

# 全国中高能核物理大会

Transverse Momentum Balance and Angular Correlation  
of  $b\bar{b}$  Dijets in Pb+Pb collisions

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In Collaboration with S. Wang, S.L Zhang, B.-W Zhang, Enke Wang

[arXiv:1806.06332]

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2019/06/23

第十八届中高能核物理, Wei Dai

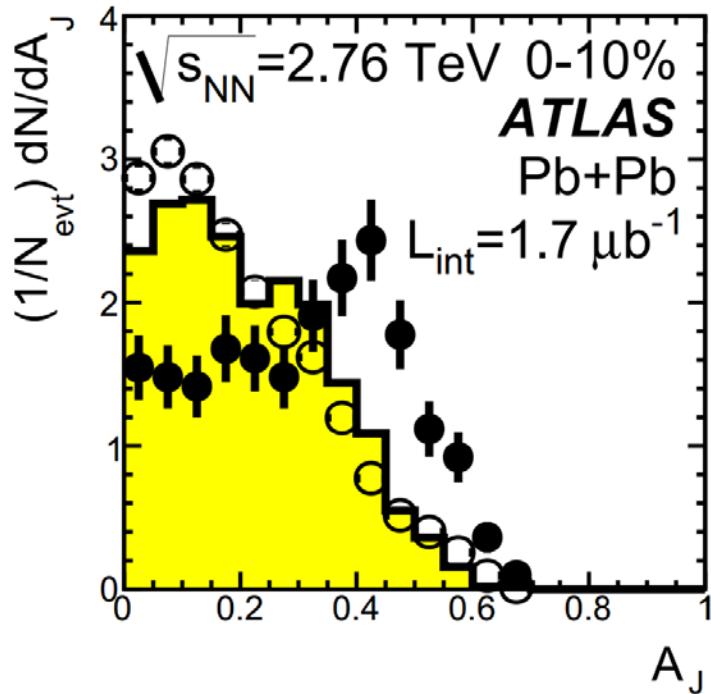


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# Outline

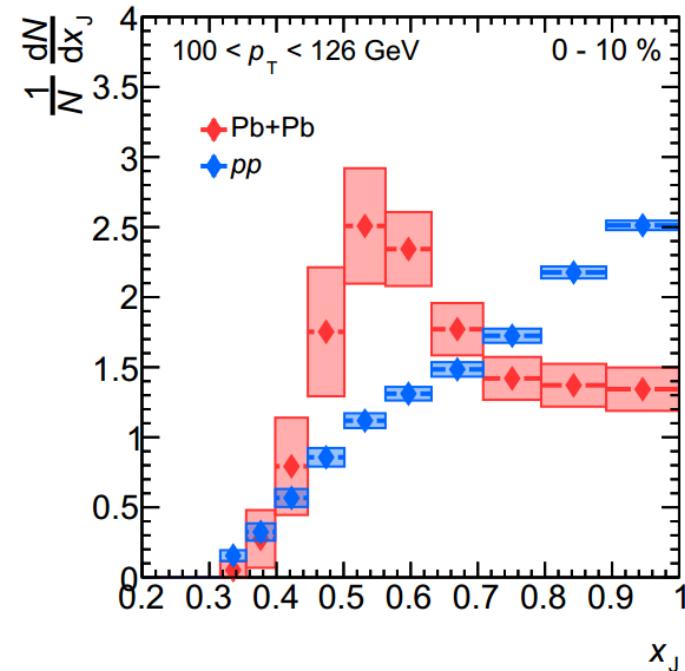
- Introduction
- p+p Baseline and Setup of Simulation
- Implementation of In-medium Evolution
- Results
  - ✓ Transverse momentum balance of  $b\bar{b}$  dijet
  - ✓ Angular correlation of  $b\bar{b}$  dijet
- Summary

# Dijet Asymmetry



$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}, \Delta\phi > \frac{\pi}{2}$$

**ATLAS Collaboration**  
**Phys. Rev. Lett. 105 (2010) 252303**

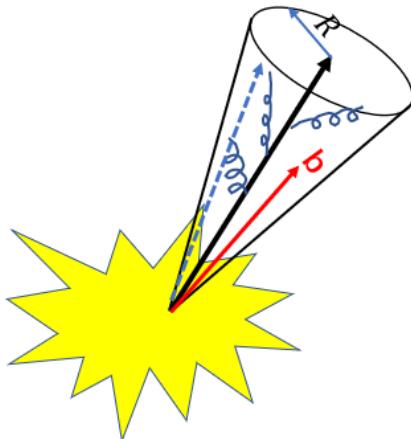


$$x_J = \frac{p_{T2}}{p_{T1}}, \Delta\phi > \frac{7\pi}{8}$$

**ATLAS Collaboration**  
**Phys. Lett. B774 (2017) 379**

Initial parton flavor dependences determined in dijets ?

# Inclusive b-jet R<sub>AA</sub>



b jet---at least one  $b$  or  $\bar{b}$  quark inside the jet cone  
with the jet radius parameter  $R$ .

Kurt Jung,  
talk on QM2017

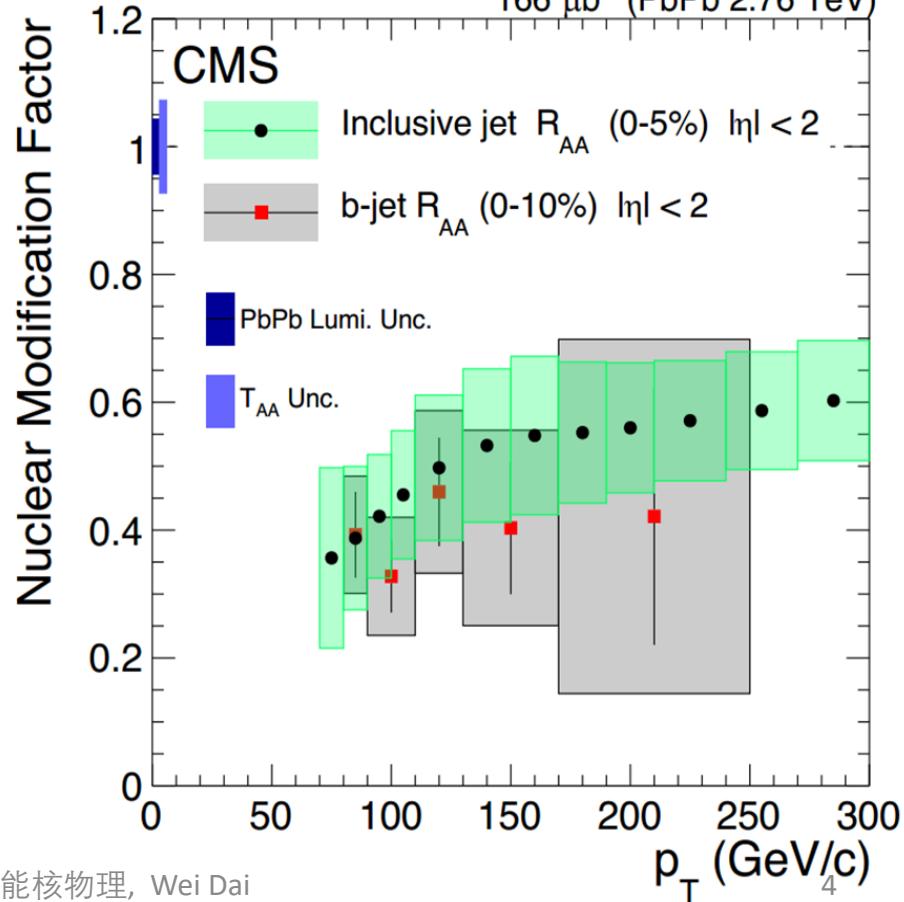
166  $\mu\text{b}^{-1}$  (PbPb 2.76 TeV)

$$R_{AA}^{\text{b-jet}}(p_T; R) = \frac{\frac{d^2\sigma^{AA}(p_T; R)}{dy dp_T}}{\langle N_{\text{bin}} \rangle \frac{d^2\sigma^{pp}(p_T; R)}{dy dp_T}}$$

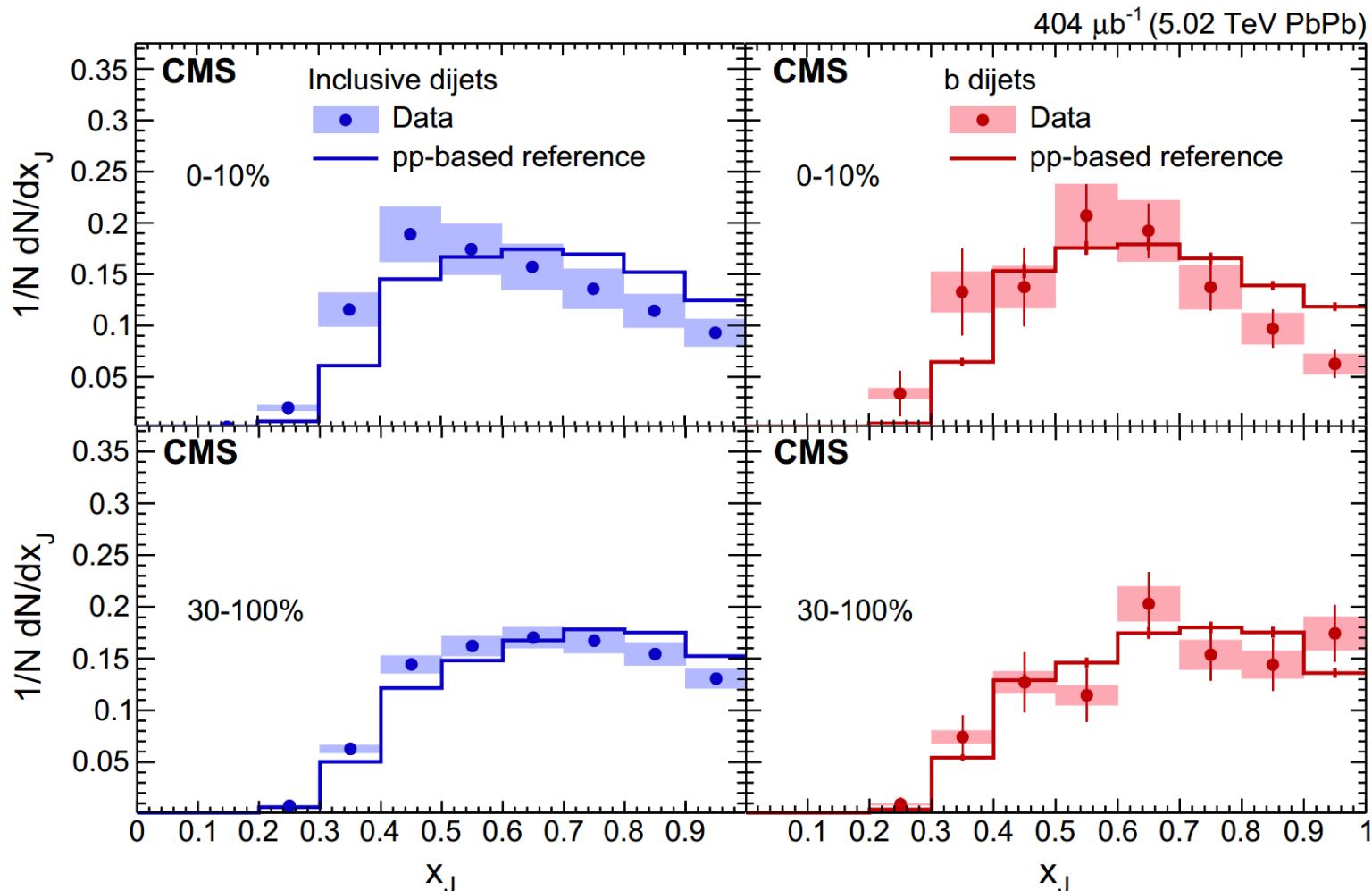
Jinrui Huang, Zhong-Bo Kang, Ivan Vitev  
PLB 726(2013)251-256

CMS collaboration , PRL 113(2014)132301

2019/06/23



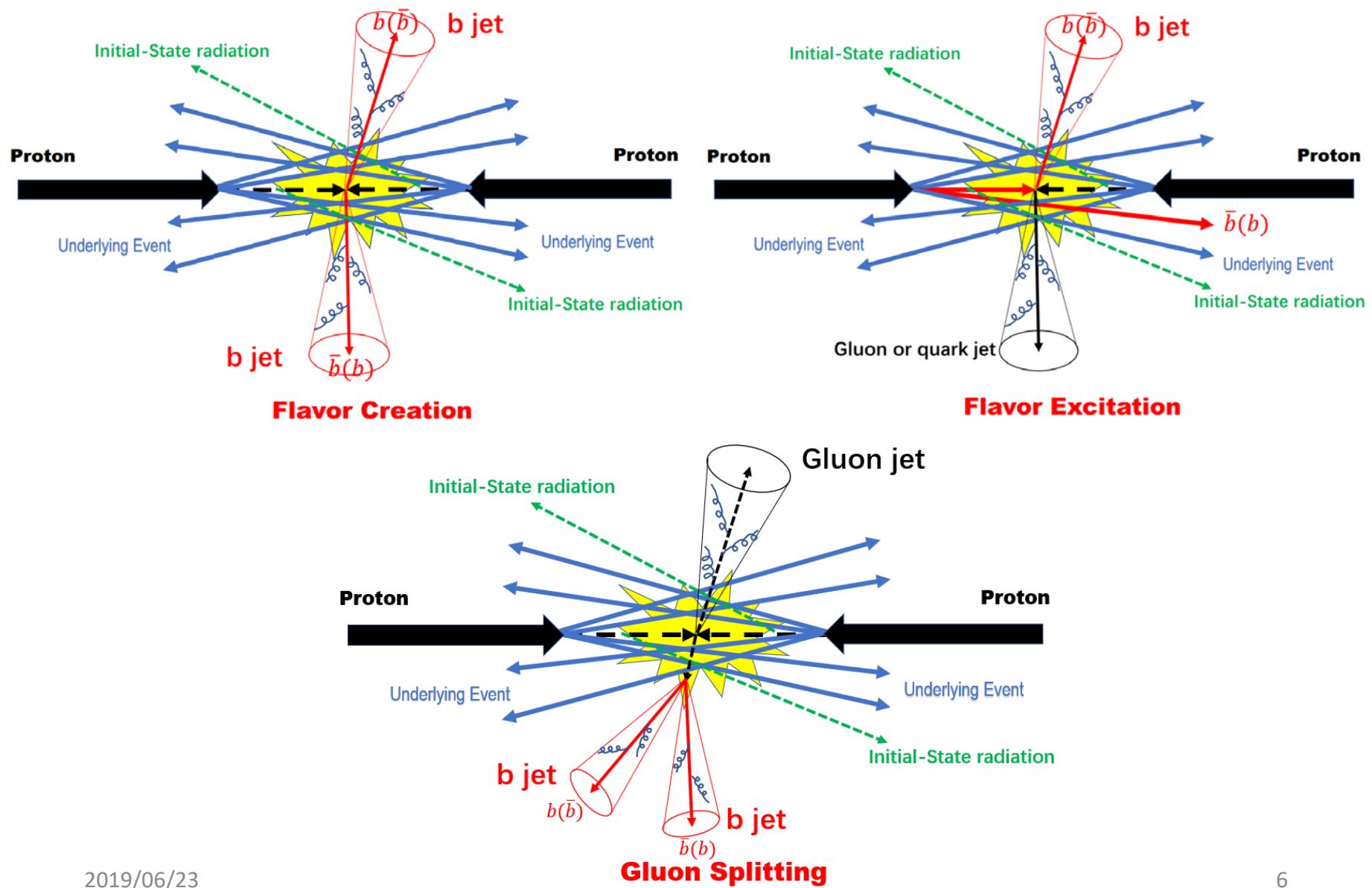
# $b\bar{b}$ Dijets Asymmetry



$$x_J = \frac{p_{T2}}{p_{T1}}, \Delta\phi > \frac{2\pi}{3}$$

CMS Collaboration JHEP 03(2018)181

# $b\bar{b}$ -Dijets Production

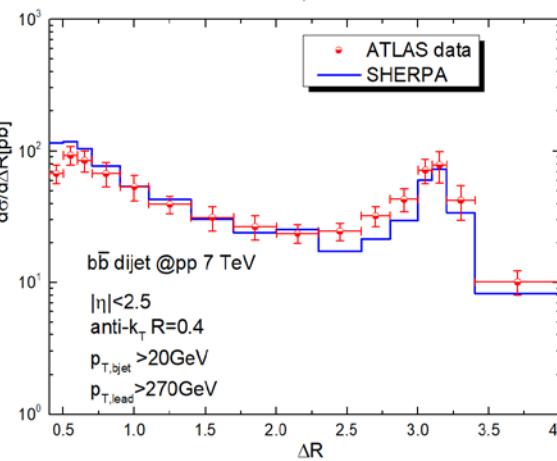
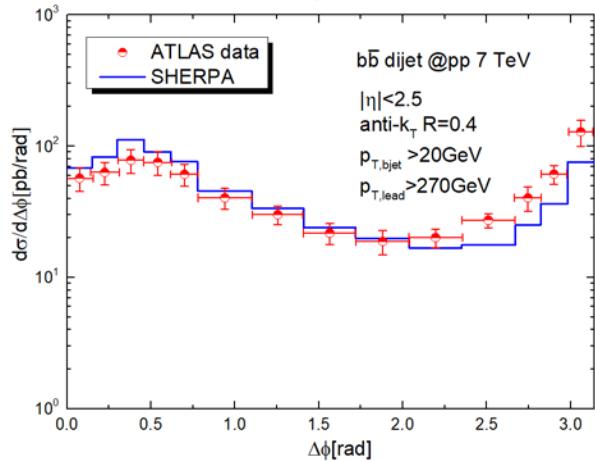
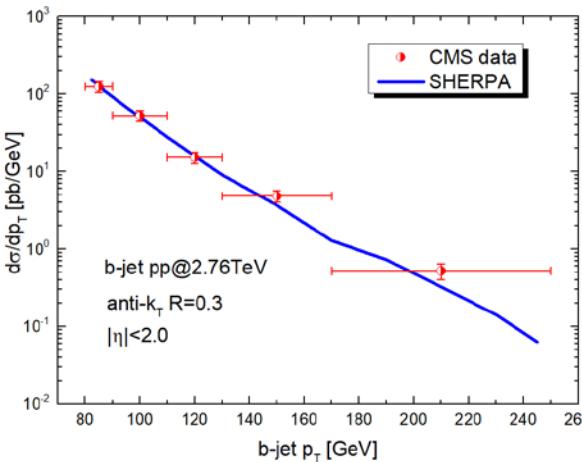
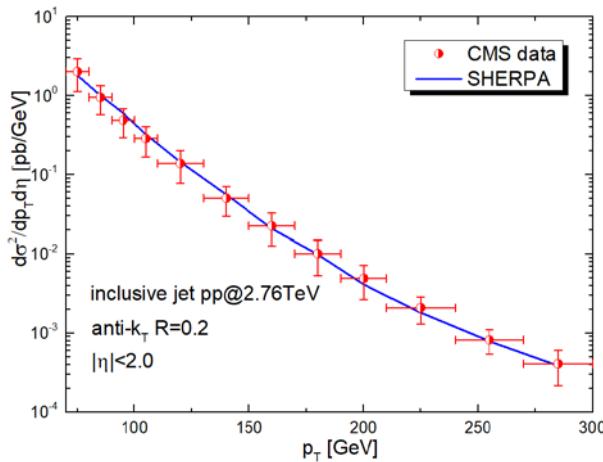


# p+p SETUP

Monte Carlo event generator: **Sherpa**

- The tree-level matrix elements---**Amegic** and **Comix**.
- The one-loop virtual corrections ---**BlackHat**.
- Parton Shower---**Catani-Seymour subtraction method** .
- Matching of NLO+PS: **MC@NLO**.
- **NNPDF3.0** NLO-5FS PDF .
- **FASTJET** ---jet reconstruction at parton level.

# p+p Baseline



Azimuthal angle difference between two b jets:

$$\Delta\phi = |\phi_{b1} - \phi_{b2}|$$

Angular distance:

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$$

A good pp baseline:  
not only on jet production but also on angular distribution  
[CMS ,PRC 96 015202(2017); CMS,PRL 113(2014)132301; ATLAS, Eur. Phys. J. C 76, no. 12, 670 (2016)]

# In-Medium E-loss: Heavy Quark

- For heavy quark, the discrete Langevin transport equations are used to describe the propagating of heavy quarks in the QGP. [[Eur. Phys. J. C 71, 1666 \(2011\); Physical Review C, 2009, 79\(5\): 054907.](#)]:

$$\vec{x}(t + \Delta t) = \vec{x}(t) + \frac{\vec{p}(t)}{E} \Delta t$$
$$\vec{p}(t + \Delta t) = \vec{p}(t) - \Gamma(p) \vec{p} \Delta t + \vec{\xi}(t) \Delta t \quad -\vec{p}_g \quad \text{Radiative E-loss}$$

The Fluctuation-Dissipation Theorem:  $\kappa = 2ET\Gamma = \frac{2T^2}{D_s}$ . Based on the Lattice calculation [[arXiv:1508.04543 \[hep-lat\]](#)],  $D_s$  is fixed as  $2\pi TD_s = 4$ .

- For light partons, the collisional energy loss is described by Hard Thermal Loop calculation. [[J.D. Bjorken, Fermilab preprint PUB-82/59-THY; Markus H.Thoma Physics Letters B 273 \( 1991\)128-132; R.B.Neufeld PRD 83 \(2011\) 065012](#)]

$$\frac{dE}{dz} = \frac{\alpha_s C_i m_D^2}{2} \ln \frac{\sqrt{ET}}{m_D}$$

# In-Medium E-loss: Radiative

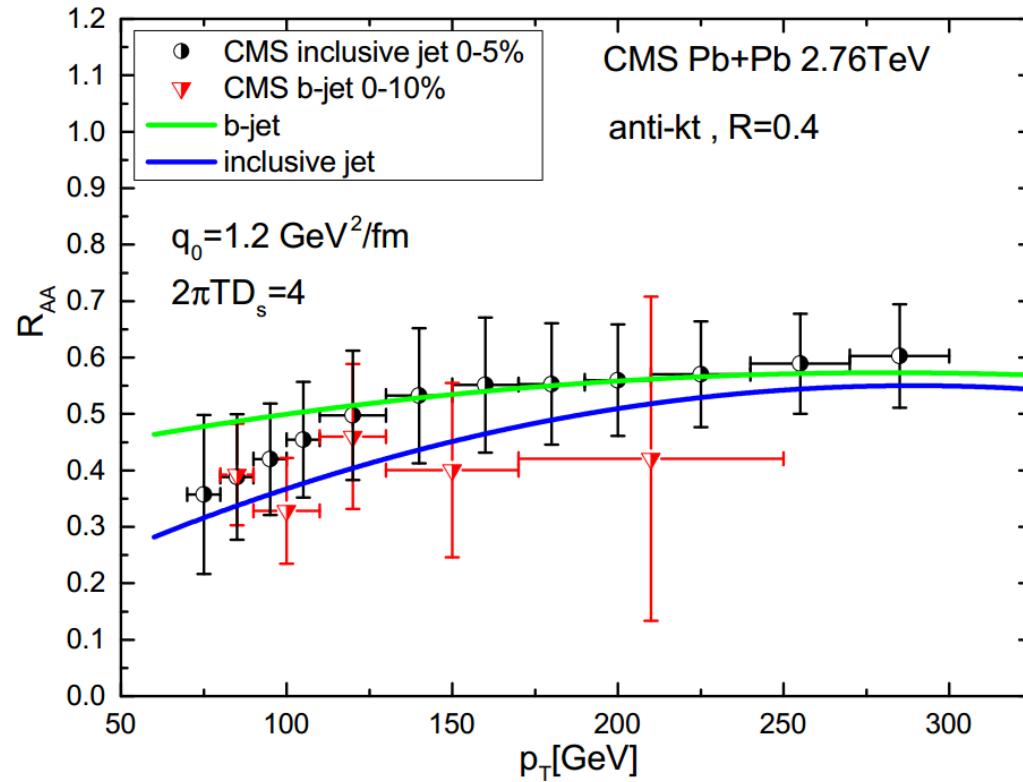
- The medium-induced radiative energy loss of final-state partons is described by the Higher-Twist scheme [X.-F. Guo and X.-N. Wang, Phys. Rev. Lett. 85, 3591(2000); B.-W. Zhang, E. Wang, and X.-N. Wang, Phys. Rev. Lett. 93, 072301 (2004); A. Majumder, Phys. Rev. D85, 014023 (2012).]

$$\frac{dN_g}{dx dk_{\perp}^2 dt} = \frac{2\alpha_s C_A P(x) \hat{q}}{\pi k_{\perp}^4} \sin^2\left(\frac{t - t_i}{2\tau_f}\right) \left(\frac{k_{\perp}^2}{k_{\perp}^2 + x^2 M^2}\right)^4$$

Jet transport parameter  $\hat{q}_0$  employed the results of recent study of re-extraction of 6 mesons  $R_{AA}$  at the RHIC and LHC. (iEBE-VISHNU Hydro)  
[G.-Y. Ma, W. Dai, B.-W. Zhang and Enke Wang, EPJC 2019]

$$\hat{q}(\tau, r) = q_0 \frac{\rho^{QGP}(\tau, r)}{\rho^{QGP}(\tau_0, 0)} \frac{p^\mu u_\mu}{p^0}$$

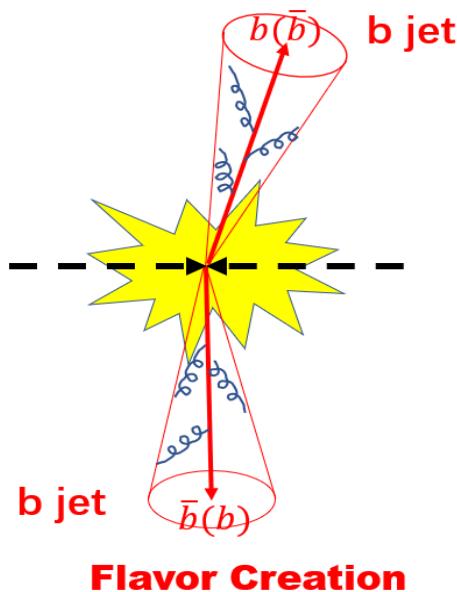
# $R_{AA}$ for Inclusive Jet and B-jet



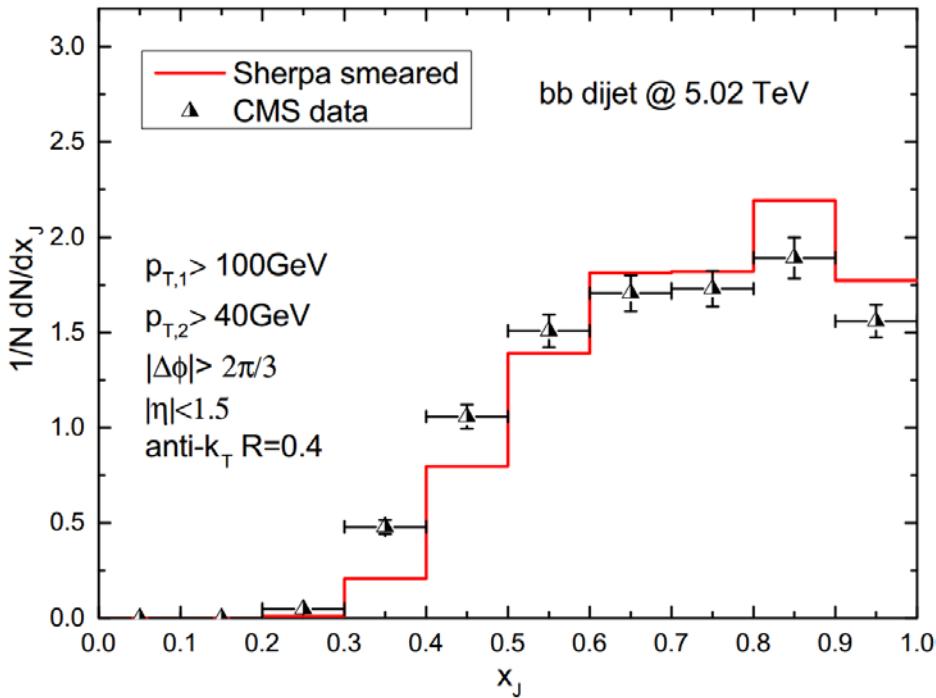
arXiv:1806.06332

- Our simulation can fairly describe the CMS data though b-jet  $R_{AA}$  slightly overestimates the data at low  $p_T$  region .
- The mass effect of the jet quenching trends to disappear.

# Transverse Momentum Imbalance- $x_J$



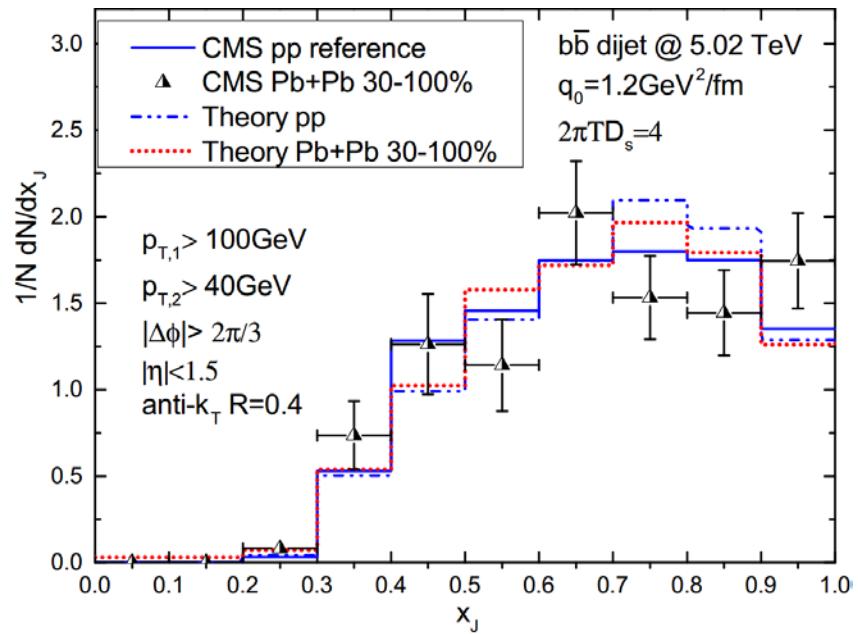
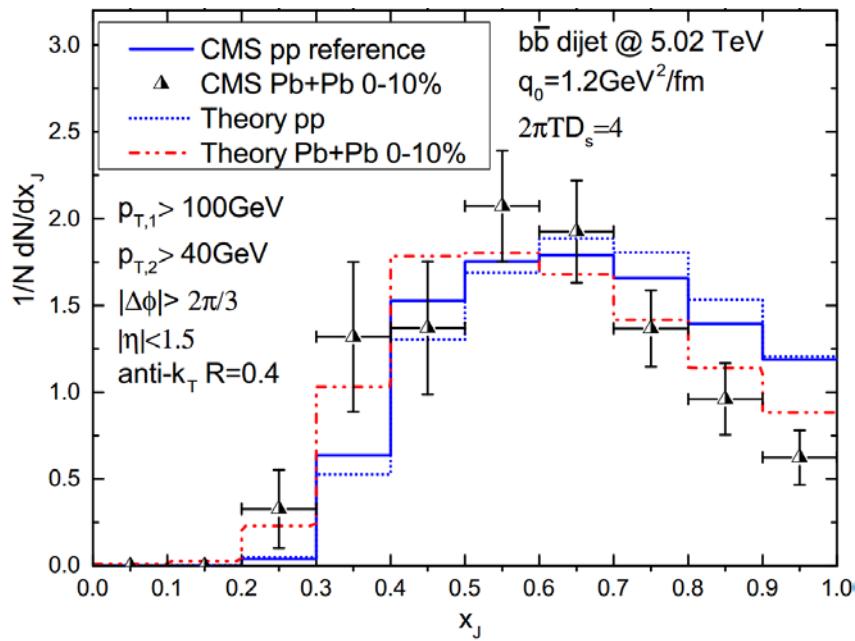
$$x_J = \frac{p_{T,2}}{p_{T,1}}$$



arXiv:1806.06332

- $|\eta| < 1.5, R = 0.4$
- Leading jet  $p_T > 100 \text{ GeV}$
- Sub-leading jet  $p_T > 40 \text{ GeV}$
- Leading and sub-leading jets must be b-tagged.
- $|\Phi_1 - \Phi_2| > \frac{2\pi}{3}$  suppresses the gluon splitting process, the two b jets are back-to-back.

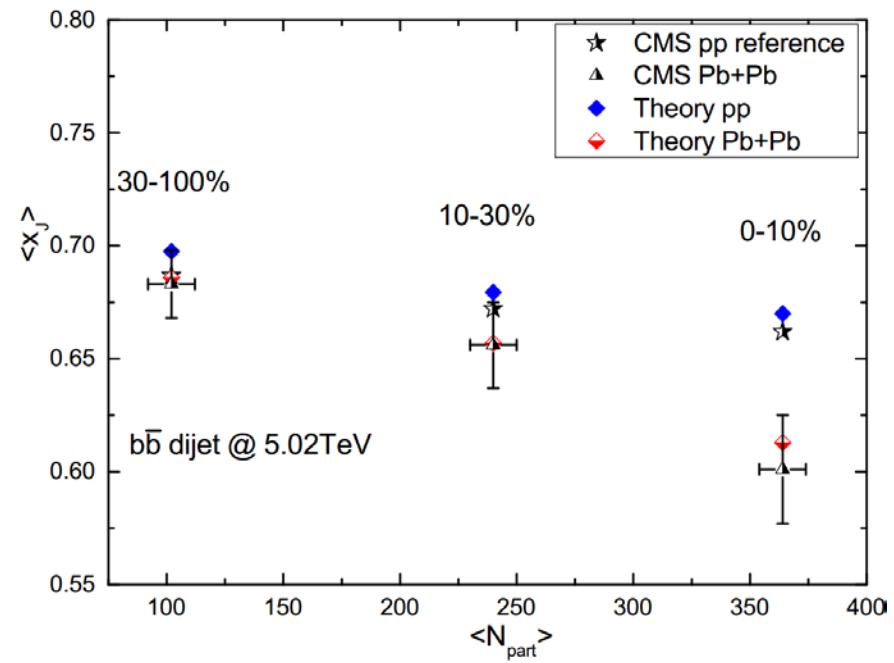
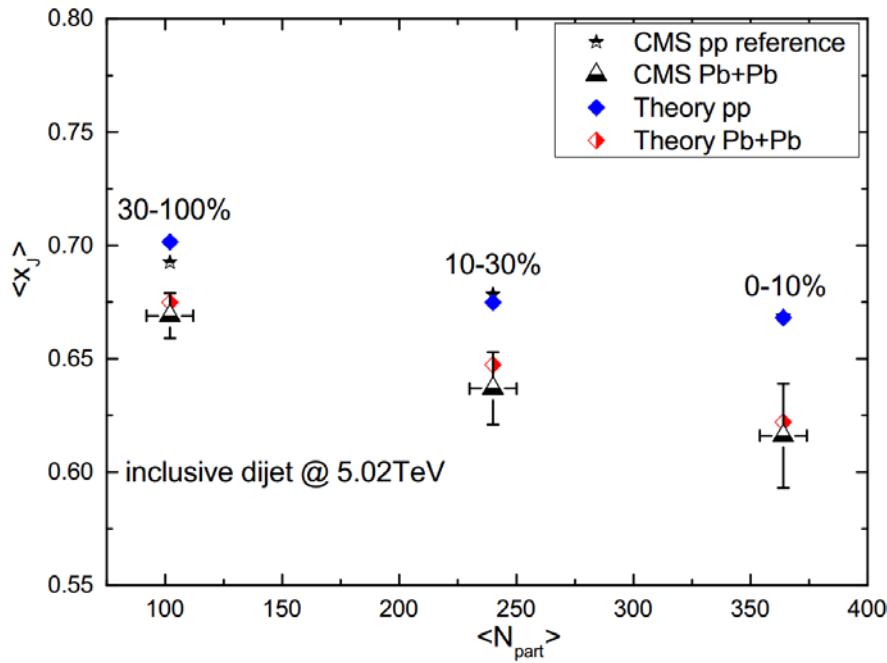
# $x_J$ for $b\bar{b}$ dijets



arXiv:1806.06332

- A visible shift of  $x_J$  towards smaller in the central Pb+Pb collision, relative to pp reference .
- Much smaller shift is observed in the peripheral Pb+Pb collision, indicating smaller energy loss.

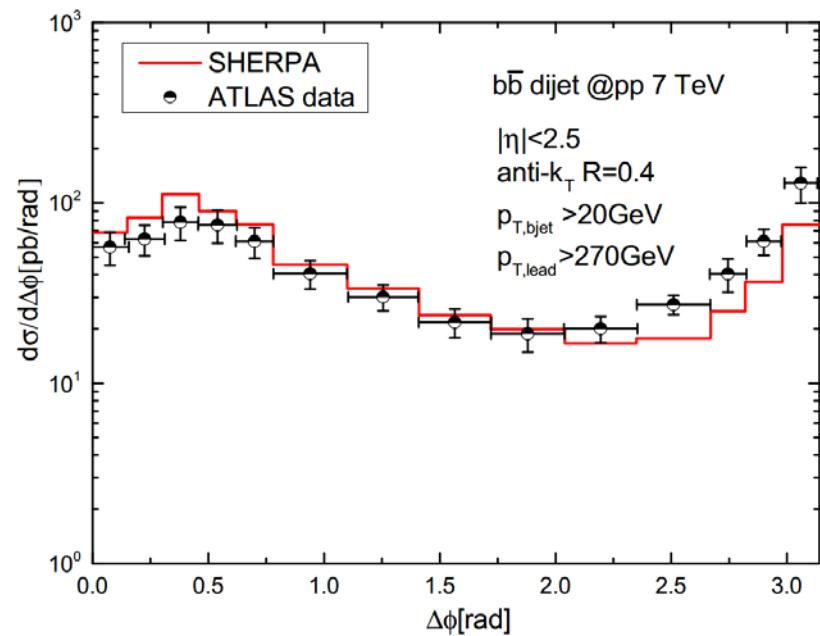
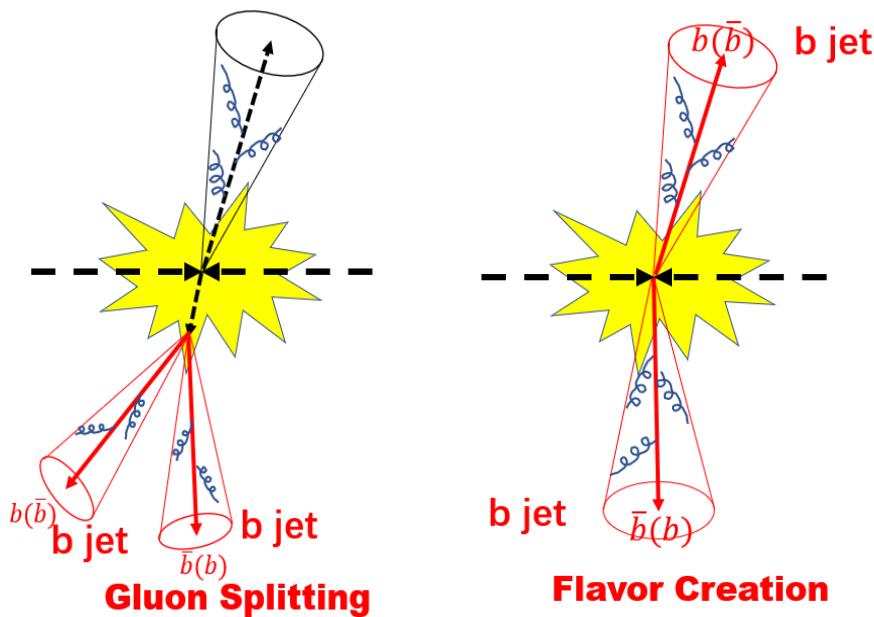
# $x_J$ for $b\bar{b}$ dijets



arXiv:1806.06332

- Imbalance of  $b\bar{b}$  dijet increasing due to jet quenching is visible in central Pb+Pb collision.
- A smaller energy loss of  $b\bar{b}$  dijets than inclusive dijet in peripheral Pb+Pb collision.

# Azimuthal Angle Correlation

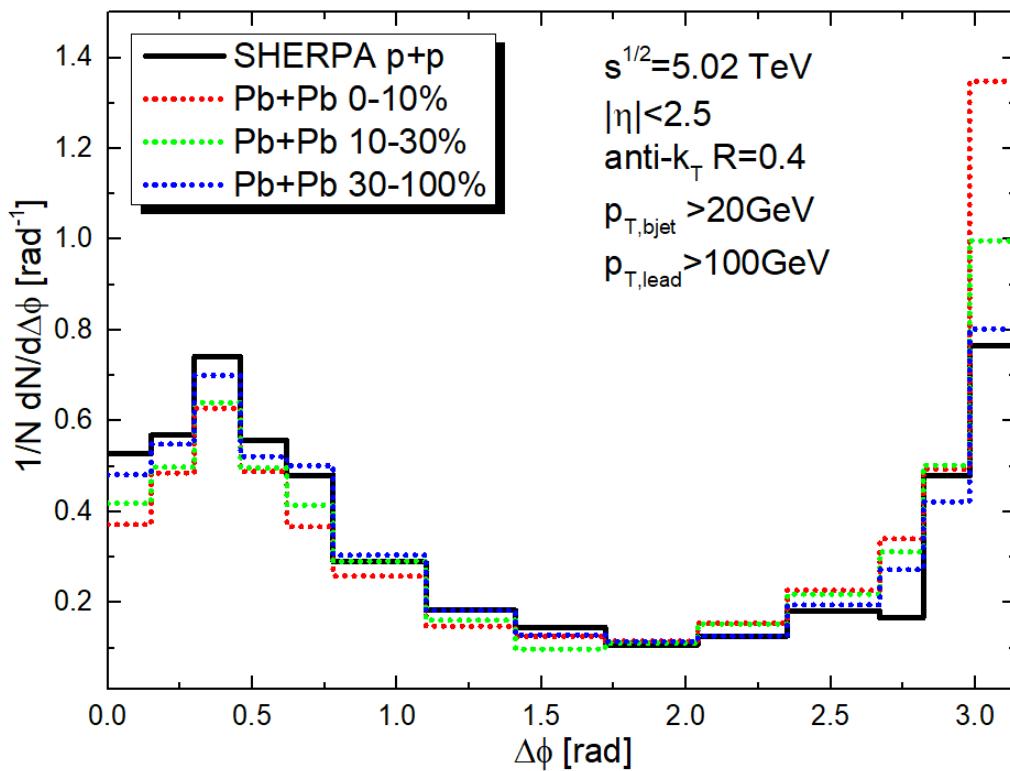


$$\Delta\phi = |\phi_{b1} - \phi_{b2}|$$

arXiv:1806.06332

- $0 < \Delta\phi < \pi, |\eta| < 2.5, R = 0.4$
- Leading jet  $p_T > 270\text{GeV}$
- $p_{T,bjet} > 20\text{GeV}$
- Leading and sub-leading jets are not required to be b-tagged.
- 2 b jets separated by  $\Delta R > 0.4$

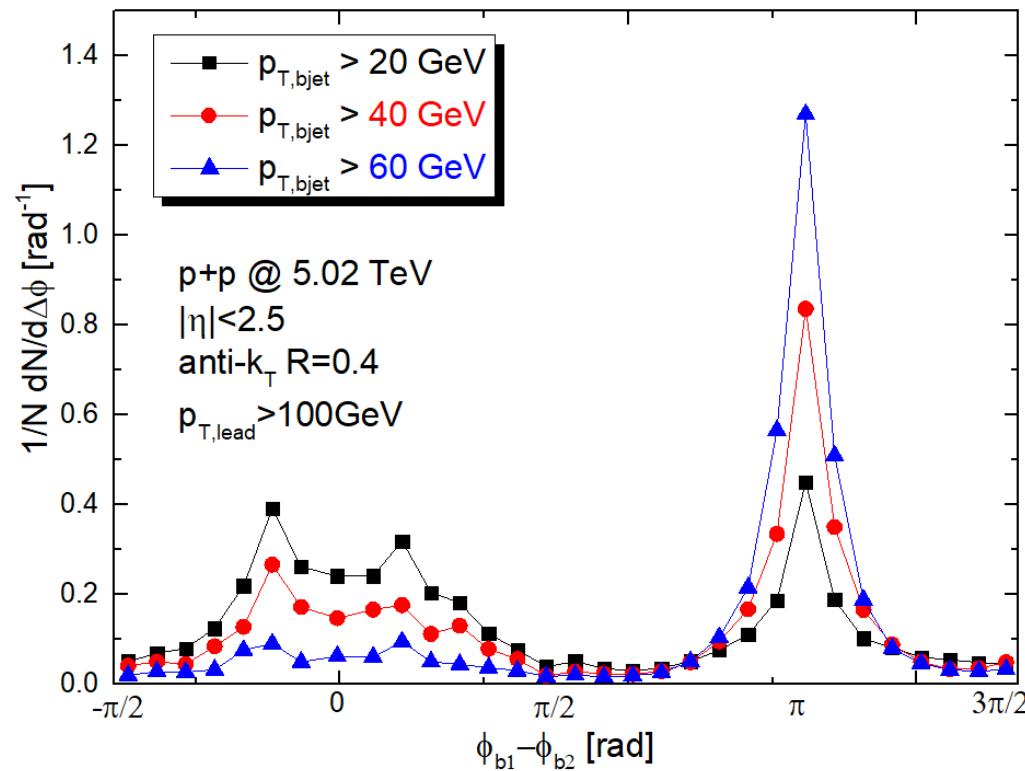
# Azimuthal Angle Distribution--- $\frac{1}{N} \frac{dN}{d\Delta\phi}$



arXiv:1806.06332

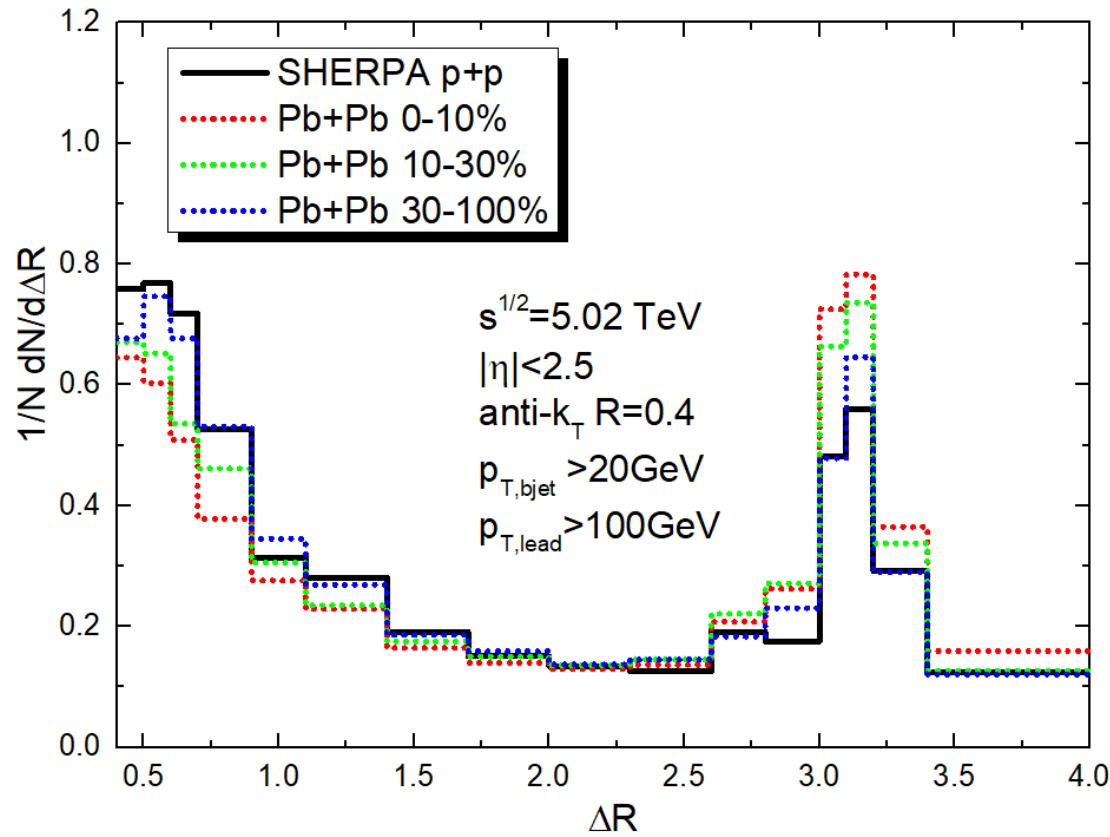
- The energy loss effect would suppress and broaden the near side (small  $\Delta\phi$ ) peak, and also enhance and sharp the away side (near  $\Delta\phi = \pi$ ) peak in the normalized  $\Delta\phi$  distribution.
- In the small angle region, it suffers stronger suppression relative to the large angle region.

# Azimuthal Angle Correlation



- The structure of azimuthal angle distribution between the two b jets was sensitive to the minimum  $p_T$  cut of b jet.
- GSP more likely produces two b jets with relatively lower  $p_T$  compared with FCR.

# Angular Distance Distribution--- $\frac{1}{N} \frac{dN}{d\Delta R}$



$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$$

- Suppression on small  $\Delta R$  and enhancement on large  $\Delta R$  with centrality dependence.

# Summary and Outlook

- SHERPA for pp baseline, light and heavy quark in-medium evolution in one framework have been implemented to study the  $b\bar{b}$  dijets production in Pb+Pb collisions.

- $x_J$

$x_J$  shift to smaller value consistent with CMS data, and a similar trend as that in dijet has been observed.

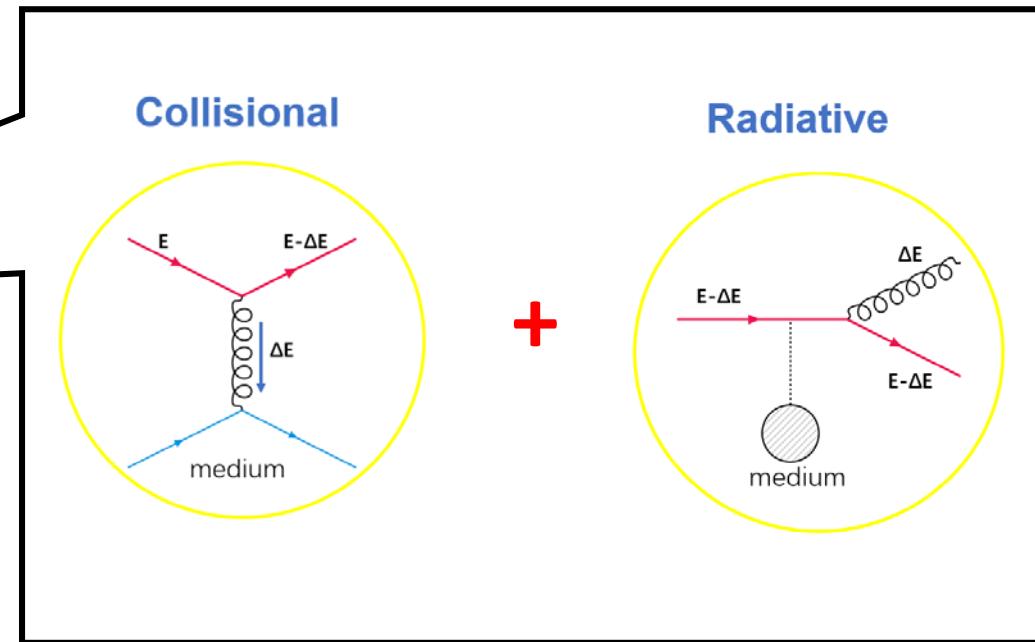
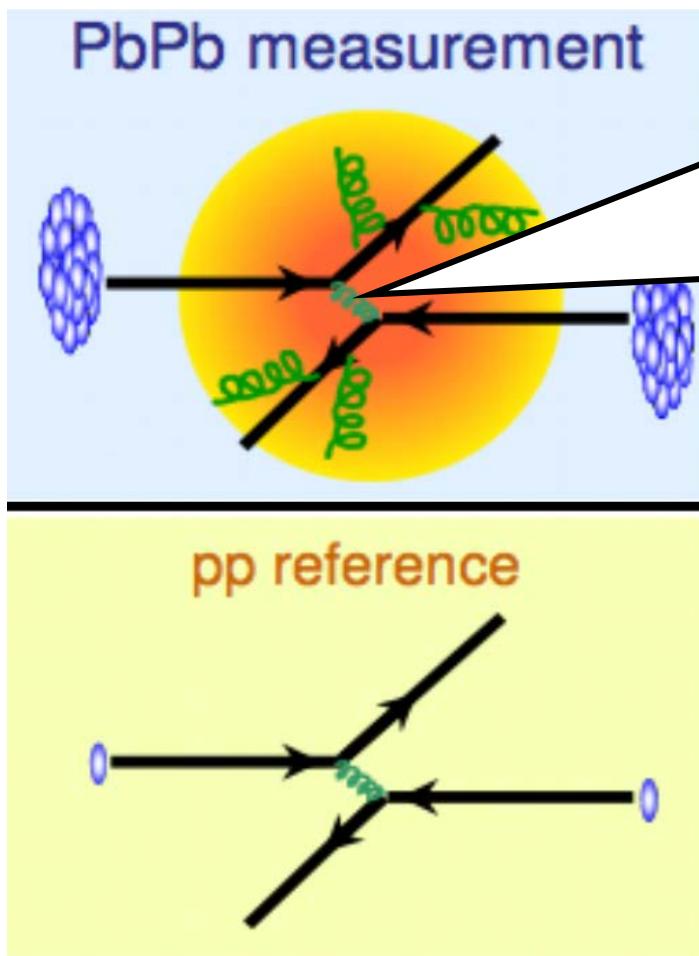
- $\Delta\phi$  and  $\Delta R$

Suppression on the near side and the enhancement on the away side in the normalized distribution are predicted.

- More exclusive channel has been investigated

$z+b\bar{b}$  production , soon to be submitted.

# In-Medium Energy Loss



# Backup-1 Gluon Sampling Method

- The average number of radiative gluon during a time step [[Phys.Rev. C94 \(2016\) no.1, 014909](#)]:

$$\langle N_g(t, \Delta t) \rangle = \Delta t \int dx dk_{\perp}^2 \frac{dN_g}{dx dk_{\perp}^2 dt}$$

- Assuming that the number of radiative gluon  $n$  obeys the Possion distribution  $P(n) = \frac{\langle N_g \rangle^n}{n!} e^{-\langle N_g \rangle}$ , the total probability for radiation during a time step  $\Delta t$ :

$$P_{rad}(t, \Delta t) = 1 - e^{-\langle N_g \rangle}$$

- Then it's available to sample the energy  $x E$  and transverse momentum  $k_{\perp}$  of the radiative gluons during every time step for arbitrary state  $(E, T, t)$  .
- The interval  $t - t_i$  would be reset when a radiation occurred .

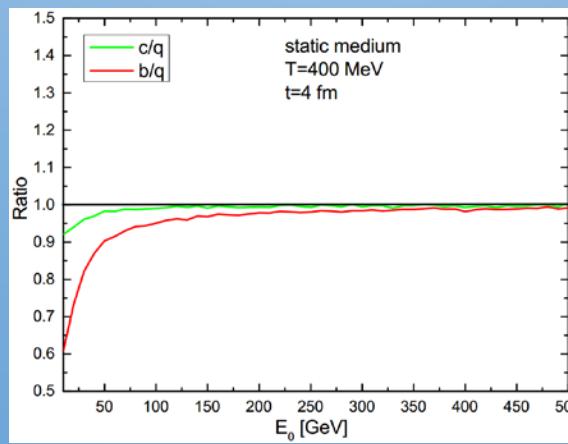
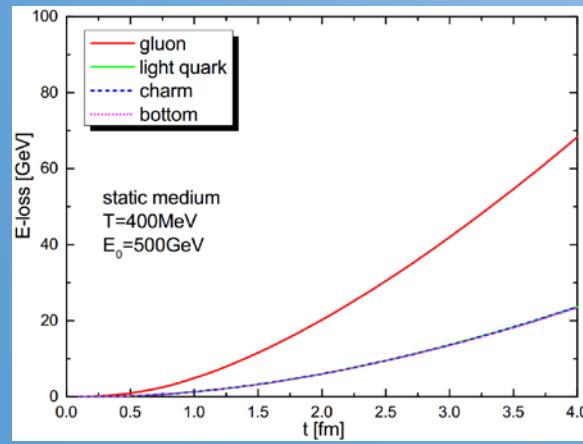
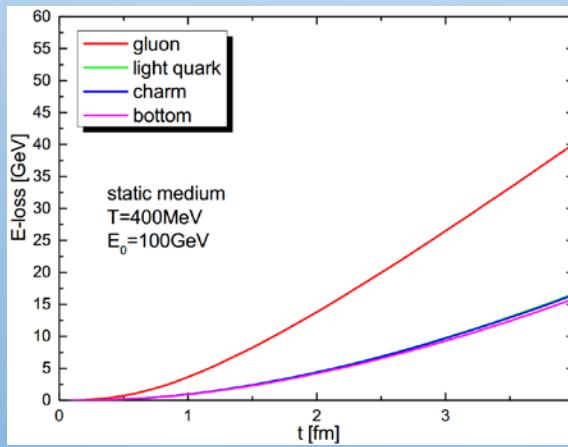
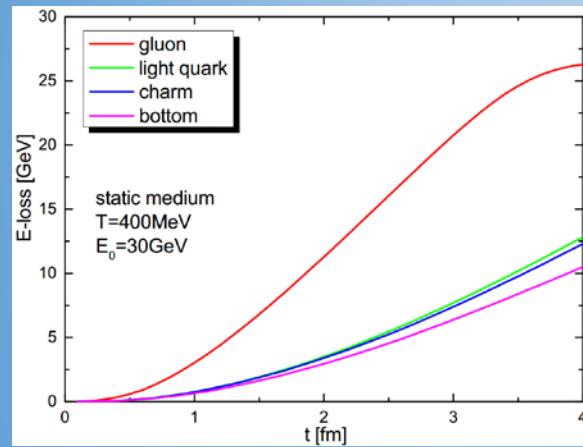
## Backup-2 Smearing parameter

The jet  $p_T$  resolution is parametrized according the following form [[arXiv:1607.03663](#)]:

$$\sigma(p_T)/p_T = \sqrt{C^2 + \frac{S^2}{p_T} + \frac{N^2}{p_T^2}}$$

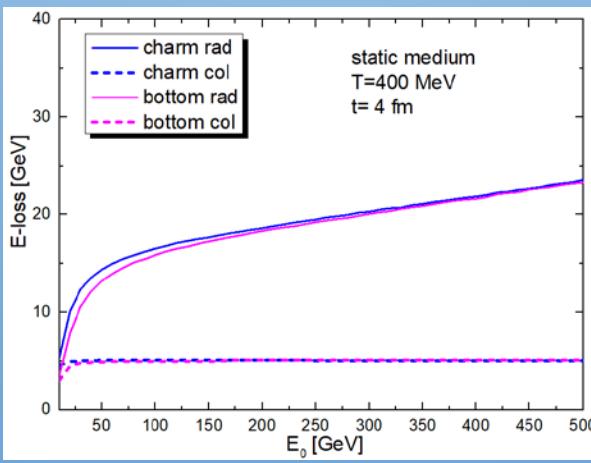
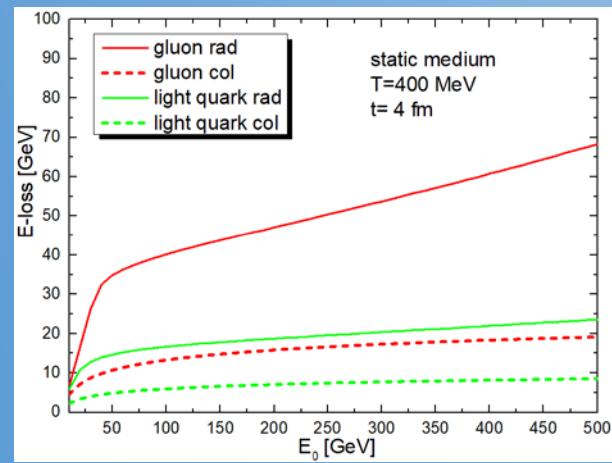
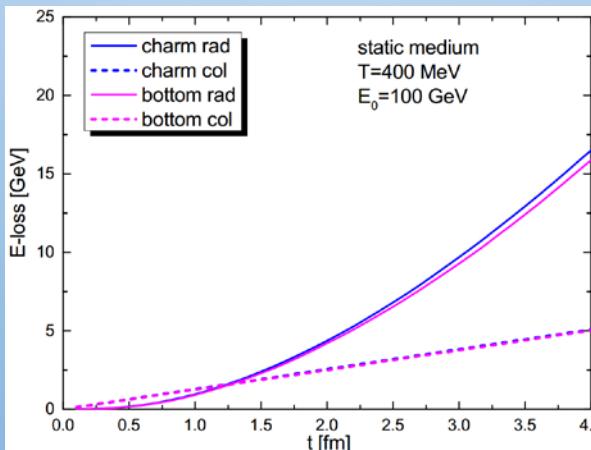
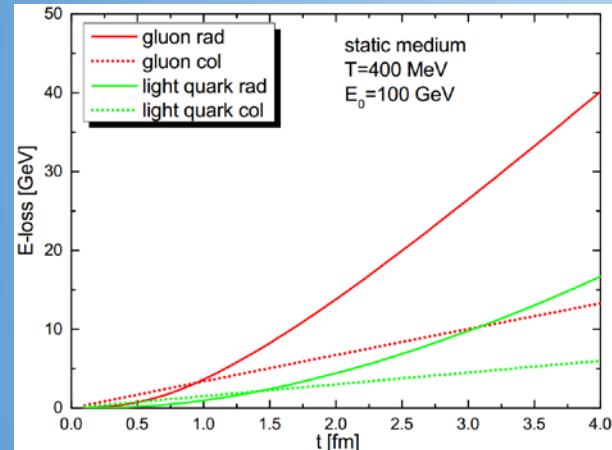
- In pp collisions, the constant C and stochastic S terms are 0.06 and  $0.8 \text{ GeV}^{-1}$
- In Pb+Pb collisions the S term is slightly larger value of  $1.0 \text{ GeV}^{-1}$ , due to the underlying event.
- The noise parameter (N) depends on collision centrality, according to  $N = 14.82 - \text{centrality}(\%)/5.40(\text{GeV}^2)$  subtraction.

# Radiative E-loss: gluon & light quark & heavy quark



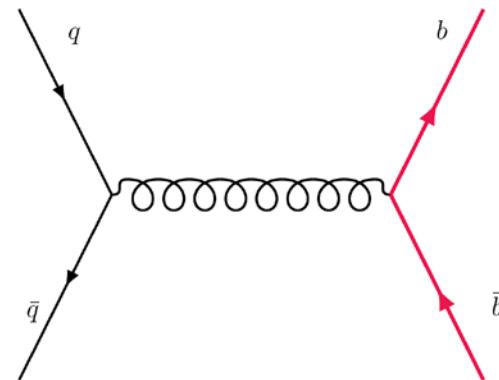
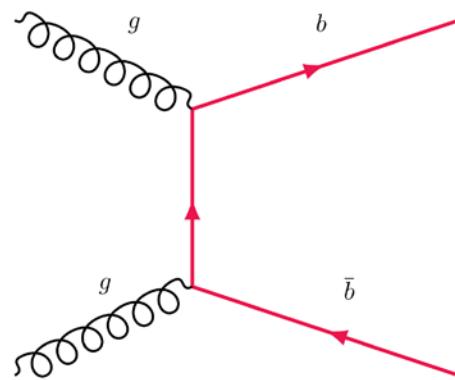
- The radiative energy loss of gluon is much larger than quarks.
- $\Delta E_g > \Delta E_q > \Delta E_c > \Delta E_b$ .
- Mass effect of heavy quarks trends to disappear as energy increasing.

# Radiative Vs. Collisional E-loss

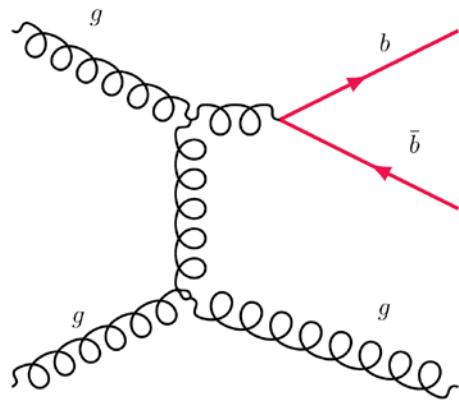


- In-medium energy loss is dominated by gluon radiation both for light and heavy flavor.
- Collisional E-loss shown no initial energy dependence.

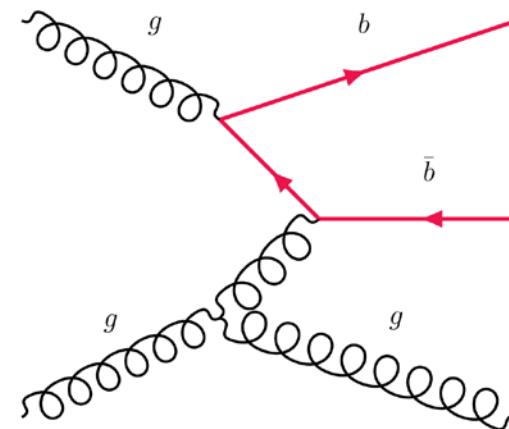
# Bottom Production



Flavor Creation (FCR)



Gluon Splitting(GSP)



Flavor Excitation(FEX)