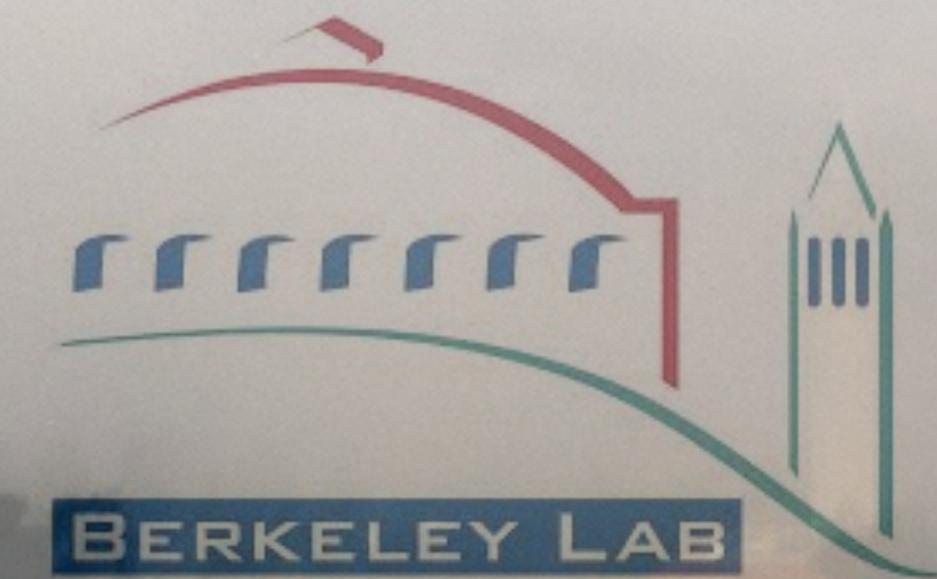


Chiral Perturbation Theory and Lattice QCD

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<https://c51.lbl.gov/~walkloud/grav/>



Effective Field Theory and Lattice QCD for Fundamental Nuclear Physics Research

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<https://c51.lbl.gov/~walkloud/grav/>



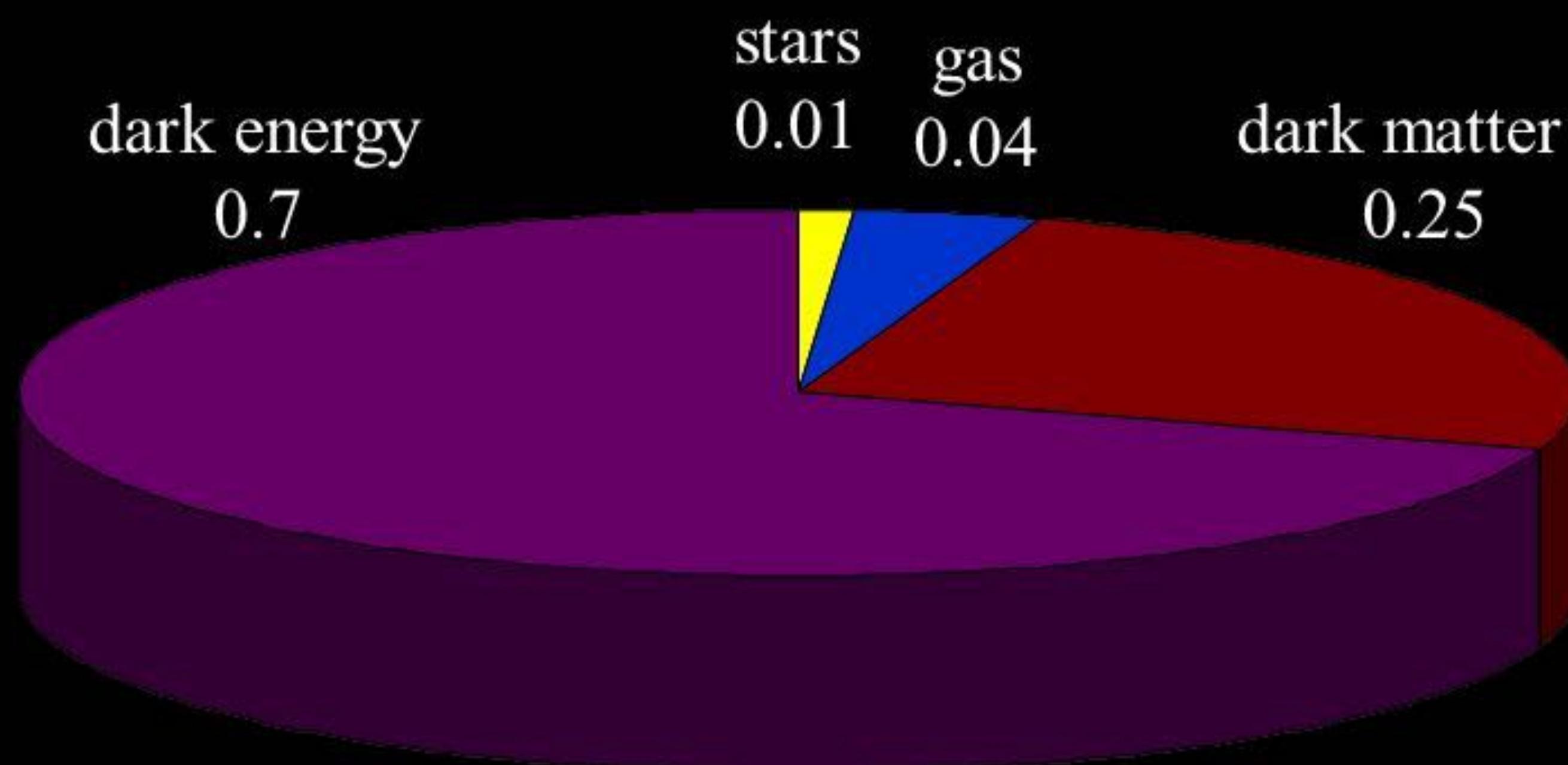
This Talk

- What are the driving science questions we are trying to understand?
- Why is it challenging to apply LQCD to nucleons and nuclei?
- Why is this an exciting time to use LQCD for basic physics?

Science Drivers

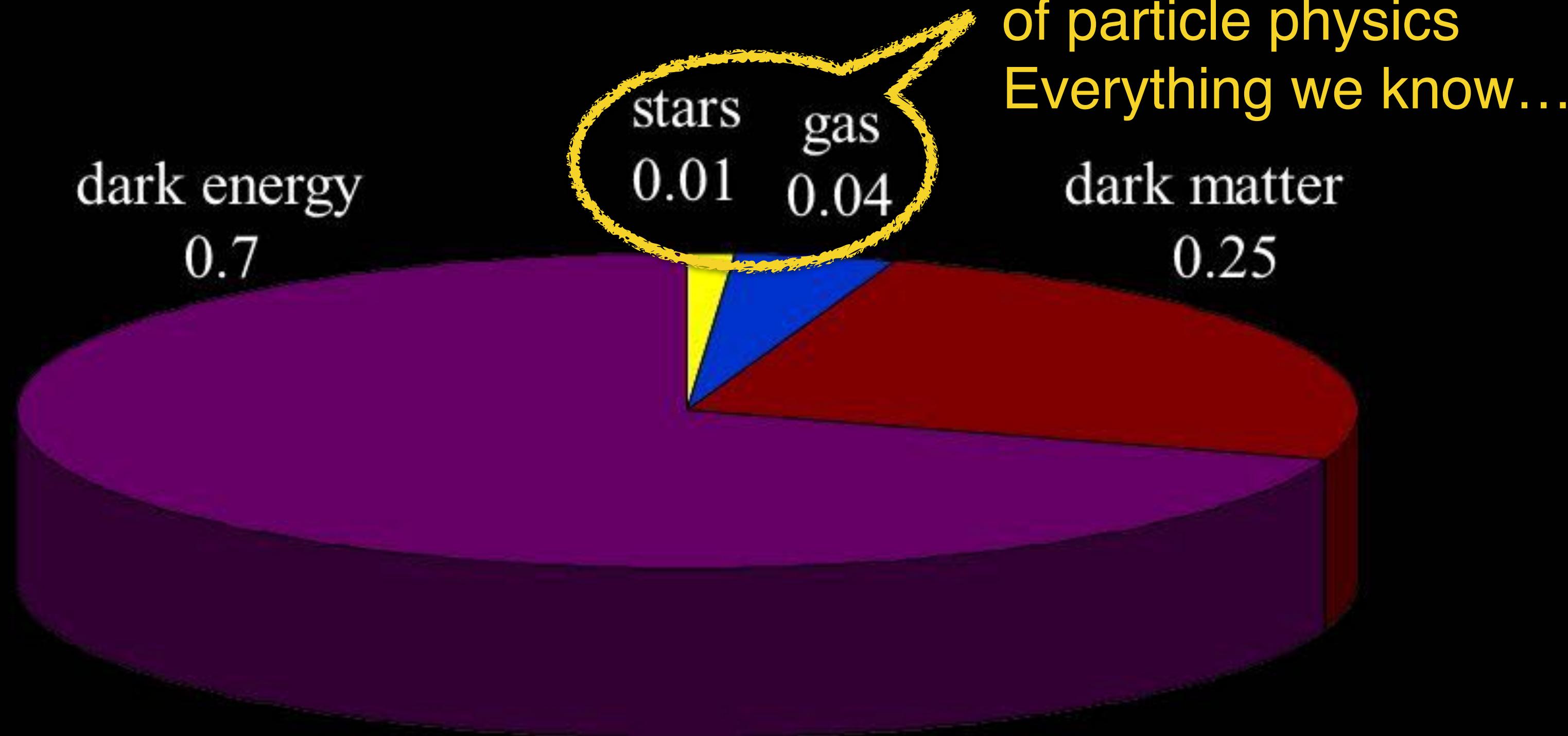
- Understanding the emergence of nuclear physics from QCD and the electroweak interactions of the Standard Model (SM)
 - How finely tuned is the universe we live in?
 - Mn-Mp, Bd, triple alpha process, ...
 - What is the pp fusion rate in the sun?
 - What is the strength of the three-neutron force?
 - ...
 - Testing the SM and its Fundamental Symmetries at low energies in nuclear environments (nuclei)
 - Why is the universe composed of matter?
 - neutrinoless double beta decay ($0\nu\beta\beta$), permanent electric dipole moments (EDMs)
 - Does Dark Matter interact with matter other than gravitationally?
 - ...

cosmic mass/energy budget



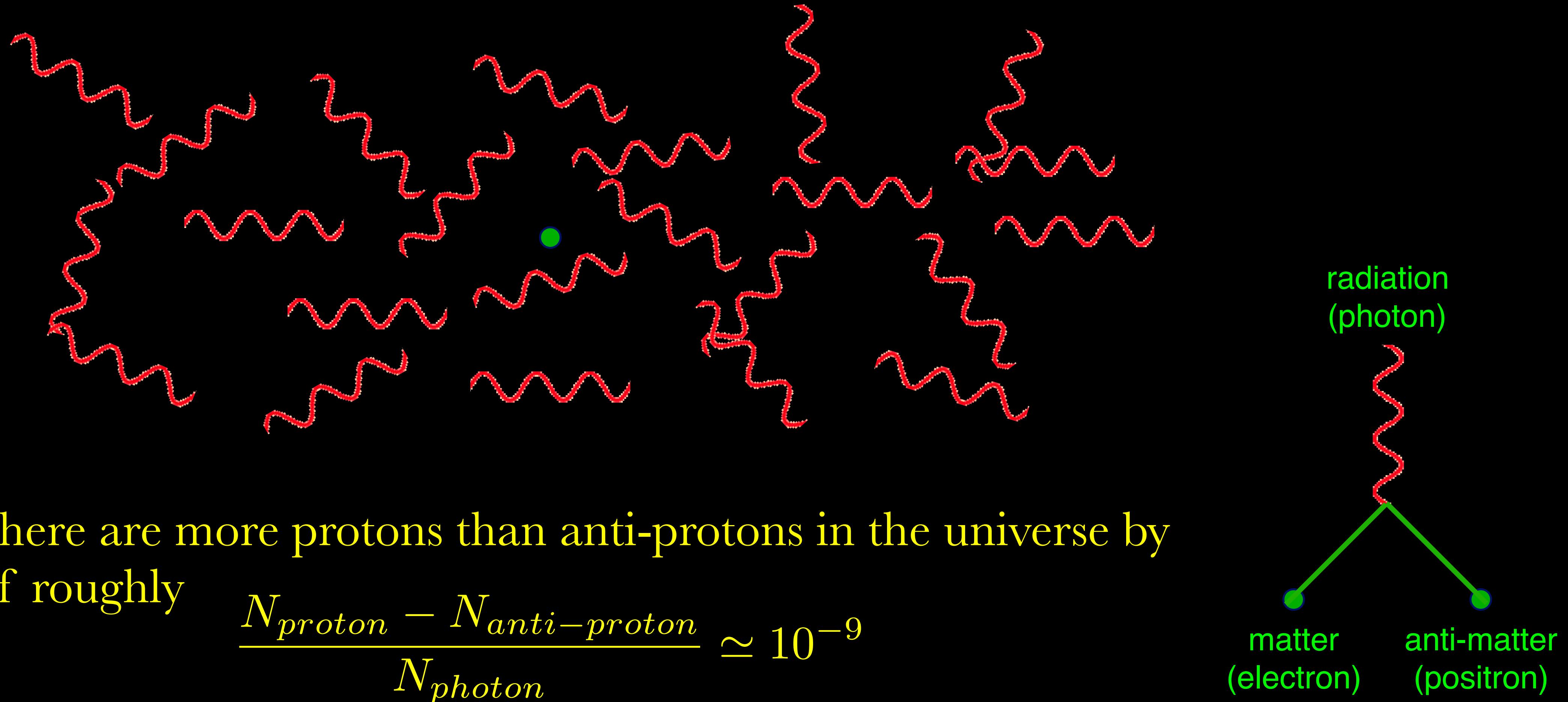
$$\Omega_{\text{mass+energy}} = 1 \text{ (or very close)}$$

cosmic mass/energy budget



$$\Omega_{\text{mass+energy}} = 1 \text{ (or very close)}$$

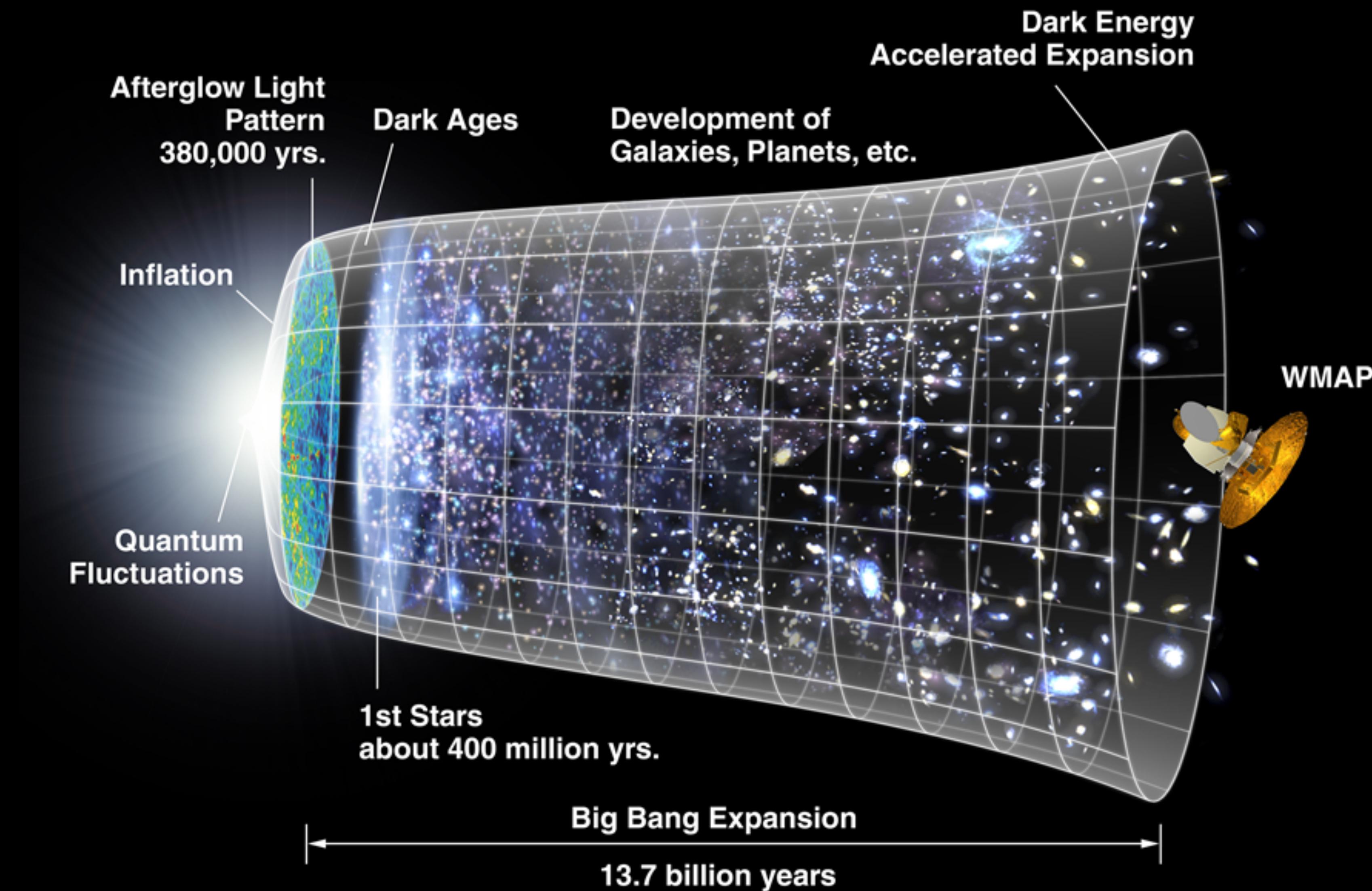
To the best of our knowledge, the Standard Model matter in the Universe is comprised entirely of matter and not anti-matter



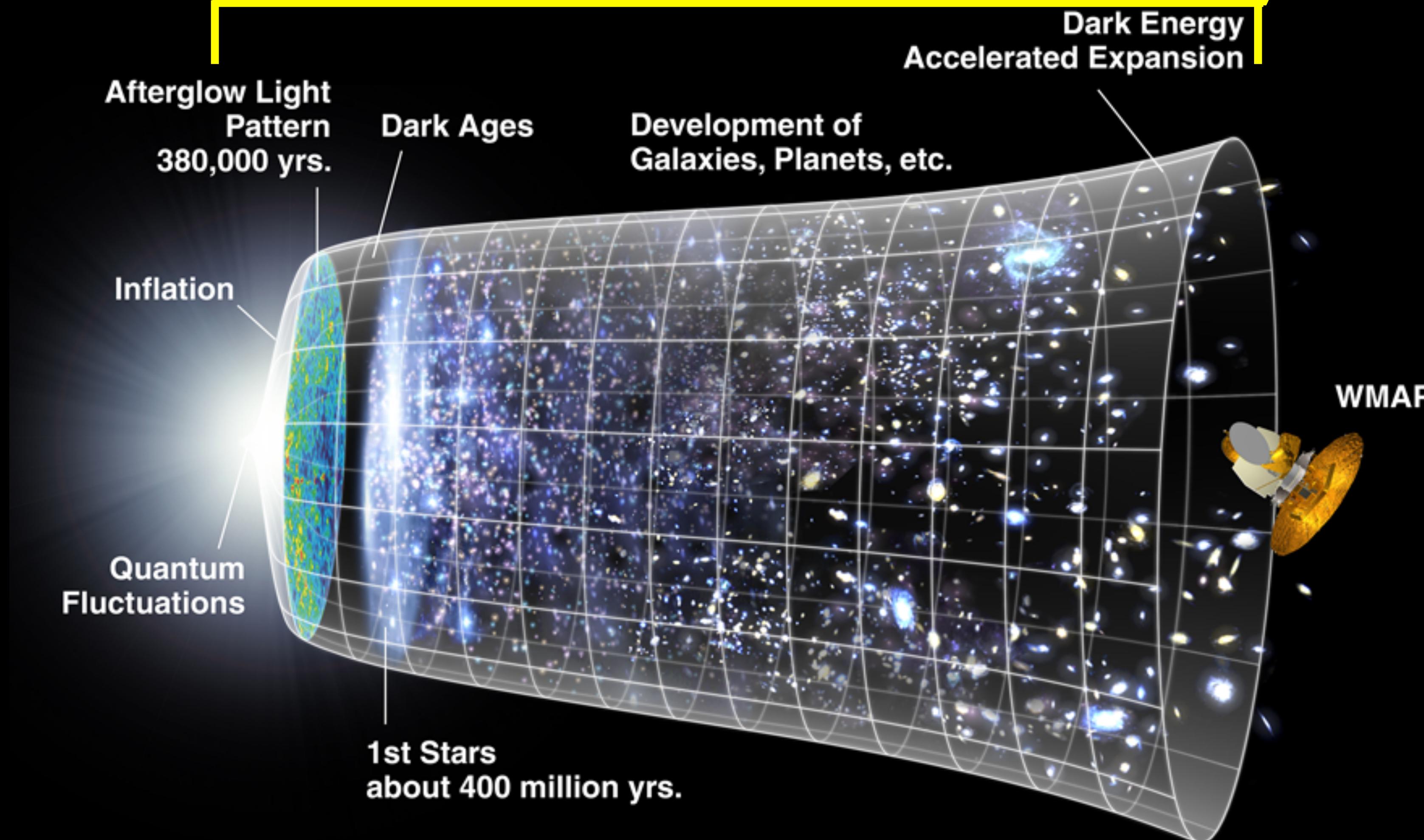
We observe there are more protons than anti-protons in the universe by an amount of roughly

$$\frac{N_{proton} - N_{anti-proton}}{N_{photon}} \simeq 10^{-9}$$

While tiny, this is still **10,000** times **or more** greater than we would predict with the Standard Model - why is there so much matter in the universe?

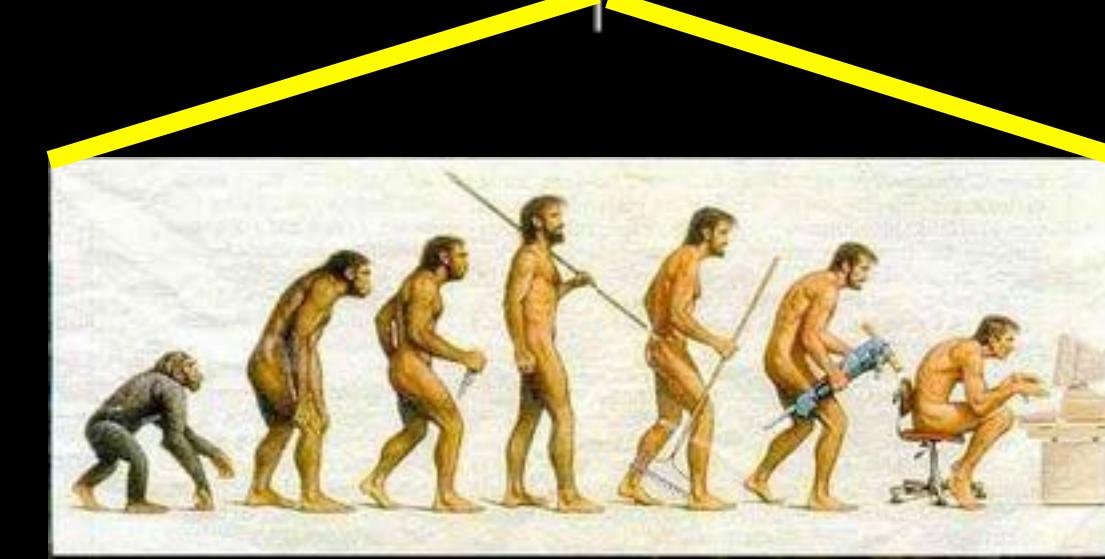


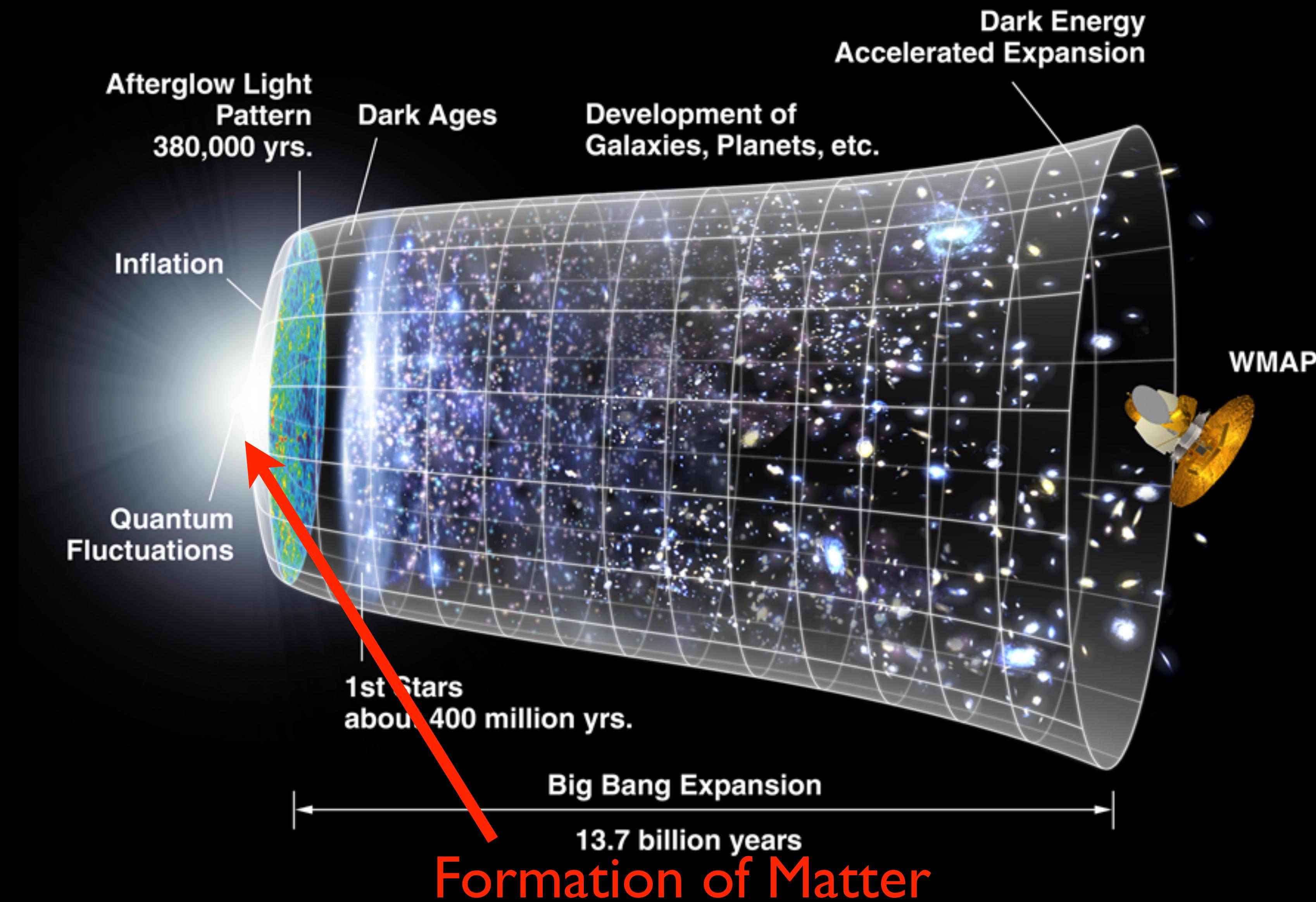
Standard Model of Particle Physics



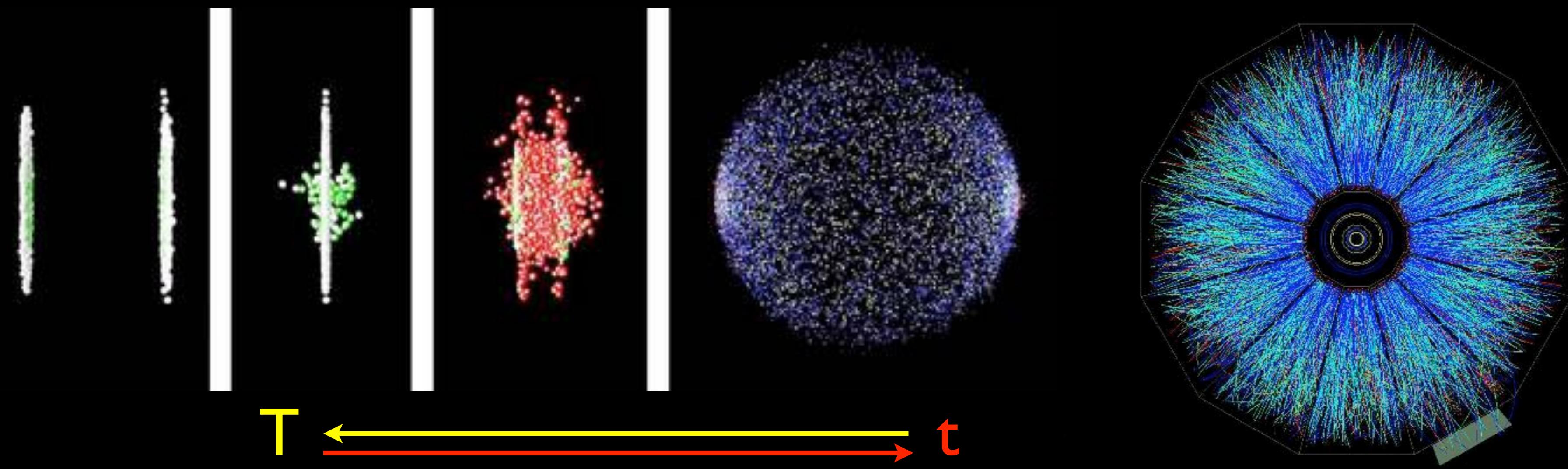
Big Bang Expansion

13.7 billion years

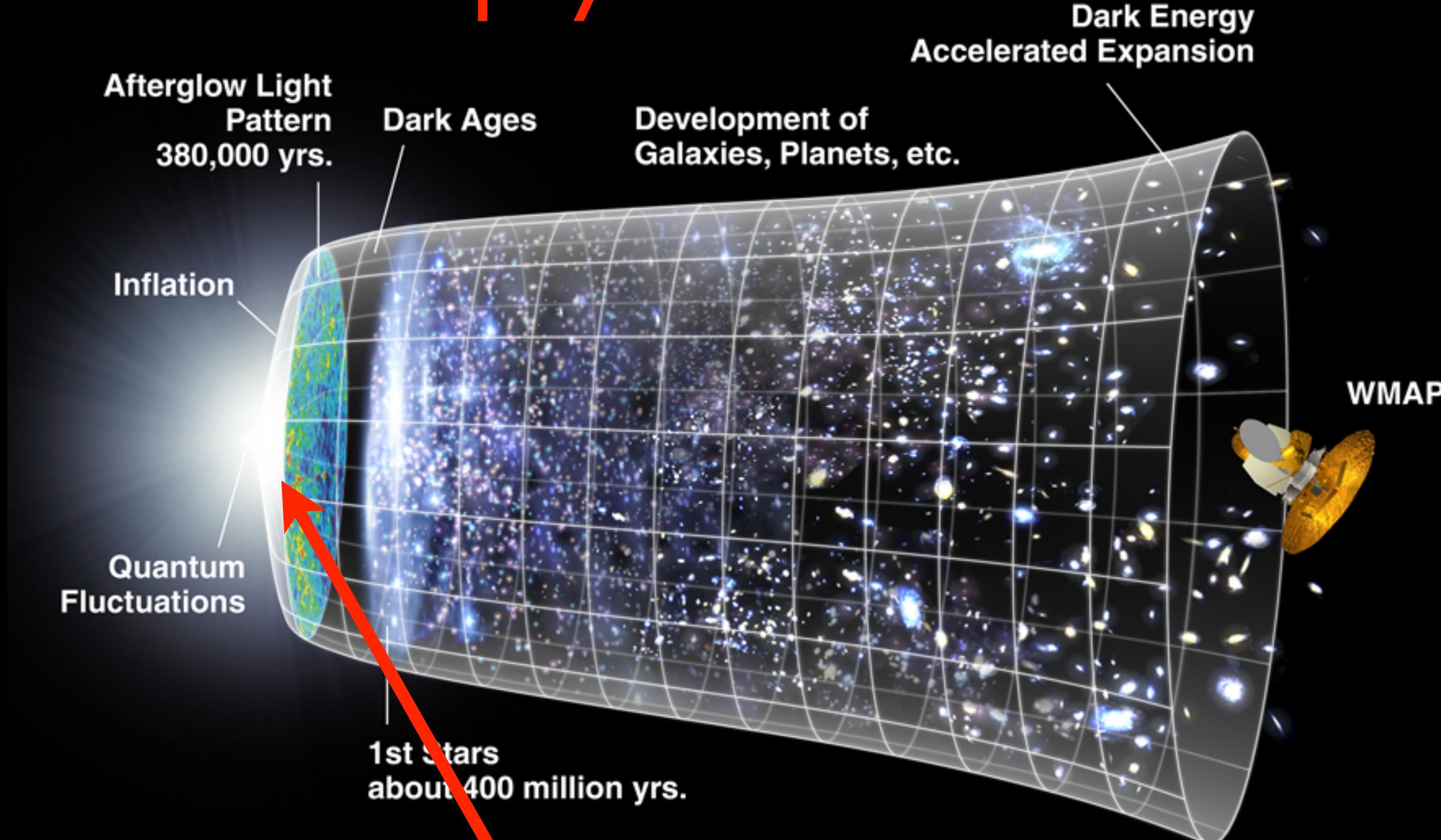




$T \approx 1 \text{ trillion K} (10^{12} \text{ K})$
 $t \approx 30 \text{ micro seconds} (3 \times 10^{-5} \text{ s})$

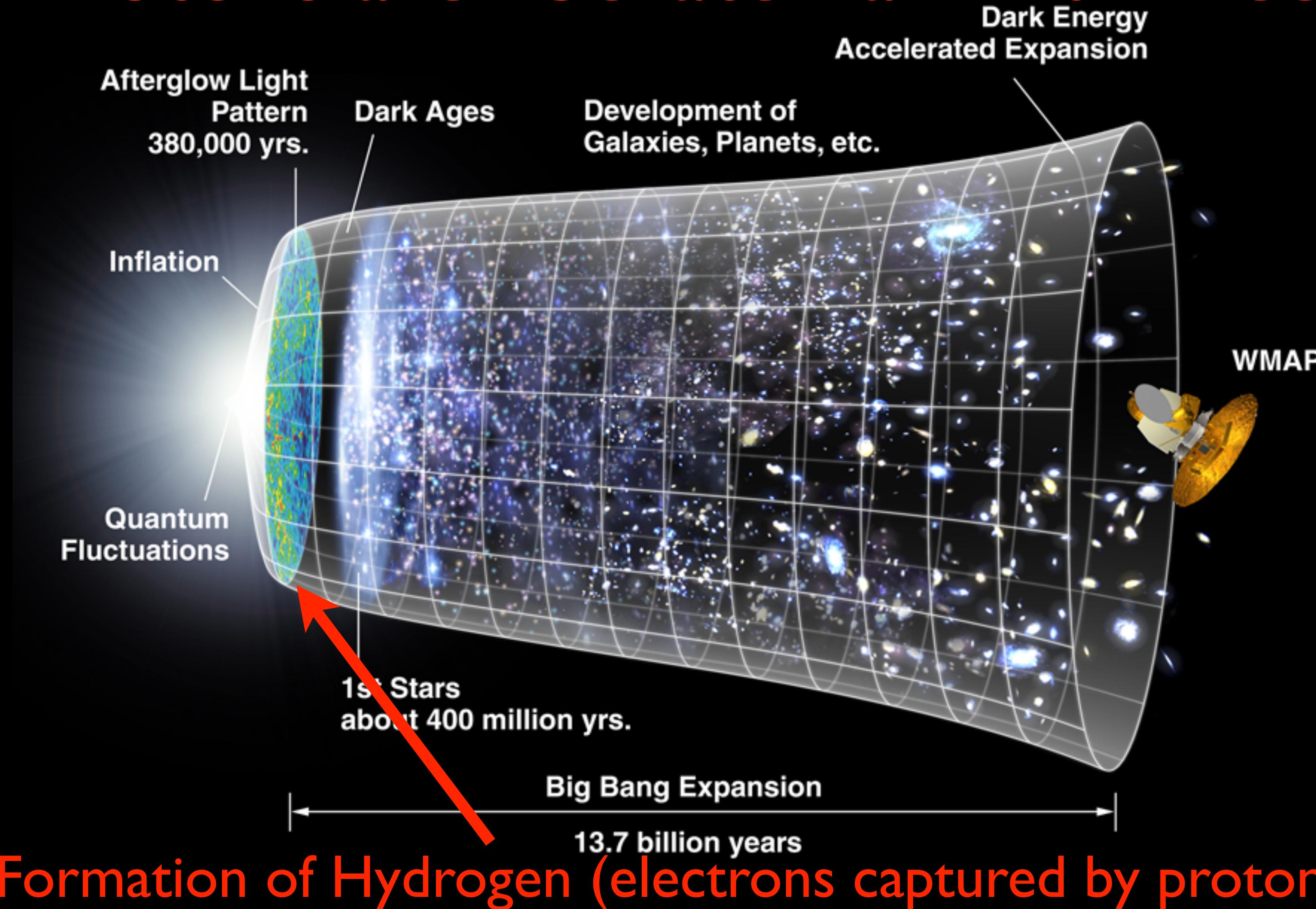


Nuclear physics to the rescue!



$T \approx 1 \text{ billion K} (10^9 \text{ K})$
 $t \approx 3 \text{ minutes}$

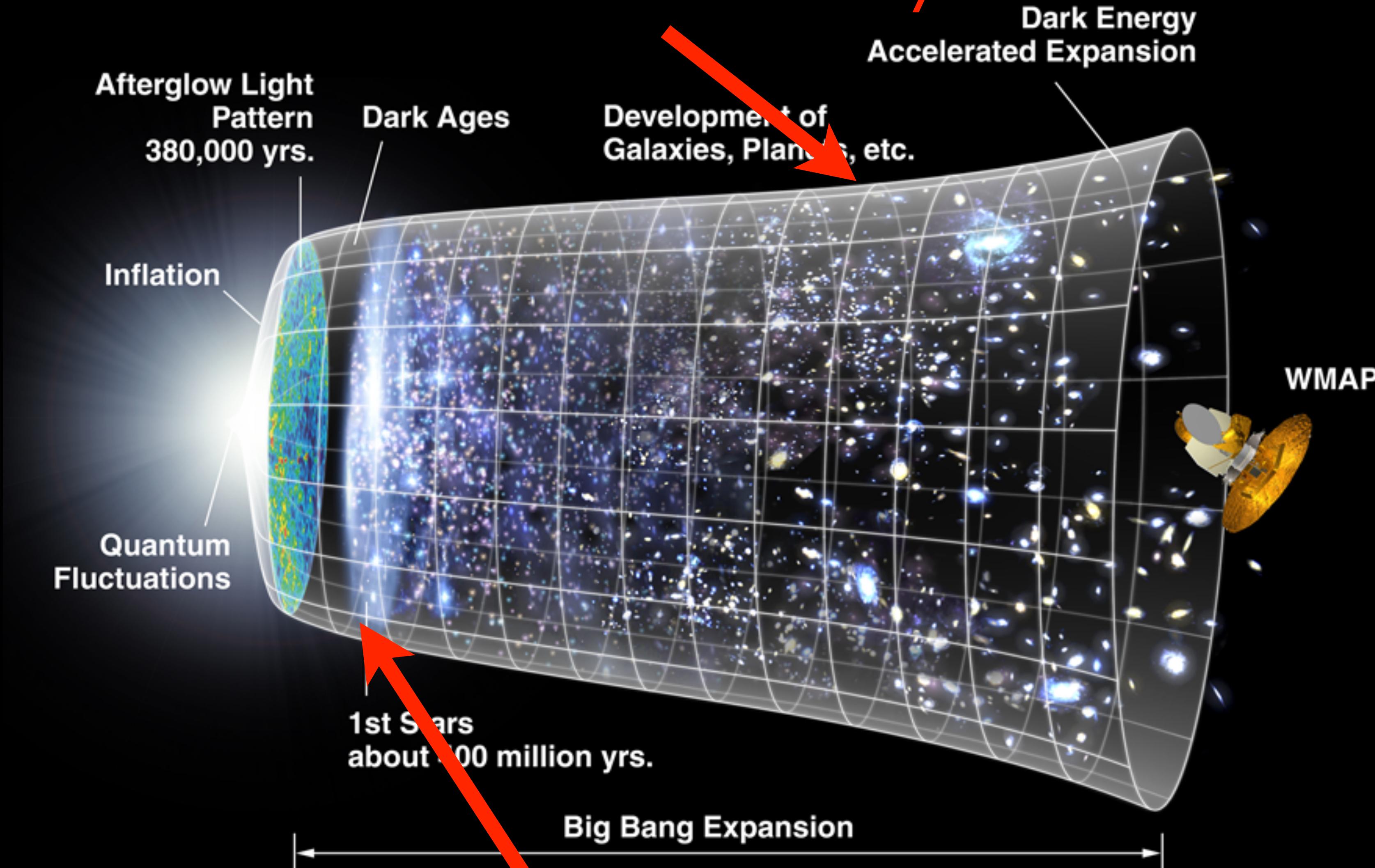
Photons are liberated and run free!



Formation of Hydrogen (electrons captured by protons)

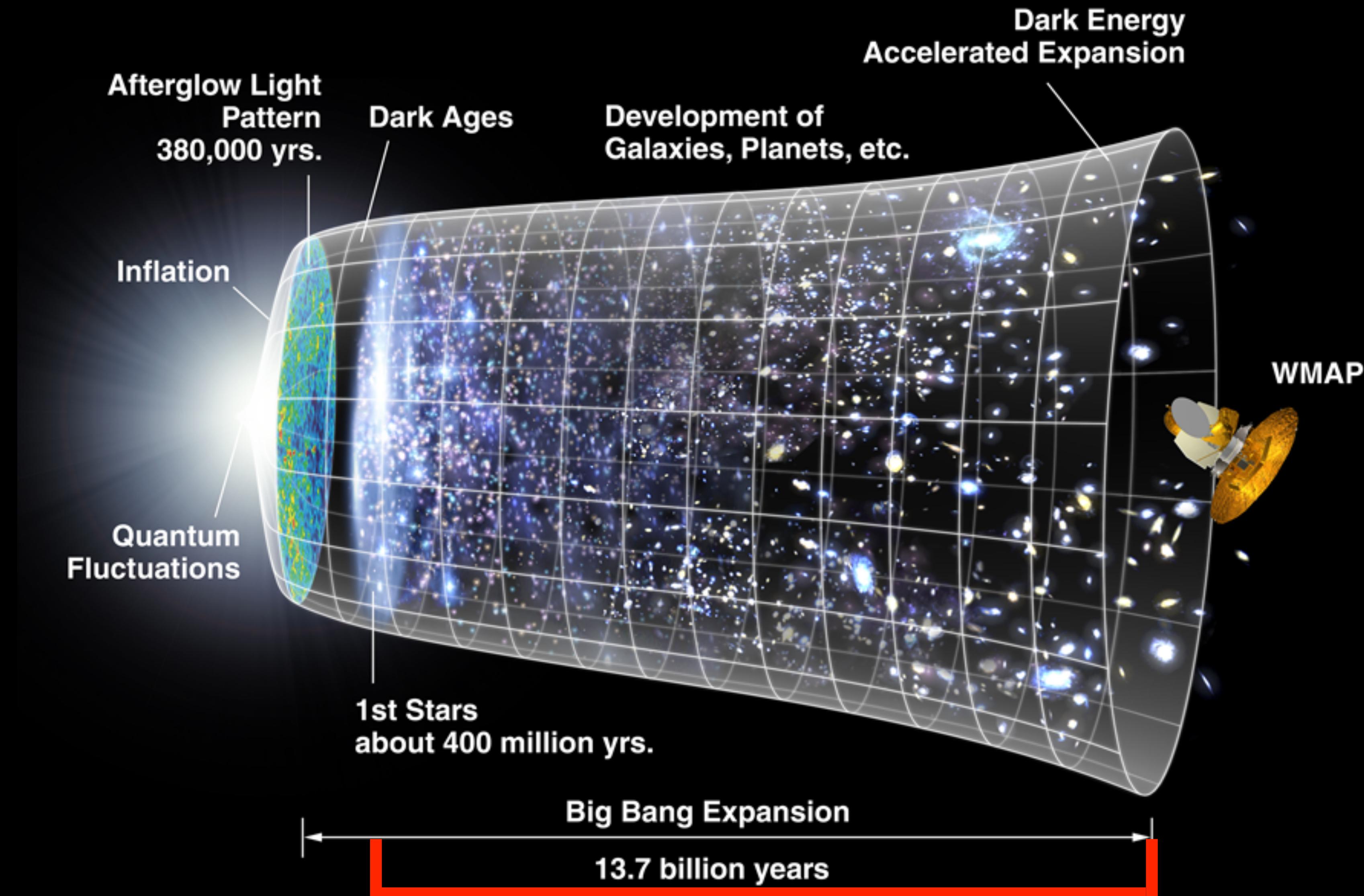
$$\begin{aligned} T &\simeq 4,000 \text{ K} \\ t &\simeq 380,000 \text{ years} \end{aligned}$$

Formation of our solar system



Formation of first stars

$T \approx 20 \text{ K}$
 $t \approx 200 \text{ Million years}$



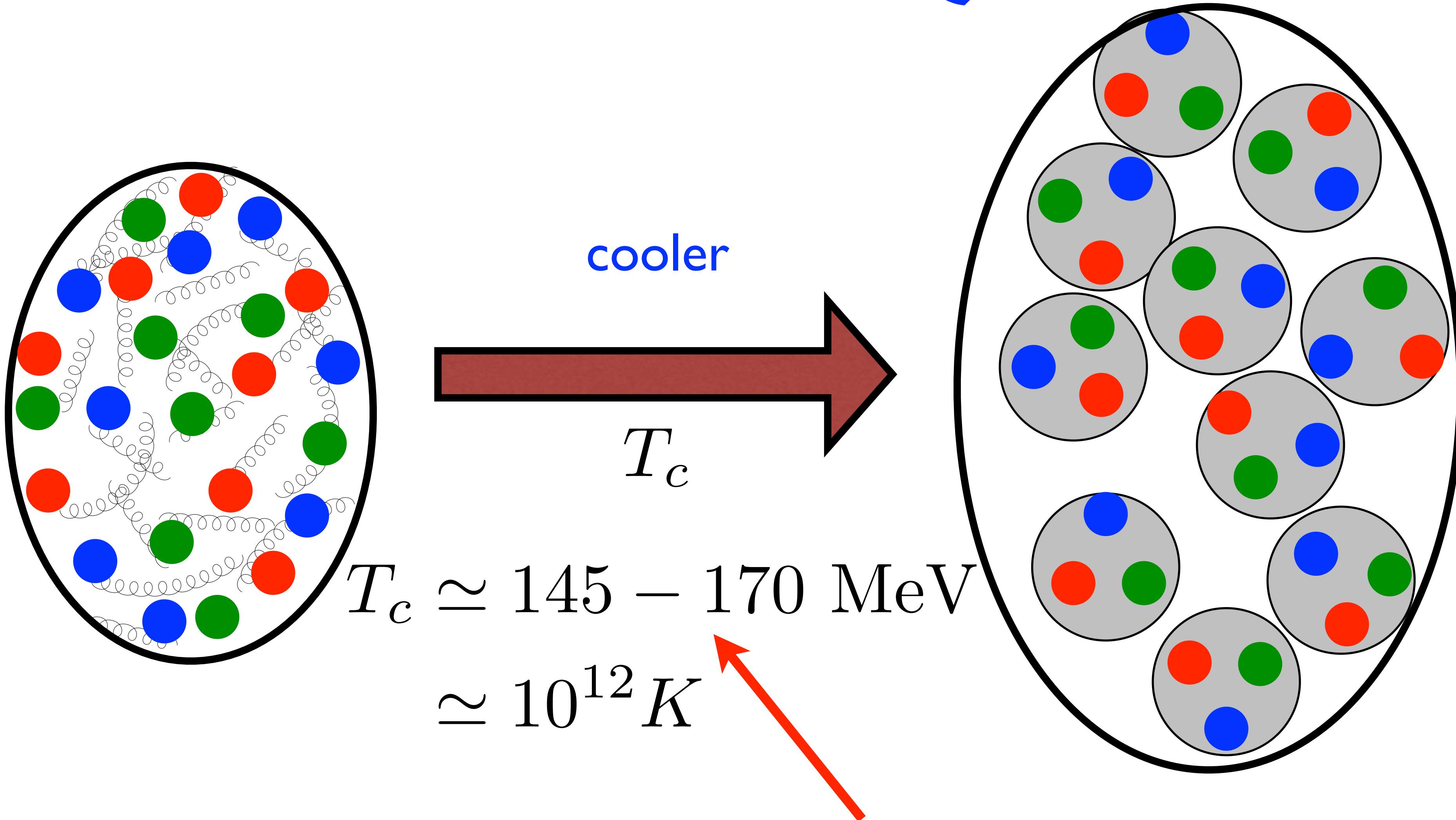
Death of stars, creation of heavy nuclei and *life*,
creation of new, *ultradense* states of nuclear matter

From Quarks to Protons and Neutrons

$T \simeq 1$ trillion K (10^{12} K)

$t \simeq 30$ micro seconds ($3.0 \times 10^{-5}s$)

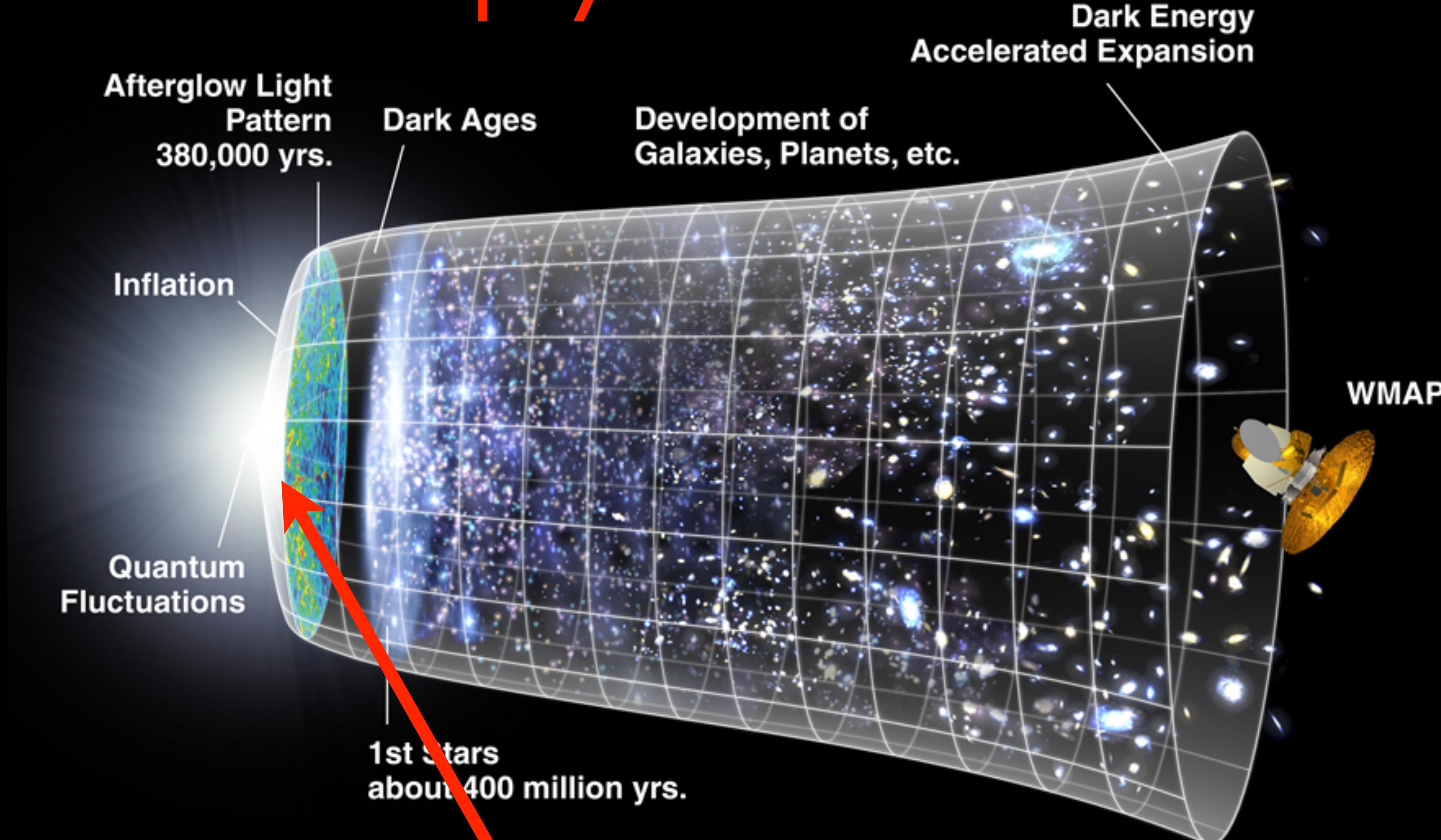
Confinement of Quarks



computed by Budapest-Wuppertal and hot-QCD Lattice
Collaborations with previous generation supercomputers

See lectures by
Peter Petreczky

Nuclear physics to the rescue!



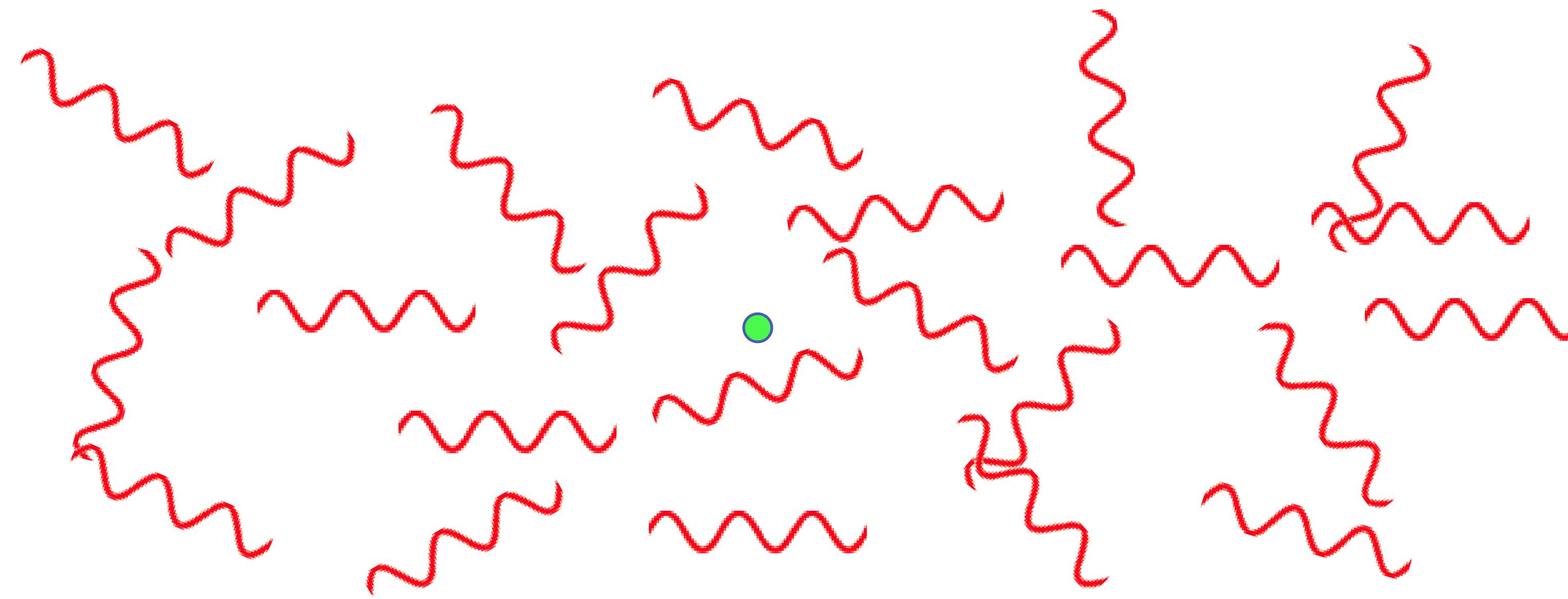
$T \approx 1 \text{ billion K} (10^9 \text{ K})$
 $t \approx 3 \text{ minutes}$

Big Bang Nucleosynthesis

$T \simeq 1 \text{ trillion K} \rightarrow 1 \text{ billion K}$

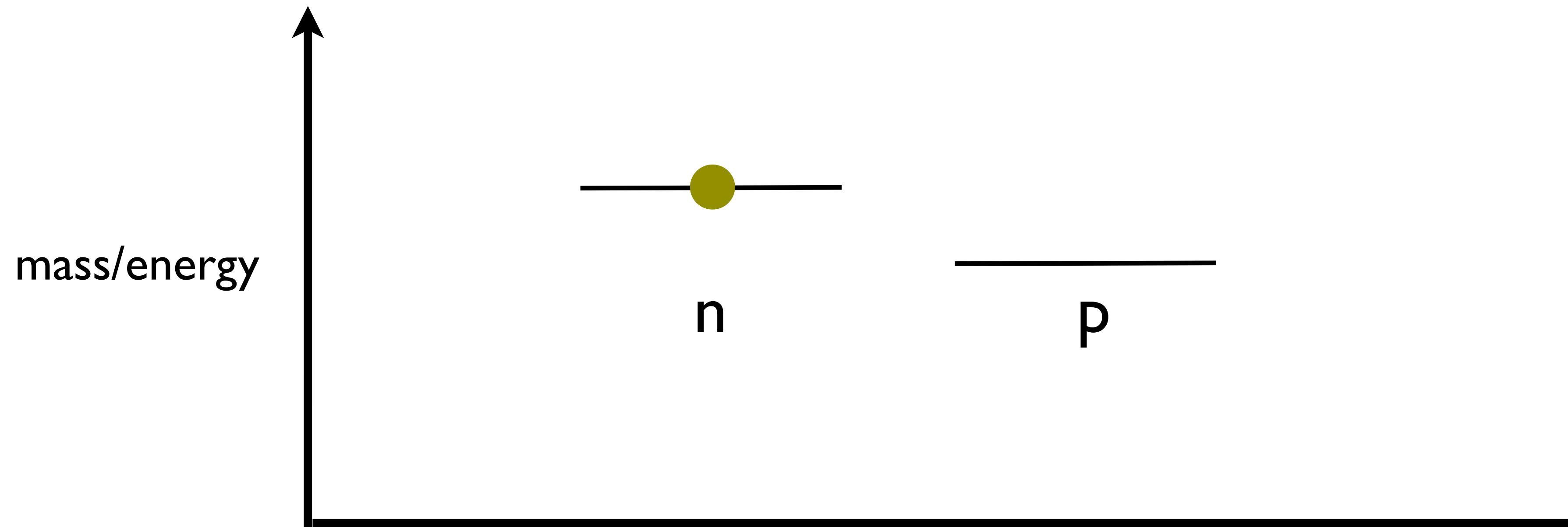
$t \simeq 3 \times 10^{-5} s \rightarrow 3 \text{ min}$

Our initial condition is a soup of radiation plus a small excess amount of matter, in the form of protons, neutrons, electrons and photons



$$\eta_{B/\gamma} \sim 10^{-9}$$

when systems cool, they settle into the lowest energy state

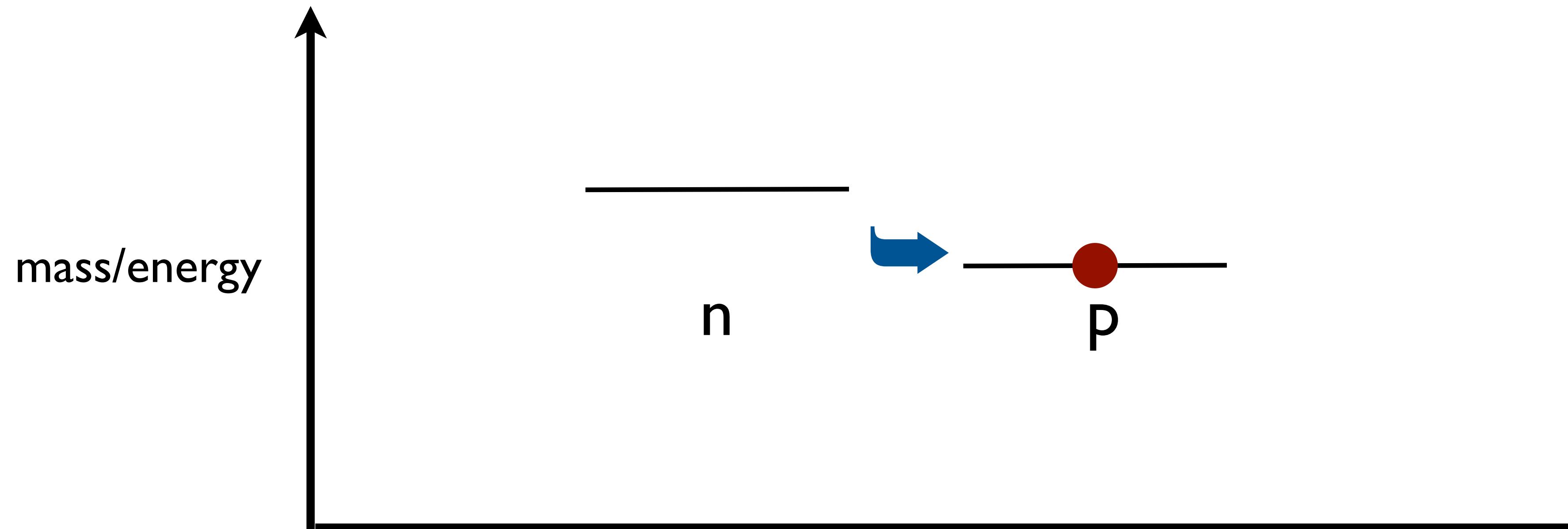


$$M_n - M_p = 1.29333217(42) \text{ MeV}$$

$$E_{\text{Hydrogen}} = 13.6 \text{ eV}$$

$$\frac{M_n + M_p}{2} = 938.9187473(58) \text{ MeV}$$

when systems cool, they settle into the lowest energy state

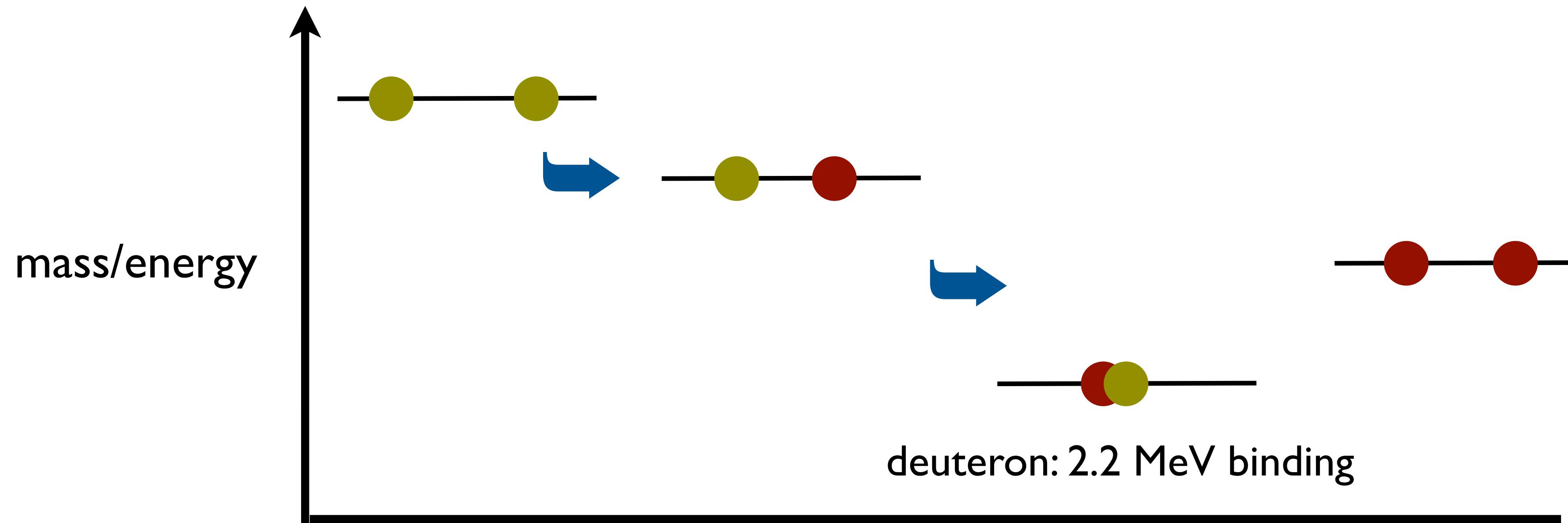


$$\tau_n \sim 15 \text{ min}$$

what prevented this from destroying all the neutrons?

if nothing else were to happen in the next few minutes,
our universe would be full of only Hydrogen

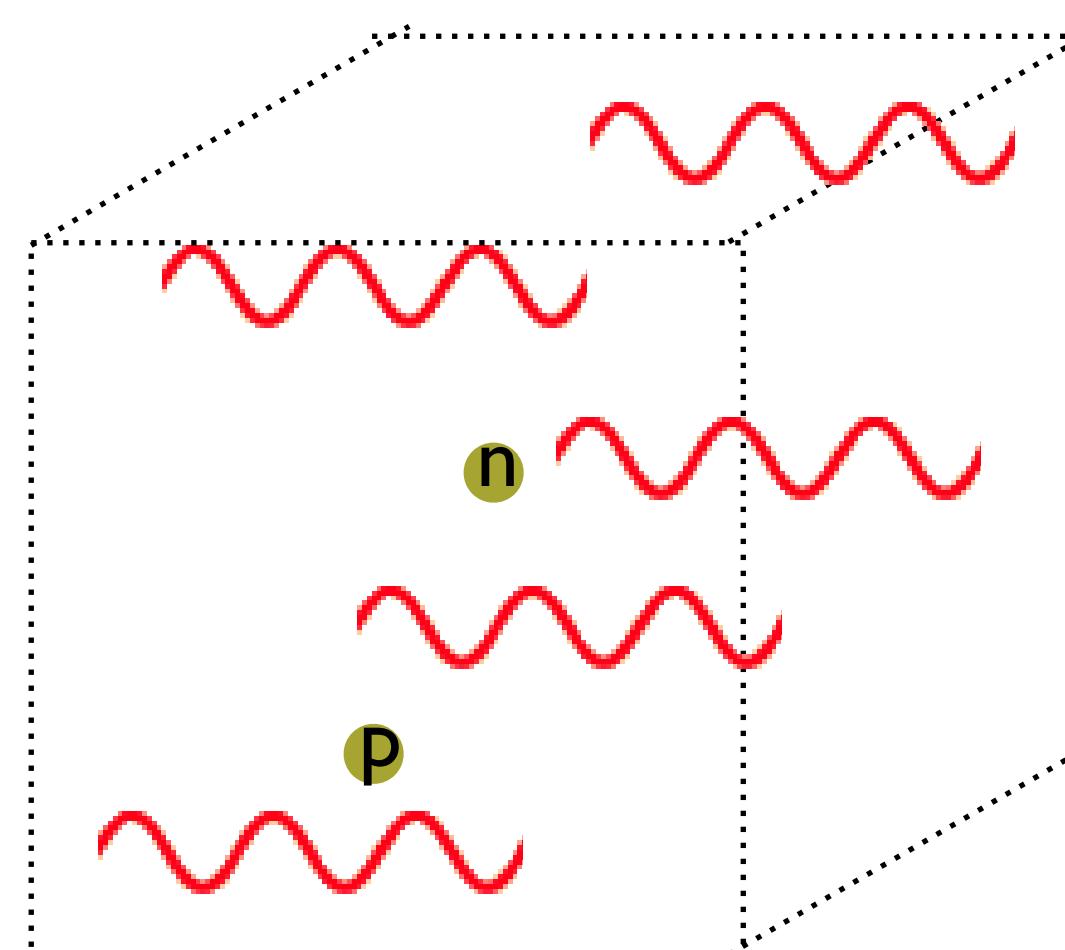
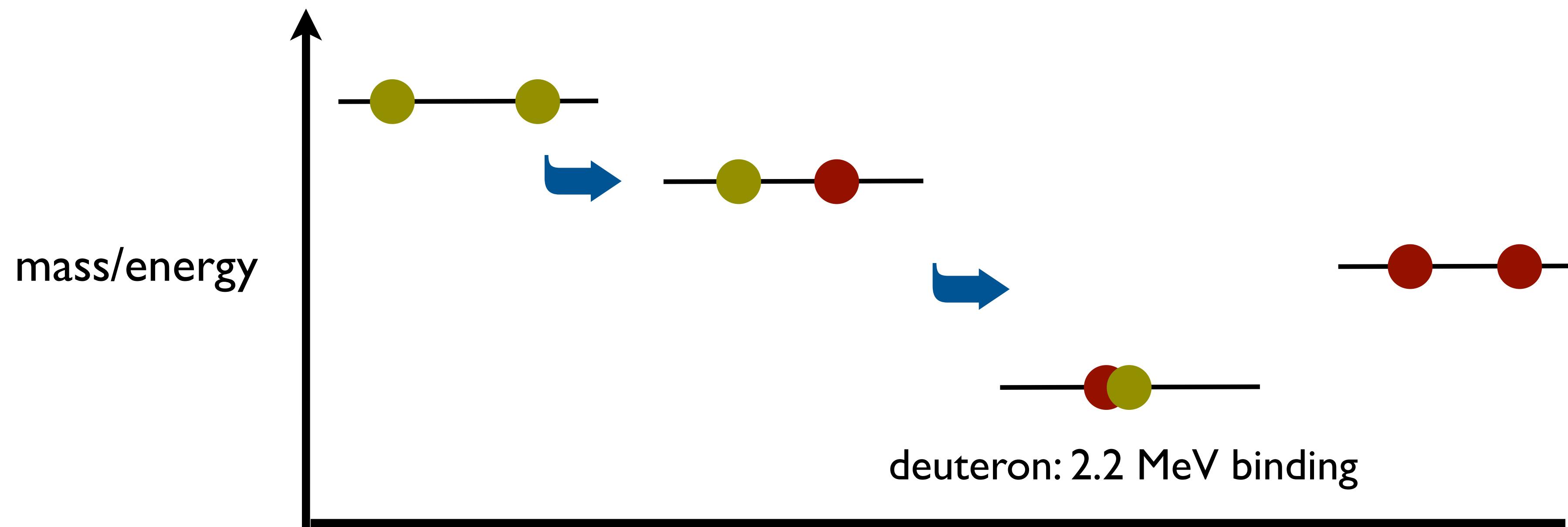
when systems cool, they settle into the lowest energy state



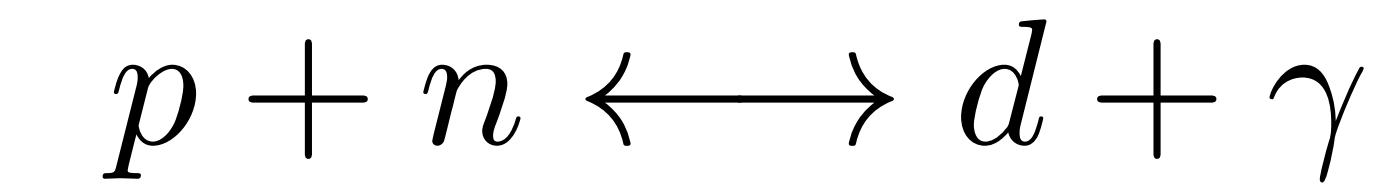
Answer: formation of nuclei

a system with protons and neutrons can collapse to a compact bound state, the **deuteron**: the attractive binding of a neutron and proton allows **neutrons to survive when embedded in nuclei**

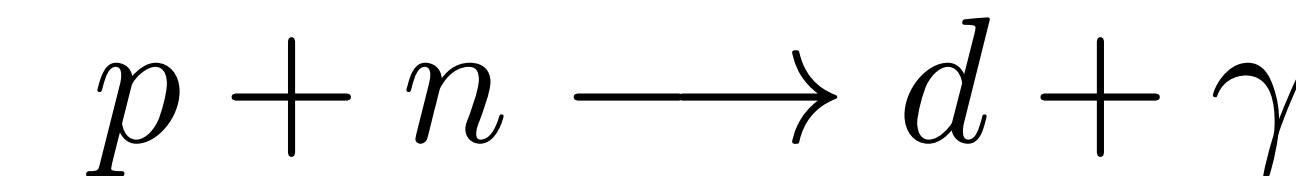
The deuterium “bottleneck”



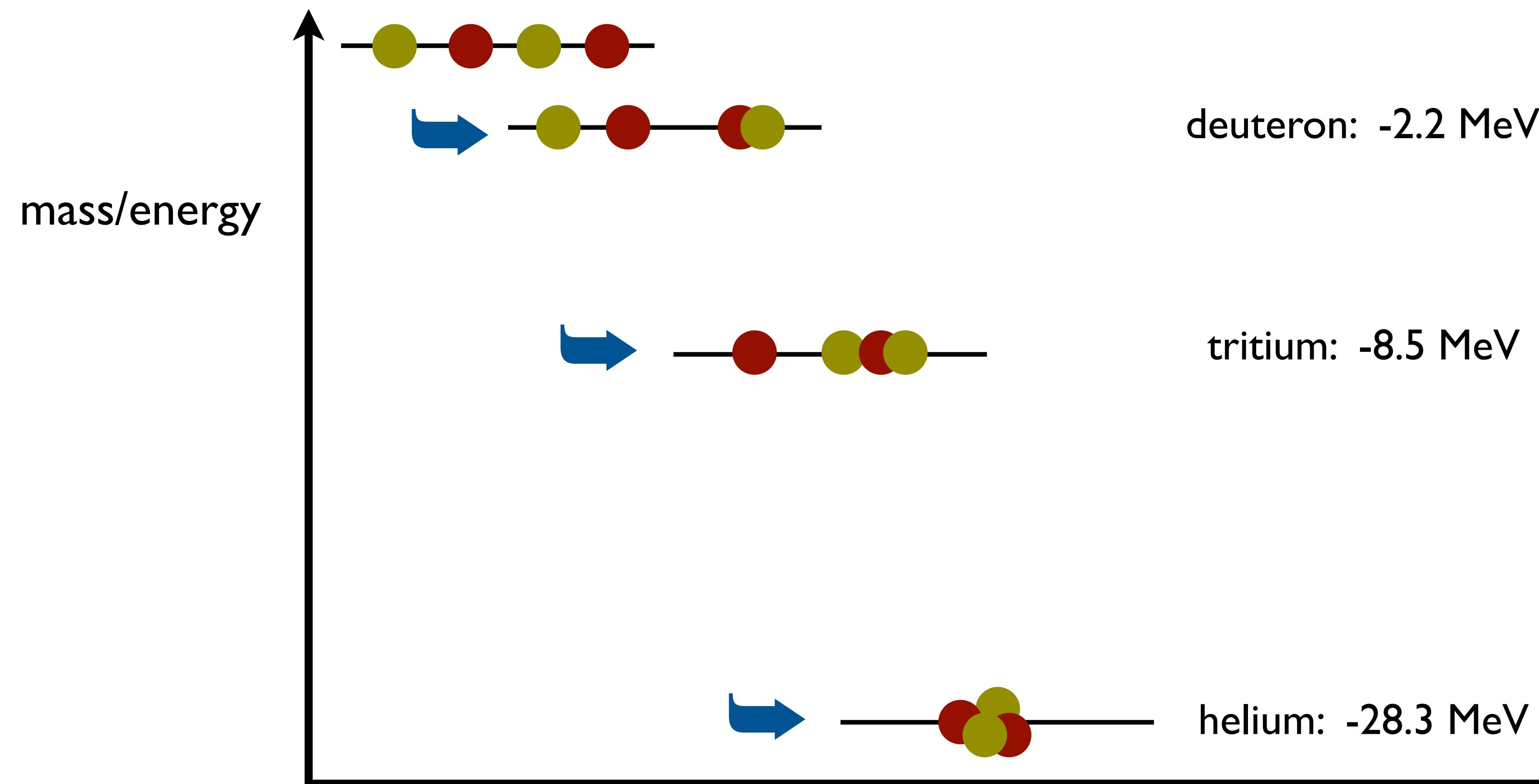
$$\eta_{B/\gamma} \sim 10^{-9}$$



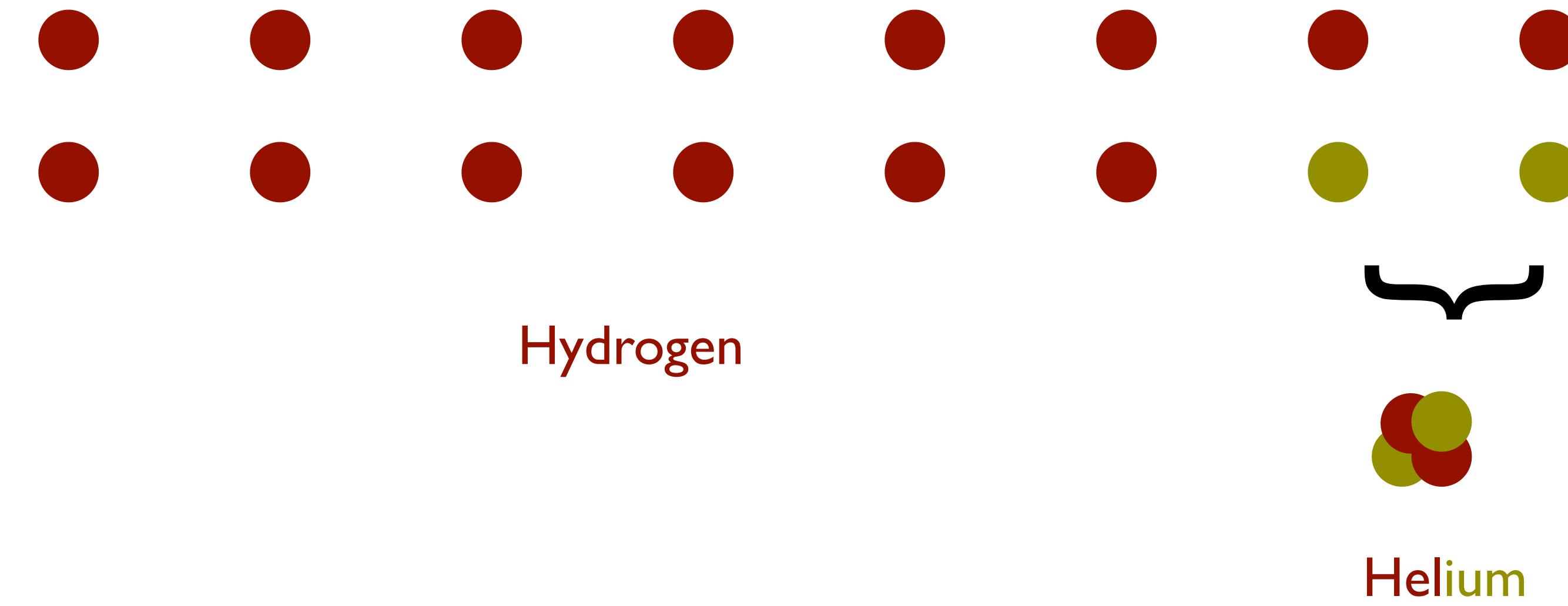
until $T \approx 100 \text{ keV}$ (1 billion K),
 $t \approx 3 \text{ min}$



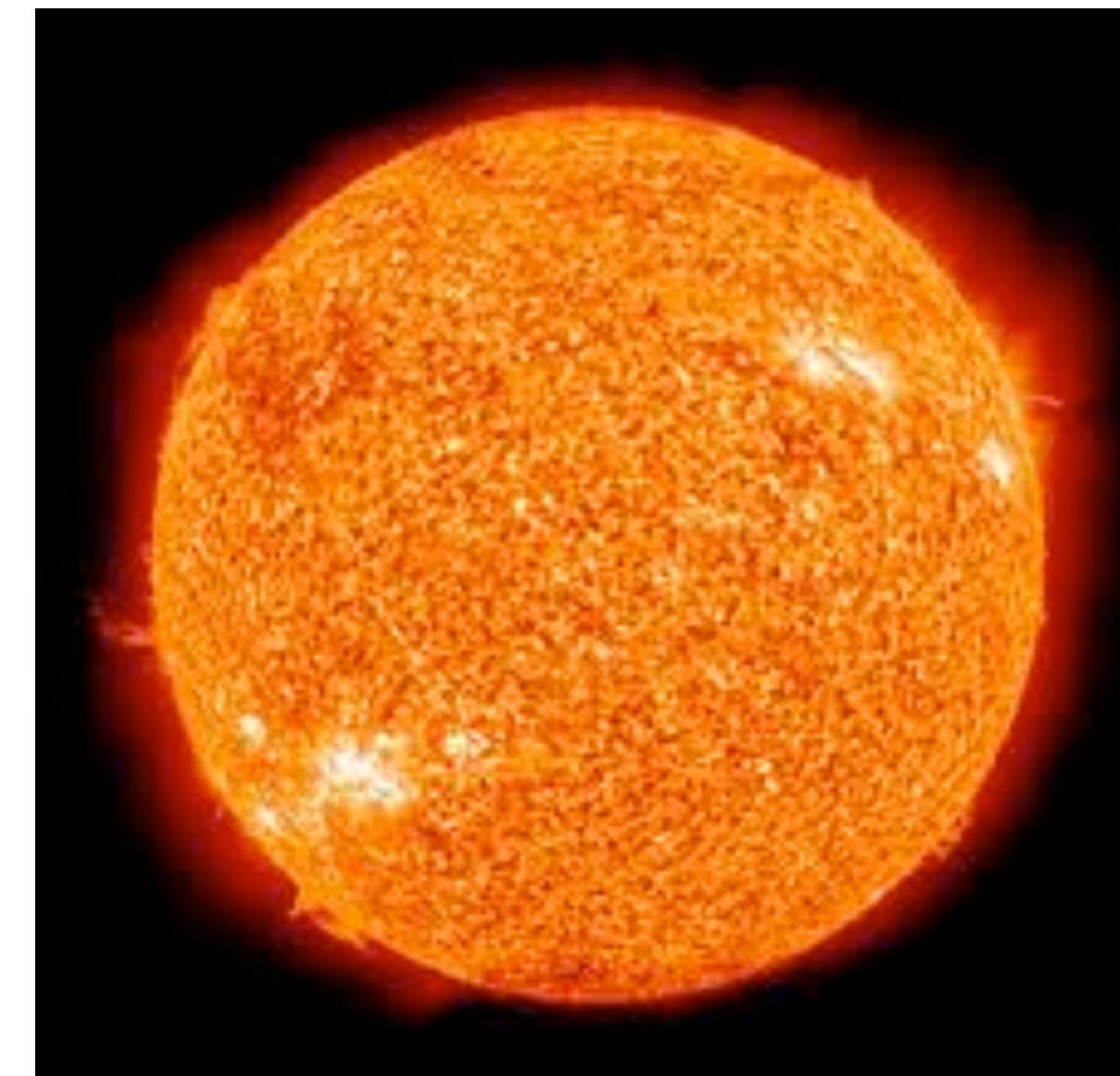
The deuterium “bottleneck” is broken, neutrons flow into He



He stability: \uparrow, \downarrow protons and \uparrow, \downarrow neutrons can be packed together



The early universe contains 75% H and 25% ${}^4\text{He}$ by mass fraction
("all" deuterium converted to ${}^4\text{He}$)



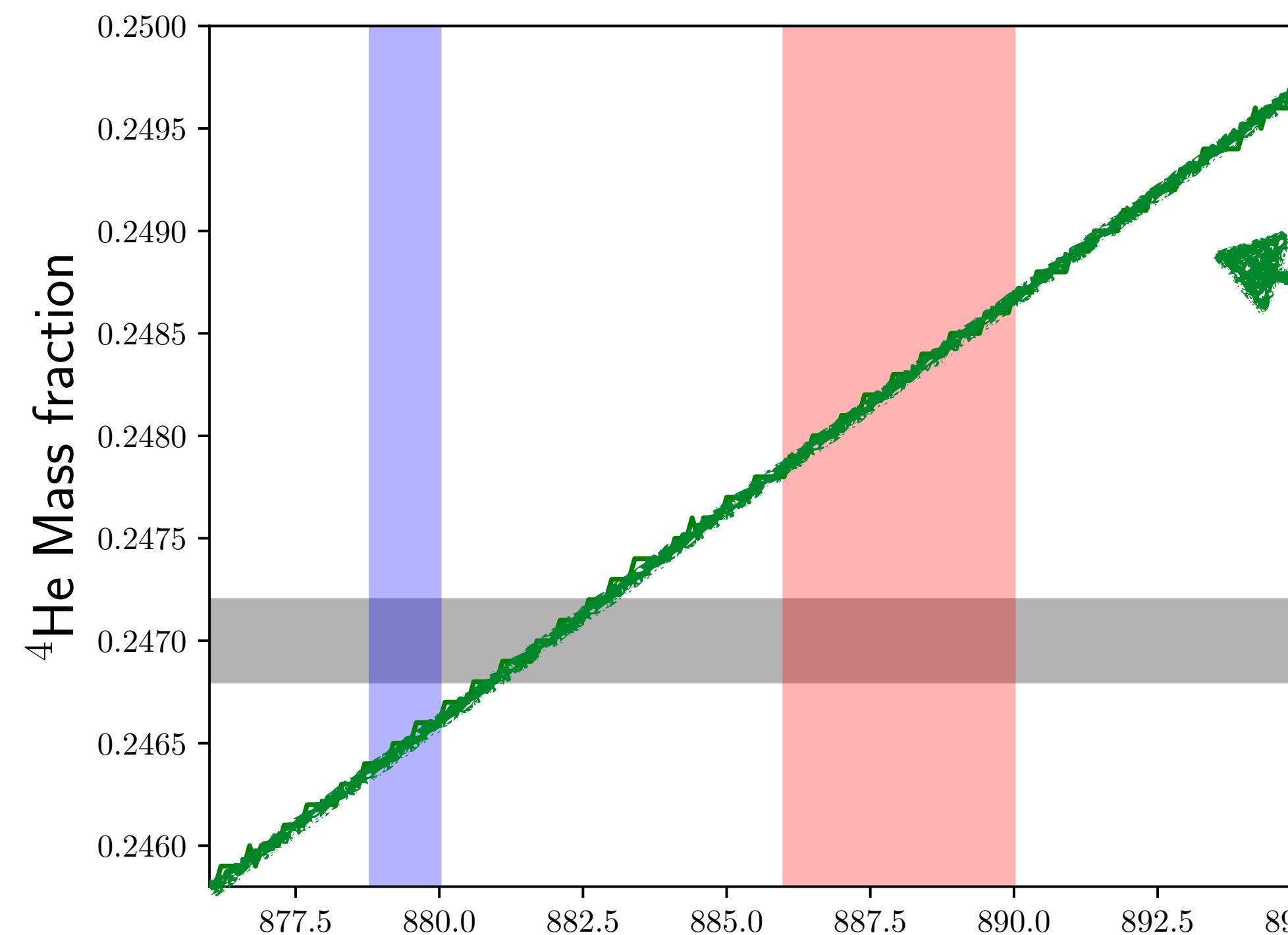
Our understanding of Big Bang Nucleosynthesis is very sensitive to a few ingredients

- The abundance of matter over antimatter $\eta_{B/\gamma} \sim 10^{-9}$
A. Nicholson
- The neutron-proton mass splitting $M_n - M_p = 1.29333217(42) \text{ MeV}$
A. Portelli, G. Martinelli
- The neutron lifetime $\frac{1}{\tau_n} = \frac{G_\mu^2 |V_{ud}|^2}{2\pi^3} m_e^5 (1 + 3g_A^2)(1 + RC) f_{V,A}$
- The deuteron binding energy $B_d \simeq 2.2 \text{ MeV}$
S. Aoki

The matter/antimatter asymmetry is an input ingredient while the other three quantities, to some degree, emerge from the Standard Model.

How sensitive are these quantities to variations of the input parameters of the SM?
With Lattice QCD and Effective Field Theory, we can answer these questions

This picture is also very sensitive to the lifetime of the neutron

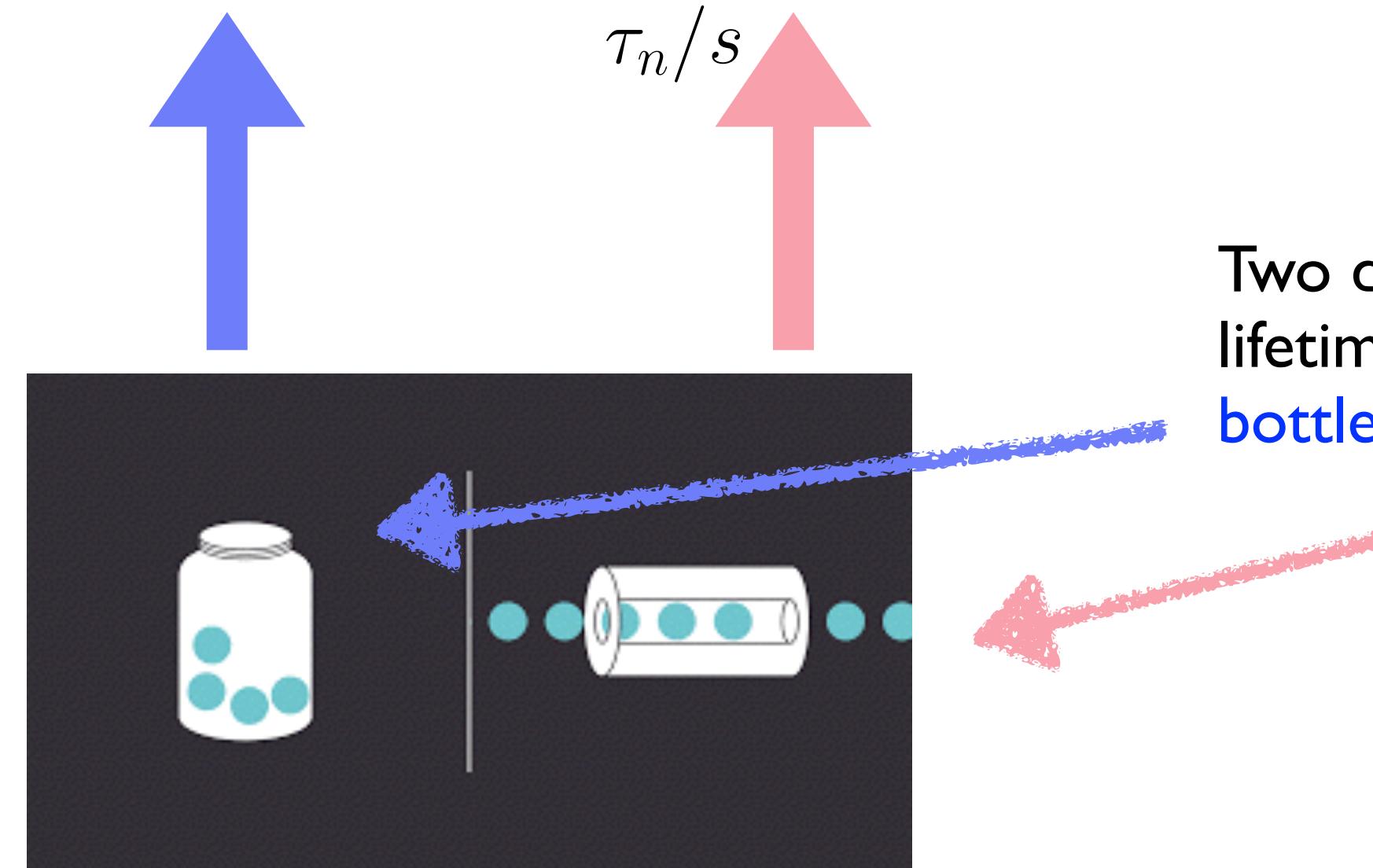


$$\tau_n = \frac{(5172.0 \pm 1.1) \text{ seconds}}{1 + 3g_A^2}$$

Czarnecki, Marciano, Sirlin

Predicted primordial ${}^4\text{He}$ mass fraction of the universe as a function of the neutron lifetime

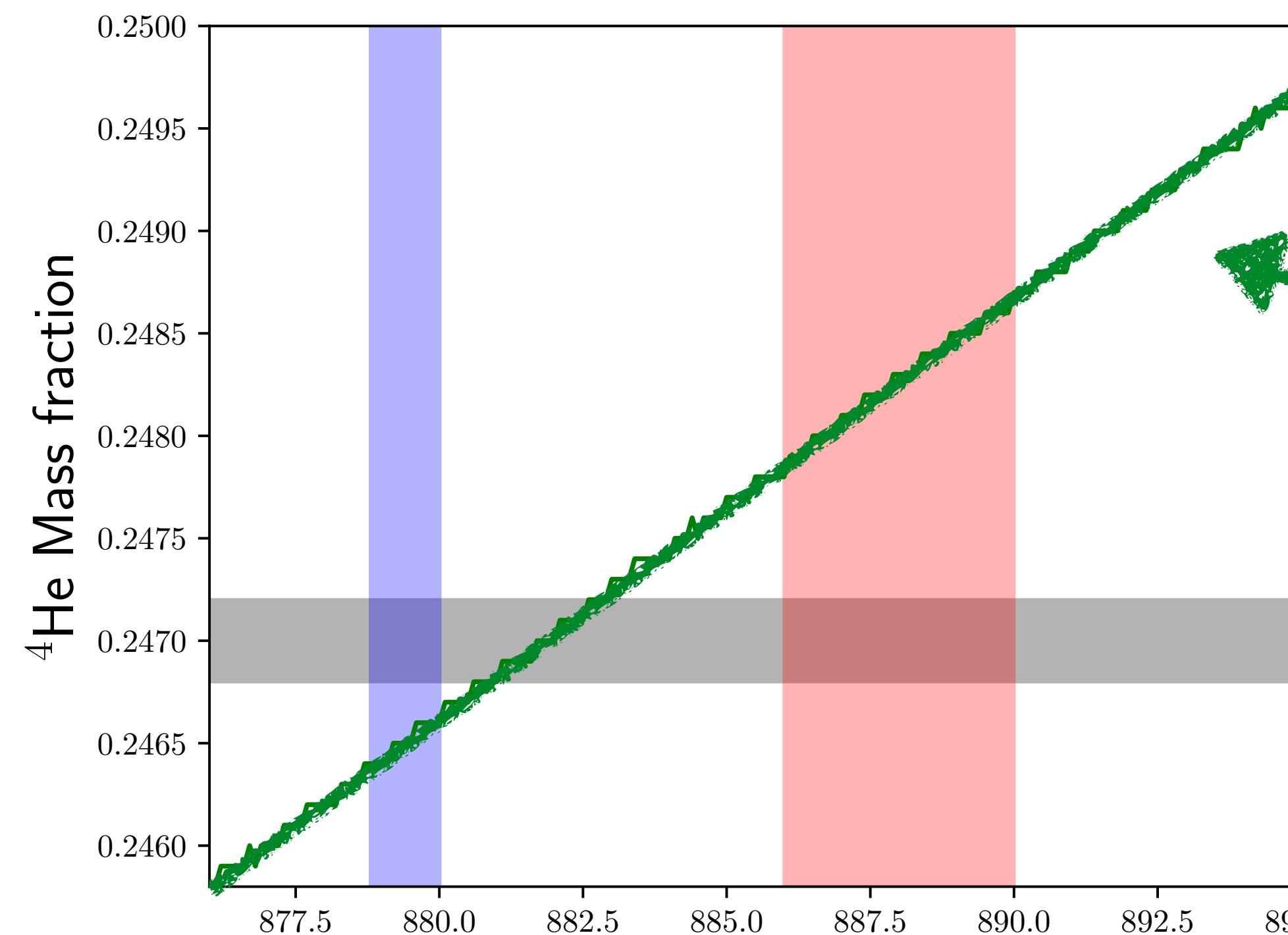
Observed primordial ${}^4\text{He}$ mass fraction of the universe



Two different methods of measuring the neutron lifetime disagree at the 99% level
bottle and **beam**

Is one of the experiments wrong? Or is there new physics hiding here?

This picture is also very sensitive to the lifetime of the neutron

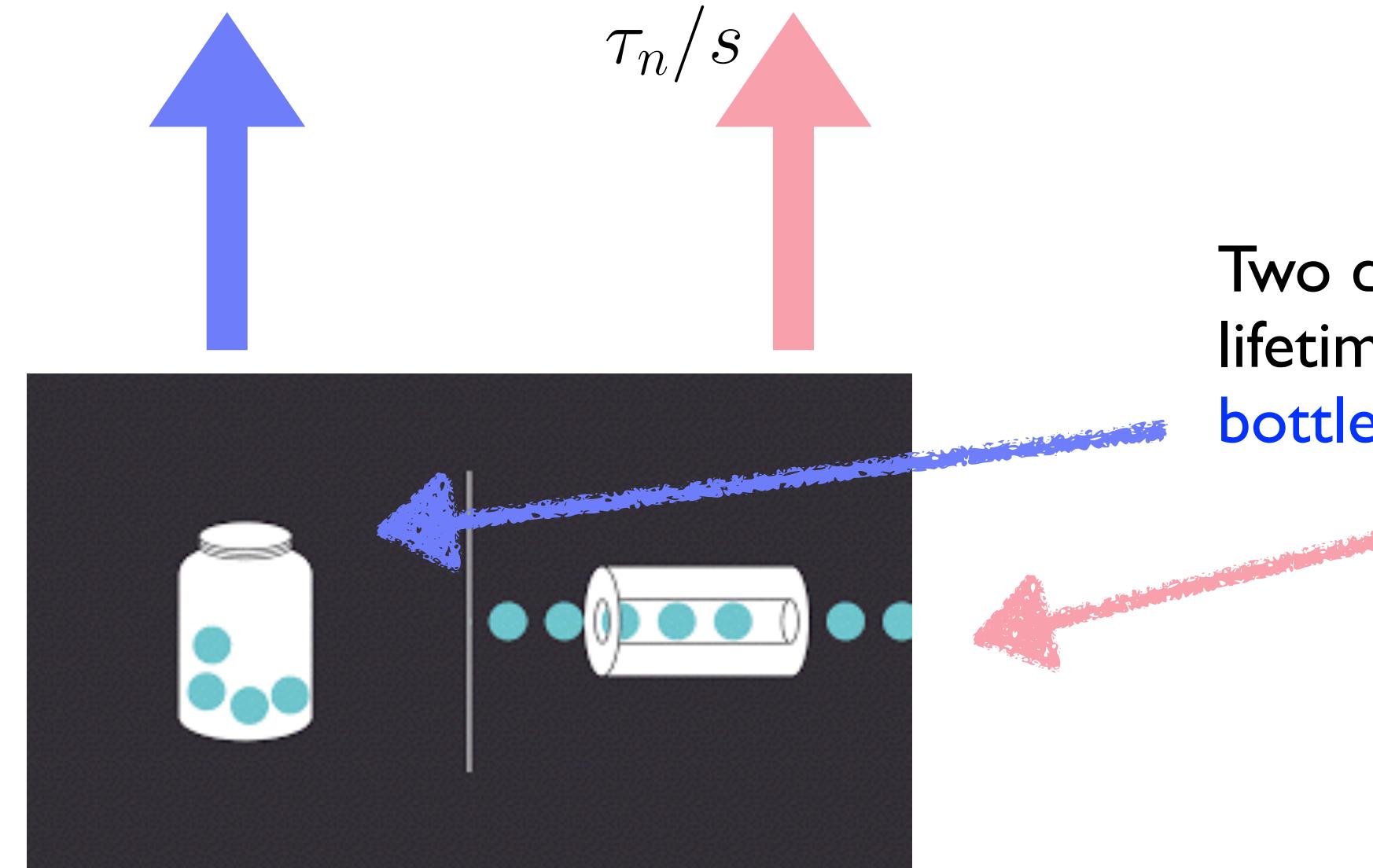


$$\tau_n = \frac{(5172.0 \pm 1.1) \text{ seconds}}{1 + 3g_A^2}$$

Czarnecki, Marciano, Sirlin

with LQCD, we can provide this number
need 0.2% uncertainty to match
precision of discrepancy 😱

Observed primordial ${}^4\text{He}$
mass fraction of the
universe



Two different methods of measuring the neutron
lifetime disagree at the 99% level
bottle and **beam**

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wrong? Or is there new
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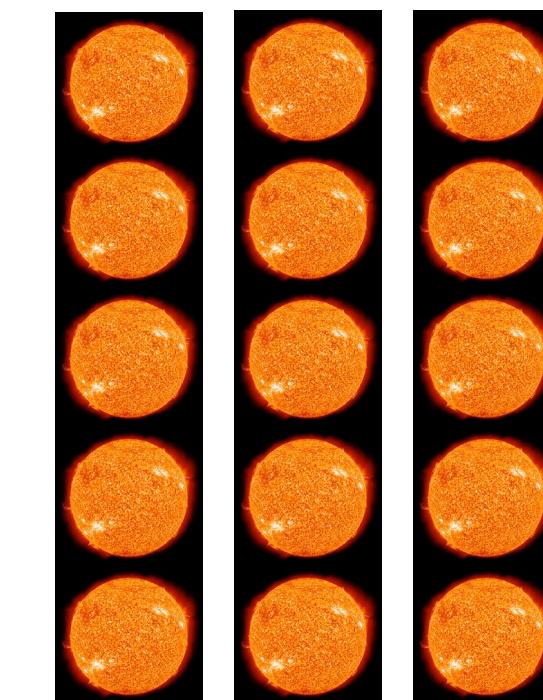
the binding energy of deuterium which is finely tuned (most nuclei have \sim 8 MeV binding per nucleon)!

$$B_d = 2.22 \text{ MeV}$$

What if

$$B_d \ll 2.22 \text{ MeV}$$

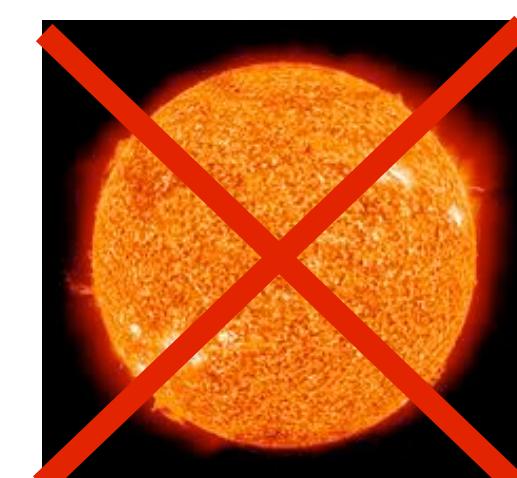
more finely tuned
all neutrons decay - no helium
mostly hydrogen stars?



Too many suns?

$$B_d \gg 2.22 \text{ MeV}$$

natural scenario
all neutrons captured in deuterium and helium - no hydrogen
no stars like ours!

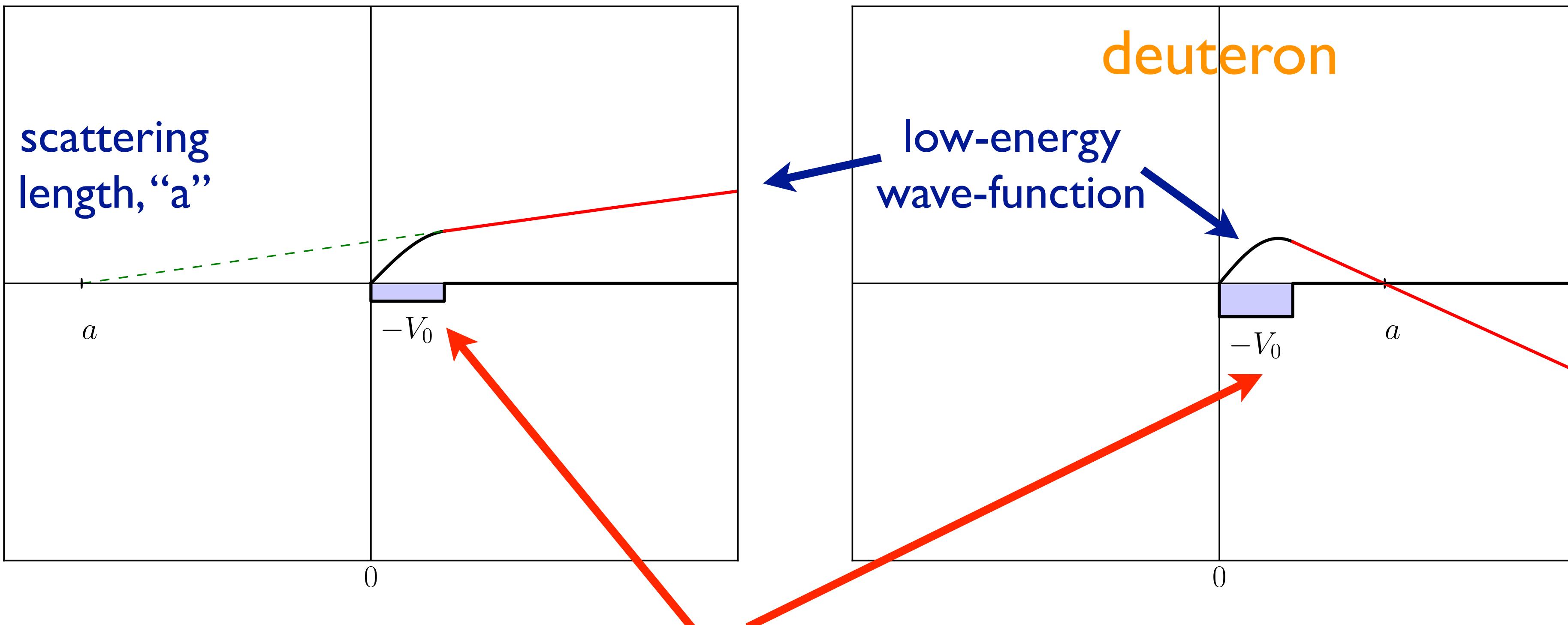


No Sun!

proton-neutron scattering at low energies

$$^1 S_0 : a \simeq -24 \text{ fm}$$

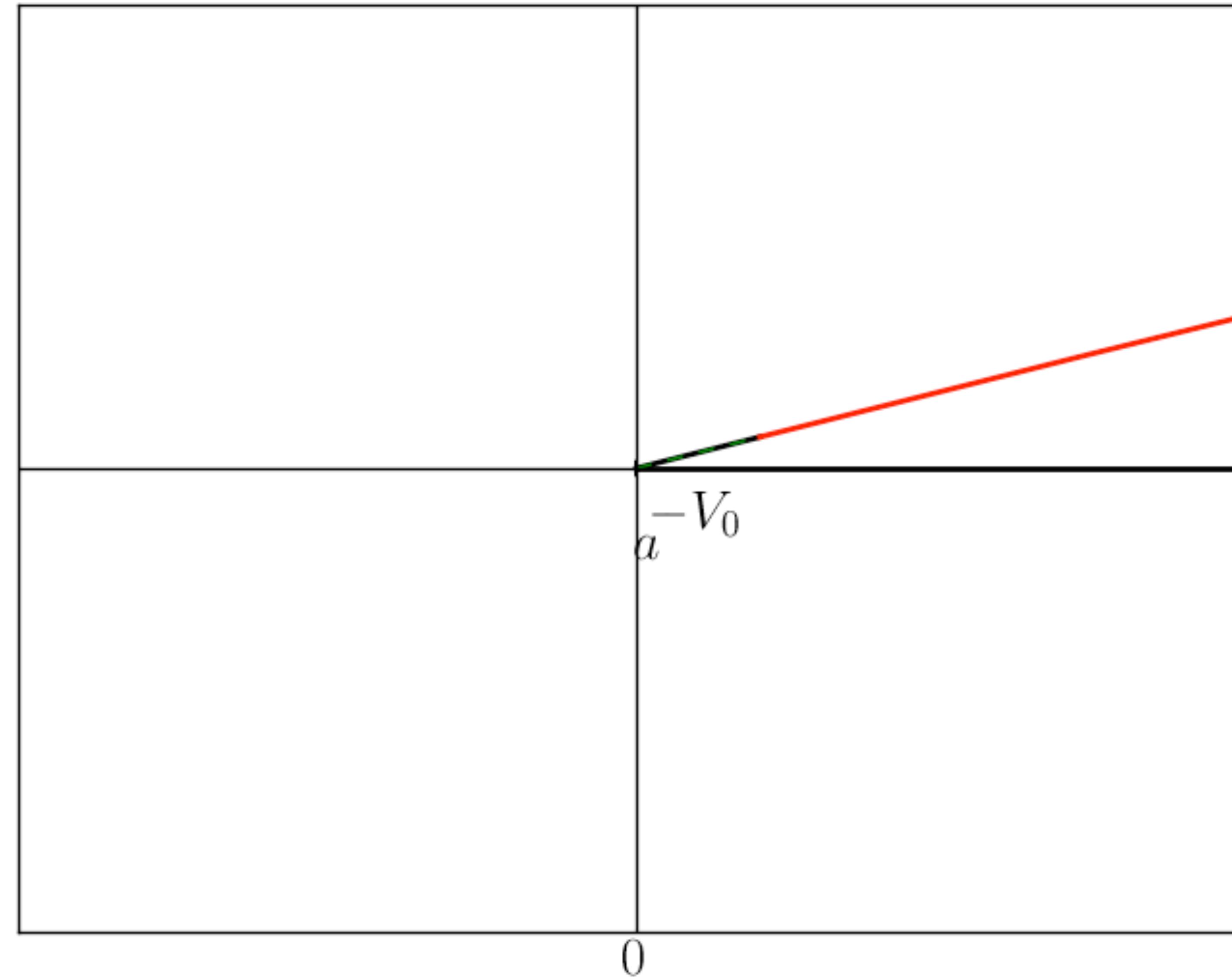
$$^3 S_1 : a \simeq 5.5 \text{ fm}$$



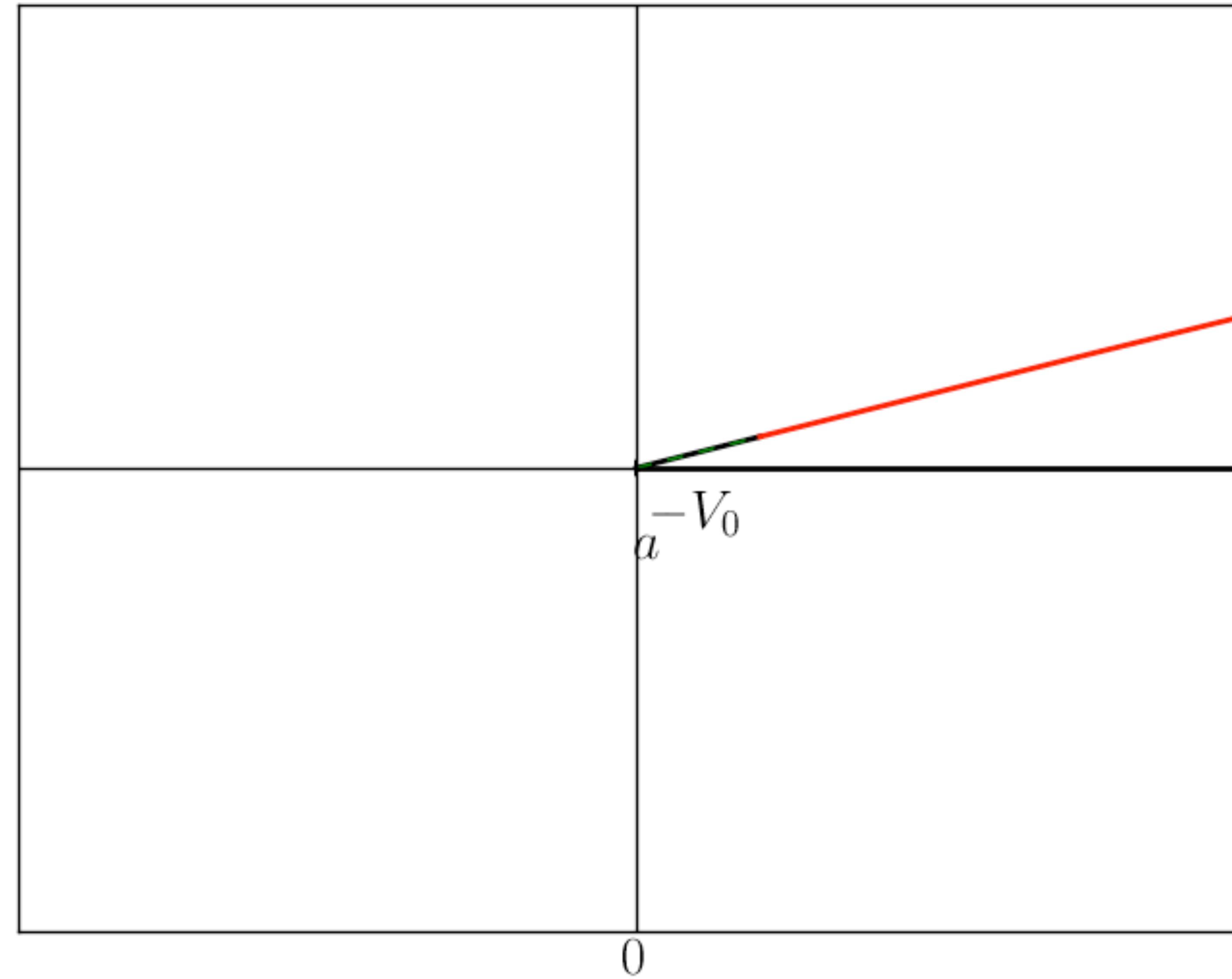
Fine tuning gives small deuteron binding energy

Solving QCD can help us determine the nature of this fine tuning

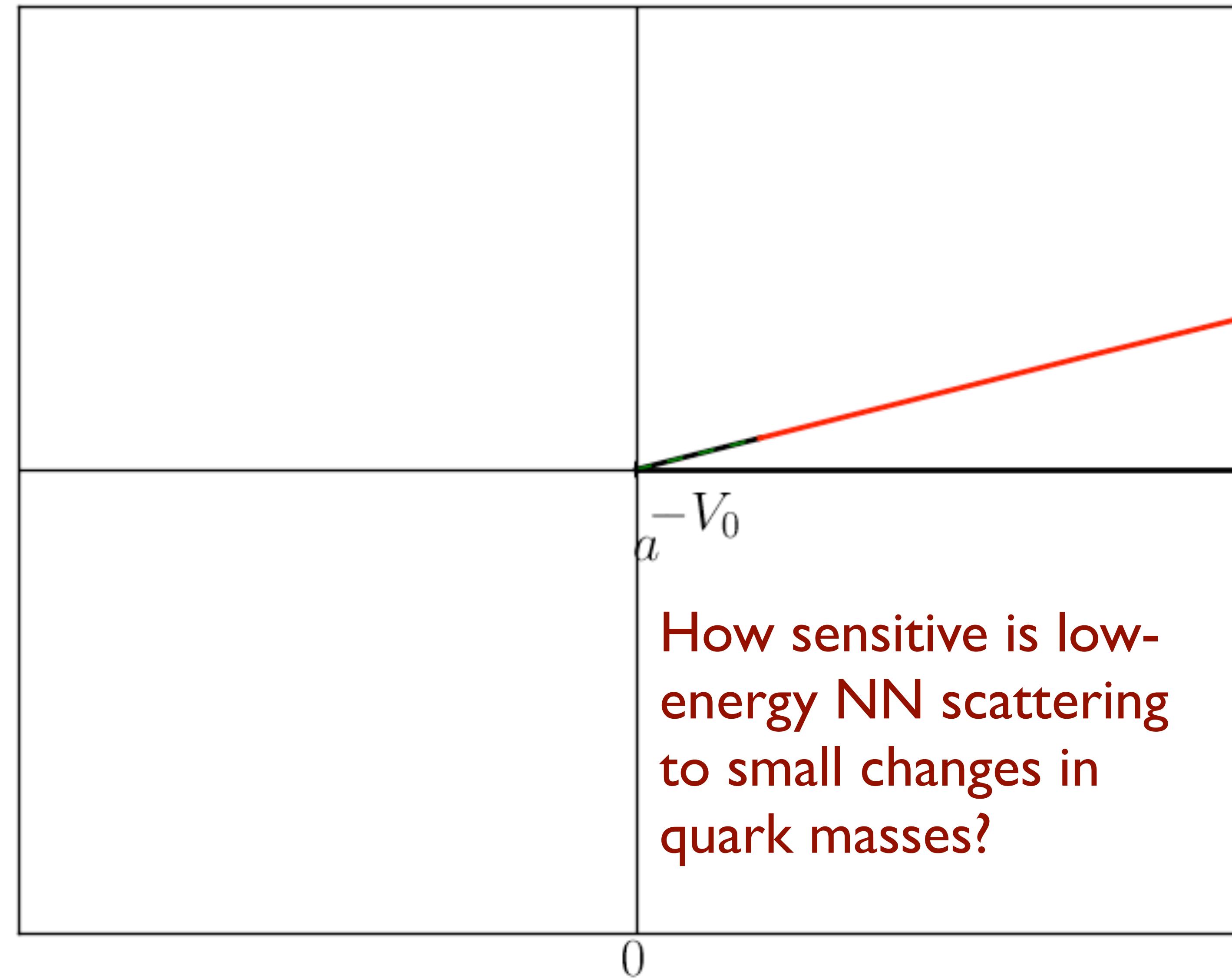
Finely tuned interactions (like in AMO systems)



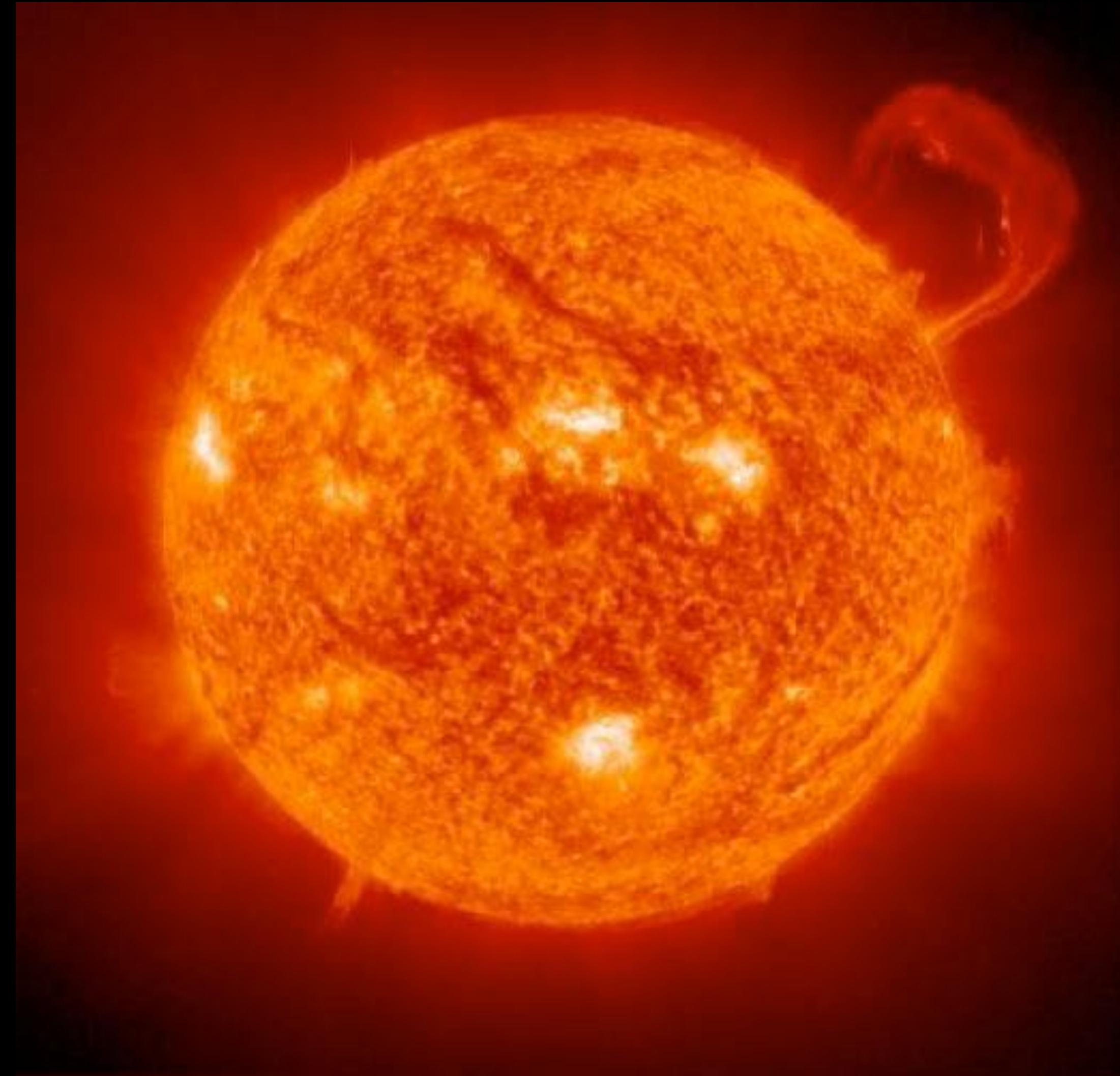
Finely tuned interactions (like in AMO systems)



Finely tuned interactions (like in AMO systems)



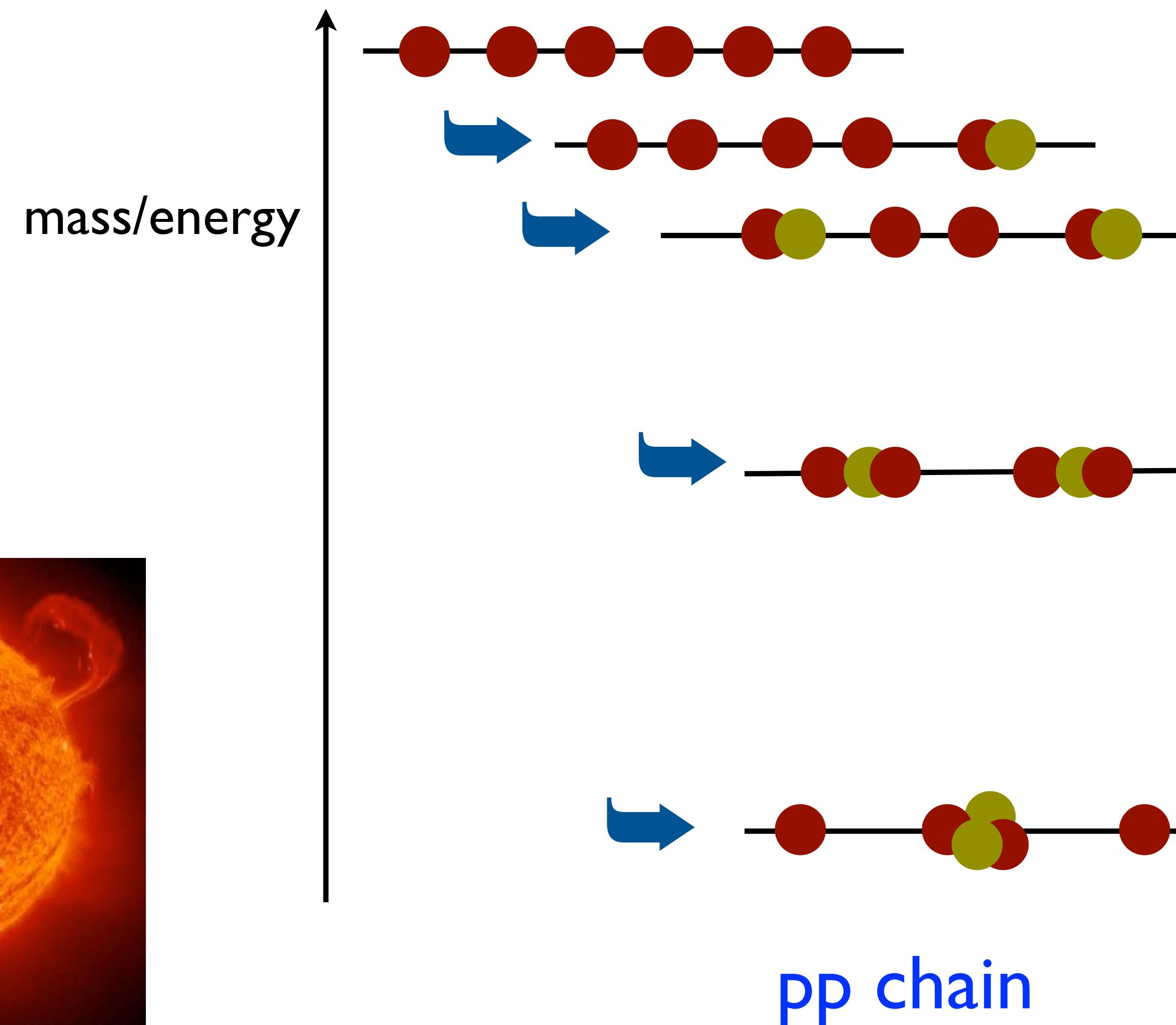
Solar Fusion



$T \simeq 20$ K
 $t \simeq 200$ Million years

One needs neutrons and protons to make new nuclei.

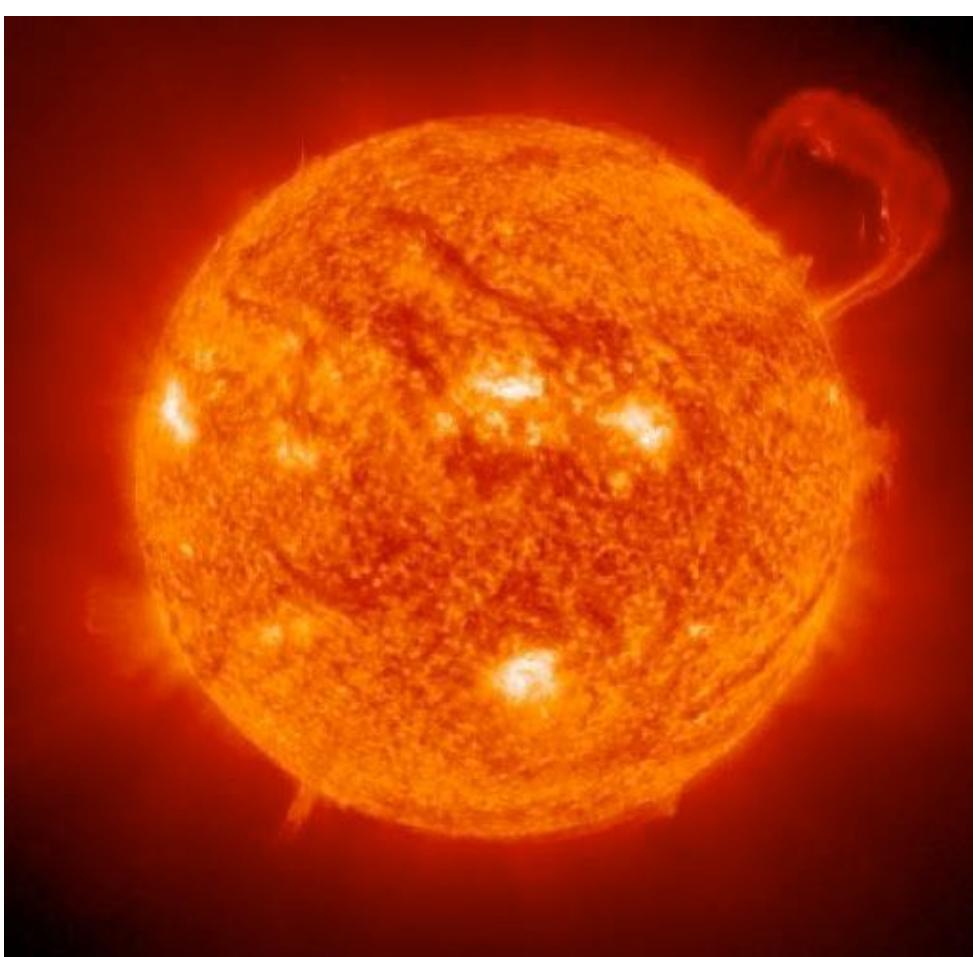
Small stars burn protons only, manufacturing the needed neutrons



$p + p \rightarrow d + \nu_e$
another bottleneck
deuteron: -1.2 MeV
2 deuterons: -2.4 MeV

2 ^3He : -12.4 MeV

helium: -28.3 MeV
thermal energy
to support star

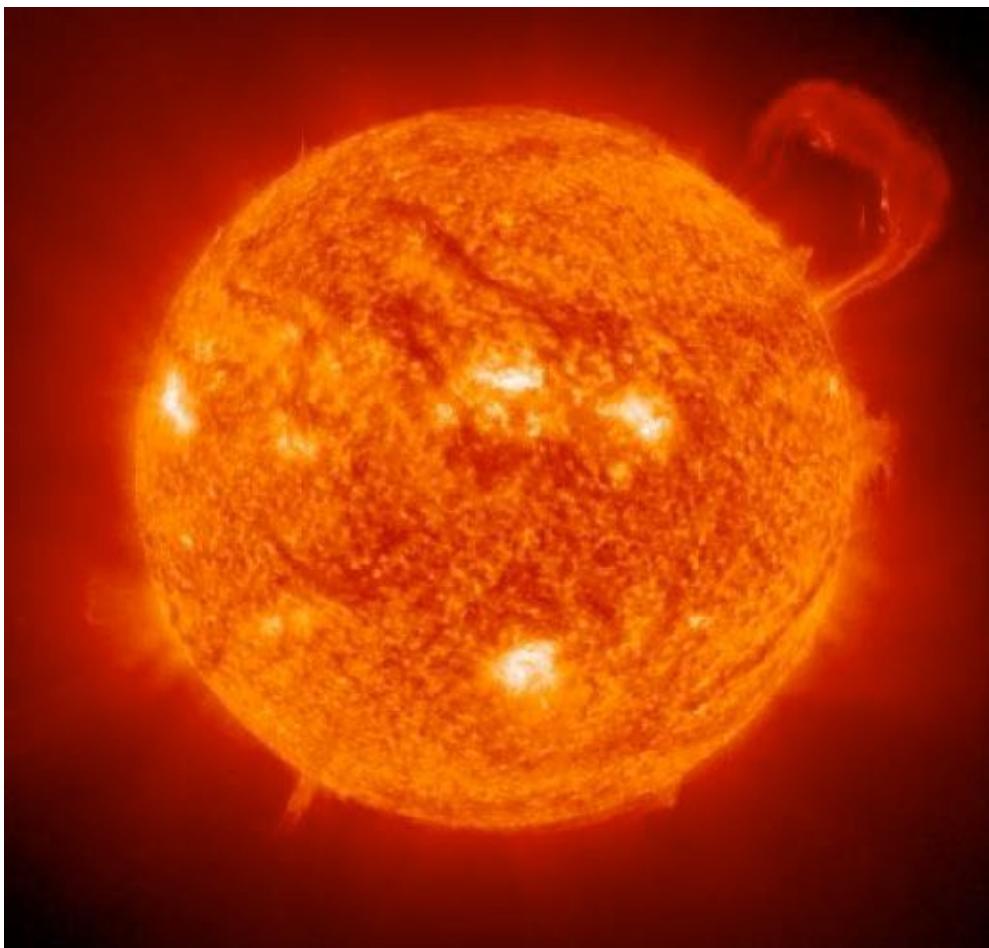


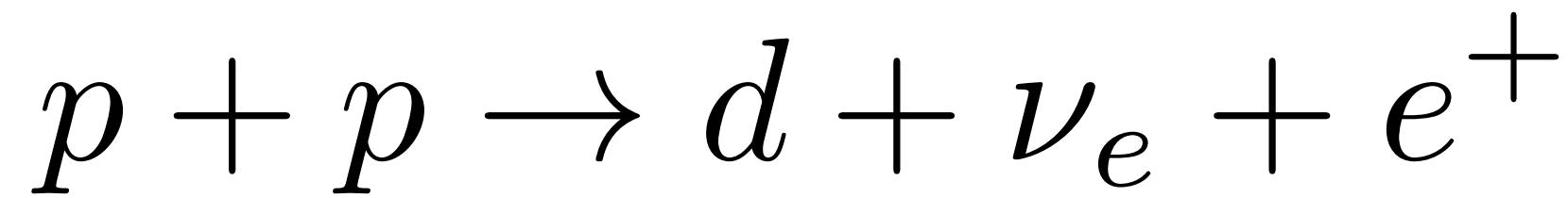
This is how our Sun generates its energy

80% of all stars generate their energy by hydrogen burning

At its very center the Sun generates 275 watts/m³ - similar to the energy generated by a **compost (garbage) heap (of the same size)!**

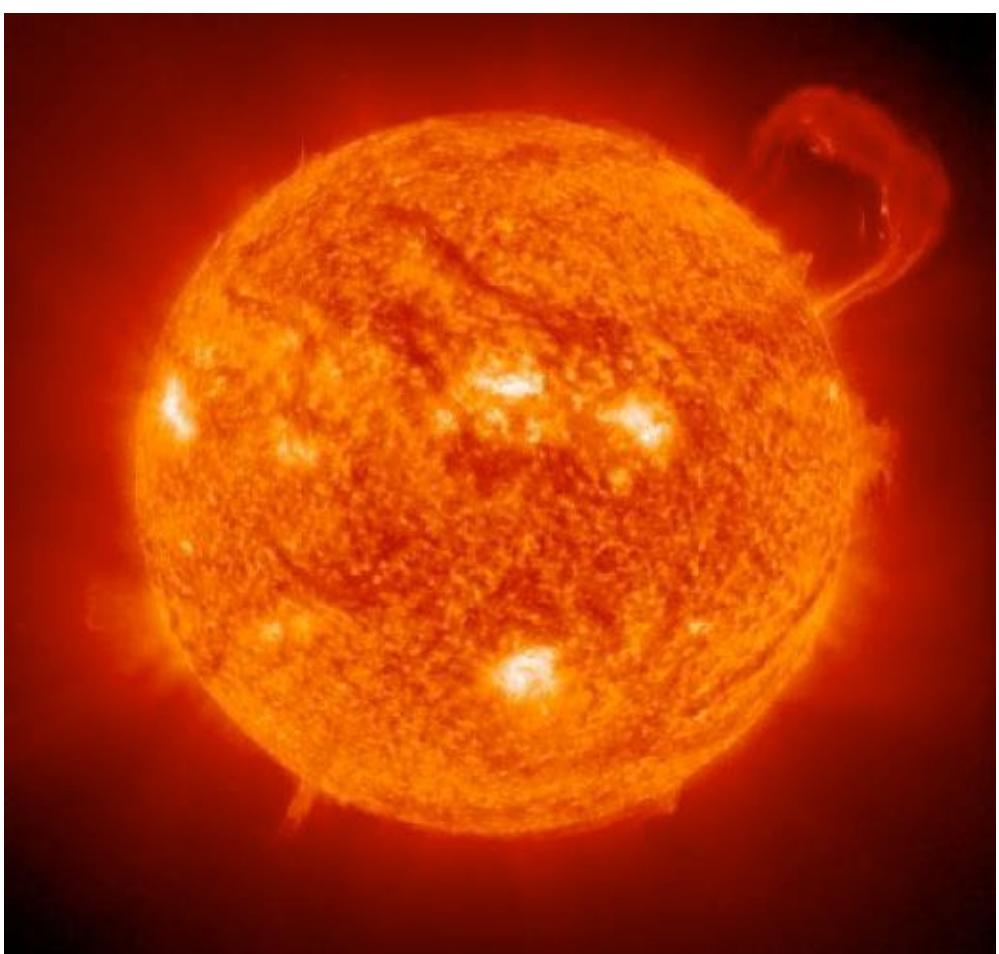
And this is why the Sun has burned for 4.6 b.y., and will burn for 5 b.y. more, fortunately -- a very big, very slow reactor

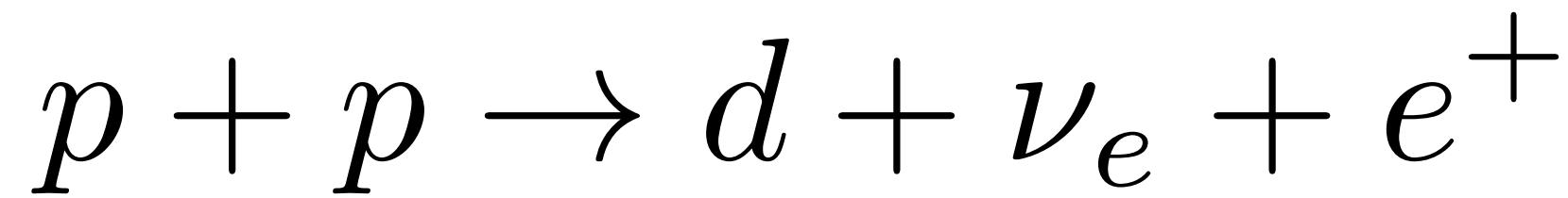




This fundamental reaction can not be measured!
(Coulomb Repulsion)

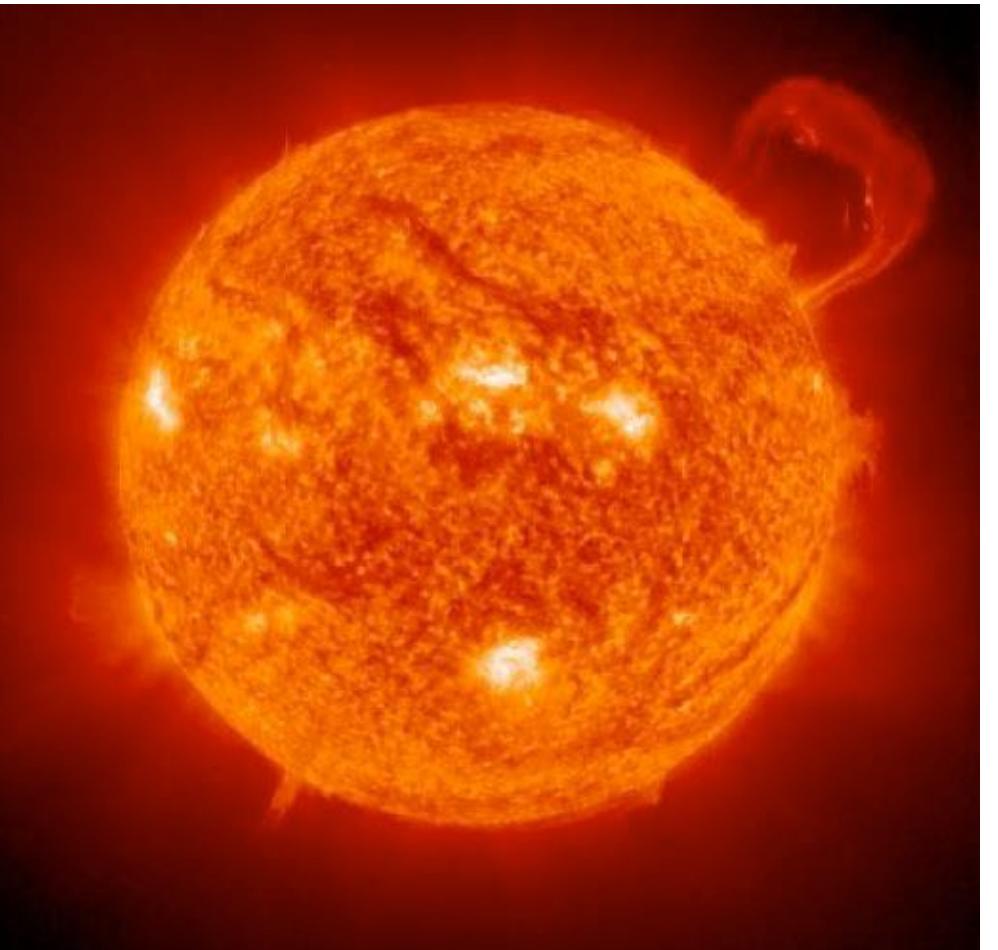
We believe we know the rate, but we have not been able to predict it directly from the fundamental theory,
can we?





This fundamental reaction can not be measured!
(Coulomb Repulsion)

We believe we know the rate, but we have not been able to predict it directly from the fundamental theory,
can we?

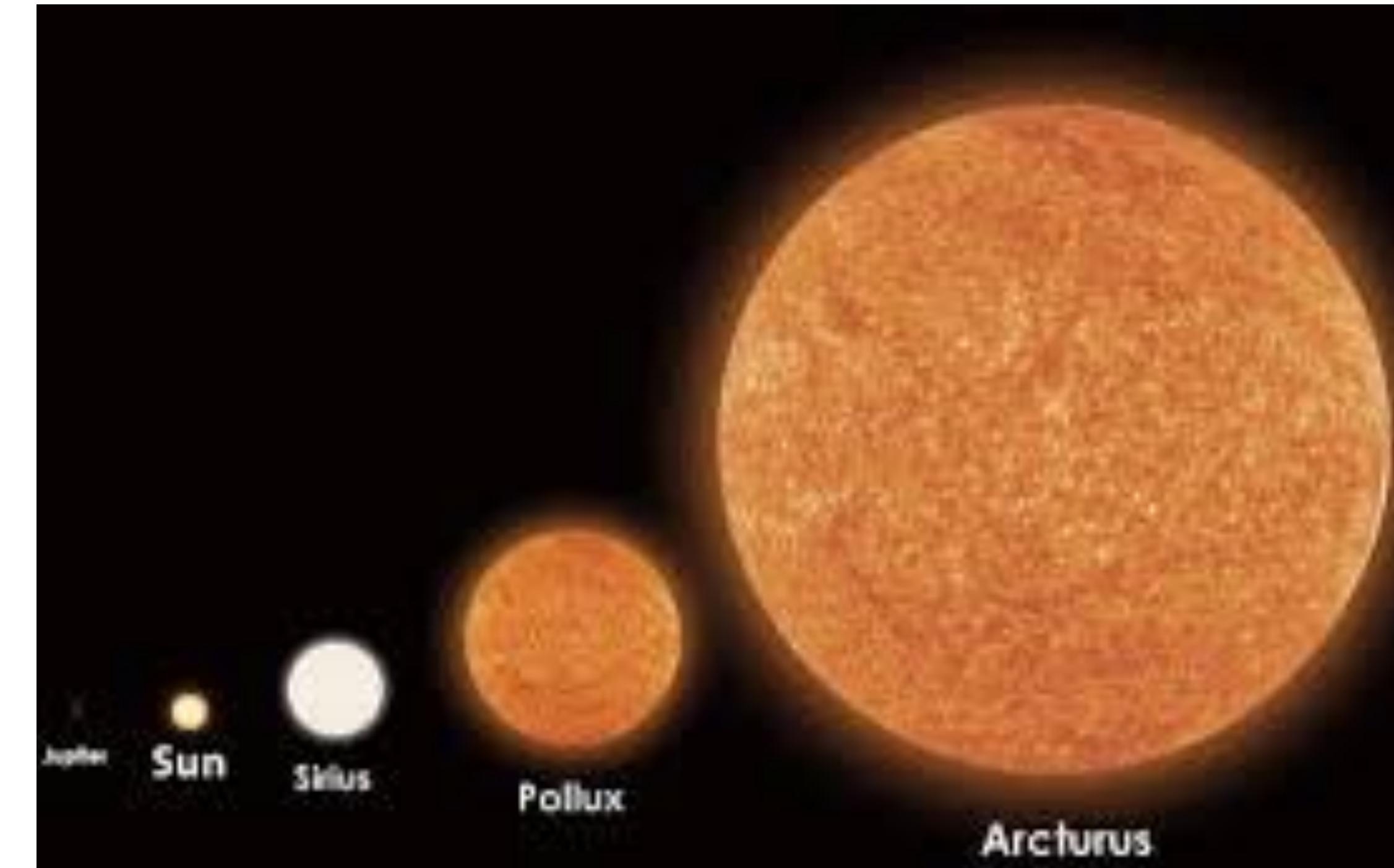
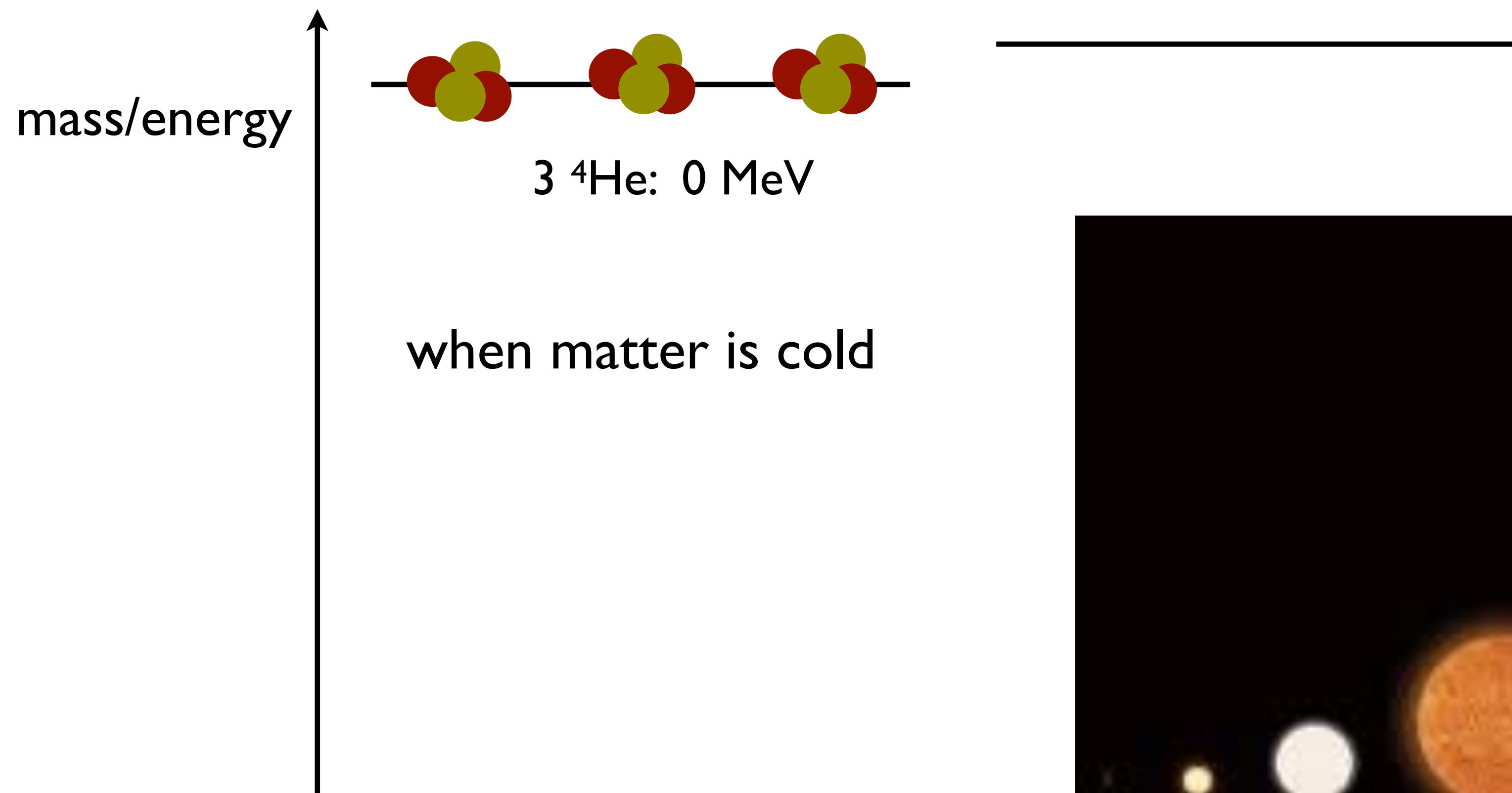


Yes!
With  !
(or equivalent pre-exascale
super computer)

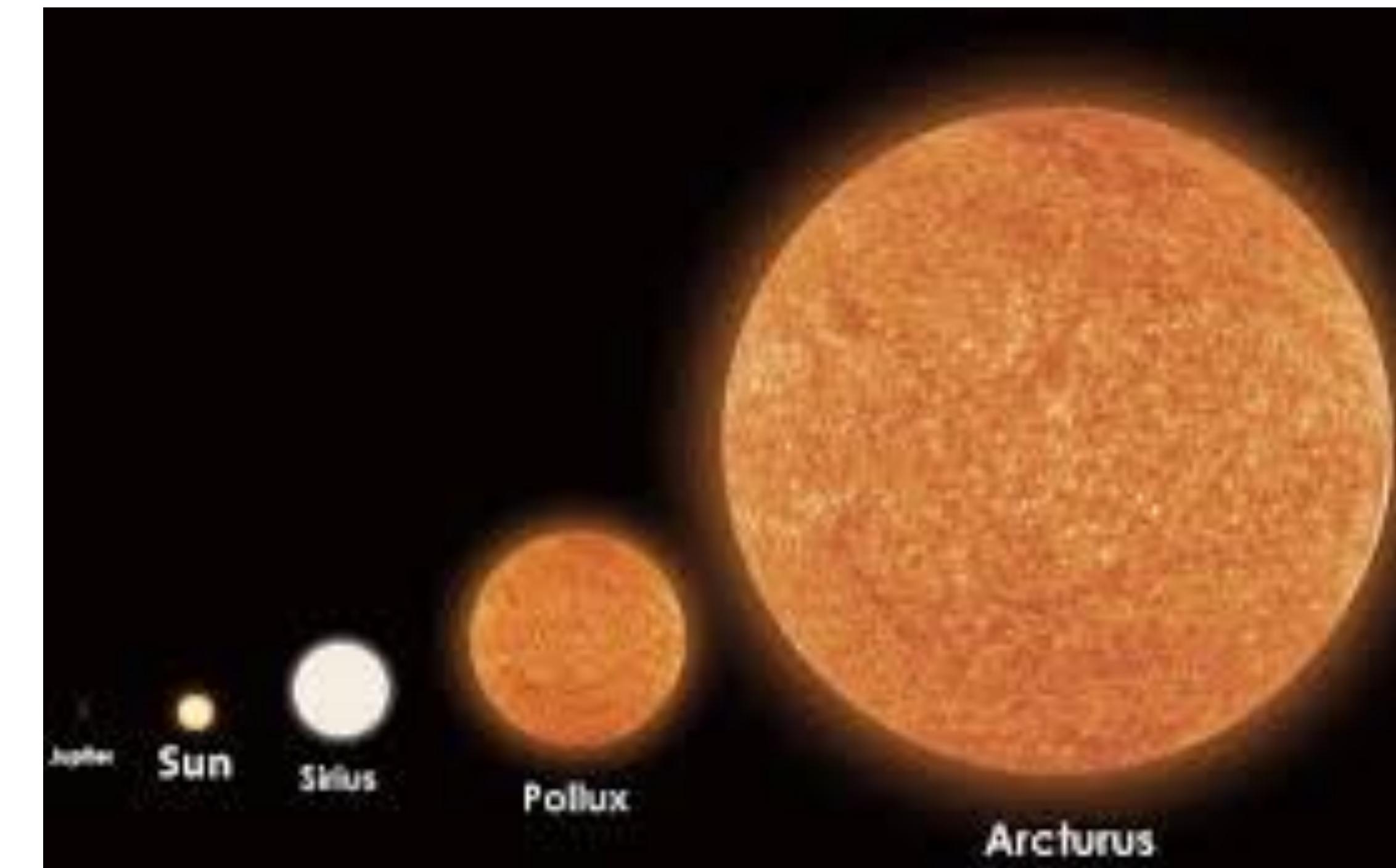
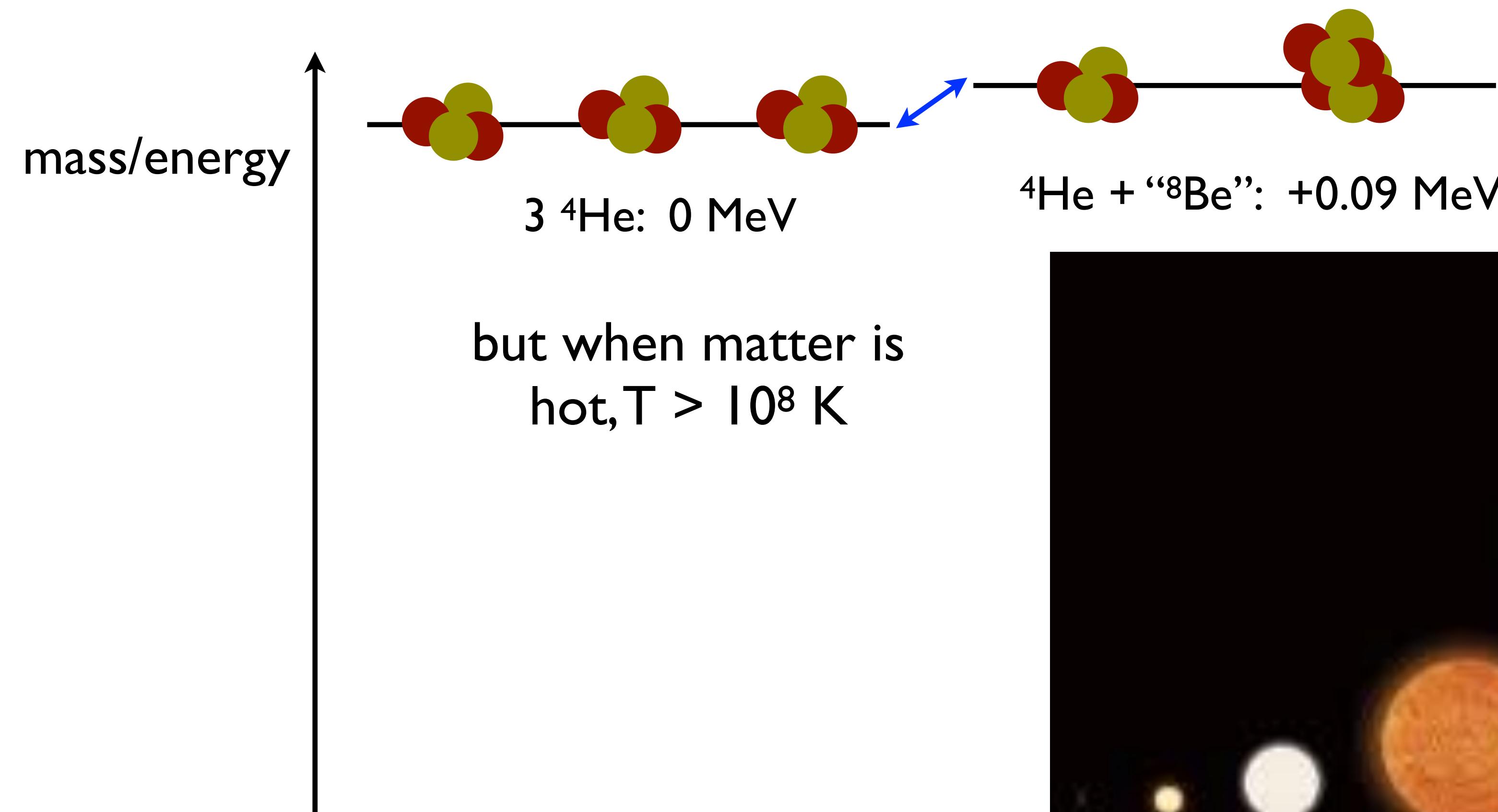
Large stars use He and neutrons to build new nuclei.

Higher temperatures and higher densities are needed.

The Big Bang could not do this because the density was too low.

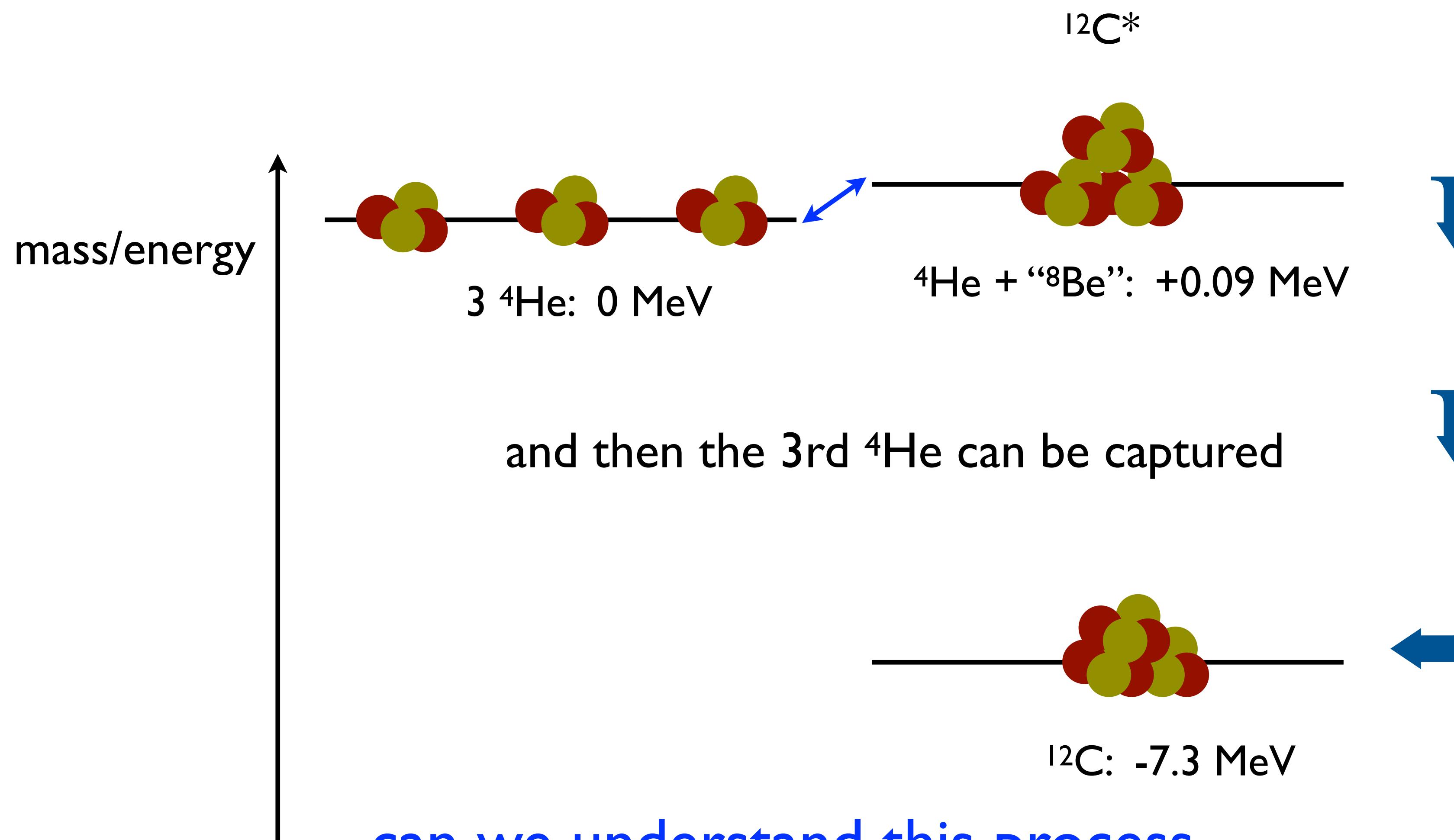


even more finely tuned



even more finely tuned - source of complex life

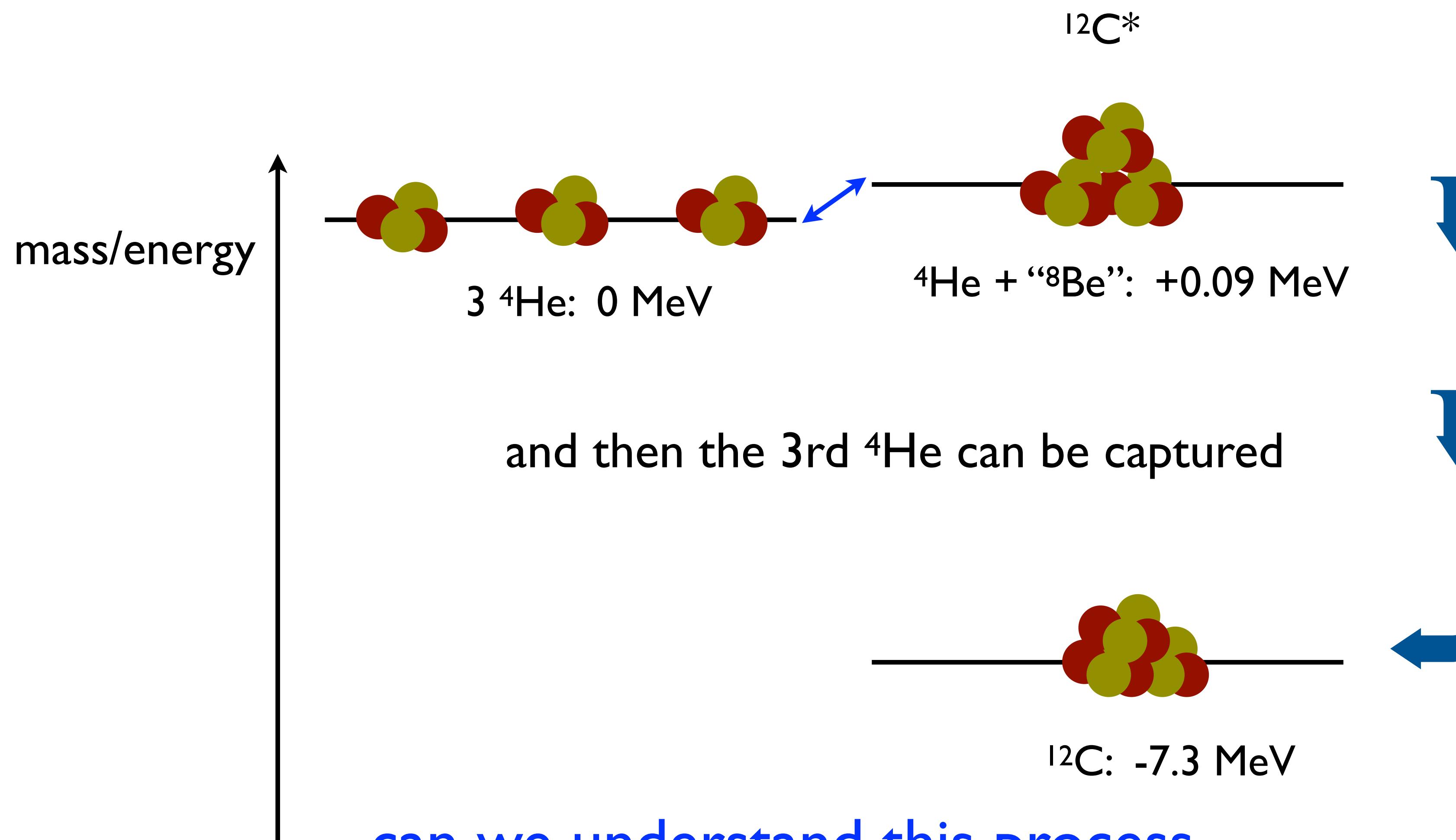
the triple- α process **Hoyle State**



can we understand this process
from the fundamental theory?

even more finely tuned - source of complex life

the triple- α process **Hoyle State**

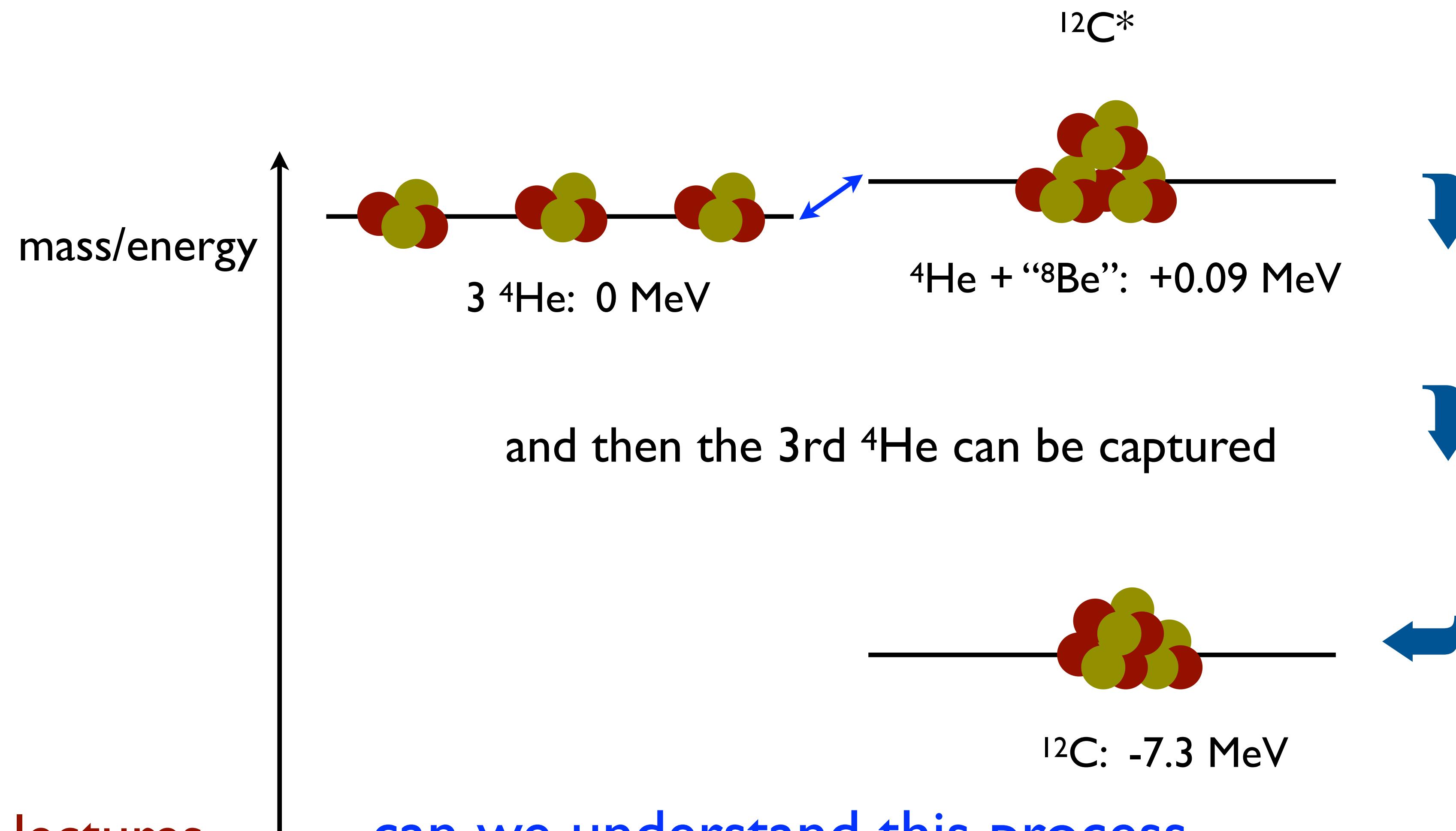


can we understand this process
from the fundamental theory?

Yes!
maybe with *summit!*

even more finely tuned - source of complex life

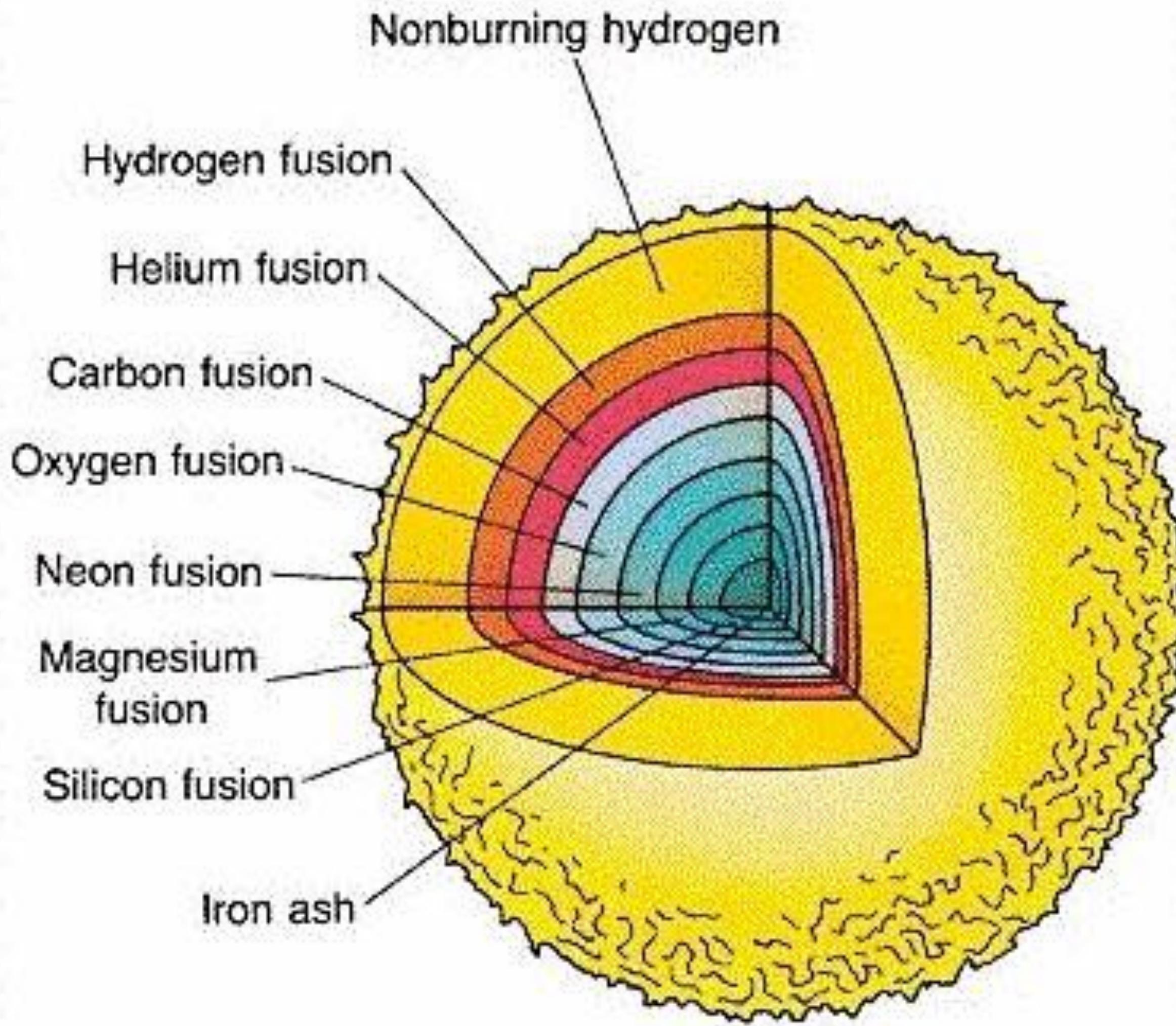
the triple- α process **Hoyle State**



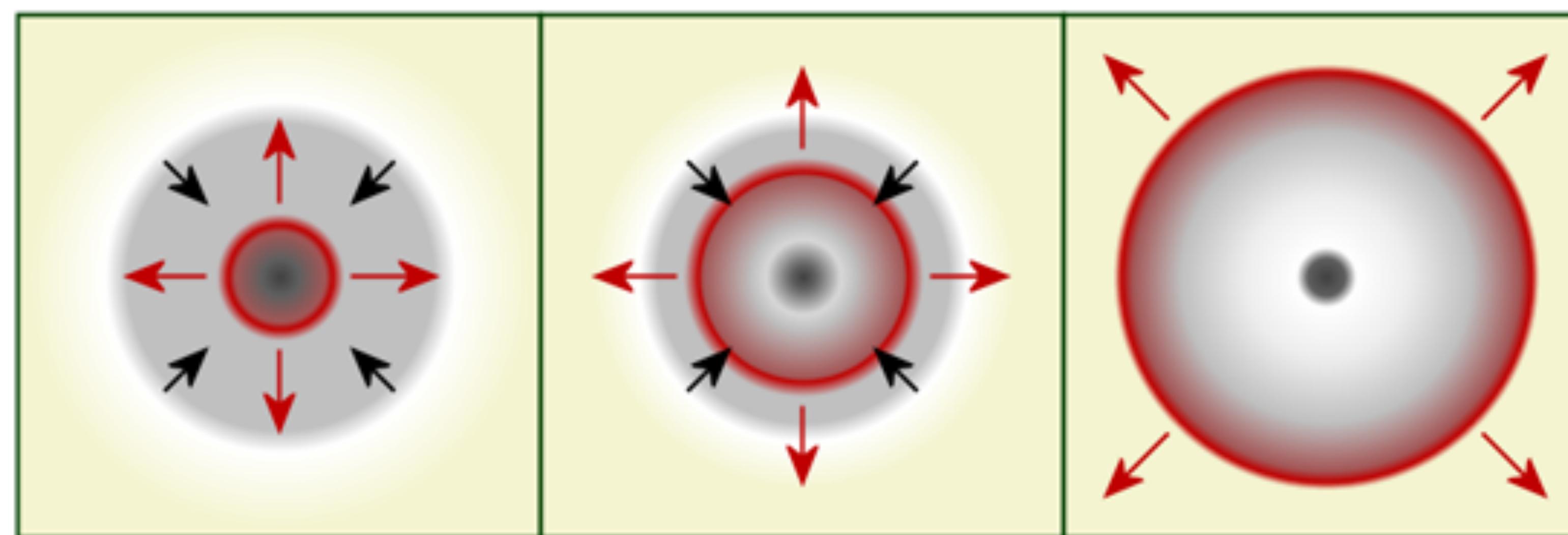
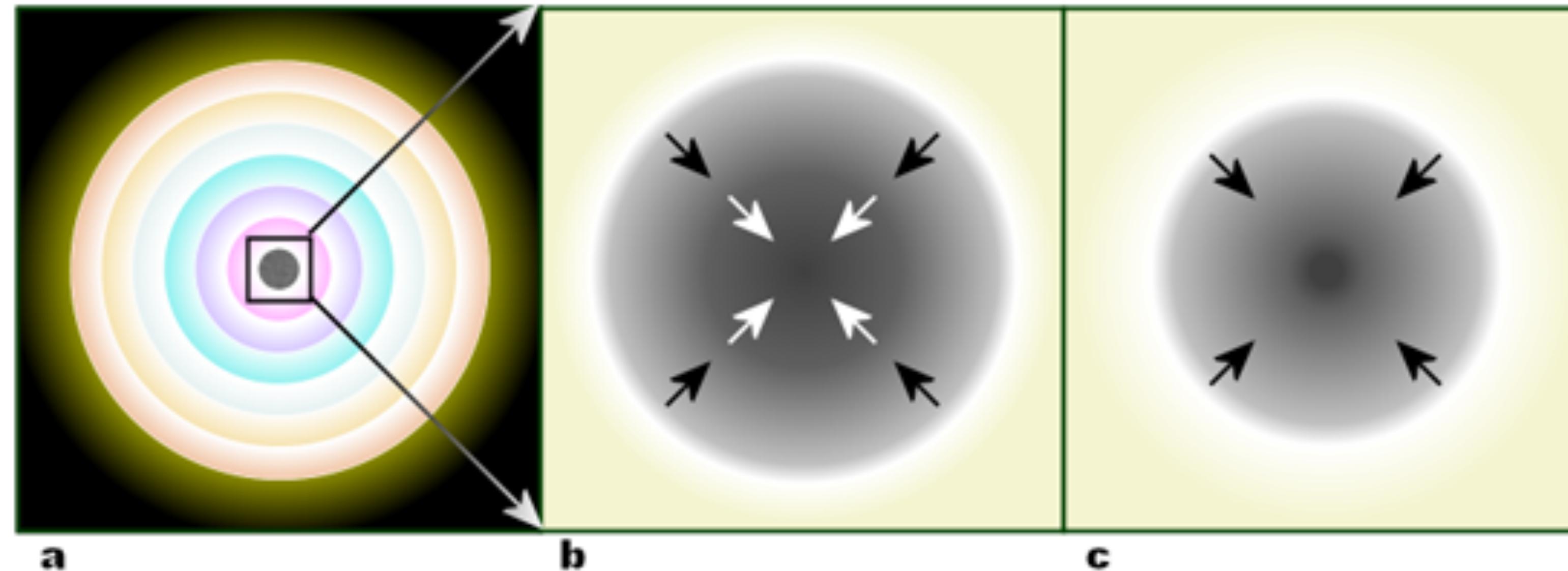
Also - see lectures
by Ulf Meissner

can we understand this process
from the fundamental theory?

Yes!
maybe with *summit!*



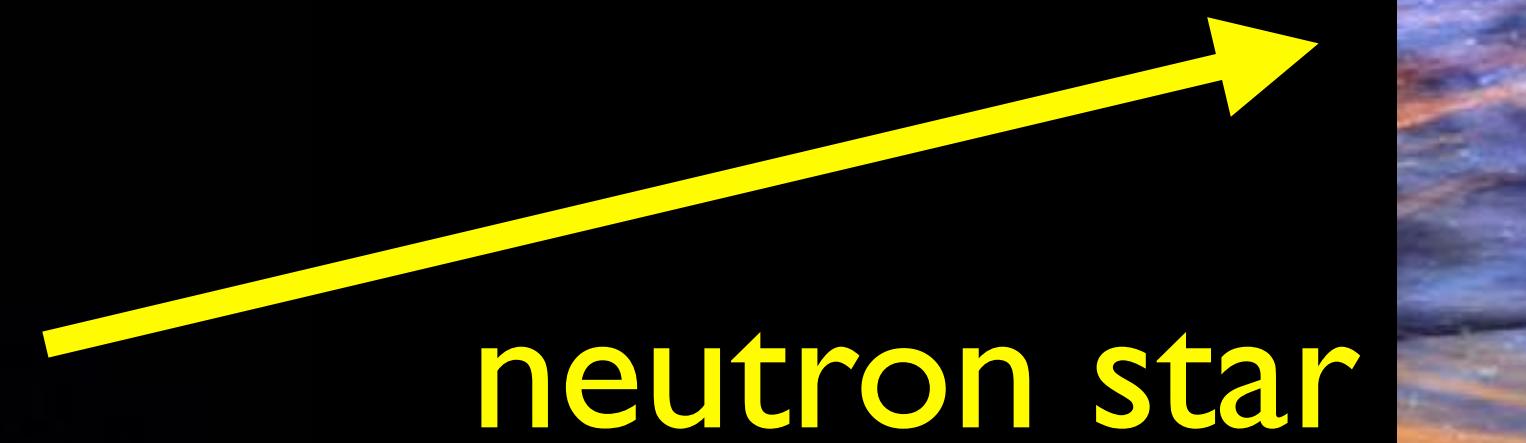
He, C, O, ... Si burning produces energy until Iron (Fe)



**core collapse supernova, shock-wave-aided ejection of
mantle**



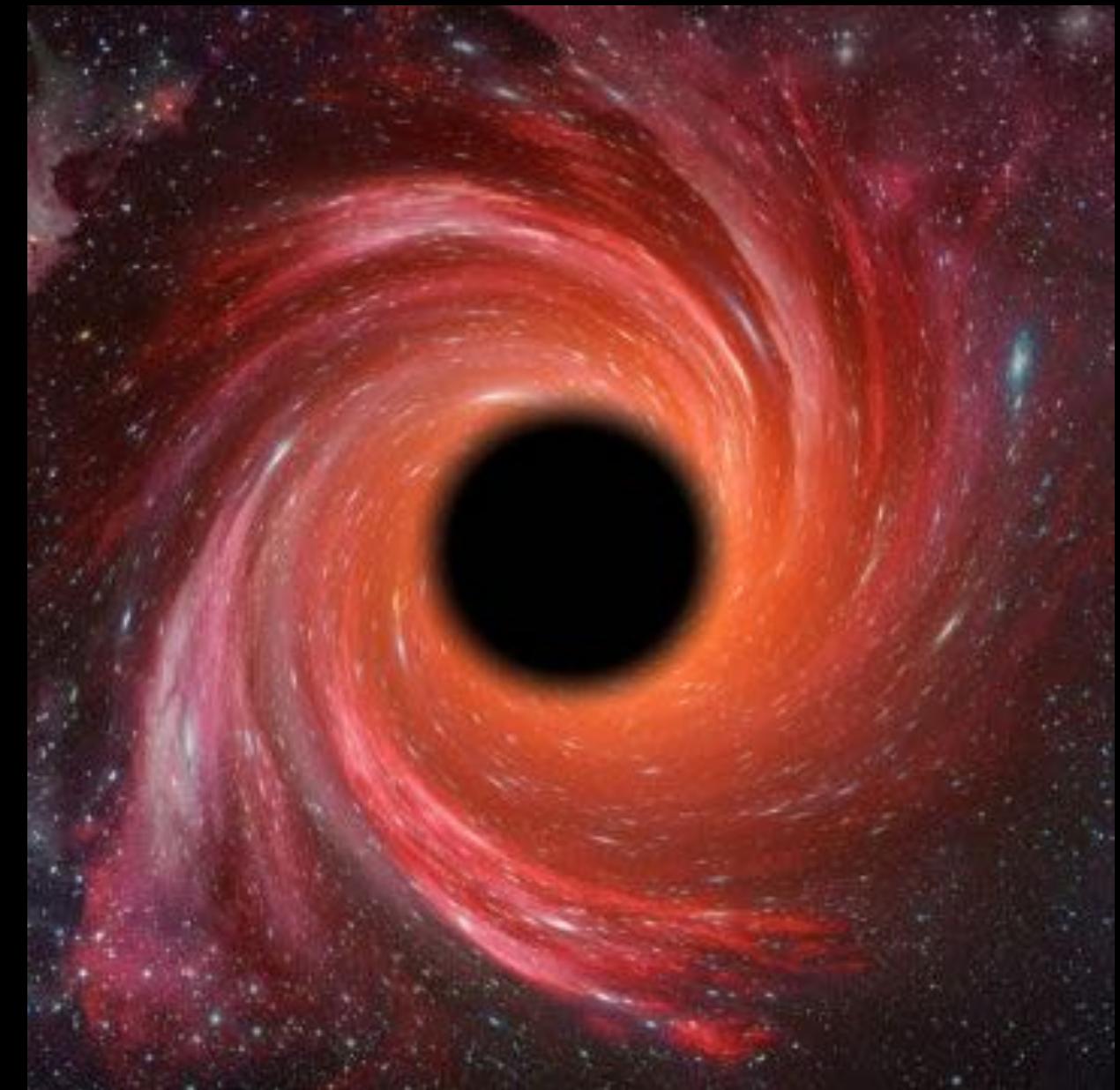
Supernova



neutron star



black hole



In Memoriam

Carl Sagan

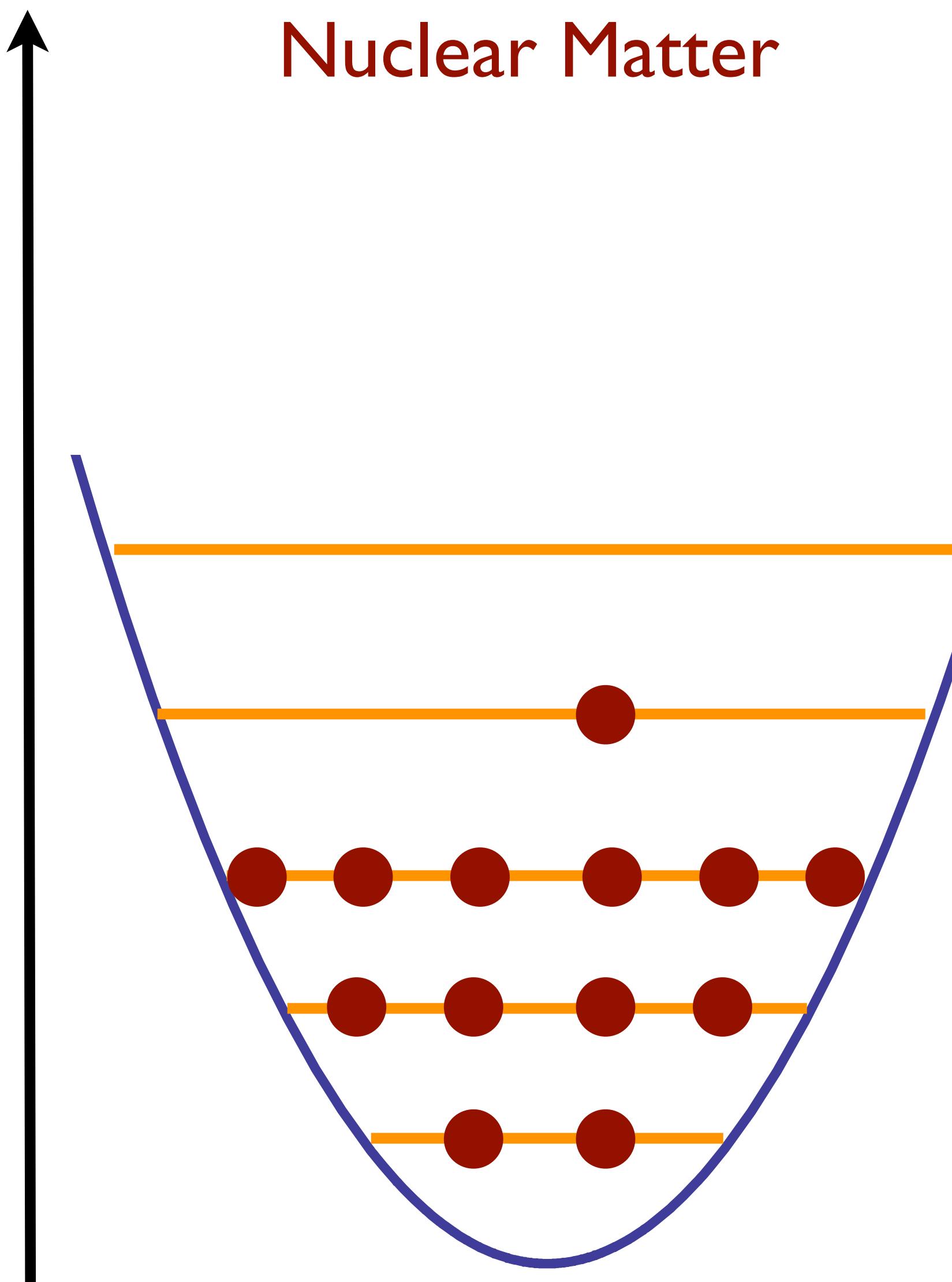
1934-1996

“We are all made of star stuff”

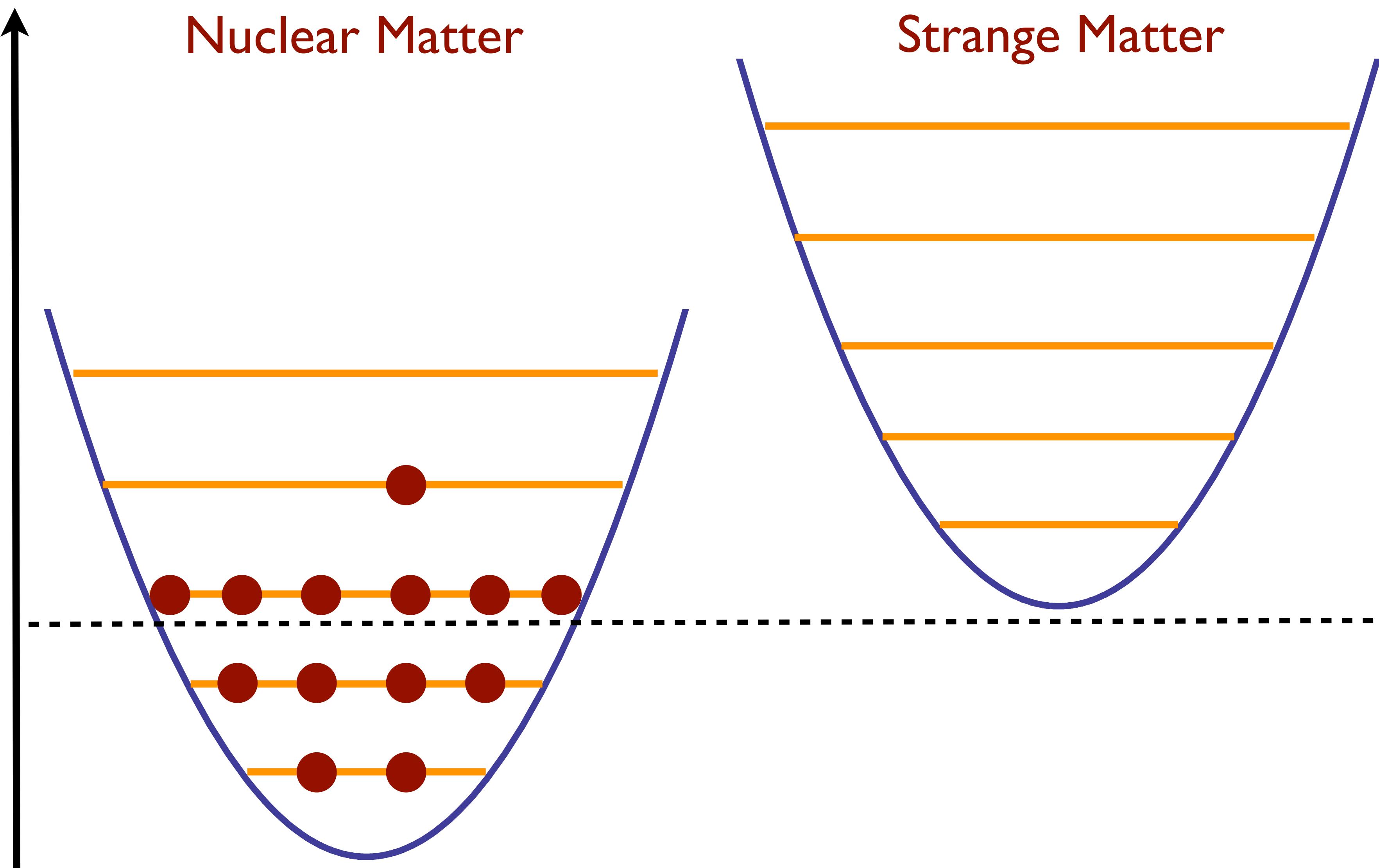
much coming from the ejecta of supernova

mass/energy

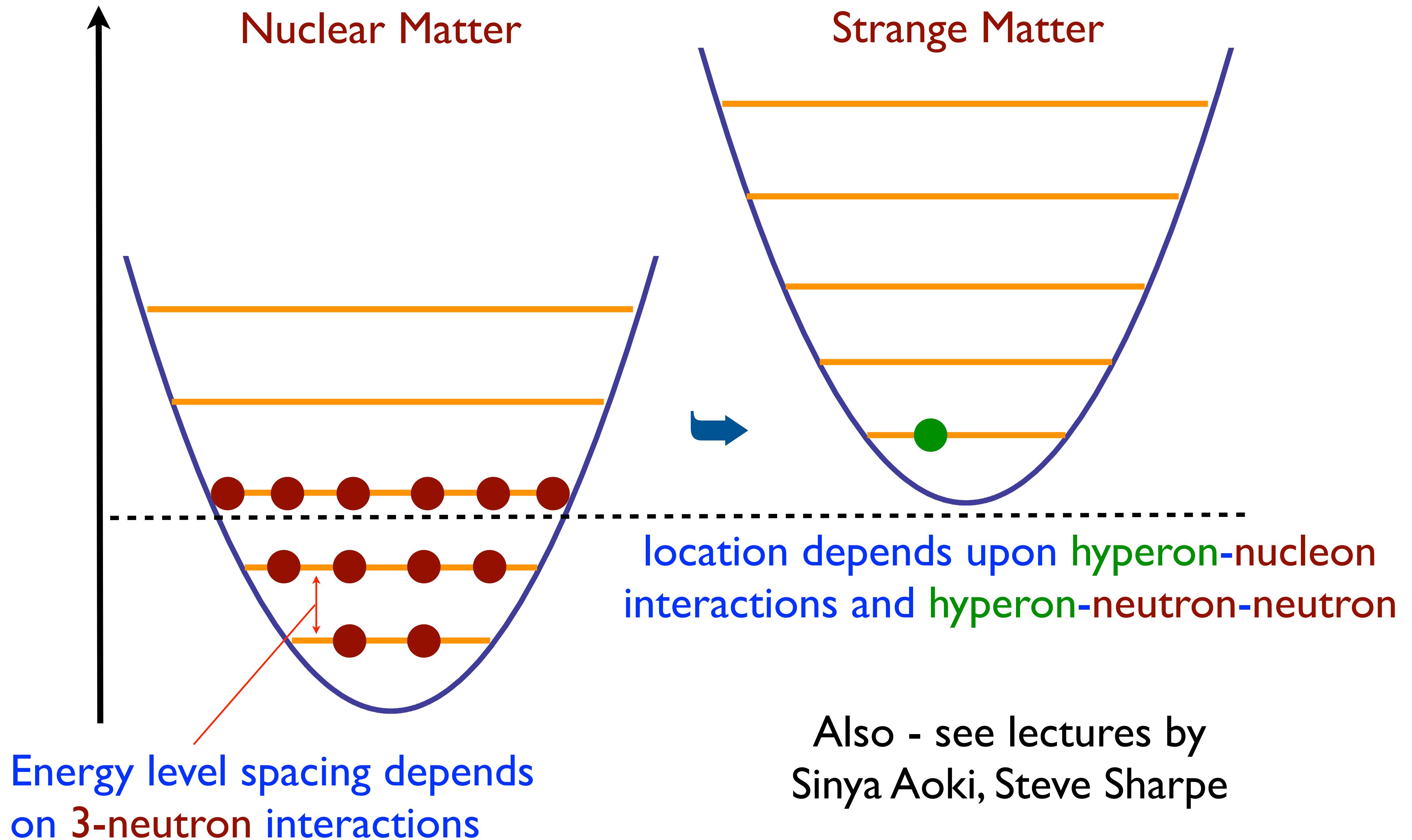
Nuclear Matter



mass/energy



mass/energy





**Can we understand properties of neutron stars
directly from the fundamental theory?**



Can we understand properties of neutron stars
directly from the fundamental theory?

Yes! With **FRONTIER**

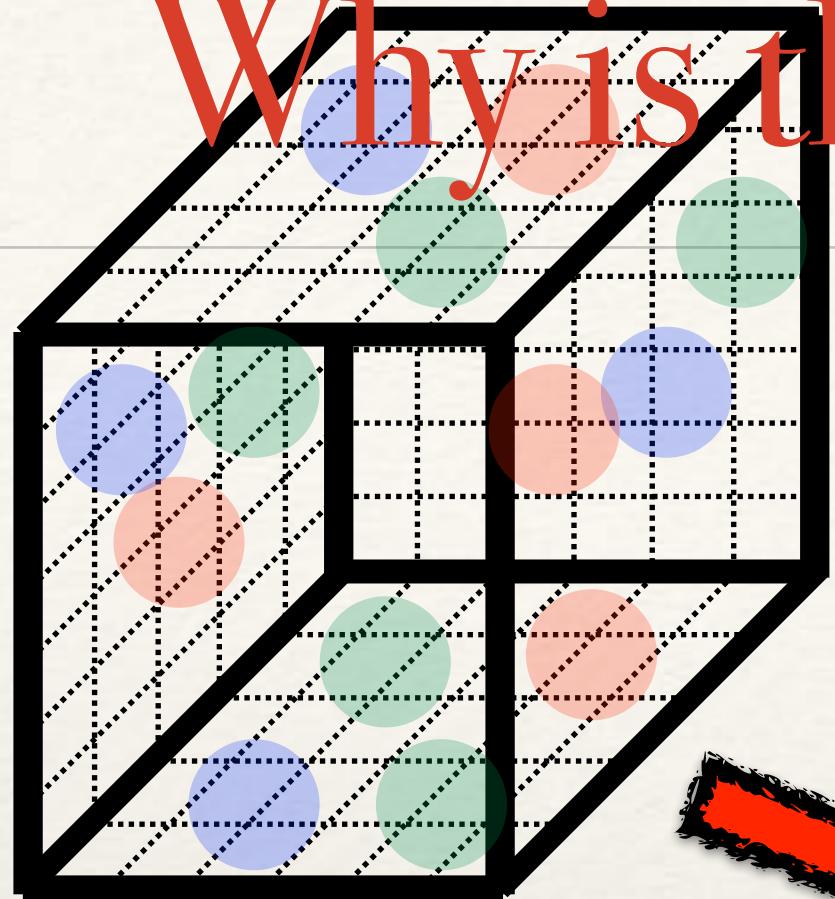
Why is the application of LQCD to NP so challenging?

Why is the application of LQCD to NP so challenging?

continuum limit

need 3 or more
lattice spacings

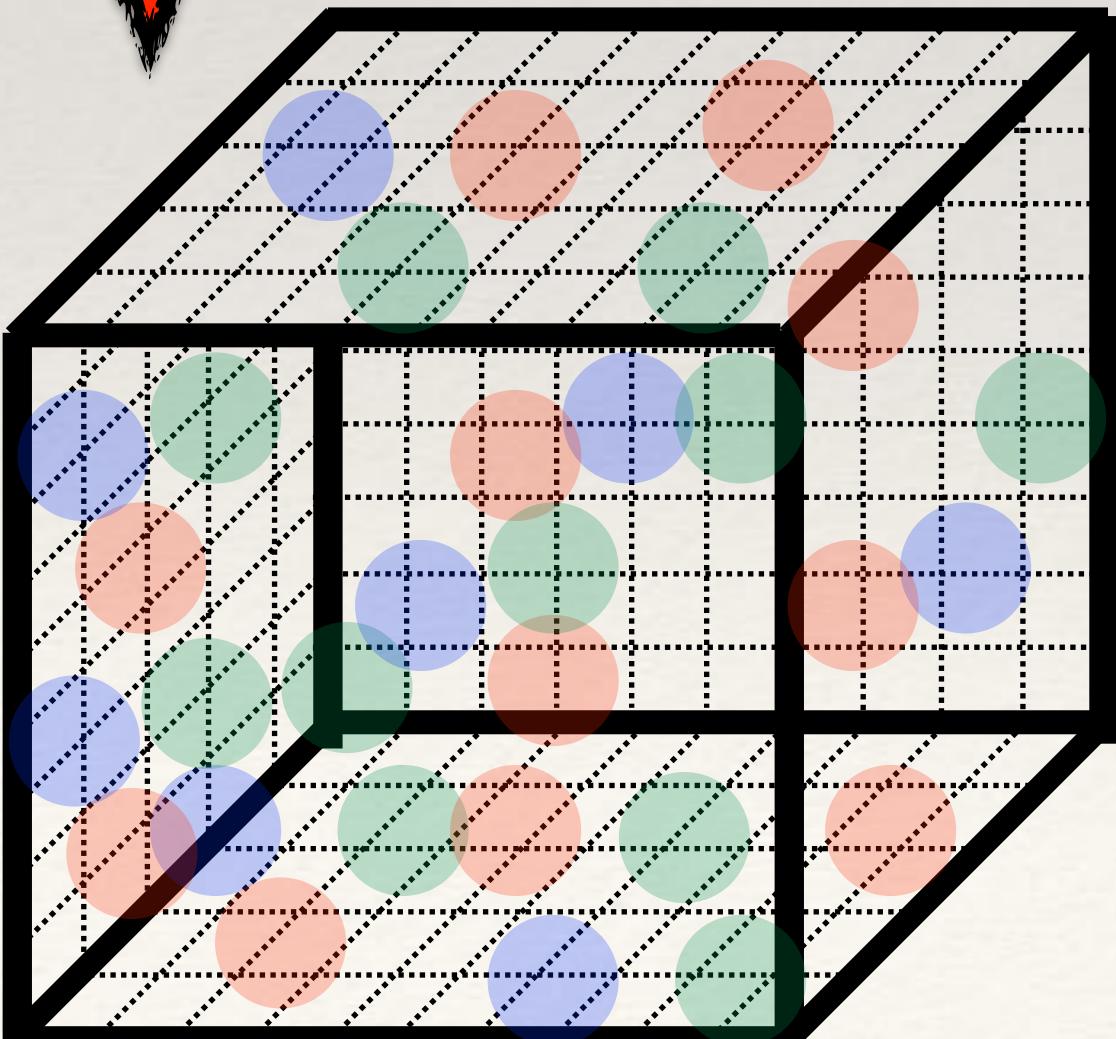
$$t_{comp} \propto \frac{1}{a^6}$$



infinite volume limit

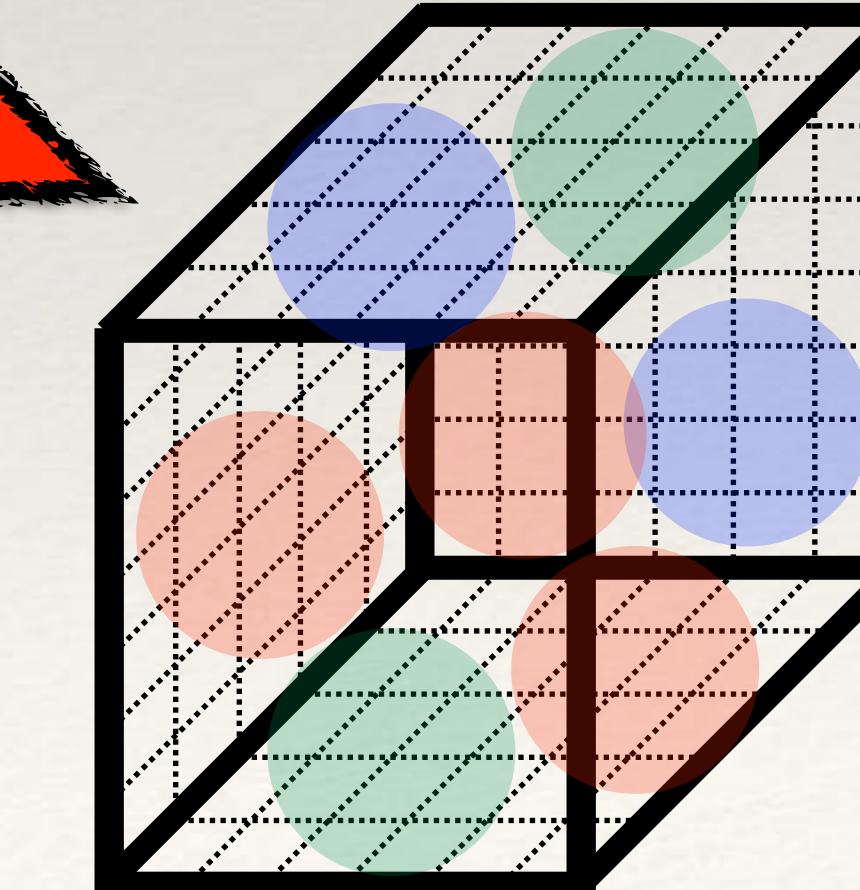
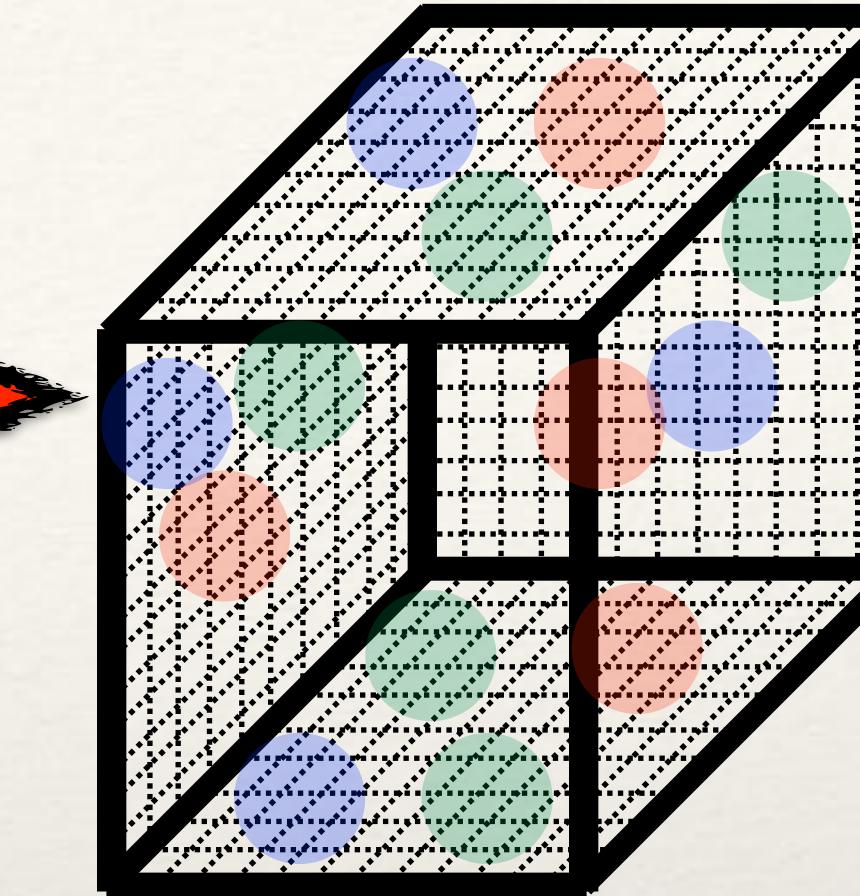
$$t_{comp} \propto V^{5/4}$$

$$V = N_L^3 \times N_T$$



physical pion masses

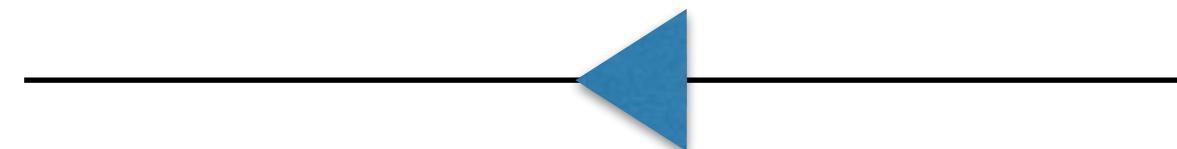
exponentially bad
signal-to-noise problem



LQCD challenges for NP

Most difficult challenge: an **exponentially bad signal-to-noise problem**

Parisi, Phys. Rep. 103 (1984) 203

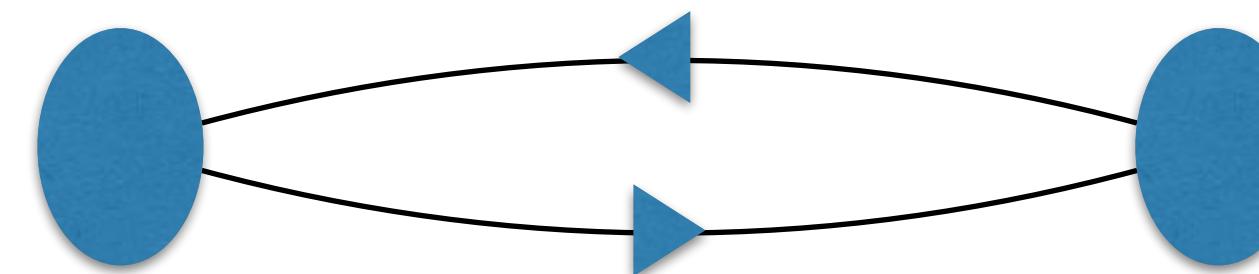

$$\sim e^{-\frac{1}{2}m_\pi t} + e^{-\frac{1}{3}m_N t} + \dots$$

Lepage, TASI 1989

Each **quark propagator** carries information about pions and nucleons
(conversations with David Kaplan)

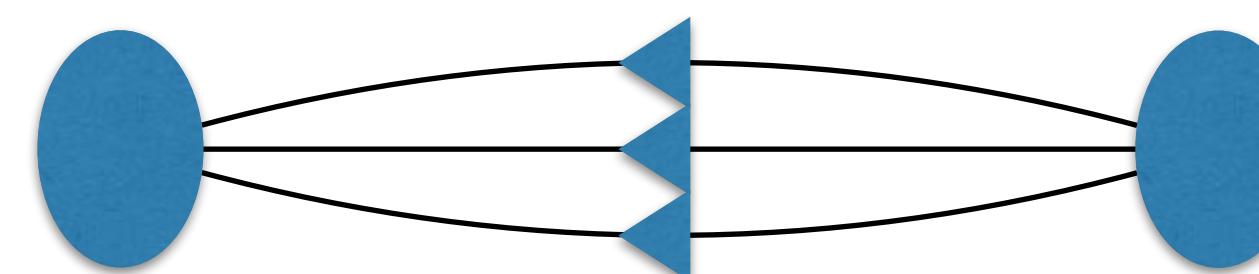
$$\lambda_\pi(t) \gg \lambda_N(t)$$

$$\lambda_i(t) \sim e^{-E_i t}$$



$$\bar{d}\gamma_5 u : C(t) = A_\pi e^{-m_\pi t} + \dots$$

Large pion eigenvalues must cancel to expose small nucleon eigenvalues



$$(u^T C \gamma_5 d) u : C(t) = A_N e^{-m_N t} + \dots$$

$$\frac{\text{Signal}}{\text{Noise}} \sim \sqrt{N} \exp \left[-A \left(m_N - \frac{3}{2} m_\pi \right) t \right] \rightarrow \text{exponential noise power-law statistics}$$

LQCD challenges for NP

2-point correlation function

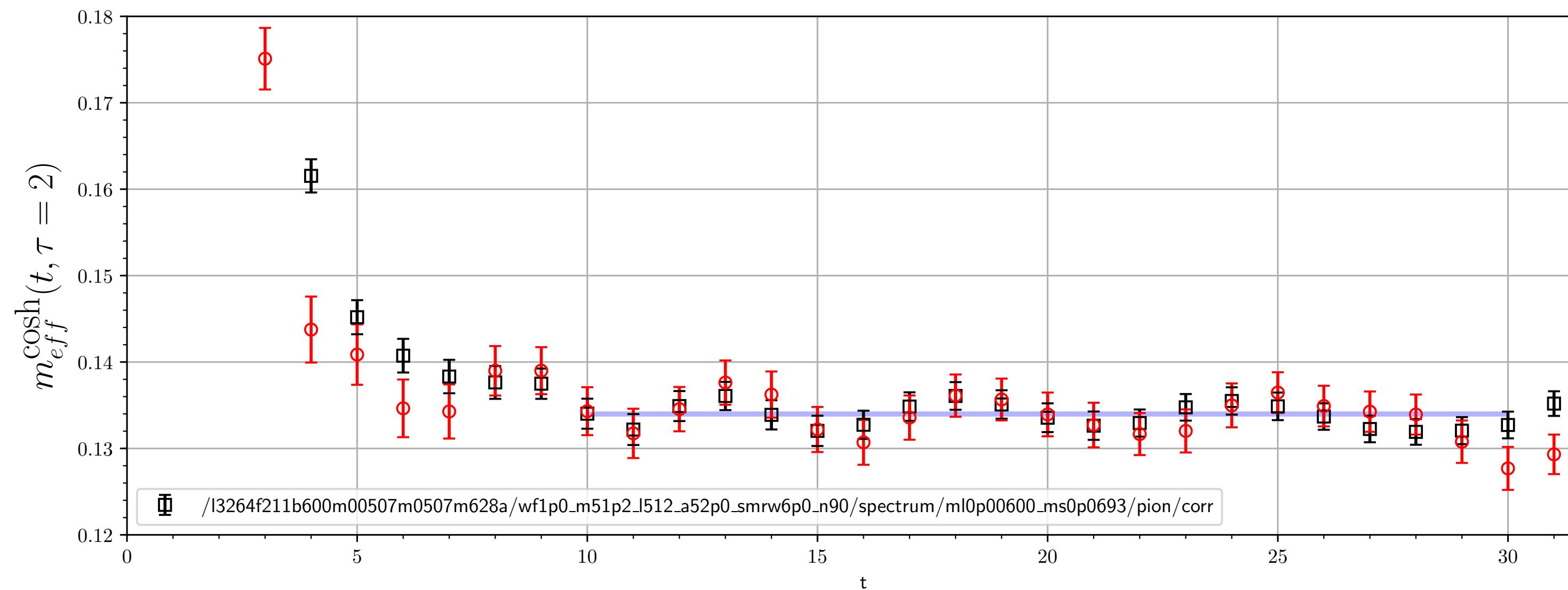
$$C(t) = \sum_n z_n z_n^\dagger e^{-E_n t}$$

$$m_{eff}(t) = \frac{1}{\tau} \ln \left(\frac{C(t)}{C(t + \tau)} \right)$$

For pions, need to consider leading finite temperature effects

$$C(t) = \sum_n z_n z_n^\dagger \left(e^{-E_n t} + e^{-E_n(T-t)} \right)$$

$$m_{eff}^{\cosh}(t, \tau) = \frac{1}{\tau} \cosh^{-1} \left(\frac{C(t + \tau) + C(t - \tau)}{2C(t)} \right)$$

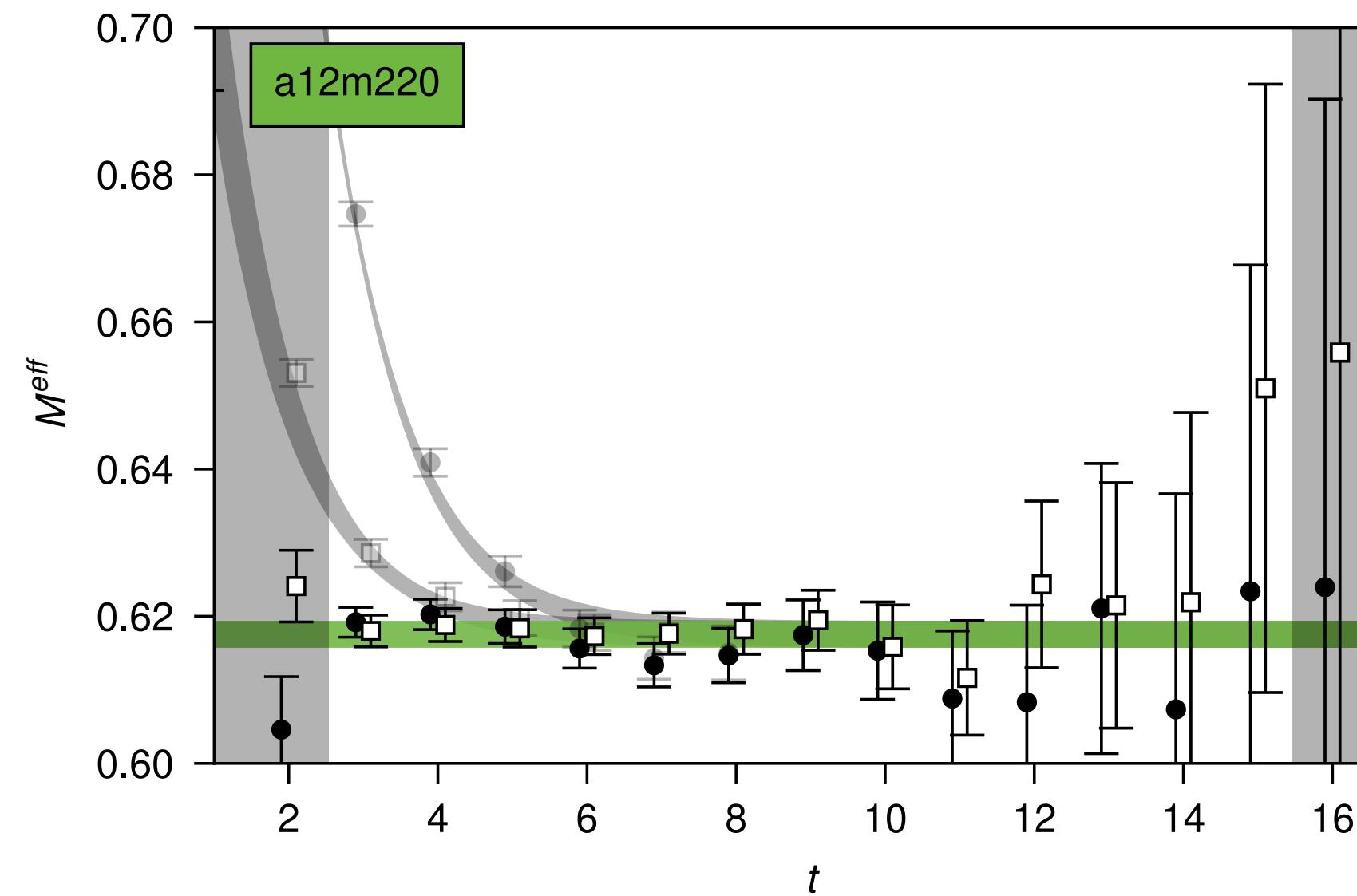


Effective mass of Pion 2-point correlation function
 red and black “data” are from different choices of *interpolating* operators

Noise is constant in time - can determine very clean ground state ([blue band](#))

LQCD challenges for NP

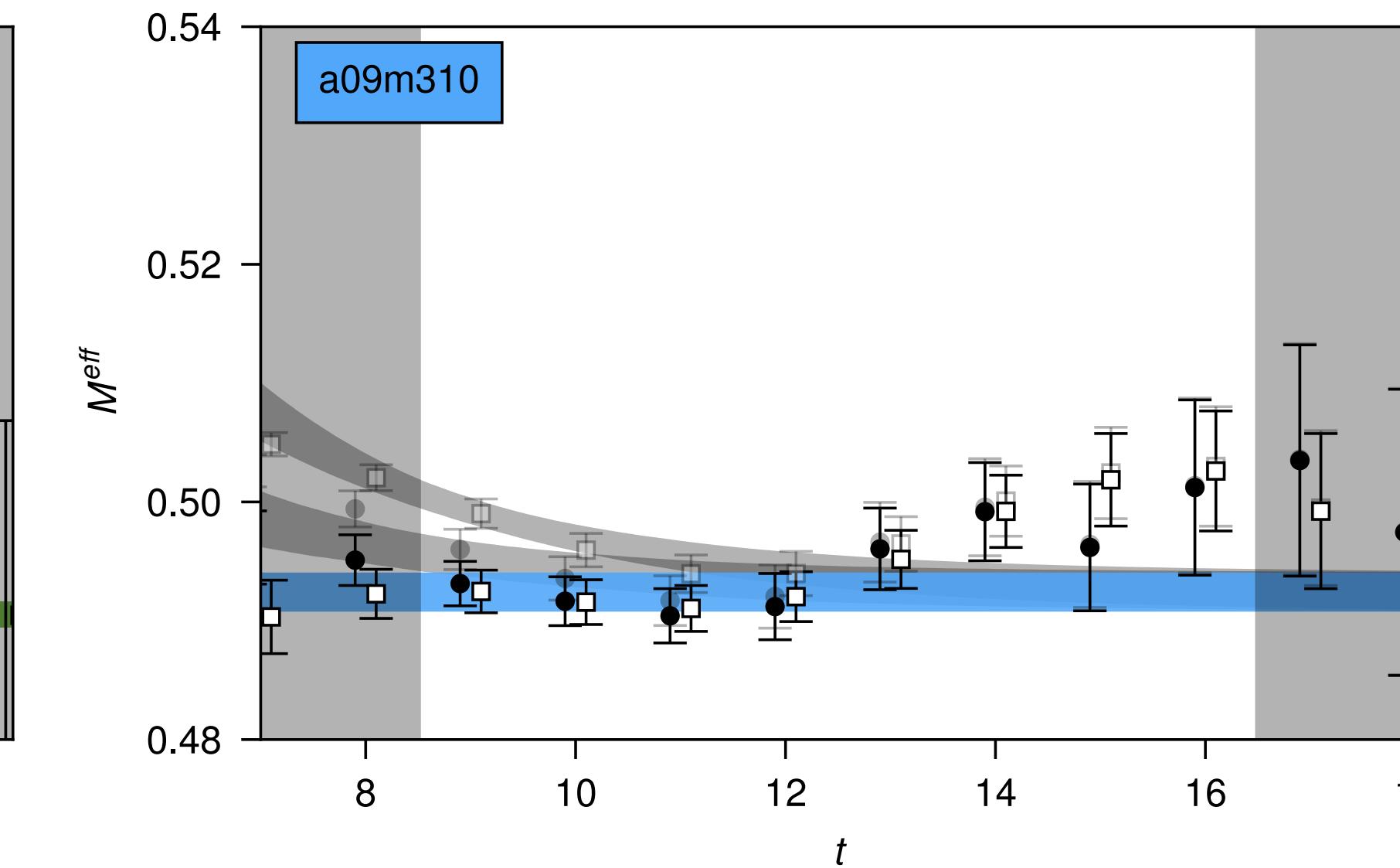
2-point correlation function



Two examples of nucleon effective mass

Noise is growing in time - can not simply go to the long-time limit without exponentially increasing the amount of statistics needed

$$\frac{\text{Signal}}{\text{Noise}} \rightarrow \sqrt{N_{\text{state}}} e^{-(m_N - \frac{3}{2} m_\pi)t}$$



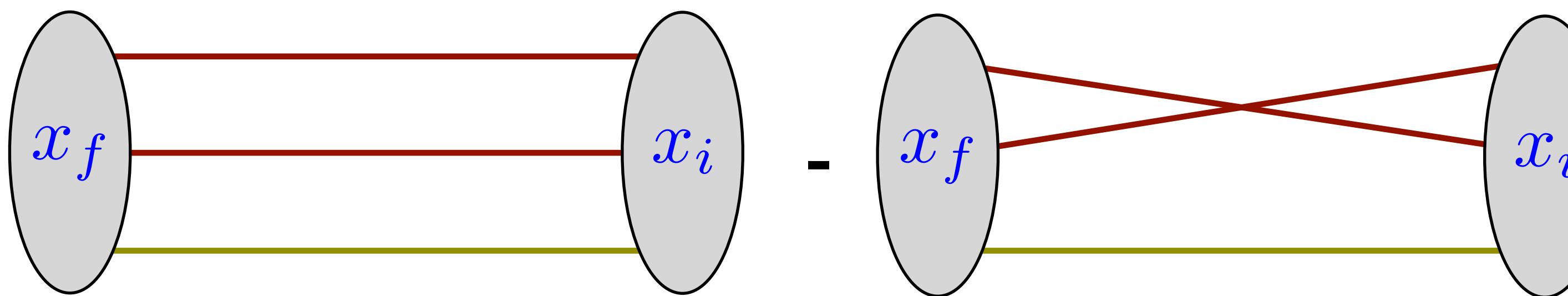
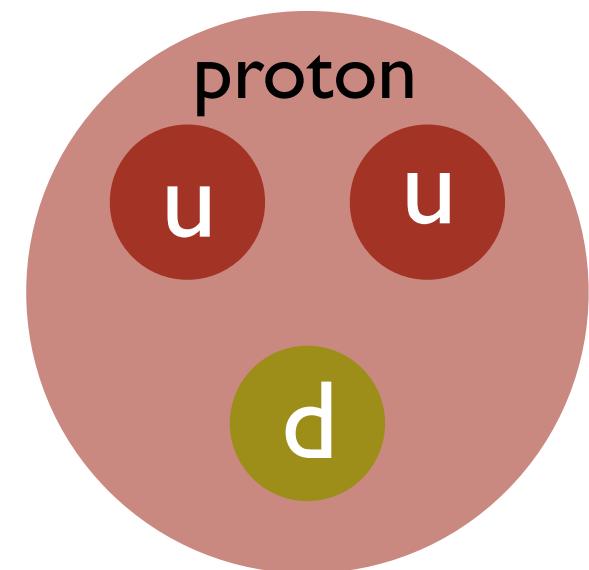
Correlated late-time fluctuations... what is the ground state?

Need sophisticated analysis to ensure you are not susceptible to correlated fluctuations

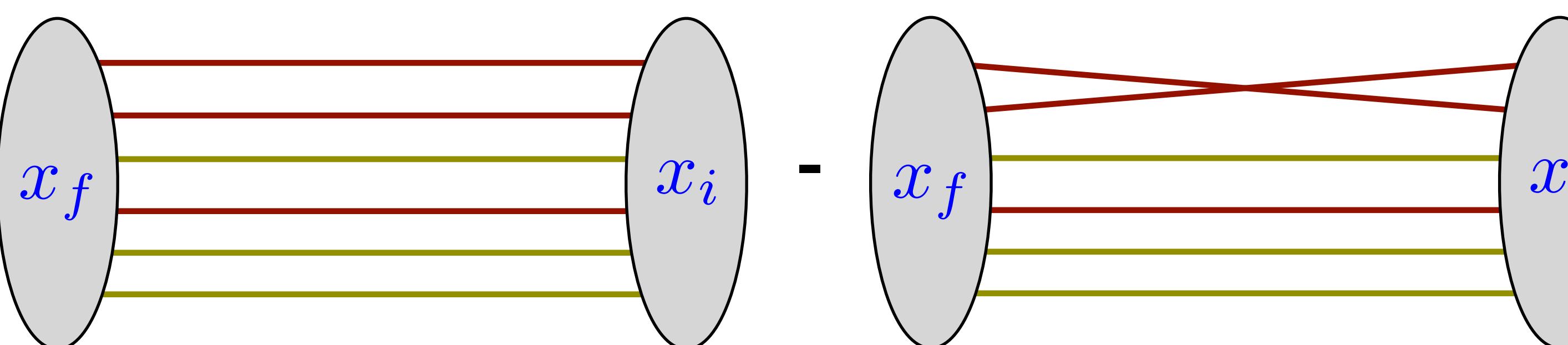
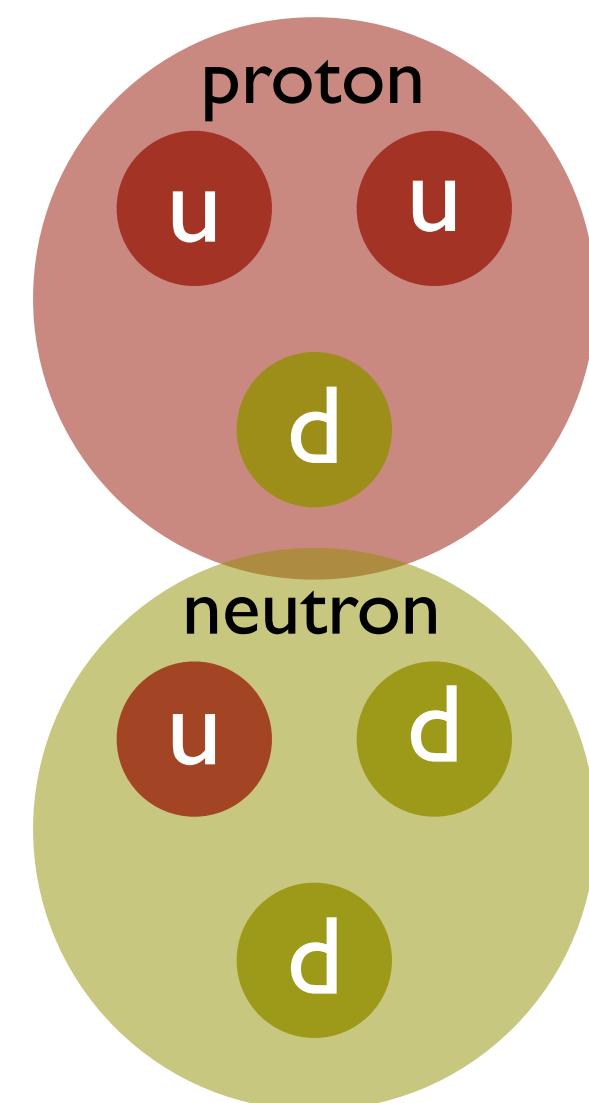
This problem is exacerbated with form-factor calculations (g_A) and 2+ nucleons

- quark contraction cost becomes dominant
- density of excited states grows significantly and gap becomes small (nuclear interaction energies instead of pion mass gap)

LQCD challenges for NP

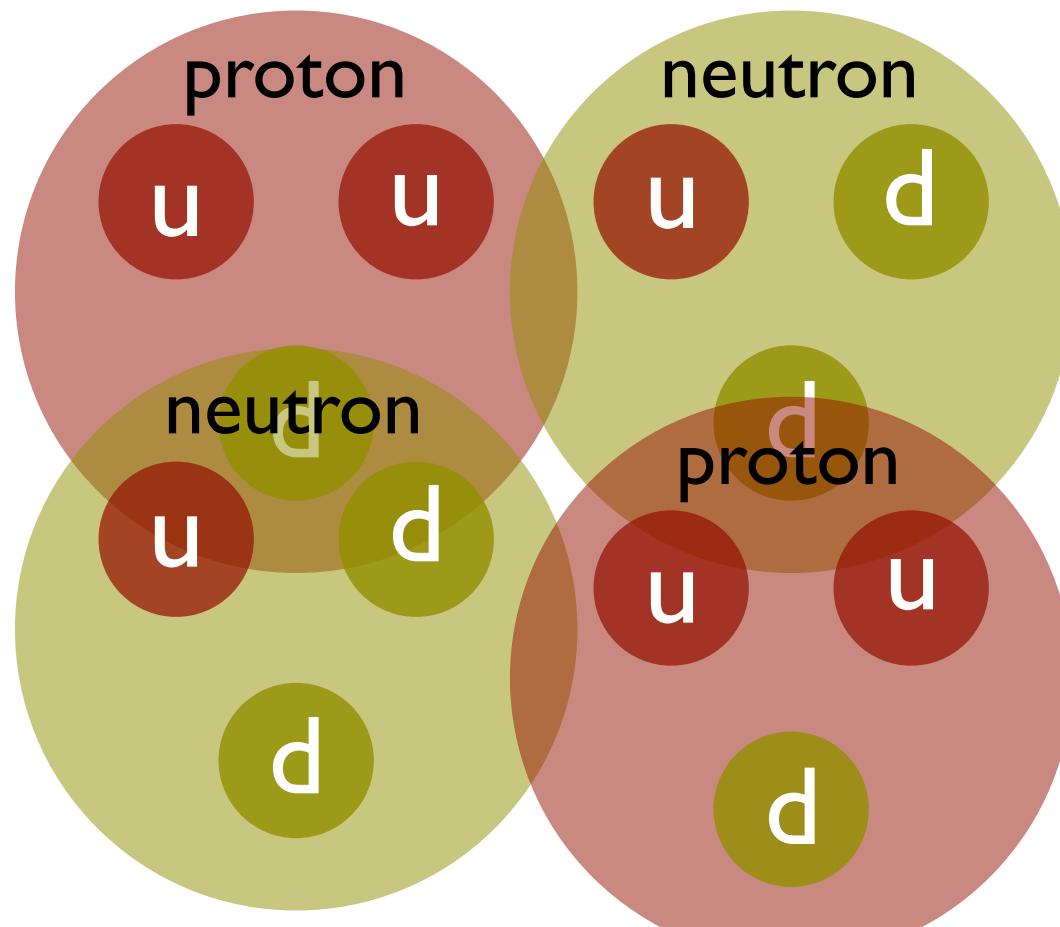


$$2! \times 1! = 2 \text{ contractions}$$



$$3! \times 3! = 36 \text{ contractions} + \dots$$

LQCD challenges for NP



$6! \times 6! = 518400$
contractions
numerical cost exceeds
HMC + props

quark-exchange diagrams are source of fermion sign problem
expensive AND noisy :(

There are clever solutions that reduce this cost significantly

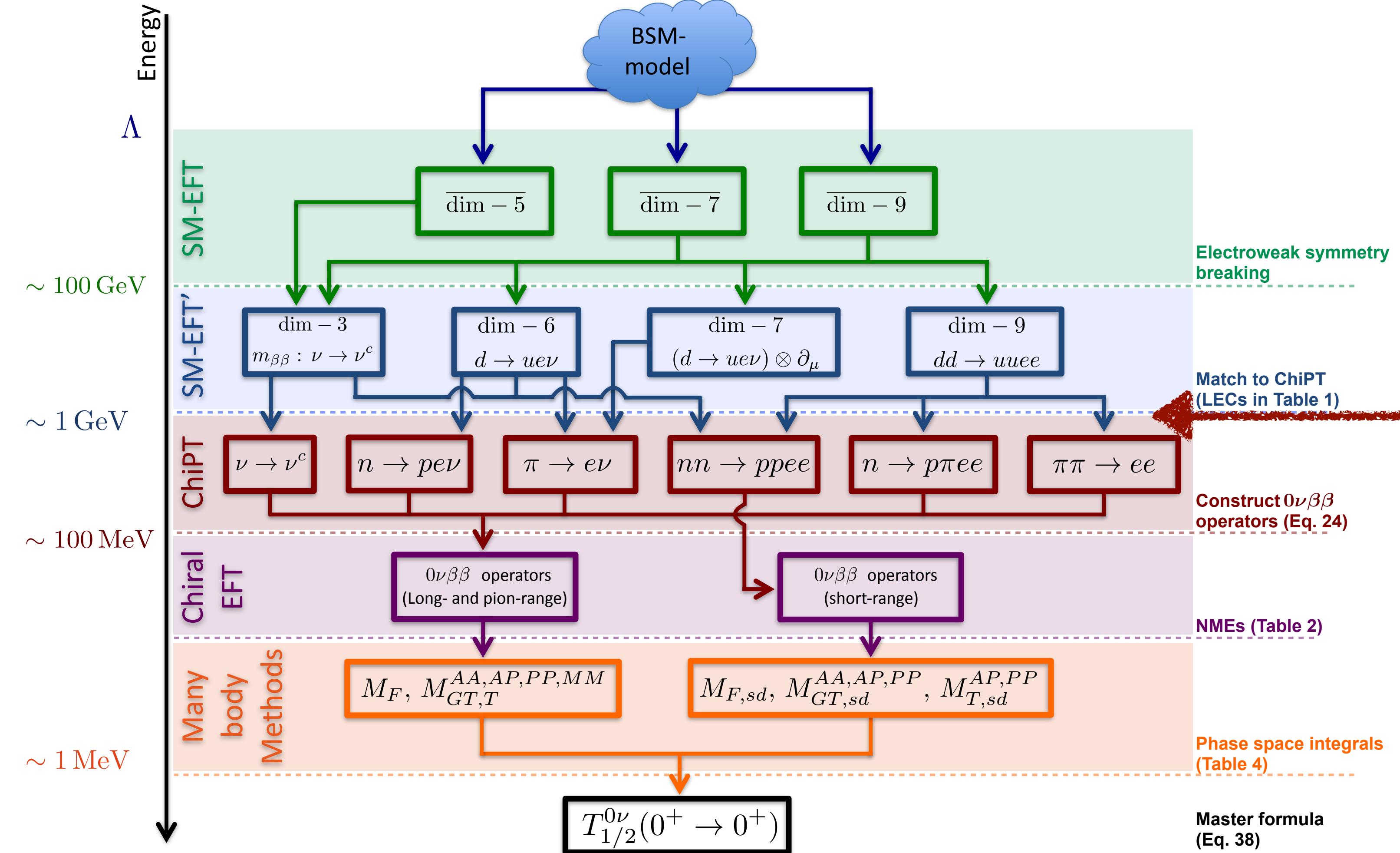
Yamazaki, Kuramashi, Ukawa arXiv:0912.1383,

Doi, Endres arXiv:1205.0585, Detmold, Orginos arXiv:1207.1452

Günther, Toth, Varnhorst arXiv:1301.4895

But unfortunately, they only work with unrealistically simplistic interpolating operators in which all quarks originate from the same spacetime point, or are otherwise, identical

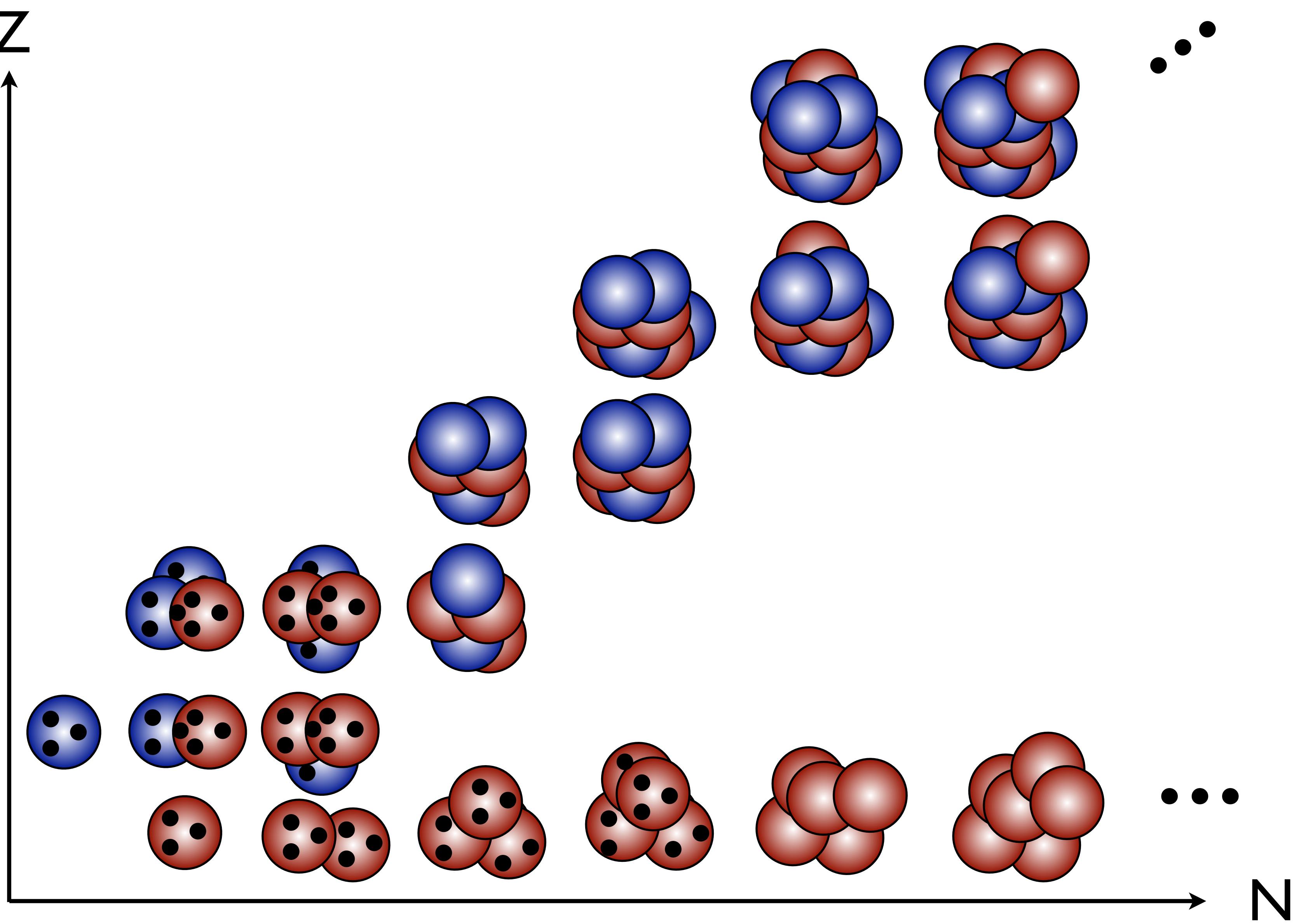
LQCD and EFT

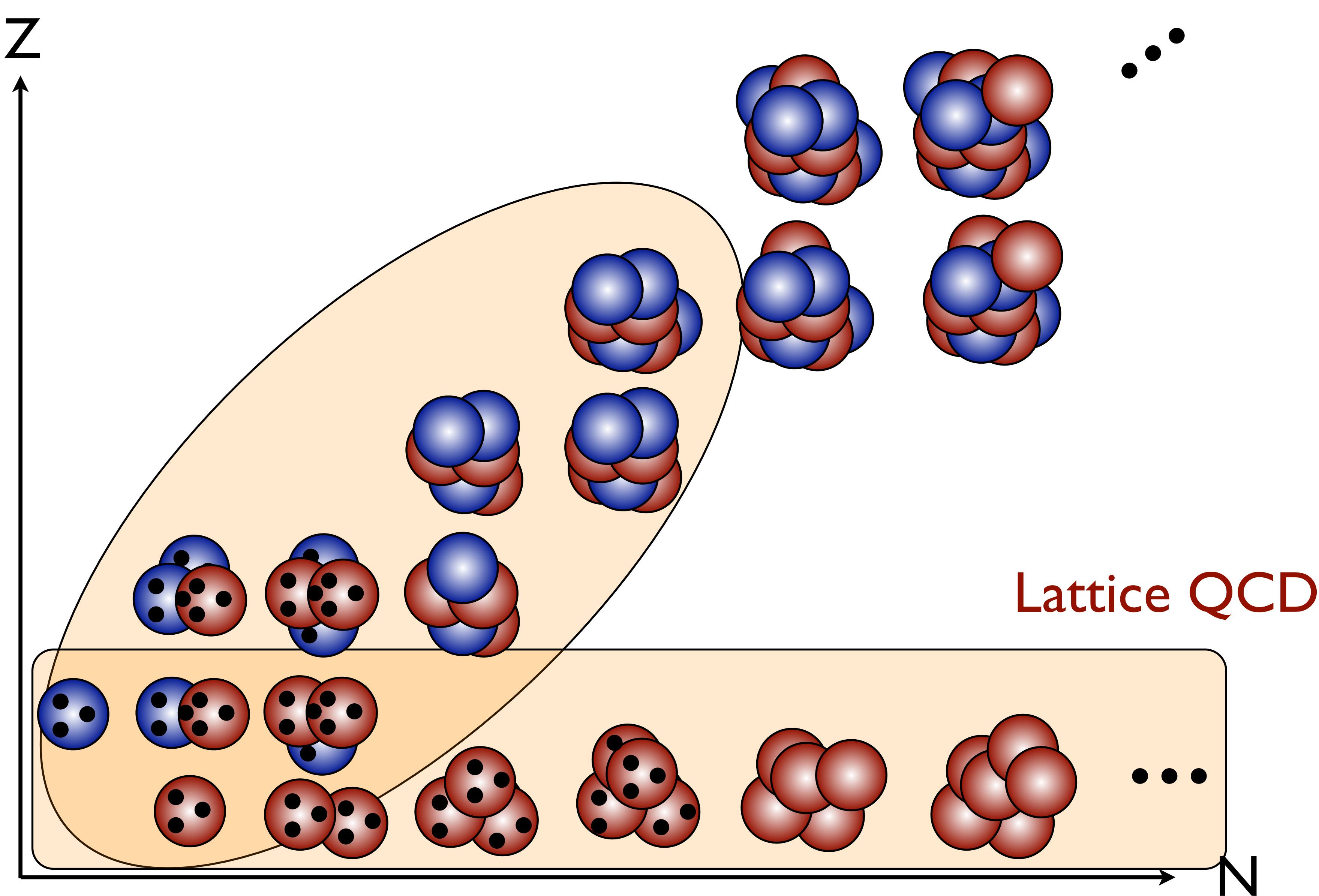


Example from $0\nu\beta\beta$

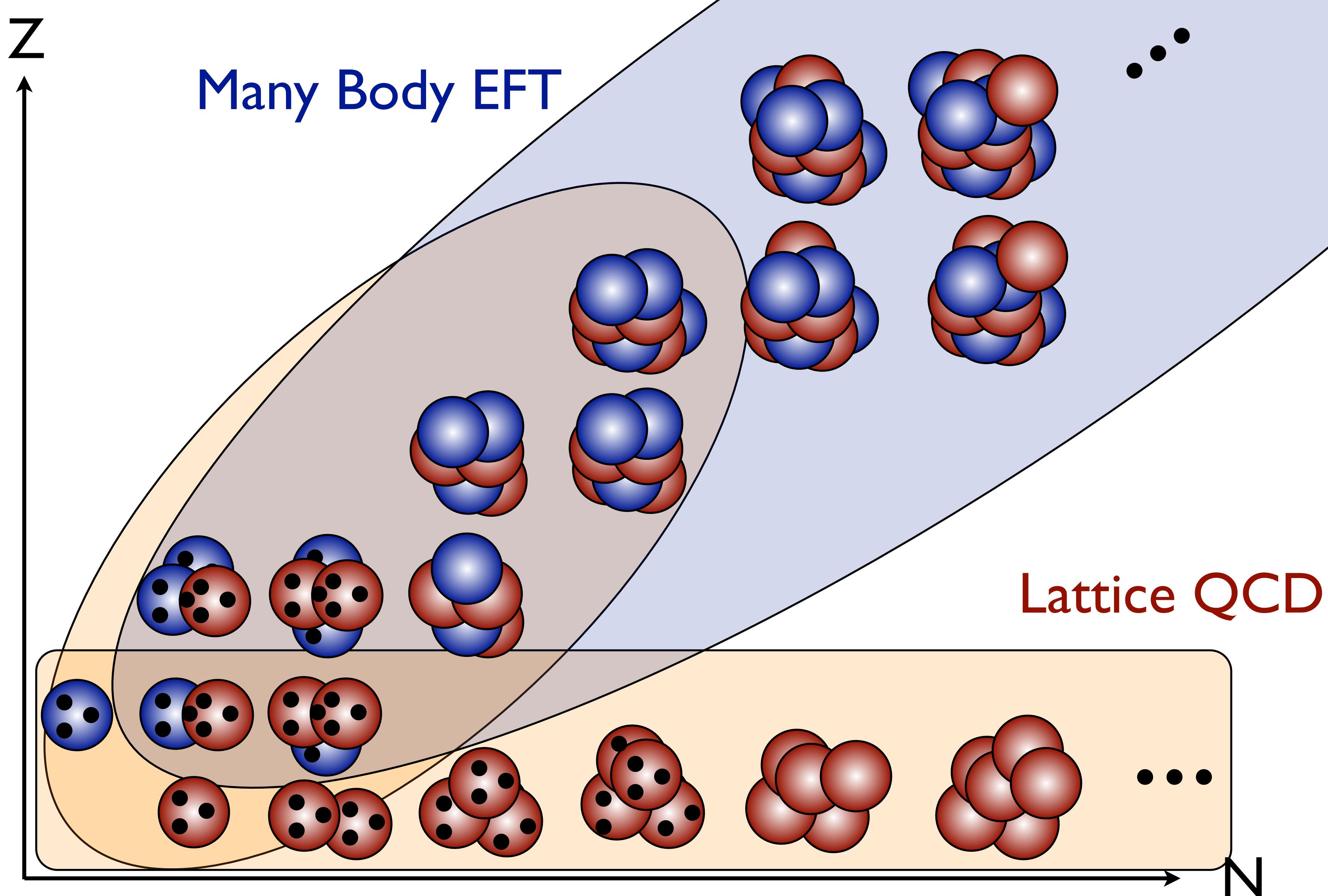
Cirigliano, Dekens, de Vries, Graesser, Mereghetti, arXiv:1806.02780

- LQCD is used at the interface between QCD-level operators and Chiral Perturbation Theory
- Combining LQCD + EFT greatly extends the reach of application and prediction





Determine 2, 3, 4 body forces directly from QCD



Determine 2, 3, 4 body forces directly from QCD
match onto many body effective field theory

Chiral Perturbation Theory Overview

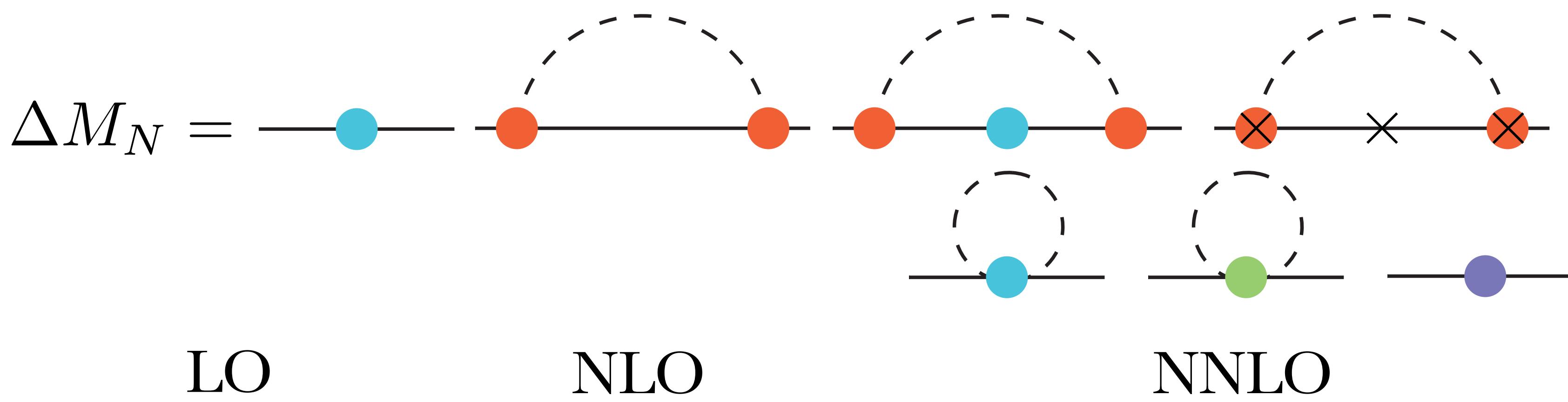
Jenkins & Manohar Phys. Lett. B 255 (1991)

○ HB χ PT for M_N up to NNLO

AWL arXiv:hep-lat/0405007

B.C. Tiburzi, AWL arXiv:hep-lat/0501018

$$\begin{aligned} \mathcal{L}_{N\pi} = & \bar{N} i v \cdot D N - \frac{\alpha_N}{4\pi F} \bar{N} \mathcal{M}_+ N - \frac{\sigma_N}{4\pi F} \bar{N} N \text{tr}(\mathcal{M}_+) + 2g_A^{\text{LO}} \bar{N} S \cdot \mathcal{A} N \\ & - \bar{N} \frac{D_\perp^2}{2M_0} N + g_A^{\text{LO}} \left(\bar{N} \frac{i \overleftrightarrow{D} \cdot S}{M_0} v \cdot \mathcal{A} N + \bar{N} v \cdot \mathcal{A} \frac{S \cdot i \overrightarrow{D}}{M_0} N \right) + \frac{b_A}{4\pi F} \bar{N} N \text{tr}(\mathcal{A} \cdot \mathcal{A}) + \frac{b_{vA}}{4\pi F} \bar{N} N \text{tr}(v \cdot \mathcal{A} v \cdot \mathcal{A}) \\ & + \frac{b_1}{(4\pi F)^3} \bar{N} \mathcal{M}_+^2 N + \frac{b_2}{(4\pi F)^3} \bar{N} N \text{tr}(\mathcal{M}_+^2) + \frac{b_3}{(4\pi F)^3} \bar{N} \mathcal{M}_+ N \text{tr}(\mathcal{M}_+) + \frac{b_4}{(4\pi F)^3} \bar{N} N [\text{tr}(\mathcal{M}_+)]^2 \end{aligned}$$

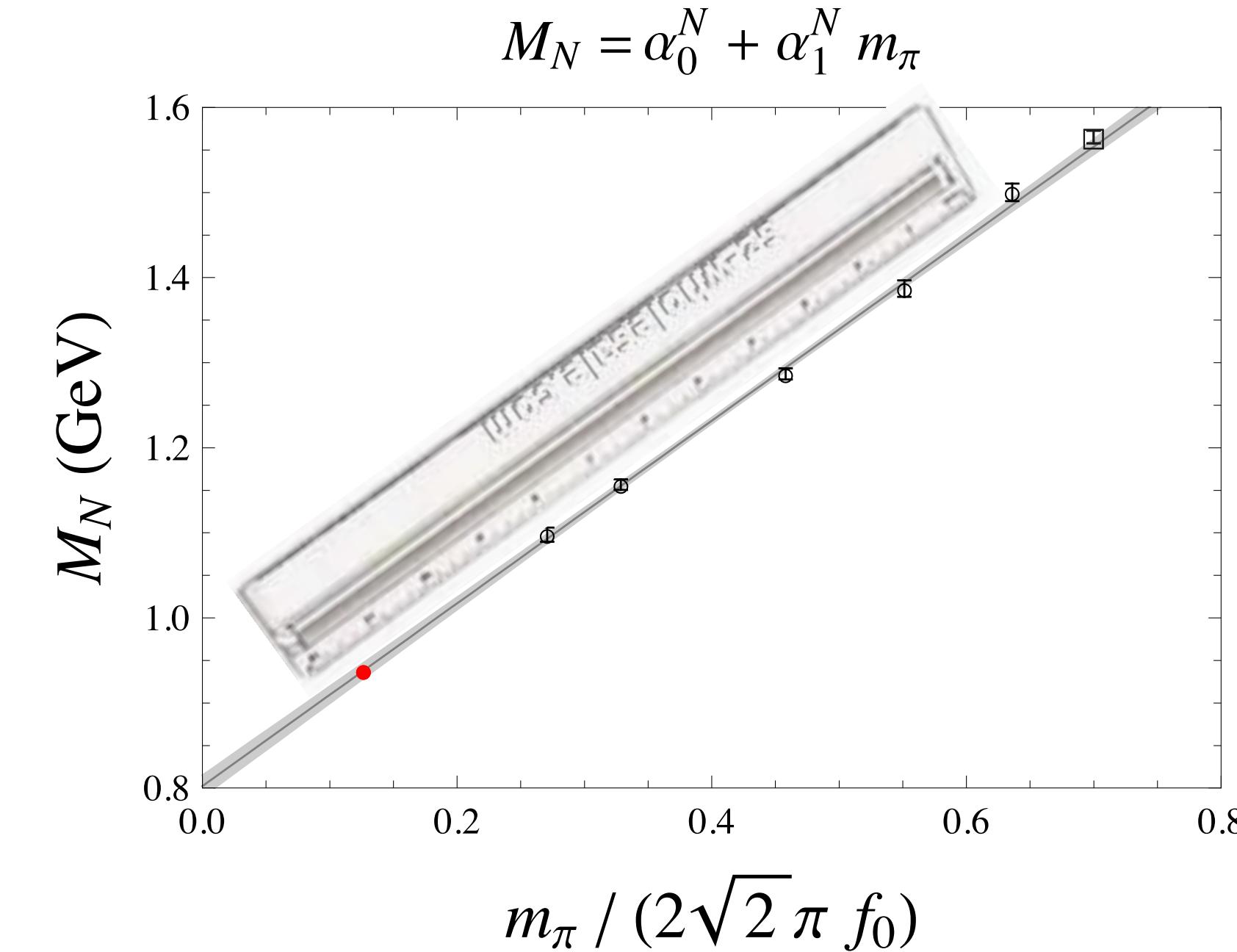
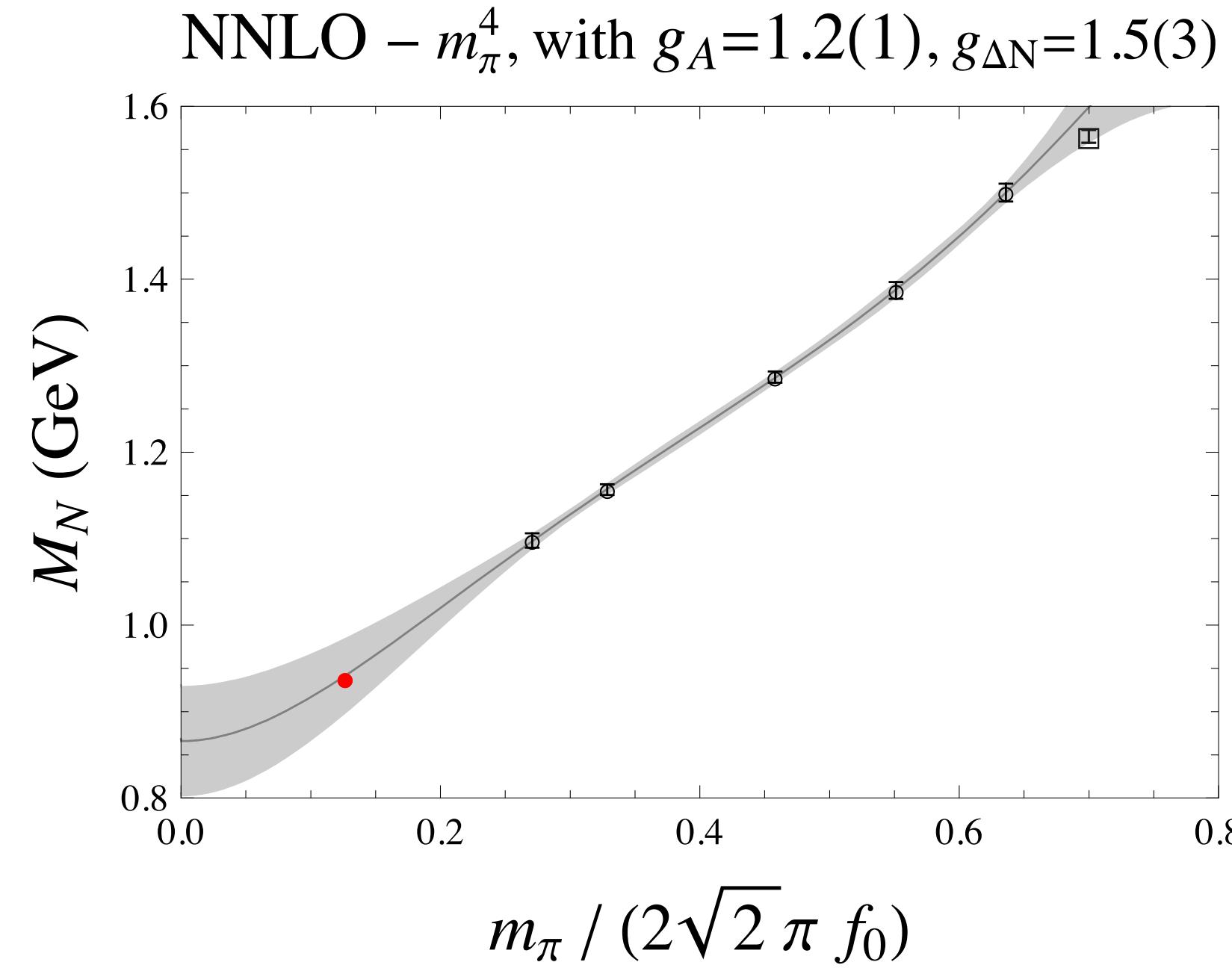


$$\frac{m_\pi^2}{4\pi F_\pi}$$

$$-\frac{3\pi g_A^2}{2} \frac{m_\pi^3}{(4\pi F_\pi)^2}$$

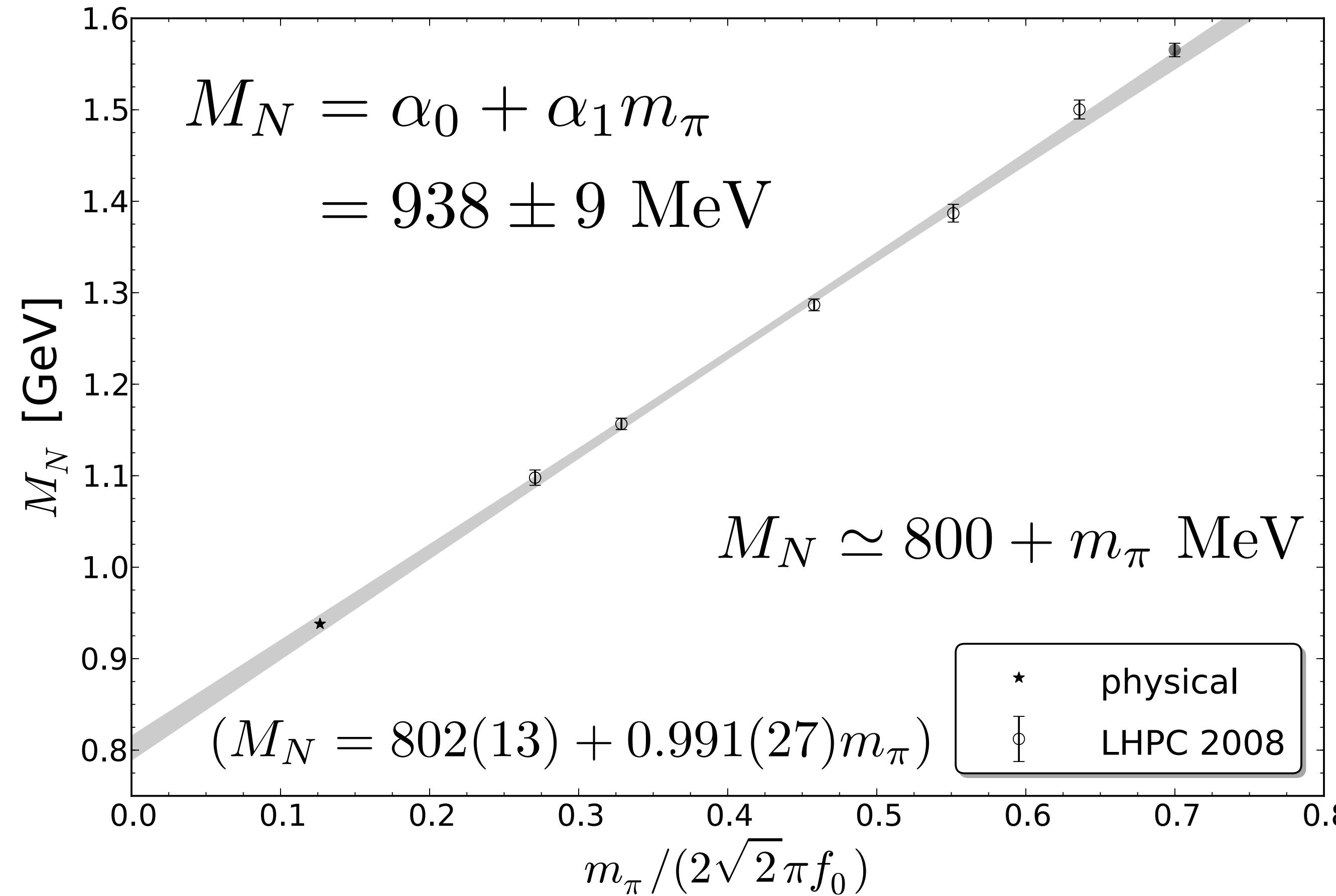
$$\frac{m_\pi^4}{(4\pi F_\pi)^3} \ln \left(\frac{m_\pi^2}{(4\pi F_\pi)^2} \right)$$

$$\frac{m_\pi^4}{(4\pi F_\pi)^3}$$



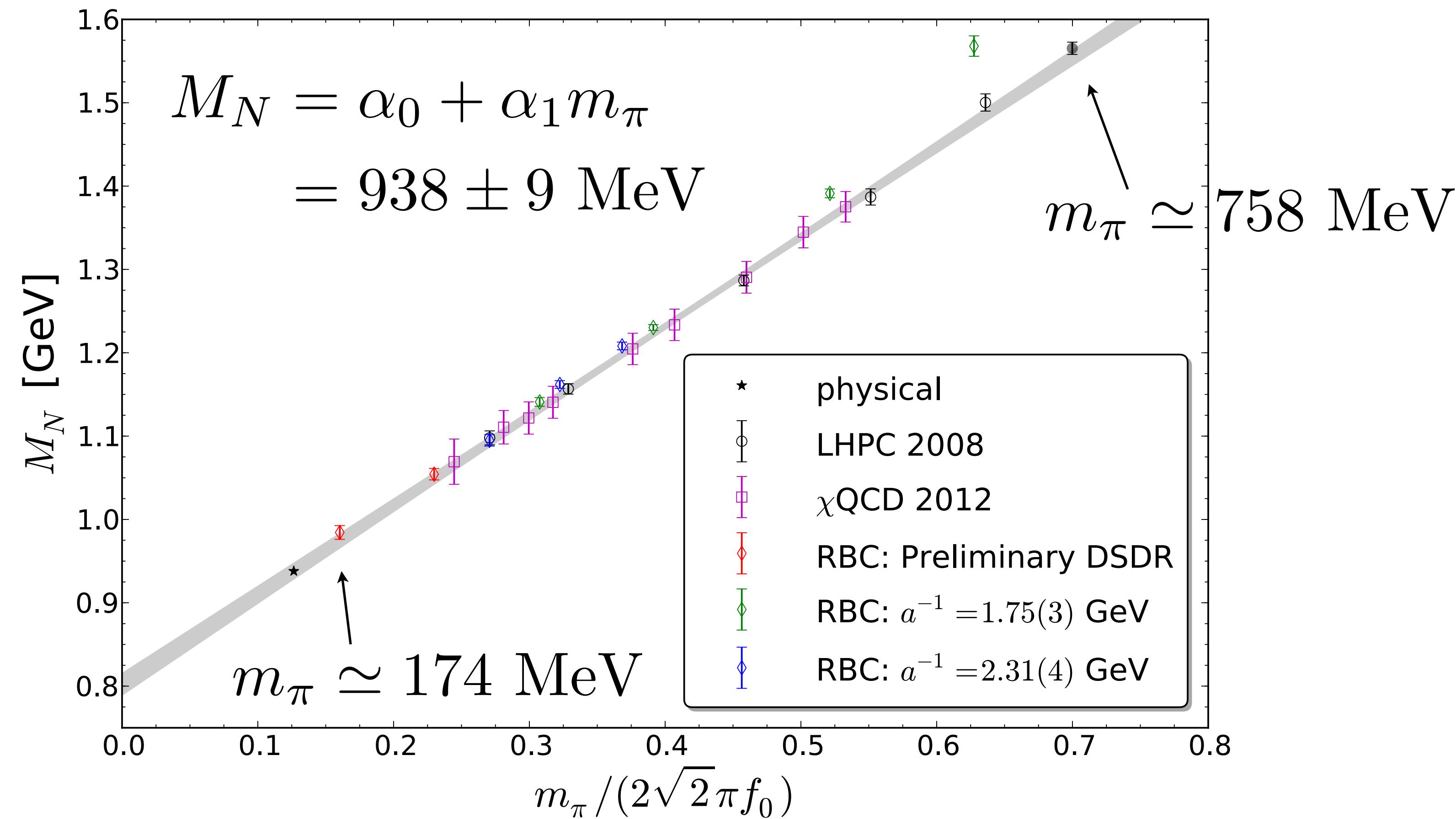
- Is this a lattice artifact? Or a conspiracy of QCD?
- All lattice calculations in 2008 with 2+1 dynamical fermion showed this linear pion mass dependence

What is the status now (2012)?



Physical point NOT included in fit

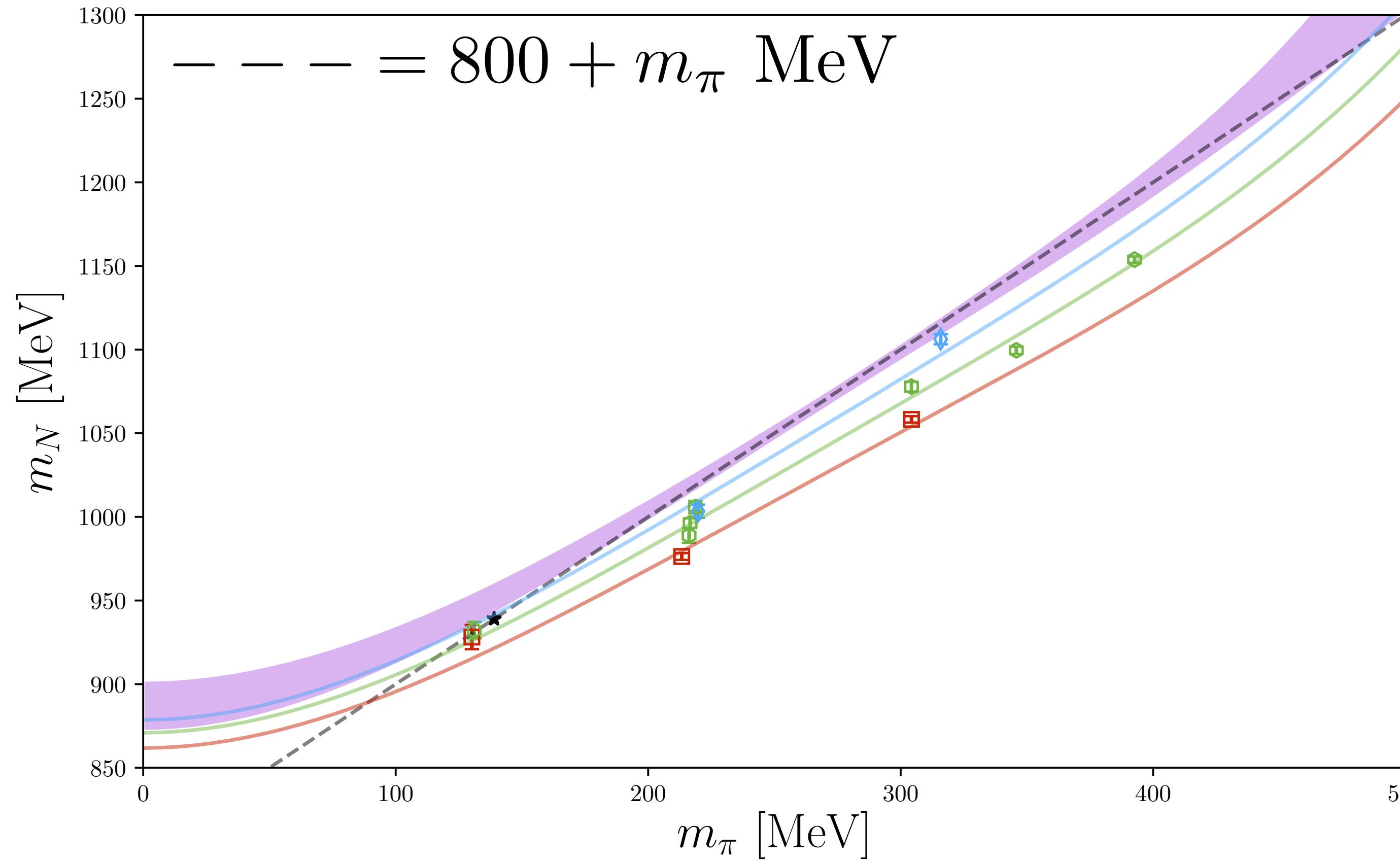
What is the status now (2012)?



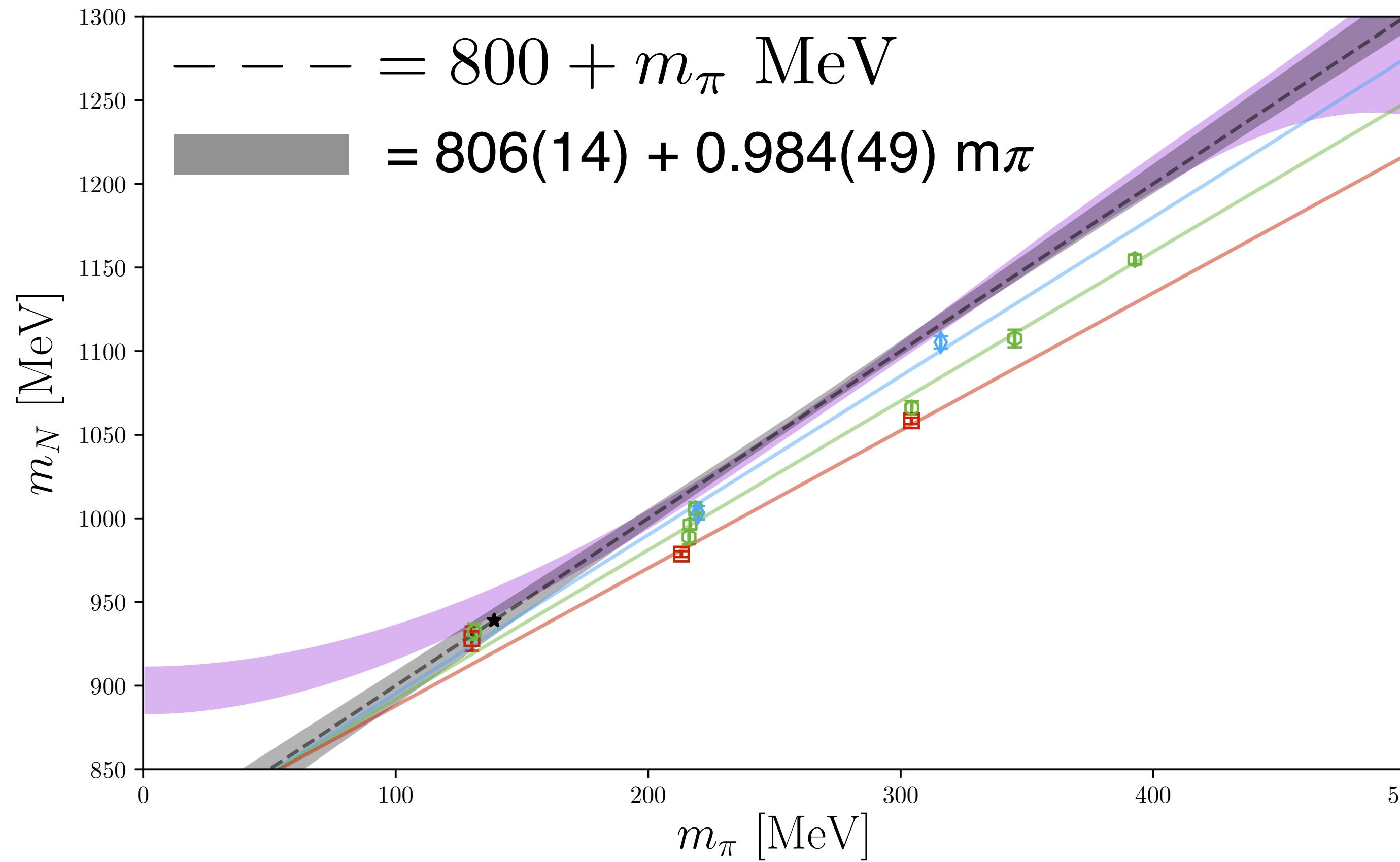
Taking this seriously yields

$$\sigma_{\pi N} = 67 \pm 4 \text{ MeV}$$

$$M_N \simeq 800 + m_\pi \text{ MeV}$$
$$(M_N = 802(13) + 0.991(27)m_\pi)$$



- NNLO fit $O(m_\pi^4) + a^2/\omega_0^2 + a^2/\omega_0^2 \times m_\pi^2/(4\pi F)$
- gA taken from our calculation
- continuum, infinite volume extrapolated mass is consistent with RULER



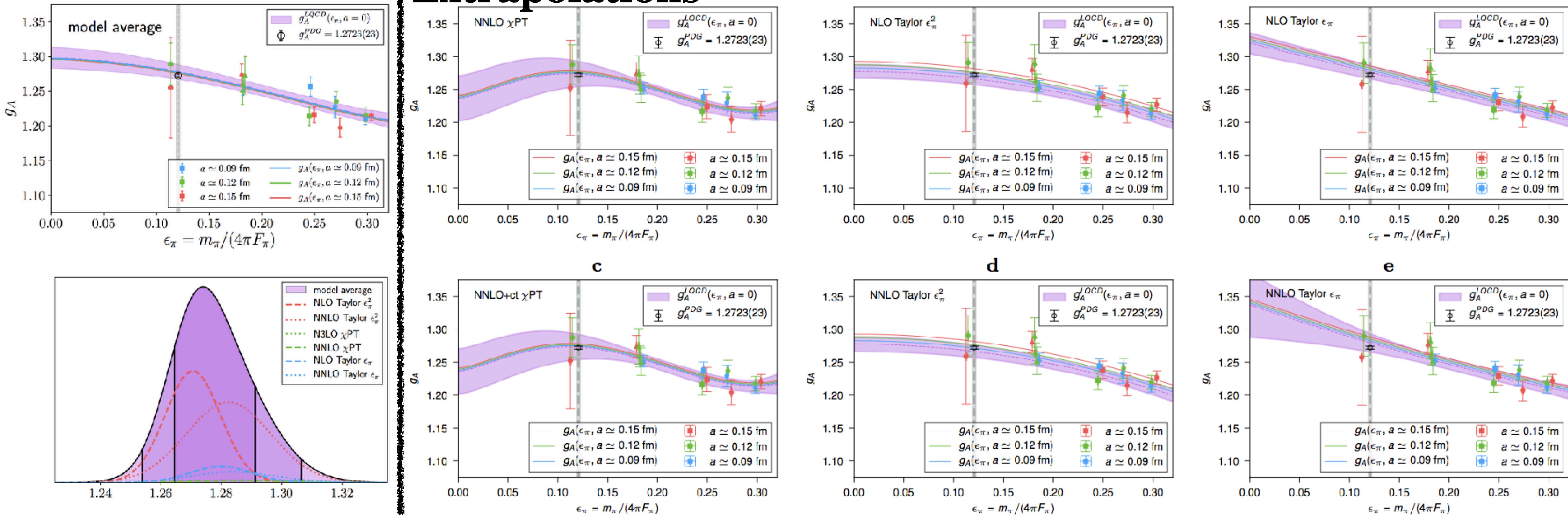
- Ruler applied to our new MDWF on HISQ results
- $M_N = c_0 + c_1 m_\pi + a_2 a^2 + c_{am,2} a^2 m_\pi$

Analysis Details

ga example

Chang et al. arXiv:1805.12130

Extrapolations

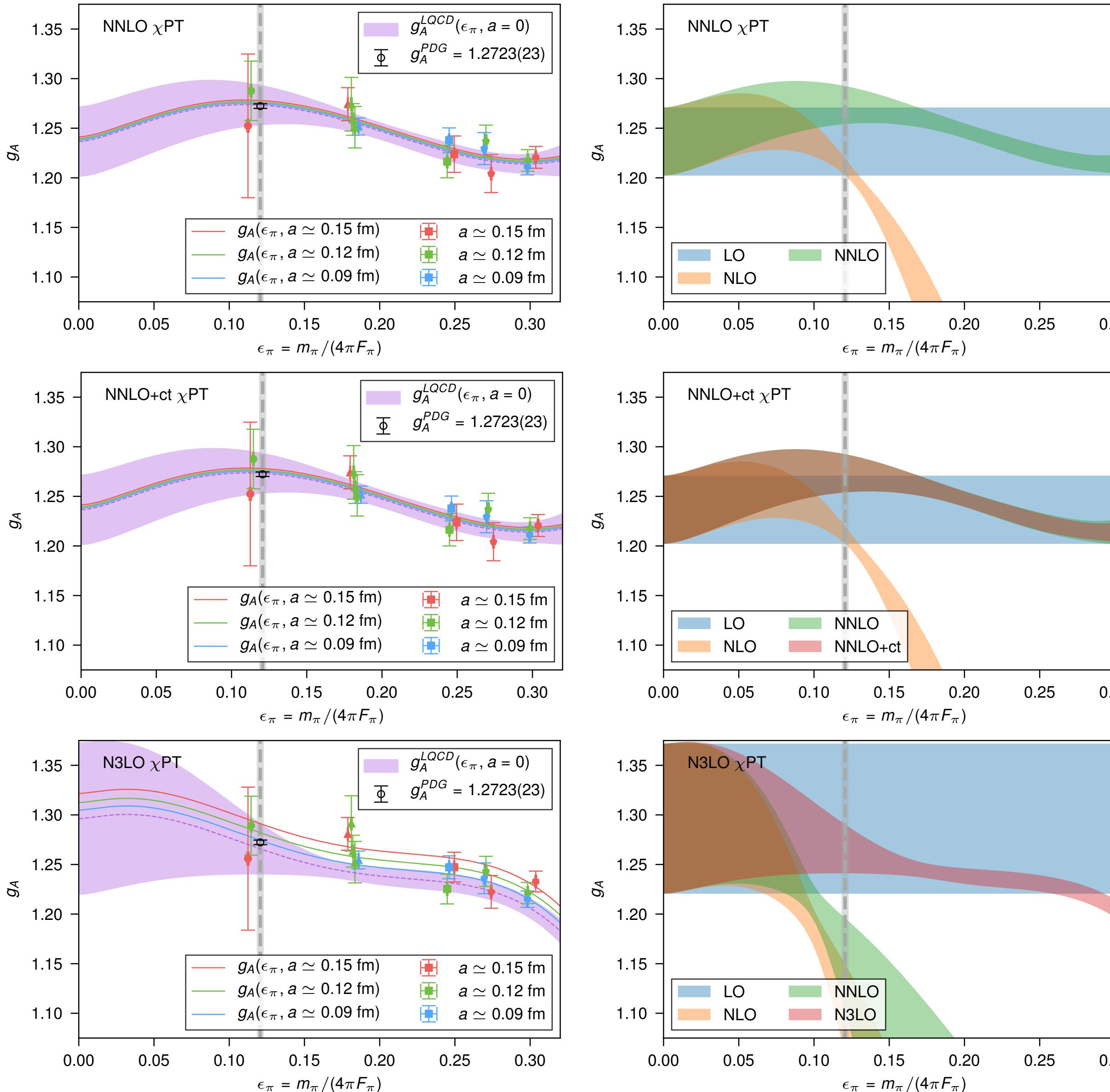


| | Fit | χ^2/dof | $\mathcal{L}(D M_k)$ | $P(M_k D)$ | $P(g_A M_k)$ |
|------------------------------|-------|---------------------|----------------------|---------------|--------------|
| NNLO χ PT | 0.727 | 22.734 | 0.033 | 1.273(19) | |
| NNLO+ct χ PT | 0.726 | 22.729 | 0.033 | 1.273(19) | |
| NLO Taylor ϵ_π^2 | 0.792 | 24.887 | 0.287 | 1.266(09) | |
| NNLO Taylor ϵ_π^2 | 0.787 | 24.897 | 0.284 | 1.267(10) | |
| NLO Taylor ϵ_π | 0.700 | 24.855 | 0.191 | 1.276(10) | |
| NNLO Taylor ϵ_π | 0.674 | 24.848 | 0.172 | 1.280(14) | |
| average | | | | 1.271(11)(06) | |

$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3$$

NNLO χ PT : Eq. (S8) + $\delta_a + \delta_L$
 NNLO+ct χ PT : Eq. (S8) + $c_4 \epsilon_\pi^4 + \delta_a + \delta_L$
 NLO Taylor ϵ_π^2 : $c_0 + c_2 \epsilon_\pi^2 + \delta_a + \delta_L$
 NNLO Taylor ϵ_π^2 : $c_0 + c_2 \epsilon_\pi^2 + c_4 \epsilon_\pi^4 + \delta_a + \delta_L$
 NLO Taylor ϵ_π : $c_0 + c_1 \epsilon_\pi + \delta_a + \delta_L$
 NNLO Taylor ϵ_π : $c_0 + c_1 \epsilon_\pi + c_2 \epsilon_\pi^2 + \delta_a + \delta_L$

convergence of the chiral expansion...



$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) \\ + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3$$

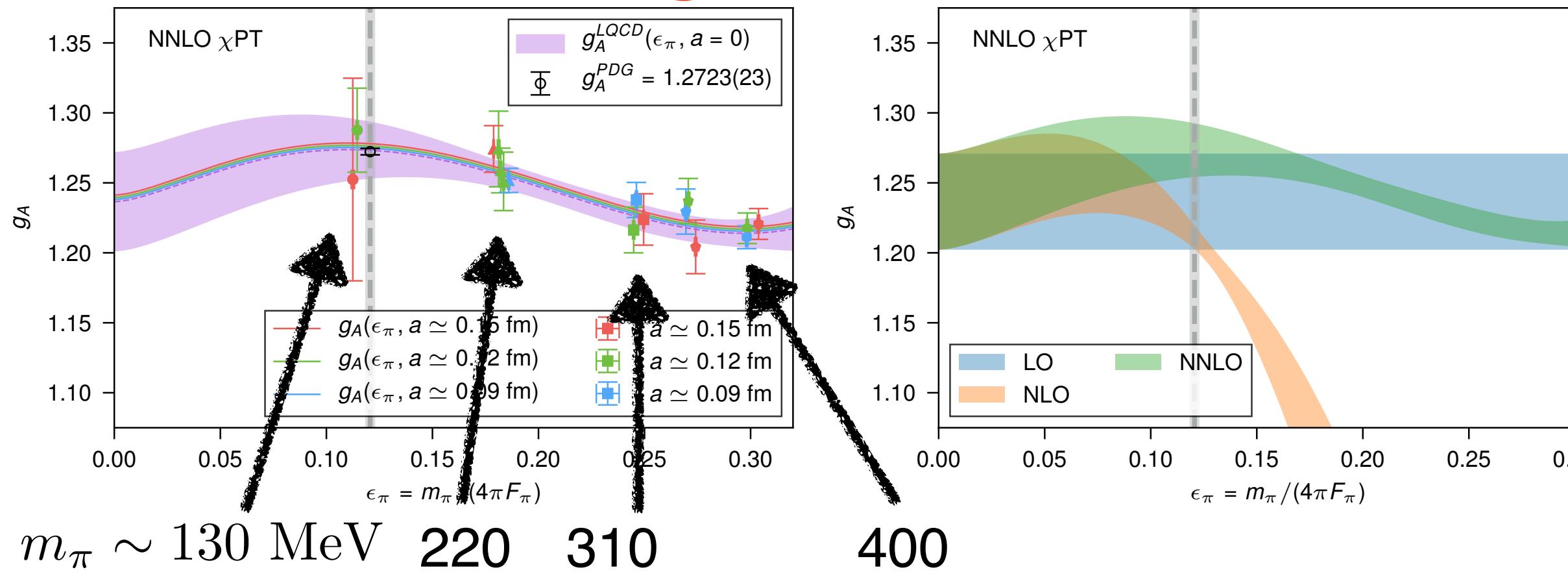
$$\epsilon_\pi = \frac{m_\pi}{4\pi F_\pi}$$

$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) \\ + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3 + c_4 \epsilon_\pi^4$$

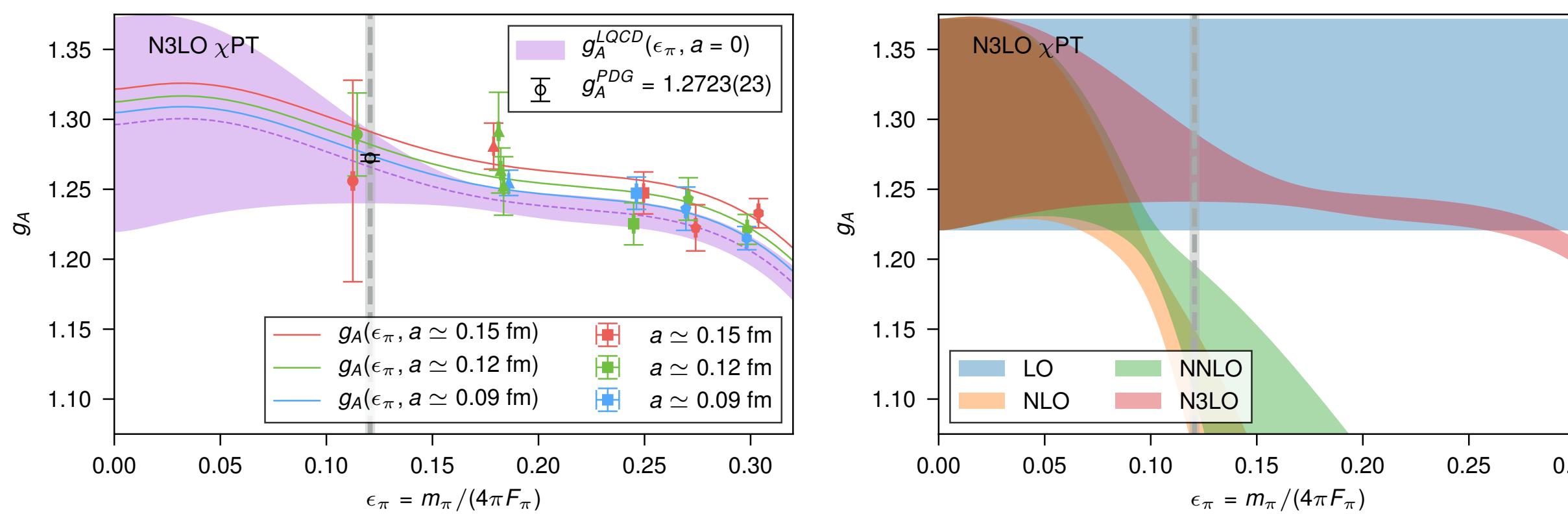
$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) \\ + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3 \\ + \epsilon_\pi^4 \left[c_4 + \tilde{\gamma}_4 \ln(\epsilon_\pi^2) \right] \\ + \left(\frac{2}{3} g_0 + \frac{37}{12} g_0^3 + 4 g_0^5 \right) \ln^2(\epsilon_\pi^2)$$

Bernard and Meissner (CD06)
 Phys.Lett.B639 [hep-lat/0605010]
 $F \longrightarrow F_\pi$

convergence of the chiral expansion...



can we trust extrapolation of quantities
with chirality-enhanced behavior?
if the single nucleon is not converging, would you
trust chiral extrapolations of two or more nucleons?



$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) \\ + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3$$

$$\epsilon_\pi = \frac{m_\pi}{4\pi F_\pi}$$

$$g_A = g_0 - \epsilon_\pi^2(g_0 + 2g_0^3) \ln(\epsilon_\pi^2) \\ + c_2 \epsilon_\pi^2 + g_0 c_3 \epsilon_\pi^3 \\ + \epsilon_\pi^4 \left[c_4 + \tilde{\gamma}_4 \ln(\epsilon_\pi^2) \right] \\ + \left(\frac{2}{3} g_0 + \frac{37}{12} g_0^3 + 4 g_0^5 \right) \ln^2(\epsilon_\pi^2)$$

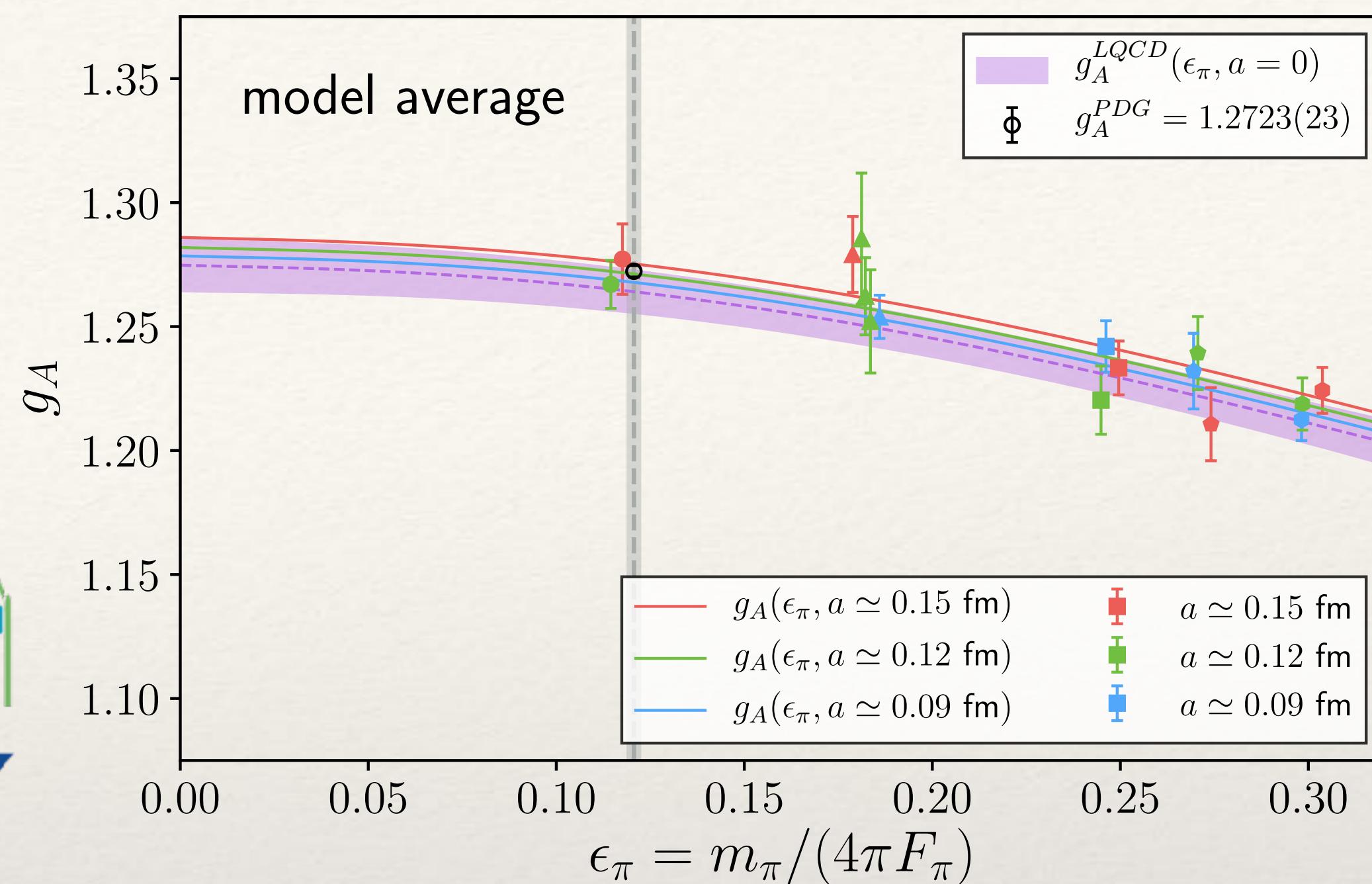
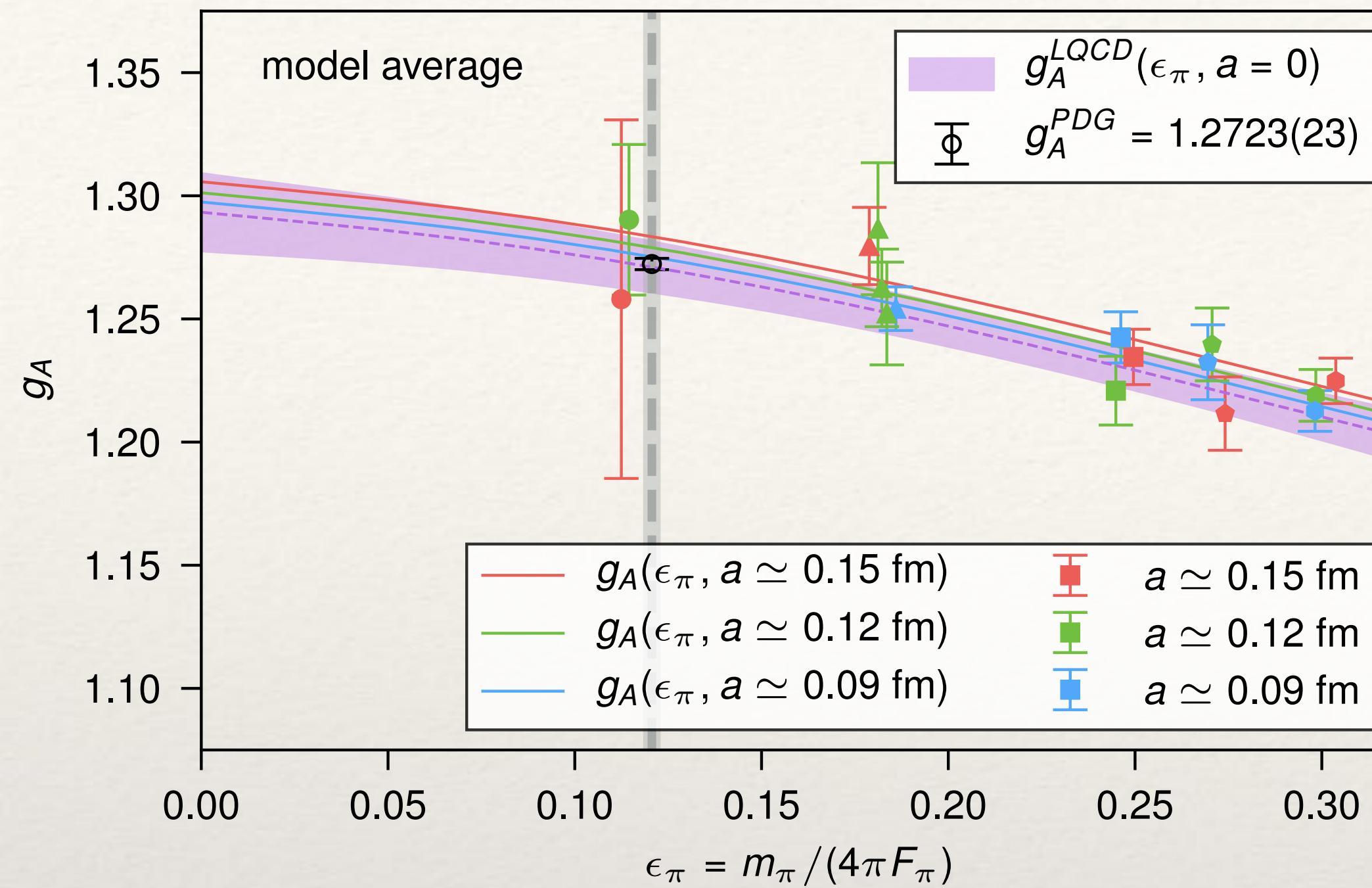
Bernard and Meissner (CD06)
 Phys.Lett.B639 [hep-lat/0605010]
 $F \longrightarrow F_\pi$

Assignment

- What is causing all the trouble for SU(2) baryon chiral perturbation theory?
- Large N_c arguments show a cancellation between nucleon and delta loops for g_A ,
BUT, they add “coherently” for M_N

- What is the range of convergence of SU(2) baryon chiral perturbation theory?

Why is this an exciting time for LQCD and NP?

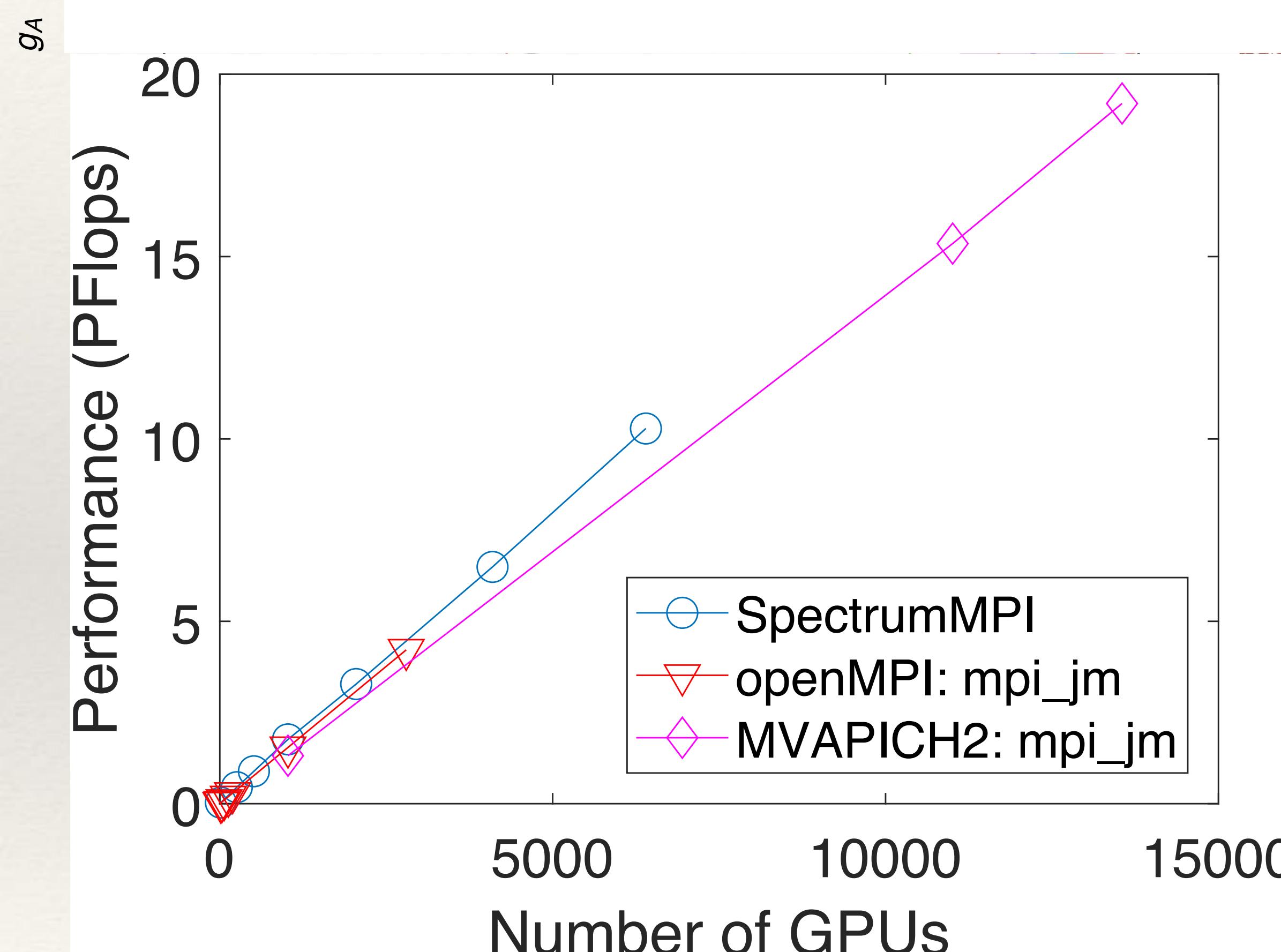


- The **a12m130** ($48^3 \times 64 \times 20$) with 3 sources cost as much as all other ensembles combined
- **2.5 weekends on Sierra** \rightarrow 16 srcs
- Now, 32 srcs (un-constrained, 3-state fit)
- We generated a new **a15m135XL** ($48^3 \times 64$) ensemble (old **a15m130** is $32^3 \times 48$)
 - $M\pi L = 4.93$ (old $M\pi L = 3.2$)
 - $L_5 = 24$, $N_{src} = 16$
- We are running $g_A(Q^2)$ on Summit this year (DOE INCITE)
- We anticipate improving g_A to $\sim 0.5\%$

$$g_A = 1.2711(125) \rightarrow 1.2641(93) [0.74\%]$$

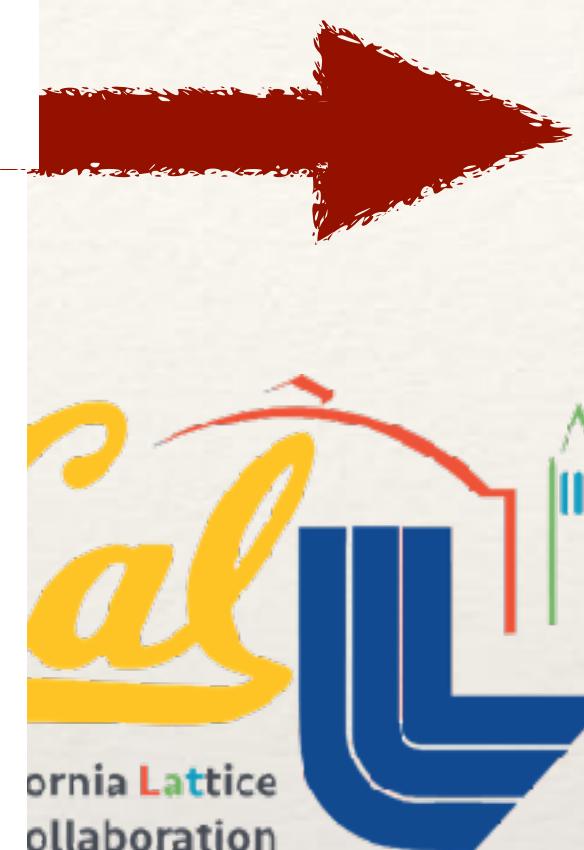
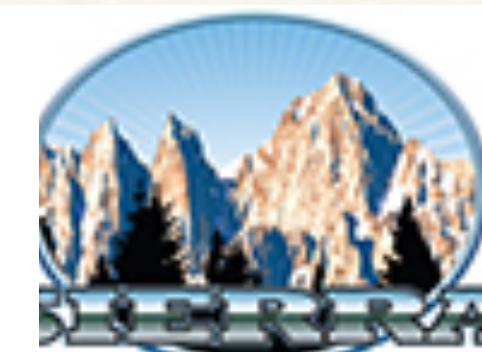
Simulating the *weak* death of the neutron in a femtoscale universe
with near-Exascale computing

Berkowitz, Clark, Gambhir, McElvain, Nicholson, Rinaldi, Vranas,
Walker-Loud, Chang, Joó, Kurth, Orginos

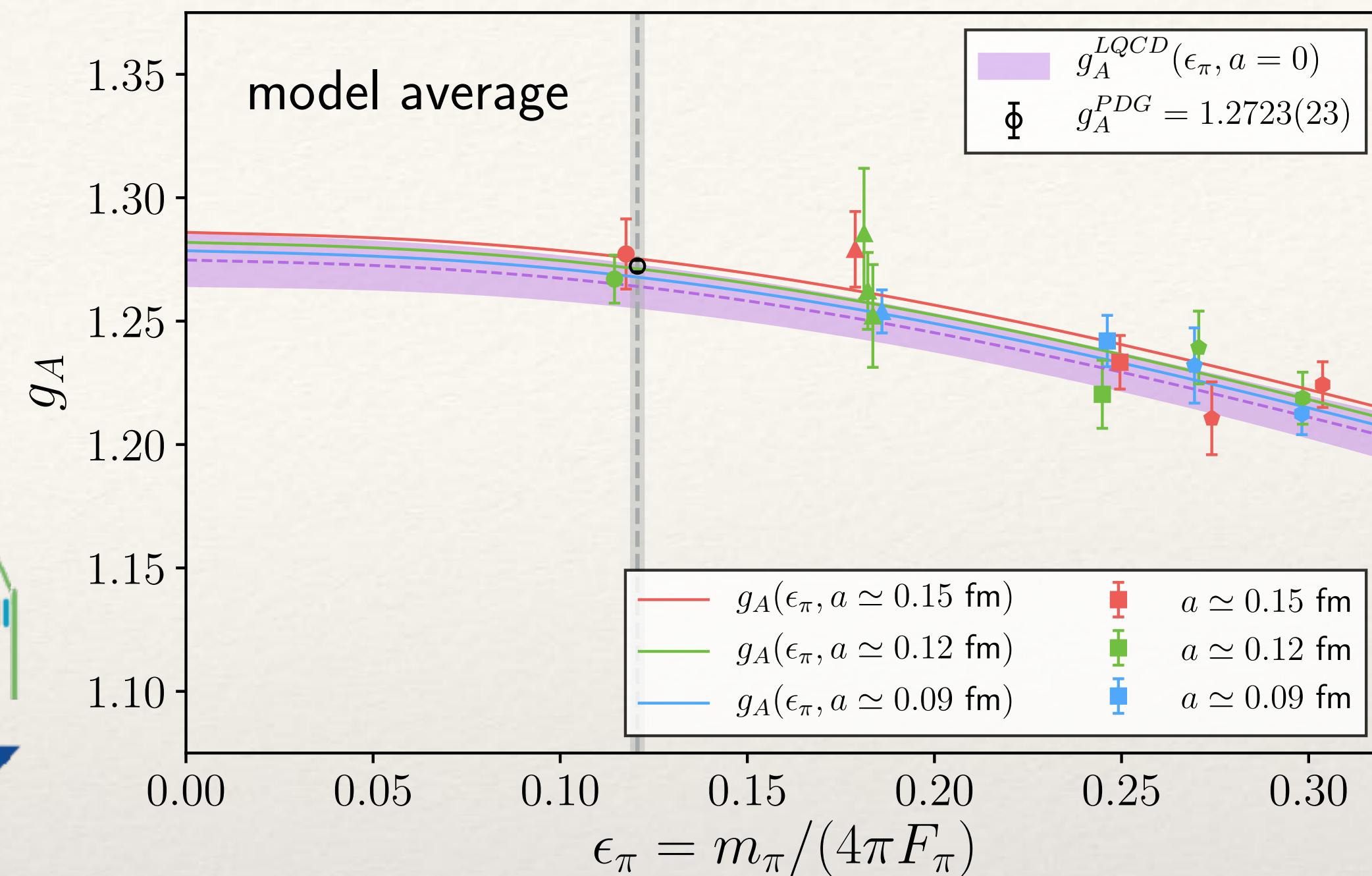


We are running $g_A(Q^2)$ on Summit this year (DOE INCITE)

We anticipate improving g_A to ~0.5%



Sierra Early Science

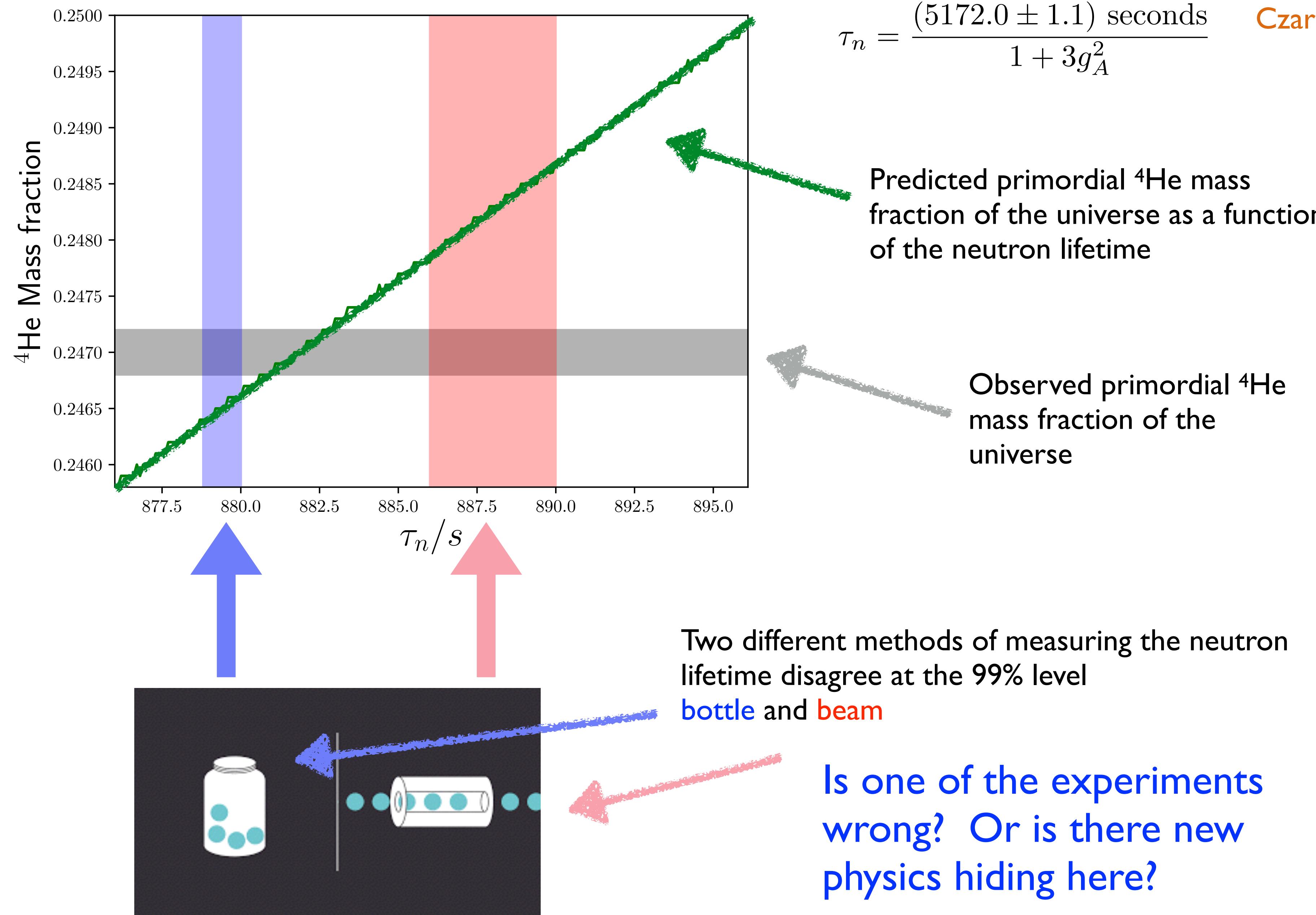


as much as all other ensembles combined

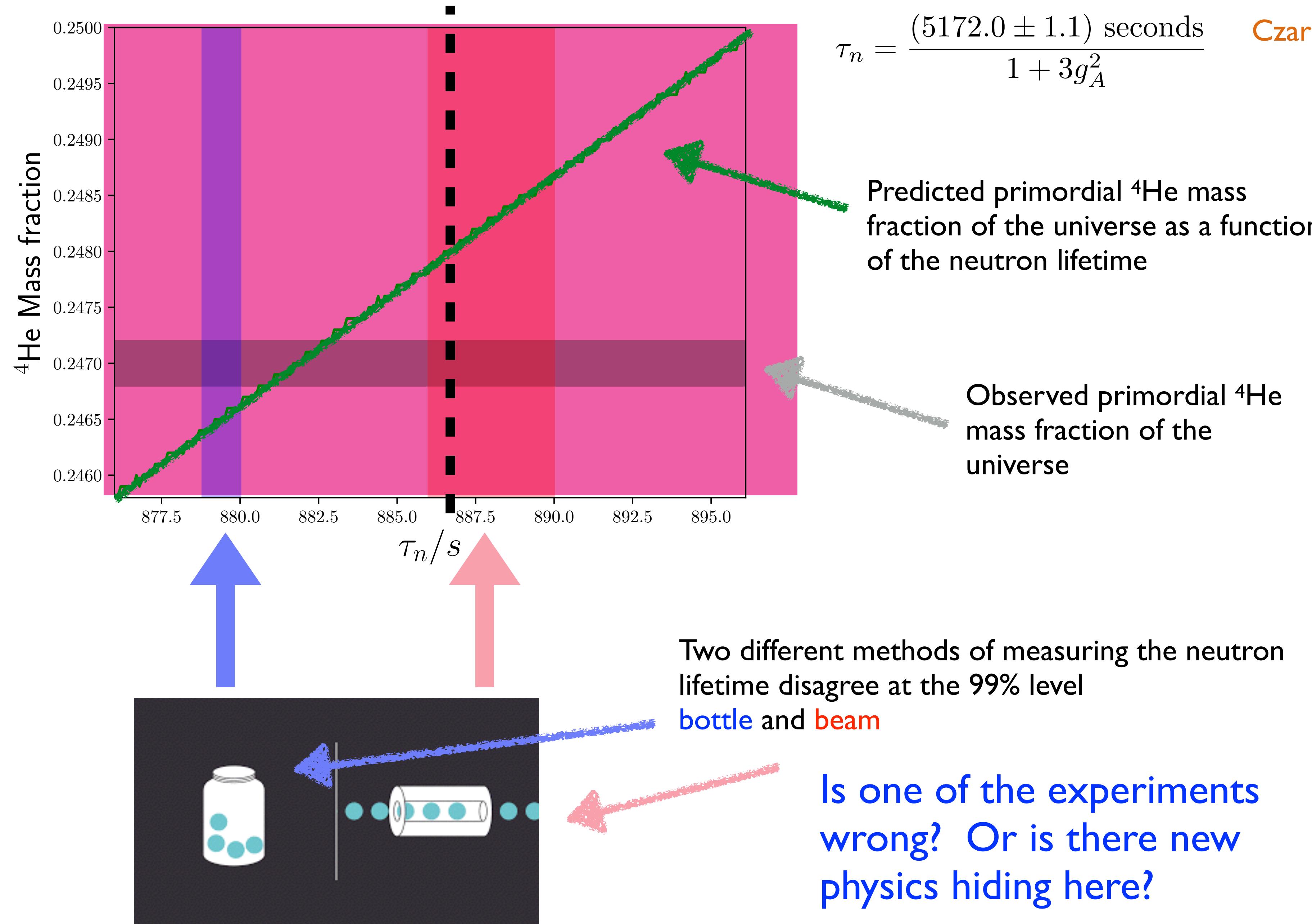
able (old a15m130 is $32^3 \times 48$)

$g_A = 1.2711(125) \rightarrow 1.2641(93) [0.74\%]$

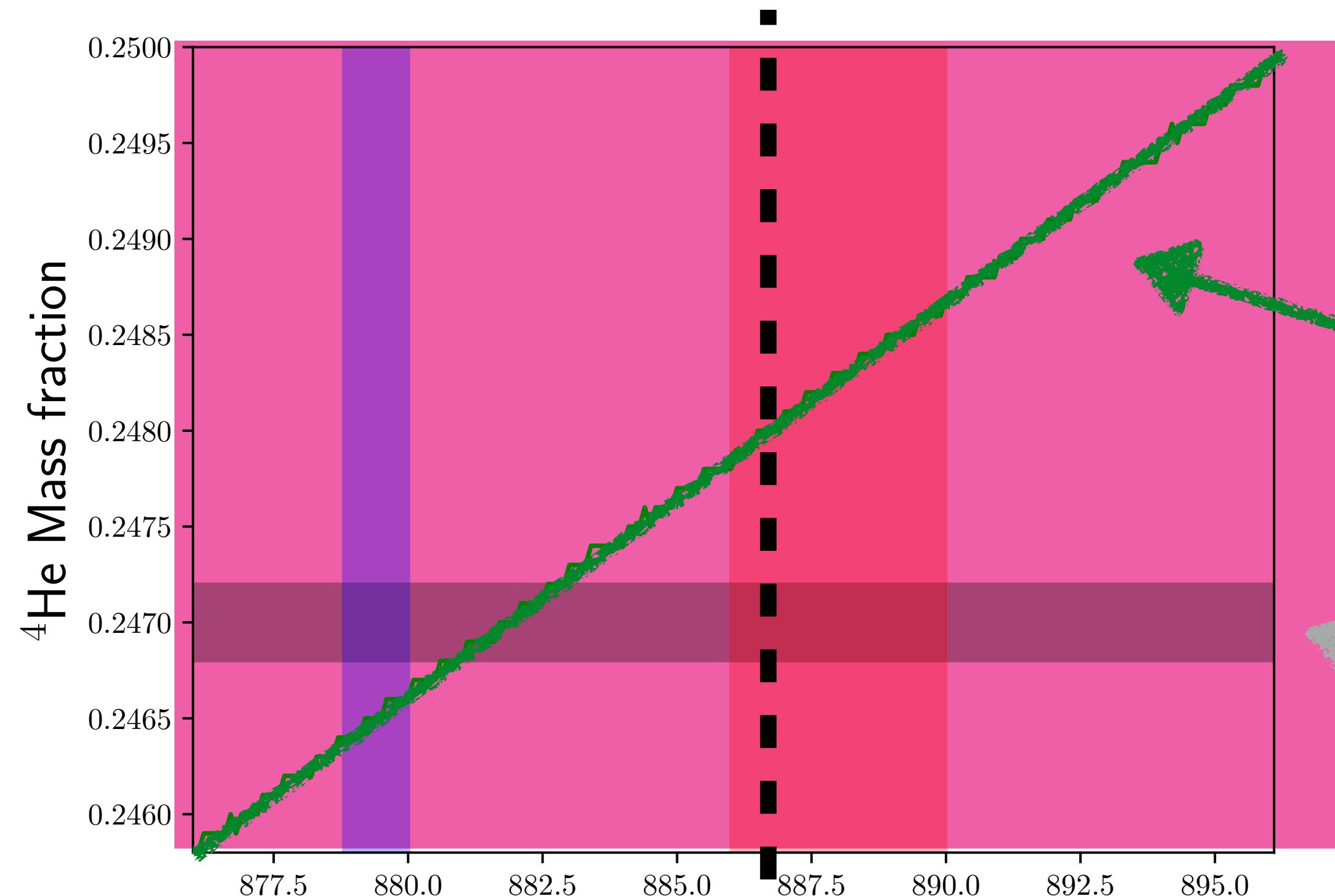
Why is understanding the neutron lifetime interesting?



Why is understanding the neutron lifetime interesting?



Why is understanding the neutron lifetime interesting?

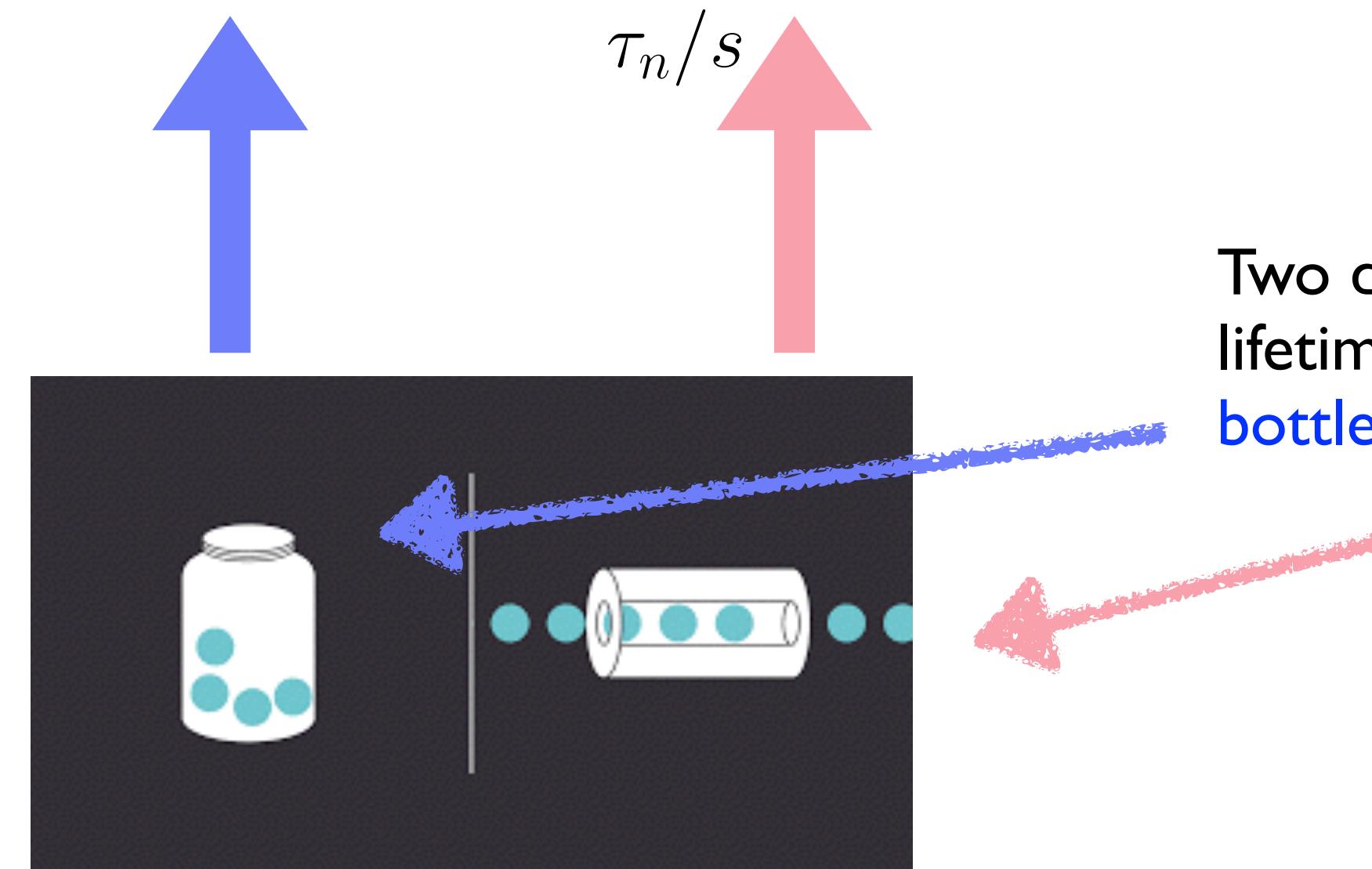


$$\tau_n = \frac{(5172.0 \pm 1.1) \text{ seconds}}{1 + 3g_A^2}$$

Czarnecki, Marciano, Sirlin

Predicted primordial ${}^4\text{He}$ mass fraction of the universe as a function of the neutron lifetime

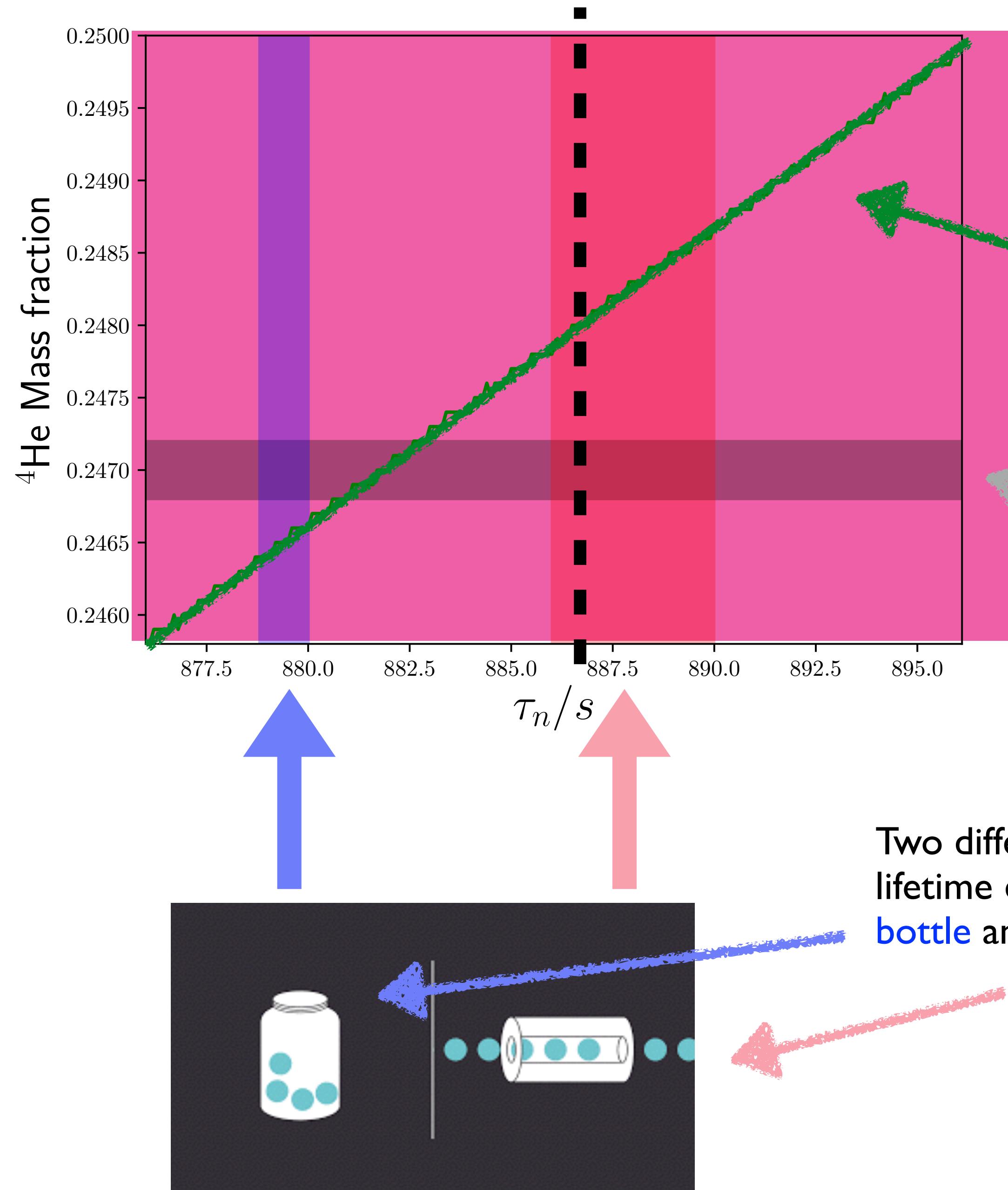
Can we reduce the uncertainty to a discriminating level?



Two different methods of measuring the neutron lifetime disagree at the 99% level
bottle and **beam**

Is one of the experiments wrong? Or is there new physics hiding here?

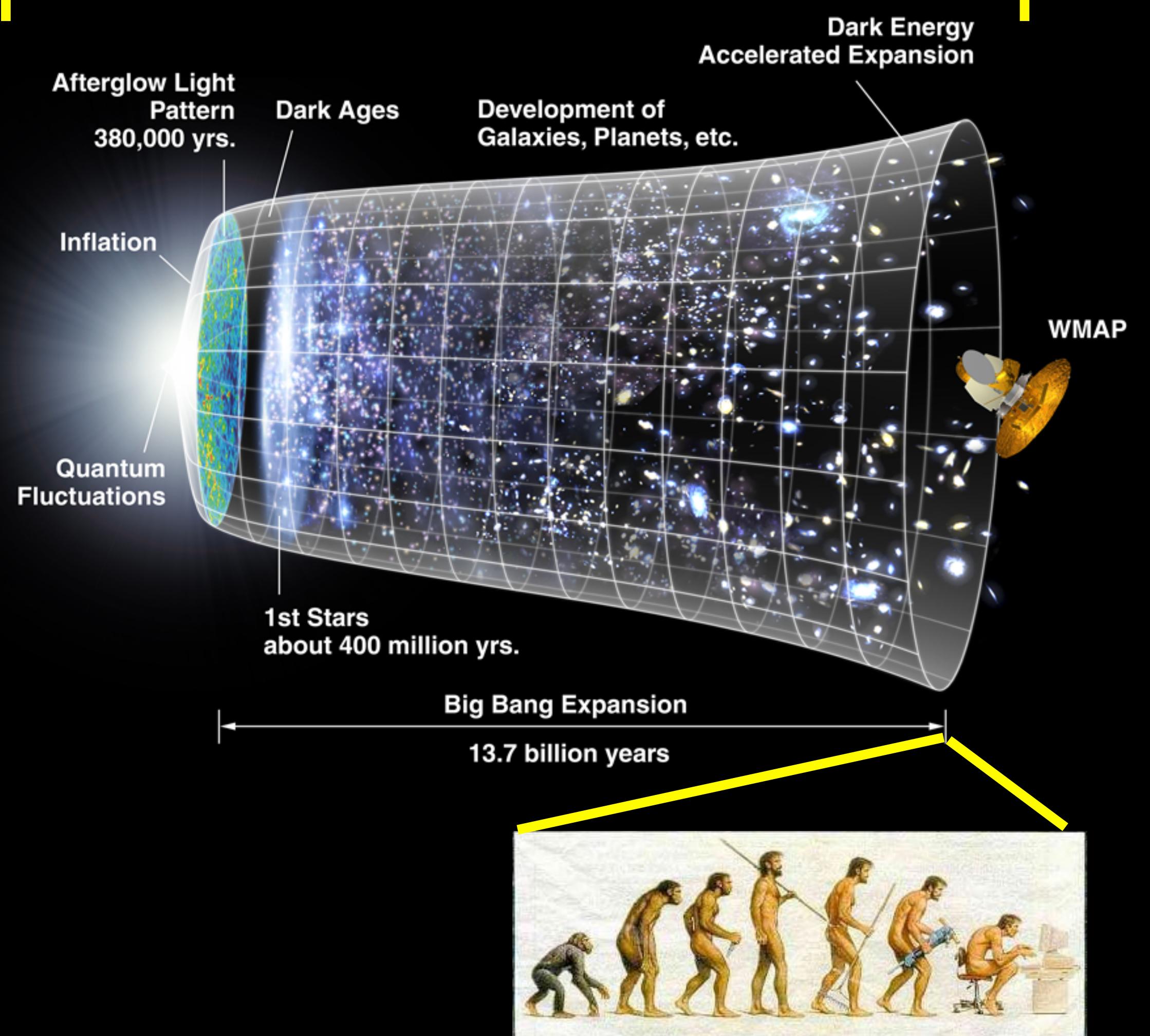
Why is understanding the neutron lifetime interesting?



$$\tau_n = \frac{(5172.0 \pm 1.1) \text{ seconds}}{1 + 3g_A^2}$$

Czarnecki, Marciano, Sirlin

Standard Model of Particle Physics



- ★ This is a very exciting time for basic science research with high-performance-computing
- ★ The application of LQCD to NP is very challenging
 - ★ Progress is very rewarding
 - ★ Necessary for a coordinated effort with Effective Theories of NP - great opportunity to interface and learn other computational methods
- ★ Summit is **disruptively** faster than Titan
 - ★ we are still learning to expand our vision of what can be accomplished
 - ★ Very exciting time to get involved: applications to NP which connect to experiment are just beginning