## CEPC Physics Workshop

## Differential measurement on Higgs

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## Introduction

$\square$ After the discovery of the SM Higgs particle a new era emerged in the so-called physics BSM;
$\square$ This comes from the fact that the neutrino oscillation experiments showed that the neutrinos possess a tiny masses, unlike the SM;
$\square$ In this context theorists and experimentalists are very keen to investigate this evidence;
$\square$ Many theories have been proposed in this regards. For instance, the 2HDM model, in which particles heavier than the SM Higgs are suggested but yet to be observed;
$\square$ Also, the crossing-symmetric of $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ showed to be promising in probing BSM scenario; and
$\square$ This is by argue that the angular asymmetries have the potential to reveal the hidden BSM physics in its differential cross section.

## Introduction

$\square$ A study showed the advantages of $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ against $H Z(\rightarrow \ell \ell)$, see arXiv:1406.1361;
$\square$ It suggested that a high-energy $\mathrm{e}^{+} e^{-}$colliders would provide a clean way to estimate the Higgs couplings; (CEPC )
$\square$ In this study we are trying to develop a generator for $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ within the CEPC framework;
$\square$ Hence, use sophisticated differential cross section analysis to do Higgs couplings measurements;

# The differential cross section $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ 



$$
\frac{\mathrm{d} \sigma}{\mathrm{~d} \cos \theta_{1} \mathrm{~d} \cos \theta_{2} \mathrm{~d} \phi}=\frac{1}{m_{H}^{2}} \mathcal{N}_{\sigma}\left(q^{2}\right) \mathcal{J}\left(q^{2}, \theta_{1}, \theta_{2}, \phi\right)
$$

$\square \mathcal{N}_{\sigma}\left(q^{2}\right)$ is the normalisation factor and it can be written in terms of the dimensionless parameters $r$ and $s$ as:

$$
\mathcal{N}_{\sigma}\left(q^{2}\right)=\frac{1}{2^{10}(2 \pi)^{3}} \frac{1}{\sqrt{r} \gamma Z} \frac{\sqrt{\lambda(1, s, r)}}{s^{2}}
$$

$\square$ The constant dimensionless parameters given by the following:

$$
s=\frac{q_{\mathrm{th}}^{2}}{m_{H}^{2}} \approx 2.98, r=\frac{m_{Z}^{2}}{m_{H}^{2}} \approx 0.53, \gamma Z=\frac{\Gamma_{Z}}{m_{H}} \approx 0.020
$$

## The differential cross section $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$

$\square \mathcal{J}\left(q^{2}, \theta_{1}, \theta_{2}, \phi\right)$ depends on nine $J_{i}$ functions expressed by:

$$
\begin{aligned}
\mathcal{J}\left(q^{2}, \theta_{1}, \theta_{2}, \phi\right) & =J_{1}\left(1+\cos ^{2} \theta_{1} \cos ^{2} \theta_{2}+\cos ^{2} \theta_{1}+\cos ^{2} \theta_{2}\right) \\
& +J_{2} \sin ^{2} \theta_{1} \sin ^{2} \theta_{2}+J_{3} \cos \theta_{1} \cos ^{2} \theta_{2} \\
& +\left(J_{4} \sin ^{2} \theta_{1} \sin ^{2} \theta_{2}+J_{5} \sin 2 \theta_{1} \sin 2 \theta_{2}\right) \sin \phi \\
& +\left(J_{6} \sin \theta_{1} \sin \theta_{2}+J_{7} \sin 2 \theta_{1} \sin 2 \theta_{2}\right) \cos \phi \\
& +J_{8} \sin ^{2} \theta_{1} \sin ^{2} \theta_{2} \sin 2 \phi+J_{9} \sin ^{2} \theta_{1} \sin ^{2} \theta_{2} \cos 2 \phi
\end{aligned}
$$

# The differential cross section <br> $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ 

$\square$ where the $J_{i}$ are given:

$$
\begin{aligned}
J_{1} & =2 r s\left(g_{A}^{2}+g_{V}^{2}\right)\left(\left|H_{1, V}\right|^{2}+\left|H_{1, A}\right|^{2}\right), \\
J_{2} & =\kappa\left(g_{A}^{2}+g_{V}^{2}\right)\left[\kappa\left(\left|H_{1, V}\right|^{2}+\left|H_{1, A}\right|^{2}\right)+\lambda \operatorname{Re}\left(H_{1, V} H_{2, V}^{*}+H_{1, A} H_{2, A}^{*}\right)\right], \\
J_{3} & =32 r s g_{A} g_{V} \operatorname{Re}\left(H_{1, V} H_{1, A}^{*}\right), \\
J_{4} & =4 \kappa \sqrt{r s \lambda} \operatorname{Re}\left(H_{1, V} H_{3, A}^{*}+H_{1, A} H_{3, V}^{*}\right), \\
J_{5} & =\frac{1}{2} \kappa \sqrt{r s \lambda}\left(g_{A}^{2}+g_{V}^{2}\right) \operatorname{Re}\left(H_{1, V} H_{3, A}^{*}+H_{1, A} H_{3, V}^{*}\right), \\
J_{6} & =4 \sqrt{r s \lambda} g_{A} g_{V}\left[4 \kappa \operatorname{Re}\left(H_{1, V} H_{2, V}^{*}\right)+\lambda \operatorname{Re}\left(H_{1, A} H_{2, V}^{*}\right)\right], \\
J_{7} & =\frac{1}{2} \sqrt{r s}\left(g_{A}^{2}+g_{V}^{2}\right)\left[2 \kappa\left(\left|H_{1, V}\right|^{2}+\left|H_{1, A}\right|^{2}\right)+\lambda \operatorname{Re}\left(H_{1, V} H_{2, V}^{*}+H_{1, A} H_{2, A}^{*}\right)\right], \\
J_{8} & =2 r s \sqrt{\lambda}\left(g_{A}^{2}+g_{V}^{2}\right) \operatorname{Re}\left(H_{1, V} H_{3, V}^{*}+H_{1, A} H_{3, A}^{*}\right), \\
J_{9} & =2 r s \sqrt{\lambda}\left(g_{A}^{2}+g_{V}^{2}\right)\left(\left|H_{1, V}\right|^{2}+\left|H_{1, A}\right|^{2}\right) .
\end{aligned}
$$

## The differential cross section <br> $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$

$$
\begin{aligned}
& H_{1, V}=-\frac{2 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{V}\left(1+\hat{\alpha}_{1}^{\mathrm{eff}}-\frac{\kappa}{r} \hat{\alpha}_{Z Z}^{\mathrm{eff}}-\frac{\kappa}{2 r} \frac{Q_{l} g_{e m}(r-s)}{s g_{V}} \hat{\alpha}_{A Z}\right), \\
& H_{1, A}=\frac{2 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{A}\left(1+\hat{\alpha}_{2}^{\mathrm{eff}}-\frac{\kappa}{r} \hat{\alpha}_{Z Z}\right), \\
& H_{2, V}=-\frac{2 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{V}\left(2 \hat{\alpha}_{Z Z}-\frac{Q_{l} g_{e m}(r-s)}{s g_{V}} \hat{\alpha}_{A Z}\right), \\
& H_{2, A}=\frac{4 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{A} \hat{\alpha}_{Z \hat{Z}}, \\
& H_{3, V}=-\frac{2 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{V}\left(2 \hat{\alpha}_{Z \tilde{Z}}+\frac{Q_{l} g_{e m}(r-s)}{s g_{V}} \hat{\alpha}_{A \tilde{Z}}\right), \\
& H_{3, A}=\frac{4 m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2} r}{r-s} g_{A} \hat{\alpha}_{Z \tilde{Z}},
\end{aligned}
$$

Where $Q_{l}=-1$.

$$
\begin{aligned}
\hat{\alpha}_{1}^{\mathrm{eff}} & \equiv \hat{\alpha}_{Z Z}^{(1)}-\frac{m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2}(r-s)}{2 \sqrt{r}} \frac{\hat{\alpha}_{\Phi l}^{V}}{g_{V}} \\
\hat{\alpha}_{2}^{\mathrm{eff}} & \equiv \hat{\alpha}_{Z Z}^{(1)}+\frac{m_{H}\left(\sqrt{2} G_{F}\right)^{1 / 2}(r-s)}{2 \sqrt{r}} \frac{\hat{\alpha}_{\Phi l}^{A}}{g_{A}}
\end{aligned}
$$

$\square$ The $\hat{\alpha}_{\Phi /}^{V}$ and $\hat{\alpha}_{\Phi /}^{A}$ curry the $d=6$ corrections into the Lagrangian.

## The differential cross section $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$

$\square J_{i}, H_{i, V}$ and $H_{i, A}$ are really hard to estimate;
$\square$ The challenge now is to figure out the $J_{i}, H_{i, V}$ and $H_{i, A}$;
$\square$ In our generator we try to kill the variables that curry $d=6$;
$\square$ This brings us back to the SM expression;
$\square$ We show results for fixing $J_{3}, J_{4}, J_{5}, J_{6}, J_{7}$ and $J_{8}$ to zero;
$\square$ Always $J_{1}=ل_{9}$, and $J_{3}$ will be changed to see how the angles affected.

## Preliminary results






## Preliminary results SM




## Preliminary results

$\operatorname{BSM}\left(J_{1}=3, J_{2}=3.5, J_{1}=1.0, J_{8}=1.0, J_{9}=1.0\right)$





## Preliminary results

$\operatorname{BSM}\left(J_{1}=3.0, J_{2}=3.5, J_{1}=1.0, J_{8}=1, J_{9}=1.0\right)$



## Summary

$\square$ We introduced the potential of the differential cross section of $e^{+} e^{-} \rightarrow H Z(\rightarrow \ell \ell)$ in probing the BSM scenario;
$\square$ The results provided here are from under developing generator for $e^{+} e^{-}$collider; and
$\square$ However, still lots of work have to be done for the estimation of the $d=6$ parameters.

## Plans

- Take background into account;
- Do a 3D fit over $\theta_{1}, \theta_{2}$ and $\phi$;
- See how the effect of the $J_{i}$ 's, and probe the sensitivity of CEPC experiment.


## Thank you!

