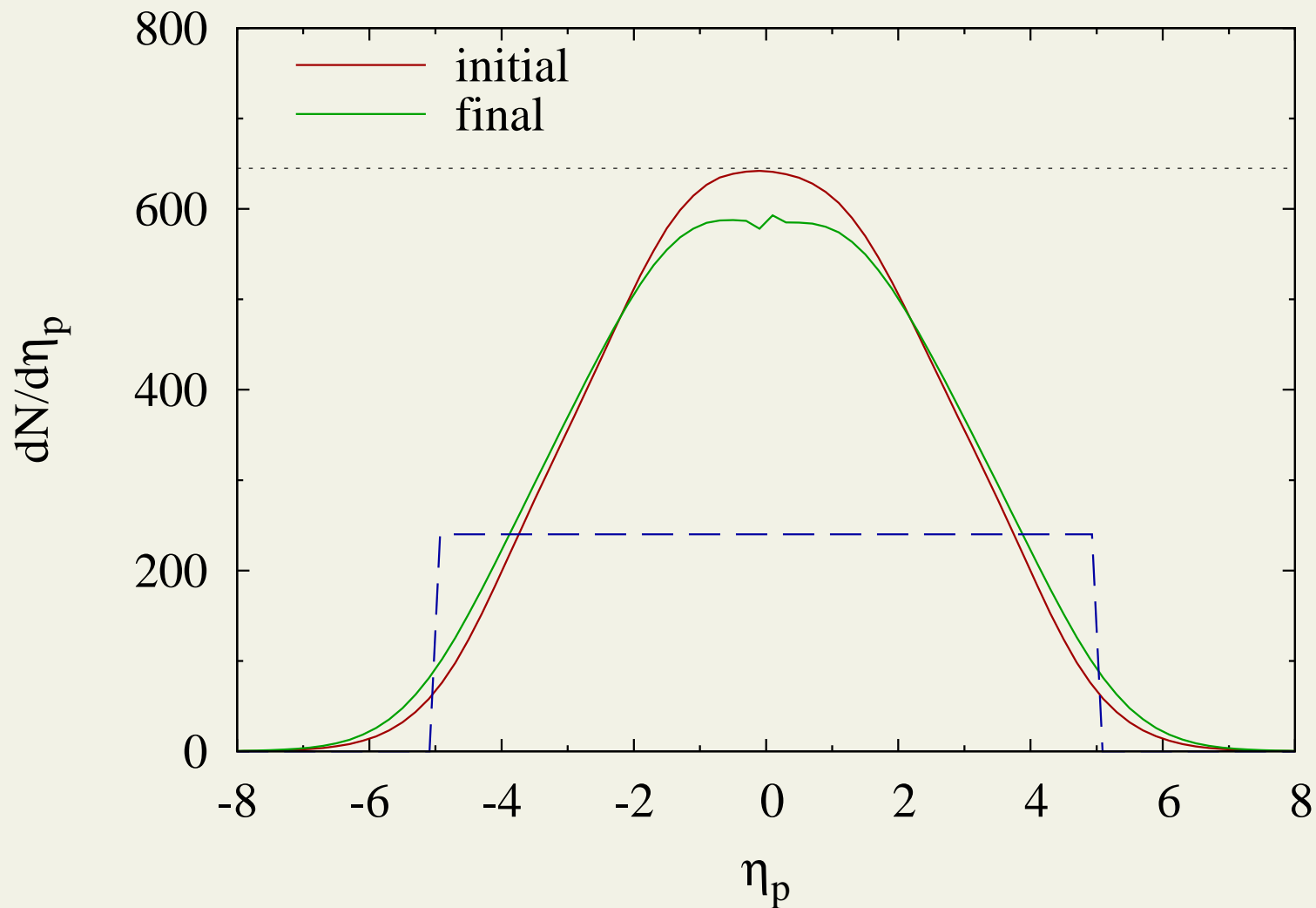
Molnar-Gyulassy study: boost-invariant fit to HIJING minijets

- massless gluons
- flat $dN_g/d\eta(b=0) = 1000 \Rightarrow dN/d\eta(b=8 \text{ fm}) \approx 240$
- locally thermal $T_0 = 0.7 \text{ GeV}$, $f = N(\vec{x}_T)e^{-m_T \cosh(\eta-y)/T_0}$
- constant formation time $\tau_0 = 0.1 \text{ fm} \sim 1/p_T$
- binary collision transverse profile for Au+Au at $b = 8 \text{ fm}$
- $d\sigma/dt \propto 1/(t - \mu_D^2)^2$, $\sigma_{gg} = 9\pi\alpha_s^2/2\mu_D^2$, with $\mu_D = 0.7 \text{ GeV}$

AMPT: v2.26t5d6

- almost massless u, d & massive s quarks
- enhanced $dN_q/d\eta \approx 2.5 \times dN_h/d\eta$ (string melting), nonuniform in η
- formation time distribution
- transverse profile close to wounded nucleons
- sizeable event-by-event fluctuations
- $\mu_D = 0.45 \text{ GeV}$, $\alpha_s = 0.33 \Rightarrow \sigma = 3 \text{ mb}$ (quark Casimirs ignored)

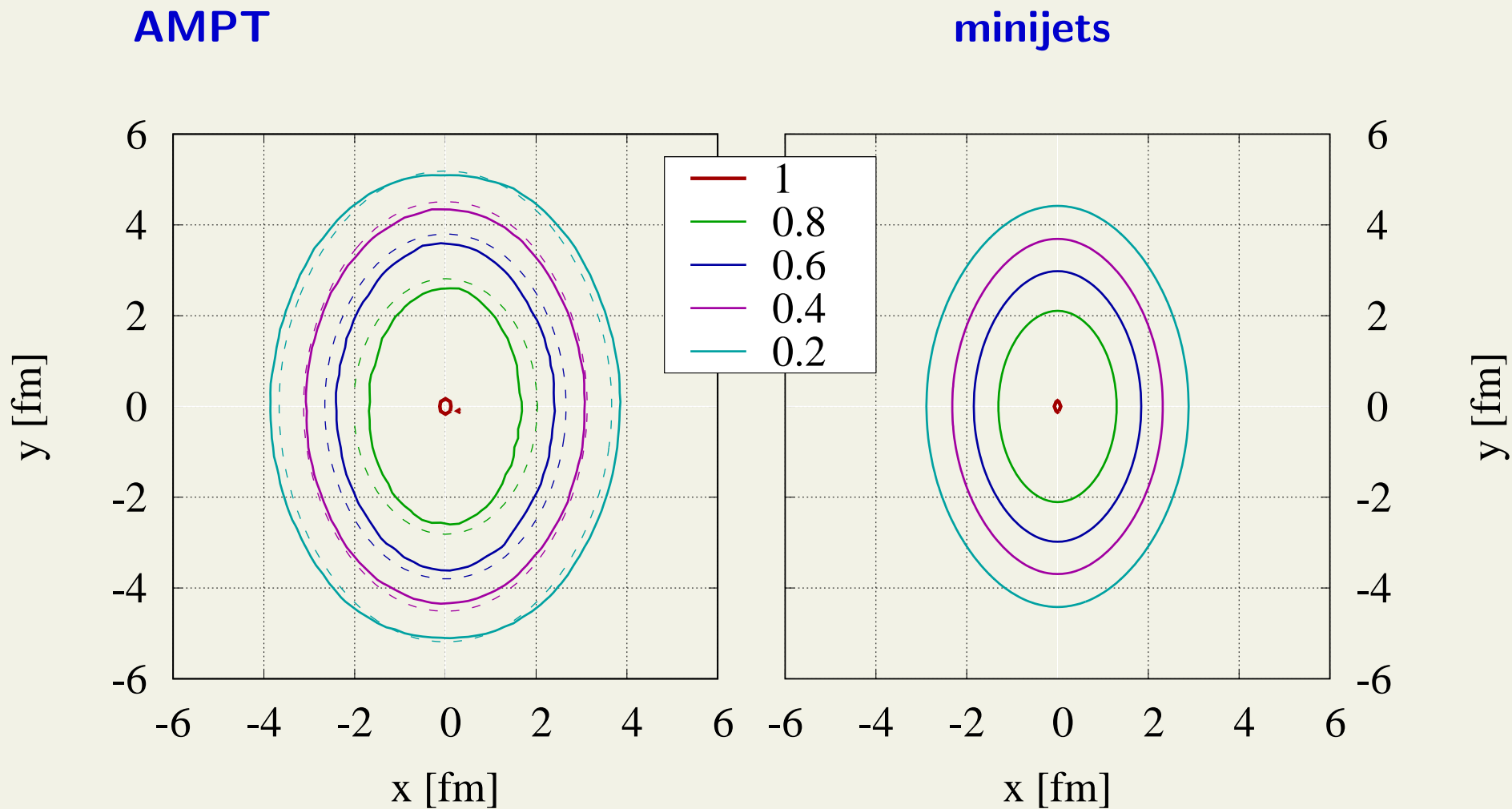
Au+Au, b=8 fm - AMPT pseudorapidity $dN/d\eta_p$



$$\eta_p \equiv \frac{1}{2} \ln \frac{|p| + p_z}{|p| - p_z}$$

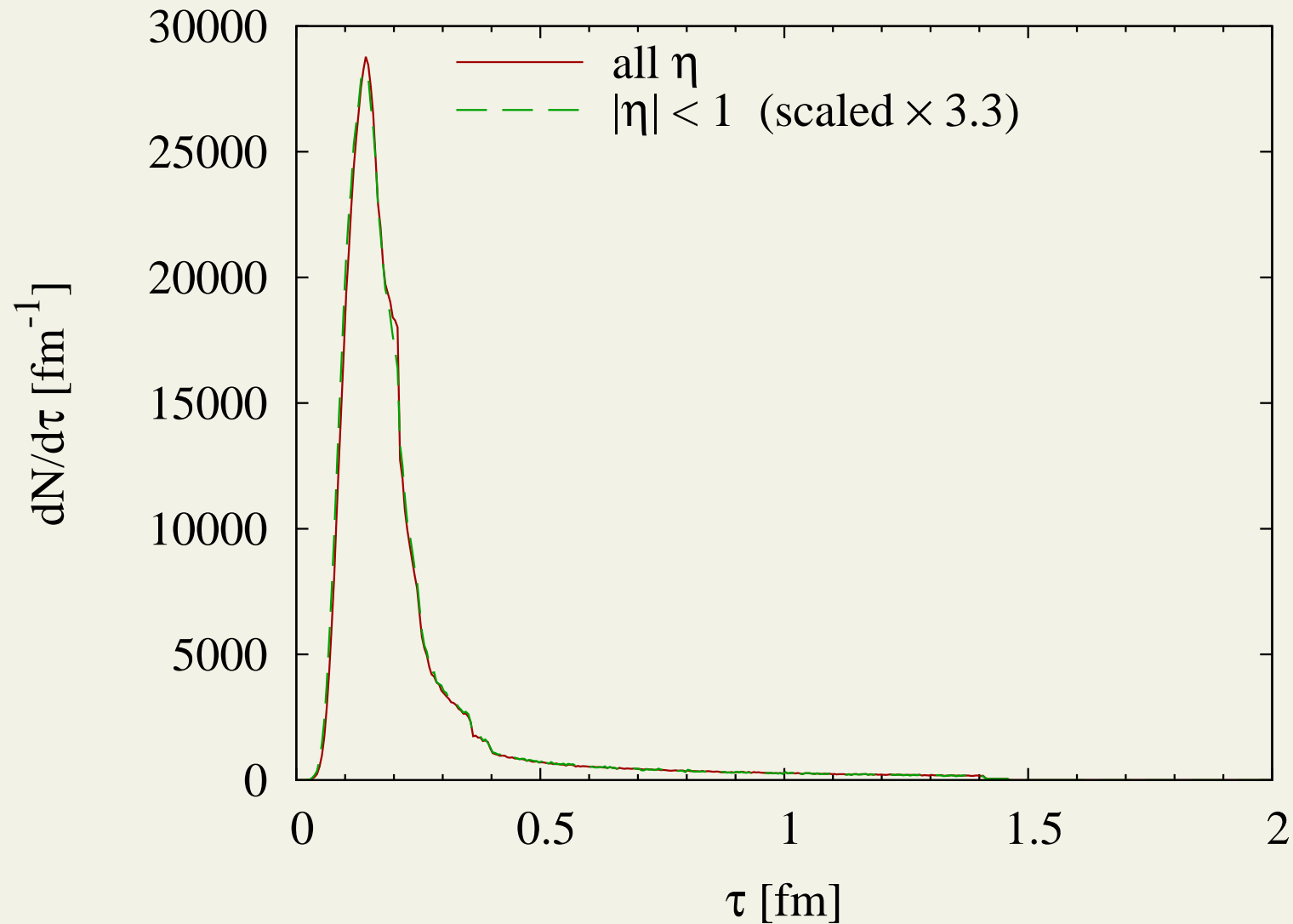
evidently not boost invariant, peak around $dN/d\eta_p \sim 650$

Au+Au, b=8 fm - transverse profile dN/d^2x_T for $|\eta_p| < 1$

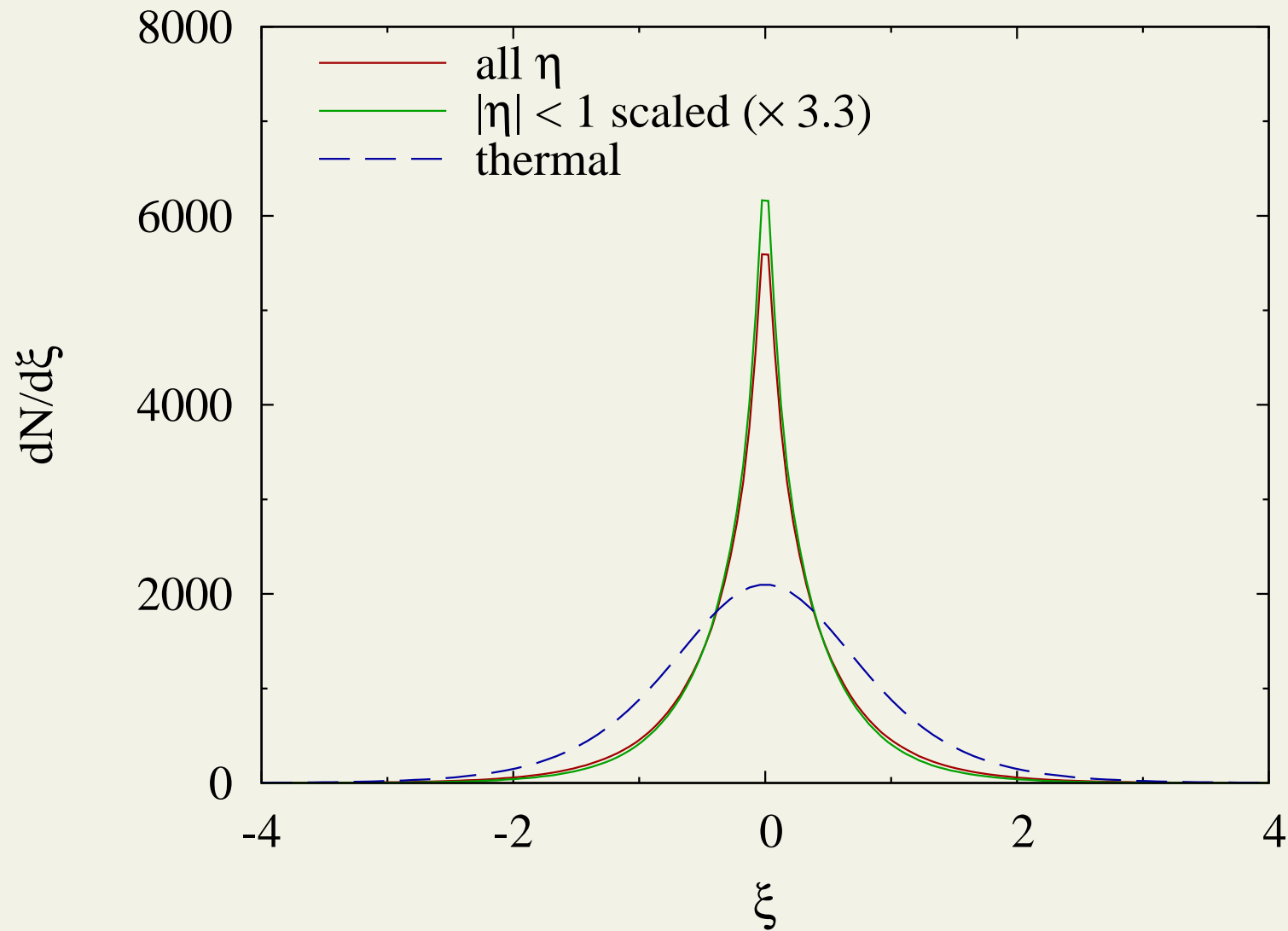


AMPT: close to wounded nucleons (dashed); minijets: binary coll profile

Au+Au, b=8 fm - AMPT formation time $\tau = \sqrt{t^2 - z^2}$

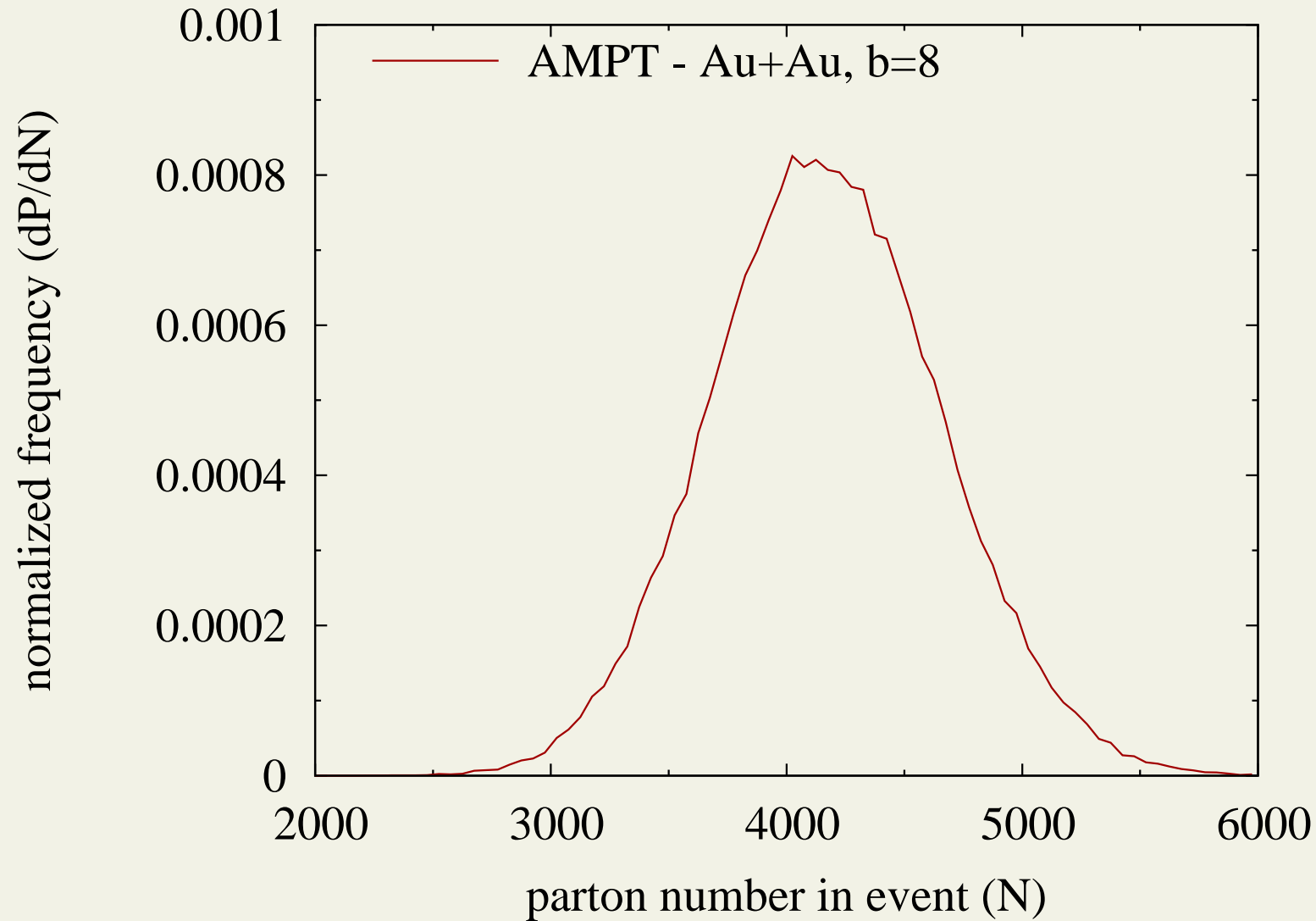


broad distribution, peaks around $\tau \sim 0.2$ fm (average $\tau \approx 0.23$ fm)



much sharper than boost-invariant thermal $\propto 1/\cosh^2 \xi$ correlation

Au+Au, b=8fm - AMPT total quark number in event

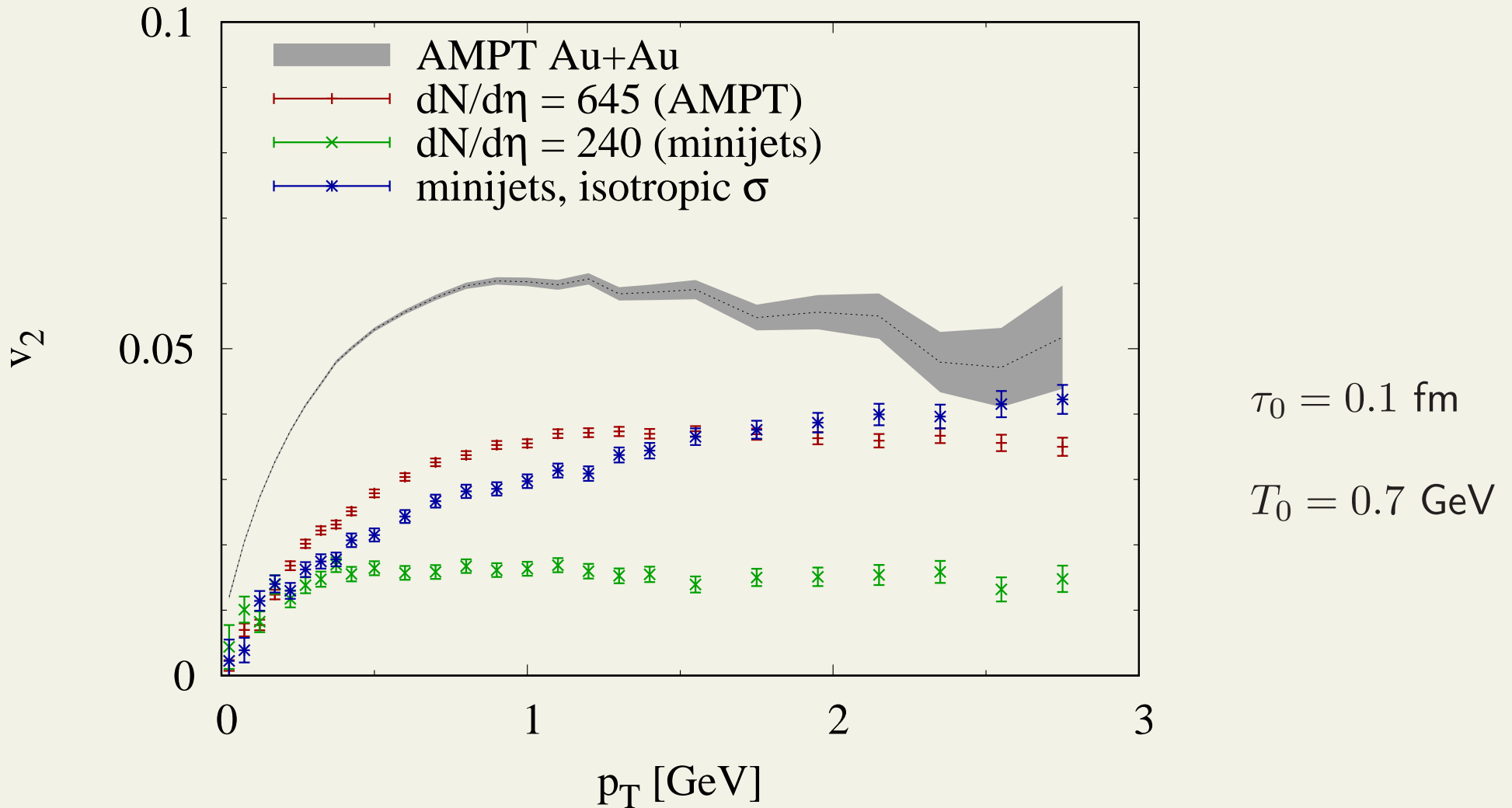


large event-by-event fluctuations, $\langle N_{quark} \rangle \approx 4200$

MPC v1.9 comparison strategy: go from simple to complicated

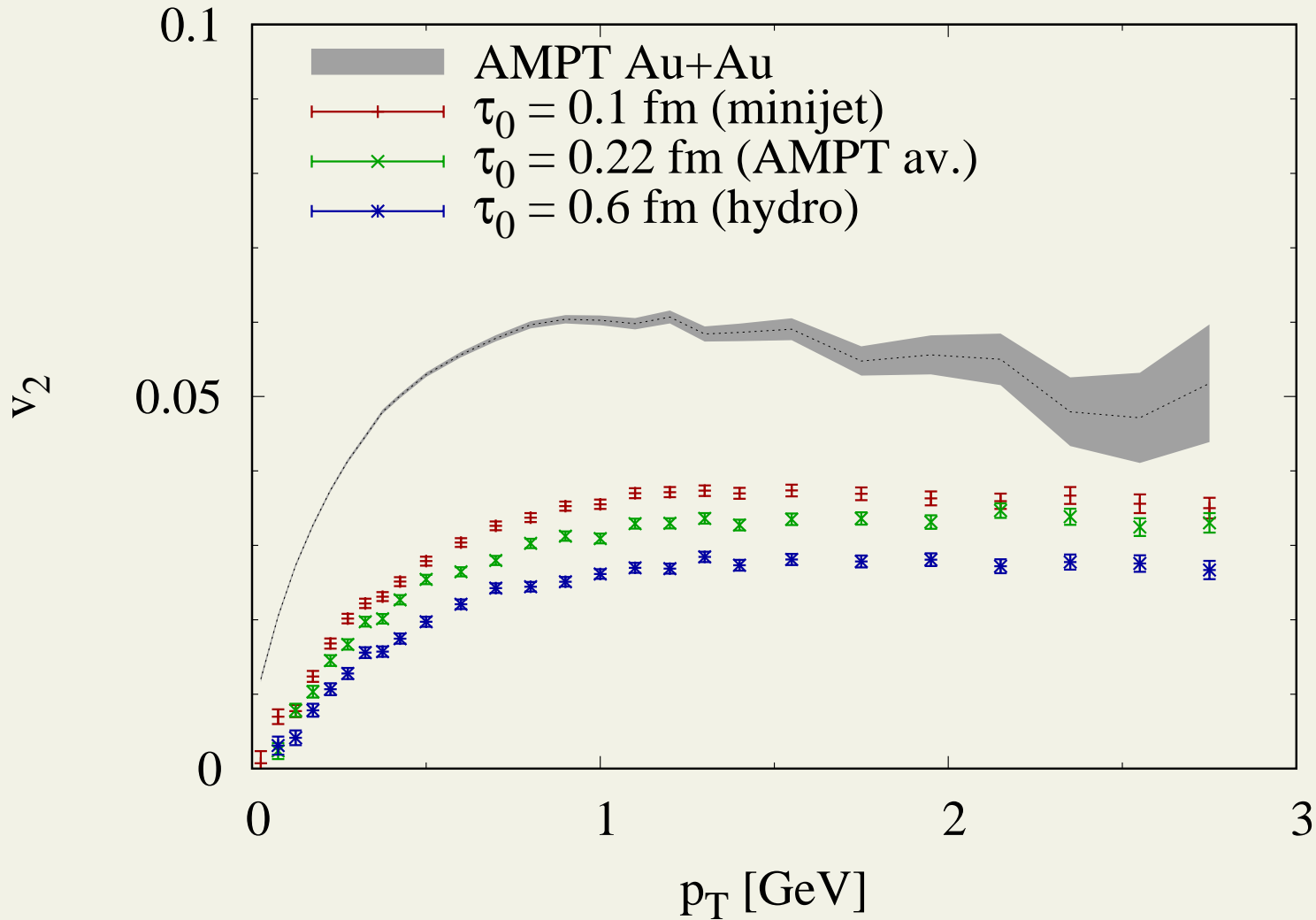
- take boost-invariant minijet study as baseline
- gradually include features of AMPT initconds
- use wounded nucleon profile, and the same $d\sigma/dt$ as AMPT

1) set $dN/\eta_p \approx 650$ to match AMPT (at midrapidity), use AMPT $d\sigma/dt$



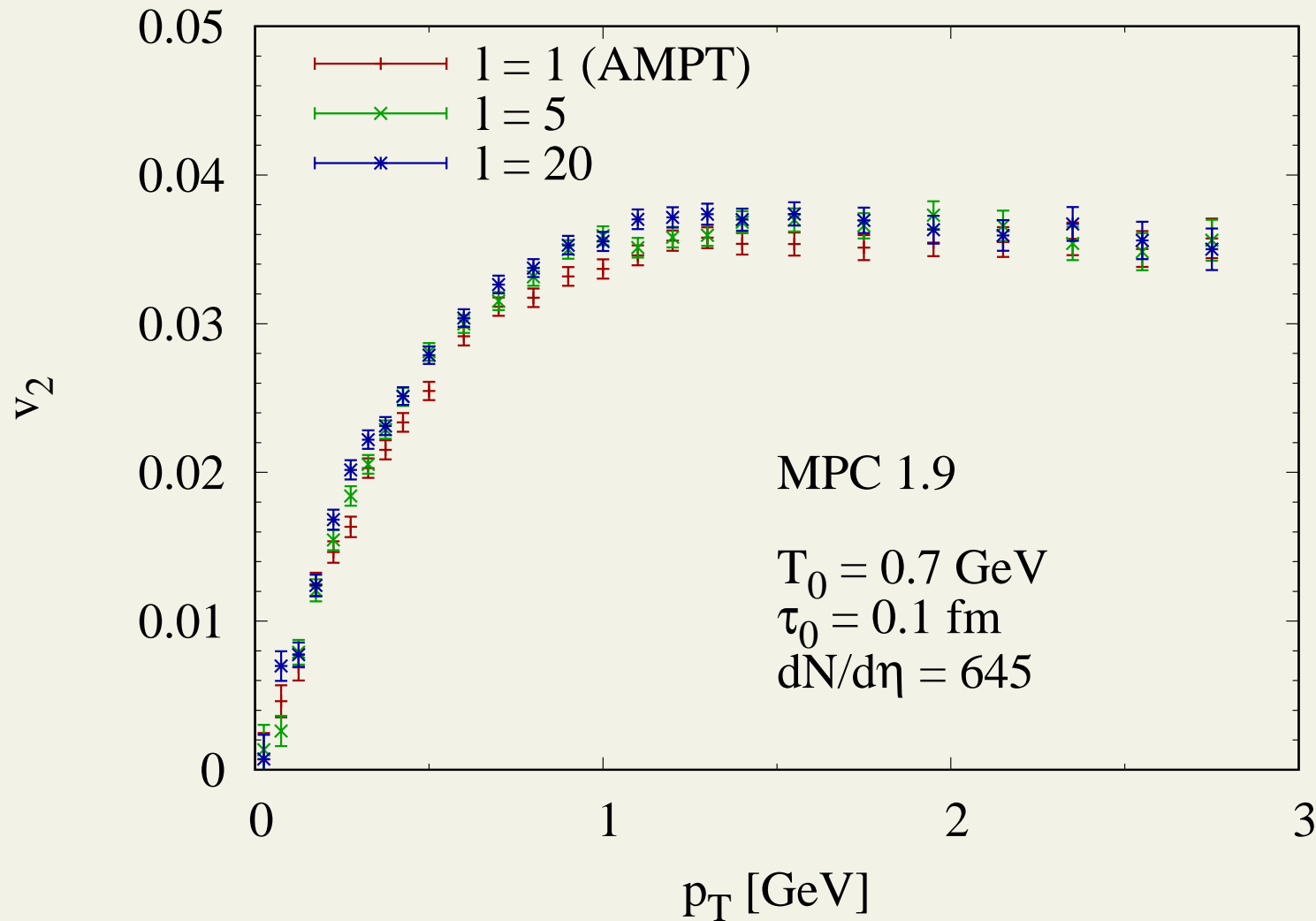
$2.7 \times$ larger density, but only gets v_2 to $\approx 3.5\%$... (need $5 \times$ more opacity)

2) match formation time to AMPT average, $\langle \tau_0 \rangle = 0.22$ fm



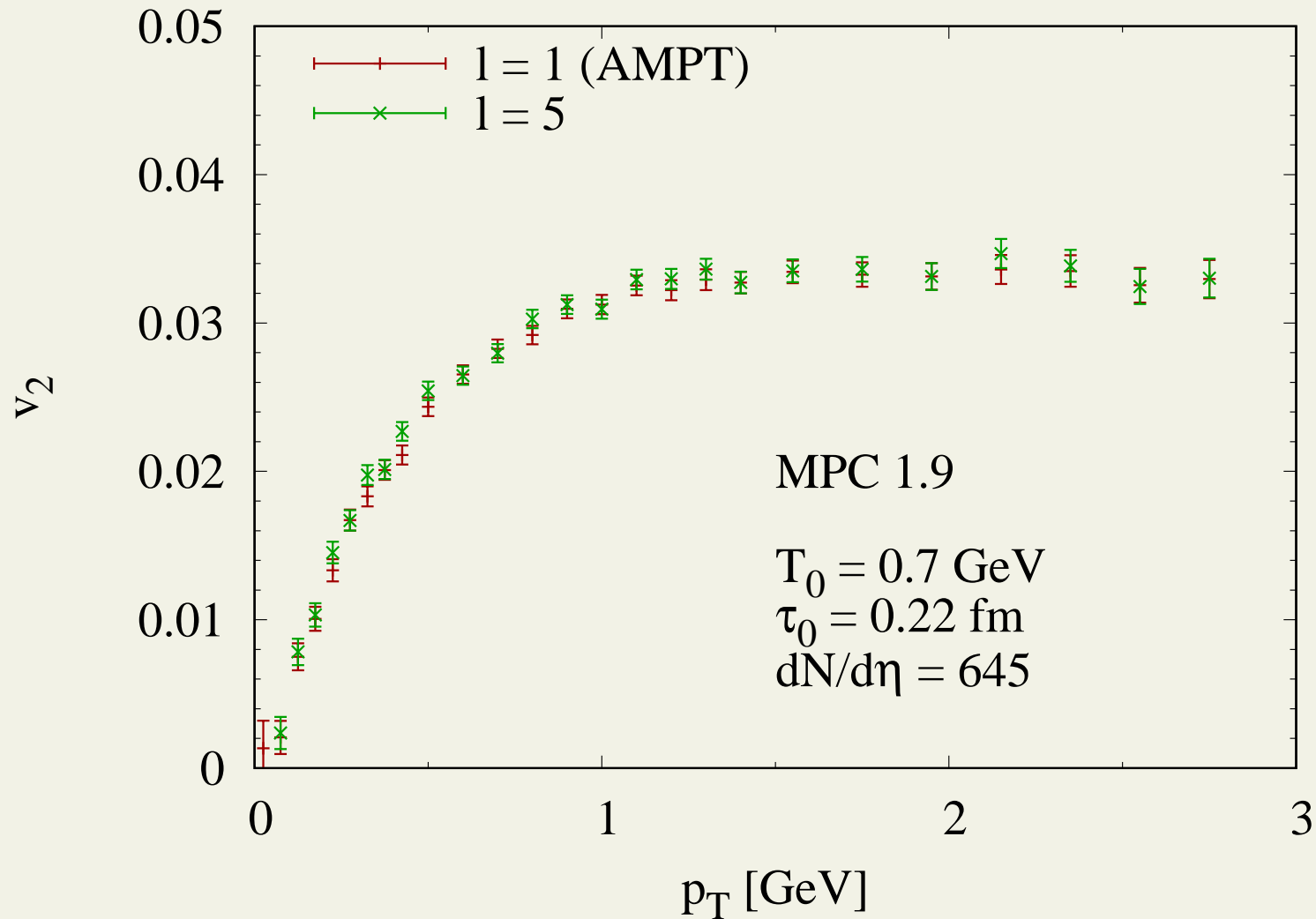
larger $\tau_0 \Rightarrow$ smaller v_2 (lower density $n \sim 1/\tau$, also shorter evolution)

3) decrease parton subdivision, down to $\ell = 1$



modest $\lesssim 10\%$ effect on $v_2(p_T)$, for short $\tau_0 = 0.1$ fm

same, but now match $\langle \tau_0 \rangle$ from AMPT



less than 5% distortion in $v_2(p_T)$, removable with subdivision $\ell = 5$

So far we...

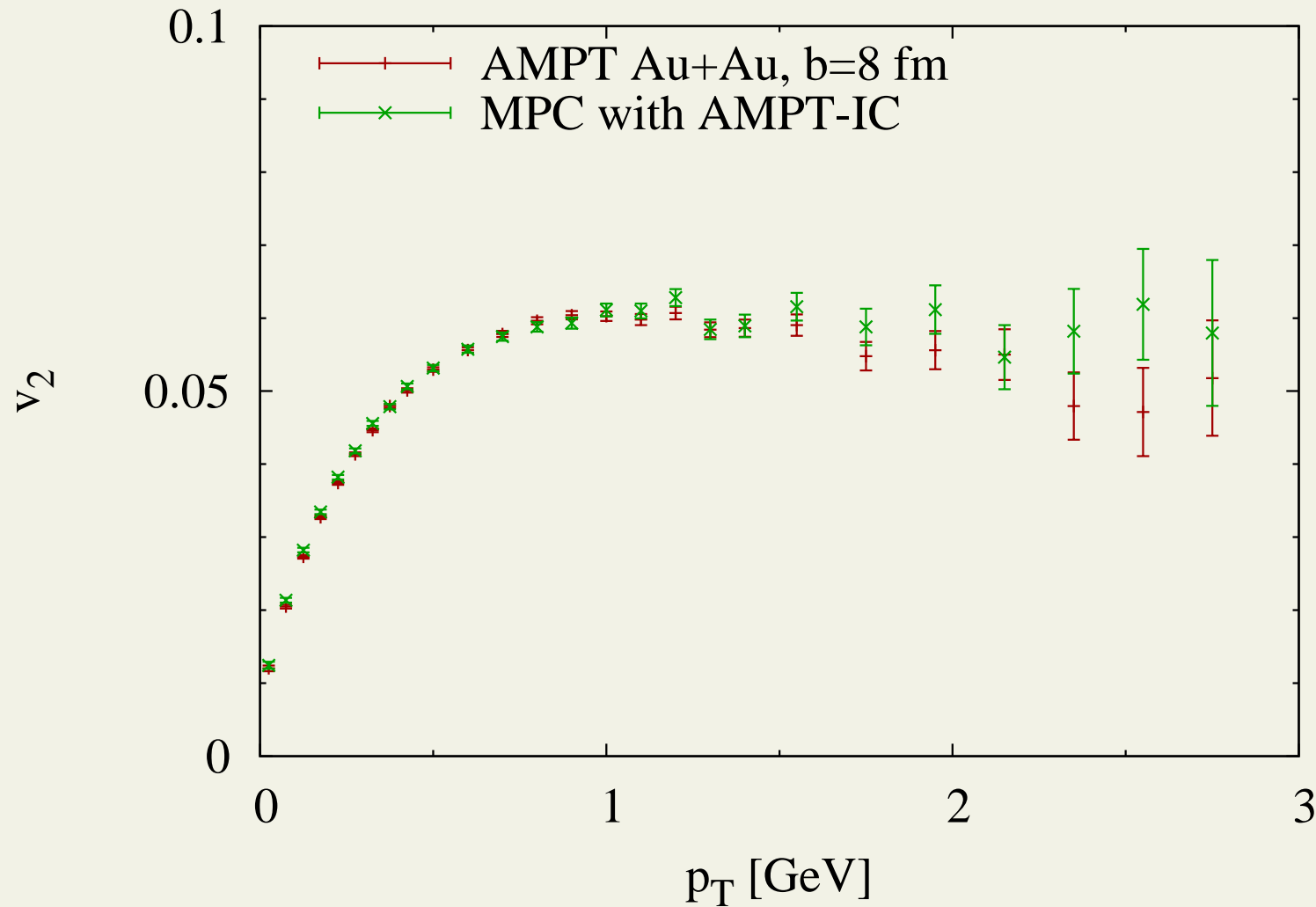
- matched initial $dN/d\eta$, formation time, transverse profile, interactions
- ruled out too low parton subdivision

(effectively, AMPT runs at subdivision ~ 10 relative to $\sigma \sim 30$ mb,
and another factor of ~ 5 would be sufficient)

still, v_2 is too small.

what if we use the exact same initial conditions?

4) MPC with full AMPT initial conditions (same \vec{p} , \vec{x} , t for each parton)



MPC with $\sigma = 3$ mb ($\ell = 1$) reproduces the elliptic flow from AMPT!

Missing link: momentum distribution

- for thermal minijets, $\langle p_T \rangle = 3\pi T/4 \approx 1.6$ **GeV** for $T_0 = 0.7$ **GeV**

$$\frac{dN}{d^2p_T dy} = E \frac{dN}{d^3p} = p_\mu d\sigma^\mu f \propto m_T \tau \cosh \xi e^{-m_T \cosh \xi / T}$$

- for AMPT quarks, $\langle p_T \rangle \approx 0.54$ **GeV** $\Rightarrow T_{eff} = 0.23$ **GeV** only

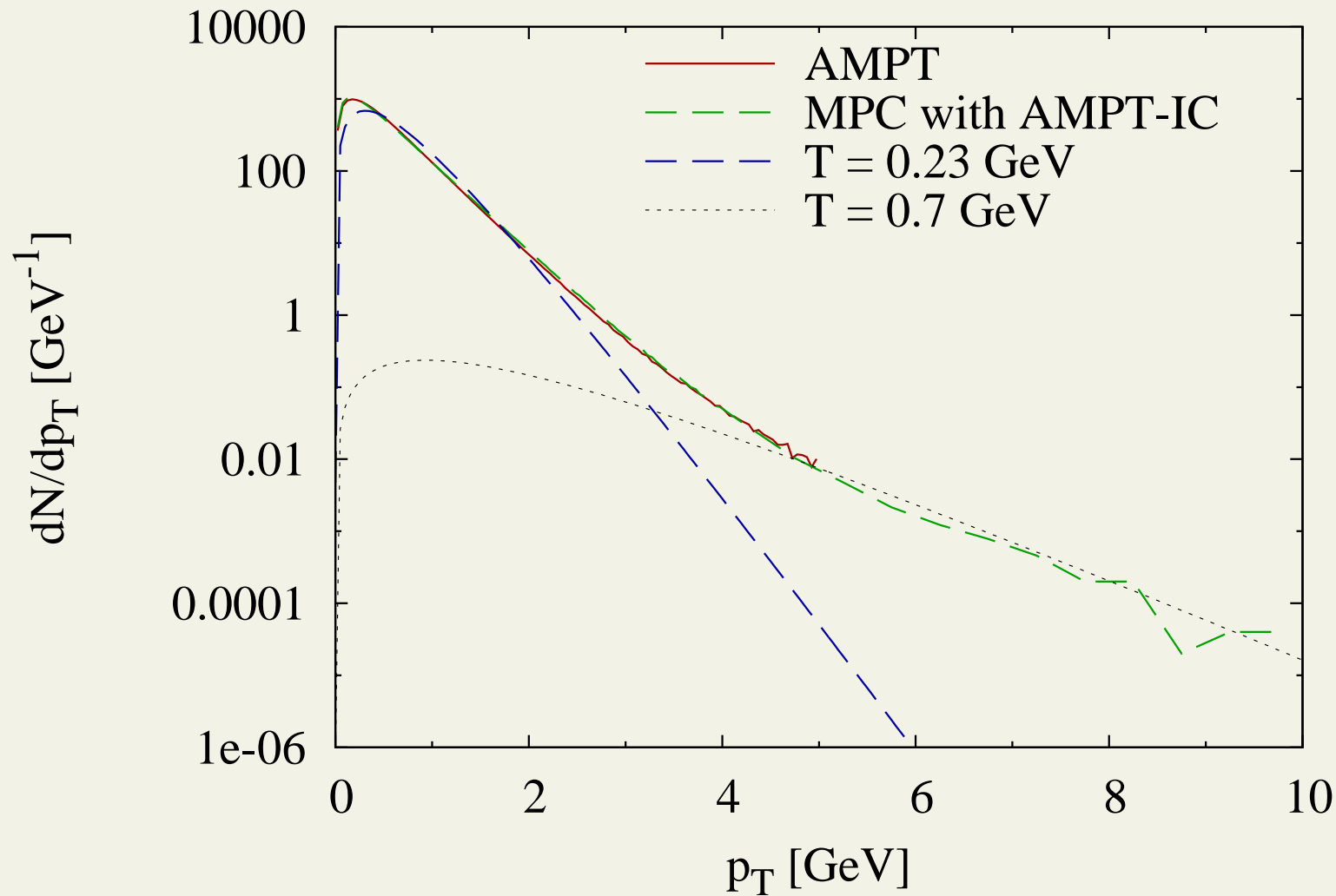
at lower temperature, the cross section is more isotropic (with μ_D fixed)

\Rightarrow more effective v_2 generation

$$\sigma_{tr} = 4\sigma_{TOT} z(1+z) \left[(2z+1) \ln \left(1 + \frac{1}{z} \right) - 2 \right], \quad z \equiv \frac{\mu^2}{s} \approx \frac{\mu^2}{18T^2}$$

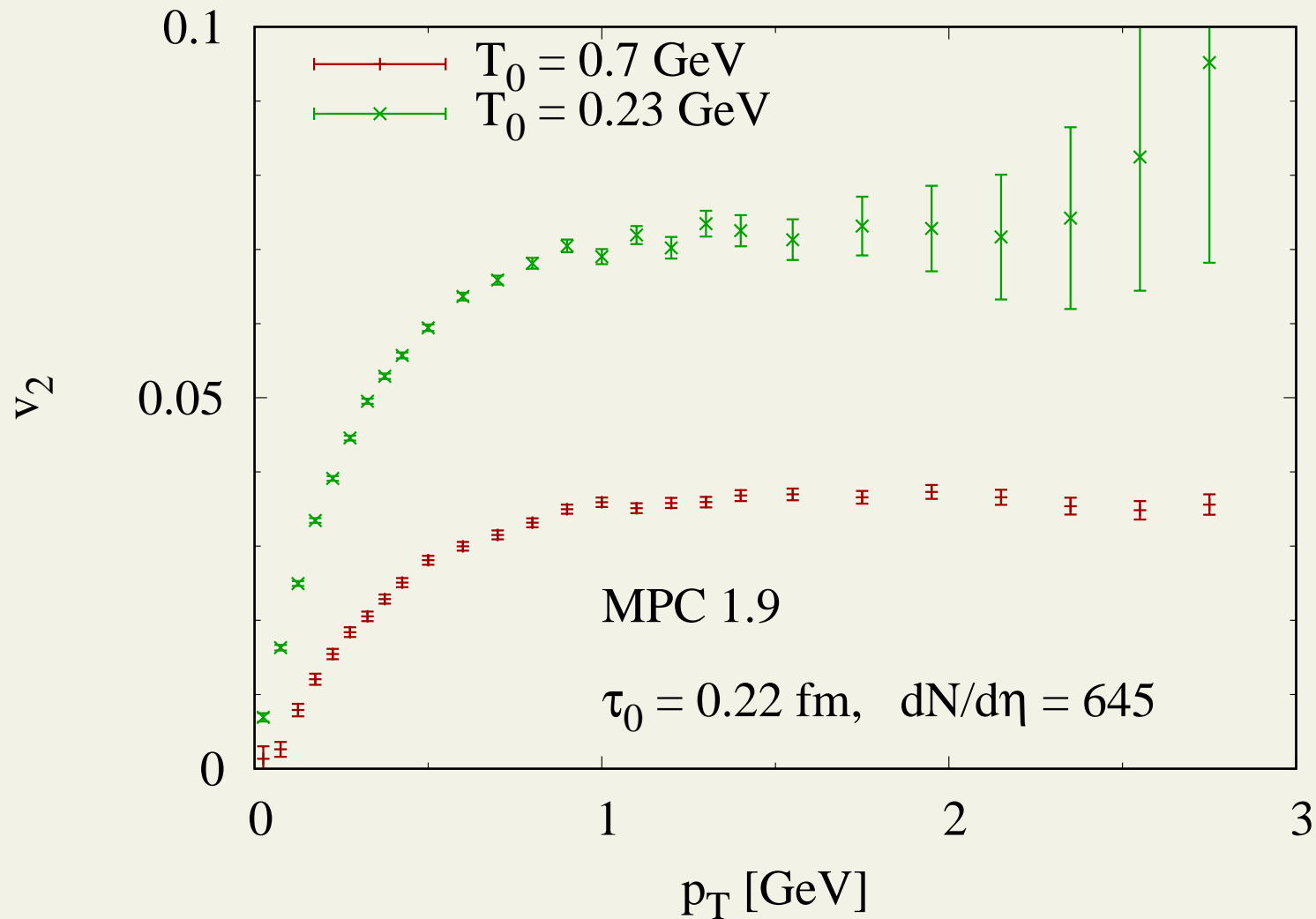
$$\text{so } \frac{\sigma_{tr}(T = 0.23 \text{ GeV})}{\sigma_{tr}(T = 0.7 \text{ GeV})} \approx 2.6 \text{ (!)}$$

initial parton spectrum in Au+Au at RHIC, $b=8$ fm ($|\eta_p| < 1$)



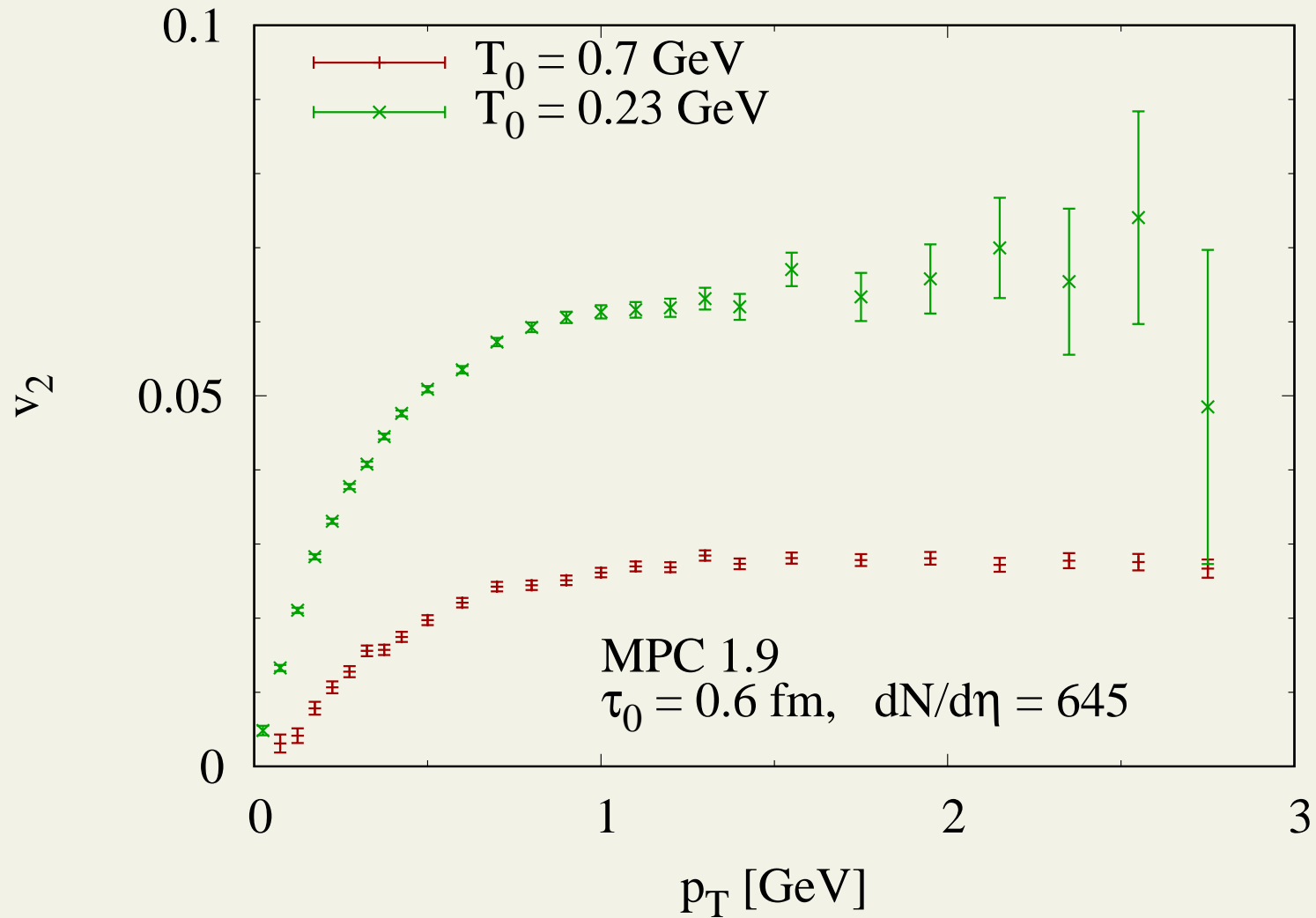
$T = 0.7$ GeV captures minijet tail, but string melting plasma is $3\times$ colder

4) now set T_{eff} to AMPT $\langle p_T \rangle$



large 6 – 7% v_2 , even slightly higher than from AMPT

same comparison but for higher $\tau_0 = 0.6$ fm



even $2.5\times$ later formation time works, v_2 practically doubles

Summary

Using the covariant transport solver MPC (Molnar's Parton Cascade), we investigated how AMPT generates enough elliptic flow (v_2) in Au+Au at RHIC with only 3-mb cross sections. Two ingredients play the most important role: i) **about $2.5\times$ higher parton densities** via string melting, and ii) **cold initial plasma** with effective temperature $T_{eff} \sim 0.2 - 0.25$ GeV, comparable to μ_D (makes forward-peaked partonic cross sections more isotropic).

The partonic stage of AMPT approximates well solutions of the covariant Boltzmann transport equation. **Artifacts due to lack of parton subdivision were below 5%** for $v_2(p_T)$ in Au+Au at RHIC. Compared to studies that generated high opacities with large $\sigma \sim 30$ mb, AMPT effectively incorporates a subdivision $\ell \sim 10$ already.

With 3 mb cross section, AMPT still generates $2 - 3\times$ smaller partonic v_2 than that of hadrons. But it makes up for the difference with quark coalescence - the algorithm of which needs to be tested in detail.

Next steps:

- include fluctuations in τ_0 , N_{quark} , initial profile
- check v_2 distribution, not just its event average