

# $H + \gamma$ resonance search and interpretation with ATLAS detector

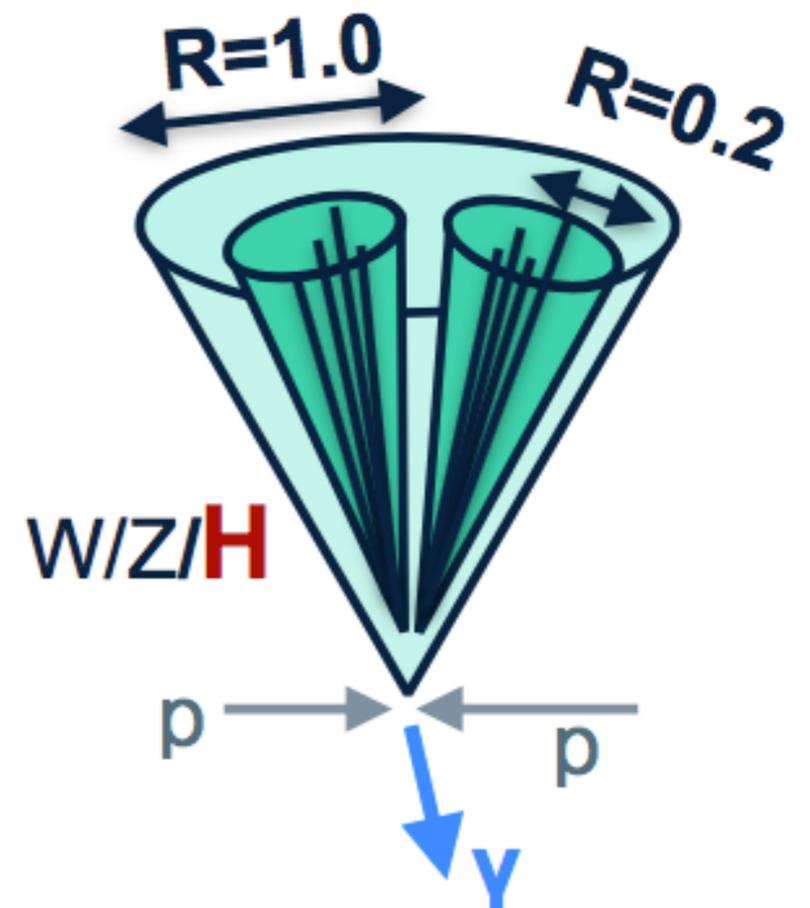
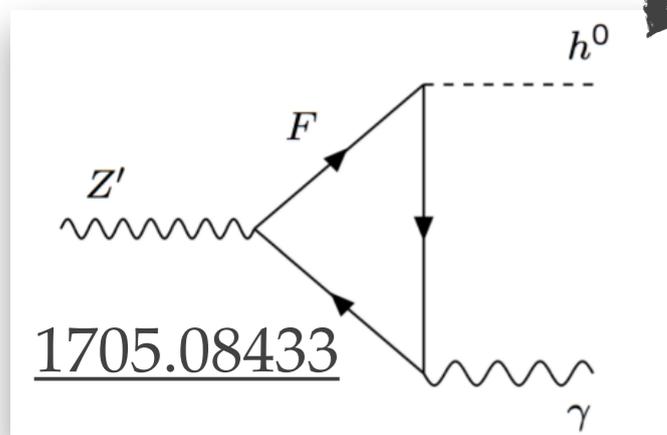
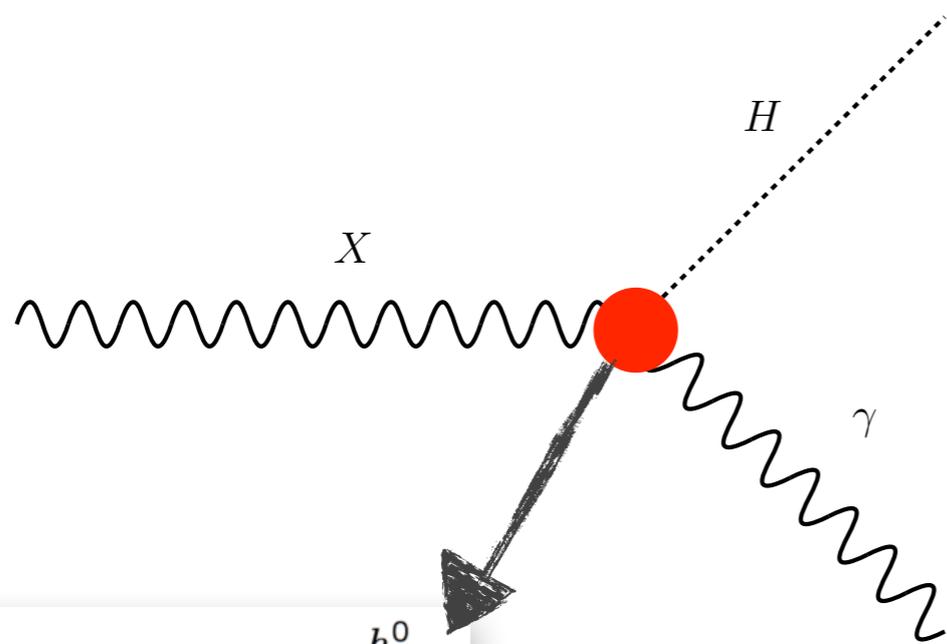
Bo Liu  
IHEP, CAS

Based on Phys. Rev. D 98 (2018) 032015  
and CPC Vol. 43, No. 4 (2019) 043001



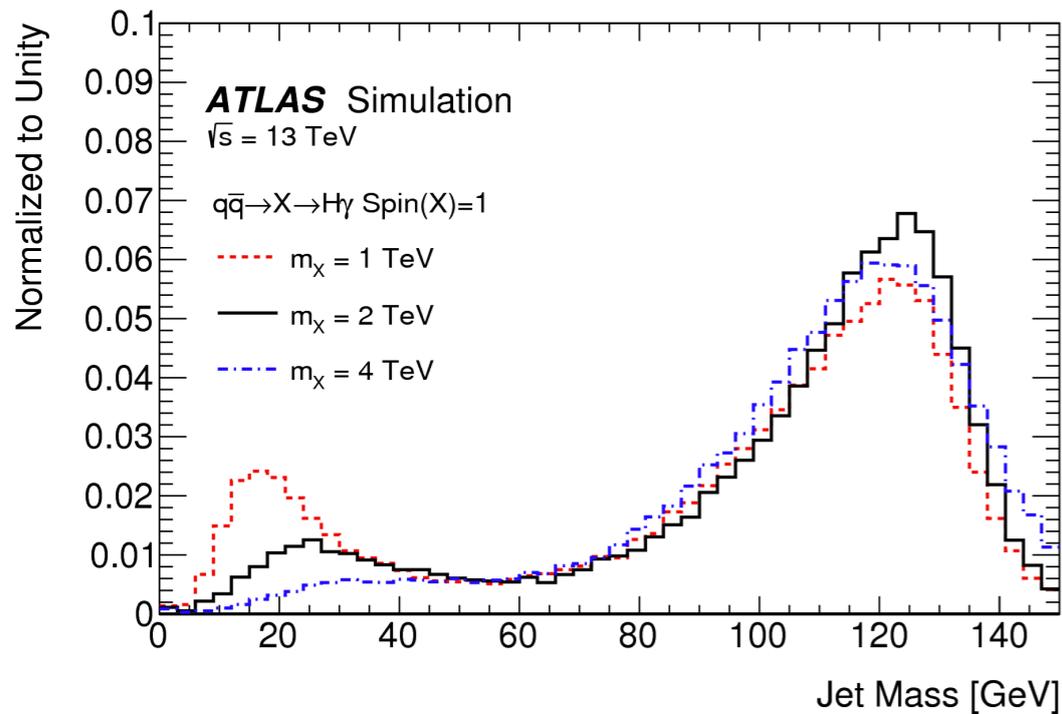
# Introduction

- ❖ Search for heavy resonance decaying to  $H+\gamma$ :
  - Only considered  $b\bar{b}$  decay mode ( $\sim 58\%$ ) and merge/boosted regime
    - Fat jet ( $R=1.0$ ) and high  $p_T$  photon in final state
- ❖ Signals considered
  - $H\gamma$  is induced by Higgs effective coupling model
- ❖ Given high  $p_T$  search, results can be further interpreted with new physics searches





# Analysis selections



## ❖ Photon:

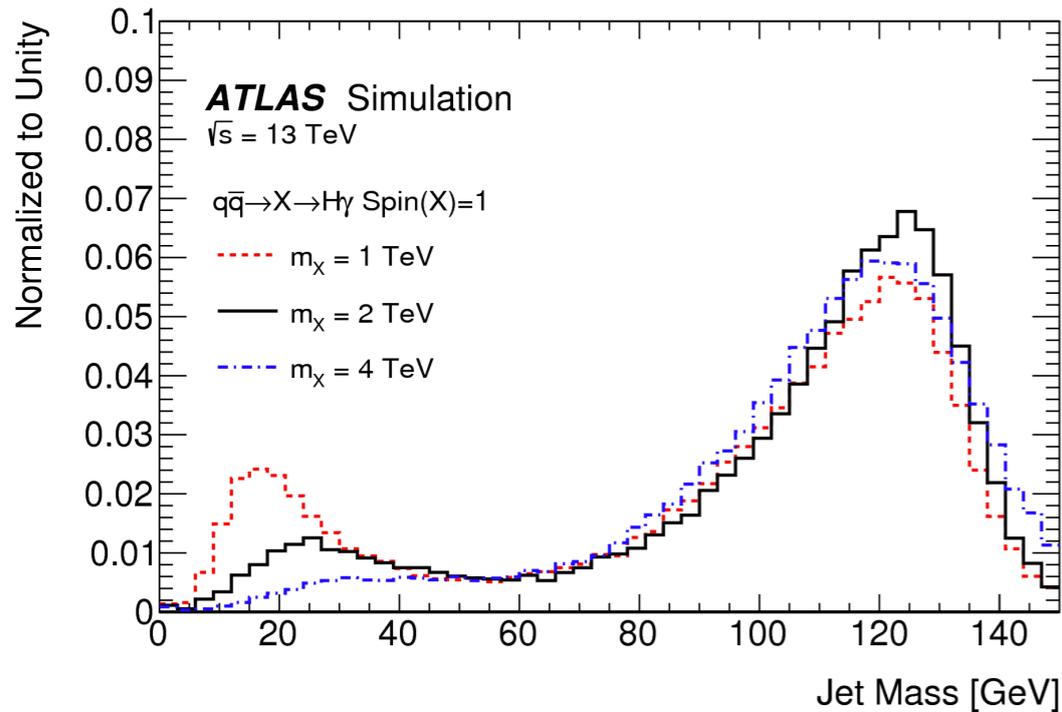
- ▶  $p_T > 250 \text{ GeV}$  and  $|\eta| < 2.37$  (without crack region [1.37, 1.52])
- ▶ Pass tight photon ID selection and tight calorimeter only isolation

## ❖ Jet:

- ▶ Anti- $k_T$  large- $R$  jet ( $R=1.0$ ), Trimmed ( $f_{\text{cut}} < 5\%$ ,  $R_{\text{sub}}=0.2$ )
- ▶  $p_T > 200 \text{ GeV}$  and  $|\eta| < 2.0$
- ▶ Apply mass window cut 93-134 GeV
- ▶ Anti- $k_T$   $R=0.2$  trackjet btagging using MV2c10 algorithm @70% efficiency.



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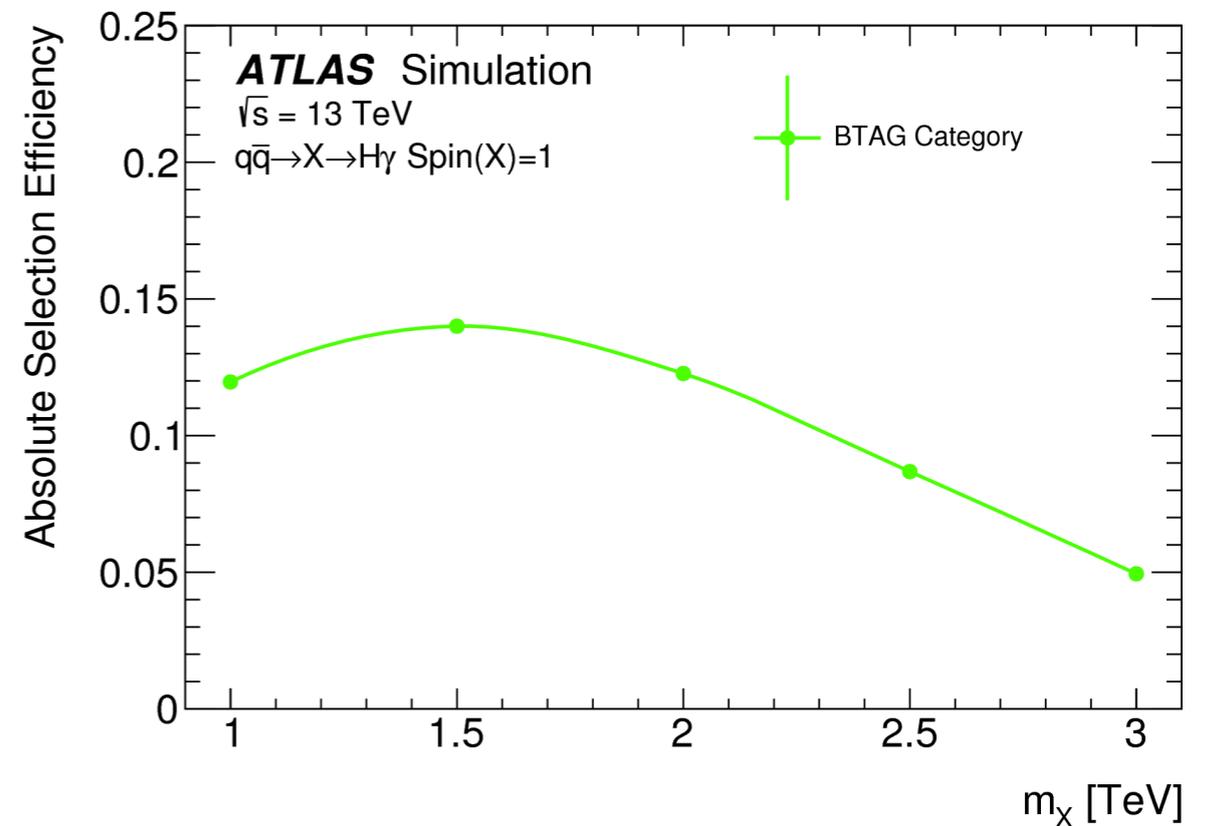
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## ❖ Preselections

- ▶ Trigger: HLT\_g140\_loose
- ▶ One good PV by require at least two associated track.
- ▶ At least one photon in barrel calorimeter ( $|\eta| < 1.37$ )
- ▶  $\Delta R(J, \gamma) > 1.0$ .
- ▶ Event should present at least one large  $R$  jet and at least one photon.
  - The leading  $p_T$  object will be used to form  $m(J\gamma)$ .

## ❖ Categorisation:

- ▶ **btag (2btagged)**



# Signal and background modelling

Signals are modelled by the *CrystalBall+Gaussian* function

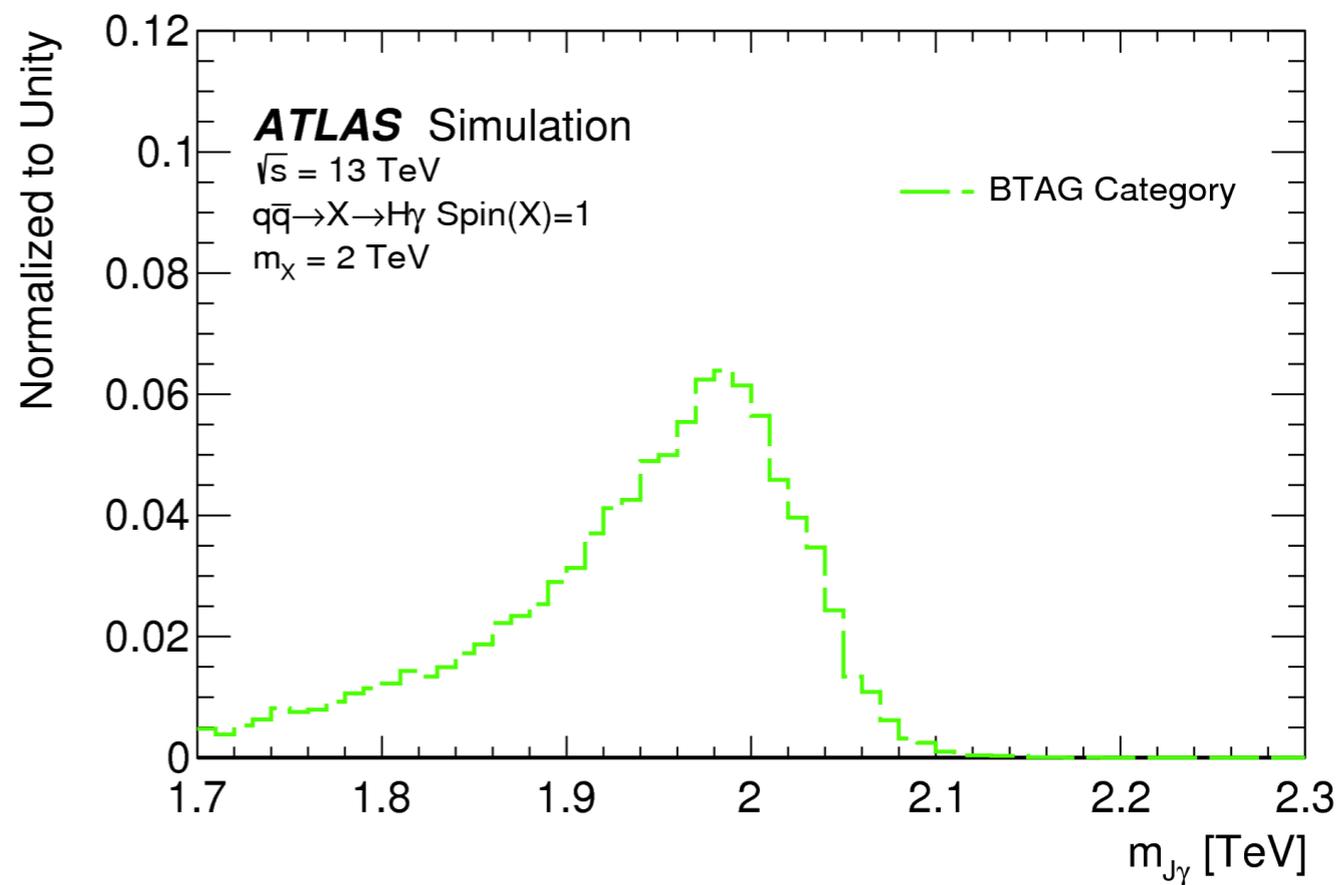
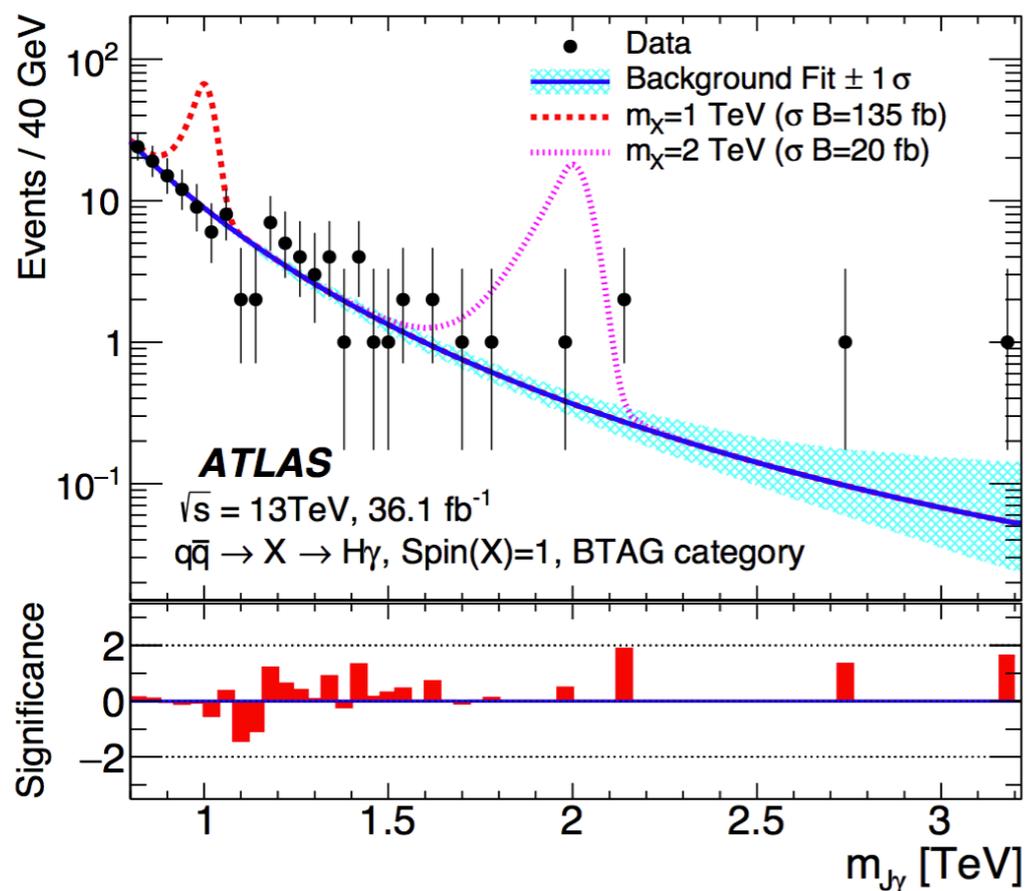
$$S(m_{J\gamma}) = f_C \cdot C(m_{J\gamma}; \mu, \sigma_C, \alpha_C, n_C) + (1 - f_C) \cdot G(m_{J\gamma}; \mu, \sigma_G)$$

Where *CrystalBall* function can be written as:

$$C(m_{J\gamma}; \mu, \sigma_C, \alpha_C, n_C) = N_C \cdot \begin{cases} \exp\left(-\frac{(m_{J\gamma}-\mu)^2}{2\sigma_C^2}\right) & \frac{m_{J\gamma}-\mu}{\sigma_C} > -\alpha_C \\ \left(\frac{n_C}{\alpha_C}\right)^{n_C} \exp\left(-\frac{\alpha_C^2}{2}\right) \left(\frac{n_C}{\alpha_C} - \alpha_C - \frac{m_{J\gamma}-\mu}{\sigma_C}\right)^{-n_C} & \frac{m_{J\gamma}-\mu}{\sigma_C} \leq -\alpha_C \end{cases}$$

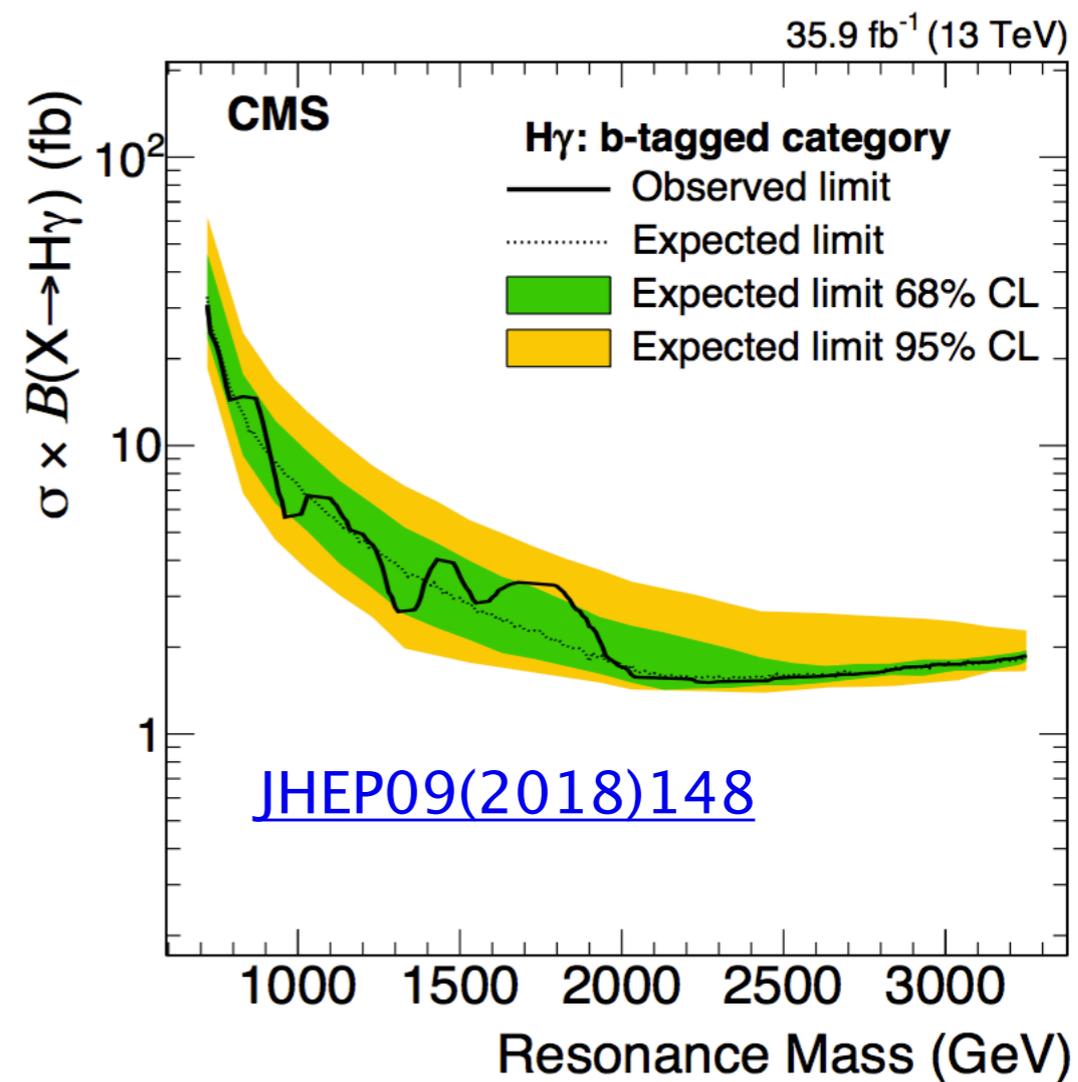
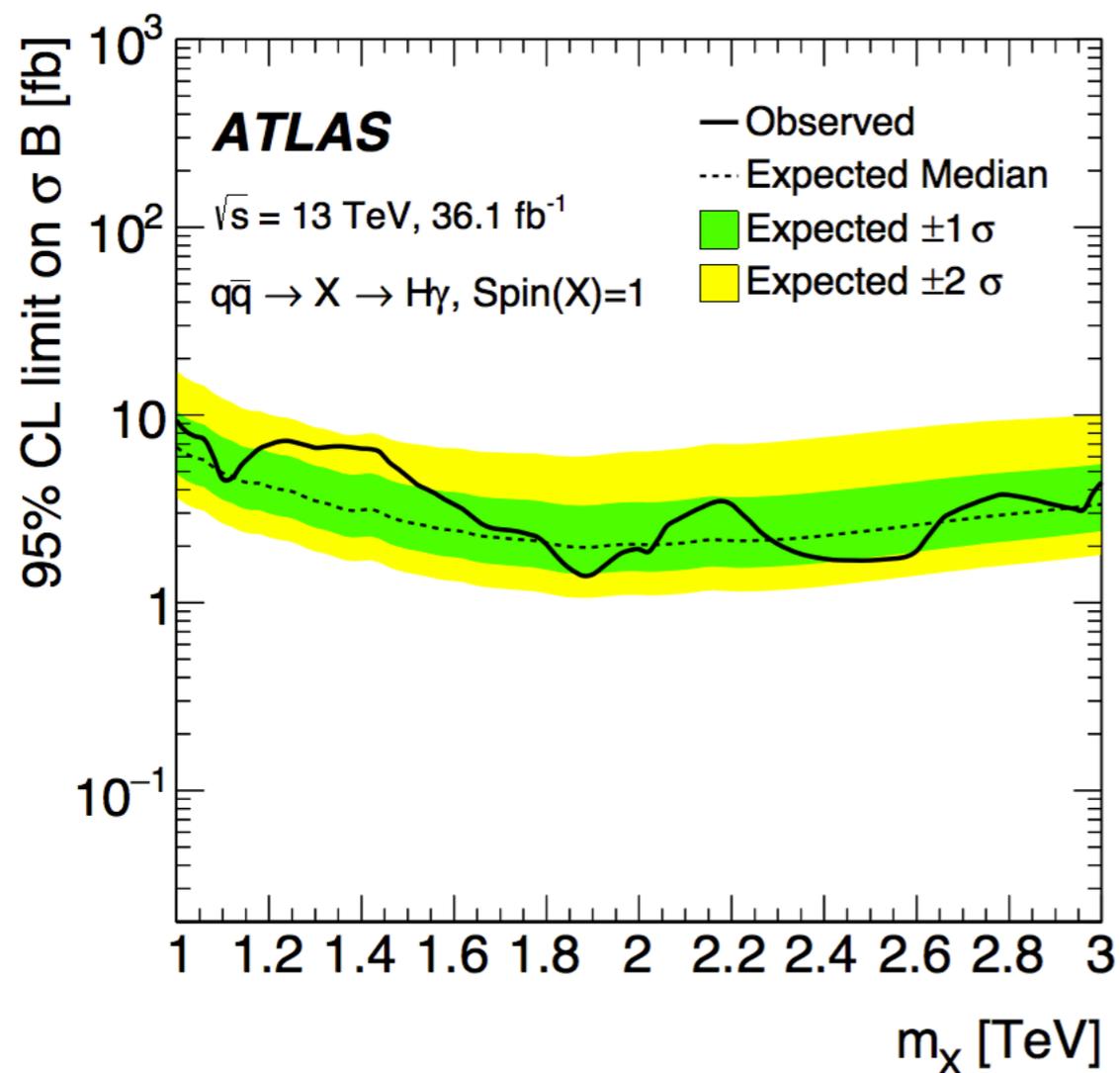
Background are modelled as:

$$B(m_{J\gamma}; p_i) = (1-x)^{p_1} x^{p_2+p_3 \log(x)+p_4 \log^2(x)+p_5 \log^3(x)}, x = m_{J\gamma}/\sqrt{s}$$





# Results



- ❖ Similar sensitivity as CMS **btagged** region (CMS also explore untagged region)
- ❖ The difference in high mass due to better Higgs Tagger performance for CMS
- ❖ More advanced Tagger has been developed for ATLAS ([ATL-PHYS-PUB-2017-010](#))
- ❖ Results to include new tagger is preparation with full run-II data



# One more story: interpretation

- ❖ Article [Phys.Lett. B773 \(2017\) 462-469](#) propose the potential way of constraint Higgs boson couplings via H+photon production channel and test new physics
- ❖ The HEFT model is used and mainly considered strongly interacting light Higgs (SILH) basis which strongly related to H+photon production

$$\begin{aligned}\mathcal{L}_{\text{SILH}} = & \frac{g_s^2 \bar{c}_g}{m_W^2} \Phi^\dagger \Phi G_{\mu\nu}^a G_a^{\mu\nu} + \frac{g'^2 \bar{c}_\gamma}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{ig' \bar{c}_B}{2m_W^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{ig \bar{c}_W}{2m_W^2} [\Phi^\dagger \sigma_k \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig \bar{c}_{HW}}{m_W^2} [D^\mu \Phi^\dagger \sigma_k D^\nu \Phi] W_{\mu\nu}^k - \frac{\bar{c}_6 \lambda}{v^2} [\Phi^\dagger \Phi]^3 \\ & + \frac{ig' \bar{c}_{HB}}{m_W^2} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} + \frac{\bar{c}_H}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{\bar{c}_T}{2v^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \\ & - \left[ \frac{\bar{c}_l}{v^2} y_l \Phi^\dagger \Phi \Phi \bar{L}_L e_R + \frac{\bar{c}_u}{v^2} y_u \Phi^\dagger \Phi \Phi^\dagger \cdot \bar{Q}_L u_R + \frac{\bar{c}_d}{v^2} y_d \Phi^\dagger \Phi \Phi \bar{Q}_L d_R + \text{h.c.} \right],\end{aligned}$$



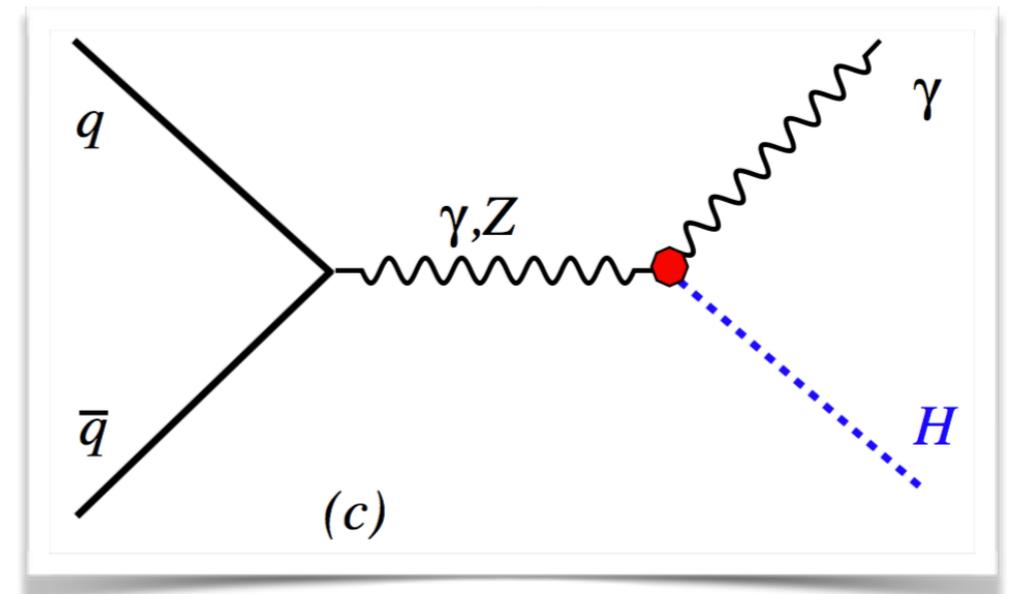
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 & - \left[ \frac{\bar{c}_l}{v^2} y_l \Phi^\dagger \Phi \Phi \bar{L}_L e_R + \frac{\bar{c}_u}{v^2} y_u \Phi^\dagger \Phi \Phi^\dagger \cdot \bar{Q}_L u_R + \frac{\bar{c}_d}{v^2} y_d \Phi^\dagger \Phi \Phi \bar{Q}_L d_R + \text{h.c.} \right],
 \end{aligned}$$

Mass basis	Gauge basis
$g_{h\gamma\gamma}$	$a_H - \frac{8gs_W^2}{m_W} \bar{c}_\gamma$
$g_{h\gamma Z}^{(1)}$	$\frac{gs_W}{c_W m_W} (\bar{c}_{HW} - \bar{c}_{HB} + 8s_W^2 \bar{c}_\gamma)$
$g_{h\gamma Z}^{(2)}$	$\frac{gs_W}{c_W m_W} (\bar{c}_{HW} - \bar{c}_{HB} - \bar{c}_B + \bar{c}_W)$
$y_u$	$y_u [1 - \frac{1}{2} \bar{c}_H + \frac{3}{2} \bar{c}_u]$
$\tilde{y}_d$	$y_d [1 - \frac{1}{2} \bar{c}_H + \frac{3}{2} \bar{c}_d]$
$g_{h\gamma uu}^{(\partial)}$	$\frac{\sqrt{2}gs_W}{m_W^2} y_u [\bar{c}_{uB} + \bar{c}_{uW}]$
$g_{h\gamma dd}^{(\partial)}$	$\frac{\sqrt{2}gs_W}{m_W^2} y_d [\bar{c}_{dB} - \bar{c}_{dW}]$

Constraints are only calculated for  $\bar{c}_{HW}, \bar{c}_{HB}, \bar{c}_\gamma$   
 Other terms are found have small effect  
 on results

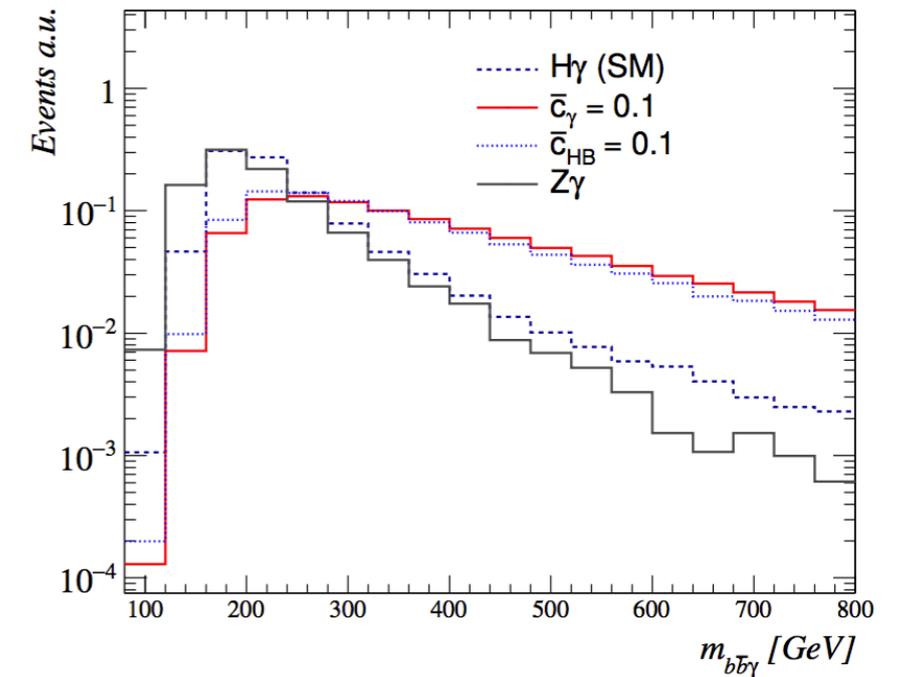


Key diagram have leading contribution 6



# Parametrization

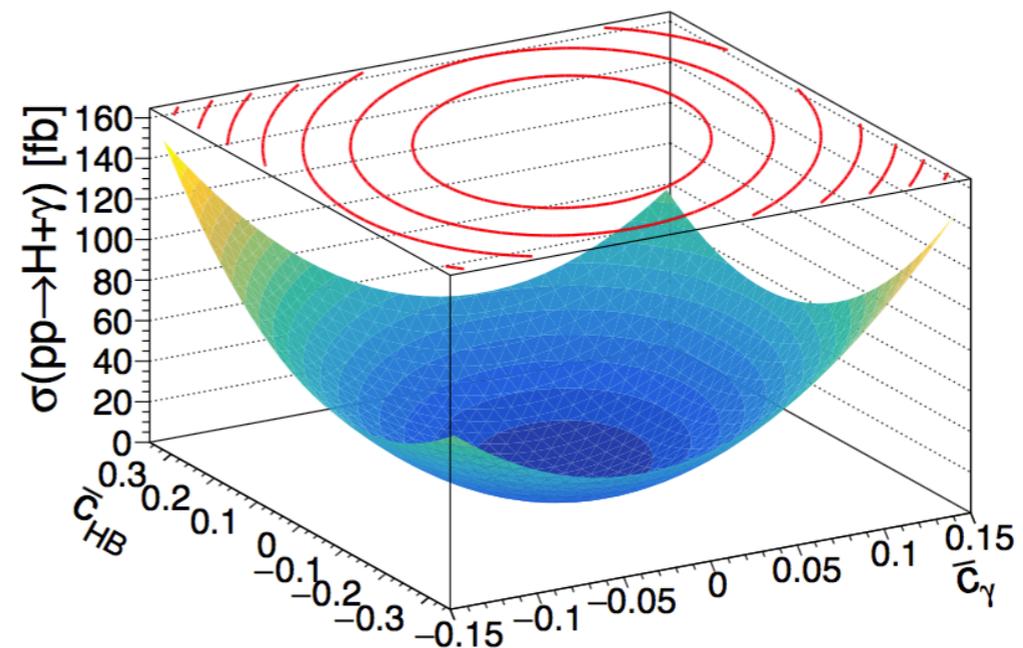
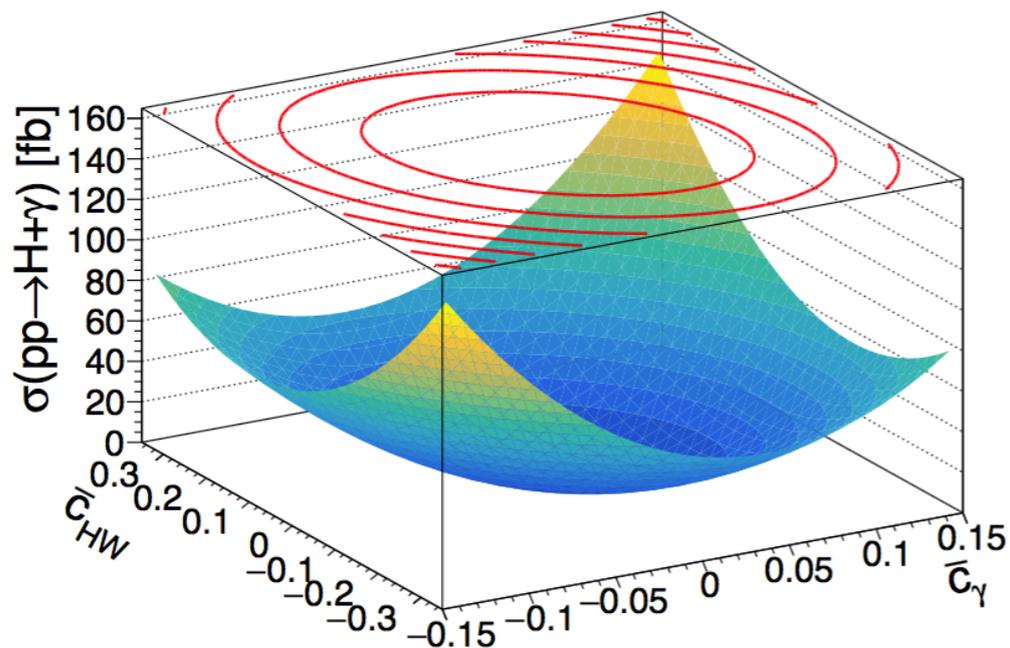
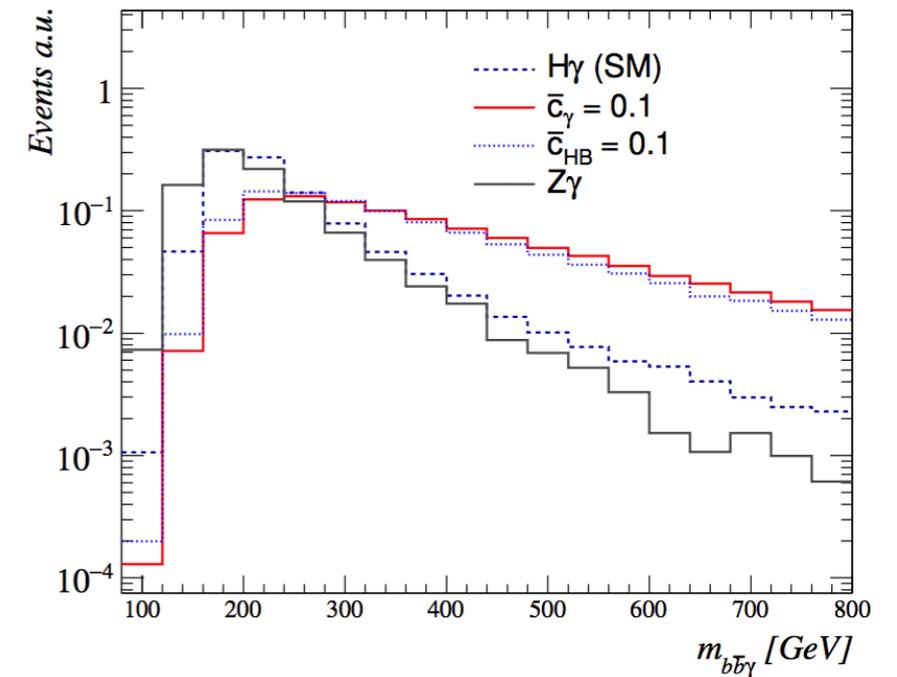
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  - Current resonance search paper only have mass distribution available





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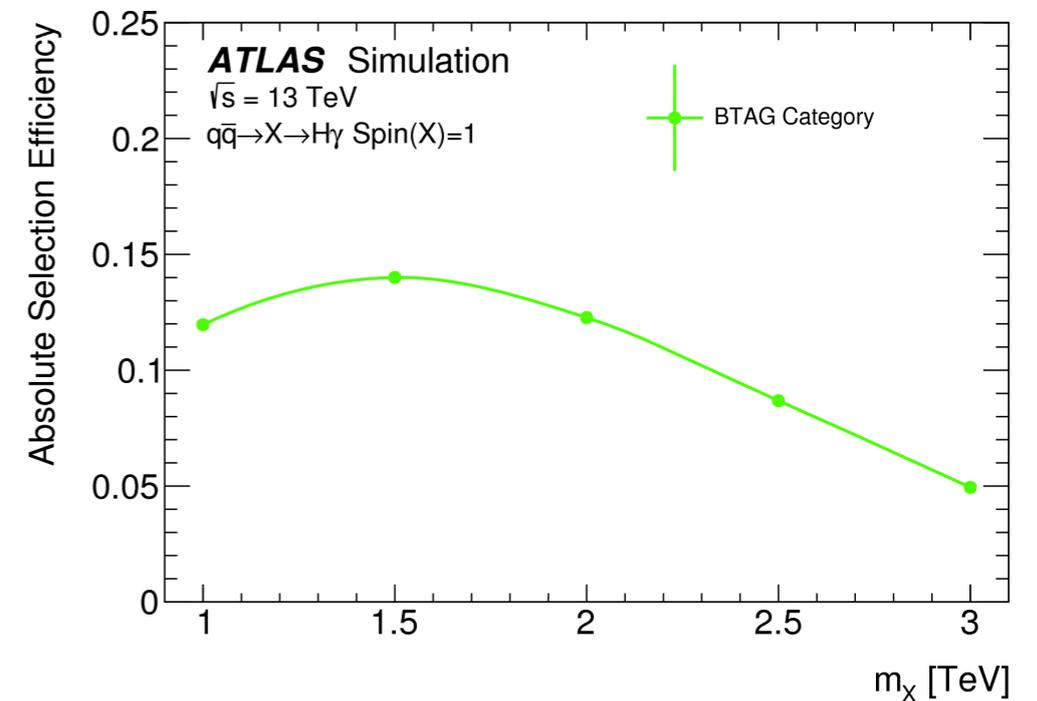
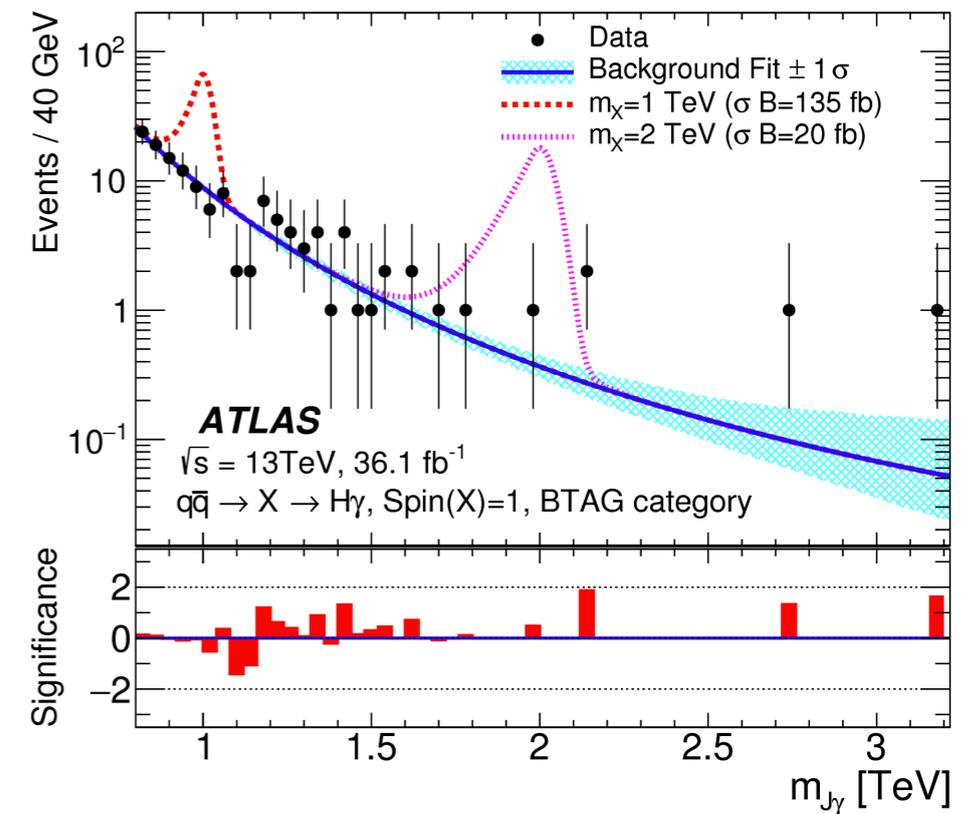
- ❖ Based on pheno paper, the invariant mass of H+photon system is one of the sensitive variable to probe new couplings
  - Current resonance search paper only have mass distribution available
- ❖ Currently, SILH is implemented along with generic HEFT model in MadGraph
- ❖ A numeric relationship between EFT coefficients and H+photon crosssection with MadGraph calculation and parametrised with polynomial function for statistical analysis





# Statistical treatment

- ❖ Since we only have truth events for EFT signals, only the number counting analysis performed for  $m_{H\gamma}$  in between 800-3200 GeV
- ❖ QCD background contribution is obtained from data fit along with uncertainty.
  - SM H+photon contribution is small and neglected
- ❖ The effects on signal efficiency due to acceptance, detector effect and selection efficiency are applied with efficiency parametrization on  $m_{H\gamma}$





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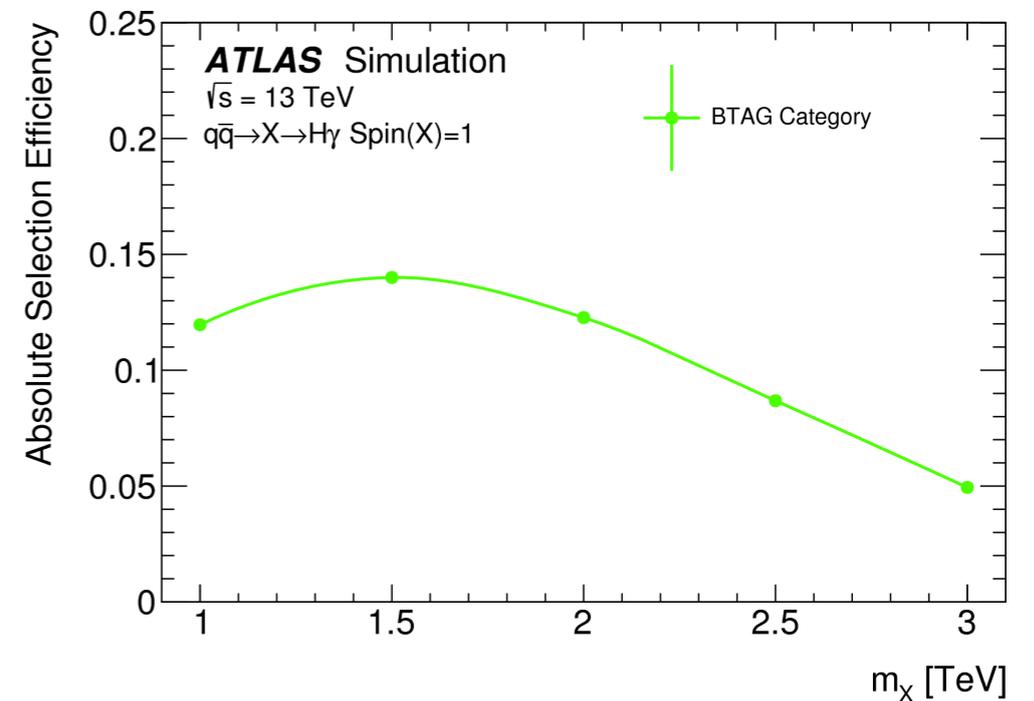
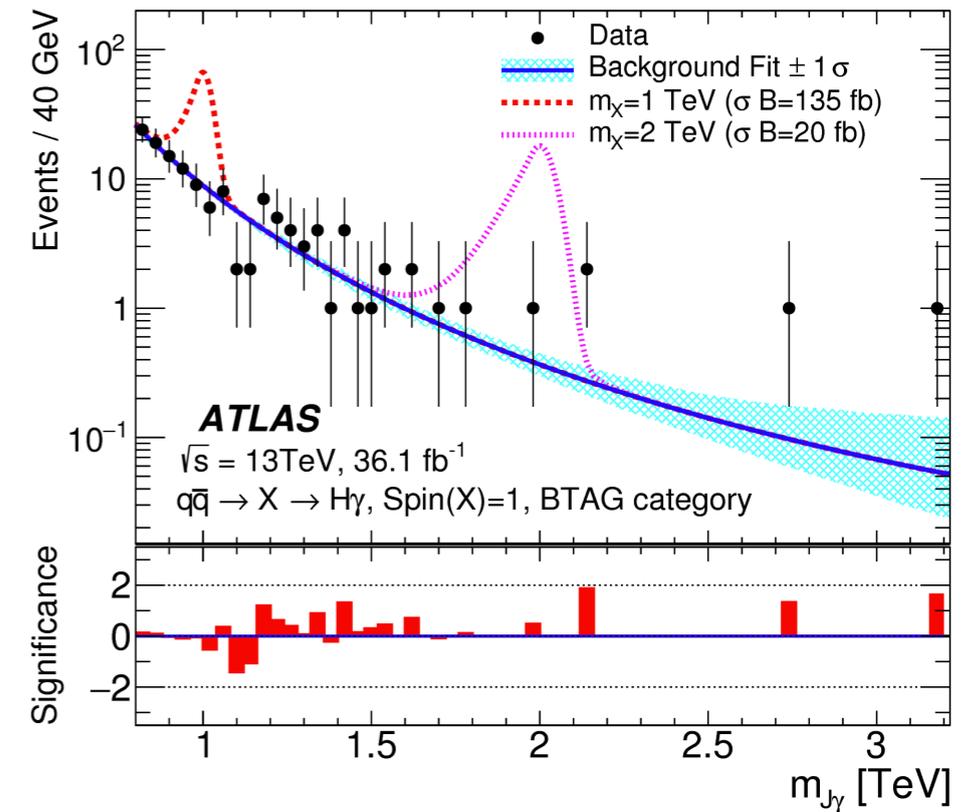
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Statistical model is built as:

$$\mathcal{L} = \text{Pois}(n|s + b) \times \text{Gaus}(b_0|b, \sigma_b) .$$

Constraints are obtained with PLR

$$\lambda(\bar{c}_i) = \frac{\mathcal{L}(\bar{c}_i, \hat{b})}{\mathcal{L}(\hat{c}_i, \hat{b})} .$$

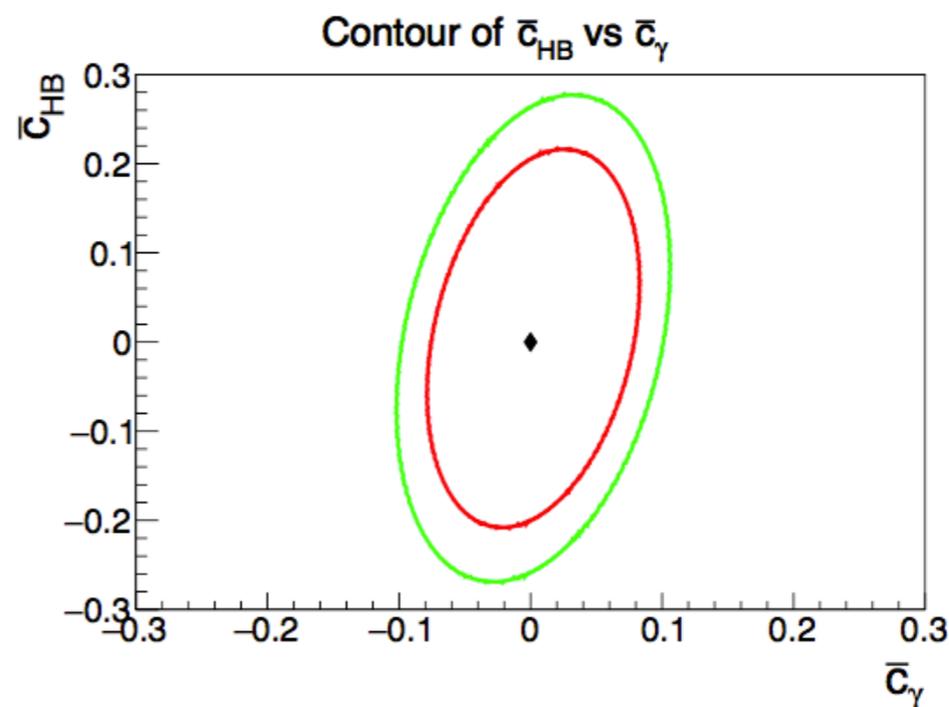
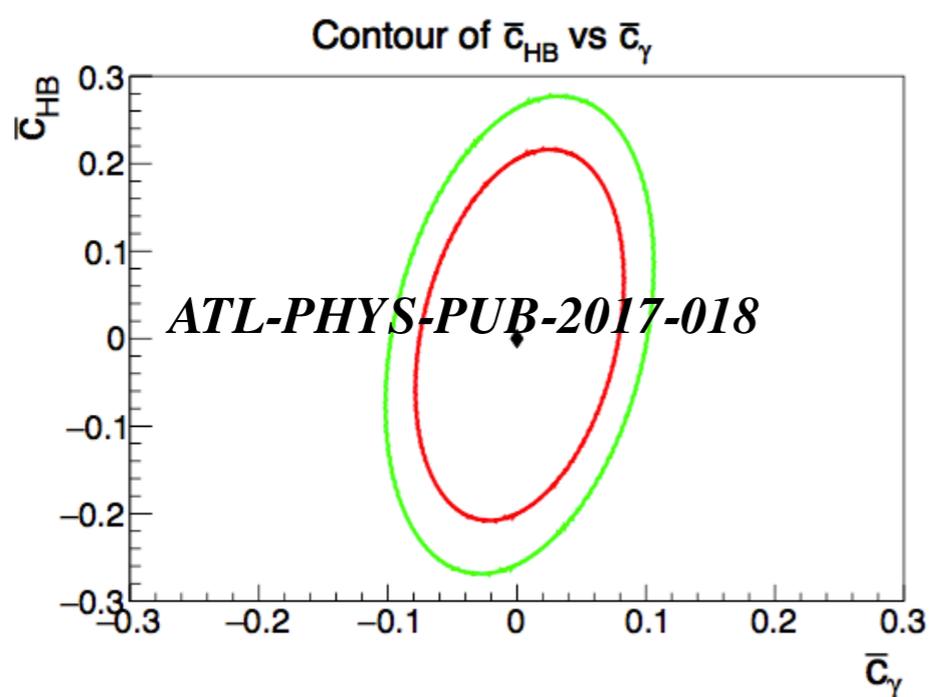
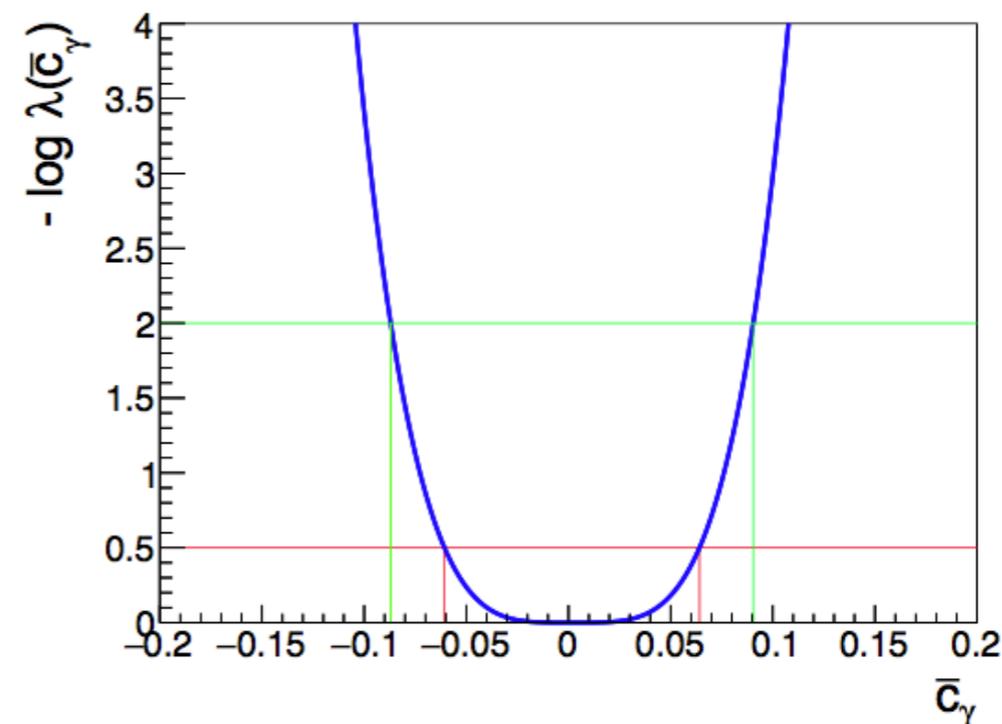




# Constraints

## ❖ Results

Parameter	68% C.L.	95% C.L.
$\bar{c}_\gamma$	[-0.061, 0.064]	[-0.087, 0.090]
$\bar{c}_{HW}$	[-0.167, 0.161]	[-0.236, 0.231]
$\bar{c}_{HB}$	[-0.162, 0.167]	[-0.230, 0.236]



❖ Future potential shape analysis can further improve the constrains

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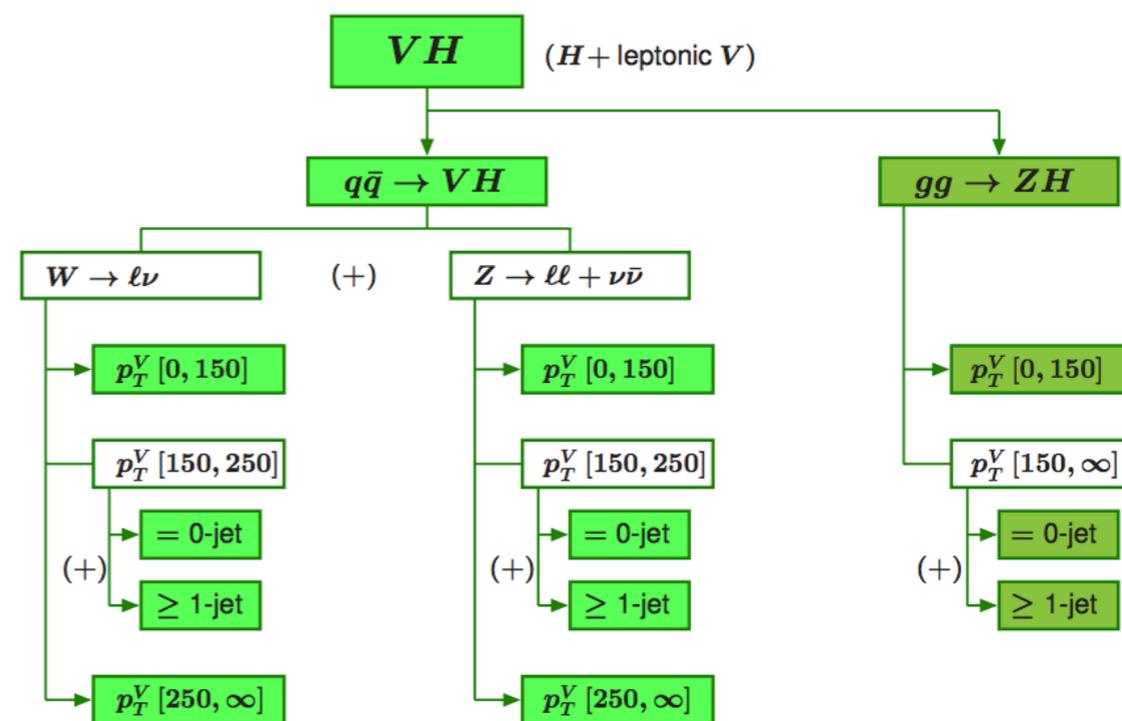
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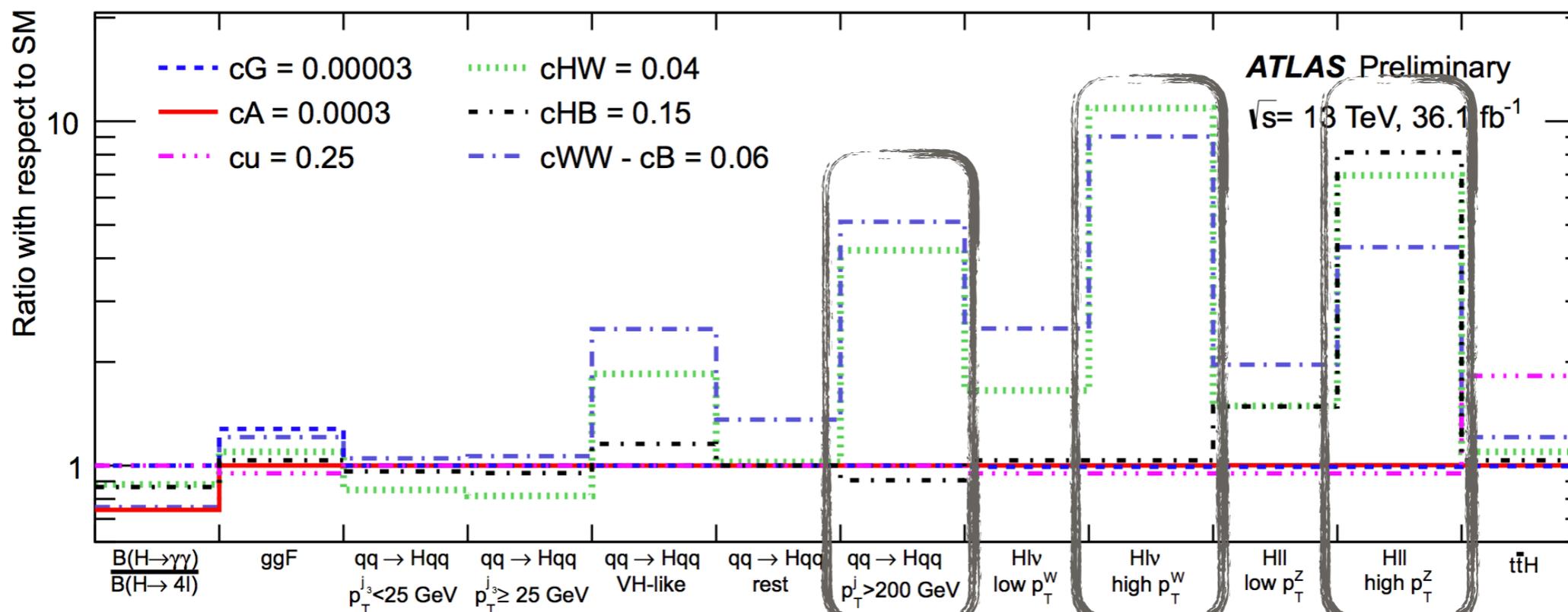
$$\bar{c}_\gamma \in [-1.5 \times 10^{-4}, 2.2 \times 10^{-4}],$$

$$\bar{c}_{HW} \in [-0.080, -0.024],$$

$$\bar{c}_{HB} \in [-0.051, 0.103].$$



From ATLAS-PHYS-PUB-2017-018 ( $H \rightarrow ZZ$  &  $H \rightarrow \gamma\gamma$  STXS)





# Summary

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- ❖ A search for heavy resonance decaying to Higgs boson and a photon has been presented
- ❖ Large-R jet is used to contain all the decay products from the Higgs boson.
  - Search range limited in between 1-3 TeV for the performance of fatjet and flavour tagging
- ❖ No significant excess observed
- ❖ In addition to the model-independent search, an EFT interpretation has been made based on public dataset
- ❖ Similar constraint power as  $H \rightarrow ZZ$  &  $H \rightarrow \gamma\gamma$  combined results for few coefficients is obtained
- ❖ Potential shape analysis can future improve constraint power
- ❖ The study to improve the sensitivity with run-II dataset is ongoing