



Study of K^* spin alignment with the toy model

Shuai Zhou

Central China Normal University

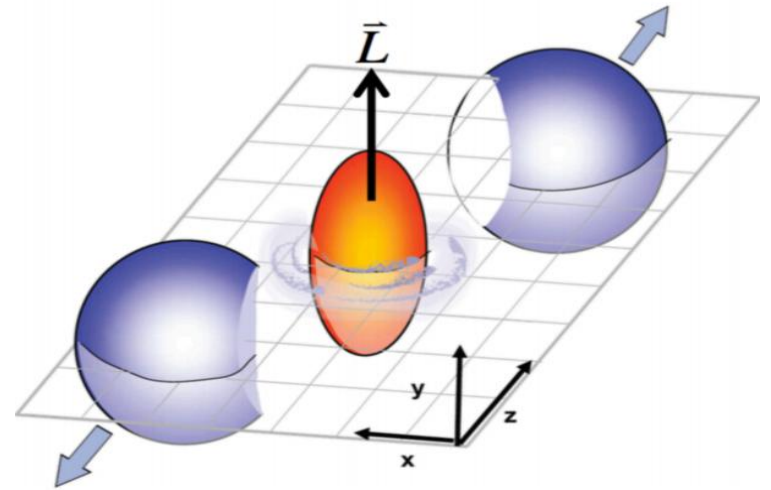
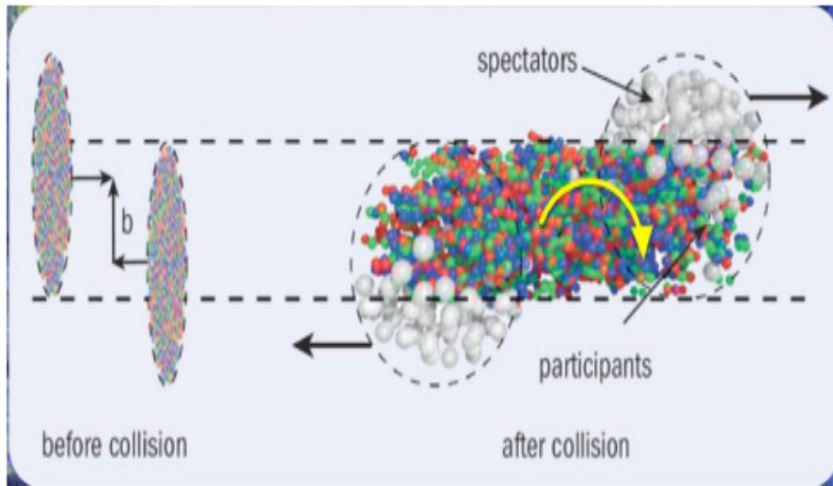


Outline

- **Introduction**
- **Analysis method**
- **Results**
- **Summary**



Introduction



Large initial angular momentum in non-central heavy ion collisions.

Due to spin-orbit coupling, it may result in net polarization of produced particles along the direction of initial angular momentum.



The information of $K^*(892)^0$

$$K^{*0} \longrightarrow K^+ + \pi^-$$

$$\overline{K}^{*0} \longrightarrow K^- + \pi^+$$

K^{*0} (The spin 1 vector meson)	
Mass	895.9 MeV
Quark Content	\bar{s} , d
Width	47.4 MeV
Lifetime	4.16 fm/c
Decay Chanel	$K^+ \pi^-$

The description of spin alignment

For spin-1 vector mesons, like the ϕ meson and K^* meson, their spin alignment can be described by a spin density matrix ρ . The ρ_{00} can be measured by angular distribution of decay daughter using:

$$\frac{dN}{d\cos\theta^*} \propto (1 + \rho_{00}) + (3\rho_{00} - 1)\cos^2\theta^*$$

θ^* is the angle between the polarization direction and the momentum direction of decay daughters in the rest frame of K^* meson.

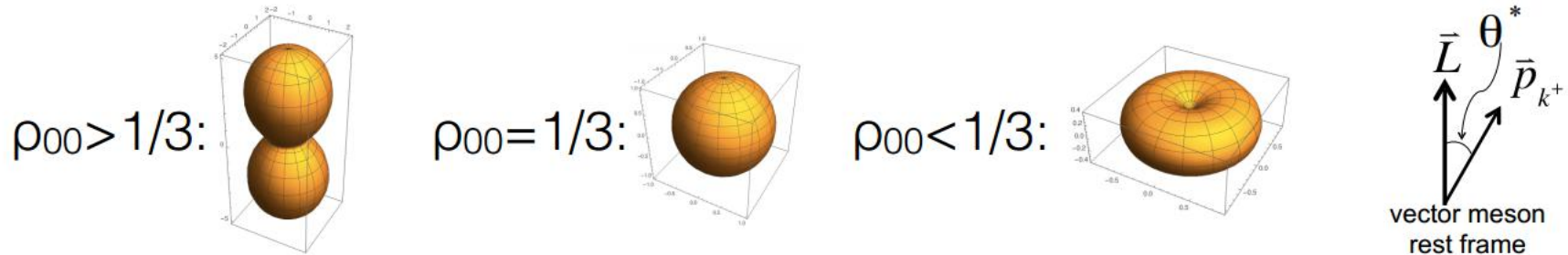
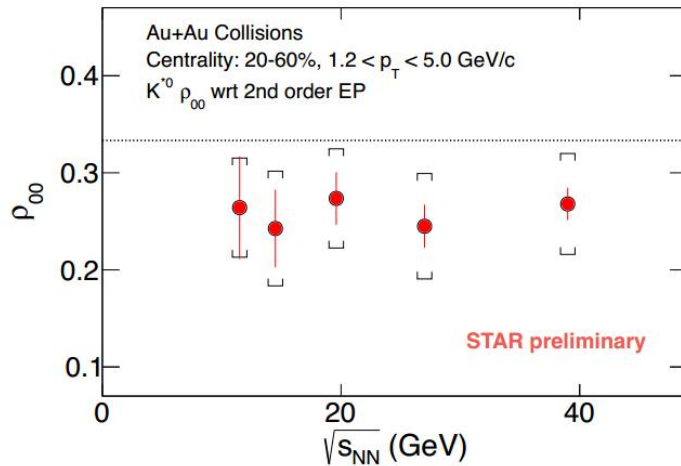


Figure: C.S. Zhou, Quark Matter 2018

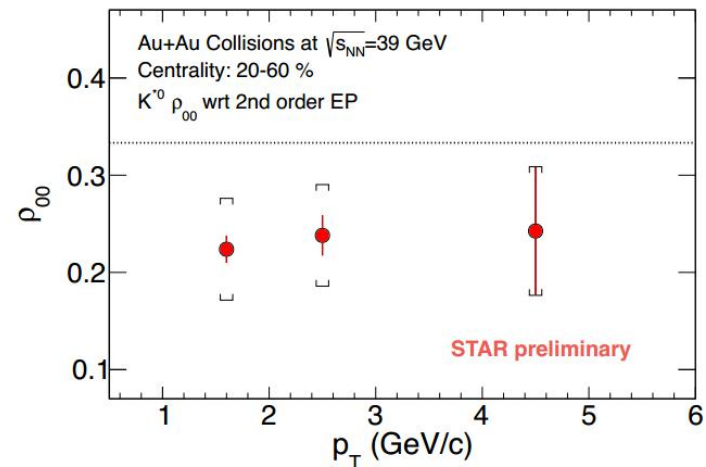


Results of K^{*0} spin alignment from STAR

Energy Dependence of $K^{*0} \rho_{00}$



p_T Dependence of $K^{*0} \rho_{00}$



C.S. Zhou, Quark Matter 2018

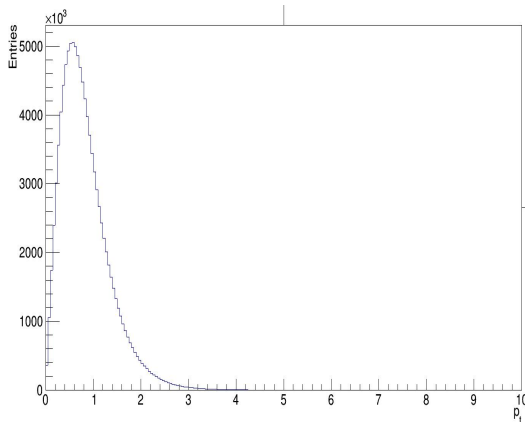
Non-1/3 ρ_{00} is observed.



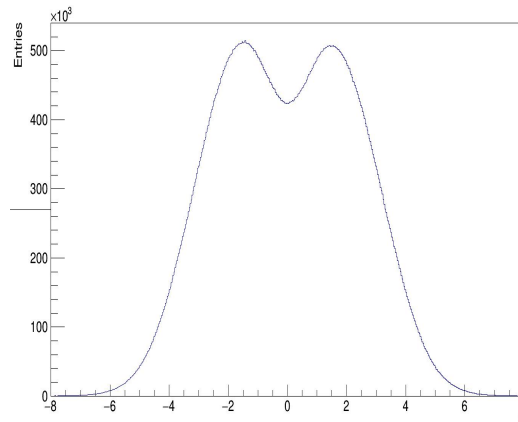
Analysis method

In the toy model, we input the p_T , η and ϕ information of K^* from AMPT (200GeV, String Melting). Phythia is used for the decay kinematics

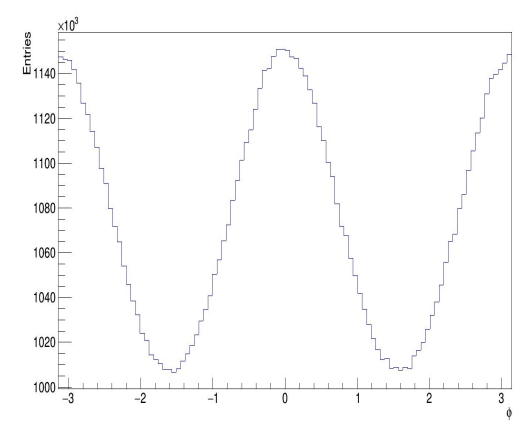
$$K^* \rightarrow K^+ + \pi^-.$$



p_T



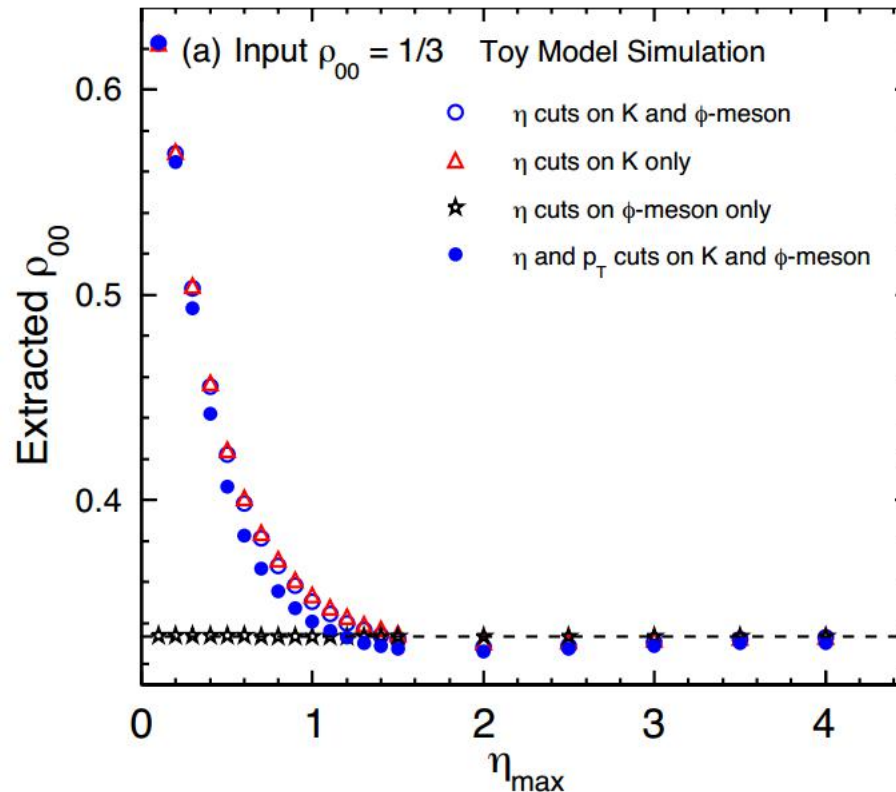
η



ϕ



The finite acceptance effect on ρ_{00} for ϕ meson

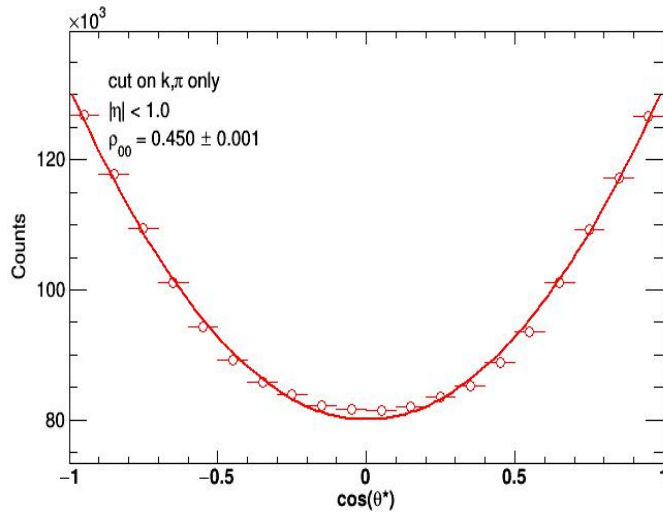


Shaowei Lan, Zi-Wei Lin Shusu Shi, Xu Sun, Phys. Lett. B 780, 319-324 (2018)

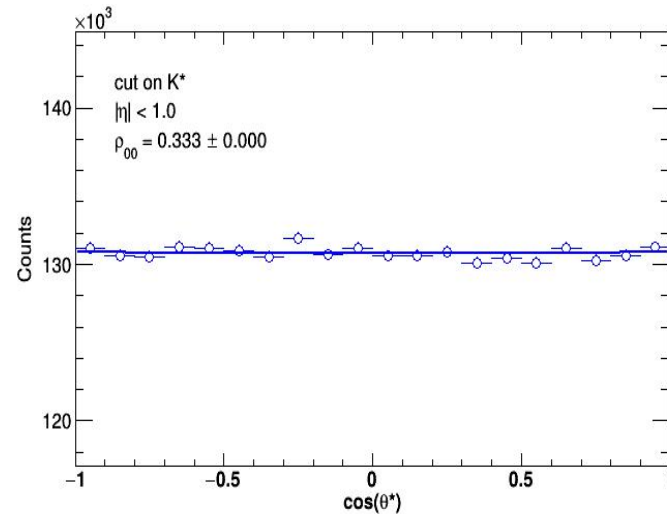
When the cut is below $|\eta| < 1$, a narrower η acceptance gives a significantly larger ρ_{00} value than the input.



Results of K^* spin alignment with toy model



a



b

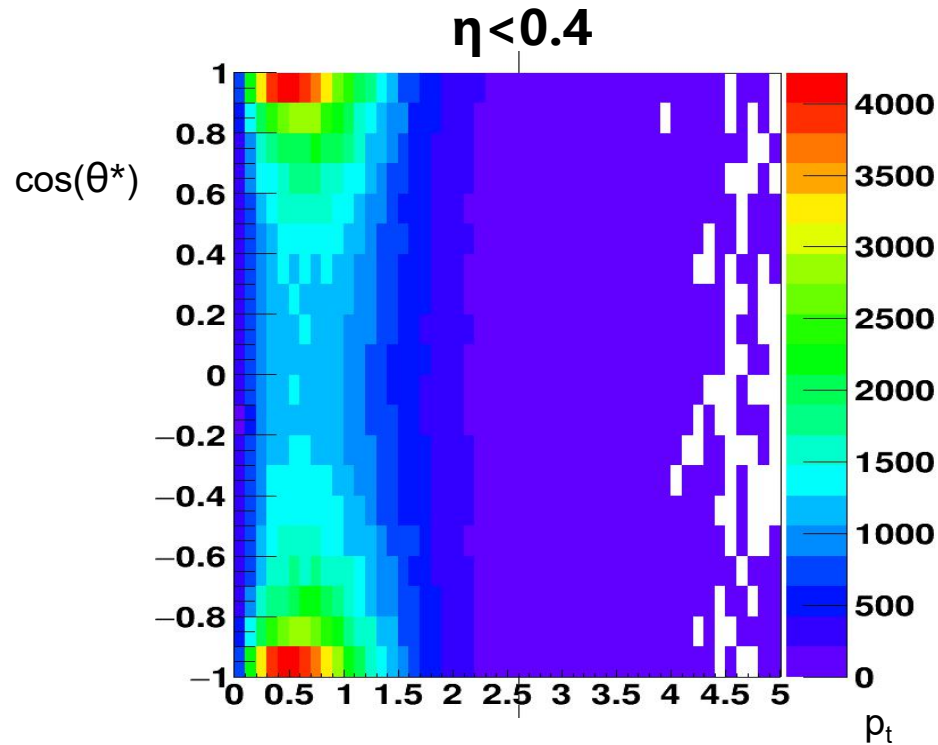
The curves represent fits based on Eq:

$$\frac{dN}{d\cos\theta^*} \propto (1 + \rho_{00}) + (3\rho_{00} - 1)\cos^2\theta^*$$

It shows that the extracted ρ_{00} is larger than the input value $1/3$ (a) when cut on daughters.



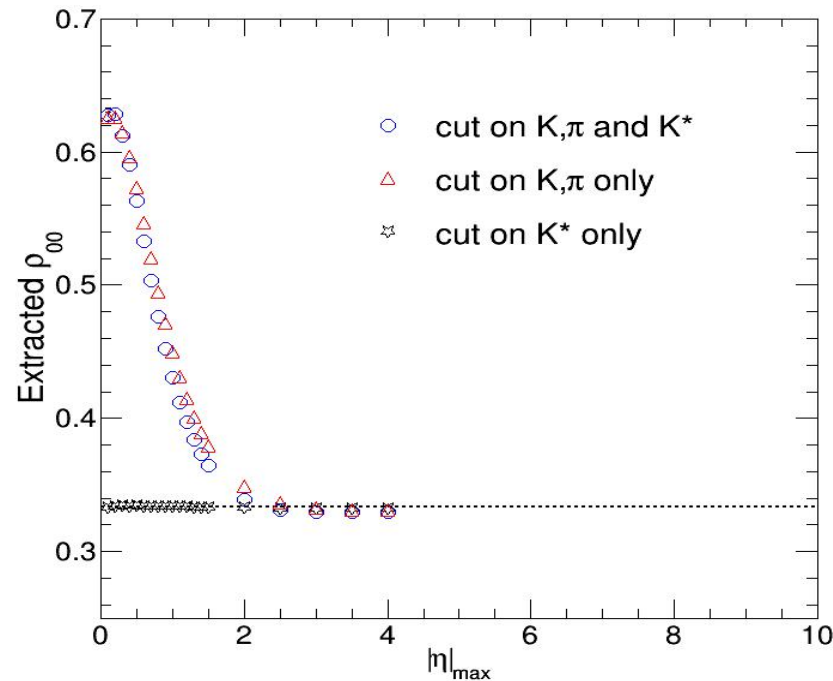
p_T versus $\cos(\theta^*)$



We can clearly observe that the η cut excludes more kaons from K^* decays around $\cos \theta^* \sim 0$.



η cut effect on ρ_{00}



It shows the extracted ρ_{00} of K^* mesons as a function of η acceptance with toy model.

The extracted ρ_{00} deviates from the input value significantly when we cut on daughter particles in a narrow η acceptance.

With in STAR η acceptance $|\eta| < 1$, the extracted $\rho_{00} = 0.425$ is larger than $1/3$.



Summary

- We used the toy model to study the effects of η cuts on the extracted polarization parameters ρ_{00} .
- We found η -cuts on daughters can lead larger value than input by using toy model.

Thank you!

Backups

The angular distribution $W(\theta, \phi) \equiv dN/d\Omega$ of the decay products is given by^[1]:

$$\begin{aligned}
 W(\theta, \phi) &= \frac{3}{4\pi} \left\{ \cos^2 \theta \rho_{00} + \sin^2 \theta (\rho_{11} + \rho_{-1-1})/2 \right. \\
 &\quad - \sin 2\theta (\cos \phi \operatorname{Re} \rho_{10} - \sin \phi \operatorname{Im} \rho_{10})/\sqrt{2} \\
 &\quad + \sin 2\theta (\cos \phi \operatorname{Re} \rho_{-10} + \sin \phi \operatorname{Im} \rho_{-10})/\sqrt{2} \\
 &\quad \left. - \sin^2 \theta [\cos(2\phi) \operatorname{Re} \rho_{1-1} - \sin(2\phi) \operatorname{Im} \rho_{1-1}] \right\}
 \end{aligned}$$

Here θ is the polar angle between the direction of motion of h and the quantization axis, ϕ is the azimuthal angle. By integrating over ϕ , we obtain

$$W(\theta) = \frac{3}{4} [(1 - \rho_{00}) + (3\rho_{00} - 1) \cos^2 \theta].$$

[1]: Physics Letters B 629 (2005) 20–26 , Zuo-Tang Liang a, Xin-Nian Wang

For hadronization scenario of constituent quark recombination in which both quarks and anti-quarks are polarized.

$$\rho_{00}^{K^*(\text{rec})} = \frac{1 - P_q P_s}{3 + P_q P_s}.$$

For Fragmentation of polarized quarks

$$\rho_{00}^{K^{*0}(\text{frag})} = \frac{f_s}{n_s + f_s} \frac{1 + \beta P_q^2}{3 - \beta P_q^2} + \frac{n_s}{n_s + f_s} \frac{1 + \beta P_s^2}{3 - \beta P_s^2}$$

$$P_q = -\frac{\pi}{4} \frac{\mu p}{E(E + m_q)} \quad \text{is the global quark polarization}$$

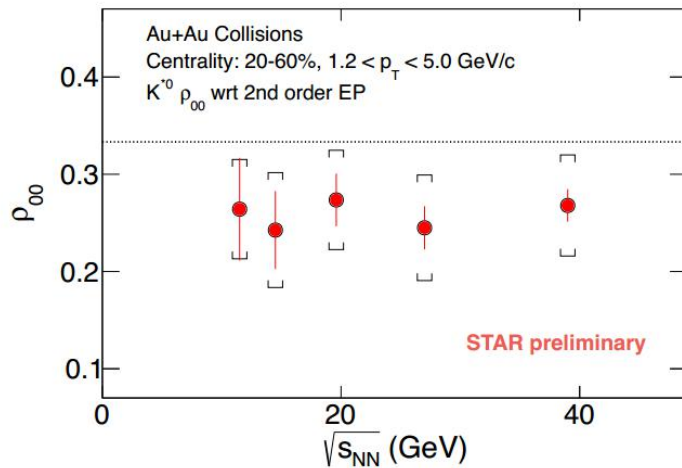
$$P_{\bar{q}}^{\text{frag}} = -\beta P_q \quad \text{is the polarization of the (anti-)quark created in the fragmentation process}$$

n_s and f_s are the strange quark abundances relative to up or down quarks in QGP and quark fragmentation, respectively.

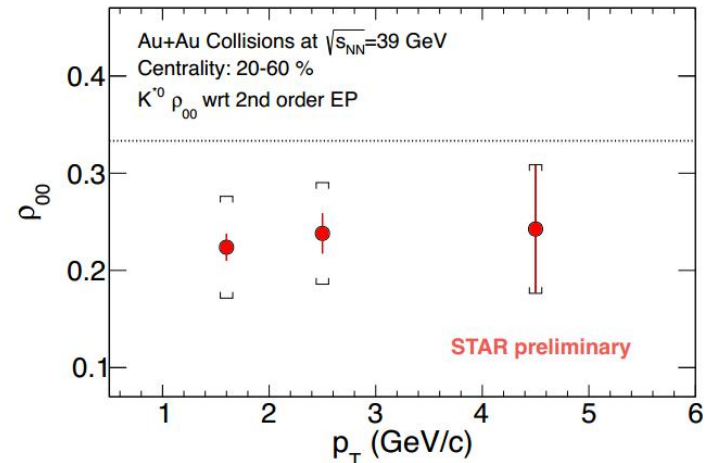


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Chensheng Zhou.Quark matter.2018

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