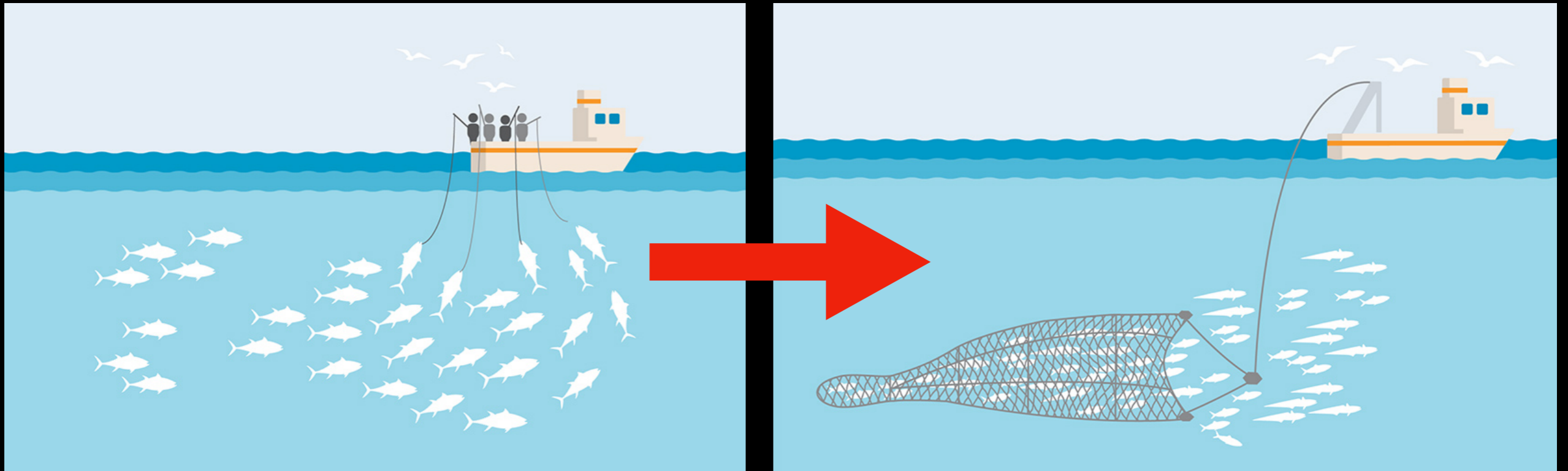


A global analysis approach

A single parameter quantifies
both Higgs Br precision and detector performance

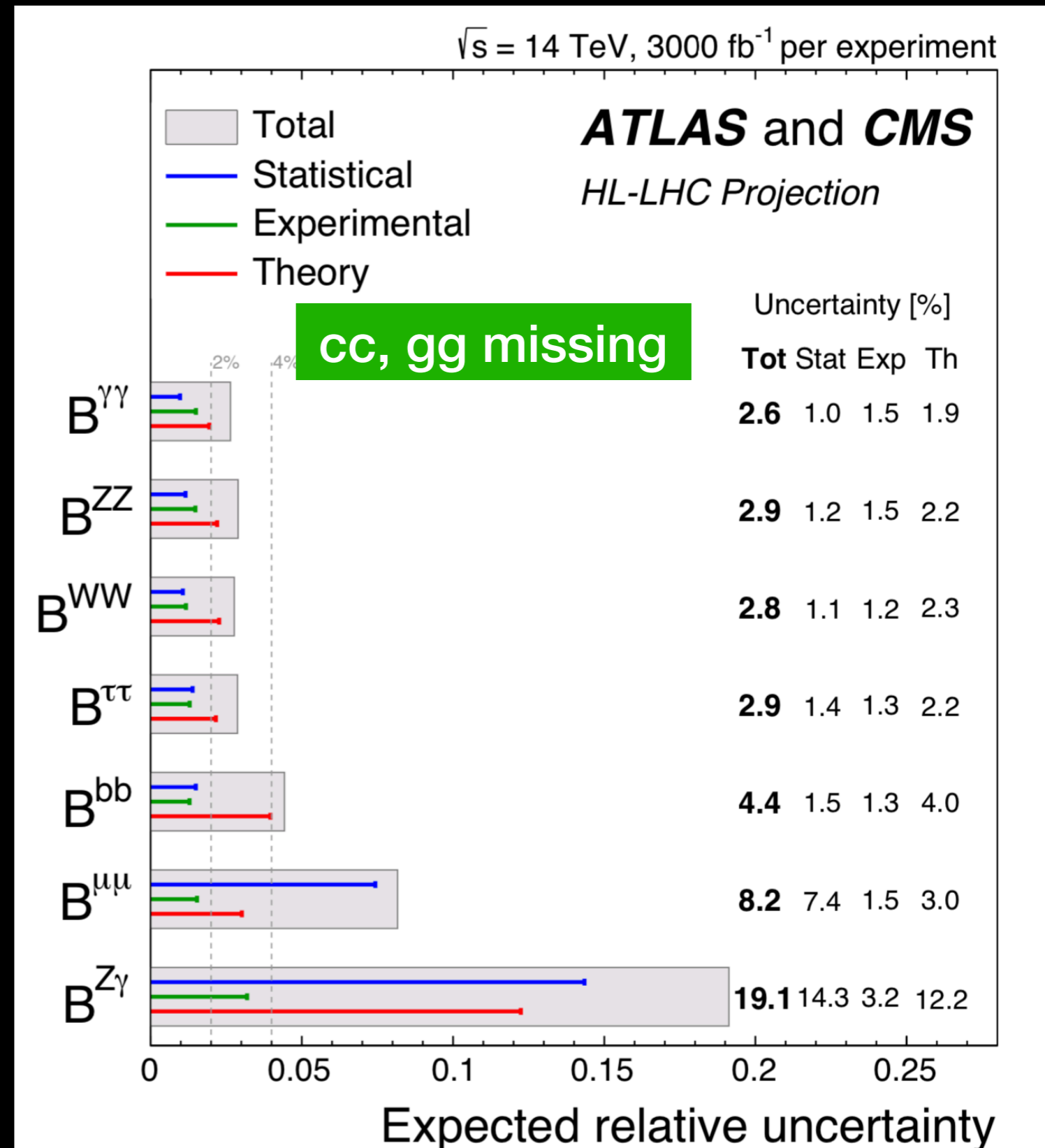


Outline

- Motivation
- Method
- Some numerical results with toy MC
- Discussion & summary

- Competition from both HL-LHC and FCC-ee

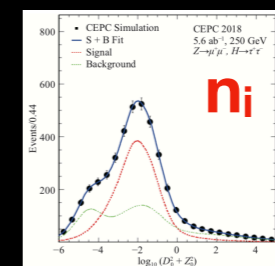
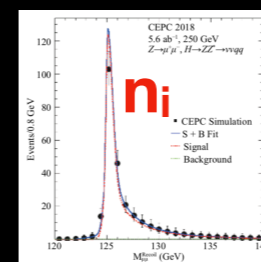
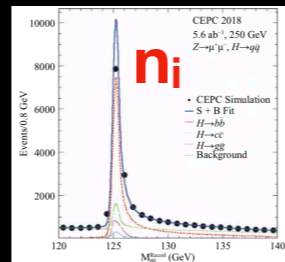
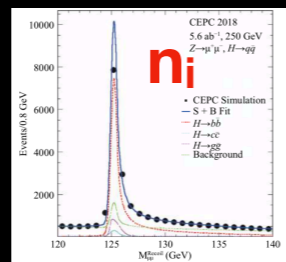
- FCC-ee
- ATLAS-CMS extrapolation range from 2 – 4%, with the exception of that on $B^{\mu\mu}$ at 8% and on $B^{Z\gamma}$ at 19%.



We possess what the LHC lacks (人无我有)

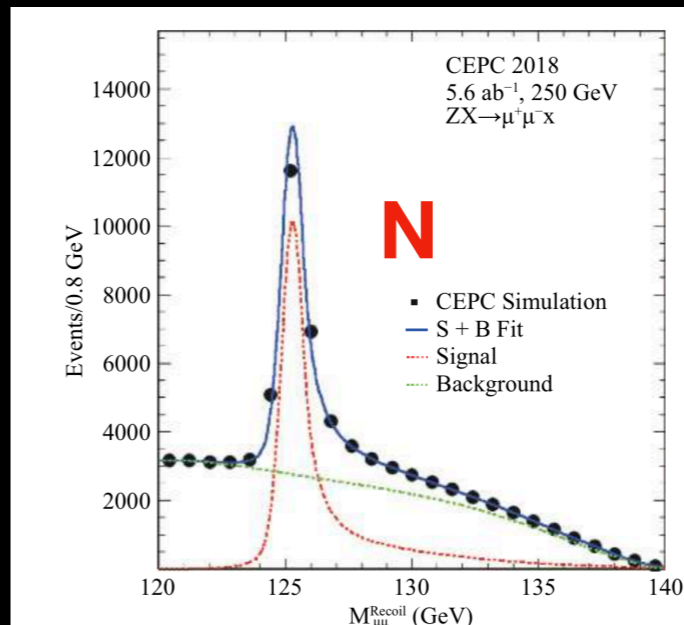
- Tagging method, absolute/model-independent
- All Higgs decays accessible except e and uds
- **Multinomial distribution: statistical constraint**
- Two types of backgrounds
 - Higgs background (crosstalk)
 - non-Higgs background (enlarge the stat. unc. of n_i)

$$B_i = \frac{n_i / \epsilon_{ii}}{N}$$



B_i =

non-Higgs background
 — subtracted with fitting
 — but enlarges σ_{n_i}



Take the simplest case as an example
– 2 decay modes

Efficiency matrix

Based on MC, no dependence on Br's

$$\begin{pmatrix} n_1 \\ n_2 \end{pmatrix} = \begin{matrix} \text{MODULATION Matrix} \\ \begin{pmatrix} \epsilon_{11} & \epsilon_{12} \\ \epsilon_{21} & \epsilon_{22} \end{pmatrix} \end{matrix} \times \begin{pmatrix} N_1 \\ N_2 \end{pmatrix}$$

Observation ↑ Production

A produced final state j reconstructed as final state i

Measurement: **DEMODULATION**

All knowns on the right

Solve N and minimize its uncertainty

DEMODULATION

$$N = E^{-1}n$$

2 decays
 $p+q=1$ — binomial distribution

Individual measurements

$$\sigma_{n_1} = \sqrt{N\epsilon_{11}}$$

$$\sigma_{n_2} = \sqrt{N\epsilon_{22}}$$



Global measurement

$$\sigma_{n_1} = \sqrt{Np(1-p)\epsilon_{11}}$$

$$\sigma_{n_2} = \sqrt{Nq(1-q)\epsilon_{22}}$$

$$\sum N_i = N$$

Based on text book, please read
https://en.wikipedia.org/wiki/Binomial_distribution
https://en.wikipedia.org/wiki/Multinomial_distribution

binomial / multinomial distributions

More on the full covariance

$$V = \begin{pmatrix} Npq & -Npq \\ -Npq & Npq \end{pmatrix}$$

100% anti-correlated between the two decays!
This can be used in data analysis to improve precisions.

Successful examples

- Precision measurement of the D^{*0} decay branching fractions by BESIII, Phys. Rev. D91 (2015) no.3, 031101
- Branching ratios of tau decays by ALEPH, Physics Reports 421 (2005) 191–284

Let's see how it happens

$$\vec{\sigma}_n = \begin{pmatrix} \vec{\sigma}_{n_1} \\ \vec{\sigma}_{n_2} \end{pmatrix},$$

$$\sigma_{n_i}^2 = \vec{\sigma}_{n_i} \cdot \vec{\sigma}_{n_i}, \quad \sigma_{n_{ij}} = \sigma_{n_{ji}} = \sigma_{n_1} \sigma_{n_2} \rho_{ij} = \vec{\sigma}_{n_i} \cdot \vec{\sigma}_{n_j},$$

$$\rho_{ij} = \frac{\vec{\sigma}_{n_i} \cdot \vec{\sigma}_{n_j}}{\sigma_{n_1} \sigma_{n_2}},$$

Matrix: compact and easy to expand to higher dimension

$$\begin{aligned} \Sigma^n &= \vec{\sigma}_n \vec{\sigma}_n^T = \begin{pmatrix} \vec{\sigma}_{n_1} \\ \vec{\sigma}_{n_2} \end{pmatrix} \begin{pmatrix} \vec{\sigma}_{n_1} \\ \vec{\sigma}_{n_2} \end{pmatrix}^T \\ &= \begin{pmatrix} \sigma_{n_1}^2 & \sigma_{n_{12}} \\ \sigma_{n_{21}} & \sigma_{n_2}^2 \end{pmatrix} \end{aligned}$$

Space transformation

$n(\text{observable}) \rightarrow N(\text{production})$

$$N = E^{-1}n$$



$$\vec{\sigma}_N = E^{-1} \vec{\sigma}_n$$

$$J_{Nn} = E^{-1} = \frac{1}{|E|} \begin{pmatrix} \epsilon_{22} & -\epsilon_{12} \\ -\epsilon_{21} & \epsilon_{11} \end{pmatrix} \equiv \frac{J_N}{|E|}$$

Space transformation

N(production) → B(branching ratios)

- Features

- ✓ Variance of B proportional to $1/(N^4|E|^2)$
- ✓ N^4 : statistical power
- ✓ $|E|^2$ proportional to the performance of Detector x Reconstruction x Analysis
- ✓ Same uncertainties for both Br's

$$J_{BN} = \frac{1}{N^2} \begin{pmatrix} N_2 & -N_1 \\ -N_2 & N_1 \end{pmatrix} = \frac{J_B}{N^2}$$

$$\vec{\sigma}_B = J_{BN} J_{Nn} \vec{\sigma}_n = \frac{\begin{pmatrix} n_2 \vec{\sigma}_1 - n_1 \vec{\sigma}_2 \\ -n_2 \vec{\sigma}_1 + n_1 \vec{\sigma}_2 \end{pmatrix}}{N^2 |E|}$$

$$\begin{aligned} \Sigma_B &= \frac{\vec{\sigma}_B \vec{\sigma}_B^T}{N^4 |E|^2} \\ &= \frac{[J_B J_N \vec{\sigma}_n][J_B J_N \vec{\sigma}_n]^T}{N^4 |E|^2} \\ &= \frac{(n_2 \sigma_{n_1} + n_1 \sigma_{n_2})^2}{N^4 |E|^2} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \end{aligned}$$

More than 2 decay modes

$$\Sigma_B = \frac{\vec{\sigma}_B \vec{\sigma}_B^T}{N^4 |E|^2}$$

=

$$\frac{1}{N^4 |E|^2} \begin{pmatrix} \left\{ \sum_j [N_1 \frac{\partial N}{\partial n_j} - N \frac{\partial N_1}{\partial n_j}] \vec{\sigma}_{n_j} \right\}^2 & \dots & \dots \\ \dots & \left\{ \sum_j [N_2 \frac{\partial N}{\partial n_j} - N \frac{\partial N_2}{\partial n_j}] \vec{\sigma}_{n_j} \right\}^2 & \dots \\ \dots & \dots & \left\{ \sum_j [N_3 \frac{\partial N}{\partial n_j} - N \frac{\partial N_3}{\partial n_j}] \vec{\sigma}_{n_j} \right\}^2 \end{pmatrix}$$

Similar features as N=2

Numerical results with toy MC

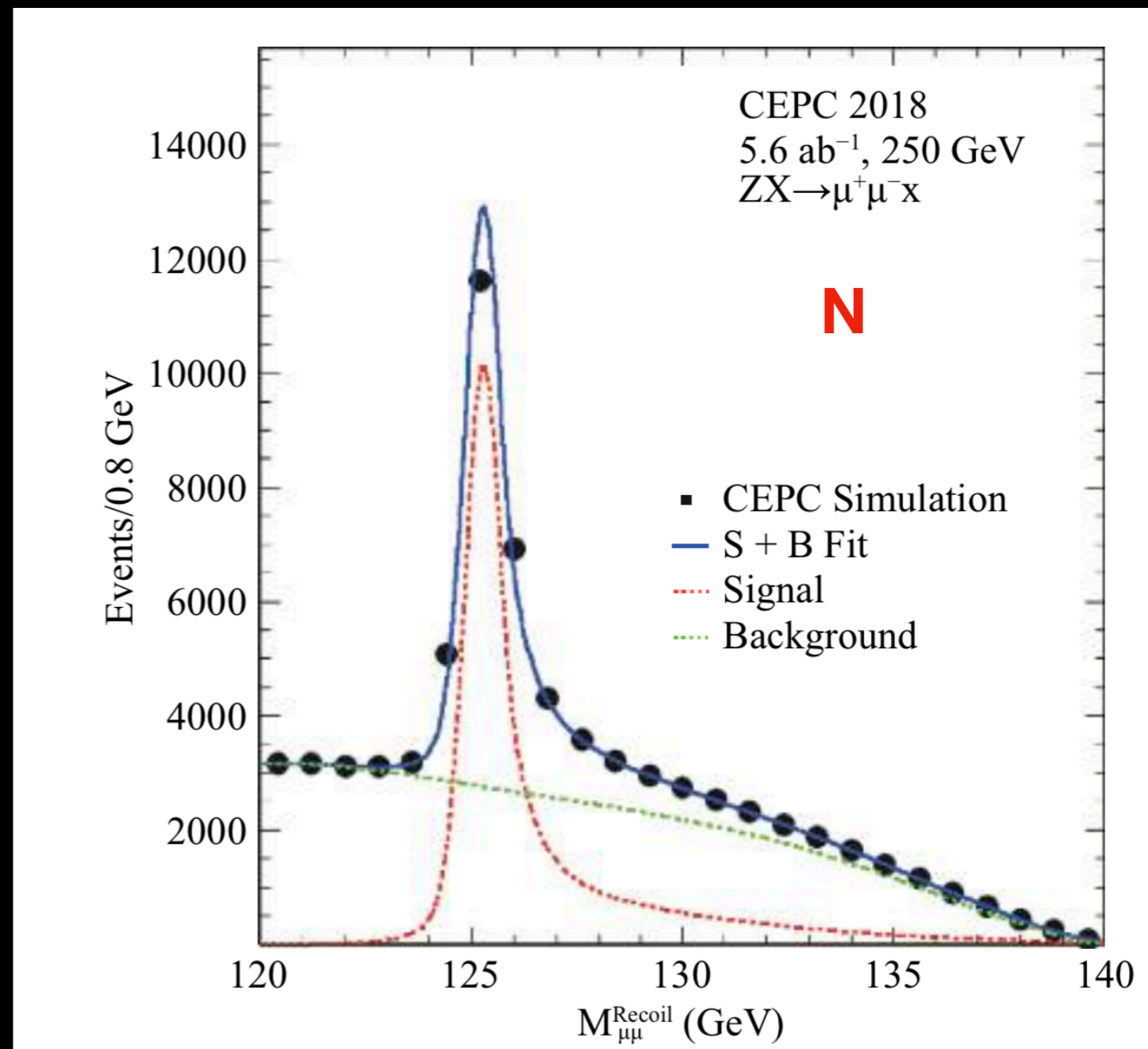
On backgrounds

- Two type of backgrounds

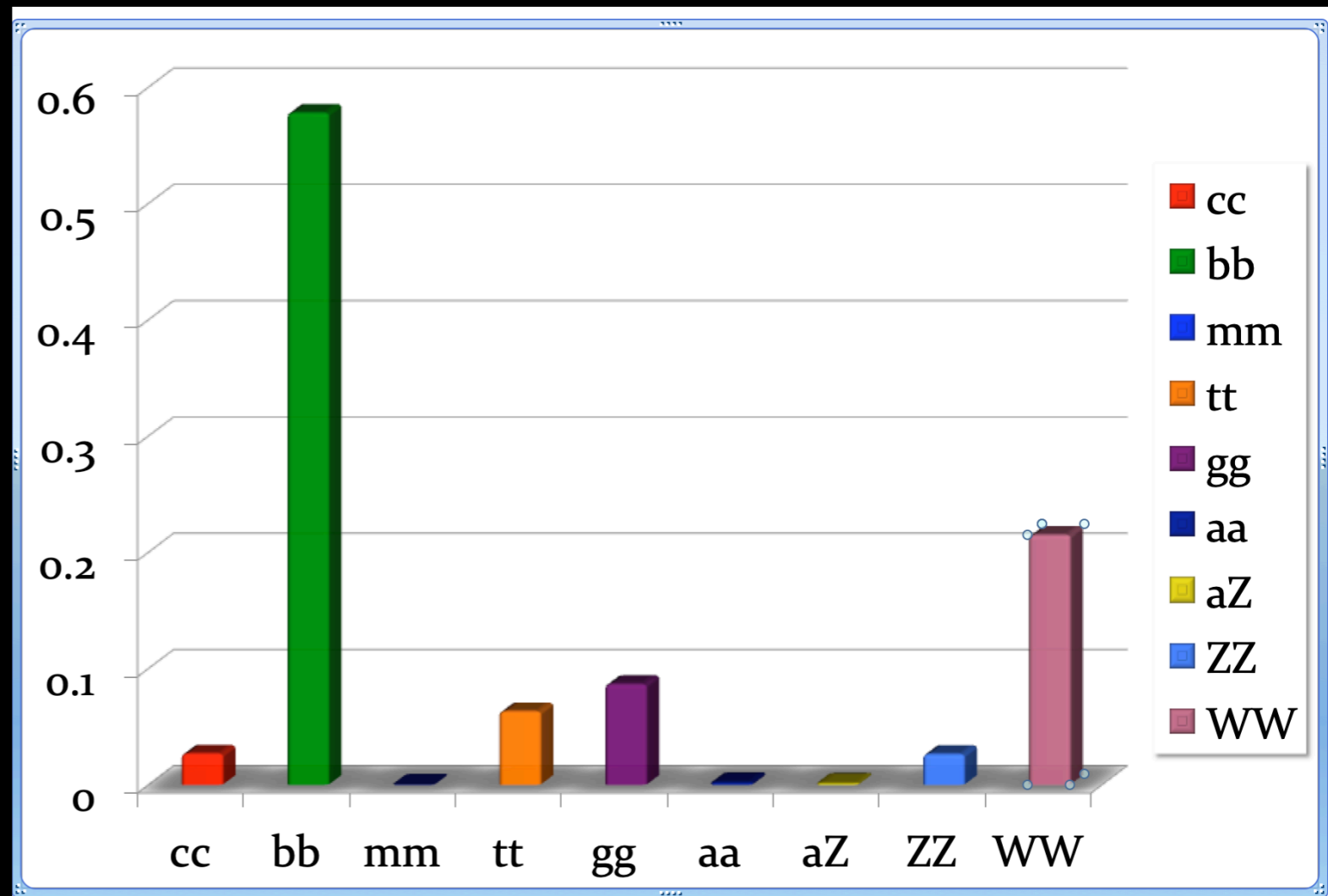
➤ Non-uuH backgrounds: subtracted by fitting, enlarging statistical uncertainty of n_i

➤ uuH backgrounds(cross talk): the efficiency matrix dealing with them

➤ N and n_i



- 9 Higgs decays accessible
- Di-muon, di-photon, and gamma Z are tiny: only 0.02%, 0.23%, and 0.15%, respectively
- cc contaminated by bb due to large bb Br
- ZZ important for Higgs Width



Solve N_i by minimizing the χ^2 with constraint

$$\chi^2 = \sum_i \frac{(\sum_j \epsilon_{ij} N_j - n_i)^2}{\sigma_{n_i}^2} + \frac{(\sum_l N_l - N)^2}{\sigma_N^2}$$

Higgs \rightarrow cc, bb, mm, tt, gg, aa, aZ, ZZ, WW

1 2 3 4 5 6 7 8 9

$$\begin{pmatrix} n_1 \\ n_2 \\ n_3 \\ n_4 \\ n_5 \\ n_6 \\ n_7 \\ n_8 \\ n_9 \end{pmatrix} = \begin{pmatrix} \epsilon_{11} & \epsilon_{12} & \epsilon_{13} & \epsilon_{14} & \epsilon_{15} & \epsilon_{16} & \epsilon_{17} & \epsilon_{18} & \epsilon_{19} \\ \epsilon_{21} & \epsilon_{22} & \epsilon_{23} & \epsilon_{24} & \epsilon_{25} & \epsilon_{26} & \epsilon_{27} & \epsilon_{28} & \epsilon_{29} \\ \epsilon_{31} & \epsilon_{32} & \epsilon_{23} & \epsilon_{34} & \epsilon_{35} & \epsilon_{36} & \epsilon_{37} & \epsilon_{38} & \epsilon_{39} \\ \epsilon_{41} & \epsilon_{42} & \epsilon_{33} & \epsilon_{44} & \epsilon_{45} & \epsilon_{46} & \epsilon_{47} & \epsilon_{48} & \epsilon_{49} \\ \epsilon_{51} & \epsilon_{52} & \epsilon_{43} & \epsilon_{54} & \epsilon_{55} & \epsilon_{56} & \epsilon_{57} & \epsilon_{58} & \epsilon_{59} \\ \epsilon_{61} & \epsilon_{62} & \epsilon_{53} & \epsilon_{64} & \epsilon_{65} & \epsilon_{66} & \epsilon_{67} & \epsilon_{68} & \epsilon_{69} \\ \epsilon_{71} & \epsilon_{72} & \epsilon_{63} & \epsilon_{74} & \epsilon_{75} & \epsilon_{76} & \epsilon_{77} & \epsilon_{78} & \epsilon_{79} \\ \epsilon_{81} & \epsilon_{82} & \epsilon_{73} & \epsilon_{84} & \epsilon_{85} & \epsilon_{86} & \epsilon_{87} & \epsilon_{88} & \epsilon_{89} \\ \epsilon_{91} & \epsilon_{92} & \epsilon_{83} & \epsilon_{94} & \epsilon_{95} & \epsilon_{96} & \epsilon_{97} & \epsilon_{98} & \epsilon_{99} \end{pmatrix} \begin{pmatrix} N_1 \\ N_2 \\ N_3 \\ N_4 \\ N_5 \\ N_6 \\ N_7 \\ N_8 \\ N_9 \end{pmatrix}$$

Neglect e and uds decays — constraint feasible

$$\sum_i N_i = N^{tag} \text{ or } \sum_i B_i = 1$$

$$B_i = \frac{N_i}{N}$$

Statistical limit

- 99% efficiency,
- no cross talk,
- no other backgrounds
- eeH and qqH as good as muuH

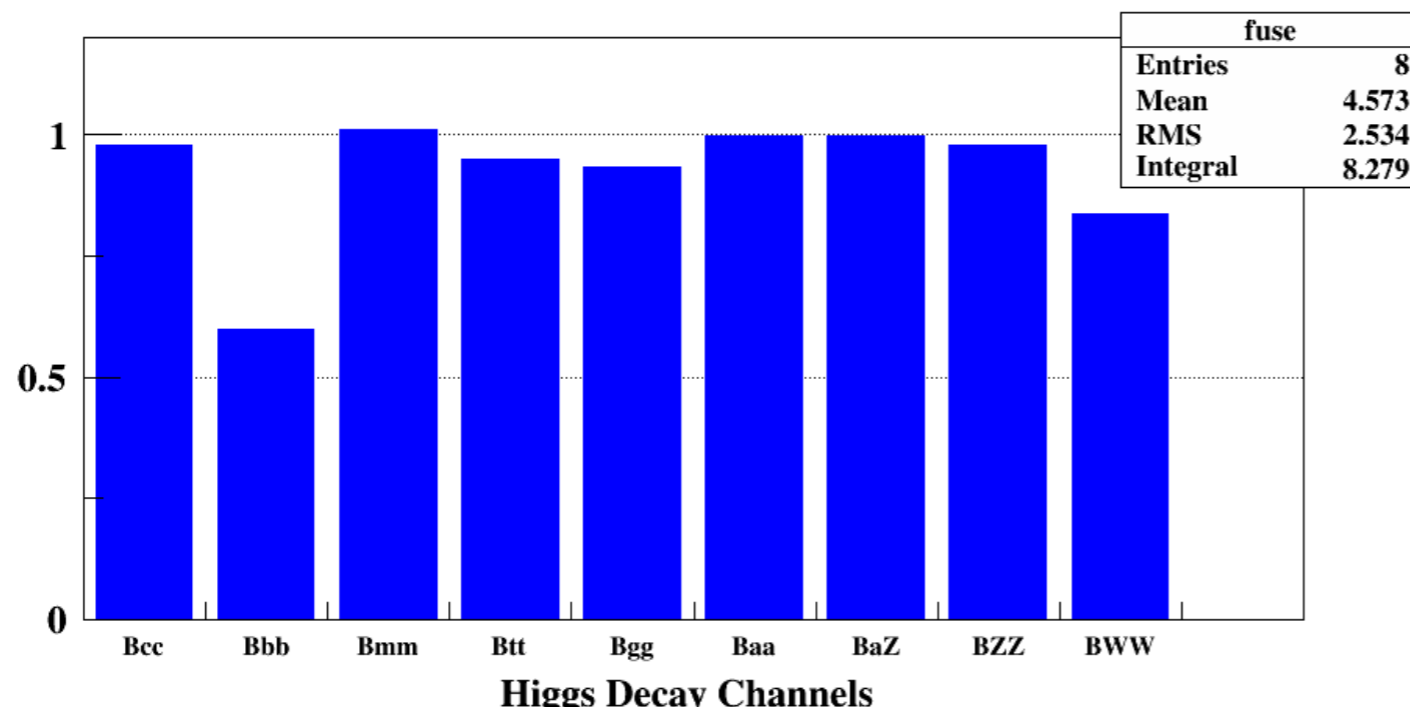
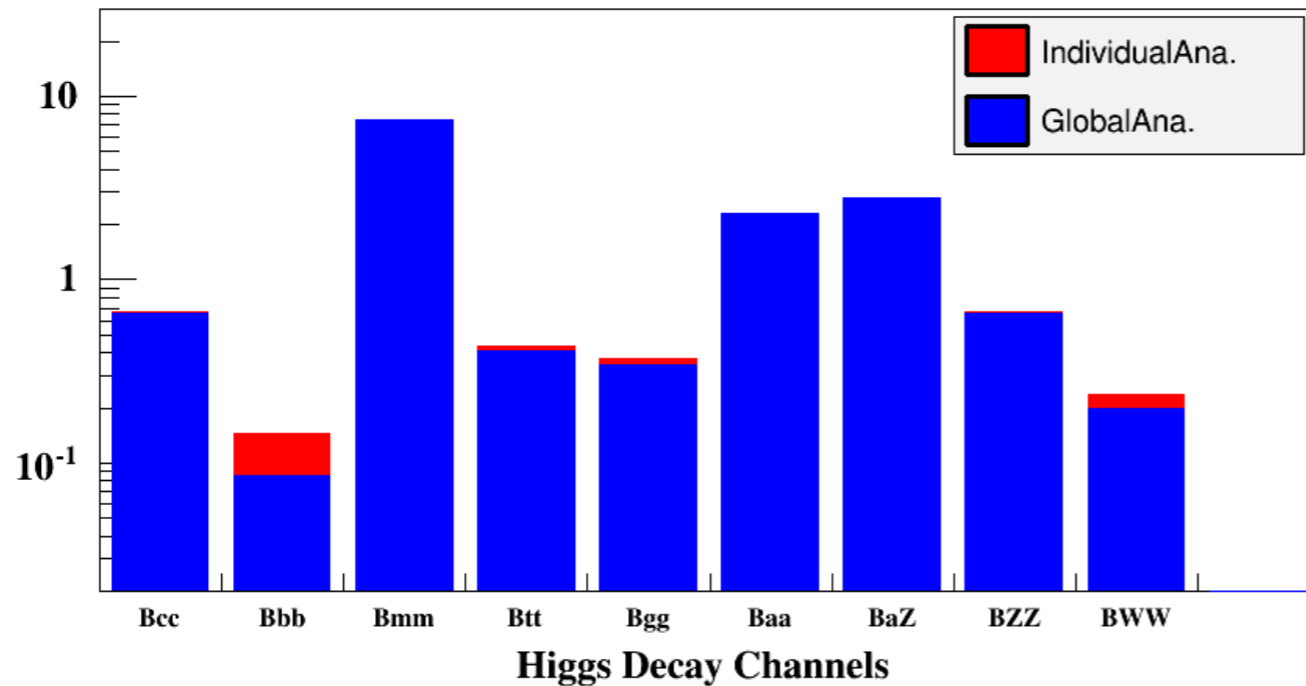
$$E = \begin{pmatrix} 0.99 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.99 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.99 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0.99 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0.99 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0.99 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0.99 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.99 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.99 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.99 \end{pmatrix}$$

$$N = L \times (\sigma_{\mu\mu H} + \sigma_{eeH} + \sigma_{qqH}) = 5600 \times (6.77 + 7.04 + 136.81) = 843,372$$

Ideal case:

eeH, qqh as good as uuH

No background, no cross talk, multinomial uncertainties, and constraint



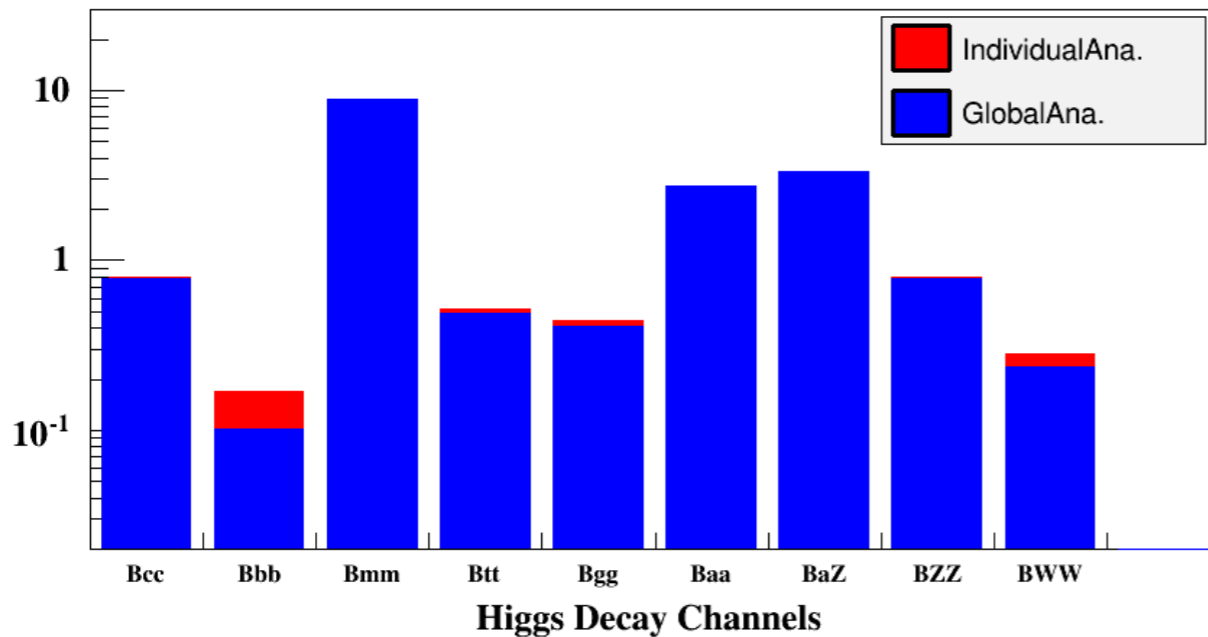
$$\sigma_{n_i} = \sqrt{Np(1-p)\epsilon_{ii}}$$

		MLT	POS
Bcc	2.713%	0.650%	0.664%
Bbb	57.799%	0.086%	0.144%
Bmm	0.023%	7.190%	7.198%
Btt	6.319%	0.413%	0.435%
Bgg	8.619%	0.347%	0.373%
Baa	0.227%	2.294%	2.299%
BaZ	0.150%	2.818%	2.822%
BZZ	2.647%	0.658%	0.673%
BWW	21.496%	0.198%	0.236%

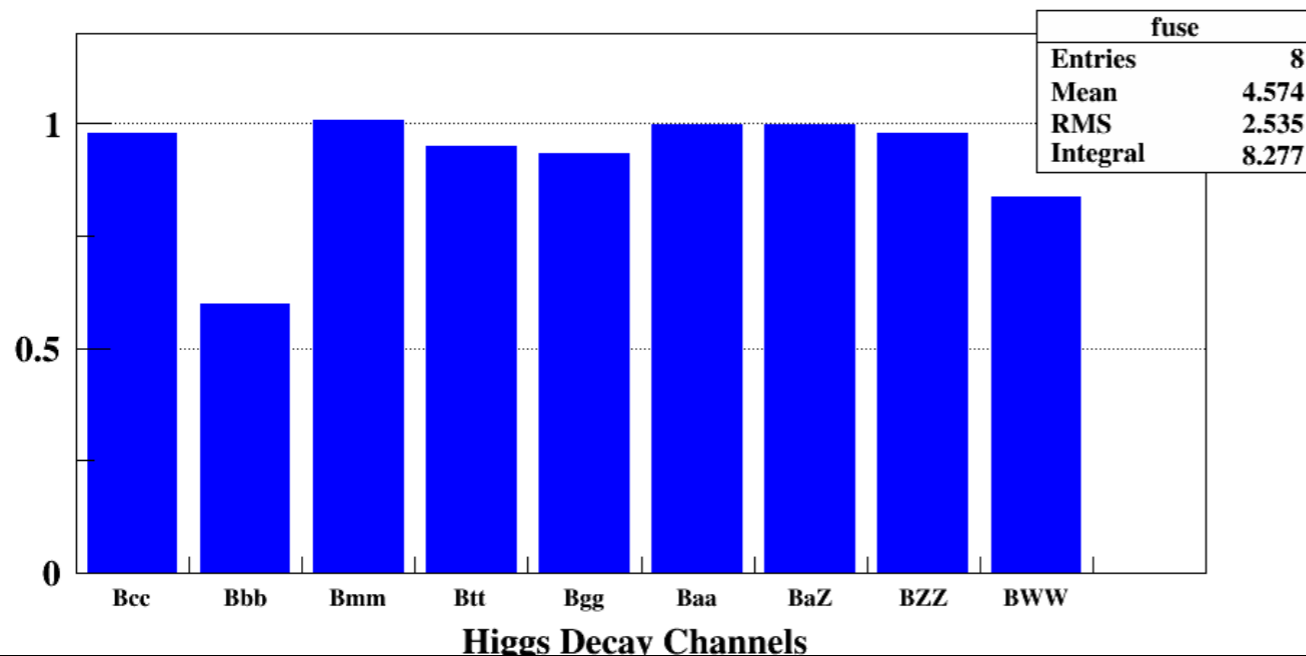
More realistic:

eeH, qqh as good as uuH

100% background, no cross talk, multinomial uncertainties, and constraint



$$\sigma_{n_i} = \sqrt{Np(1-p)\epsilon_{ii}}$$



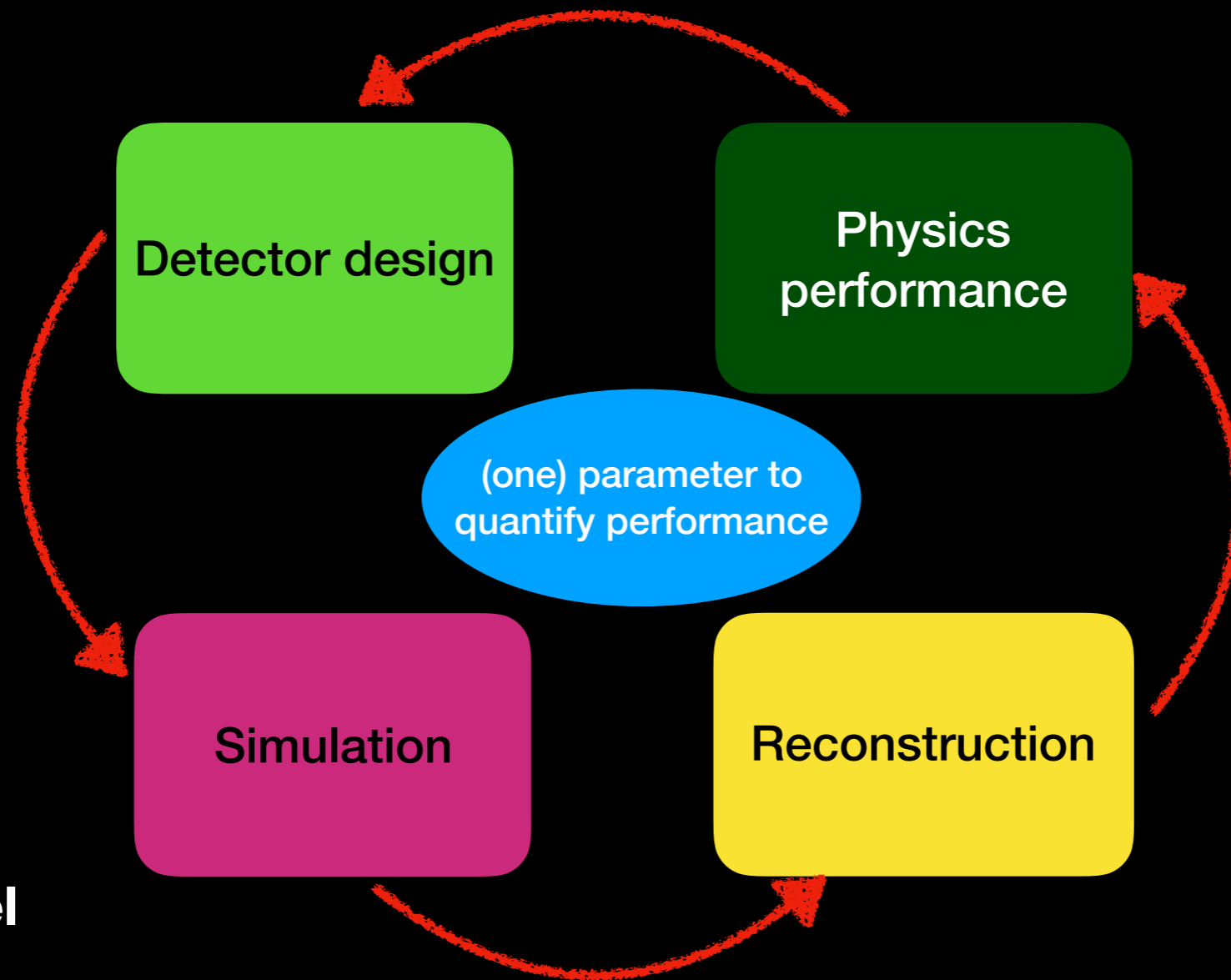
MLT

POS

Bcc	2.713%	0.773%	(0.779%	0.790%)
Bbb	57.799%	0.102%	(0.111%	0.171%)
Bmm	0.023%	8.547%	(8.559%	8.560%)
Btt	6.319%	0.492%	(0.501%	0.518%)
Bgg	8.619%	0.413%	(0.424%	0.443%)
Baa	0.227%	2.728%	(2.731%	2.734%)
BaZ	0.150%	3.350%	(3.353%	3.356%)
BZZ	2.647%	0.783%	(0.789%	0.800%)
BWW	21.496%	0.235%	(0.249%	0.281%)

Short discussion

- This approach can improve Higgs branching ratio measurement and set a statistical limit
- qqH and eeH not good as uuH , but much more statistics
- Degrading in real analysis and lots of compromises
- No full cross talk information in current analyses



- Good design:
- ★ From top level
 - ★ Break-down

Detector design & Optimization

Multi-purpose optimization: a bunch of benchmarks —
A single parameter is favored, which means single-purpose optimization

Physics performance parameterized
as a function of several parameters,
or precision of a set of benchmark processes
or determinant of efficiency matrix $|E|$

Difficult

$$P = f(\sigma_p, \sigma_{E_\gamma}, PID, JID, JER, \dots)$$
$$= |E|^2 \propto \frac{1}{|\Sigma_B|^2}$$

Easy to minimize

Now problem successfully becomes
how to Maximize $|E|^2$

Again on efficiency matrix

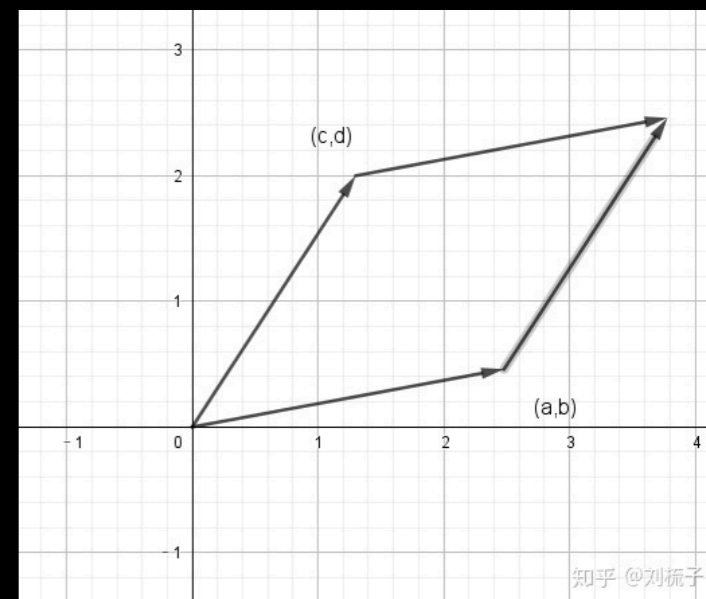
- *Not necessary to know the branching ratios of Higgs decays*
- *Quantifies the detector/software/analysis performance with a single parameter $\det|E|$*
- *It could be useful for detector optimization*

A single purpose optimization instead of that of a bunch of benchmarks

$$\Sigma_B \propto \frac{1}{|E|^2}$$

Geometrical interpretation of the efficiency matrix

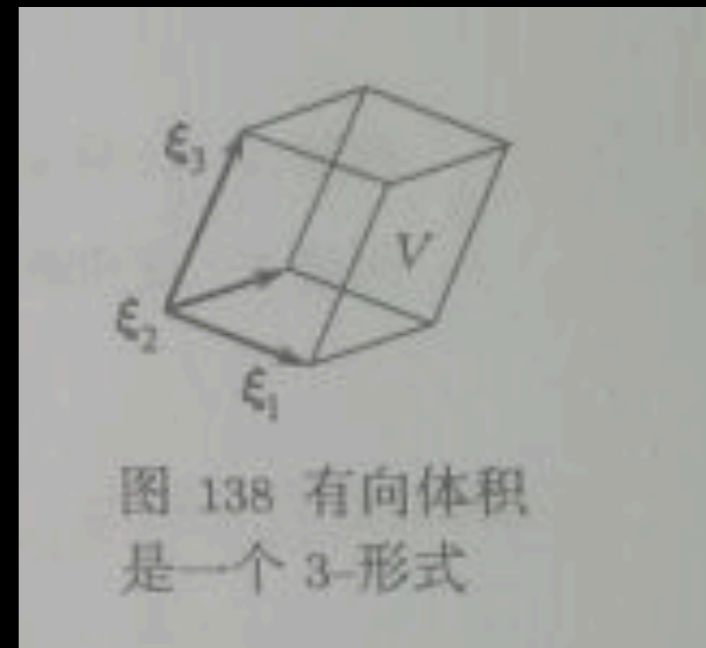
- For a matrix $E = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$, $\det(E)$ is the area of a quadrilateral
- $S = \det(E) = ad - bc$
- For $n=3$, $\det(E)$ is volume of a parallelepiped
- For $N>3$, $\det(E)$ is hyper volume ...



例 1 有向欧氏空间 \mathbb{R}^n 中的以 ξ_1, \dots, ξ_n 为棱的平行体的有向体积是一个 n -形式 (图 138).

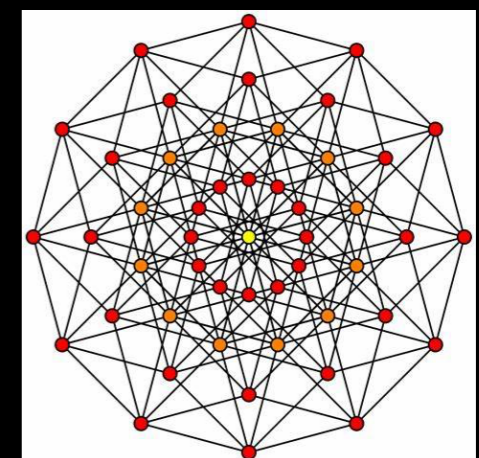
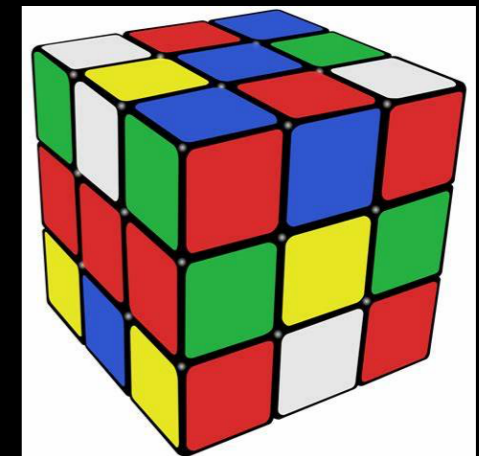
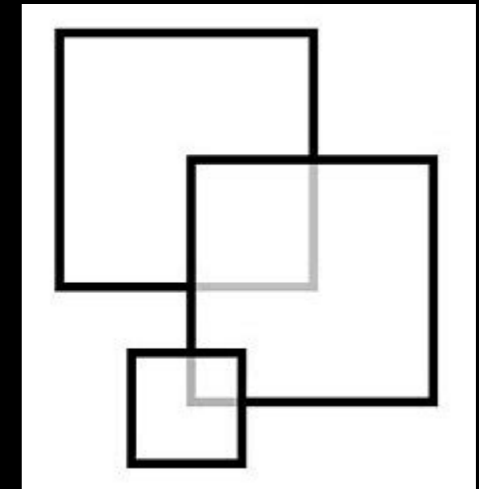
$$V(\xi_1, \dots, \xi_n) = \begin{vmatrix} \xi_{11} & \dots & \xi_{1n} \\ \dots & \dots & \dots \\ \xi_{n1} & \dots & \xi_{nn} \end{vmatrix}.$$

这里 $\xi_i = \xi_{i1}e_1 + \dots + \xi_{in}e_n$, e_1, \dots, e_n 是 \mathbb{R}^n 的一个基底.



Maximize $|E|^2$

- $N=2$, the maximum $|E| \rightarrow$ area of a square
- $N=3$, the maximum $|E| \rightarrow$ volume of a cube
- $N>3$, ... \rightarrow volume of a HyperCube
- Hypercube efficiency matrix has the ideal/best performance, this is the dream of experimentalists



Detector: HC
HyperCube or HiggsCube

Summary

- CEPC dedicated for Higgs study, the Br's important
- Global analysis with constraint improves the precision
- Global analysis of $e^+e^- \rightarrow u^+u^-H$, $H \rightarrow$ all 9 decay modes serves as a “benchmark” to optimize detector, software, and analysis,
- Advantage : single parameter, easy to optimize, easily to realize in ML
- Using fast simulation + global analysis + machine learning to maximize $|E|$ — fast iteration
- Including eeH and qqH much better but difficult, possible to do ..., not very necessary at present